

New results from Double Chooz



Leonidas N. Kalousis (IPHC Strasbourg)
for the Double Chooz Collaboration
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The hunt for θ_{13}

- Double Chooz is a two-detector short baseline experiment searching for the disappearance of reactor electron antineutrinos.
- Tight constraints from single reactor experiments performed in the 1990's (CHOOZ, Palo Verde).
- CHOOZ limit: $\sin^2(2\theta_{13}) < 0.20$ at 90% C.L. assuming $\Delta m^2_{\text{atm}} = 2.0 \cdot 10^{-3} \text{ eV}^2$.

Predominant searches:

- Disappearance mode: $\bar{\nu}_e \longrightarrow ?$
 - (Reactor exp.) Double Chooz, RENO and Daya Bay
- Appearance channel: $\nu_\mu \longrightarrow \nu_e$ and $\bar{\nu}_\mu \longrightarrow \bar{\nu}_e$
 - (Accelerator exp.) MINOS, T2K and Nova
- Note though, that lower bounds (hints) can also be obtained through three flavour analyses to Solar and Atmospheric data.
- The value of θ_{13} is important for the observation of genuine three flavour effects like CP violation in the lepton sector.

Neutrino mixing and oscillations

- The neutrino mass (energy) eigenstates ν_1, ν_2, ν_3 , do not coincide with the flavor (interaction) eigenstates ν_e, ν_μ, ν_τ .
- Leptonic mixing matrix (in the minimal scenario):

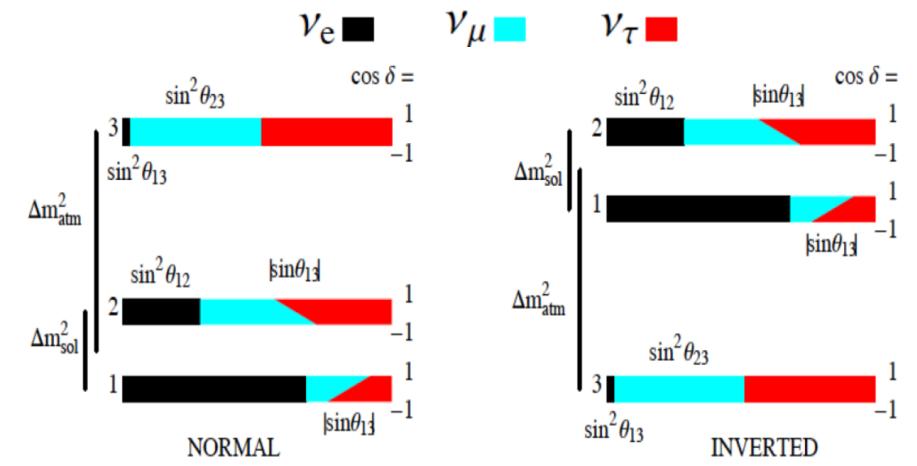
$$|\nu_\alpha\rangle = \sum_{i=1}^3 U_{\alpha i}^* |\nu_i\rangle$$

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_{23} & \sin \theta_{23} \\ 0 & -\sin \theta_{23} & \cos \theta_{23} \end{pmatrix} \times \begin{pmatrix} \cos \theta_{13} & 0 & e^{-i\delta} \sin \theta_{13} \\ 0 & 1 & 0 \\ -e^{i\delta} \sin \theta_{13} & 0 & \cos \theta_{13} \end{pmatrix} \times \begin{pmatrix} \cos \theta_{12} & \sin \theta_{12} & 0 \\ -\sin \theta_{12} & \cos \theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \times \text{diag}\{e^{i\alpha_1/2}, e^{i\alpha_2/2}, 1\}$$



Lepton flavour violation through
the mechanism of neutrino oscillations

- The oscillation patterns are driven by the mass square differences, $\Delta m_{ij}^2 = m_i^2 - m_j^2$
- The oscillation formula is a function of L/E, (distance)/(energy).

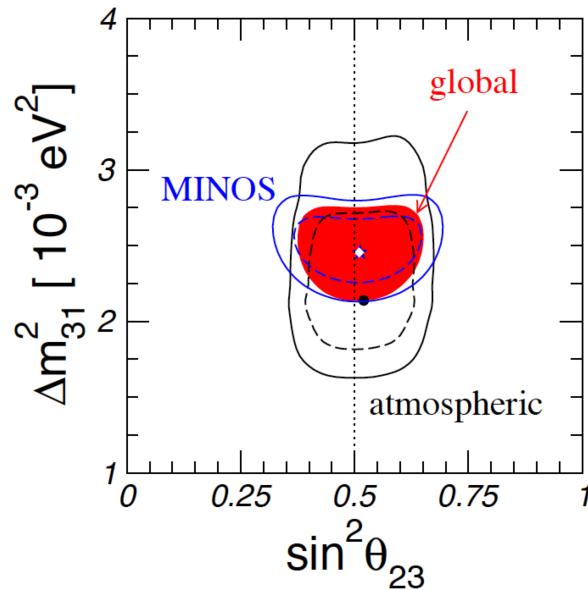


O. Mena and S. Parke Phys. Rev. D 69, 117301 (2004)

$$p(\nu_\alpha \rightarrow \nu_\beta) = \sum_{i,j} U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^* e^{-i \left(\frac{\Delta m_{ij}^2 L}{2E} \right)}$$

Current Understanding

Atmospheric sector



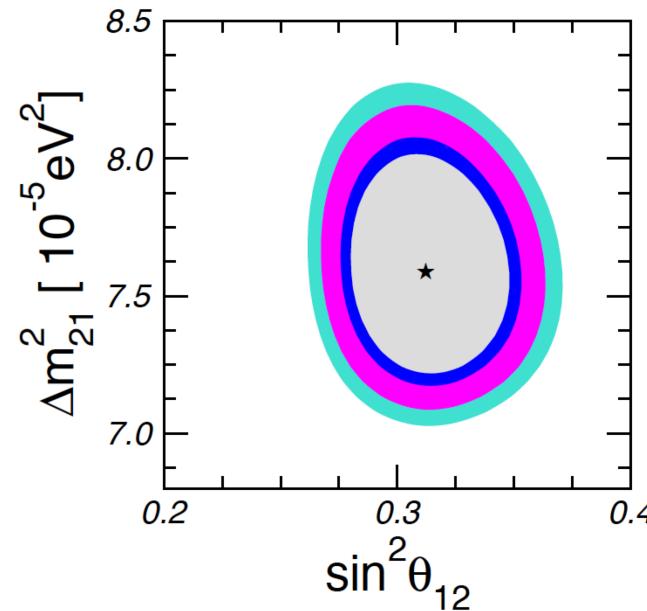
Super-K and MINOS

$$|\Delta m_{31}^2| \approx 2.4 \times 10^{-3} \text{ eV}^2$$

$$\sin(2\theta_{23}) \approx 1.0$$

(almost maximal mixing)

Solar sector



Solar exp. and KamLAND

$$\Delta m_{21}^2 \approx 7.6 \times 10^{-5} \text{ eV}^2$$

$$\sin(2\theta_{12}) \approx 0.85$$

(large mixing angle)

1-3 sector

- Reactor experiments
 - CHOOZ
 - Palo Verde
 - Solar exp. and KamLAND
 - Atmospheric exp.
 - K2K
 - MINOS
- Results from global fits point towards a non-zero θ_{13} scenario.

No measurement before 2011

$$\sin(2\theta_{13}) \approx 0.1$$

θ_{13} is surprisingly small

Oscillation search with reactor antineutrinos

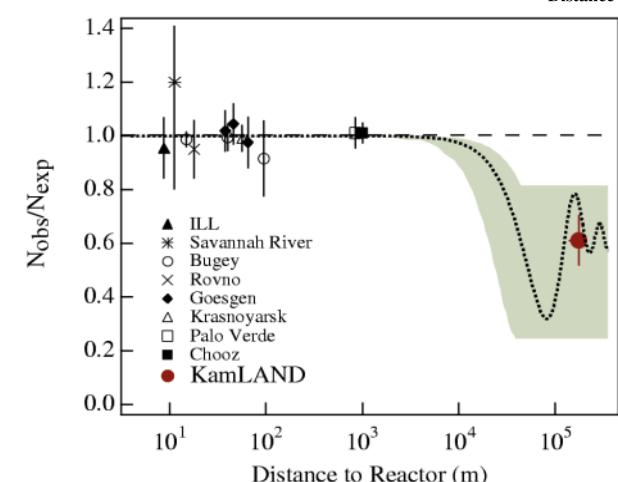
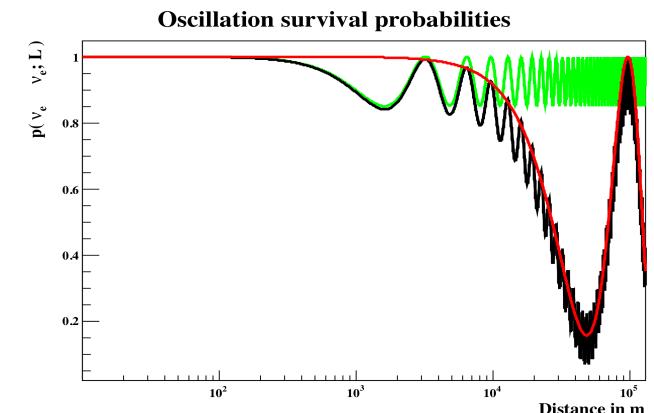
- Reactors are copious sources of $\bar{\nu}_e$.
 - An 1 GW_{th} power reactor emits $2 \cdot 10^{20} \bar{\nu}_e/\text{sec.}$
- This intense flux permitted:
 - The observation of neutrino (Reines and Cowan)
 - Study of the electroweak sector of the SM.
- Flux characteristics:
 - Pure composition
 - Isotropic flux
 - Low energy antineutrinos, $\langle E_\nu \rangle = 3 - 4 \text{ MeV}$.
- $\bar{\nu}_e$ disappearance channel:

(L in km, E in GeV and Δm^2 in eV²)

$$p(\nu_e \rightarrow \nu_e) \simeq 1 - \sin^2(2\theta_{13}) \sin^2 \left(1.27 \Delta m_{13}^2 \frac{L}{E} \right) - \cos^4 \theta_{13} \sin^2(2\theta_{12}) \sin^2 \left(1.27 \Delta m_{12}^2 \frac{L}{E} \right)$$

- No CP-violating terms,
- Negligible matter effects and small dependence on Δm_{12}^2 at $\approx 1 \text{ km.}$

$\langle E_\nu \rangle = 3.0 \text{ MeV}$
 $\Delta m_{13}^2 = 2.30 \cdot 10^{-3} \text{ eV}^2 \text{ and } \sin^2(2\theta_{13}) = 0.160$
 $\Delta m_{12}^2 = 7.65 \cdot 10^{-5} \text{ eV}^2 \text{ and } \sin^2(2\theta_{12}) = 0.304$



Reactor systematics

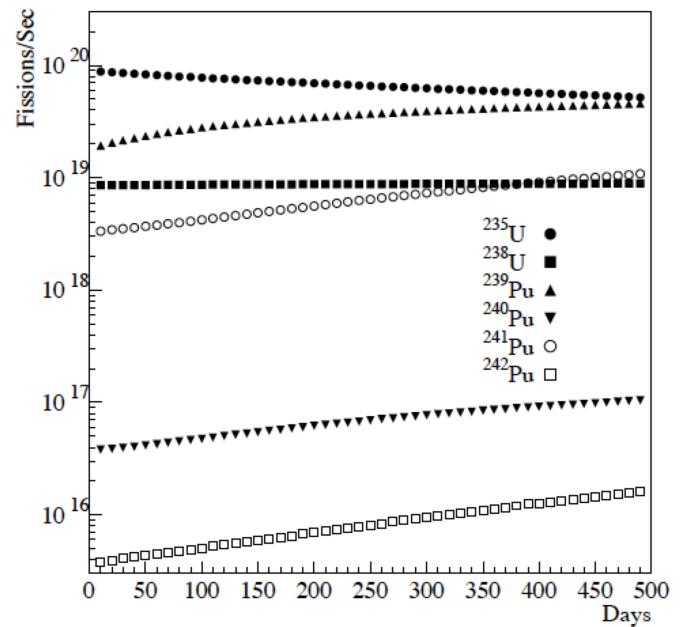
NUMBER OF FISSIONS:

- Fissile isotopes contribution
 - ^{235}U , ^{239}Pu , ^{241}Pu and ^{238}U .
 - The uncertainty on these parameters is small.
- Thermal power
 - Monitored in the heat balance of the steam generator
 - The $\delta P_{\text{th}}/P_{\text{th}}$ is small.

NEUTRINO YIELD:

- Electron antineutrino per MeV per fission
 - Conversion to the ILL reference spectrum
 - Drives the main uncertainty.
- Prediction: $\sigma_f = 5.824 \cdot 10^{-19} \text{ b/fission} \pm 2.7 \%$
- Bugey-4: $\sigma_f = 5.752 \cdot 10^{-19} \text{ b/fission} \pm 1.4 \%$
(Phys. Lett. B338 (1994) 383-389)

C. Bemporad, G.Gratta and P. Vogel
Rev. Mod. Phys. 74 (2002) 297



$$\sigma_f = \int_0^\infty S(E_\nu) \sigma_{V-A}(E_\nu) dE_\nu$$

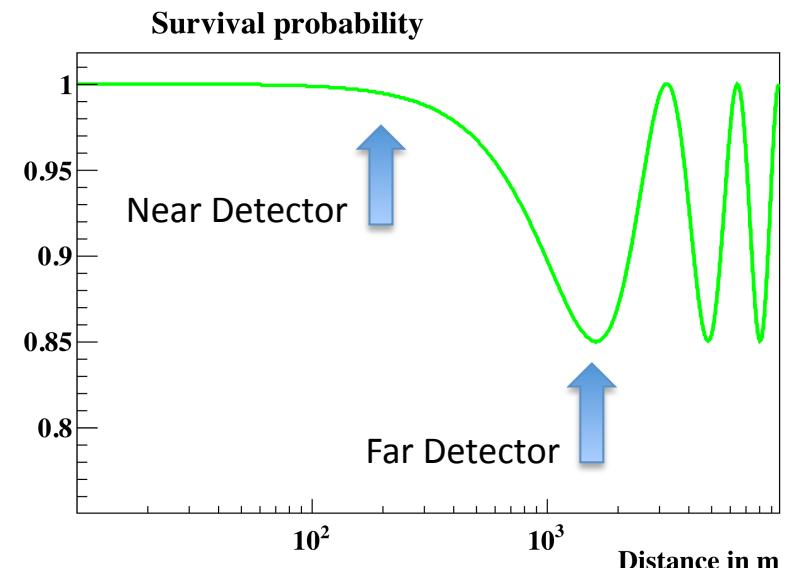
$$n = \frac{1}{4\pi R^2} \frac{P_{\text{th}}}{\langle E_f \rangle} N_p \sigma_f \epsilon$$

The novel two detector concept

$$p(\nu_e \rightarrow \nu_e; L) \simeq 1 - \sin^2(2\theta_{13}) \sin^2 \left(1.27 \Delta m_{atm}^2 \frac{L}{E} \right)$$

$$\frac{n_f}{n_n} = \left(\frac{L_n}{L_f} \right)^2 \left(\frac{N_{p,f}}{N_{p,n}} \right) \left(\frac{\epsilon_f}{\epsilon_n} \right) \left[\frac{p(E, L_f, \theta_{13})}{p(E, L_n, \theta_{13})} \right]$$

$\langle E_\nu \rangle = 3.0 \text{ MeV}$
 $\Delta m_{13}^2 = 2.30 \cdot 10^{-3} \text{ eV}^2$ and $\sin^2(2\theta_{13}) = 0.160$



- In the 2000's it was realized that a two, or more, detector disappearance experiment can further improve the sensitivity on θ_{13} by an order of magnitude.
- Reduce the uncertainties in both the production and detection of electron antineutrinos.
- “*A new nuclear reactor experiment to measure ϑ_{13}* ”, [arxiv.org/0402041](https://arxiv.org/abs/0402041).
- The near detector will monitor the reactor flux and spectrum.
- Identical detectors, as much as possible.

The Double Chooz collaboration

Brazil: CBPF, UNICAMP, UFABC

France: APC Paris, CEA/Irfu Saclay, Subatech Nantes, IPHC
Strasbourg

Germany: EKU Tübingen, MPIK Heidelberg, RWTH Aachen
TU München, U. Hamburg

Japan: Tohoku U., Tokyo Inst. Tech., Tokyo Metro. U., Niigata U.,
Kobe U., Tohoku Gakuin U., Hiroshima Inst Tech.

