

# Some recent results from lattice QCD

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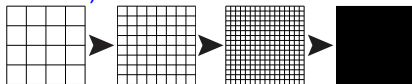
# Purpose of lattice QCD

- QCD fundamental objects: quarks and gluons
- QCD observed objects: protons, neutrons ( $\pi$ , K, ...)
- ! Huge discrepancy: not even the same particles observed as in the Lagrangean
- Perturbation theory has no chance
- Need to solve low energy QCD to:
  - Compute hadronic and nuclear properties  
“people who love QCD”
    - Masses, decay widths, scattering lengths, thermodynamic properties, ...
  - Compute hadronic background  
“people who hate QCD”
    - Non-leptonic weak MEs, quark masses, g-2, ...

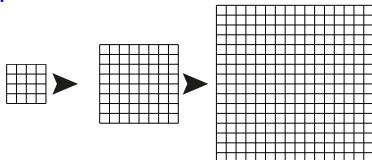
# Lattice

Lattice QCD=QCD when

- Cutoff removed (continuum limit)

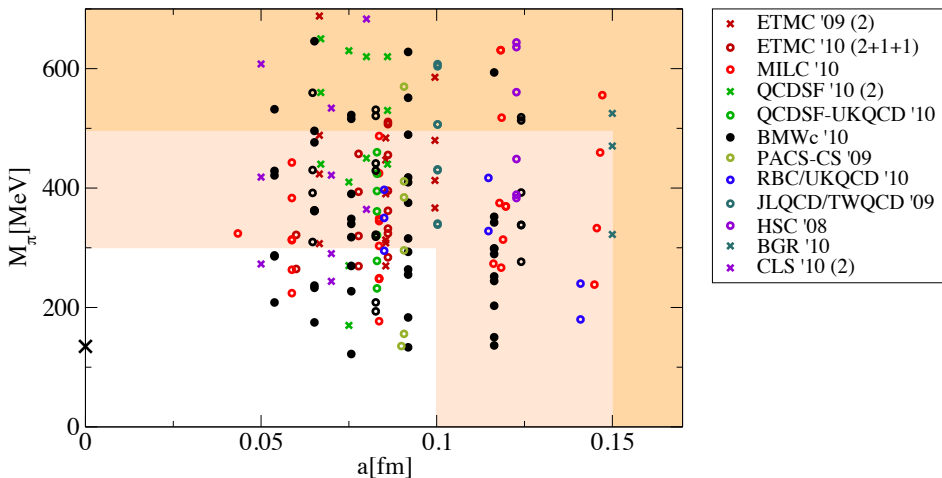


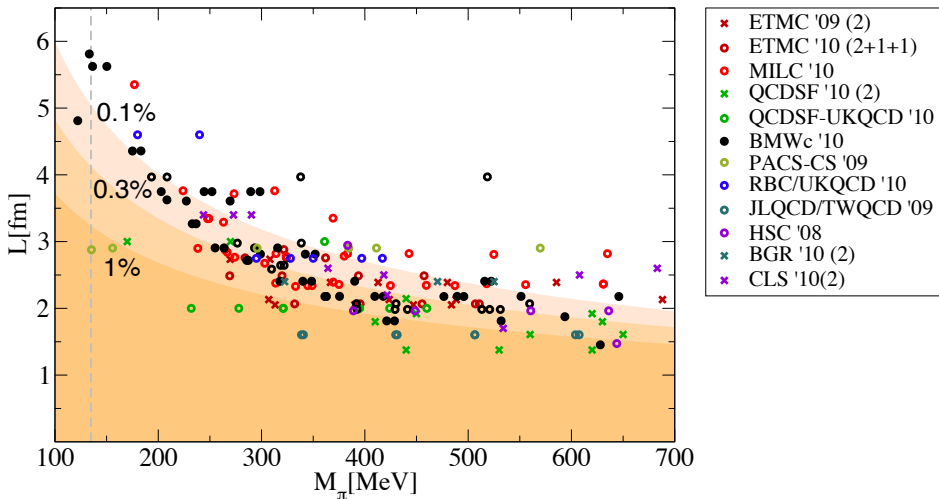
- Infinite volume limit taken

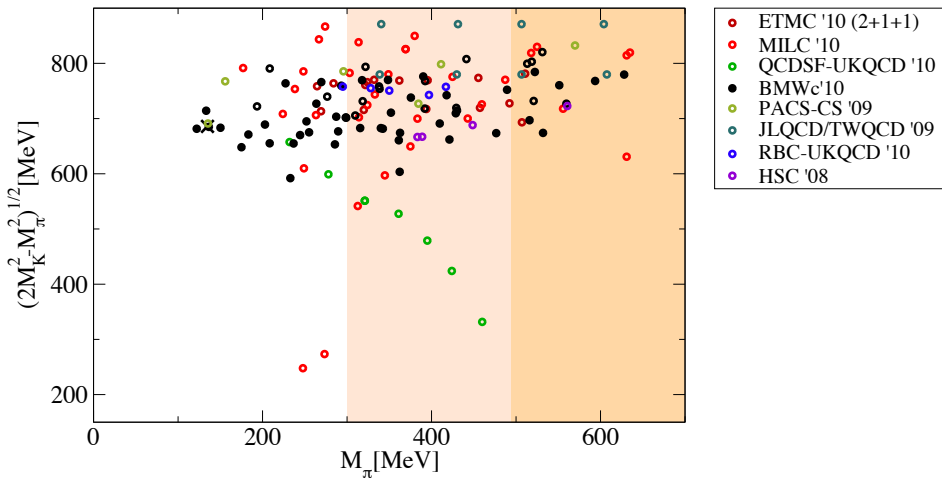


- At physical hadron masses (Especially  $\pi$ )
  - Numerically challenging to reach light quark masses

Statistical error from stochastic estimate of the path integral

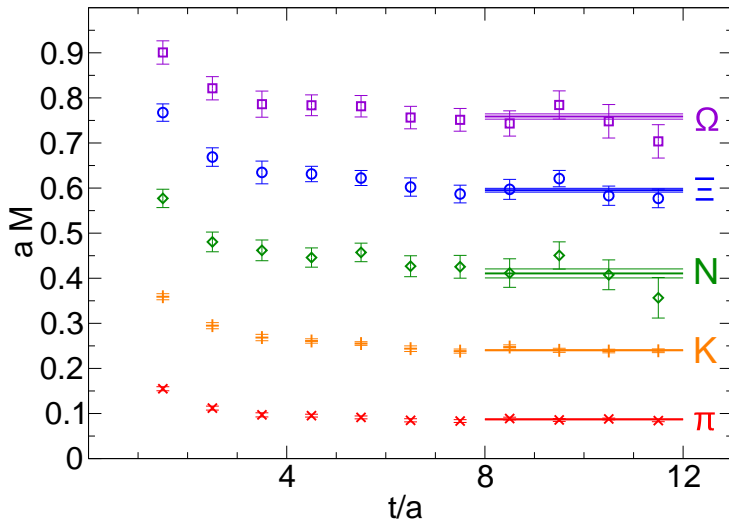
Landscape  $M_\pi$  vs.  $a$ 

Landscape  $L$  vs.  $M_\pi$ 

Landscape  $M_K$  vs.  $M_\pi$ 

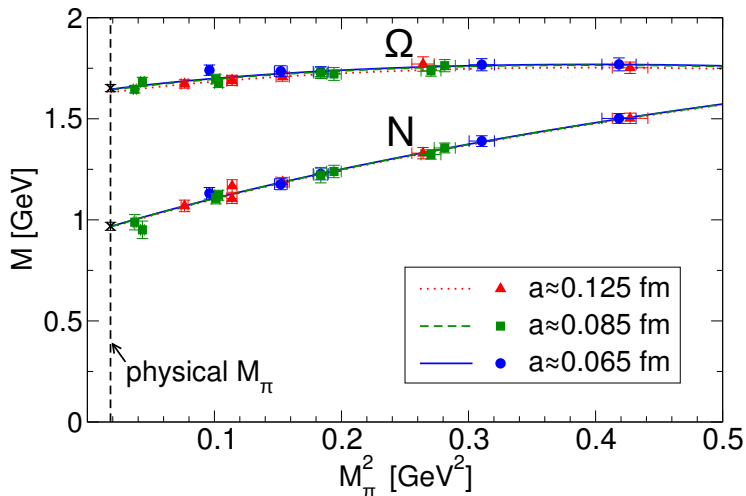
We have done our homework

## Effective masses and correlated fits



We have done our homework

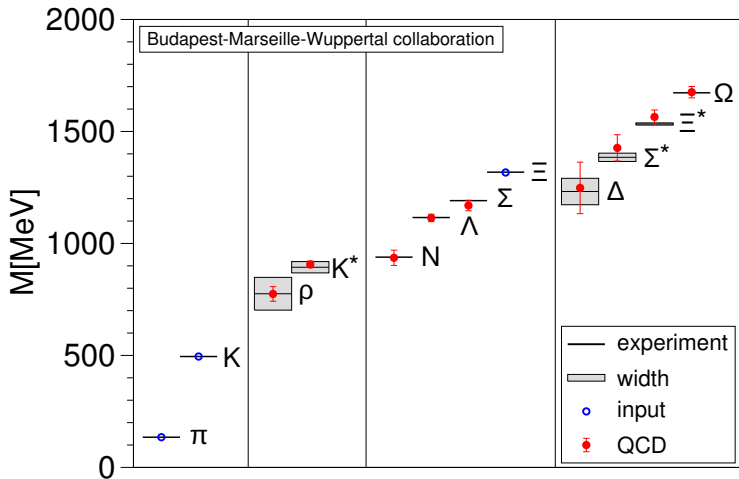
## Chiral fit





We have done our homework

## The light hadron spectrum



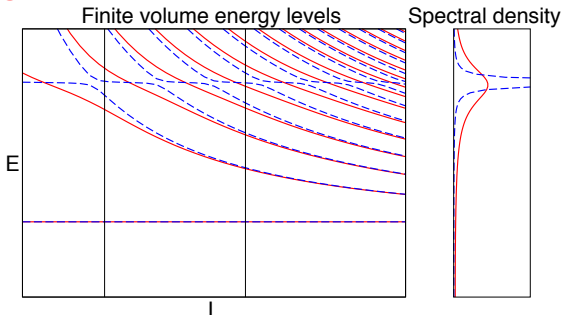
(Science 322:1224,2008)

