

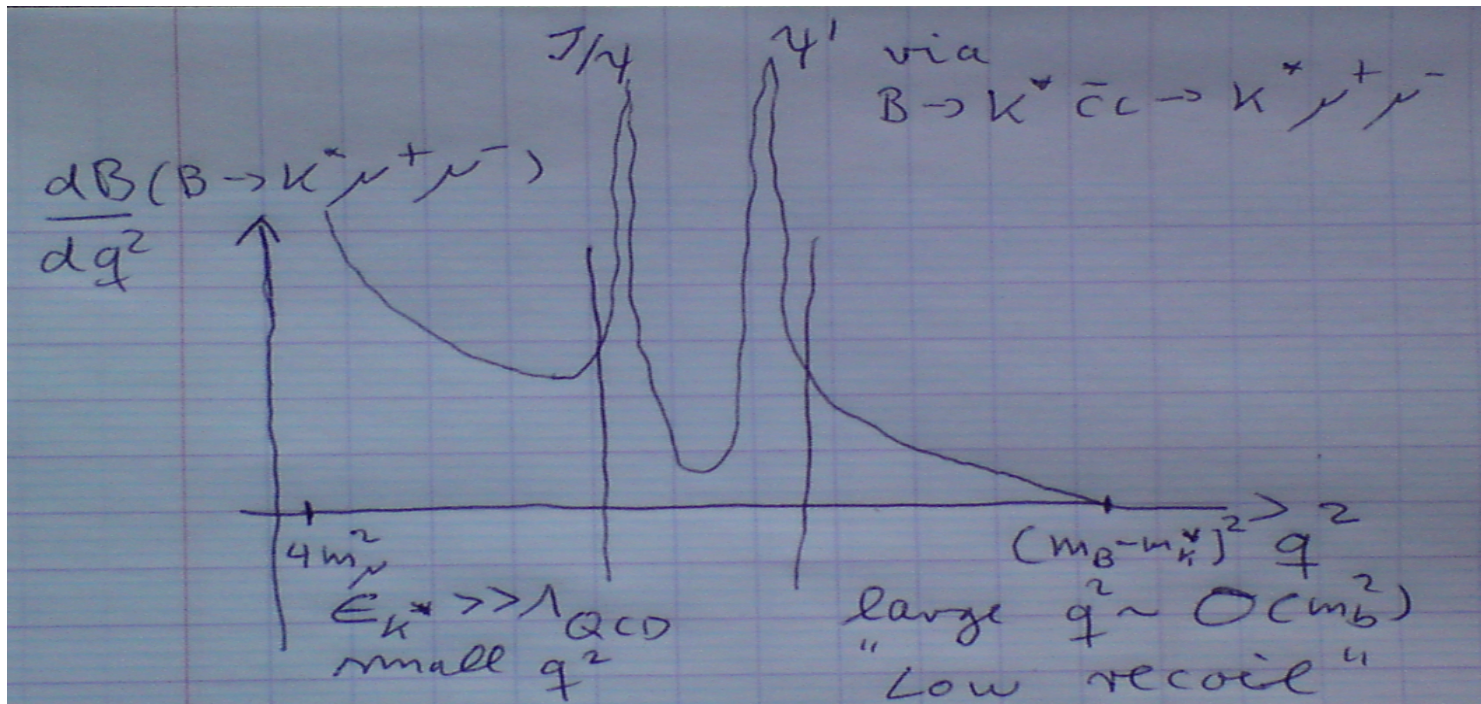
Opportunities & Interplay

(two comments)