Russia: INR RAS, IPC RAS, RRC Kurchatov

Spain: CIEMAT - Madrid

USA: U. Alabama, ANL, U. Chicago, Columbia U., UC Davis,
Drexel U., IIT, KSU, LLNL, MIT, U. Notre Dame, U. Tennessee

Spokesperson: H. de Kerret (CNRS/IN2P3 - APC)

Project manager: C. Veyssiére (CEA/Irfu - Saclay)



Double Chooz: A Search for the Neutrino Mixing Angle θ_{13}

F. Ardellier ¹⁹	I. Barabanov ⁴⁰	J. C. Barriére ¹⁹	F. Beibel ¹
S. Berridge ²³	L. Bezrukov ¹⁰	A. Bernstein ¹⁴	T. Bolton ¹²
N.S. Bowden ²⁰	Ch. Buck ¹⁶	B. Bugg ²³	J. Busenitz ²
E. Caden ⁶	C. Cattadori ^{7,17}	S. Cazaux ¹⁹	M. Cerrada ¹
H. Cohn ²³	J. Coleman ¹⁵	S. Cormon ²¹	B. Court ¹
M. Cribier ^{4,19}	N. Danilov ¹¹	S. Dazeley ¹⁵	A. Di Vacri ⁷
Y. Efremenko ²³	A. Etenko ¹³	M. Fallot ²¹	V. Hernandez-Bedoya ⁵
F. von Feilitzsch ²²	Y. Foucher ²¹	T. Gabaldon ¹	P. Ghislain ⁴
I. Gil Botella ⁵	G. Giurgiu ³	M. Goeger-Beck ^{4,17}	M. Goodman ^{3*}
D. Greiner ²⁴	Ch. Grieb ²²	V. Guarino ³	A. Guertin ²¹
C. Hagner ⁸	W. Hampel ¹⁶	T. Jandler ¹	F. X. Hartmann ¹⁶
G. Horton-Smith ¹²	P. Huber ^{22†}	J. Ichum ²⁴	Y. Kamyshev ²³
D. M. Kaplan ⁹	H. de Kerret ⁴	T. Kirchner ²¹	V. Kopeikin ¹³
J. Kopp ²²	A. Kozlov ²³	T. Kotter ⁴	Yu. S. Krylov ¹¹
T. Lachenmaier ²⁴	C. Lane ⁶	T. Lasserre ^{4,19}	D. Kryn ⁴
A. Letourneau ¹⁹	D. Lhuillier ¹⁹	C. Lendvai ²²	Y. Liu ²
I. Machulin ¹³	F. Marie ¹¹	M. Lindner ²²	J. LoSecco ¹⁸
F. Meignier ¹⁹	G. Martino ¹⁹	J. Martino ²¹	D. McKee ²
A. Milsztajn ¹⁹	J. P. Meyer ¹⁹	W. Metcalf ¹⁵	R. McNeil ¹⁵
M. Obolensky ¹	C. Palomares ⁵	D. Motta ¹⁹	L. Mikaelyan ¹³
J. Reicher-Bacher ³	B. Reinhold ¹	P. Perrin ¹⁹	L. Oberauer ²²
L. Romeo ¹²	S. Roth ¹	D. Reyna ³	W. Potzel ²²
L. Scola ¹⁹	V. Sinev ^{13,19}	M. Skorokhvatov ¹³	M. Rolinec ²²
N. Stanton ¹²	S. Sukhotin ^{4,13}	U. Schwan ¹⁶	T. Schwetz ²²
A. Tonazzo ⁴	D. Underwood ³	R. Svoboda ^{14,15}	A. Tang ¹²
D. Vincent ⁴	W. Winter ^{22†}	F.J. Valdivia ⁵	D. Vignaud ⁴
		K. Zbiri ²¹	R. Zimmermann ⁸

20 June 2006 (revised 26 October 2006)

The site

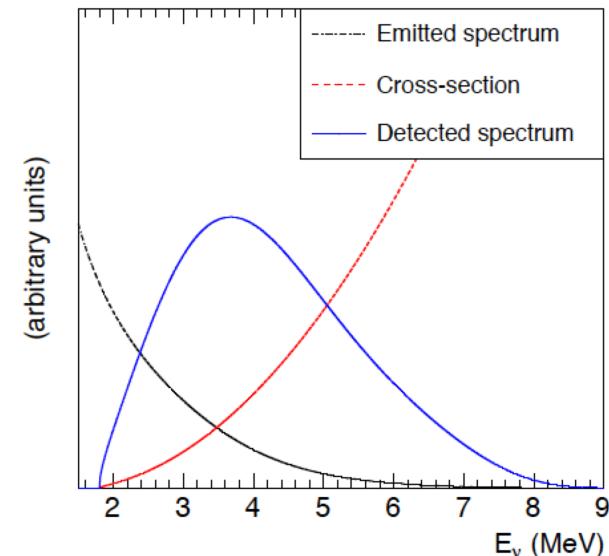
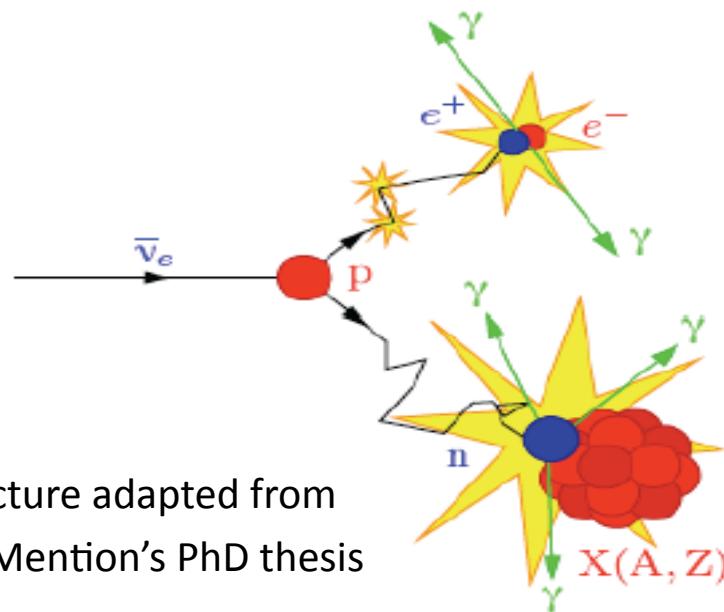
- The Chooz Nuclear power plant is located at the North-East part of France, near the border with Belgium.
- Two nuclear reactors $P_{th}=4.27$ GW and $P_e=1.5$ GW each.
- The Chooz-far laboratory was constructed by the EDF during the 1990's for the CHOOZ experiment.
- 300 mwe shielding from cosmics of 2.8 gr/cm^3 rock.
- The construction of near laboratory started at the 29th of April 2011.
- Artificial overburden of 120 mwe and flat topology.
- According to schedule the near detector would be ready to operate at 2013.
- Currently only one detector is taking data in stable conditions since April 2011.



Neutrino detection

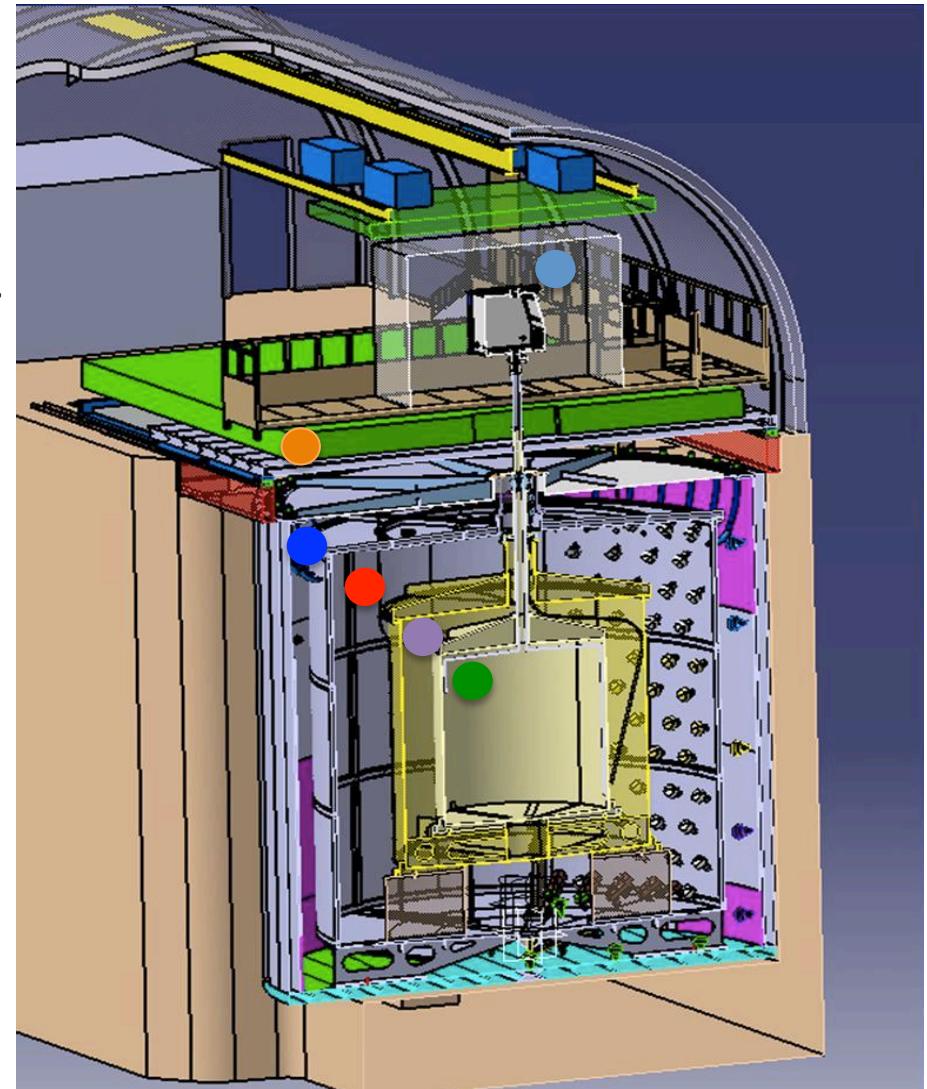
- Antineutrinos are detected through inverse beta decay (IBD) on the free protons of the scintillator, $\bar{\nu}_e + p \longrightarrow n + e^+$
- High cross section but threshold at 1.8 MeV.
- Prompt signal: e^+ thermalization and the subsequent annihilation.
- Delayed signal: n-capture on Gd; a number of gammas around 8 MeV.
- Time coincidence with $\tau_{\text{mean}} = 30 \mu\text{s}$ and,
- Strong space correlation ($< 1 \text{ m}$).

T. Meuller et al.,
Phys. Rev. C 83, 054615 (2011)



The Double Chooz Detector

- Four concentric cylindrical detector volumes.
- A four-layer russian doll.
- **Target volume** : 10.3m^3 of liquid scintillator
(PXE and dodecane, 20:80) doped with 0.1% Gd.
- **Gamma catcher** : 22.3m^3 of liquid scintillator.
- **Buffer**: 110 m^3 of mineral oil.
- 390 R7081 Hamamatsu PMT mounted
in the Buffer vessel (10 inches diameter).
- Photocathode coverage $\approx 13\%$.
- **Inner Veto** : 90 m^3 of scintillator.
- 78 8-inched PMT in the Inner Veto
(R1408 Hamamatsu)
- Steel shielding above the detector.
- **Outer Veto** : plastic scintillator strips.
- **Calibration box**.





Muon Veto:

- 1 ms away the muon signature.
- 0.5 s dead-time after a $E_{ID} > 600$ MeV muon
- No OV coincidence in the prompt

Prompt signal:

- $0.7 \text{ MeV} < E_p < 12.2 \text{ MeV}$
- $Q_{\max}/Q_{\text{tot}} < 0.09$ and $\text{RMS}(T_s) < 40 \text{ ns}$

Delayed signal:

- $0.7 \text{ MeV} < E_d < 12.0 \text{ MeV}$
- $Q_{\max}/Q_{\text{tot}} < 0.055$ and $\text{RMS}(T_s) < 40 \text{ ns}$

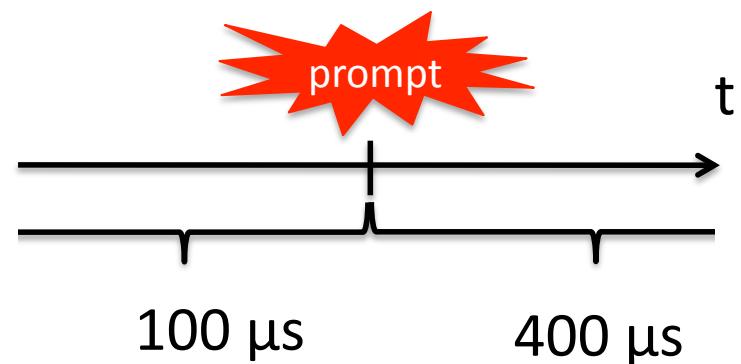
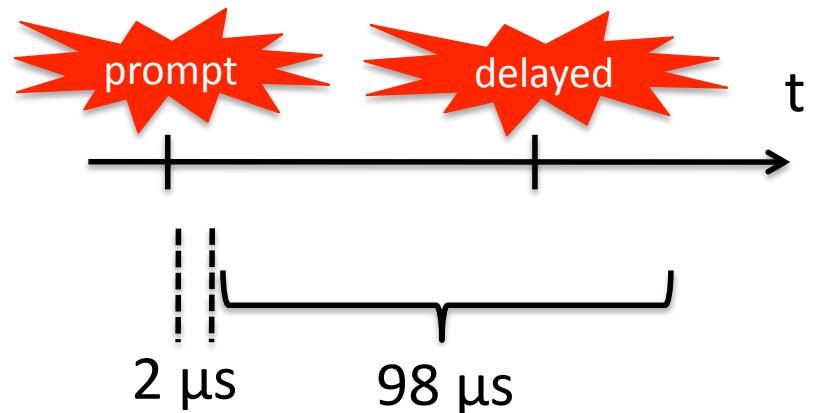
Coincidence window:

- $[2, 100] \mu\text{s}$

Multiplicity isolation cuts:

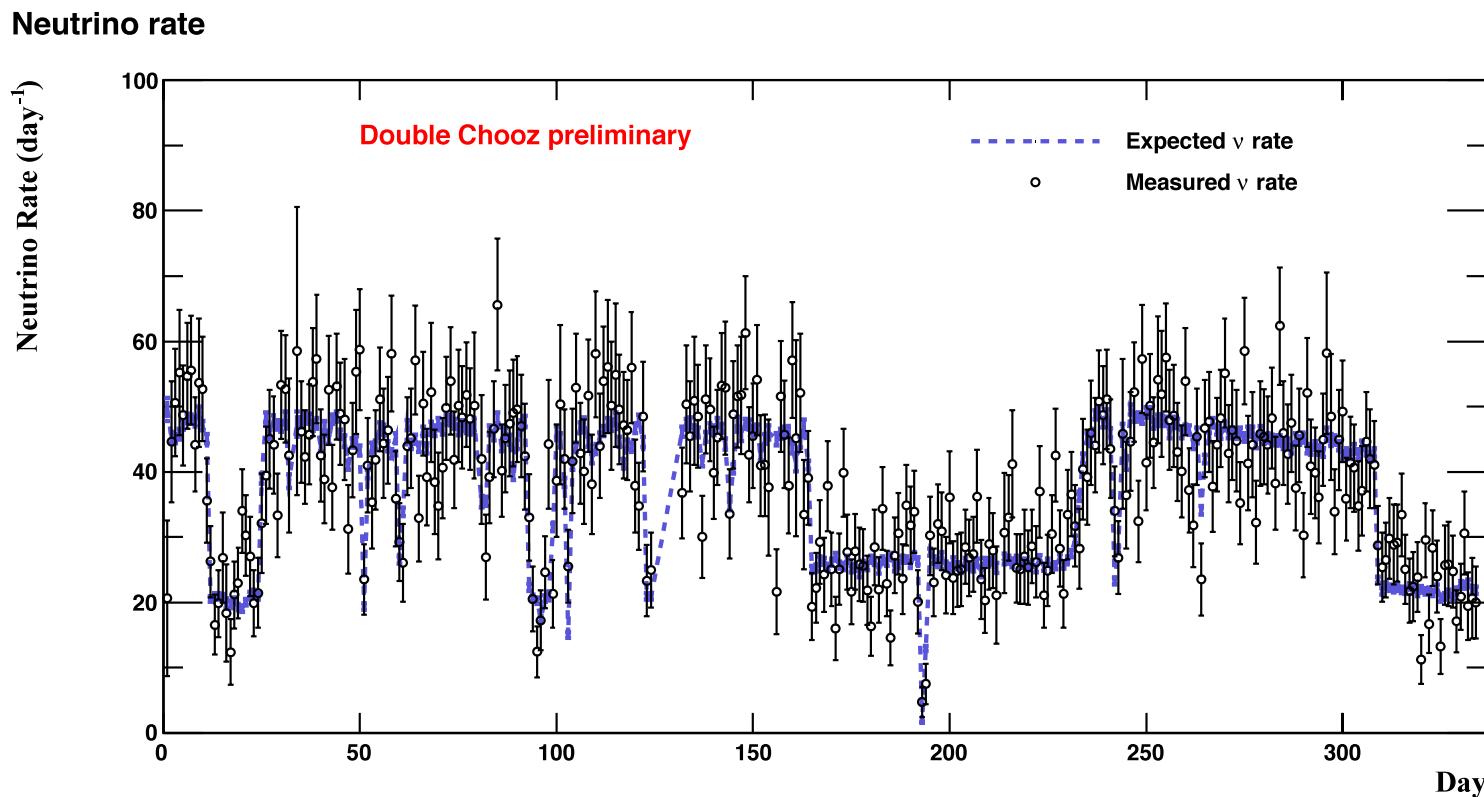
- Valid trigger: $E > 0.5 \text{ MeV}$,
 $Q_{\max}/Q_{\text{tot}} < 0.09$ and $\text{RMS}(T_s) < 40 \text{ ns}$.
- No valid trigger $100 \mu\text{s}$ before the prompt
- No valid trigger $400 \mu\text{s}$ after the prompt