We have done our homework

# Excited states

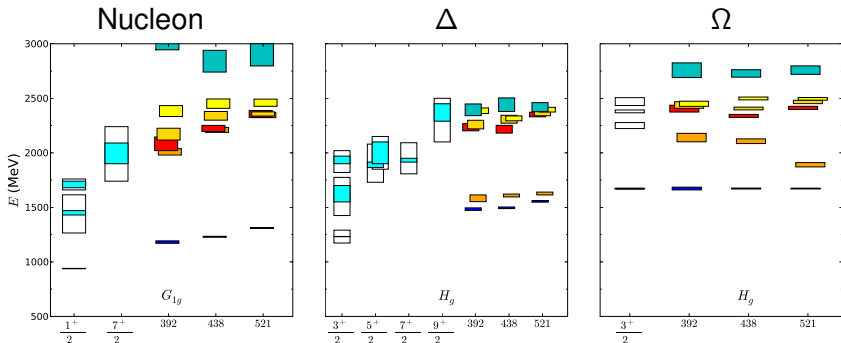
Extracting excited states is much tougher:

- ☞ Extraction of energy levels is harder:
  - Die out at large  $t$   $\Rightarrow$  need to use small  $t$  correlators
- ☞ Once extracted, relation to  $V \rightarrow \infty$  is nontrivial:
  - Disentangle resonances and scattering states at finite volume



We have done our homework

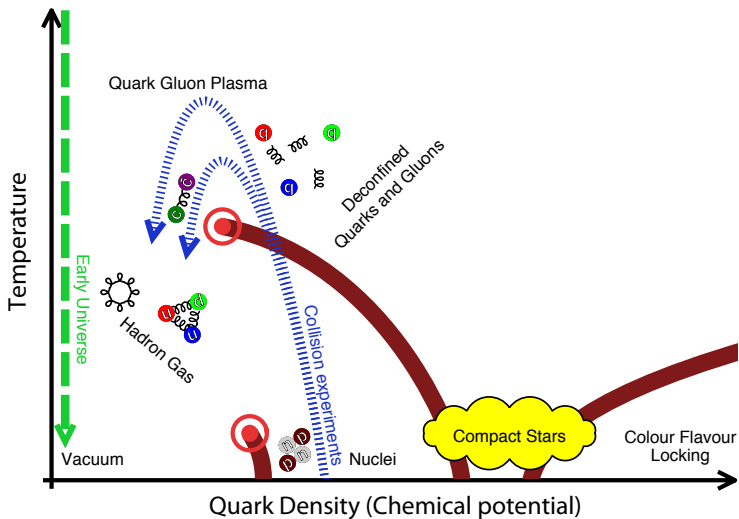
# High lying resonances



(Bulava, et al., 2010)

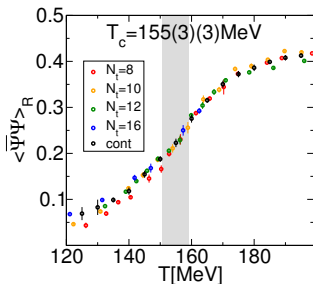
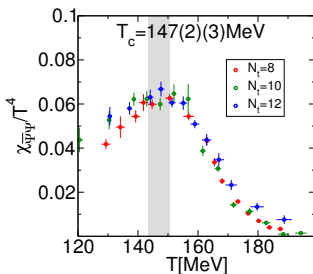
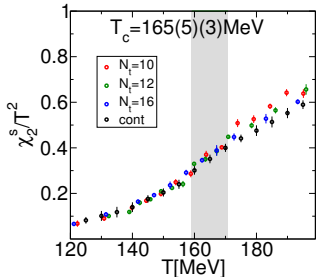
- ✓ Qualitative understanding of experimental spectrum
- ✗ No extrapolation to physical point, continuum

# QCD phase diagram



Crossover at  $\mu = 0$ QCD at  $\mu = 0$ 

- No phase transition but crossover (Aoki et. al. '06)
  - At physical quark masses
- Spread in pseudocritical temperatures for different observables (Borsanyi et. al. '10)



Crossover at  $\mu = 0$ 

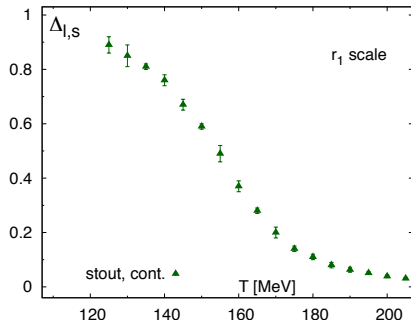
# Transition temperature

## Discrepancies in $T_c$

- Cheng et. al. '06:
  - Unique  $T_c$
  - $T_c = 192(7)(4)\text{MeV}$
- Aoki et. al. '06:
  - Spread in  $T_c$
  - $T_c(\chi_{\bar{\psi}\psi}) = 151(3)(3)\text{MeV}$
  - $T_c(\chi_2^S) = 175(2)(4)\text{MeV}$

Has been resolved:

- Lattice artifacts
- Unphysical quark masses



(HotQCD (Bazavov et. al.) 2012)

Crossover at  $\mu = 0$ 

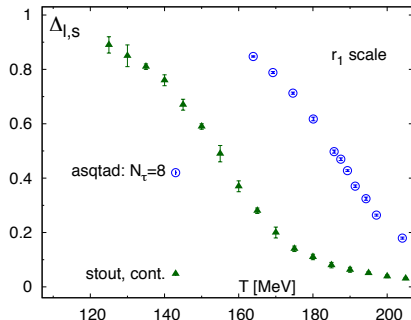
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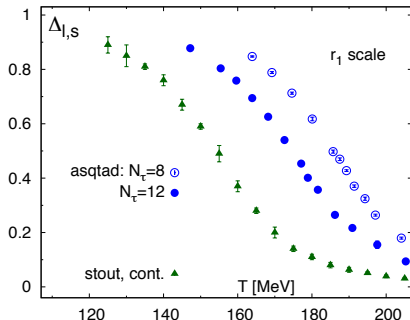
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Crossover at  $\mu = 0$ 

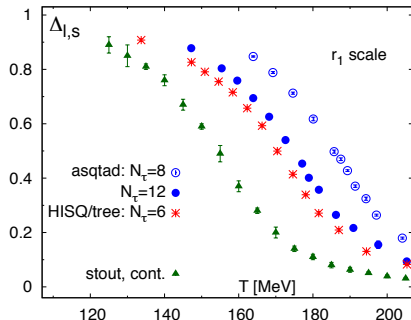
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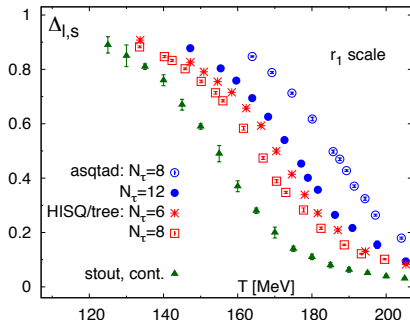
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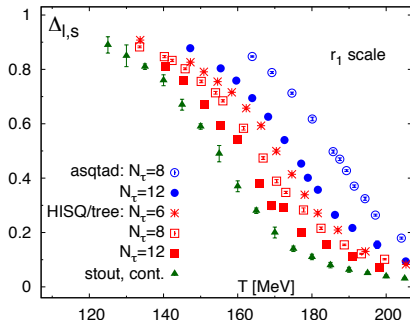
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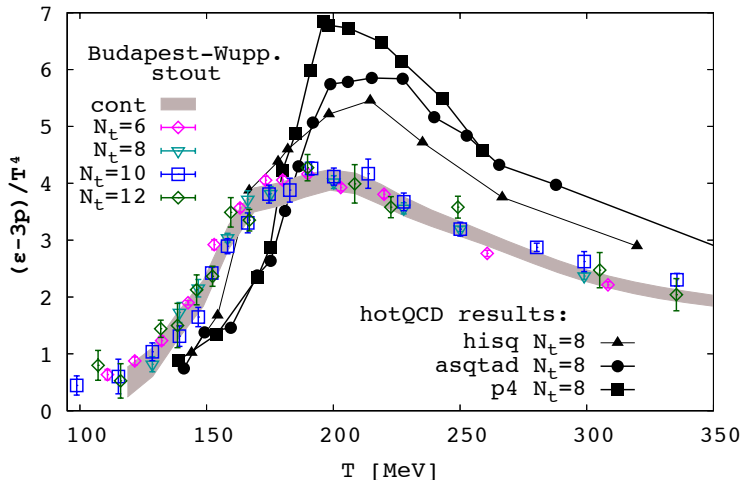
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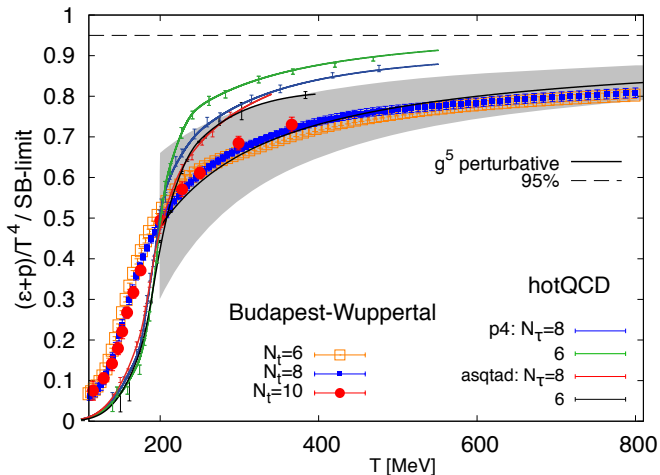
Has been resolved:

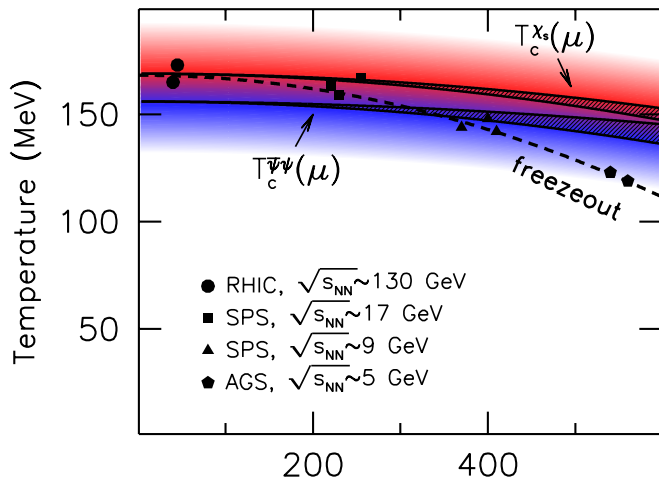
- Lattice artifacts
- Unphysical quark masses



(HotQCD (Bazavov et. al.) 2012)

Crossover at  $\mu = 0$ QCD equation of state at  $\mu = 0$ 

Crossover at  $\mu = 0$ Entropy at  $\mu = 0$ 

Transition at  $\mu > 0$ Curvature of the transition line at  $\mu = 0$ 

Baryonic chemical potential (MeV) (Endrodi et al., 2011)

# Light quark masses

## Goal:

- Compute light quark masses ab initio

## Method:

- Go to the physical point
- Read off input quark masses and renormalize

## Challenge:

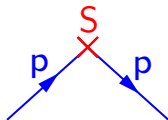
- Minimize and control **all** systematics
  - 2+1 dynamical fermion flavors
  - Physical quark masses
  - Continuum extrapolation
  - Infinite volume
  - Nonperturbative renormalization

# Renormalization

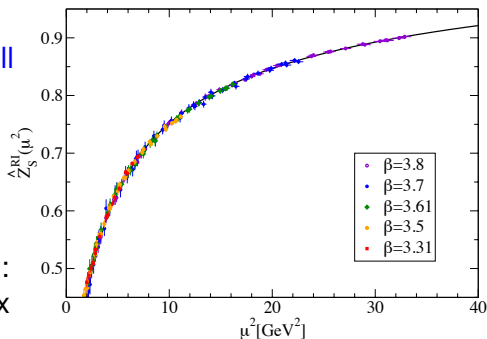
- Quark masses logarithmically divergent ( $a \rightarrow 0$ )  $\rightarrow$  renormalization
- Usually  $\overline{\text{MS}}$  scheme: only perturbatively defined

## RI-MOM scheme

- matrix elements of off-shell quarks in fixed gauge

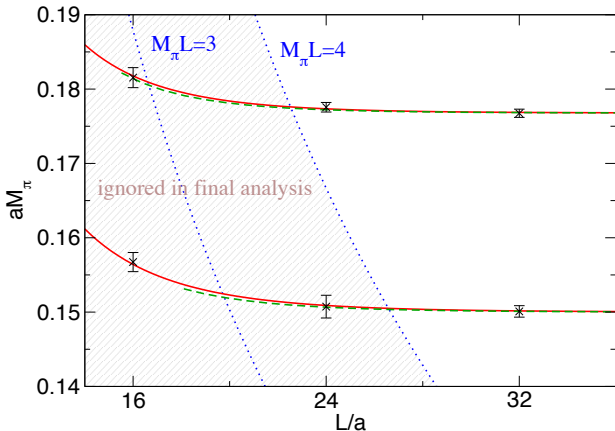


- Renormalization condition: at  $p^2 = \mu^2$  tree level matrix element



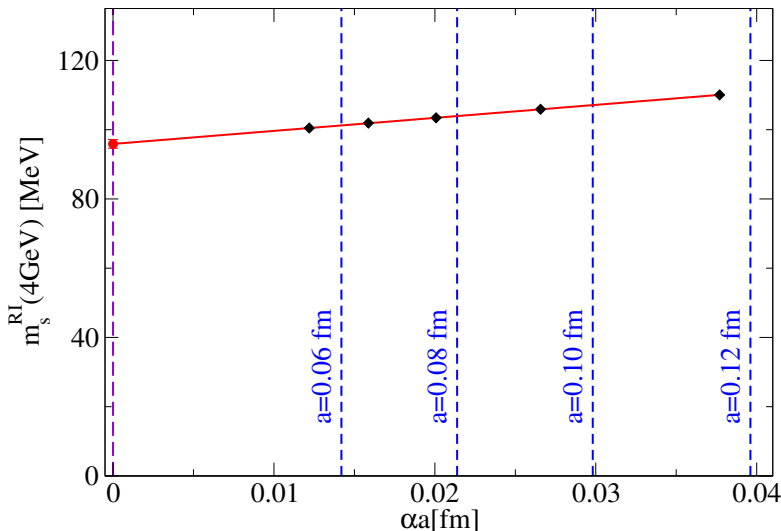


## Tiny finite volume effects



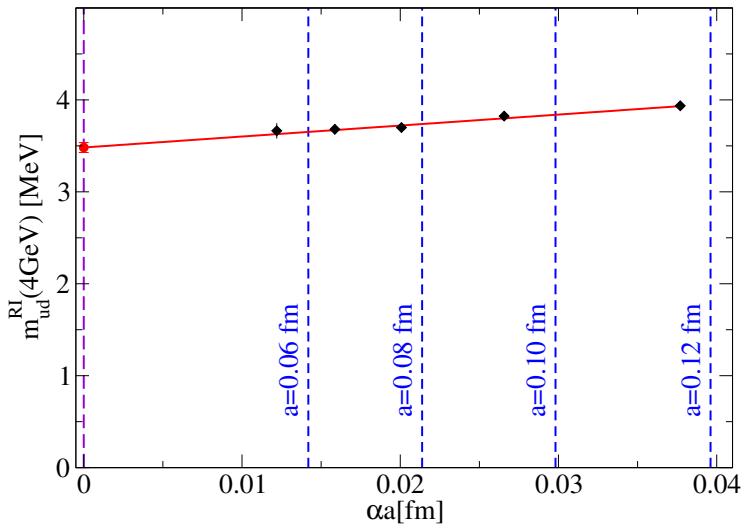
- FV effects tiny
- Dedicated FV runs
- Perfect agreement with FV  $\chi$ PT (Colangelo et. al. 2005)

# Strange quark mass



$m_{ud}$  and  $m_s$ 

## Light quark masses



# Individual $m_u$ and $m_d$

- **Goal:**
  - Compute  $m_u$  and  $m_d$  separately
- **Method:**
  - Need QED and isospin breaking effects in principle
  - Alternative: use dispersive input -Q from  $\eta \rightarrow \pi\pi\pi$ 

$$Q^2 = \frac{1}{2} \left( \frac{m_s}{m_{ud}} \right)^2 \frac{m_d - m_u}{m_{ud}}$$
  - ✓ Transform precise  $m_s/m_{ud}$  into  $(m_d - m_u)/m_{ud}$
  - We use the conservative  $Q = 22.3(8)$  (Leutwyler, 2009)

# Systematic error treatment

- Goal:
  - Reliably estimate total systematic error
- Method:
  - 288 full analyses (2000 bootstrap on each)
    - 2 plateau regions
    - 2 continuum forms:  $O(\alpha_s a)$ ,  $O(a^2)$
    - 3 chiral forms:  $2 \times SU(2)$ , Taylor
    - 2 chiral ranges:  $M_\pi < 340, 380$  MeV
    - 3 renormalization matching procedures
    - 2 NP continuum running forms
    - 2 scale setting procedures
  - All analyses weighted by fit quality
    - Mean gives final result
    - Stdev gives systematic error
  - Statistical error from 2000 bootstrap samples

## Final result

	RI @ 4 GeV	RGI	$\overline{\text{MS}}$ @ 2 GeV
$m_s$	96.4(1.1)(1.5)	127.3(1.5)(1.9)	95.5(1.1)(1.5)
$m_{ud}$	3.503(48)(49)	4.624(63)(64)	3.469(47)(48)
$m_s/m_{ud}$		27.53(20)(8)	
$m_u$	2.17(04)(10)	2.86(05)(13)	2.15(03)(10)
$m_d$	4.84(07)(12)	6.39(09)(15)	4.79(07)(12)

