

# Mass hierarchy and physics beyond the Standard Model

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New Frontiers in Physics ICNFP 2013

Kolymbari, Crete, Greece, 28 August - 5 September 2013

- Low energy SUSY and 126 GeV Higgs
- Live with the hierarchy
- Low scale strings and extra dimensions

## $H^0$ (Higgs Boson)

particle  
listing

The observed signal is called a Higgs Boson in the following, although its detailed properties and in particular the role that the new particle plays in the context of electroweak symmetry breaking need to be further clarified. The signal was discovered in searches for a Standard Model (SM)-like Higgs. See the following section for mass limits obtained from those searches.

### $H^0$ MASS

<u>VALUE (GeV)</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>125.9±0.4 OUR AVERAGE</b>			
125.8±0.4±0.4	<sup>1</sup> CHATRCHYAN 13J	CMS	<i>pp</i> , 7 and 8 TeV
126.0±0.4±0.4	<sup>2</sup> AAD 12AI	ATLS	<i>pp</i> , 7 and 8 TeV
● ● ● We do not use the following	data for averages, fits, limits, etc. ● ● ●		
126.2±0.6±0.2	<sup>3</sup> CHATRCHYAN 13J	CMS	<i>pp</i> , 7 and 8 TeV
125.3±0.4±0.5	<sup>4</sup> CHATRCHYAN 12N	CMS	<i>pp</i> , 7 and 8 TeV

[HTTP://PDG.LBL.GOV](http://pdg.lbl.gov)

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Created: 7/31/2013

summary tables

$H^0$  Mass  $m = 125.9 \pm 0.4$  GeV

$H^0$  signal strengths in different channels  $[\mu]$

Combined Final States =  $1.07 \pm 0.26$  (S = 1.4)

$W W^*$  Final State =  $0.88 \pm 0.33$  (S = 1.1)

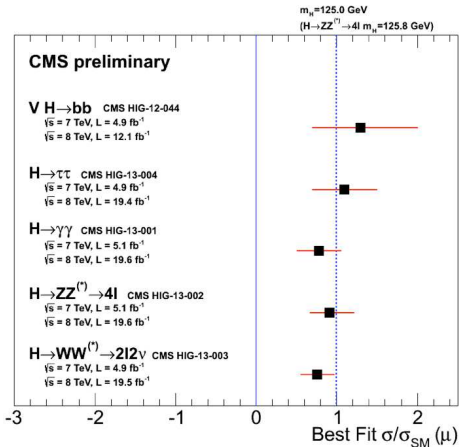
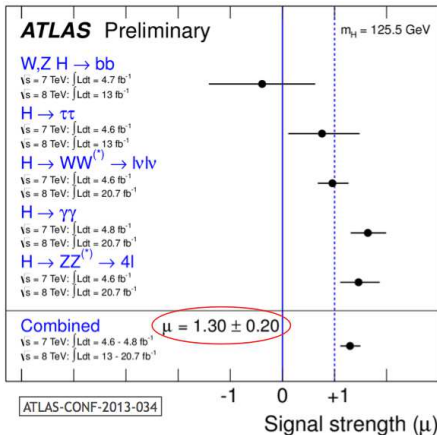
$Z Z^*$  Final State =  $0.89^{+0.30}_{-0.25}$

$\gamma\gamma$  Final State =  $1.65 \pm 0.33$

$b\bar{b}$  Final State =  $0.5^{+0.8}_{-0.7}$

$\tau^+\tau^-$  Final State =  $0.1 \pm 0.7$

# Couplings of the new boson vs SM



exclusion : spin 2 and pseudoscalar at  $\gtrsim 95\%$  CL

Agreement with Standard Model expectation at  $\sim 2\sigma$

# Beyond the Standard Theory of Particle Physics: driven by the mass hierarchy problem

Standard picture: low energy supersymmetry

## Advantages:

- natural elementary scalars
- gauge coupling unification
- LSP: natural dark matter candidate
- radiative EWSB

## Problems:

- too many parameters: soft breaking terms
- MSSM : already a % - %<sub>00</sub> fine-tuning 'little' hierarchy problem

Natural framework: Heterotic string (or high-scale M/F) theory

ATLAS SUSY Searches\* - 95% CL Lower Limits (Status: March 26, 2013)