Last Talk of the Workshop, September 11, 2012

Gudrun Hiller, Dortmund

Opportunity: Different Kinematics/Theory

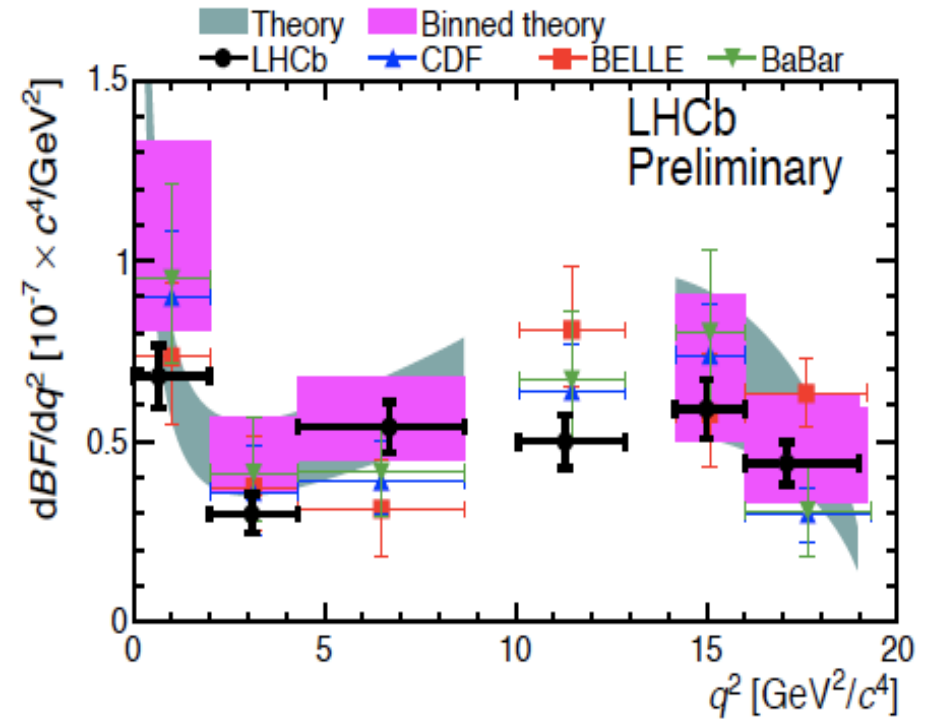
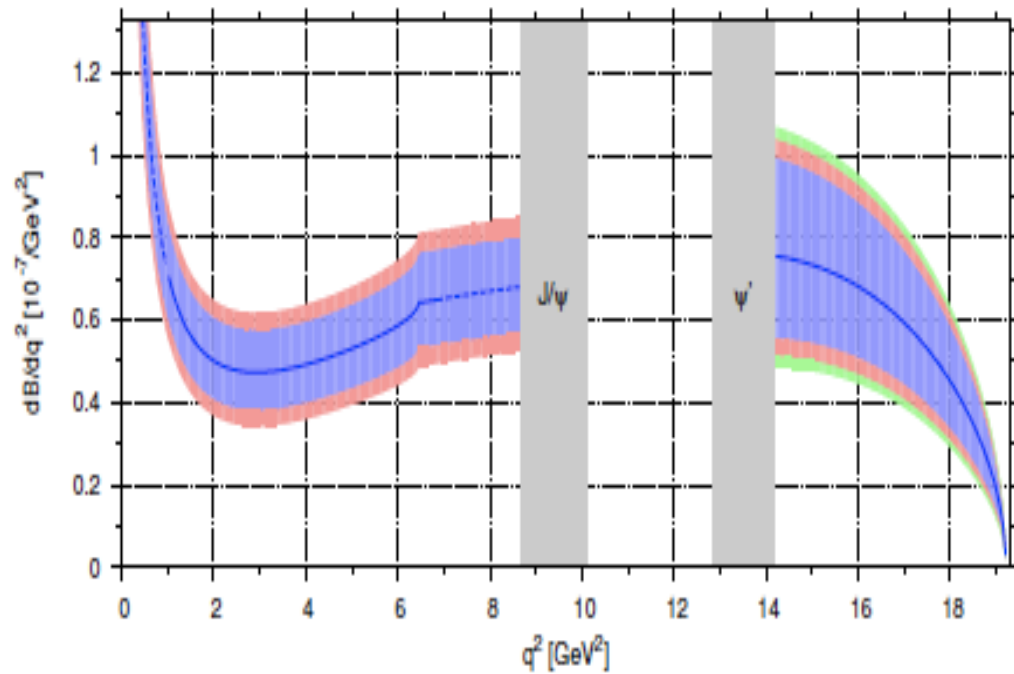