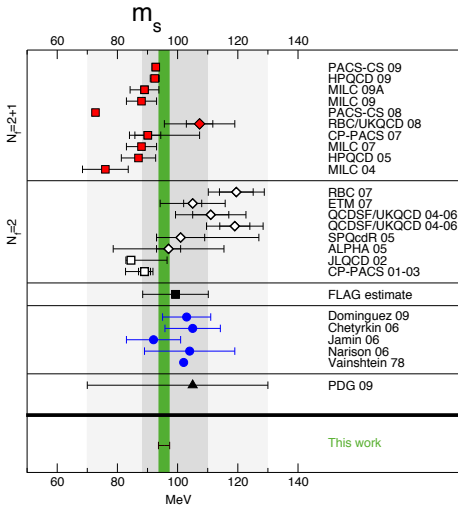
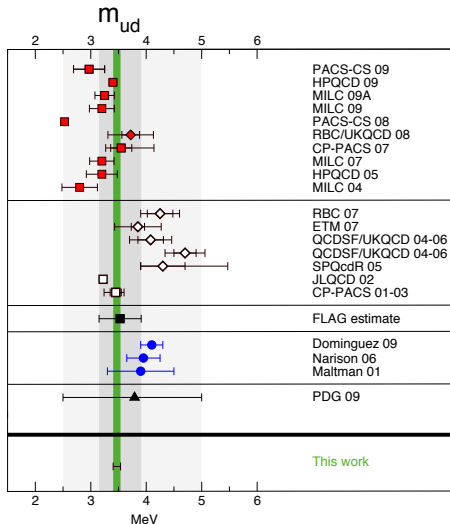
## Relative error budget:

	stat.	plateau	scale set	mass	renorm.	cont.
$m_s$	0.702	0.148	0.004	0.064	0.061	0.691
$m_{ud}$	0.620	0.259	0.027	0.125	0.063	0.727
$m_s/m_{ud}$	0.921	0.200	0.078	0.125	—	0.301

(JHEP 1108:148,2011; PLB 701:265,2011)

## Results

## Comparison



# Future perspective

- **Goal:**
  - Decrease precision below percent level
- **Method:**
  - Take light masses non-degenerate
  - Include QED effects
  - Include charm quark effects
- **Challenge:**
  - Stability near  $m_u < m_{ud}$
  - Finding the physical point
  - Non gauge invariant final states
  - Quark mass ratios renormalize



# Standard model neutral K mixing

- **Goal:**

- Check SM CP violation in neutral K system

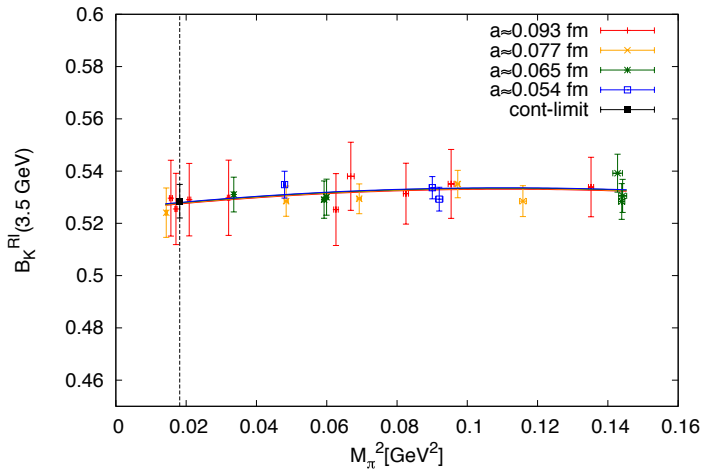
- **Method:**

- Compute effective weak matrix element  $B_K = \frac{3\langle \bar{K} | O_{\Delta S=2} | K \rangle}{8f_K^2 M_K^2}$
- Relate kaon CP violation to CKM phase  
→ from CKM unitarity:  $\hat{B}_K = 0.83_{-0.15}^{+0.21}$  (CKMfitter, 2011)

- **Challenge:**

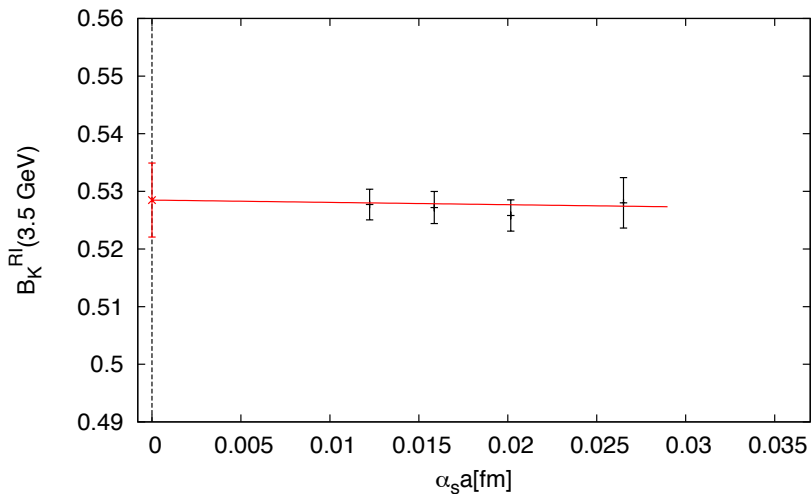
- Minimize and control **all** systematics
  - 2+1 dynamical fermion flavors
  - Physical quark masses
  - Mixing of unphysical operators
  - Continuum
  - Infinite volume

## Physical point



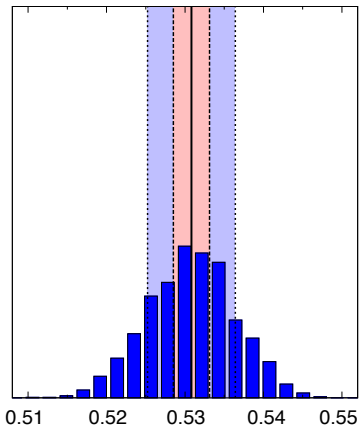


# Continuum extrapolation

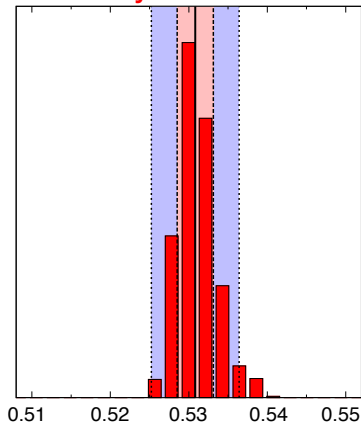


## Errors

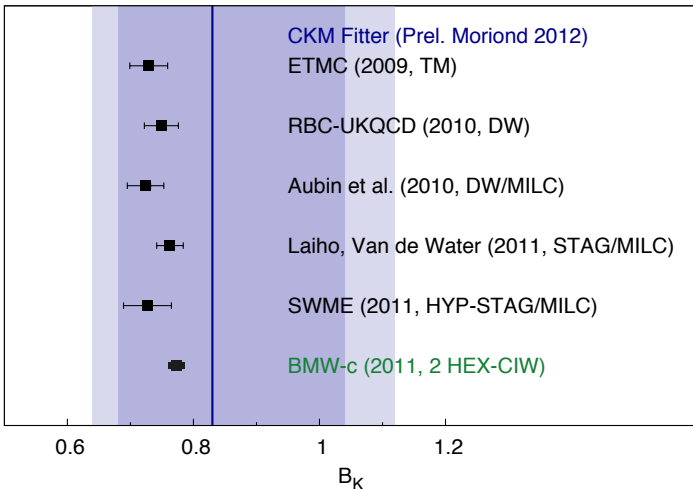
statistical



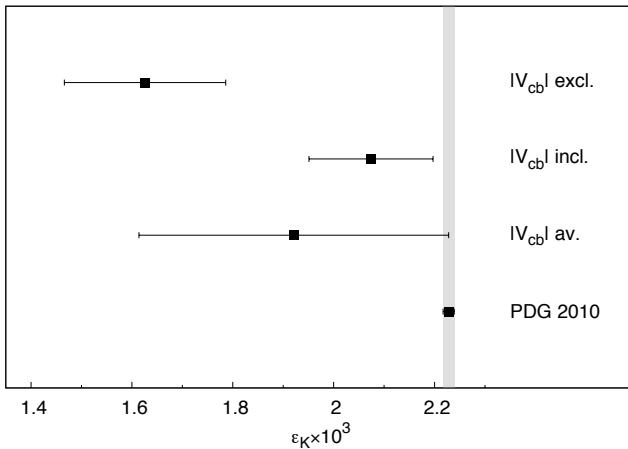
systematic

 $B_K^{\text{RI}}(3.5 \text{ GeV})$

## Comparison



# Implications

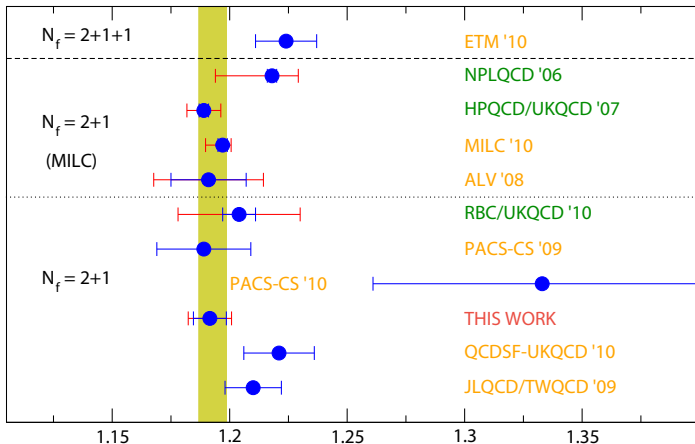


# Pseudoscalar decay constant ratio

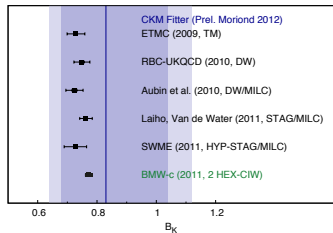
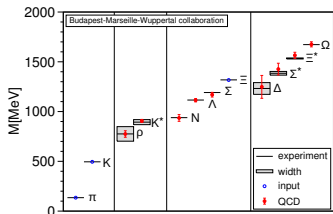
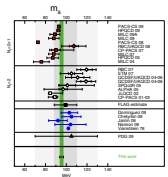
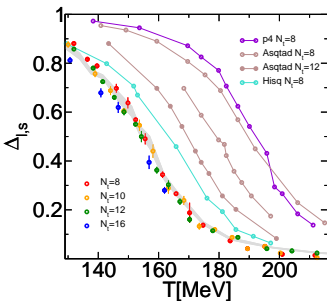
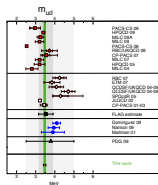
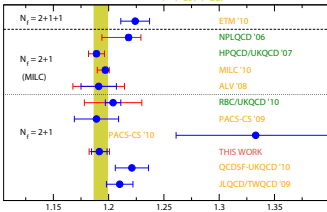
- **Goal:**
  - Check first row unitarity of CKM matrix
- **Method:**
  - Compute  $F_K/F_\pi$
  - Perturbative relation to  $|V_{us}|^2/|V_{ud}|^2$  with 0.4% accuracy
- **Challenge:**
  - Minimize and control **all** systematics
    - 2+1 dynamical fermion flavors
    - Physical quark masses
    - Continuum
    - Infinite volume