Neutrino selection cuts



Neutrino candidates

8249 neutrino candidates

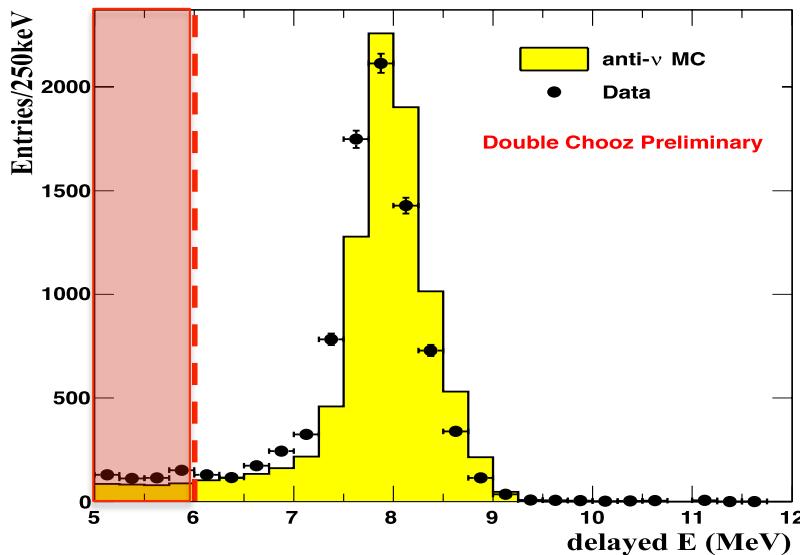


Attention, in this plot:

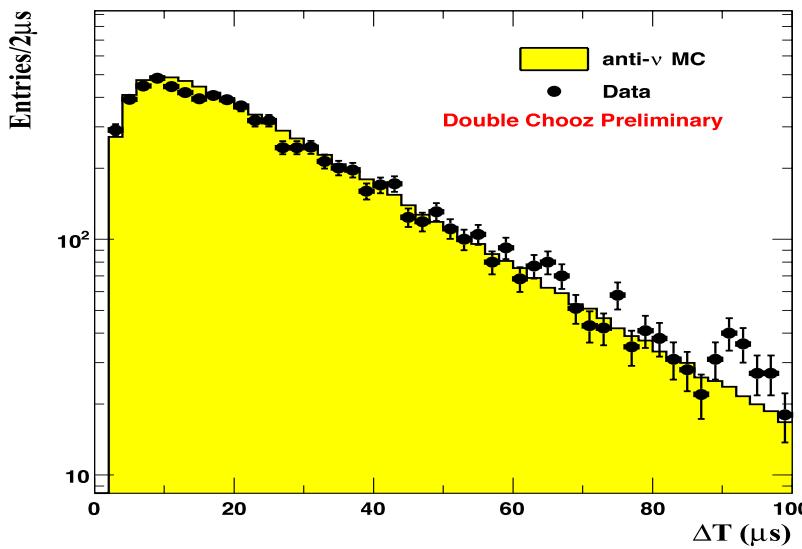
- No high energy muon veto
- No Outer Veto veto
- No background subtraction

Neutrino imprint

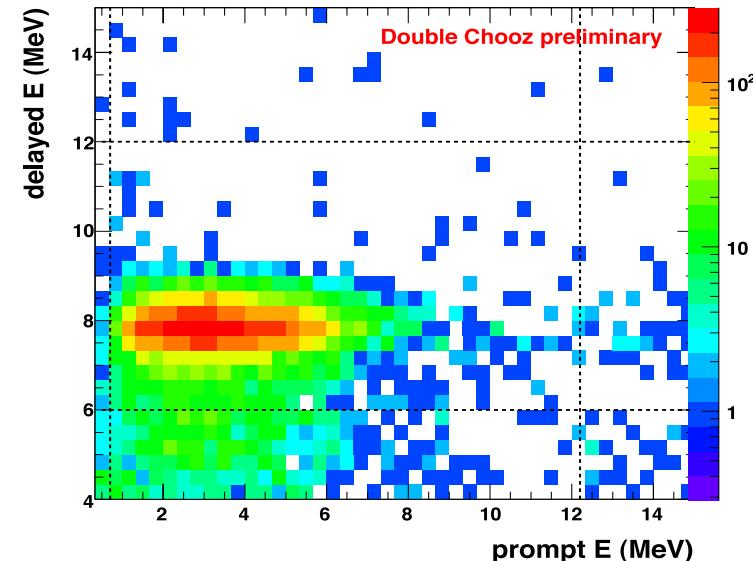
Delayed energy



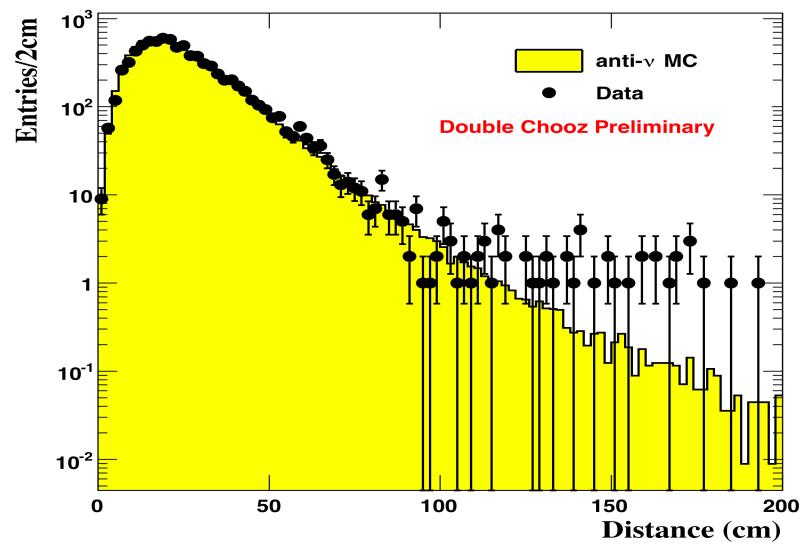
Prompt-delay time difference



E_{prompt} vs E_{delayed}

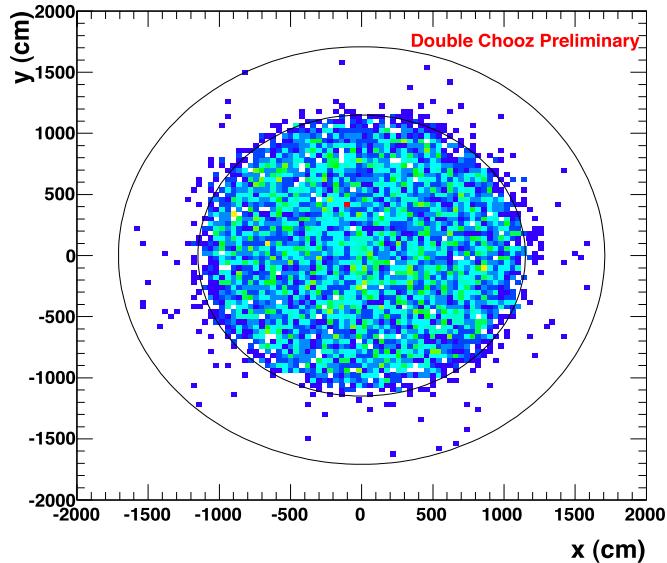


Distance

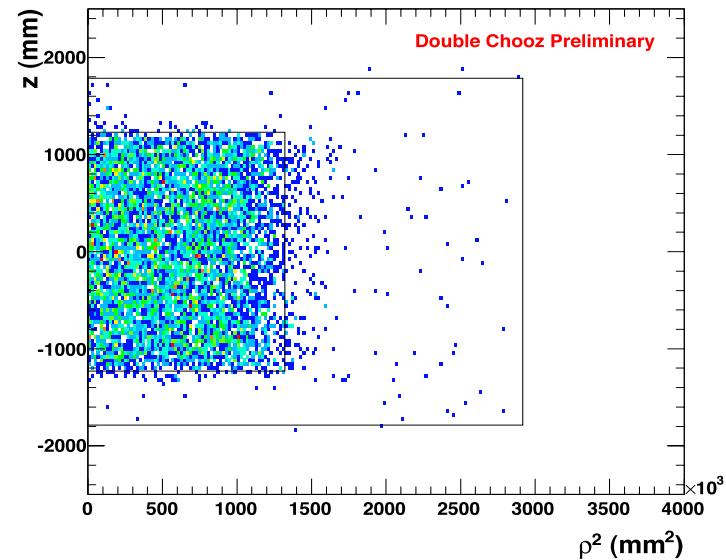


Reconstructed vertices

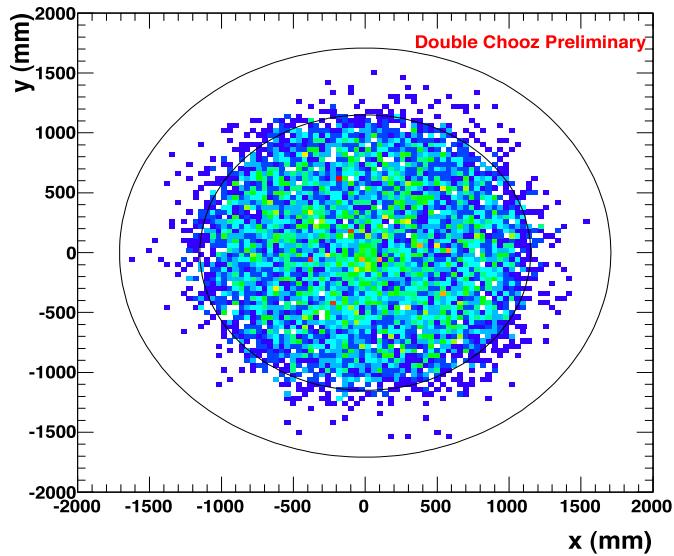
Prompt vertex XY position



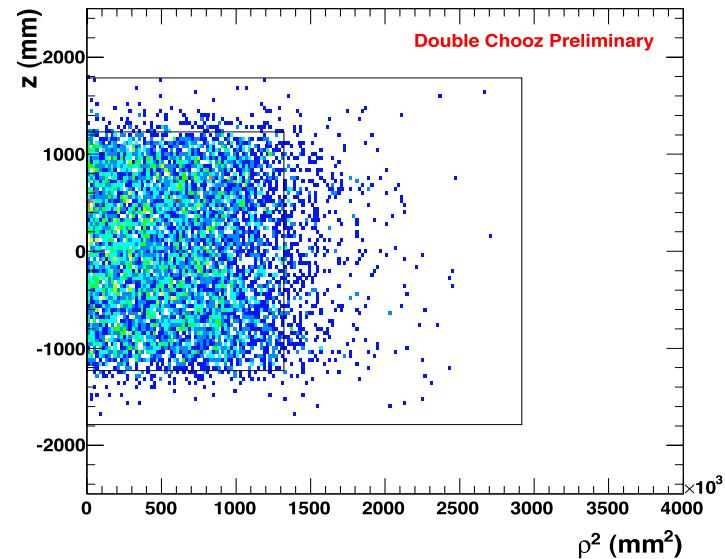
Prompt vertex Z ρ^2 position



Delayed vertex XY position

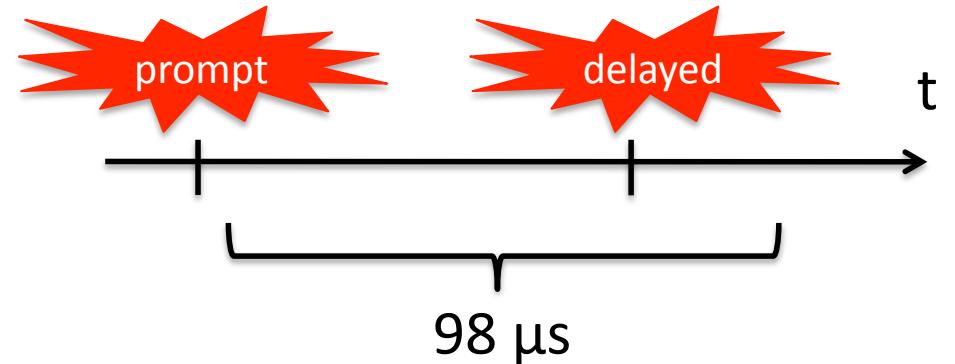


Delayed vertex Z ρ^2 position



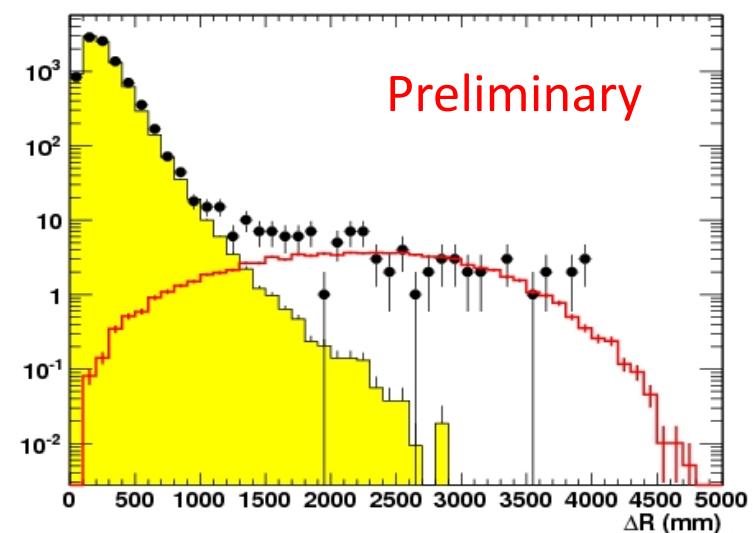
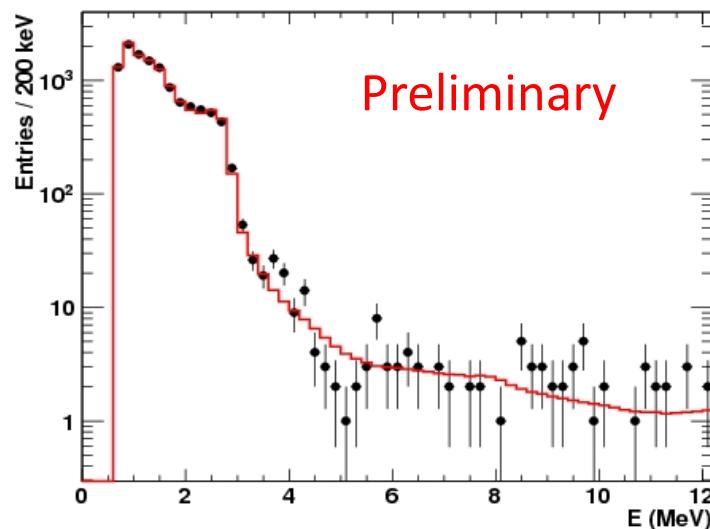
Accidentals

- Uncorrelated events.
- Prompt and delayed signals: natural radioactivity, mu-untagged spallation neutrons, etc ...



Selection cuts:

- The same criteria used already for neutrino selection but,
- Altering the coincidence window: [1s+500*i+2μs, 1s+500*i+100μs], $i = 0, \dots, 197$
- Rate: 0.261 ± 0.002 events per day (Stable).