Different TH at **low q^2** QCDF; BBNS, Beneke, Feldmann, Seidel'01,04 and **high q^2 /low recoil**

OPE in $1/m_b$ Grinstein, Pirjol '04, Beylich, Buchalla, Feldmann '11; Low recoil $B \rightarrow K^{(*)} \mu^+ \mu^-$ predictions/pheno

Bobeth, GH, vanDyk, Wacker '10,11 Binned data needed. New developments at low recoil in theory pheno+lattice greatly support exploitation of today's and tomorrow's data. E.g., Preliminary unquenched lattice $B \rightarrow K^{(*)}$ form factors by Liu et al [1101.2726 \[hep-ph\]](#).

Dilepton Mass Spectra in $B \rightarrow K^* \mu^+ \mu^-$



left-hand Fig. from 1006.5013 [hep-ph] Blue band: form factor uncertainties, red: $1/m_b$ right-hand Fig. from LHCb-CONF-2012-008

Biggest source of TH uncertainty: the $B \rightarrow K^*$ form factors.

Opportunity: Angular Analysis $B \rightarrow V(\rightarrow PP)\mu^+\mu^-$

$$d\Gamma^4 \sim J dq^2 d \cos \Theta_l d \cos \Theta_{K^*} d\Phi \text{ hep-ph/9907386}$$

$$\begin{aligned} J(q^2, \theta_l, \theta_{K^*}, \phi) = & J_1^s \sin^2 \theta_{K^*} + J_1^c \cos^2 \theta_{K^*} + (J_2^s \sin^2 \theta_{K^*} + J_2^c \cos^2 \theta_{K^*}) \cos 2\theta_l \\ & + J_3 \sin^2 \theta_{K^*} \sin^2 \theta_l \cos 2\phi + J_4 \sin 2\theta_{K^*} \sin 2\theta_l \cos \phi + J_5 \sin 2\theta_{K^*} \sin \theta_l \cos \phi \\ & + J_6 \sin^2 \theta_{K^*} \cos \theta_l + J_7 \sin 2\theta_{K^*} \sin \theta_l \sin \phi \\ & + J_8 \sin 2\theta_{K^*} \sin 2\theta_l \sin \phi + J_9 \sin^2 \theta_{K^*} \sin^2 \theta_l \sin 2\phi, \end{aligned} \quad (2.3)$$

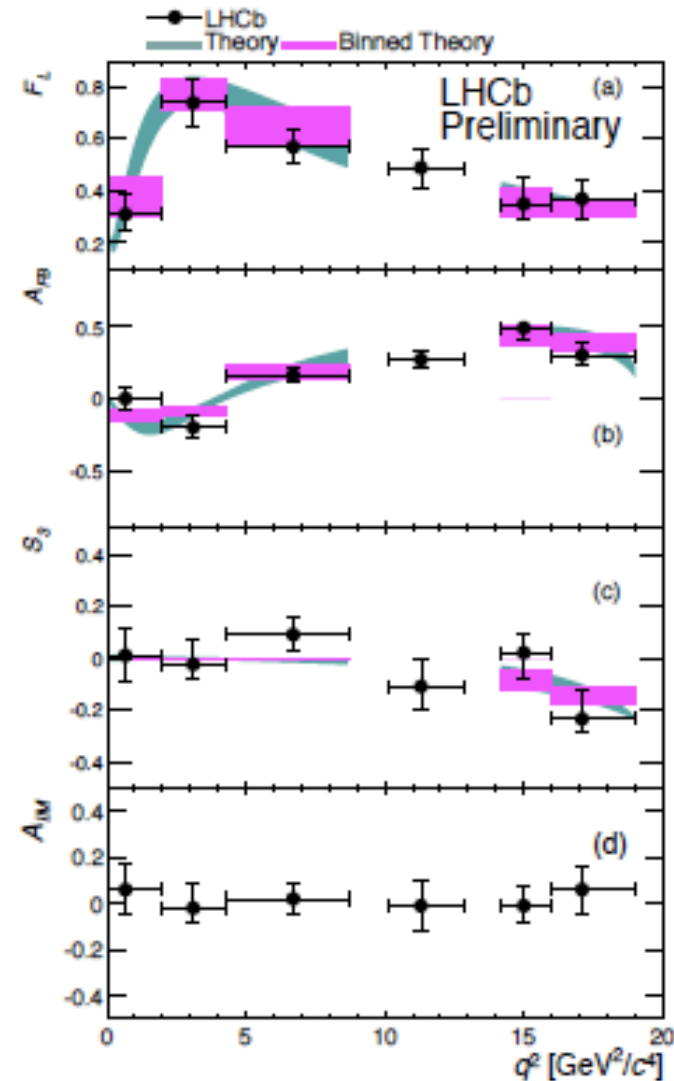
Θ_l : angle between l^- and \bar{B} in dilepton CMS (warning: different conventions in literature)