## Decay constants

## Decay constant ratio

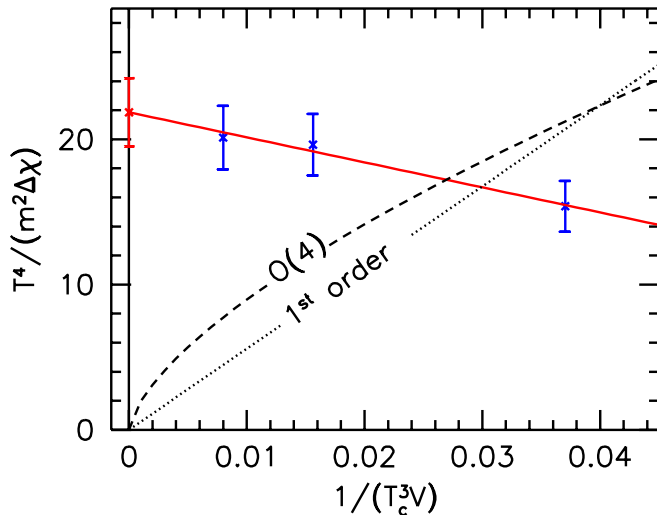
Prediction from CKM unitarity ( $|V_{us}|/|V_{ud}|$ )



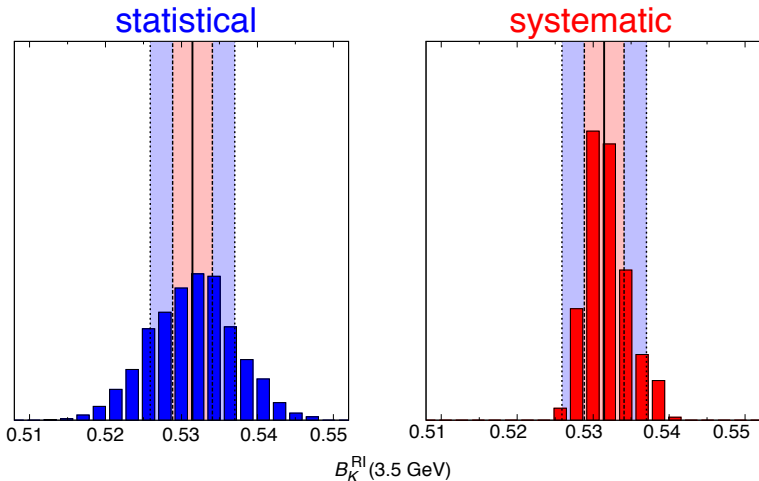
Prediction from CKM unitarity ( $|V_{us}|/|V_{ud}|$ )

# BACKUP

# Order of QCD phase transition



## Errors - unit weight



# Action details

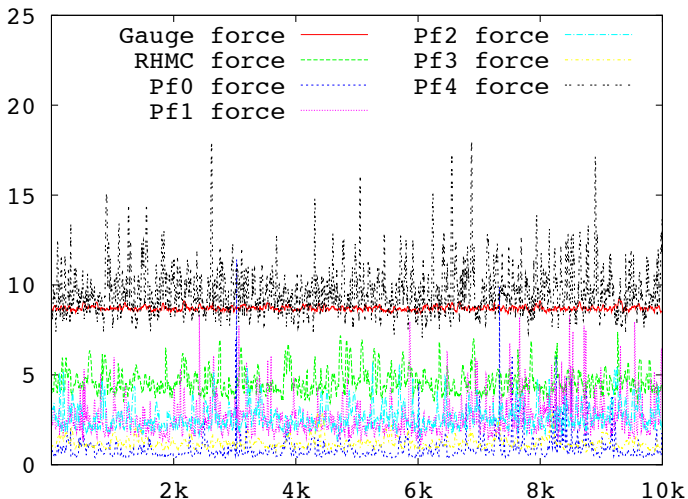
## Goal:

- Optimize physics results per CPU time
- Conceptually clean formulation

## Method:

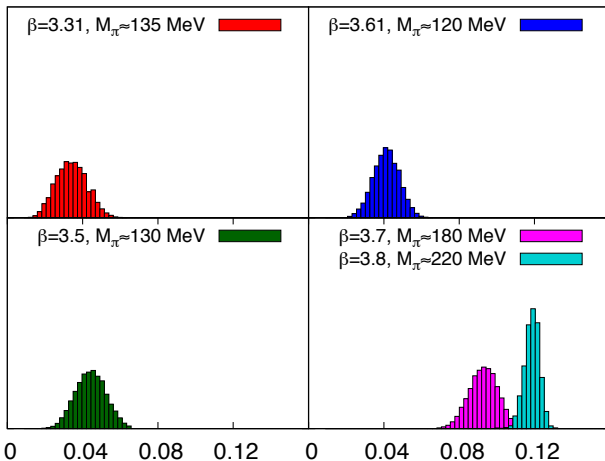
- Dynamical 2 + 1 flavor, Wilson fermions at physical  $M_\pi$
  - 3-5 lattice spacings  $0.053 \text{ fm} < a < 0.125 \text{ fm}$
  - Tree level  $O(a^2)$  improved gauge action (Lüscher, Weisz, 1985)
  - Tree level  $O(a)$  improved fermion action (Sheikholeslami, Wohlert, 1985)
    - Why not go beyond tree level?
      - Keeping it simple (parameter fine tuning)
      - No real improvement, UV mode suppression took care of this
    - This is a crucial advantage of our approach
  - UV filtering (APE coll. 1985; Hasenfratz, Knechtli, 2001; Capitani, Durr, C.H., 2006)
- Discretization effects of  $O(\alpha_s a, a^2)$
- ✓ We include both  $O(\alpha_s a)$  and  $O(a^2)$  into systematic error

# Algorithm stability



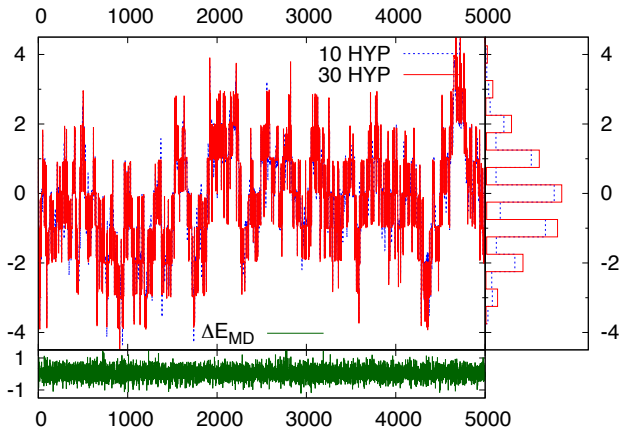
# No exceptional configs

Inverse iteration count ( $1000/N_{\text{CG}}$ )



# Topological sector sampling

Topological charge  $\beta=3.8$ ,  $m_{ud}=-0.02$ ,  $m_s=0$



worst case

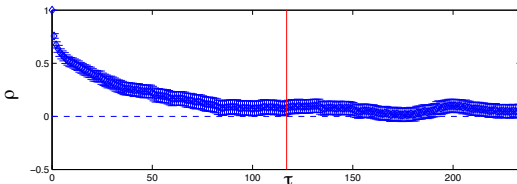


# Autocorrelation time (finest lattice, small mass)

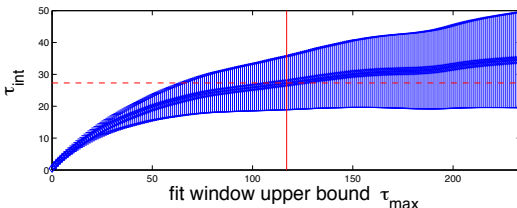
$$\tau_{\text{int}} = 27.3(7.4)$$

(MATLAB code from Wolff,  
2004-7)

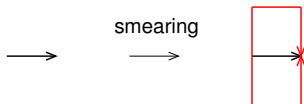
normalized autocorrelation for  $lq^{\text{ren}}$  at  $\beta=3.8$ ,  $m_{ud}=-0.02$ ,  $m_s=0$



$\tau_{\text{int}}$  with statistical errors for  $lq^{\text{ren}}$  at  $\beta=3.8$ ,  $m_{ud}=-0.02$ ,  $m_s=0$



# Locality properties



- locality in position space:

$|D(x, y)| < \text{const } e^{-\lambda|x-y|}$  with  $\lambda = O(a^{-1})$  for all couplings.