ATLAS  
Preliminary

Inclusive searches

3rd gen. gluino mediated

3rd gen. squarks direct production

EW direct

Long-lived particles

RPV

- MSUGRA/CMSSM : 0 lep + j's + E
- MSUGRA/CMSSM : 1 lep + j's + E
- Pheno model : 0 lep + j's + E
- Pheno model : 0 lep + j's + E
- Glauino med.  $\tilde{\chi}^0$  ( $\tilde{g} \rightarrow q\tilde{\chi}^0$ ) : 1 lep + j's + E
- GMSB ( $\tilde{l}$  NLSP) : 2 lep (OS) + j's + E
- GMSB ( $\tilde{\tau}$  NLSP) : 1-2  $\tau$  + j's + E
- GGM (bino NLSP) :  $\gamma\gamma$  + E
- GGM (wino NLSP) :  $\gamma$  + lep + E
- GGM (higgsino-bino NLSP) :  $\gamma$  + b + E
- GGM (higgsino NLSP) : Z + jets + E
- Gravitino LSP : 'monojet' + E
- $\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}^0$  : 0 lep + 3 b-jets + E
- $\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}^0$  : 2 SS-lep + (0-3b- $\tilde{\chi}$ )'s + E
- $\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}^0$  : 0 lep + multi-j's + E
- $\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}^0$  : 0 lep + 3 b-jets + E
- $\tilde{b}\tilde{b}, \tilde{b}_1 \rightarrow b\tilde{b}\tilde{\chi}^0$  : 0 lep + 2-b-jets + E
- $\tilde{b}\tilde{b}, \tilde{b}_1 \rightarrow t\tilde{b}\tilde{\chi}^0$  : 2 SS-lep + (0-3b- $\tilde{\chi}$ )'s + E
- $\tilde{t}\tilde{t}$  (light),  $\tilde{t}_1 \rightarrow b\tilde{t}\tilde{\chi}^0$  : 1/2 lep + (b-jet) + E
- $\tilde{t}\tilde{t}$  (medium),  $\tilde{t}_1 \rightarrow b\tilde{t}\tilde{\chi}^0$  : 1 lep + b-jet + E
- $\tilde{t}\tilde{t}$  (heavy),  $\tilde{t}_1 \rightarrow b\tilde{t}\tilde{\chi}^0$  : 2 lep + E
- $\tilde{t}\tilde{t}$  (heavy),  $\tilde{t}_1 \rightarrow t\tilde{t}\tilde{\chi}^0$  : 1 lep + b-jet + E
- $\tilde{t}\tilde{t}$  (heavy),  $\tilde{t}_1 \rightarrow t\tilde{t}\tilde{\chi}^0$  : 0 lep + 6(2b)-jets + E
- $\tilde{t}\tilde{t}$  (natural GMSB) : Z(-ll) + b-jet + E
- $\tilde{t}_1 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 Z$  : Z(-ll) + 1 lep + b-jet + E
- $\tilde{t}_1 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 \tilde{b}\tilde{\chi}^0$  : 2 lep + E
- $\tilde{\chi}_2^0 \tilde{\chi}_2^0 \rightarrow \tilde{\nu}(\tilde{\nu})$  : 2 lep + E
- $\tilde{\chi}_2^0 \tilde{\chi}_2^0 \rightarrow \tilde{\nu}(\tilde{\nu})\tilde{\nu}(\tilde{\nu})$  : 2  $\tau$  + E
- $\tilde{\chi}_1^0 \tilde{\chi}_2^0 \rightarrow \tilde{\nu}(\tilde{\nu})\tilde{\nu}(\tilde{\nu})$ ,  $\tilde{\nu}(\tilde{\nu})\tilde{\nu}(\tilde{\nu})$  : 3 lep + E
- $\tilde{\chi}_1^0 \tilde{\chi}_2^0 \rightarrow W^+ \tilde{\chi}_1^+ Z^0 \tilde{\chi}_1^0$  : 3 lep + E
- Direct  $\tilde{\chi}_1^0$  pair prod. (AMSB) : long-lived  $\tilde{\chi}_1^0$
- Stable  $\tilde{g}$ , R-hadrons : low  $\beta$ , low GMSB, stable  $\tilde{\tau}$ , low  $\beta$
- GMSB,  $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$  : non-pointing photons
- $\tilde{\chi}_1^0 \rightarrow q\bar{q}$  (RPV) :  $\mu$  + heavy displaced vertex
- LFV :  $pp \rightarrow \tilde{\nu} + X, \tilde{\nu}_e \rightarrow e + \text{resonance}$
- LFV :  $pp \rightarrow \tilde{\nu} + X, \tilde{\nu}_e \rightarrow e(\mu) + \tau$  resonance
- Bilinear RPV CMSSM : 1 lep + 7 j's + E
- $\tilde{\chi}_1^0 \tilde{\chi}_2^0 \rightarrow W^0_{\mu\nu} \tilde{\chi}_1^0 \rightarrow e\tilde{\nu}_\mu e\tilde{\nu}_\mu$  : 4 lep + E
- $\tilde{\chi}_1^0 \tilde{\chi}_2^0 \rightarrow W^0_{\mu\nu} \tilde{\chi}_1^0 \rightarrow \tilde{\nu}_\mu \tilde{\nu}_\mu, e\tilde{\nu}_\mu$  : 3 lep + 1  $\tau$  + E
- $\tilde{g} \rightarrow q\bar{q}$  : 3-jet resonance pair
- $\tilde{g} \rightarrow t\bar{t}$ ,  $\tilde{t} \rightarrow bs$  : 2 SS-lep + (0-3b- $\tilde{\chi}$ )'s + E
- Scalar gluon : 2-jet resonance pair
- WIMP interaction (D5, Dirac  $\tilde{\chi}$ ) : 'monojet' + E