Fast Neutrons and Stopping Muons

Fast neutrons

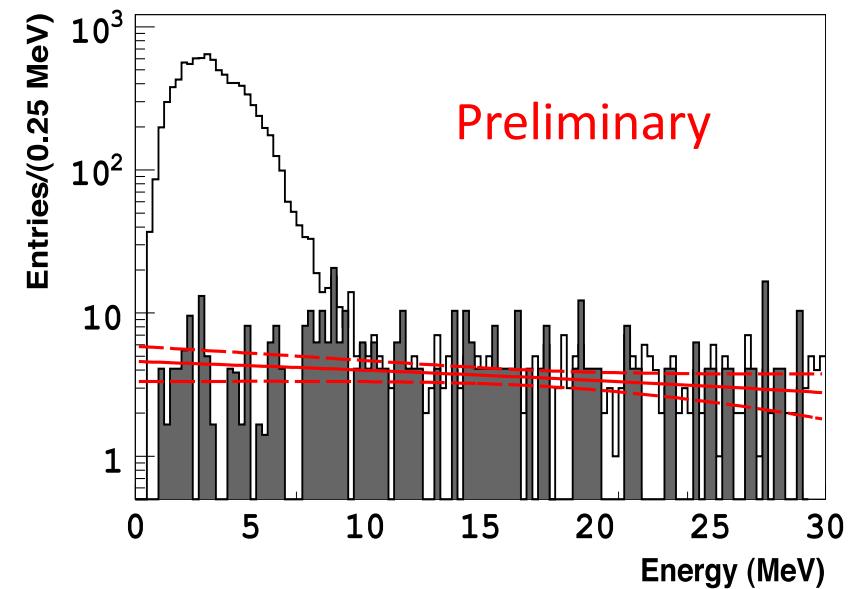
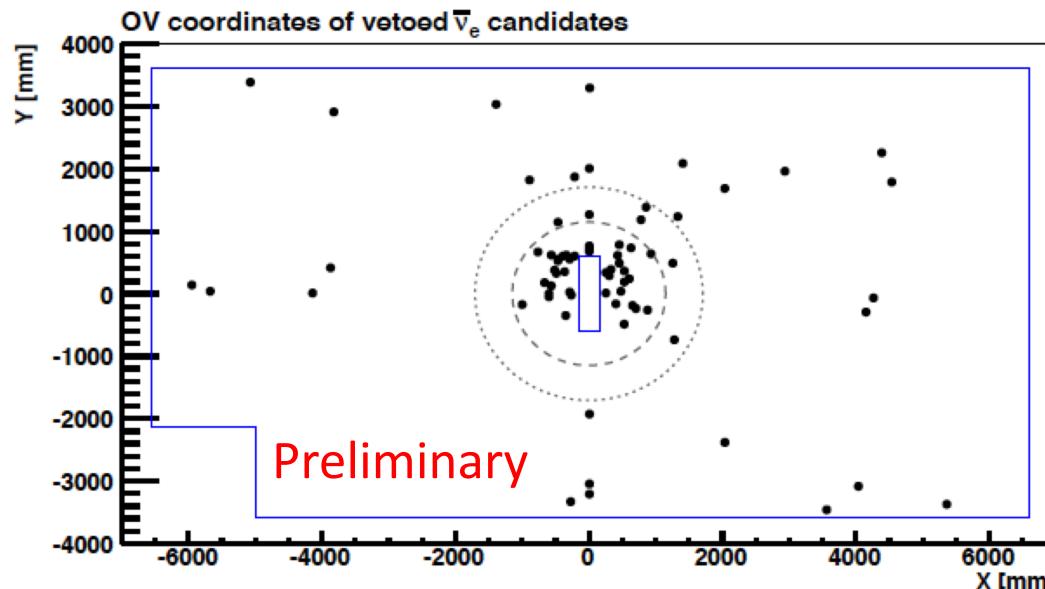
- Prompt signal: proton recoil, and delayed: normal neutron capture

Stopping muons

- Muon michel decay: $\mu \longrightarrow e + \nu_\mu + \bar{\nu}_e, \tau \approx 2.2 \mu\text{s}$

Selection cuts:

- The same criteria used already for neutrino selection but,
- Altering the prompt energy window to [0.7, 30.0] MeV.
- Rate: 0.67 ± 0.20 event/day

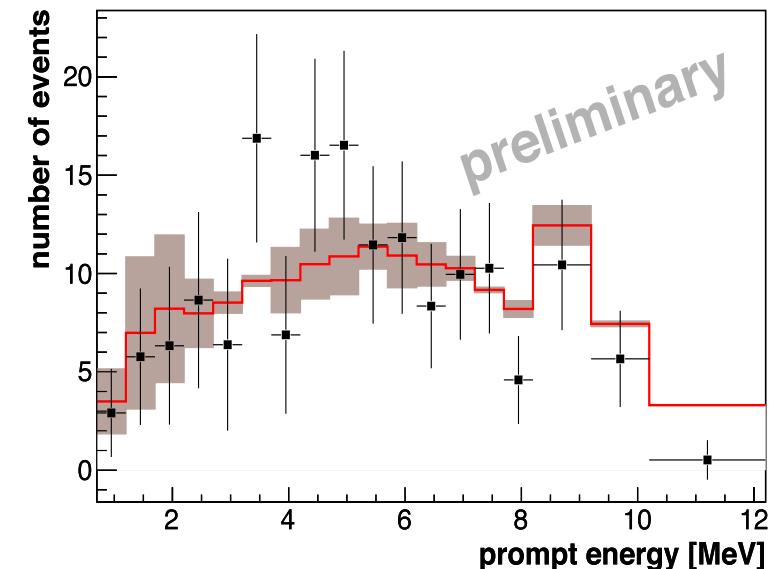
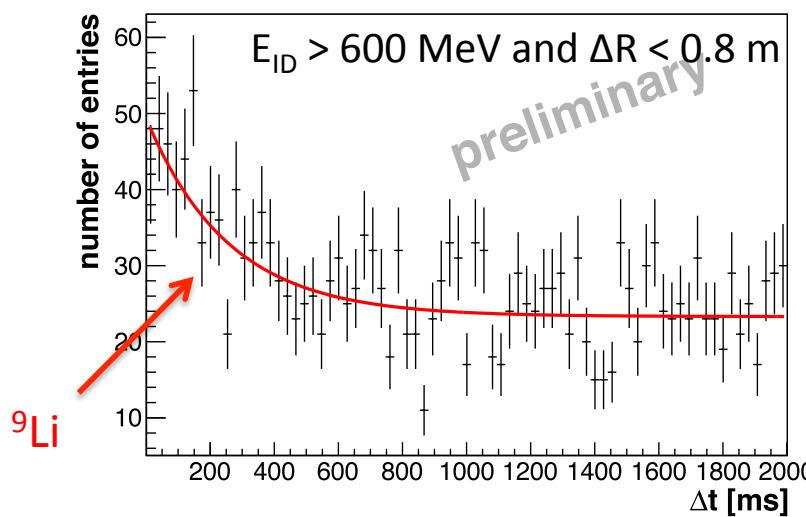


Cosmogenic isotopes

- High energy muons are able to create long-lived isotopes from spallation processes ([arxiv.org/0907.0066](https://arxiv.org/abs/0907.0066)).
- Large life-times; can not be tagged by the muon veto.
- ${}^8\text{He}$ and ${}^9\text{Li}$ have electron and neutron decay cascade modes that mimic the IBD signature.

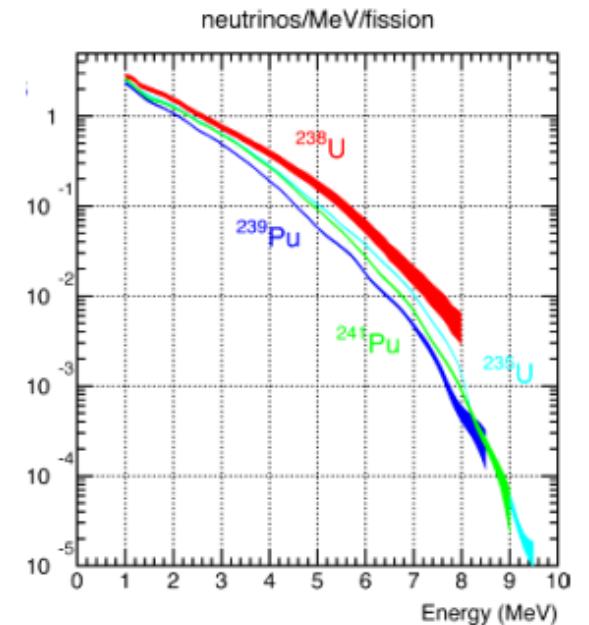
Estimation:

- $E_{\text{ID}} > 600 \text{ MeV}$ rejected, estimate the residual ${}^9\text{Li}$ contribution.
- Separately in the $275 \text{ MeV} < E_{\text{ID}} < 600 \text{ MeV}$ and $E_{\text{ID}} < 275 \text{ MeV}$ regions
- Fits to the DT distributions
- Rate: $1.20+0.59/-0.48$ évents per day



Reactor input

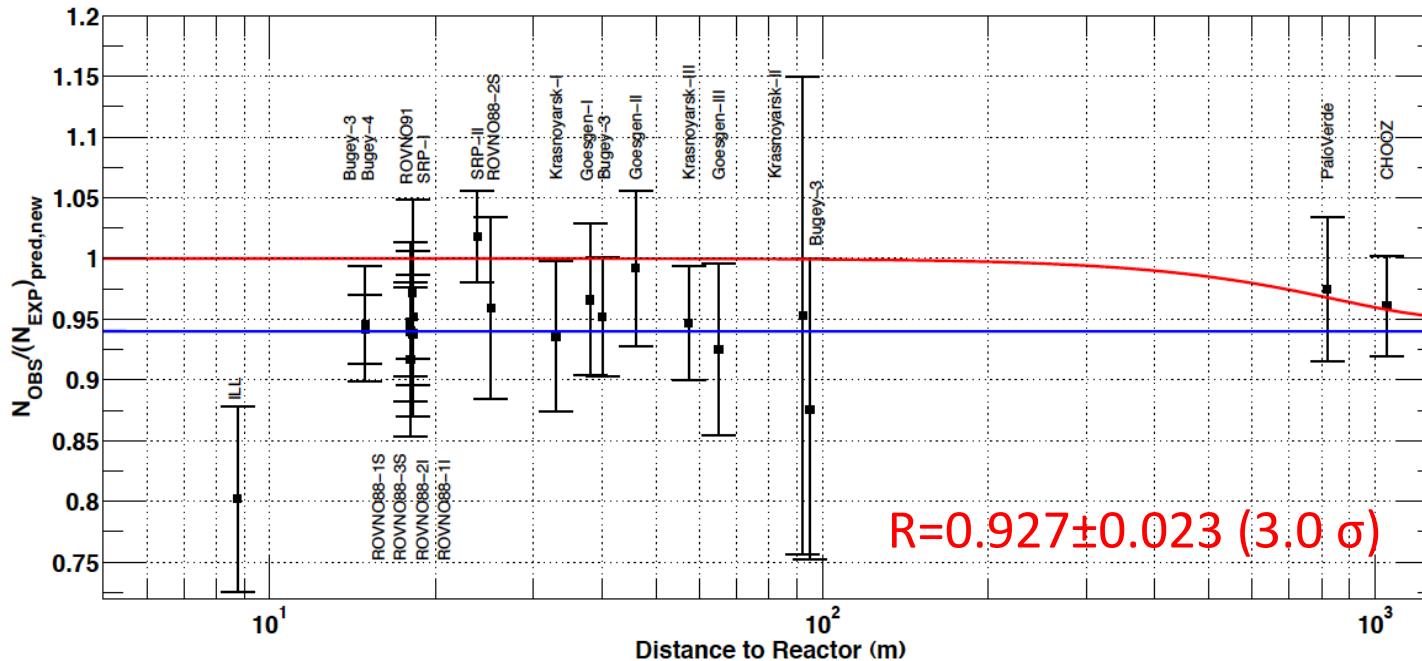
- Thermal power monitored by EDF from the balance at the steam generator.
- For the reactor evolution two simulation packages were used: the MURE simulation package and the DRAGON code, C. Jones et al. arxiv.org/1109.5379
- Lots of information from EDF:
 - Geometrical parameters
 - Fuel composition
 - Boron concentration, etc...
- JEFF3.1 nuclear database was used.
- Fissions per energy were taken from V. Kopeikin et al., hep-ph/0410100v1.
- Neutrino yield: a new conversion method to the ILL reference spectrum developed by Saclay and Subatech.
- Net $\approx 3.0\%$ upward shift in energy-averaged neutrino fluxes, Meuller et al., Phys. Rev. C 83, 054615 (2011)
- Confirmed by P. Huber, Phys. Rev. C 84, 024617 (2011).
- IBD cross-section also increased, neutron lifetime.



Additional information:
«Reactor Flux Calculations»
D. Lhuillier, Neutrino 2012

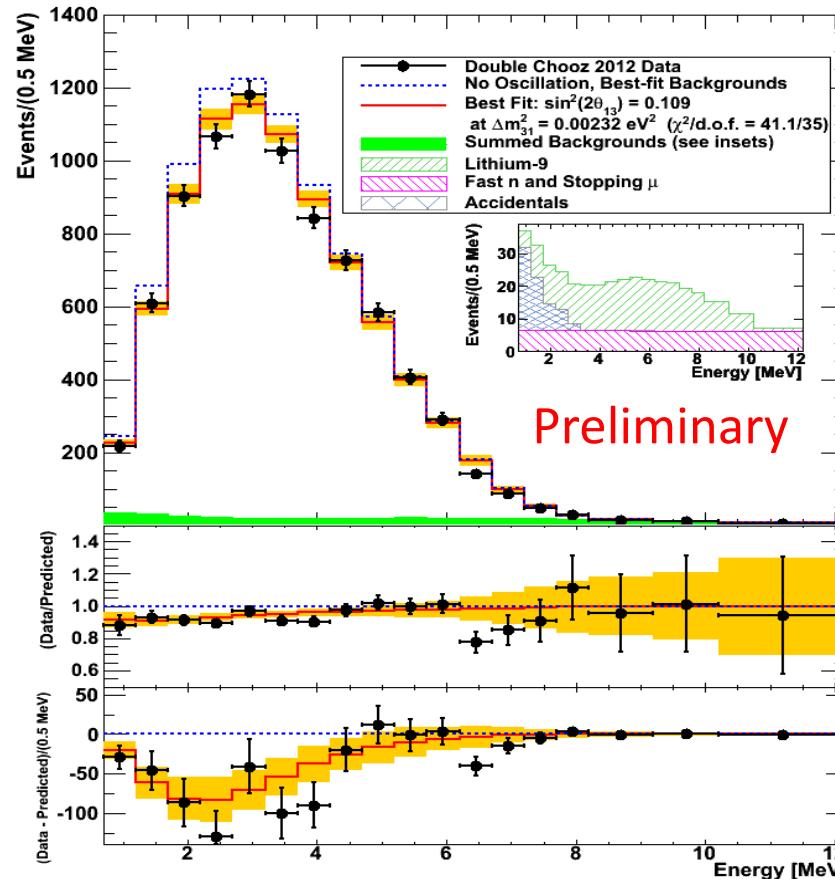
Reactor anomaly

(G. Mention et al. arxiv.org/1101.2755)



- A common shift to all exp. with baseline smaller than 100 m.
- Possible explanations:
 - Wrong estimation of the antineutrino spectrum
 - A common bias in all experiments!!!
 - Hint for new physics... a forth sterile oscillation mode?
- To avoid oscillations from possible sterile modes Double Chooz uses the absolute normalization of Bugey-4 correcting for the burn-up differences.
- Further info: P. Huber et al., arxiv.org/1204.5379 and Th. Lasserre's talk at Neutrino 2012

Oscillation analysis



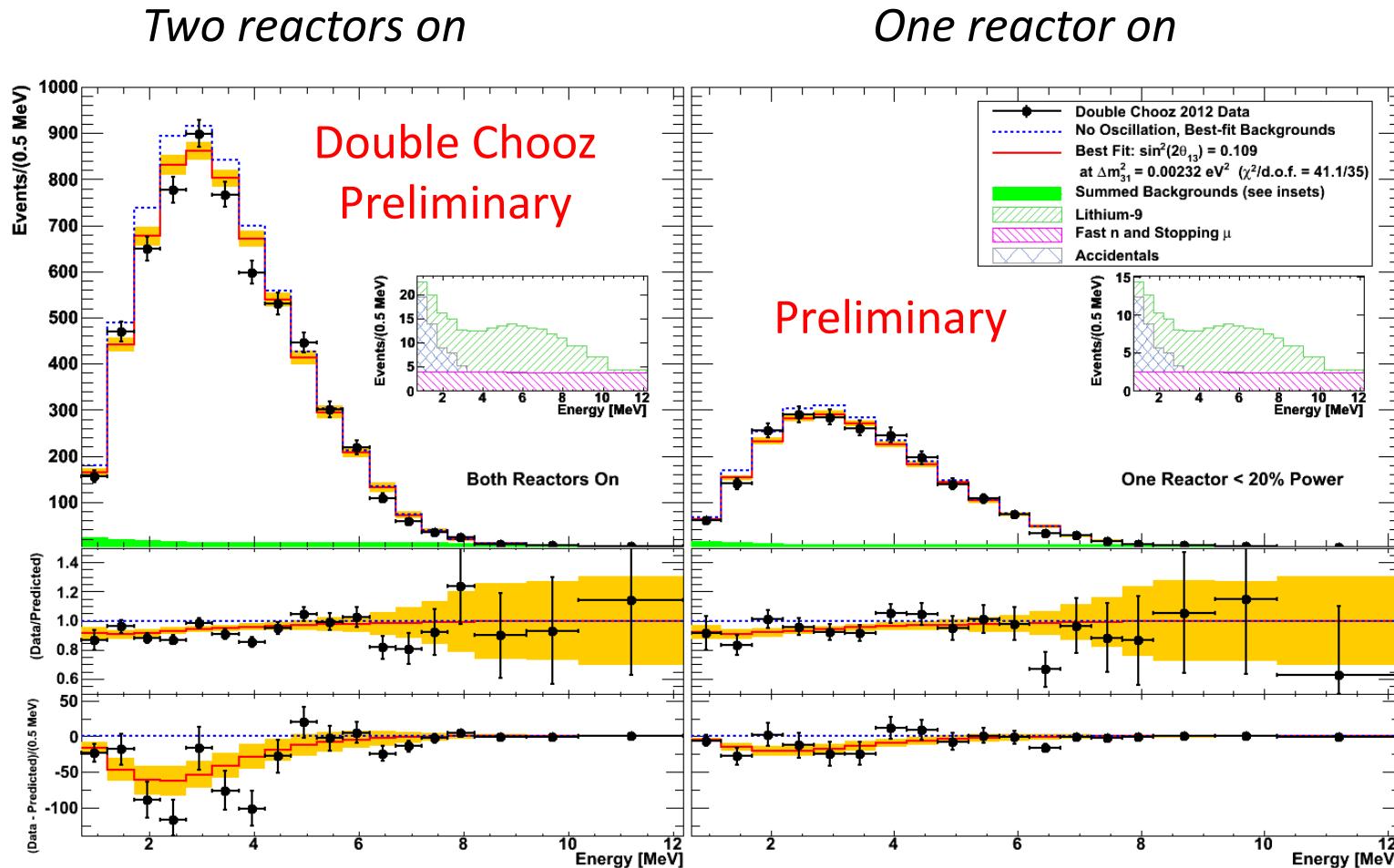
NOTE:

In a fit, dataset is divided into two periods based on reactor power (next slide)

Rate + Shape Fit: $\sin^2 2\theta_{13} = 0.109 \pm 0.030(\text{stat.}) \pm 0.025(\text{syst.})$

Rate Only Fit: $\sin^2 2\theta_{13} = 0.170 \pm 0.035(\text{stat.}) \pm 0.040(\text{syst.})$

Oscillation analysis (cont'd)



Rate + Shape Fit: $\sin^2 2\theta_{13} = 0.109 \pm 0.030(\text{stat.}) \pm 0.025(\text{syst.})$

Rate Only Fit: $\sin^2 2\theta_{13} = 0.170 \pm 0.035(\text{stat.}) \pm 0.040(\text{syst.})$

Summary

- Double Chooz has already delivered some very interesting physics:
 - A novel detector design
 - Hard radiopurity R&D and stable scintillator soluble
 - Improved reactor spectrum, “Reactor anomaly”.
 - Important oscillation analyses
- Double Chooz first result presented by H. de Kerret at LowNu 2011
(Phys. Rev. Lett. 108 (2012) 131801)
- Updated result doubling the statistics and reducing systematics:

$$\sin^2 2\theta_{13} = 0.109 \pm 0.030(\text{stat.}) \pm 0.025(\text{syst.})$$

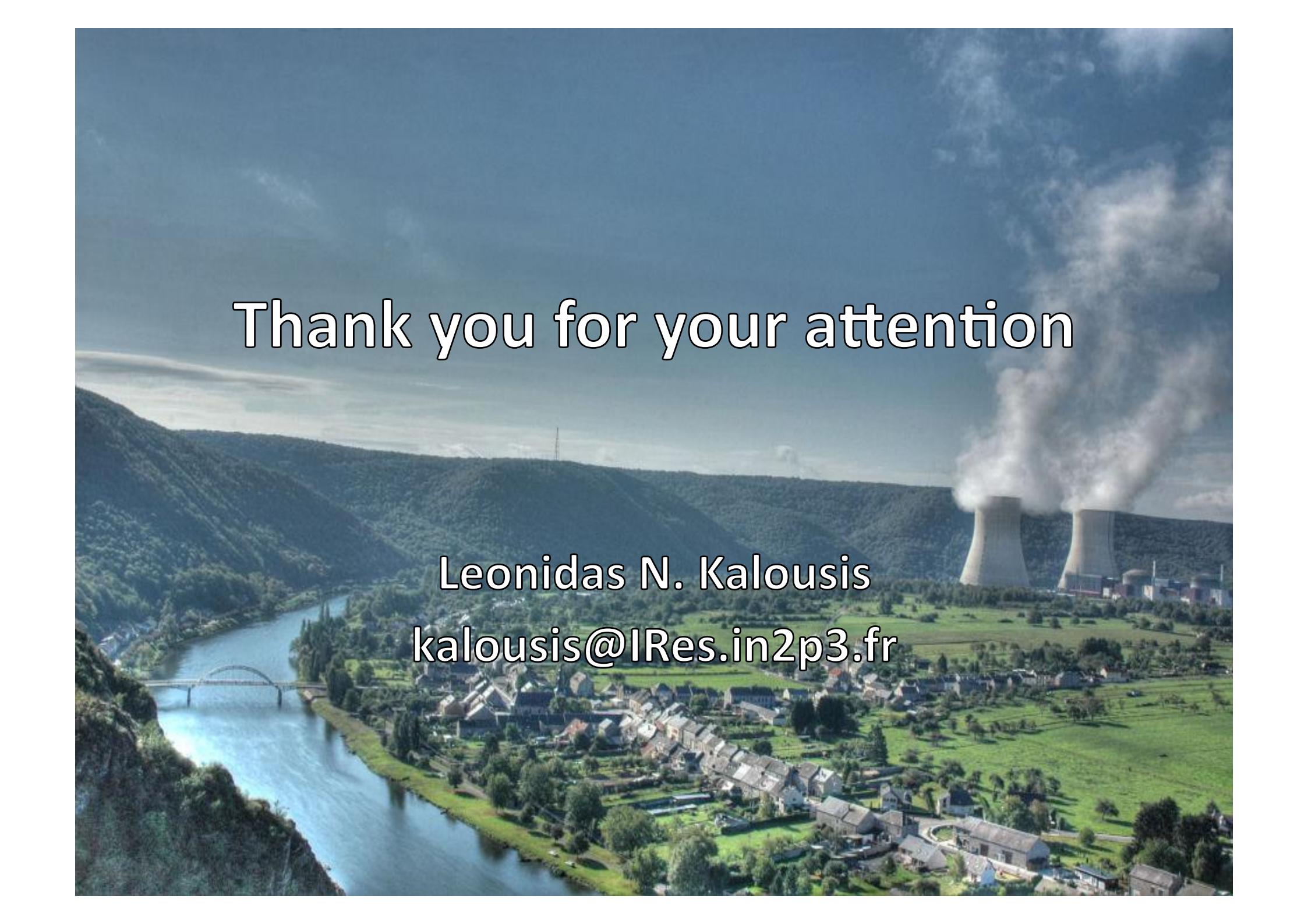
Non-oscillation scenario excluded at a 99.9 C.L., 3.1σ

RENO, Neutrino 2012

Daya Bay, Neutrino 2012

$$\sin^2 2\theta_{13} = 0.113 \pm 0.013(\text{stat.}) \pm 0.019(\text{syst.}) \quad \sin^2 2\theta_{13} = 0.089 \pm 0.010(\text{stat.}) \pm 0.005(\text{syst.})$$

- Important news: almost a week of reactor OFF-OFF data!
- The near detector is expected to start data taking late 2013.

The background image shows a panoramic view of a valley. A river flows through the center, with a bridge crossing it. On the right side, there is a large industrial complex, specifically a nuclear power plant, characterized by two tall cooling towers emitting thick plumes of white steam into the sky. The surrounding area is a mix of green fields and some forested hills. The overall scene is a blend of natural beauty and industrial activity.

Thank you for your attention

Leonidas N. Kalousis
kalousis@IRes.in2p3.fr

Additional plots and information

BACKUP SLIDES

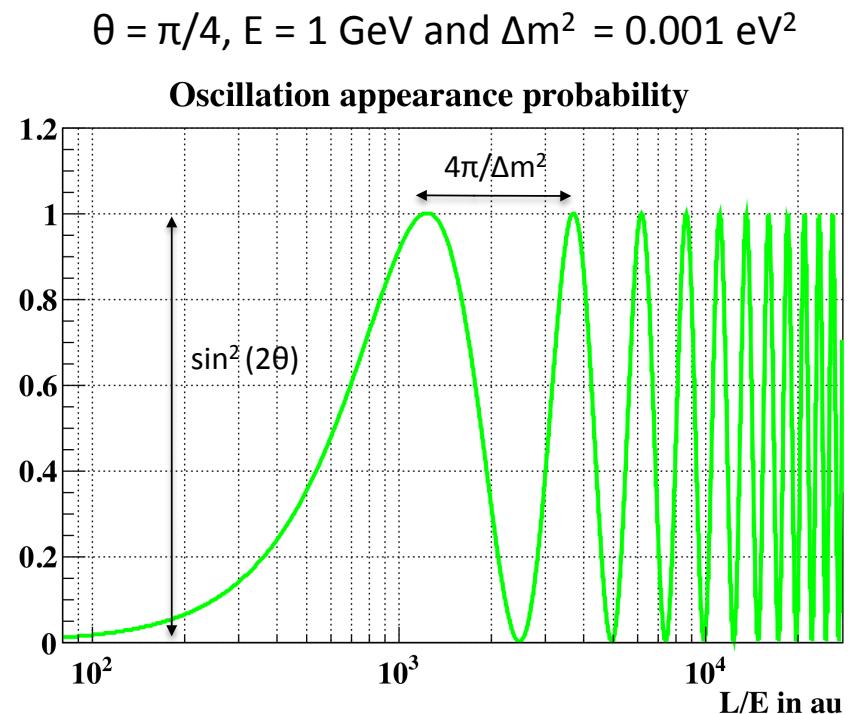
Two flavour approximation

- Three flavour are highly suppressed since $|\Delta m^2_{31}| \ll \Delta m^2_{21}$ and $\cos^2(2\theta_{13}) \approx 1.0$
- Dominant oscillations are well described by effective two-flavour oscillations.
- One mixing angle no complex phase.

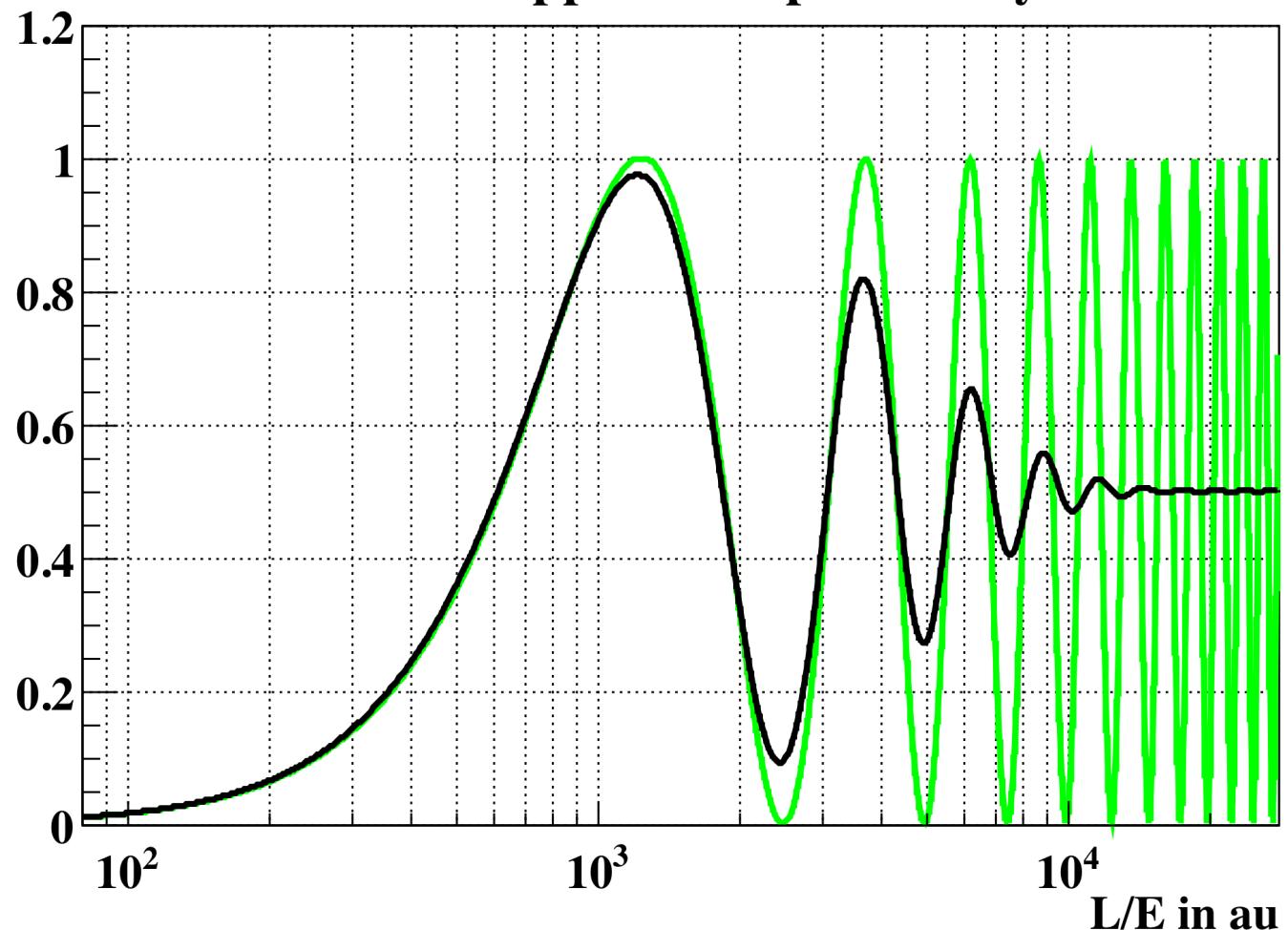
$$\mathcal{U} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}$$

$$p(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta \sin^2 \frac{\Delta m^2 L}{4E}$$

Appearance probability
Typical oscillatory behaviour



Oscillation appearance probability



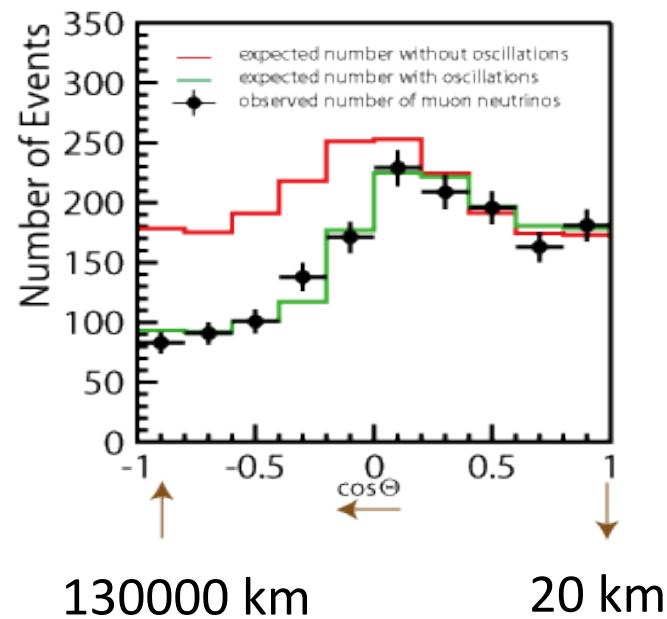
$\langle E \rangle = 1 \text{ GeV}$ and $\Delta E = 0.1 \text{ GeV}$

Oscillation patterns

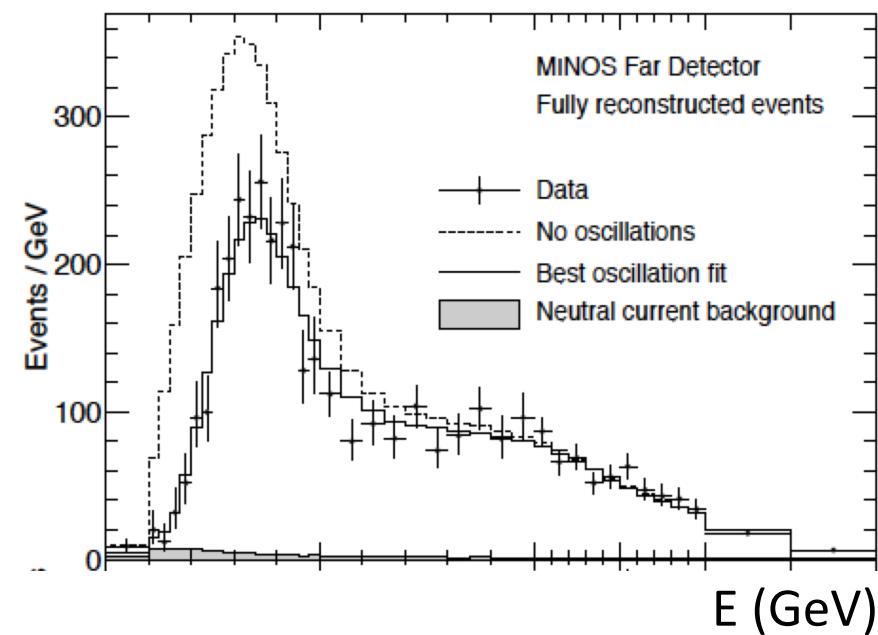
- Oscillation patterns driven by the mass square differences, $\Delta m_{ij}^2 = m_i^2 - m_j^2$
- The oscillation formula is a function of (distance)/(energy).
- Reduction in the neutrino flux and a combined energy spectrum distortion.