Θ_{K^*} : angle between K and \bar{B} in K^* -CMS

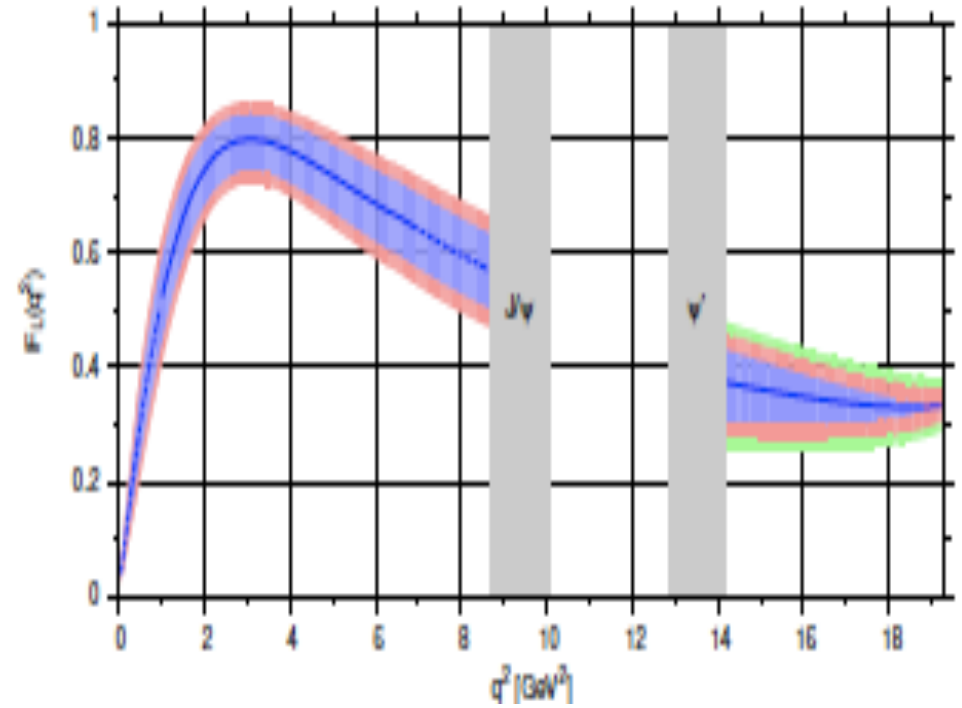
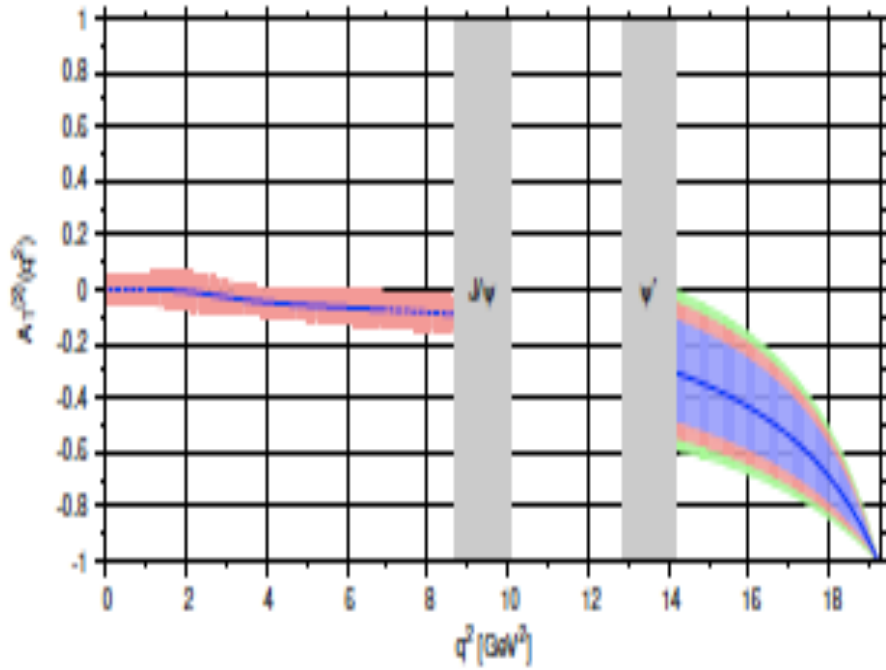
Φ : angle between normals of the $K\pi$ and l^+l^- plane