L45.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-148]	1.80 TeV	$\tilde{g} = \tilde{g}$ mass	
L45.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-144]	1.24 TeV	$\tilde{g} = \tilde{g}$ mass	
L45.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-139]	1.18 TeV	$\tilde{g}$ mass ( $m(\tilde{g}) < 2 \text{ TeV}, \text{light } \tilde{\chi}_1^0$ )	
L45.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-139]	1.28 TeV	$\tilde{g}$ mass ( $m(\tilde{g}) < 2 \text{ TeV}, \text{light } \tilde{\chi}_1^0$ )	
L4.7 fb <sup>-1</sup> , 7 TeV [1208.4688]	900 GeV	$\tilde{g}$ mass ( $m(\tilde{g}) < 200 \text{ GeV}, m(\tilde{\chi}_2^0) = \frac{1}{2}m(\tilde{\chi}_1^0) + m(\tilde{g})$ )	
L4.7 fb <sup>-1</sup> , 7 TeV [1208.4688]	1.24 TeV	$\tilde{g}$ mass ( $\tan\beta < 15$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [1210.1314]	1.40 TeV	$\tilde{g}$ mass ( $\tan\beta > 18$ )	
L4.8 fb <sup>-1</sup> , 7 TeV [1209.0753]	1.07 TeV	$\tilde{g}$ mass ( $m(\tilde{g}) > 50 \text{ GeV}$ )	
L4.8 fb <sup>-1</sup> , 7 TeV [ATLAS-CONF-2012-144]	619 GeV	$\tilde{g}$ mass	
L4.8 fb <sup>-1</sup> , 7 TeV [1211.1167]	900 GeV	$\tilde{g}$ mass ( $m(\tilde{\chi}_2^0) < 220 \text{ GeV}$ )	
L25.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-132]	890 GeV	$\tilde{g}$ mass ( $m(\tilde{g}) > 200 \text{ GeV}$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-147]	845 GeV	$F^2$ scale ( $m(\tilde{g}) > 10^4 eV$ )	
L12.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-145]	1.24 TeV	$\tilde{g}$ mass ( $m(\tilde{g}_1) < 200 \text{ GeV}$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-007]	900 GeV	$\tilde{g}$ mass (any $m(\tilde{\chi}_2^0)$ )	
L45.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-143]	1.00 TeV	$\tilde{g}$ mass ( $m(\tilde{g}_1) < 300 \text{ GeV}$ )	
L12.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-145]	1.18 TeV	$\tilde{g}$ mass ( $m(\tilde{g}_1) < 200 \text{ GeV}$ )	
L12.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-105]	620 GeV	$\tilde{b}$ mass ( $m(\tilde{g}_1) < 120 \text{ GeV}$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-007]	430 GeV	$\tilde{b}$ mass ( $m(\tilde{g}_1) = 2m(\tilde{g}_2)$ )	
L4.7 fb <sup>-1</sup> , 7 TeV [1208.4395, 1209.2162]	167 GeV	$\tilde{t}$ mass ( $m(\tilde{g}_1) = 55 \text{ GeV}$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-037]	160-410 GeV	$\tilde{t}$ mass ( $m(\tilde{g}_1) = 0 \text{ GeV}, m(\tilde{g}_2) = 150 \text{ GeV}$ )	
L13.0 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-167]	160-640 GeV	$\tilde{t}$ mass ( $m(\tilde{g}_1) = 0 \text{ GeV}, m(\tilde{g}_2) = 10 \text{ GeV}$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-037]	200-610 GeV	$\tilde{t}$ mass ( $m(\tilde{g}_1) = 0$ )	
L20.5 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-024]	320-660 GeV	$\tilde{t}$ mass ( $m(\tilde{g}_1) = 0$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-025]	500 GeV	$\tilde{t}$ mass ( $m(\tilde{g}_1) > 150 \text{ GeV}$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-025]	520 GeV	$\tilde{t}_2$ mass ( $m(\tilde{g}_1) = m(\tilde{g}_2) + 180 \text{ GeV}$ )	
L4.7 fb <sup>-1</sup> , 7 TeV [1208.2894]	85-195 GeV	$\tilde{t}$ mass ( $m(\tilde{g}_1) = 0$ )	
L4.7 fb <sup>-1</sup> , 7 TeV [1208.2894]	116-340 GeV	$\tilde{\chi}_2^0$ mass ( $m(\tilde{g}_1) < 10 \text{ GeV}, m(\tilde{\nu}_\tau) = \frac{1}{2}m(\tilde{g}_1) + m(\tilde{\chi}_2^0)$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-028]	180-330 GeV	$\tilde{\chi}_2^0$ mass ( $m(\tilde{g}_1) < 10 \text{ GeV}, m(\tilde{\nu}_\tau) = \frac{1}{2}m(\tilde{g}_1) + m(\tilde{\chi}_2^0)$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-035]	600 GeV	$\tilde{\chi}_2^0$ mass ( $m(\tilde{g}_1) = m(\tilde{g}_2), m(\tilde{g}_2) = 0, m(\tilde{\nu}_\tau) \text{ as above}$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-035]	315 GeV	$\tilde{\chi}_2^0$ mass ( $m(\tilde{g}_1) = m(\tilde{g}_2), m(\tilde{g}_2) = 0$ , sleptons decoupled)	
L4.7 fb <sup>-1</sup> , 7 TeV [1210.2852]	220 GeV	$\tilde{\chi}_1^0$ mass ( $1 < v(\tilde{\chi}_1^0) < 10 \text{ ns}$ )	
L4.7 fb <sup>-1</sup> , 7 TeV [1211.1597]	985 GeV	$\tilde{g}$ mass	
L4.7 fb <sup>-1</sup> , 7 TeV [1211.1597]	300 GeV	$\tilde{\tau}$ mass ( $5 < \tan\beta < 20$ )	
L4.7 fb <sup>-1</sup> , 7 TeV [ATLAS-CONF-2013-016]	230 GeV	$\tilde{\chi}_1^0$ mass ( $0.4 < v(\tilde{\chi}_1^0) < 2 \text{ ns}$ )	
L4.8 fb <sup>-1</sup> , 7 TeV [1210.1251]	790 GeV	$\tilde{q}$ mass (1 mm $< c\tau < 1 \text{ m}, \tilde{g}$ decoupled)	
L4.8 fb <sup>-1</sup> , 7 TeV [1212.1272]	1.61 TeV	$\tilde{\nu}_\tau$ mass ( $\lambda_{311} = 0.10, \lambda_{322} = 0.05$ )	
L4.8 fb <sup>-1</sup> , 7 TeV [1212.1272]	1.10 TeV	$\tilde{\nu}_\tau$ mass ( $\lambda_{311} = 0.10, \lambda_{322} = 0.05$ )	
L4.7 fb <sup>-1</sup> , 7 TeV [ATLAS-CONF-2012-140]	1.2 TeV	$\tilde{g}$ mass ( $c\tau_{333} < 1 \text{ mm}$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-036]	760 GeV	$\tilde{\chi}_1^0$ mass ( $m(\tilde{g}_1) > 300 \text{ GeV}, \lambda_{333} > 0$ )	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-036]	350 GeV	$\tilde{\chi}_1^0$ mass ( $m(\tilde{g}_1) > 80 \text{ GeV}, \lambda_{333} > 0$ )	
L4.8 fb <sup>-1</sup> , 7 TeV [1210.4813]	666 GeV	$\tilde{g}$ mass	
L20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-007]	880 GeV	$\tilde{g}$ mass (any $m(\tilde{g}_1)$ )	
L4.8 fb <sup>-1</sup> , 7 TeV [1210.4826]	100-287 GeV	sgluon mass (incl. limit from 1110.2693)	
L19.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-141]	794 GeV	$M^*$ scale ( $m_* < 80 \text{ GeV}$ , limit of $< 687 \text{ GeV}$ for D8)	