Super-Kamiokande

<http://www-sk.icrr.u-tokyo.ac.jp/sk/physics/atmnu-e.html>

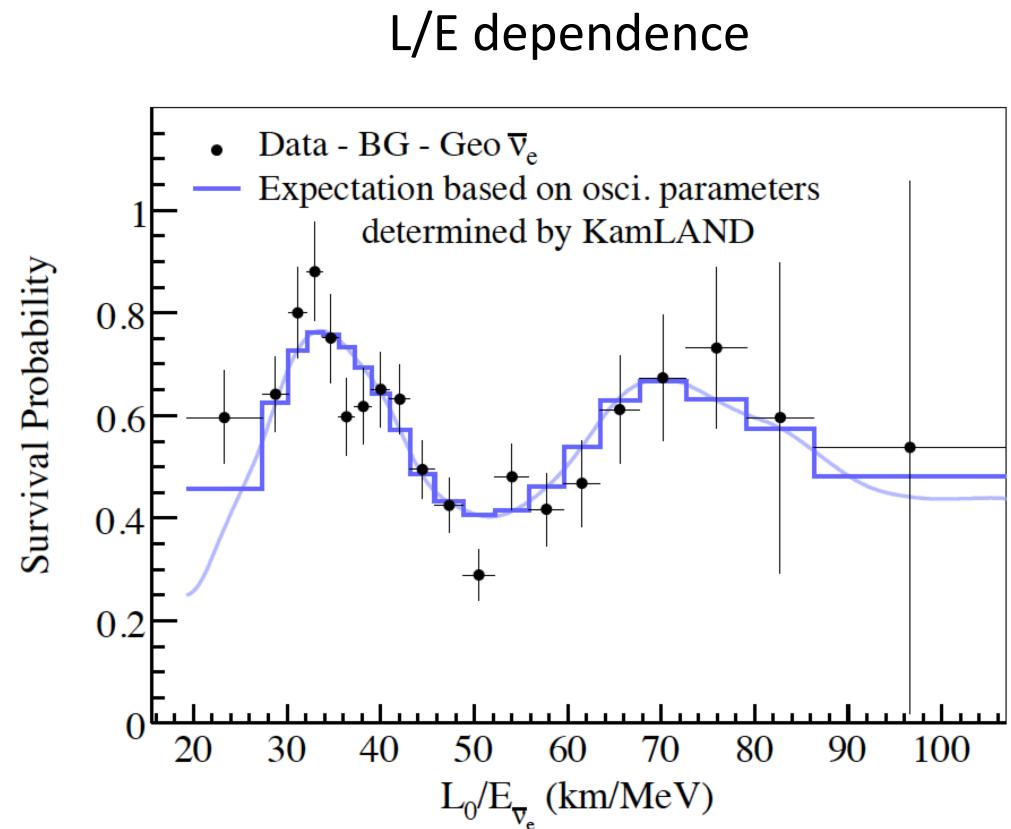
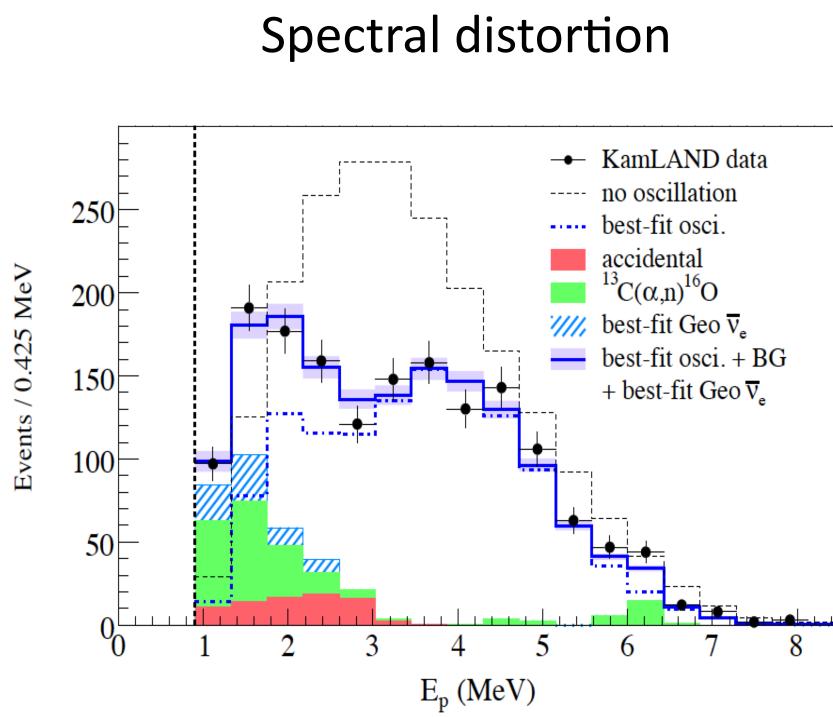


MINOS, Phys. Rev. Lett. 106, 181801 (2011)



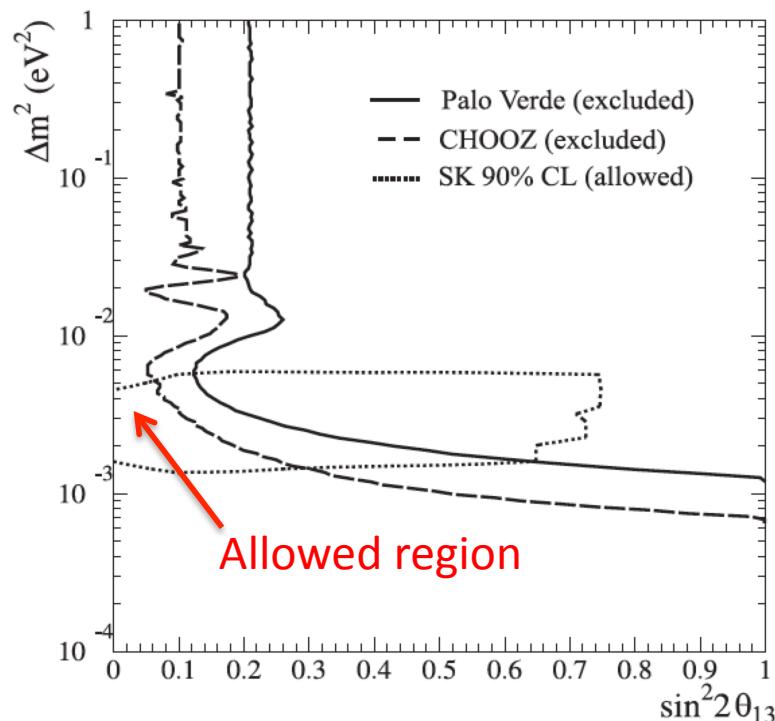
Precision measurements of solar oscillation parameters with KamLAND

Phys. Rev. Lett. 100, 221803 (2008)



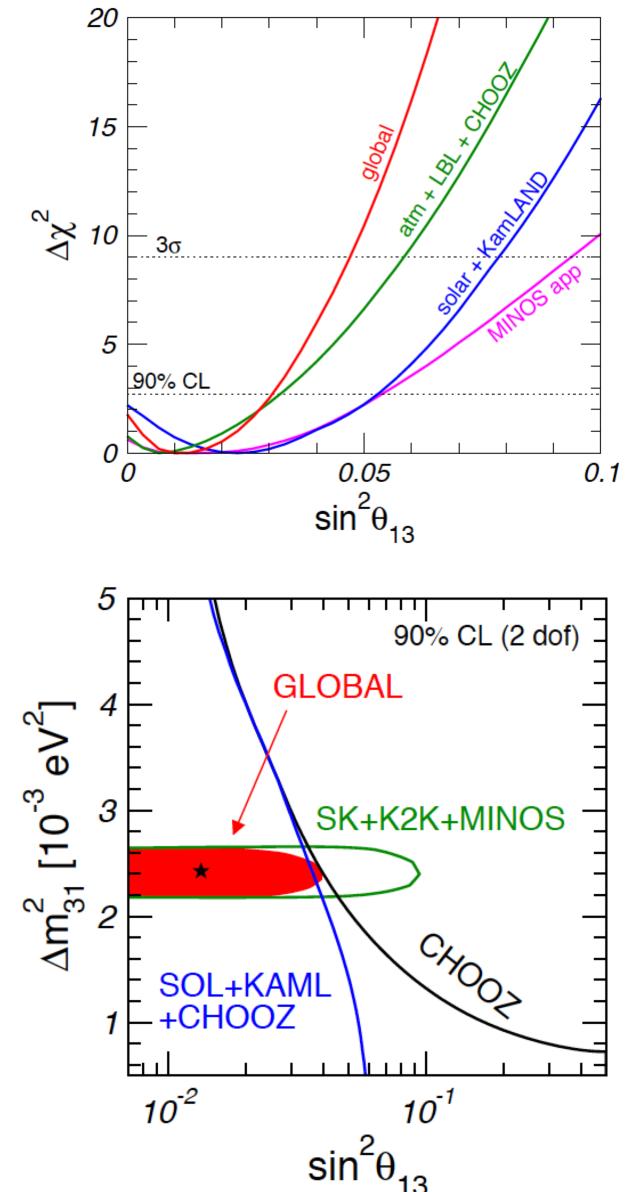
Prehistory of θ_{13}

Palo Verde, Phys. Rev. D 64, 112001 (2001)



Results from CHOOZ, Palo Verde and Super-K suggest that θ_{13} is surprisingly small!

Old CHOOZ limit, Eur. Phys. J. C27 (2003) 331



Global fits with the inclusion of Solar, Atmospheric exp., K2K and MINOS results suggest $\theta_{13} \neq 0$

MINOS first appearance result:

Phys. Rev. Lett. 103, 261802 (2009)

Ref: M. Mezzetto and T. Schwetz. arxiv.org/1003.5800

ν_e appearance

- Complicated formula:

$$\begin{aligned}
 p(\nu_\mu \rightarrow \nu_e) &\simeq \sin^2(2\theta_{13}) \sin^2(\theta_{23}) \frac{\sin^2(A - 1)\Delta}{(1 - A)^2} \\
 &+ \alpha \sin^2(2\theta_{13}) \sin^2(2\theta_{12}) \sin^2(2\theta_{13}) \cos(\Delta + \delta) \frac{\sin(A\Delta)}{A} \frac{\sin^2(A - 1)\Delta}{(1 - A)} \\
 &+ \alpha^2 \sin^2(2\theta_{12}) \cos^2(\theta_{23}) \frac{\sin^2(A\Delta)}{A^2}
 \end{aligned}$$

$$\alpha = \frac{\Delta m_{21}^2}{\Delta m_{31}^2}$$

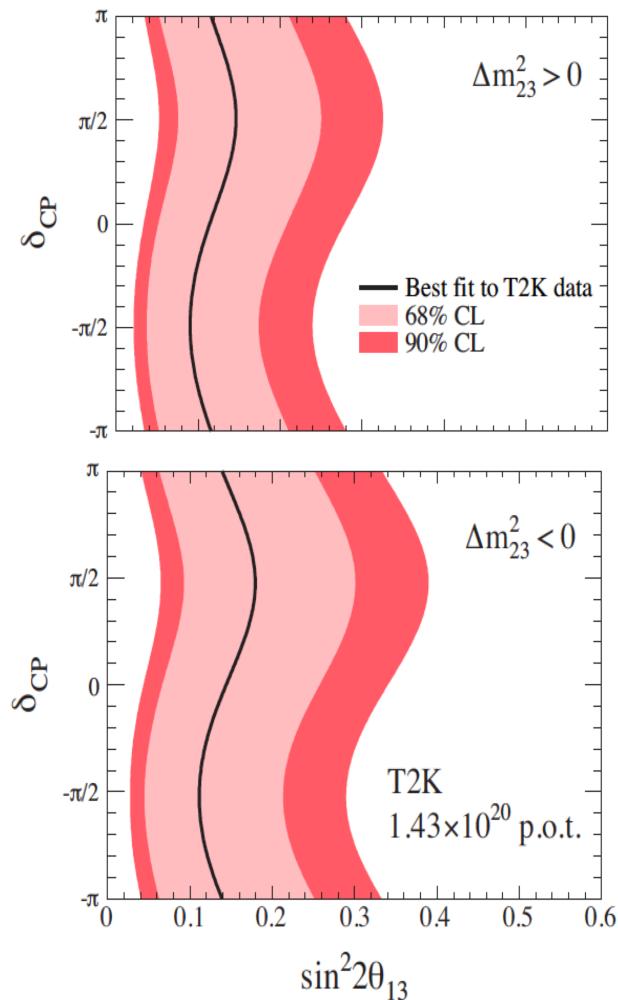
$$A = \frac{2EV}{\Delta m_{31}^2}$$

$$\Delta = \frac{\Delta m_{31}^2 L}{4E}$$

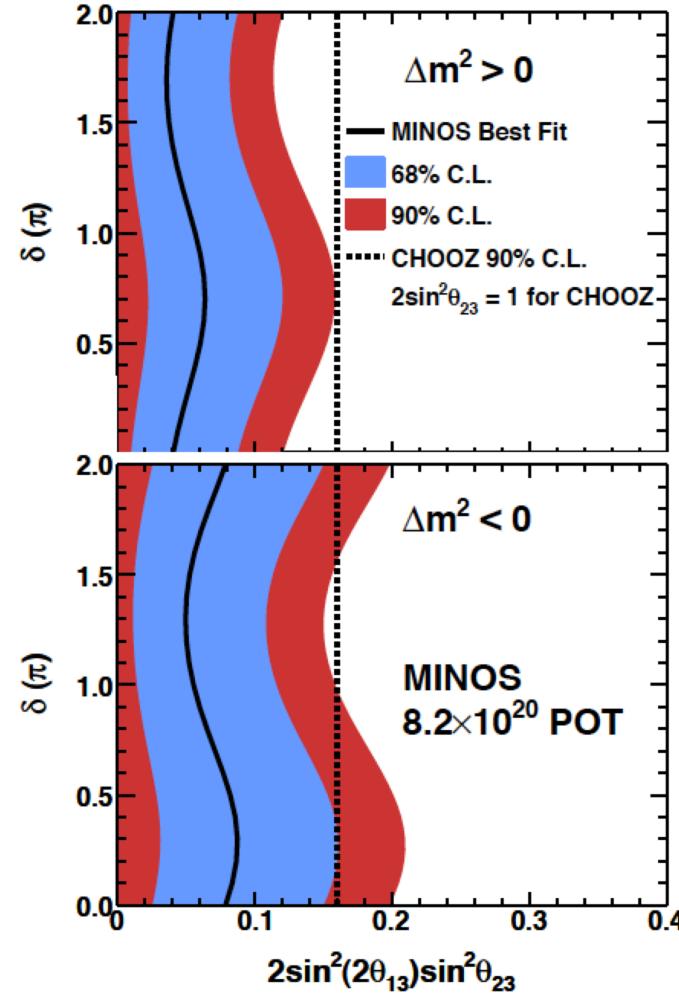
- Dependence on CP-violating terms, mass hierarchy.
- An overall eight-fold correlation ($\theta_{13}, \delta_{CP}, |\Delta m_{31}^2|, \theta_{23}, \pi/2 - \theta_{23}$)
- The bibliography is very rich:
 - *G.L. Fogli, E. Lisi, Phys. Rev. D54 (1996) 3667*
 - *J. Burguet-Castell et al., Nucl. Phys. B608 (2001) 301*
 - *H. Minakata, H. Nunokawa, JHEP 10 (2001) 001*
 - *V. Barger, D. Marfatia, K. Whisnant, Phys. Rev. D65 (2002) 073023*
 - *P. Huber, M. Lindner, W. Winter, Nucl. Phys. B645 (2002) 3*
 - *And many more...*

T2K and MINOS

T2K, Phys. Rev. Lett. 107, 041801 (2011) MINOS, Phys. Rev. Lett. 107, 181802 (2011)



NH: $0.03 < \sin^2(2\theta_{13}) < 0.28$
IH: $0.04 < \sin^2(2\theta_{13}) < 0.28$ at 90% CL

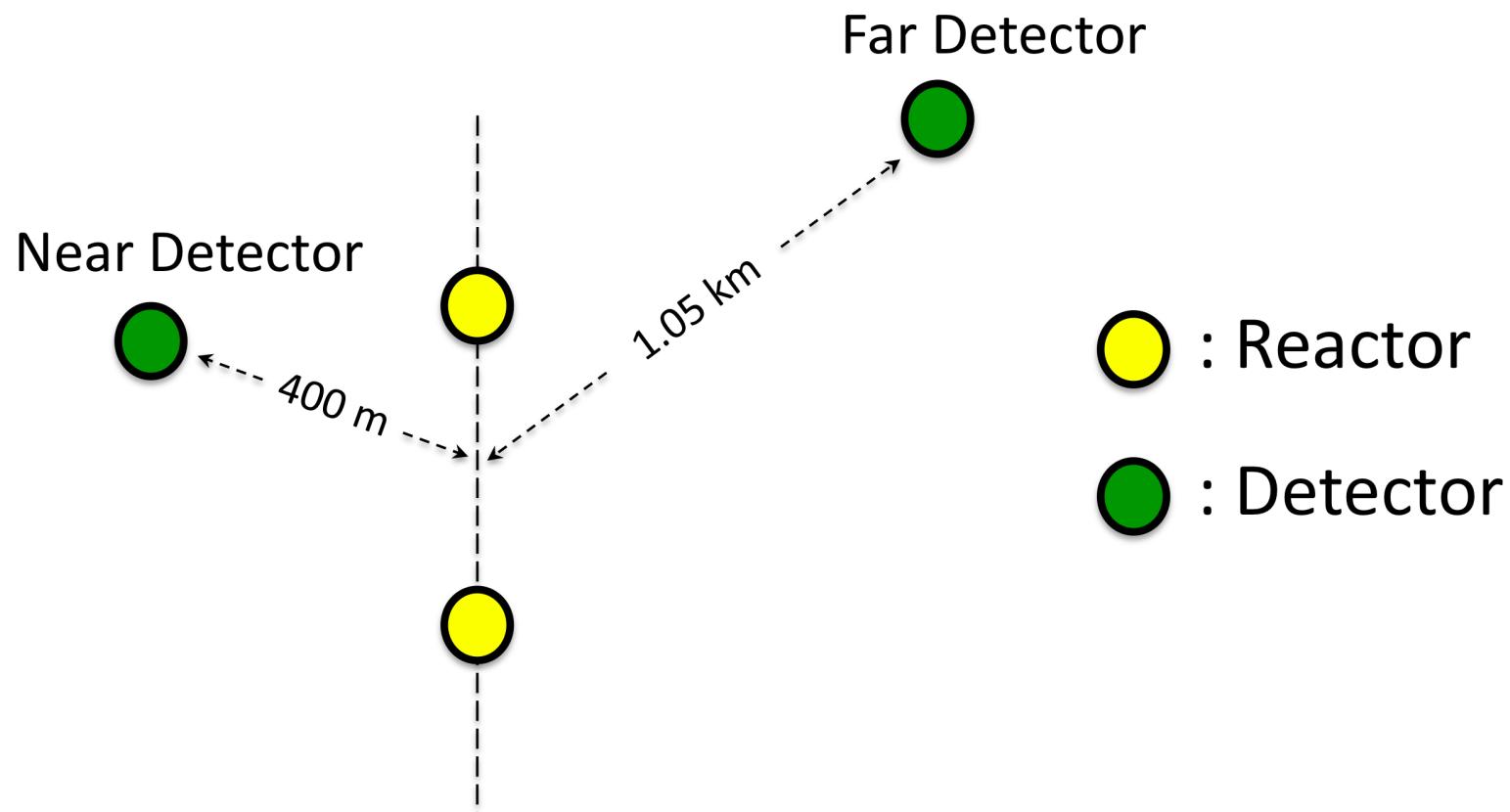


NH: $0.0 < \sin^2(2\theta_{13}) < 0.12$
IH: $0.0 < \sin^2(2\theta_{13}) < 0.19$ at 90% CL