Our case:  $D(x, y) = 0$  as soon as  $|x - y| > 1$

(despite smearing)

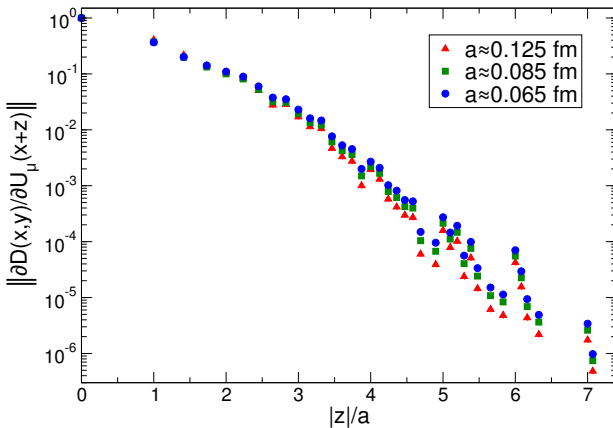
- locality of gauge field coupling:

$|\delta D(x, y)/\delta A(z)| < \text{const } e^{-\lambda|(x+y)/2-z|}$  with  $\lambda = O(a^{-1})$  for all couplings.

Our case:  $\delta D(x, x)/\delta A(z) < \text{const } e^{-\lambda|x-z|}$  with  $\lambda \simeq 2.2a^{-1}$   
for  $2 \leq |x - z| \leq 6$

# Gauge field coupling locality

6-stout case:



# Light hadron spectrum

- **Goal:**
  - Firmly establish (or invalidate?) QCD as the theory of strong interaction in the low energy region
- **Method:**
  - Post-diction of light hadron spectrum
    - Octet baryons
    - Decuplet baryons
    - Vector mesons
- **Challenge:**
  - Minimize and control **all** systematics
    - 2+1 dynamical fermion flavors
    - Physical quark masses
    - Continuum
    - Infinite volume (treatment of resonant states)

# Scale setting

## Goal:

- Unambiguous, precise scale setting

## Method:

- We set the scale via a baryon mass
- Desirable properties:
  - experimentally well known
  - small lattice error (Octet better than Decuplet)
  - independent of light quark mass  $\rightarrow$  large strange content
- Best candidates:
  - $\Xi$ : largest strange content of the octet
  - $\Omega$ : member of the decuplet, but no light quarks

# Quark mass dependence

## Goal:

- Extra-/Interpolate  $M_X$  (baryon/vector meson mass) to physical point ( $M_\pi$ ,  $M_K$ )

## Method:

- Fundamental parameters:  $g$ ,  $m_{ud}$ ,  $m_s$ 
  - Experimentally inaccessible (confinement!)
  - Must be set via 3 experimentally accessible quantities
- Use  $M_\Xi$  or  $M_\Omega$  and  $M_\pi$ ,  $M_K$  to set parameters
- Variables to parametrize  $M_\pi^2$  and  $M_K^2$  dependence of  $M_X$ :
  - Use bare masses  $aM_y$ ,  $y \in \{X, \pi, K\}$  and  $a$  (bootstrapped)
  - Use dimensionless ratios  $r_y := \frac{M_y}{M_{\Xi/\Omega}}$  (cancellations)

We use both procedures  $\rightarrow$  systematic error

# Quark mass dependence (ctd.)

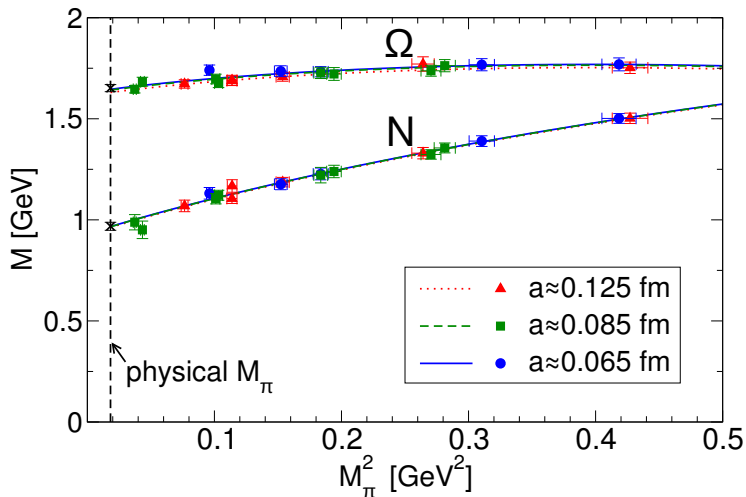
## Method (ctd.):

- Parametrization:  $M_X = M_X^{(0)} + \alpha M_\pi^2 + \beta M_K^2 + \text{higher orders}$ 
  - Leading order sufficient for  $M_K^2$  dependence
  - We include higher order term in  $M_\pi^2$ 
    - Next order  $\chi$ PT (around  $M_\pi^2 = 0$ ):  $\propto M_\pi^3$
    - Taylor expansion (around  $M_\pi^2 \neq 0$ ):  $\propto M_\pi^4$

Both procedures fine  $\rightarrow$  systematic error  
No sensitivity to any order beyond these
- Vector mesons: higher orders not significant
- Baryons: higher orders significant
  - Restrict fit range to further estimate systematics:
    - full range,  $M_\pi < 550/450\text{MeV}$

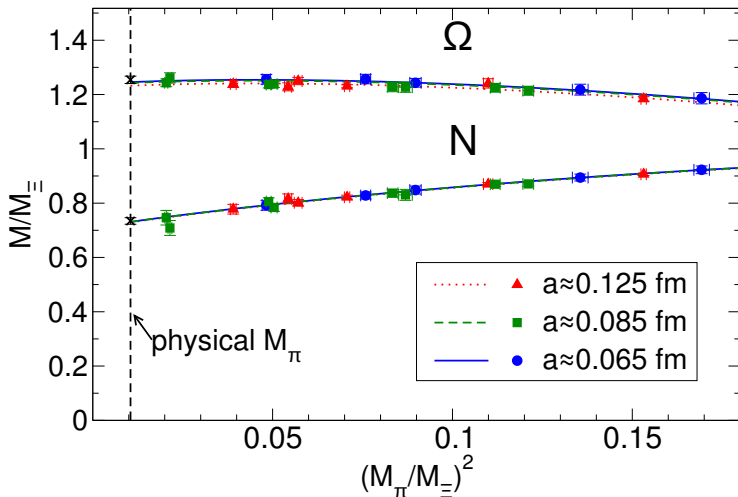
We use all 3 ranges  $\rightarrow$  systematic error

## Chiral fit





# Chiral fit using ratios



# Continuum extrapolation

## Goal:

- Eliminate discretization effects

## Method:

- Formally in our action:  $O(\alpha_s a)$  and  $O(a^2)$
  - Discretization effects are tiny
    - Not possible to distinguish between  $O(a)$  and  $O(a^2)$
- include both in systematic error

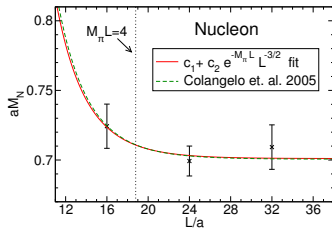
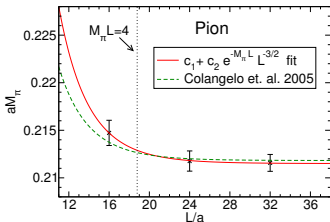
# Finite volume effects from virtual pions

## Goal:

- Eliminate virtual pion finite  $V$  effects
  - Hadrons see mirror charges
  - Exponential in lightest particle (pion) mass

## Method:

- Best practice: use large  $V$ 
  - Rule of thumb:  $M_\pi L \gtrsim 4$
  - Leading effects  $\frac{M_X(L) - M_X}{M_X} = cM_\pi^{1/2} L^{-3/2} e^{M_\pi L}$  (Colangelo et. al., 2005)



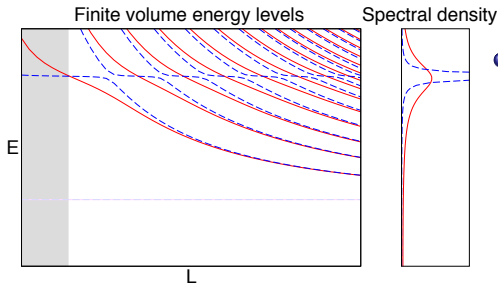
# Finite volume effects in resonances

## Goal:

- Eliminate spectrum distortions from resonances mixing with scattering states

## Method:

- Stay in region where resonance is ground state
  - Otherwise no sensitivity to resonance mass in ground state



- Treatment as scattering problem

(Lüscher, 1985-1991)

- Parameters: mass and coupling (width)
- Alternative approaches suggested

# Systematic uncertainties

## Goal:

- Accurately estimate total systematic error

## Method:

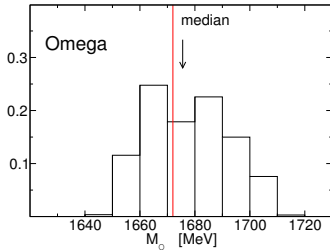
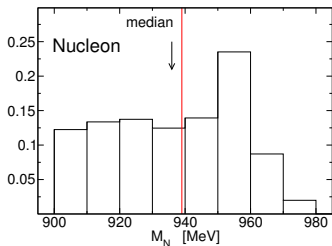
- We account for all the above mentioned effects
- When there are a number of sensible ways to proceed, we take them: Complete analysis for each of
  - 18 fit range combinations
  - ratio/nonratio fits ( $r_X$  resp.  $M_X$ )
  - $O(a)$  and  $O(a^2)$  discretization terms
  - NLO  $\chi$ PT  $M_\pi^3$  and Taylor  $M_\pi^4$  chiral fit
  - 3  $\chi$  fit ranges for baryons:  $M_\pi < 650/550/450$  MeV

resulting in 432 (144) predictions for each baryon (vector meson) mass with each 2000 bootstrap samples for each  $\Xi$  and  $\Omega$  scale setting