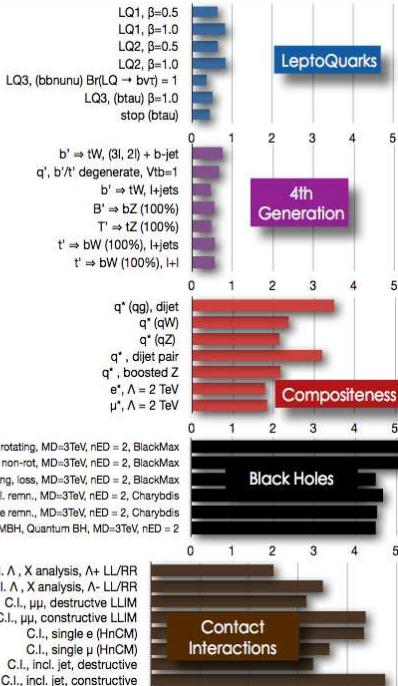
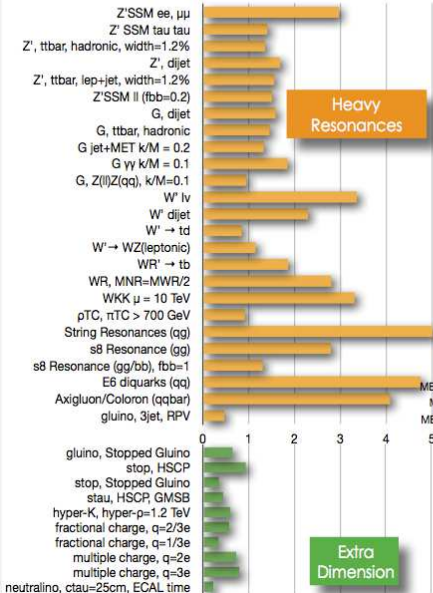
$\int L dt = (4.4 - 20.7) \text{ fb}^{-1}$   
 $\sqrt{s} = 7, 8 \text{ TeV}$

8 TeV, all 2012 data  
8 TeV, partial 2012 data  
7 TeV, all 2011 data

\*Only a selection of the available mass limits on new states or phenomena shown. All limits quoted are observed minus 1  $\sigma$  theoretical signal cross section uncertainty

# CMS EXOTICA

95% CL EXCLUSION LIMITS (TeV)



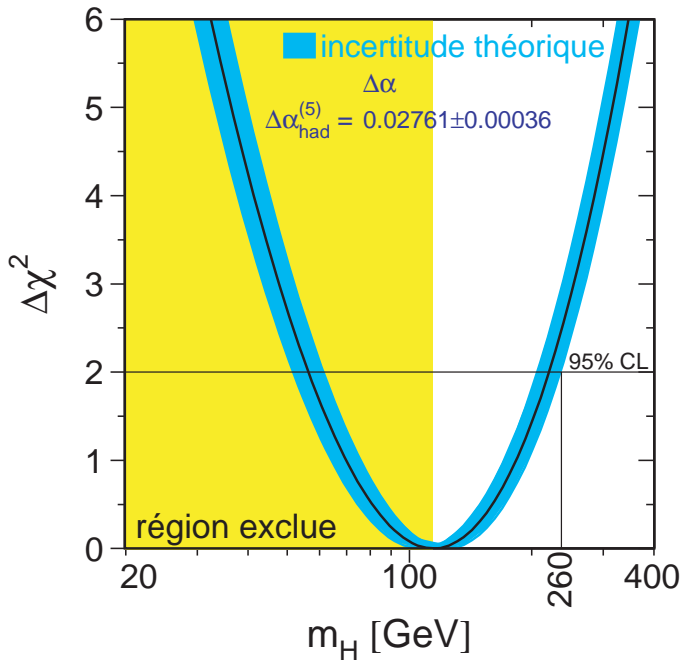
# Remarks on the value of the Higgs mass $\sim 126$ GeV

- consistent with expectation from precision tests of the SM
- favors perturbative physics      quartic coupling  $\lambda = m_H^2/v^2 \simeq 1/8$
- 1st elementary scalar in nature signaling perhaps more to come

## Window to new physics ?

- compatible with supersymmetry  
but appears fine-tuned in its minimal version [10]  
early to draw a general conclusion before LHC13/14 [11]  
e.g. an extra singlet or split families can alleviate the fine tuning [12]
- very important to measure its properties and couplings  
any deviation of its couplings to top, bottom and EW gauge bosons  
implies new light states involved in the EWSB altering the fine-tuning





# Fine-tuning in MSSM

Upper bound on the lightest scalar mass:

$$m_h^2 \lesssim m_Z^2 \cos^2 2\beta + \frac{3}{(4\pi)^2} \frac{m_t^4}{v^2} \left[ \ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{A_t^2}{m_{\tilde{t}}^2} \left( 1 - \frac{A_t^2}{12m_{\tilde{t}}^2} \right) \right] \lesssim (130 \text{ GeV})^2$$

$$m_h \simeq 126 \text{ GeV} \Rightarrow m_{\tilde{t}} \simeq 3 \text{ TeV or } A_t \simeq 3m_{\tilde{t}} \simeq 1.5 \text{ TeV}$$

$\Rightarrow$  % to a few ‰ fine-tuning

$$\text{minimum of the potential: } m_Z^2 = 2 \frac{m_1^1 - m_2^2 \tan^2 \beta}{\tan^2 \beta - 1} \sim -2m_2^2 + \dots$$

$$\text{RG evolution: } m_2^2 = m_2^2(M_{\text{GUT}}) - \frac{3\lambda_t^2}{4\pi^2} m_{\tilde{t}}^2 \ln \frac{M_{\text{GUT}}}{m_{\tilde{t}}} + \dots \quad [20]$$

$$\sim m_2^2(M_{\text{GUT}}) - \mathcal{O}(1)m_{\tilde{t}}^2 + \dots \quad [8]$$

# Reduce the fine-tuning

- minimize radiative corrections

$M_{\text{GUT}} \rightarrow \Lambda$  : low messenger scale (gauge mediation)

$$\delta m_{\tilde{t}} = \frac{8\alpha_s}{3\pi} M_3^2 \ln \frac{\Lambda}{M_3} + \dots$$

- extend the MSSM

extra fields beyond LHC reach  $\rightarrow$  effective field theory approach

- ...