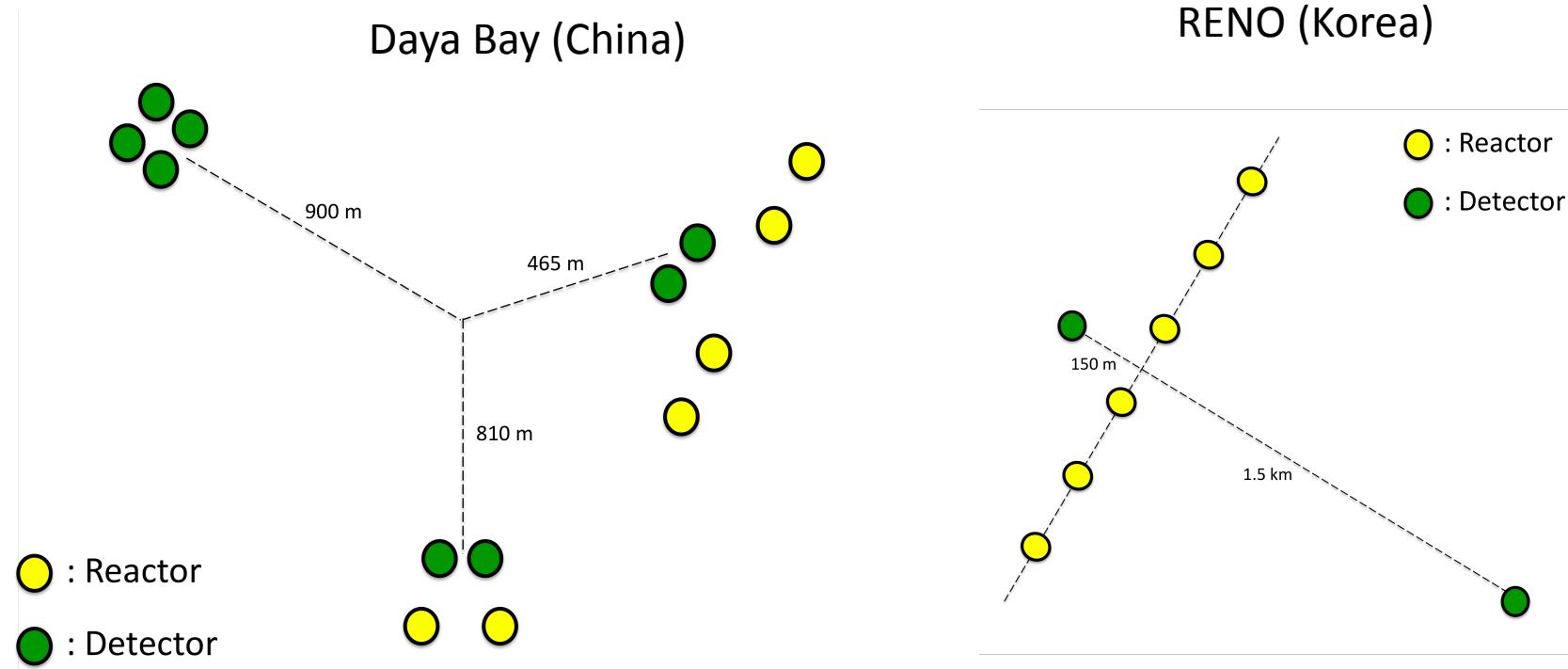
Improving CHOOZ

- Statistics:
 - Larger detector volume
 - Longer exposure
- Antineutrino flux and cross-section:
 - Detectors at different distances from the reactors
 - Identical detectors to face possible inter-detector systematics
- Detector associated errors:
 - Improve detector design
 - Better detector understanding
 - More careful calibration campaign
- Backgrounds:
 - Reduce the accidental rate: detector design and shielding
 - Aggressive radiopurity efforts
 - Improve background knowledge

Detectors configuration



Daya Bay and RENO

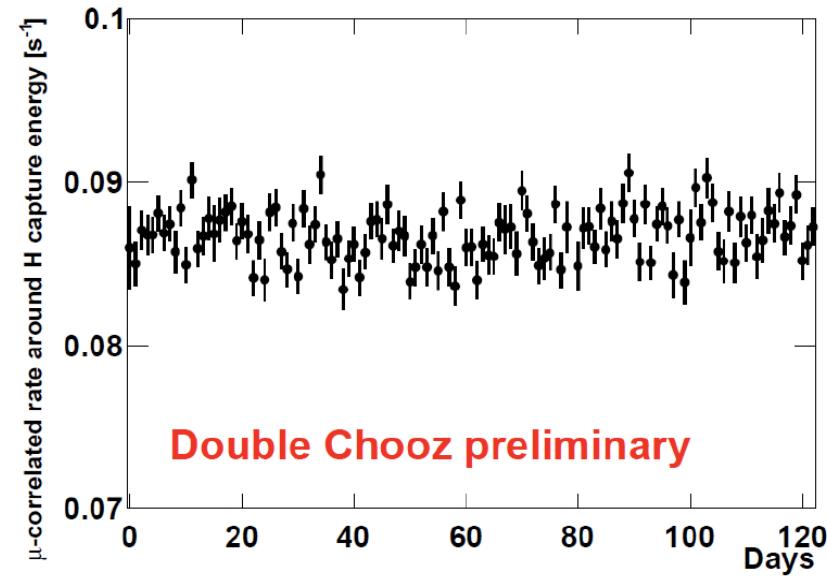
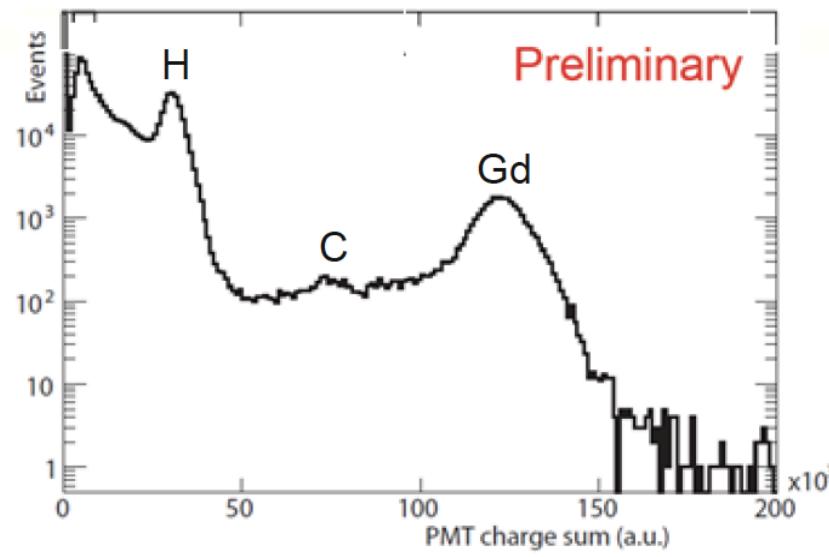


See this in connection with:

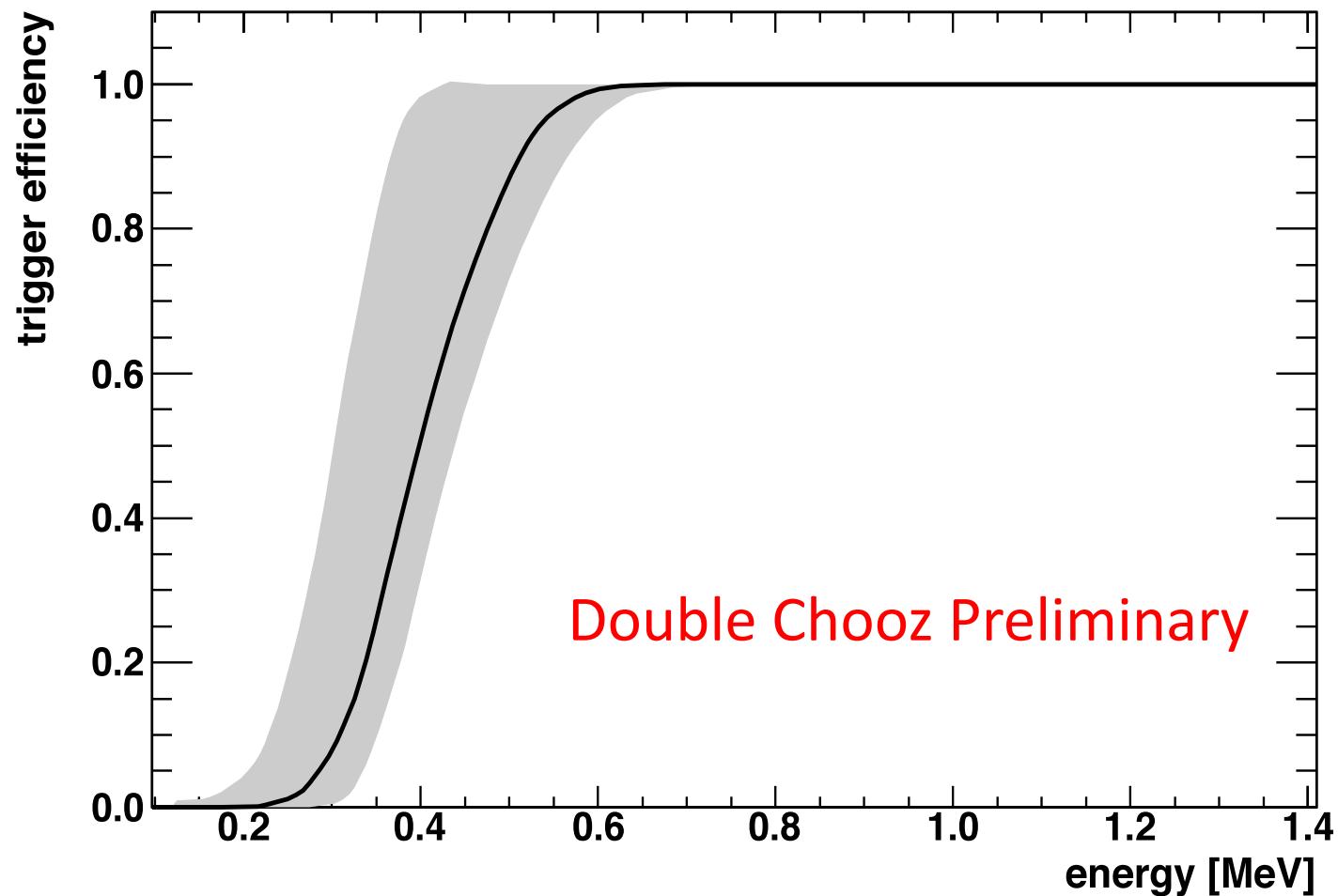
“A unified analysis of the reactor neutrino program towards the measurement of the ϑ_{13} mixing angle”, G. Mention et al., arxiv.org/0704.0498

Spallation neutrons

- Muons create copiously spallation products.
- Among them neutrons can be absorbed in hydrogen and gadolinium several μs after the prompt muon signal.
- Neutron capture peaks are very pronounced and can be used to calibrate the detector (set the energy scale) and,
- Monitor its stability.

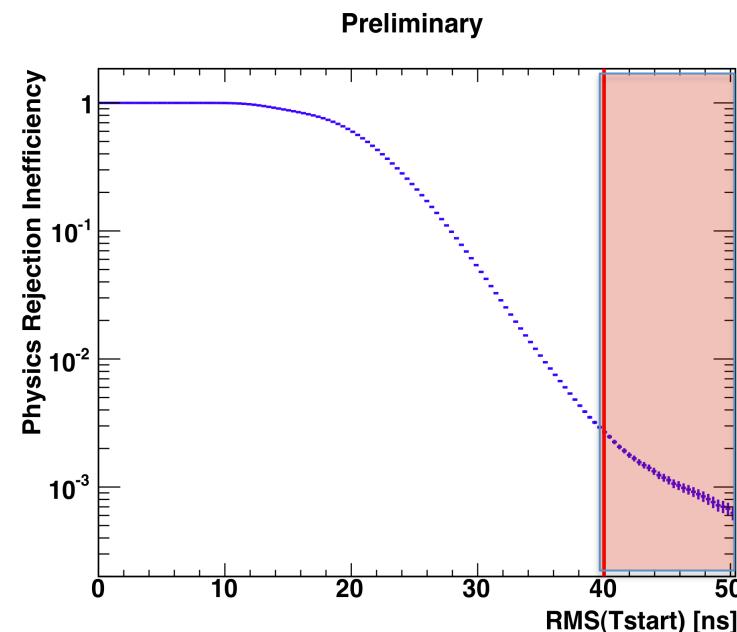
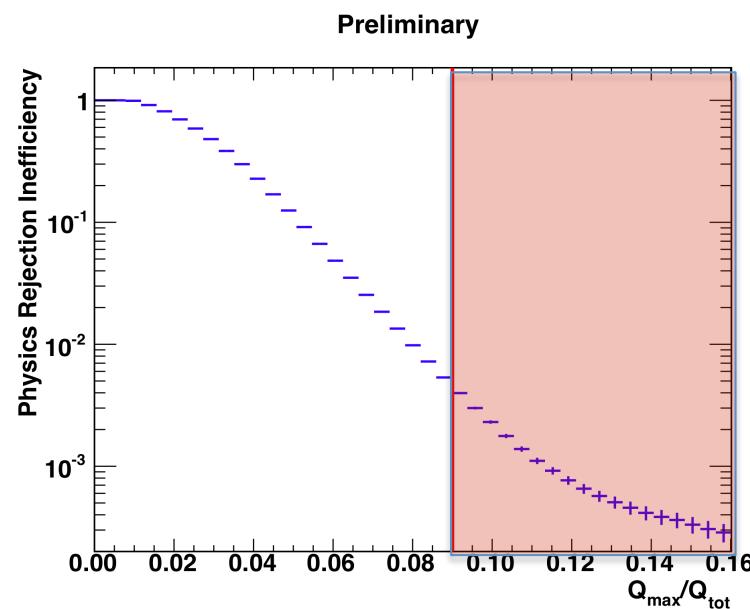
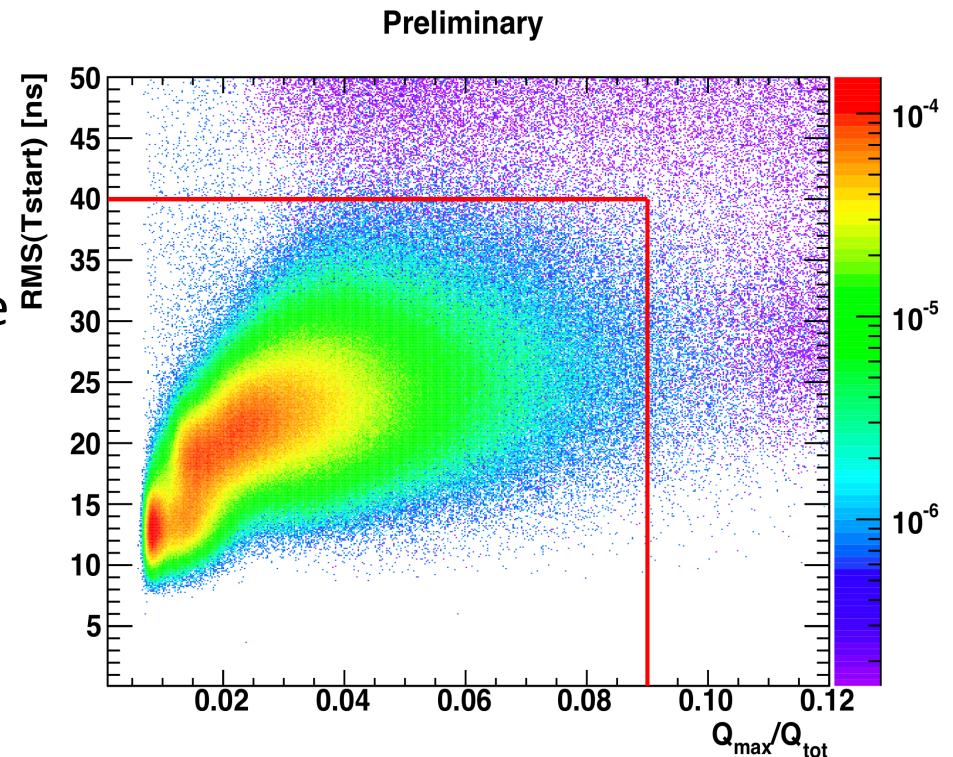