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$ Angular Analysis Results

- 4D fit to 3 angles and mass
- Larger data sample enables measurements of S_3 and A_{IM}
- Error bars include systematic uncertainties
- Data points at average q^2 of signal candidates in data
- These are the **most precise measurements** to-date [preliminary]
- The results are consistent with the SM prediction [1]



A_T^2 and F_L in $B \rightarrow K^* \mu^+ \mu^-$: SM



Figs. from 1006.5013 [hep-ph] Blue band: form factor uncertainties, red: $1/m_b$

Benefits of $B \rightarrow K^*$ at low recoil

At low hadronic recoil: OPE in $1/m_b$ Grinstein, Pirjol '04, Beylich et al '10

$$A_i^{L,R} \propto C^{L,R} f_i + \mathcal{O}(\alpha_s \Lambda/m_b, C_7/C_9 1/m_b), \quad i = \perp, \parallel, 0 \quad \text{Bobeth et al '10}$$

$C^{L,R}$: universal short-dist.-physics; $C^{L,R} = (C_9^{\text{eff}} \mp C_{10}) + \kappa \frac{2\hat{m}_b}{\hat{s}} C_7^{\text{eff}}$

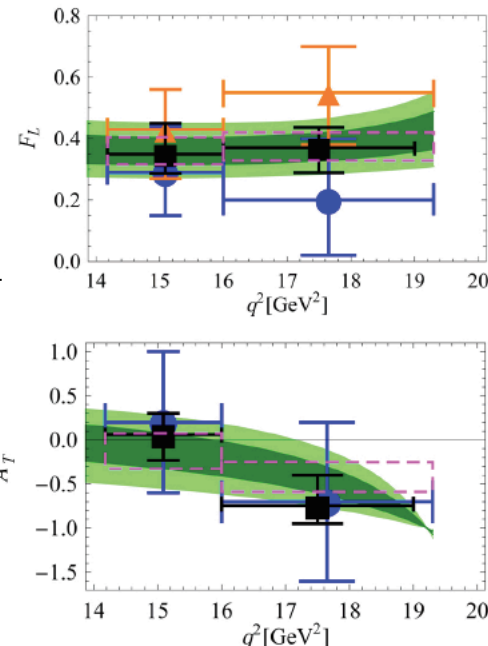
f_i : generalized form factors: $f_\perp \propto V$, $f_\parallel \propto A_1$, $f_0 \sim A_1, \lambda_{kin} A_2$