# MSSM with dim-5 and 6 operators

I.A.-Dudas-Ghilencea-Tziveloglou '08, '09, '10

parametrize new physics above MSSM by higher-dim effective operators

relevant super potential operators of dimension-5:

$$\mathcal{L}^{(5)} = \frac{1}{M} \int d^2\theta (\eta_1 + \eta_2 S) (H_1 H_2)^2$$

$\eta_1$  : generated for instance by a singlet

$$W = \lambda \sigma H_1 H_2 + M \sigma^2 \quad \rightarrow \quad W_{\text{eff}} = \frac{\lambda^2}{M} (H_1 H_2)^2$$

Strumia '99 ; Brignole-Casas-Espinosa-Navarro '03

Dine-Seiberg-Thomas '07

$\eta_2$  : corresponding soft breaking term      spurion  $S \equiv m_S \theta^2$

# Physical consequences of $MSSM_5$ : Scalar potential

$$\begin{aligned}\mathcal{V} = & m_1^2 |h_1|^2 + m_2^2 |h_2|^2 + B\mu(h_1 h_2 + \text{h.c.}) + \frac{g_2^2 + g_Y^2}{8} (|h_1|^2 - |h_2|^2)^2 \\ & + (|h_1|^2 + |h_2|^2) (\eta_1 h_1 h_2 + \text{h.c.}) + \frac{1}{2} [\eta_2 (h_1 h_2)^2 + \text{h.c.}] \\ & + \eta_1^2 |h_1 h_2|^2 (|h_1|^2 + |h_2|^2)\end{aligned}$$

- $\eta_{1,2} \Rightarrow$  quartic terms along the D-flat direction  $|h_1| = |h_2|$
- tree-level mass can increase significantly
- bigger parameter space for LSP being dark matter

Bernal-Blum-Nir-Losada '09

- last term  $\sim \eta_1^2$  : guarantees stability of the potential

but requires addition of dim-6 operators

# MSSM Higgs with dim-6 operators

**dim-6 operators can have an independent scale from dim-5**

Classification of all dim-6 contributing to the scalar potential

(without SUSY)  $\Rightarrow$

large  $\tan \beta$  expansion:  $\delta_6 m_h^2 = f v^2 + \dots$

constant receiving contributions from several operators

$f \sim f_0 \times (\mu^2/M^2, m_S^2/M^2, \mu m_S/M^2, v^2/M^2)$

$m_S = 1 \text{ TeV}, M = 10 \text{ TeV}, f_0 \sim 1 - 2.5$  for each operator

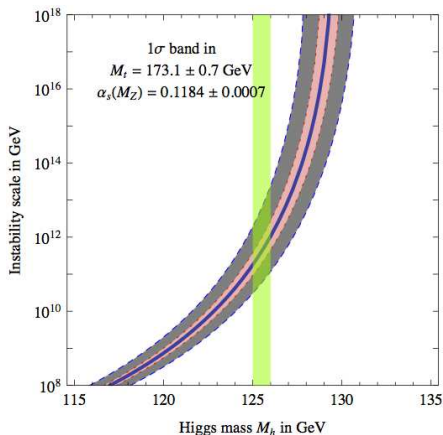
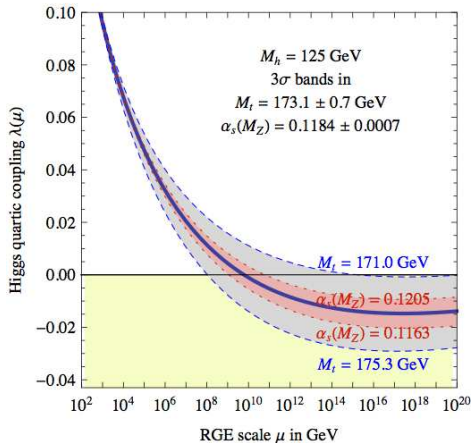
$\Rightarrow m_h \simeq 103 - 119 \text{ GeV}$

$\Rightarrow$  MSSM with dim-5 and dim-6 operators:

possible resolution of the MSSM fine-tuning problem

# Can the SM be valid at high energies?

Degrassi-Di Vita-Elias Miró-Espinosa-Giudice-Isidori-Strumia '12



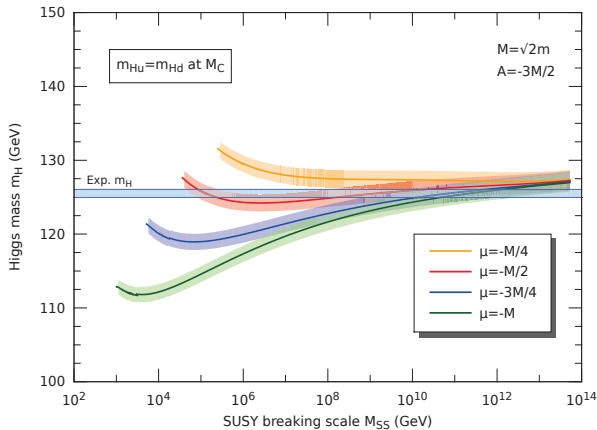
Instability of the SM Higgs potential  $\Rightarrow$  metastability of the EW vacuum

SUSY :  $\lambda = 0 \Rightarrow \sin \beta = 1$

$$H_{SM} = \sin \beta H_u - \cos \beta H_d^* \quad \lambda = \frac{1}{8}(g_2^2 + g'^2) \cos^2 2\beta$$