# Systematic uncertainties II

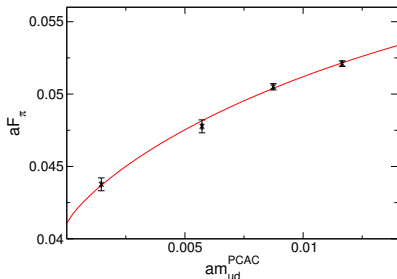
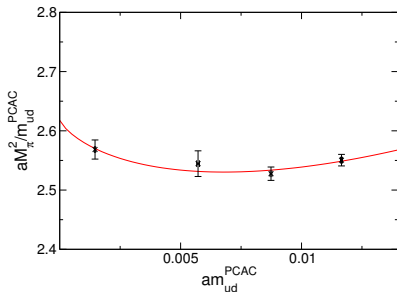
## Method (ctd.):

- Weigh each of the 432 (144) central values by fit quality  $Q$ 
  - Median of this distribution  $\rightarrow$  final result
  - Central 68%  $\rightarrow$  systematic error
- Statistical error from bootstrap of the medians



# Chiral interpolation

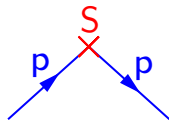
- Simultaneous fit to NLO  $SU(2)$   $\chi$ PT (Gasser, Leutwyler, 1984)
- Consistent for  $M_\pi \lesssim 400$  MeV



- We use 2 safe chiral interpolation ranges:  
 $M_\pi < 340, 380$  MeV
- We use  $SU(2)$   $\chi$ PT and Taylor interpolation forms

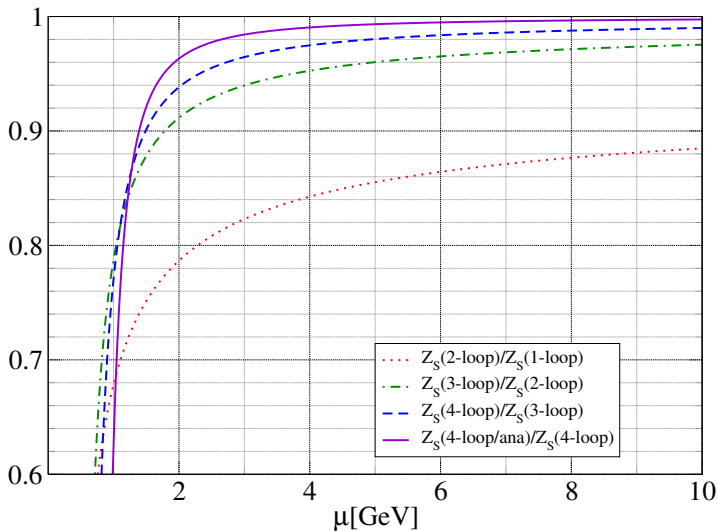
# Renormalization strategy

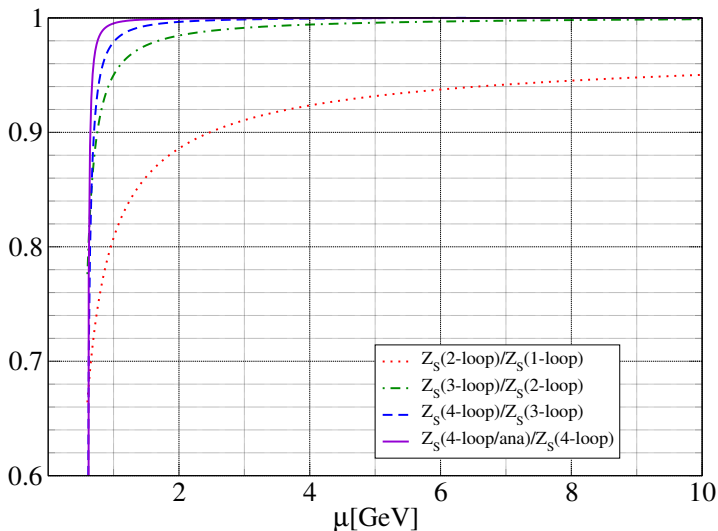
- **Goal:**
  - Full nonperturbative renormalization
  - Optional accurate conversion to perturbative scheme
- **Method:**
  - We use RI-MOM scheme (Martinelli et. al., 1993)
    - $O(a)$  correction (Maillard, Niedermayer, 2008)
  - Compute  $m_q$  at low scale  $\mu \ll 2\pi/a \sim 11 - 24$  GeV
    - $\mu = 2.1$  GeV
    - $\mu = 1.3$  GeV
  - Do continuum non-perturbative running to high scale  $\mu' \gg \Lambda_{\text{QCD}}$
  - Further conversion in 4-loop PT



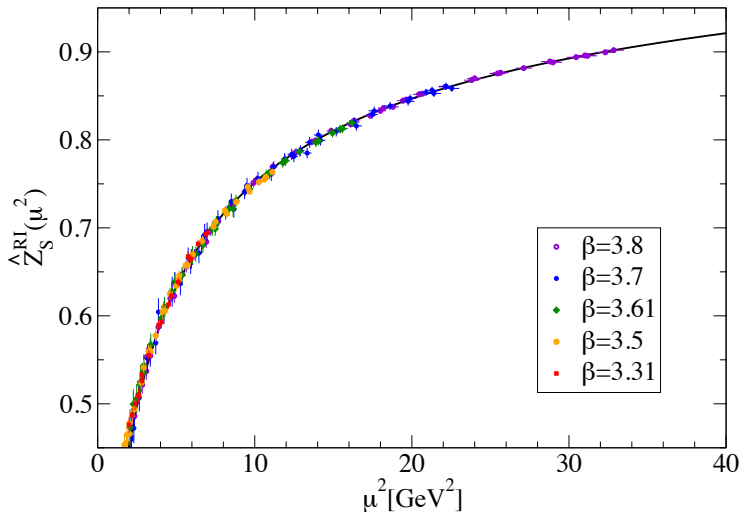


# Desired scale in RI-MOM scheme

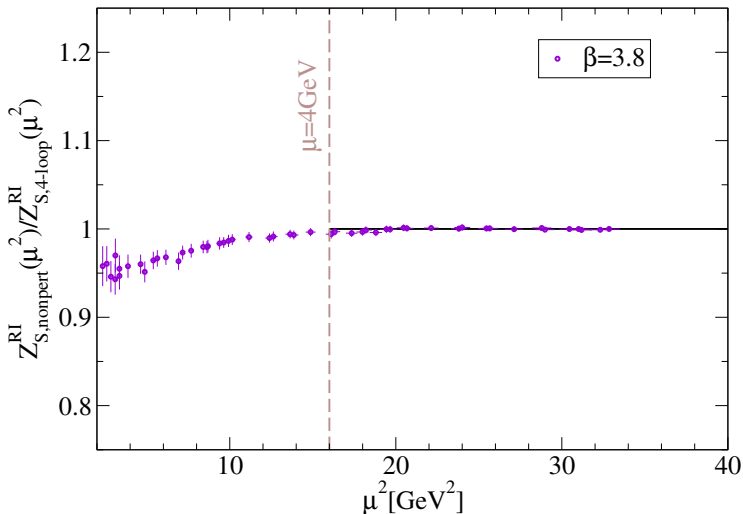


Optional conversion to  $\overline{\text{MS}}$ 

# Nonperturbative running



# Reaching the perturbative regime



# Quark mass definitions

- Lagrangian mass  $m^{\text{bare}}$

- $m^{\text{ren}} = \frac{1}{Z_S} (m^{\text{bare}} - m_{\text{crit}}^{\text{bare}})$

Better use

- $d = m_s^{\text{bare}} - m_{ud}^{\text{bare}}$

- $d^{\text{ren}} = \frac{1}{Z_S} d$

- $m_s^{\text{ren}} = \frac{1}{Z_S} \frac{rd}{r-1}$

- $m^{\text{PCAC}}$  from  $\frac{\langle \partial_0 A_0(t) P(0) \rangle}{\langle P(t) P(0) \rangle}$

- $m^{\text{ren}} = \frac{Z_A}{Z_P} m^{\text{PCAC}}$

- $r = m_s^{\text{PCAC}} / m_{ud}^{\text{PCAC}}$

- $r^{\text{ren}} = r$

and reconstruct

- $m_{ud}^{\text{ren}} = \frac{1}{Z_S} \frac{d}{r-1}$

✓ No additive mass renormalization and ambiguity in  $m_{\text{crit}}$

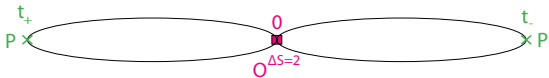
✓ Only  $Z_S$  multiplicative renormalization (no pion poles)

👉 Works with  $O(a)$  improvement (we use this)

## Observable

Matrix element of the effective weak operator  $\langle \bar{K} | O | K \rangle$ :

$$\langle P^\dagger(t_+) O(0) P(t_-) \rangle \xrightarrow{t_\pm \rightarrow \pm\infty} \frac{|\langle K | P | 0 \rangle|^2}{(2M_K)^2} \langle \bar{K} | O | K \rangle e^{-M_K(t_+ - t_-)}$$