$C^{L,R}$ drops

out in ratios:

$$F_L = \frac{|A_0^L|^2 + |A_0^R|^2}{\sum_{X=L,R} (|A_\perp^X|^2 + |A_\parallel^X|^2)} = \frac{f_0^2}{f_\perp^2 + f_\parallel^2}$$

$$A_T^{(2)} = \frac{|A_\perp^L|^2 + |A_\perp^R|^2 - |A_\parallel^L|^2 - |A_\parallel^R|^2}{|A_\perp^L|^2 + |A_\perp^R|^2 + |A_\parallel^L|^2 + |A_\parallel^R|^2} = \frac{f_\perp^2 - f_\parallel^2}{f_\perp^2 + f_\parallel^2}$$

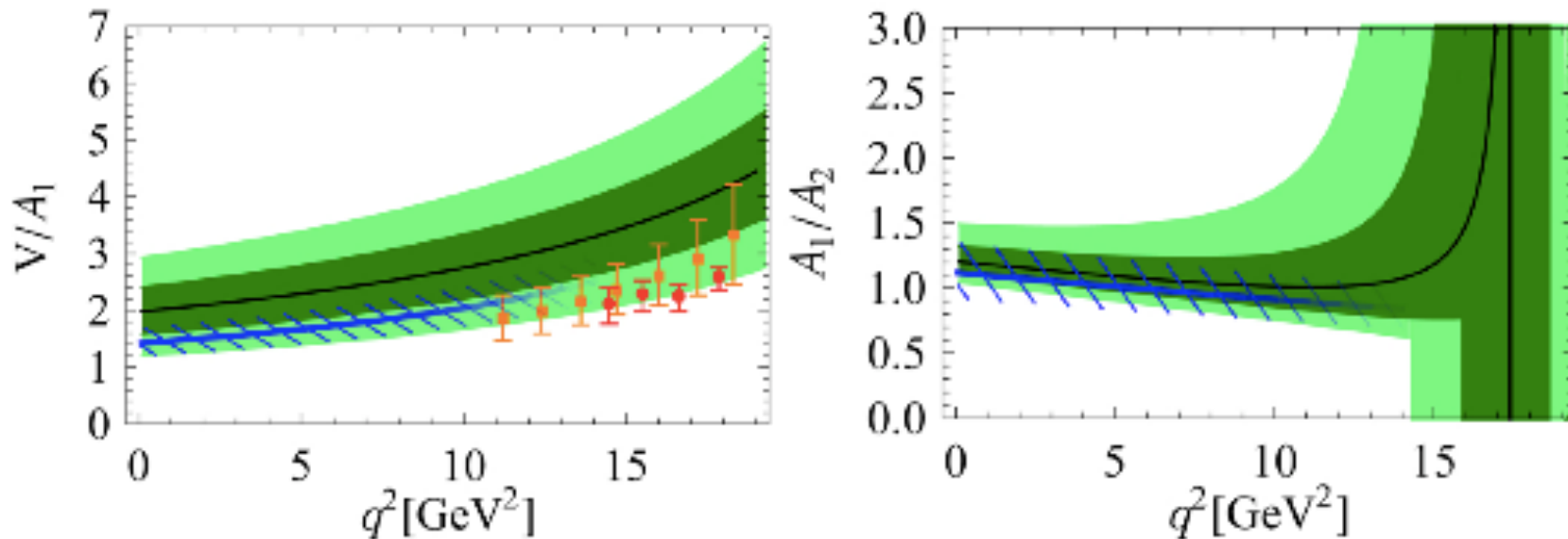


Extracting $B \rightarrow K^*$ form factors from data

Series expansion $z(t) \equiv z(t, t_0) = \frac{\sqrt{t_+ - t} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - t} + \sqrt{t_+ - t_0}}$,

$$\hat{f}_i(t) = \frac{(\sqrt{-z(t, 0)})^m (\sqrt{z(t, t_-)})^l}{B(t) \varphi_f(t)} \sum_k \alpha_{i,k} z^k(t),$$

The best-fit results: $\alpha_{\parallel}/\alpha_{\perp} = 0.43_{-0.08}^{+0.11}$, $\alpha_0/\alpha_{\perp} = 0.15_{-0.02}^{+0.03}$