$\lambda = 0$  at a scale  $\geq 10^{10}$  GeV  $\Rightarrow m_H = 126 \pm 3$  GeV

Ibanez-Valenzuela '13



e.g. for universal  $\sqrt{2}m = M = M_{SS}$ ,  $A = -3/2M$



If the weak scale is tuned  $\Rightarrow$  split supersymmetry is a possibility

Arkani Hamed-Dimopoulos '04, Giudice-Romanino '04

- natural splitting: gauginos, higgsinos carry R-symmetry, scalars do not
- main good properties of SUSY are maintained
  - gauge coupling unification and dark matter candidate
- also no dangerous FCNC, CP violation, ...
- experimentally allowed Higgs mass  $\Rightarrow$  'mini' split [20]

$m_S \sim \text{few - thousands TeV}$

gauginos: a loop factor lighter than scalars ( $\sim m_{3/2}$ )

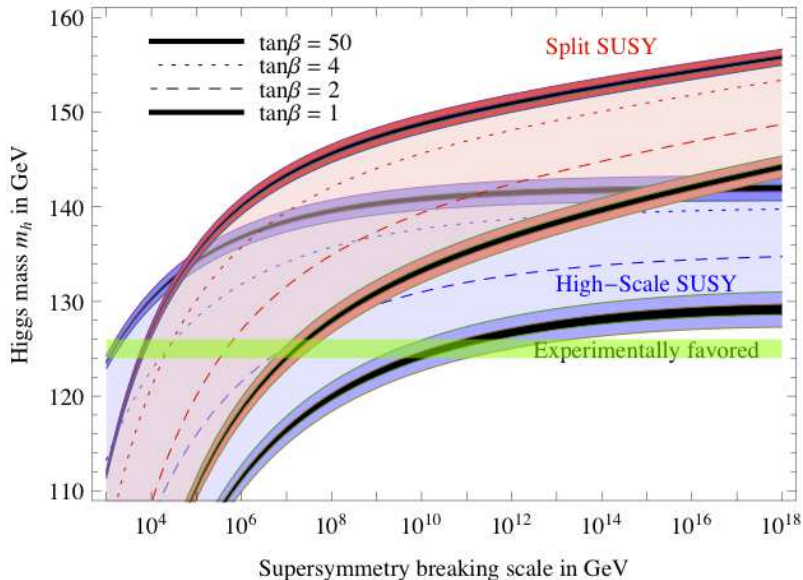
- natural string framework: intersecting (or magnetized) branes

IA-Dimopoulos '04

D-brane stacks are supersymmetric with massless gauginos

intersections have chiral fermions with broken SUSY & massive scalars

## Predicted range for the Higgs mass



# An extra $U(1)$ can also cure the instability problem

Anchordoqui-IA-Goldberg-Huang-Lüst-Taylor-Vlcek '12

usually associated to known global symmetries of the SM:  $B, L, \dots$

- $B$  anomalous and superheavy
- $B - L$  massless at the string scale (no associated 6d anomaly)  
but broken at TeV by a scalar VEV with the quantum numbers of  $N_R$
- $L$ -violation from higher-dim operators suppressed by the string scale
- $U(3)$  unification,  $Y$  combination  $\Rightarrow$  2 parameters: 1 coupling +  $m_Z$
- perturbativity  $\Rightarrow 0.5 \lesssim g_{U(1)_R} \lesssim 1$
- interesting LHC phenomenology and cosmology

## Alternative answer: Low UV cutoff $\Lambda \sim \text{TeV}$

- low scale gravity  $\Rightarrow$  extra dimensions: large flat or warped
- low string scale  $\Rightarrow$  low scale gravity, ultra weak string coupling

$M_s \sim 1 \text{ TeV} \Rightarrow$  volume  $R_{\perp}^n = 10^{32} l_s^n$  ( $R_{\perp} \sim .1 - 10^{-13} \text{ mm}$  for  $n = 2 - 6$ )

- spectacular model independent predictions
- radical change of high energy physics at the TeV scale

Moreover no little hierarchy problem:

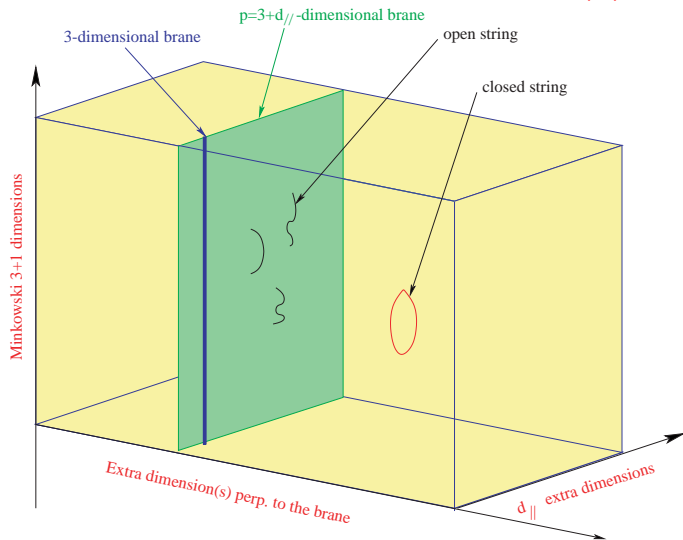
radiative electroweak symmetry breaking with no logs [10]

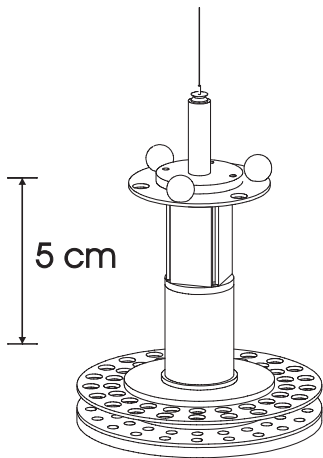
$\Lambda \sim$  a few TeV and  $m_H^2 =$  a loop factor  $\times \Lambda^2$  [23]