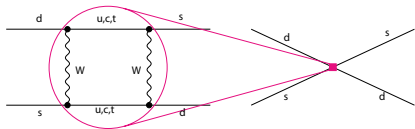
where

$$P = [\bar{s}d]_P = \bar{s}\gamma_5 d$$

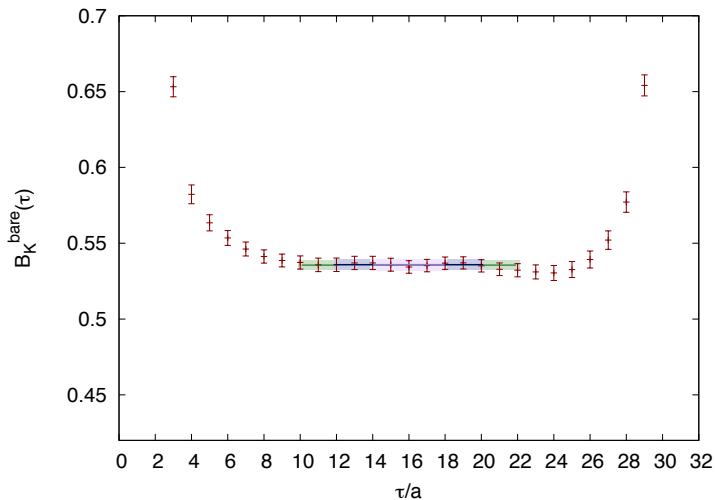
$$O_{\Delta S=2} = [\bar{s}d]_{V-A} [\bar{s}d]_{V-A}$$

Norm from:

$$\langle P^\dagger(t) P(0) \rangle \xrightarrow{t \rightarrow \infty} \frac{|\langle K | P | 0 \rangle|^2}{2M_K} e^{-M_K t}$$



# Signal



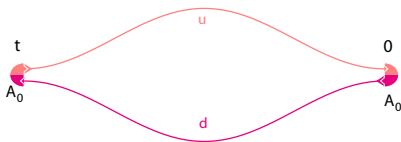
# Observable

With the **axial vector current**

$$A_\mu(t) = \sum_{\vec{x}} \left( \bar{\Psi}^d \gamma_\mu \gamma_5 \Psi^u \right) (\vec{x}, t)$$

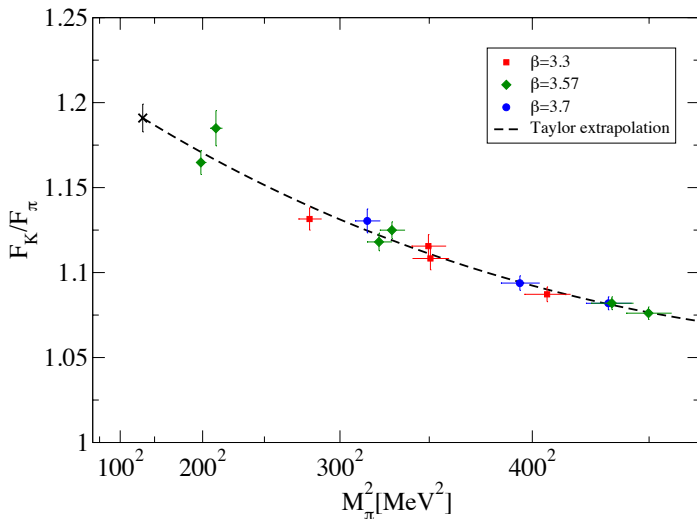
one obtains

$$\langle A_0^\dagger(t) A_0(0) \rangle \xrightarrow{t \rightarrow \infty} \frac{|\langle \pi | A_0 | 0 \rangle|^2}{2M_\pi} e^{-M_\pi t} = \frac{M_\pi^2 F_\pi^{\text{bare}2}}{2M_\pi} e^{-M_\pi t}$$

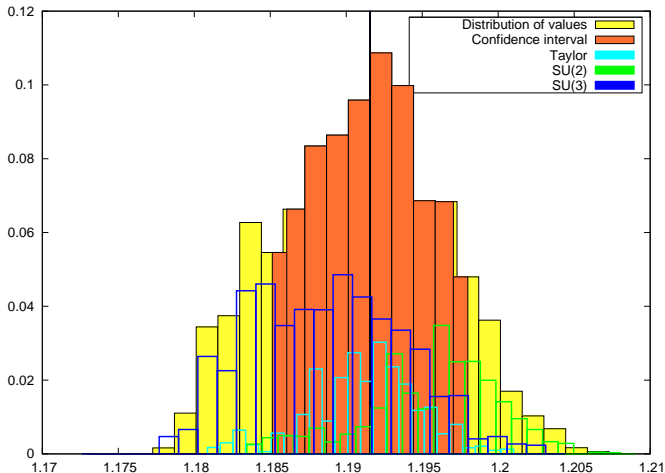




# Chiral extrapolation

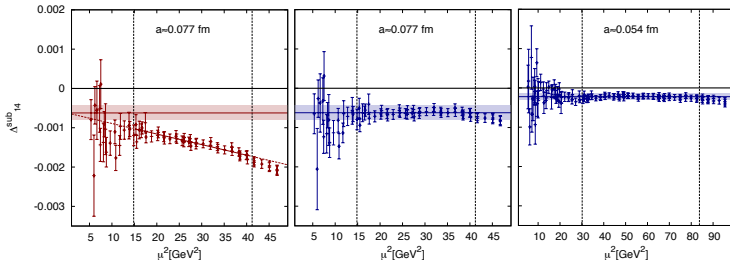
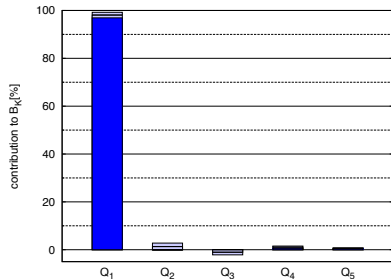


# Errors

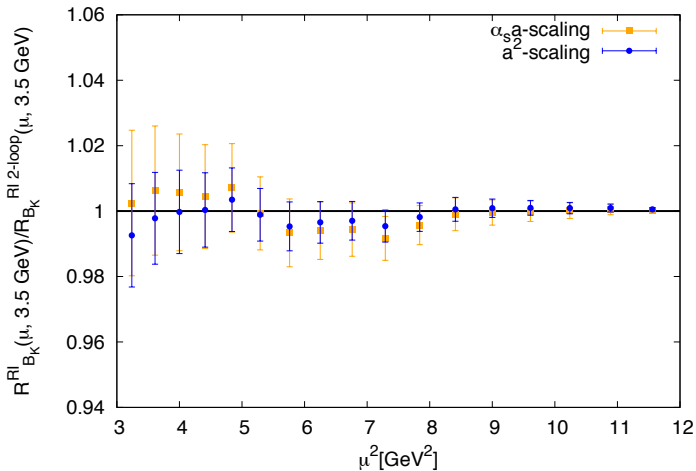


# Unphysical operator mixing

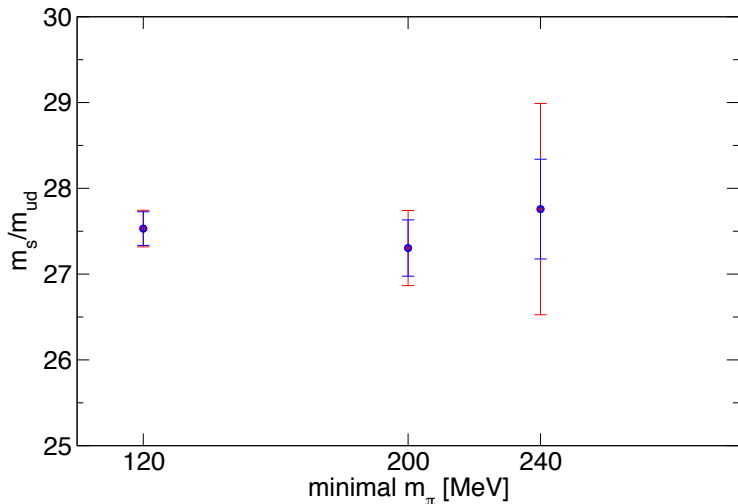
- ☞  $\chi$ SB induces mixing with 4 unphysical operators
- ☞ Mixing terms chirally enhanced
- ✓ Small even below physical  $m_\pi$
- ✓ Good chirality of our action



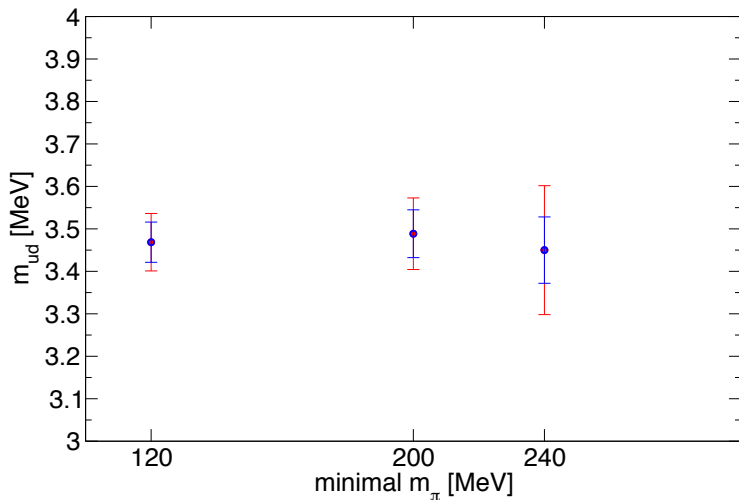
# Running



# Chiral cuts



# Chiral cuts



# Chiral cuts