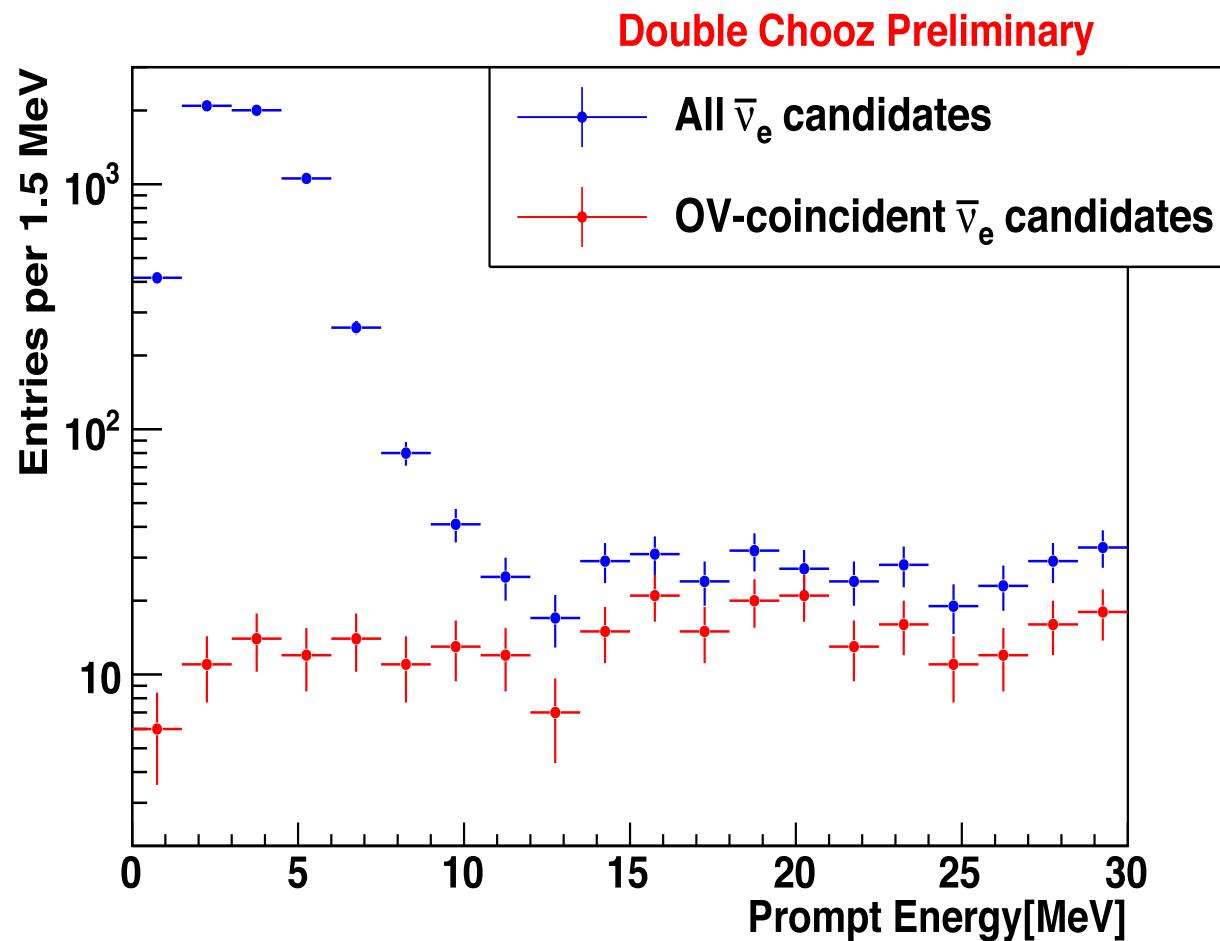
Trigger efficiency



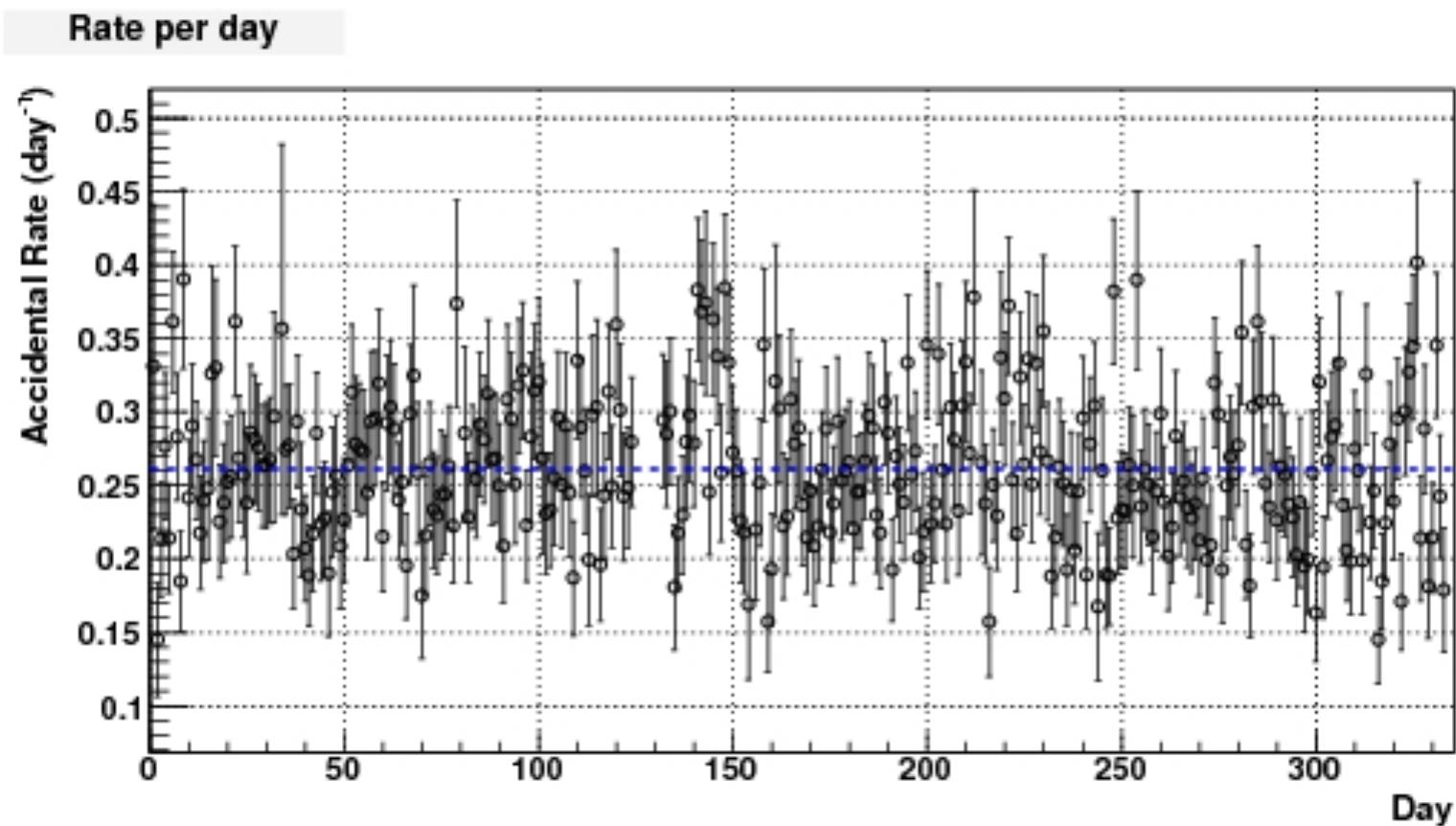
Spontaneous light emission

- Parasitic light emitted by some photomultiplier bases.
- 14 phototubes switched off.
- Software cuts:
 - $Q_{\max}/Q_{\text{tot}} < 0.09$ (prompt), 0.055 (delayed)
 - $\text{RMS}(T_s) < 40.0 \text{ ns}$
- This issue is well under control.



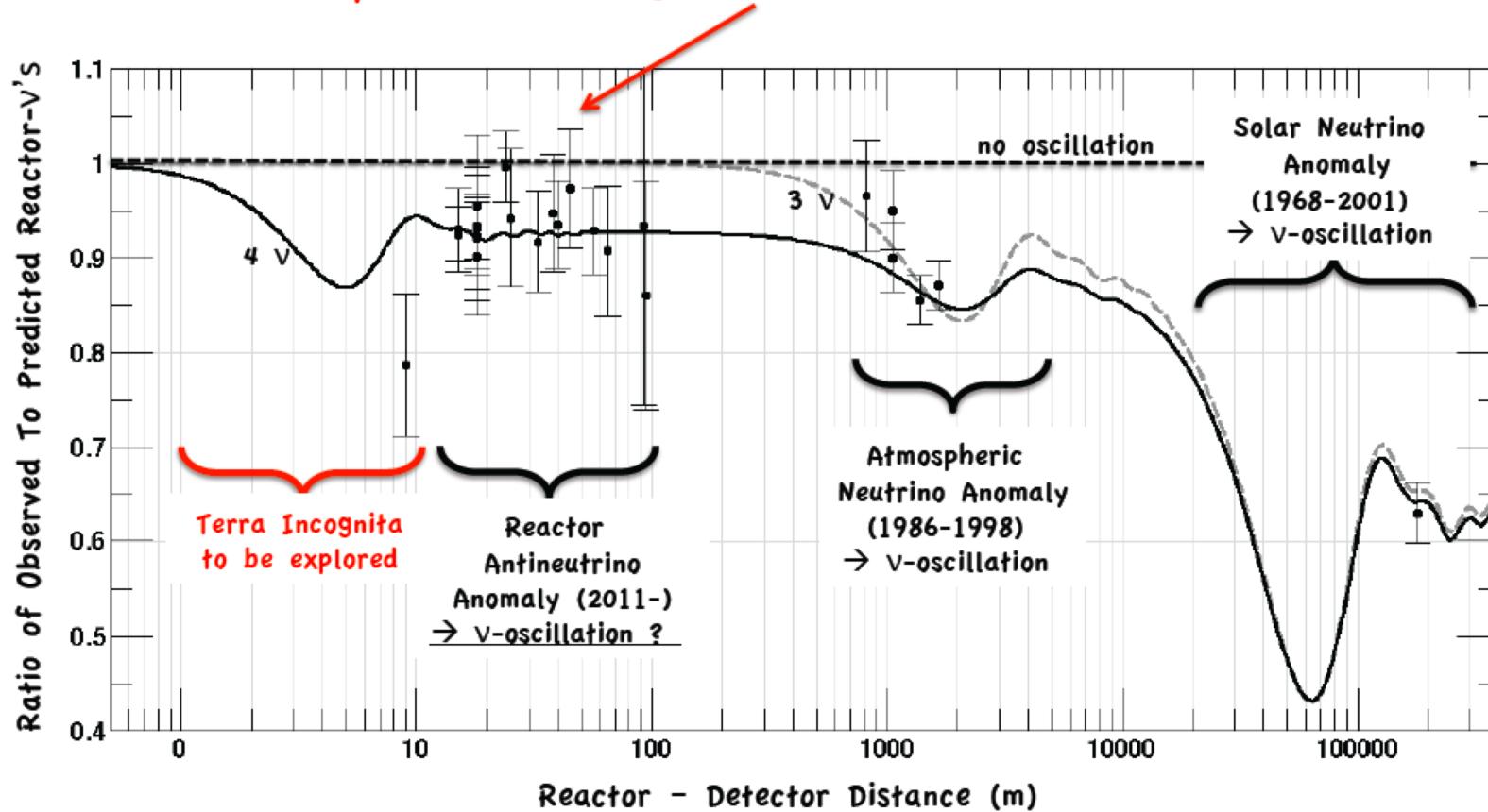


Accidental rate



Terra incognita

- Observed/predicted averaged event ratio: $R=0.927\pm0.023$ (3.0σ)



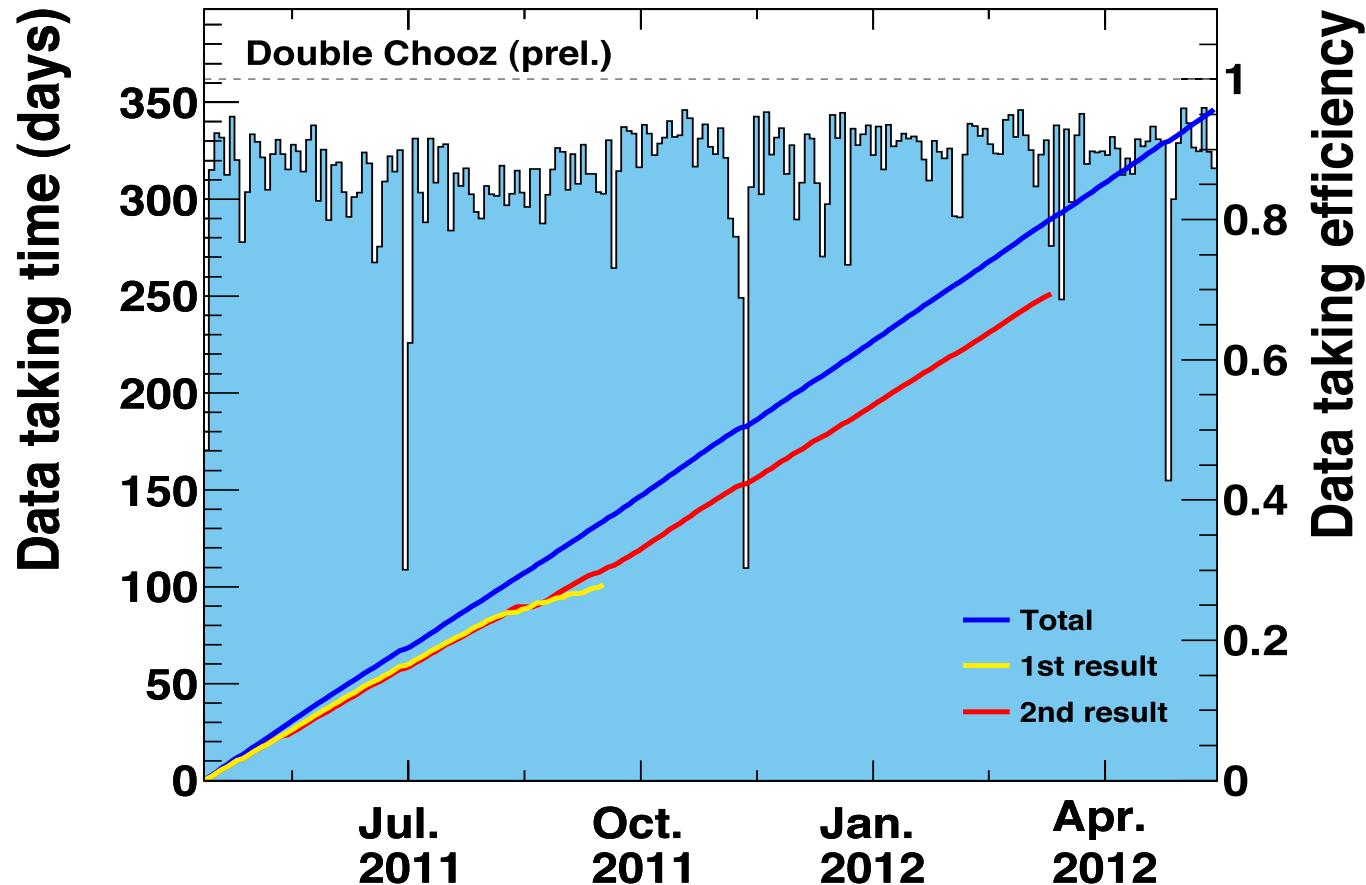
Th. Lasserre Neutrino 2012

Error budget

(M. Ishitsuka, Neutrino 2012)

Source	Uncertainty w.r.t. signal (previous analysis)
Statistics	1.1% (1.6%)
Flux	1.7%
Detector	Energy response
	0.3% (1.7%)
	E_{delay} containment
	0.7%
	Gd fraction
	0.3%
	Δt cut
Background	0.5%
	Spill in/out
	0.3%
	Trigger efficiency
	<0.1%
Background	Target H
	0.3%
	Accidental
Background	<0.1%
	Fast neutron + stop μ
	0.5% (0.9%)
	1.4% (2.8%)

Data taking summary



- Run-time: **251.3 days** (13 April 2011 - 15 March 2012)
- Data taking efficiency in total: **87.0 %**
- Data taking efficiency for physics: **79.8 %**