But unification has to be probably dropped

New Dark Matter candidates e.g. in the extra dims

- 2 types of compact extra dimensions:
- parallel ( $d_{\parallel}$ ):  $\lesssim 10^{-16}$  cm (TeV)
  - transverse ( $\perp$ ):  $\lesssim 0.1$  mm (meV)





$R_{\perp} \lesssim 45 \mu\text{m}$  at 95% CL

- dark-energy length scale  $\approx 85 \mu\text{m}$

# Accelerator signatures: 4 different scales

- Gravitational radiation in the bulk  $\Rightarrow$  missing energy

present LHC bounds:  $M_* \gtrsim 3 - 5$  TeV

- Massive string vibrations  $\Rightarrow$  e.g. resonances in dijet distribution [25]

$$M_j^2 = M_0^2 + M_s^2 j \quad ; \quad \text{maximal spin : } j + 1$$

higher spin excitations of quarks and gluons with strong interactions

present LHC limits:  $M_s \gtrsim 5$  TeV

- Large TeV dimensions  $\Rightarrow$  KK resonances of SM gauge bosons I.A. '90

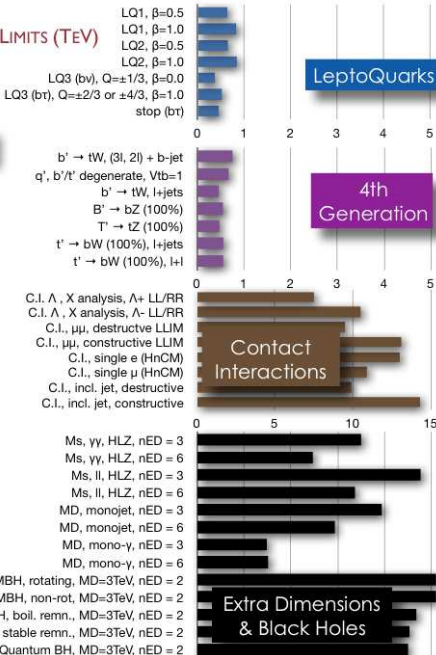
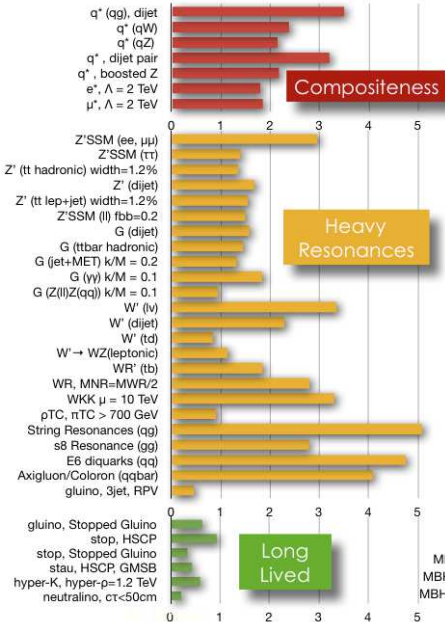
$$M_k^2 = M_0^2 + k^2/R^2 \quad ; \quad k = \pm 1, \pm 2, \dots$$

experimental limits:  $R^{-1} \gtrsim 0.5 - 4$  TeV (UED - localized fermions) [27]

- extra  $U(1)$ 's and anomaly induced terms

masses suppressed by a loop factor from  $M_s$  [30]

# CMS EXOTICA 95% CL EXCLUSION LIMITS (TeV)



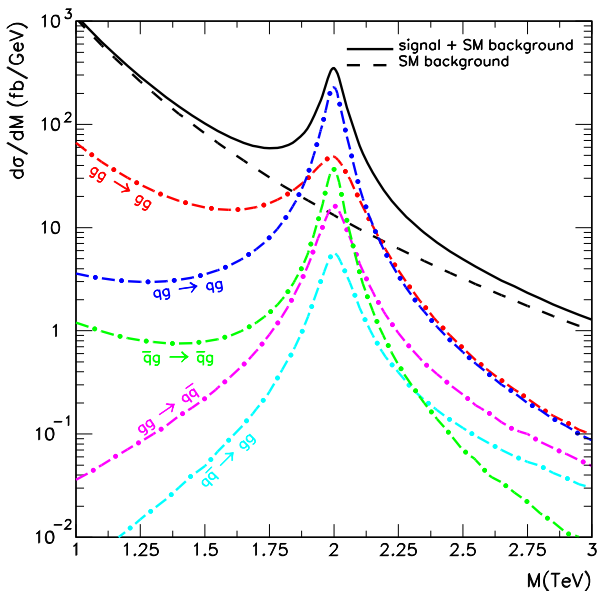


**Universal** deviation  
from Standard Model  
in jet distribution

$M_s = 2$  TeV

Width = 15-150 GeV

Anchordoqui-Goldberg-  
Lüst-Nawata-Taylor-  
Stieberger '08

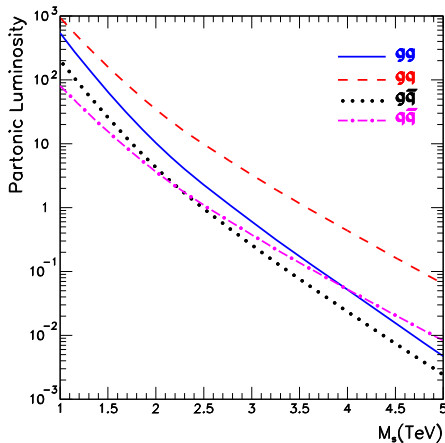


Tree level superstring amplitudes involving at most 2 fermions and gluons:  
 model independent for any compactification, # of susy's, even none  
 no intermediate exchange of KK, windings or graviton emission  
 Universal sum over infinite exchange of string (Regge) excitations

Parton luminosities in pp above TeV  
 are dominated by  $gq$ ,  $gg$

⇒ model independent

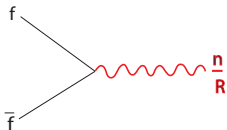
$gq \rightarrow gq$ ,  $gg \rightarrow gg$ ,  $gg \rightarrow q\bar{q}$



## Localized fermions (on 3-brane intersections)

⇒ single production of KK modes

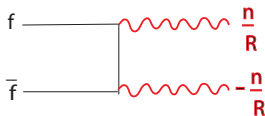
I.A.-Benakli '94



- strong bounds indirect effects
- new resonances but at most  $n = 1$

## Otherwise KK momentum conservation

⇒ pair production of KK modes (universal dims)



- weak bounds
- no resonances
- lightest KK stable ⇒ dark matter candidate

Servant-Tait '02

# UED hadron collider phenomenology

- large rates for KK-quark and KK-gluon production
- cascade decays via KK- $W$  bosons and KK-leptons  
determine particle properties from different distributions
- missing energy from LKP: weakly interacting escaping detection
- phenomenology similar to supersymmetry

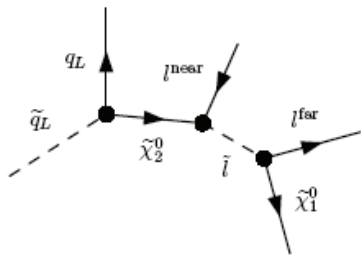
spin determination important for distinguishing SUSY and UED [23]

gluino	1/2	KK-gluon	1
squark	0	KK-quark	1/2
chargino	1/2	KK- $W$ boson	1
slepton	0	KK-lepton	1/2
neutralino	1/2	KK- $Z$ boson	1

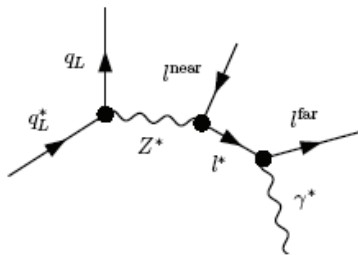
# SUSY vs UED signals at LHC

Example: jet dilepton final state

SUSY



UED



# Extra $U(1)$ 's and anomaly induced terms

masses suppressed by a loop factor

usually associated to known global symmetries of the SM

(anomalous or not) such as (combinations of)

Baryon and Lepton number, or PQ symmetry

Two kinds of massive  $U(1)$ 's:

I.A.-Kiritsis-Rizos '02

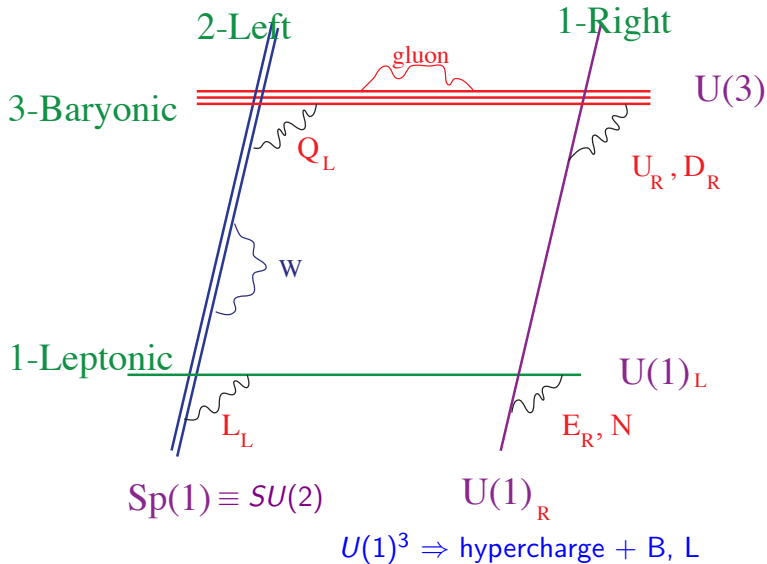
- 4d anomalous  $U(1)$ 's:  $M_A \simeq g_A M_s$

- 4d non-anomalous  $U(1)$ 's: (but masses related to 6d anomalies)

$$M_{NA} \simeq g_A M_s V_2 \leftarrow (6d \rightarrow 4d) \text{ internal space} \Rightarrow M_{NA} \geq M_A$$

or massless in the absence of such anomalies

# Standard Model on D-branes : SM<sup>++</sup>




- $B$  and  $L$  become massive due to anomalies

Green-Schwarz terms

- the global symmetries remain in perturbation

- Baryon number  $\Rightarrow$  proton stability

- Lepton number  $\Rightarrow$  protect small neutrino masses

no Lepton number  $\Rightarrow \frac{1}{M_s} LLHH \rightarrow$  Majorana mass:  $\frac{\langle H \rangle^2}{M_s} LL$   


- $B, L \Rightarrow$  extra  $Z'$ 's

with possible leptophobic couplings leading to CDF-type  $Wjj$  events

$Z' \simeq B$  lighter than 4d anomaly free  $Z'' \simeq B - L$



# Conclusions

- Confirmation of the EWSB scalar at the LHC:  
important milestone of the LHC research program
  - Precise measurement of its couplings is of primary importance
  - Hint on the origin of mass hierarchy and of BSM physics
    - natural or unnatural SUSY?
    - low string scale in some realization?
    - something new and unexpected?
- all options are still open
- LHC enters a new era with possible new discoveries

# The LHC timeline

## LS1 Machine Consolidation

## LS2 Machine upgrades for high Luminosity

- Collimation
- Cryogenics
- Injector upgrade for high intensity (lower emittance)
- Phase I for ATLAS : Pixel upgrade, FTK, and new small wheel

## LS3 Machine upgrades for high Luminosity

- Upgrade interaction region
- Crab cavities?
- Phase II: full replacement of tracker, new trigger scheme (add L0), readout electronics.



*Europe's top priority should be the exploitation of the full potential of the LHC, including the high-luminosity upgrade of the machine and detectors with a view to collecting ten times more data than in the initial design, by around 2030.*