Yellow, red points; lattice QCD; blue bands: QCD sum rules Ball, Zwicky '05; green bands: 1, 2 σ fit [1204.4444 \[hep-ph\]](#), PRL'12

-
1. Its great to have (even more) data.
 2. With (even one) more bins the sensitivity in the fits to the q^2 -shape increases.
 3. If you(lattice, sum rules,..) calculate form factors, please provide also ratios (with uncertainties)

Interplay

Precision tests from global fits C_7, C_9, C_{10}

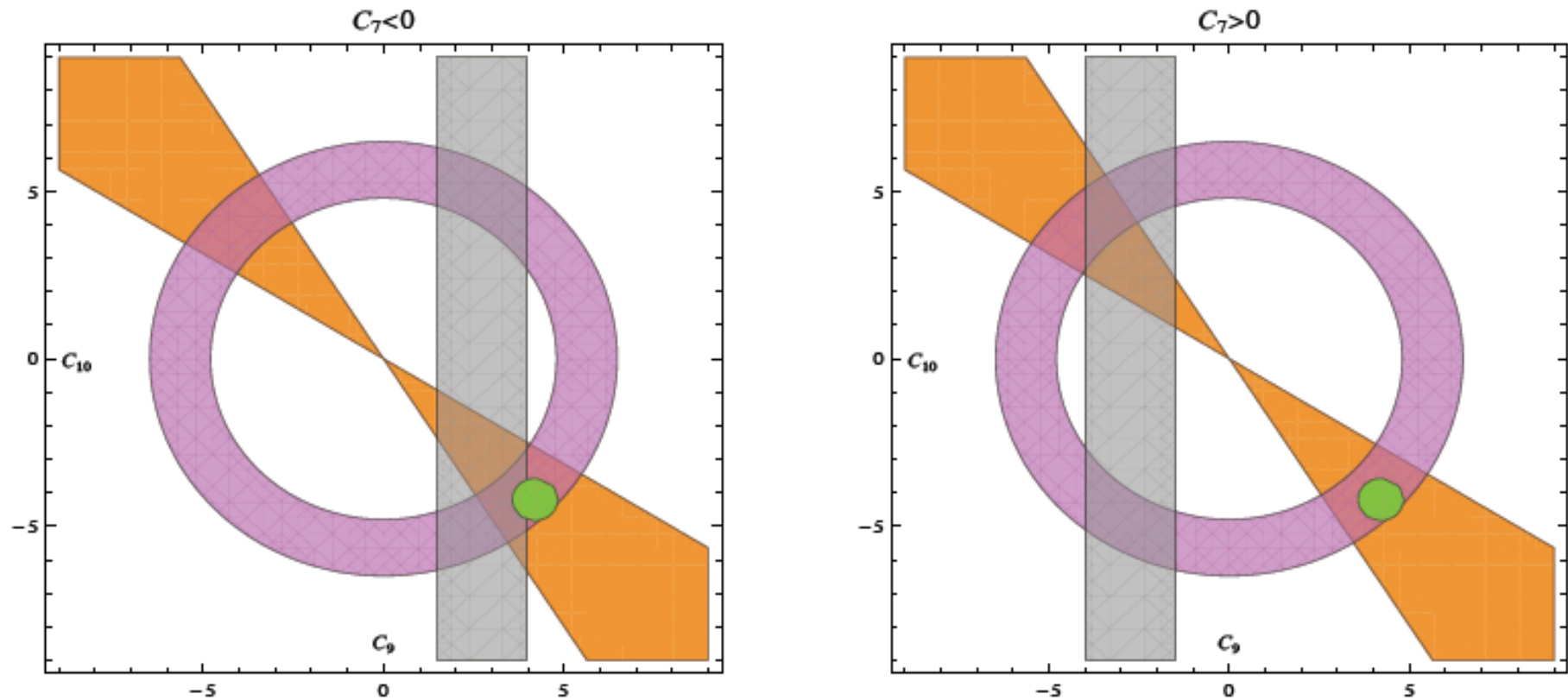
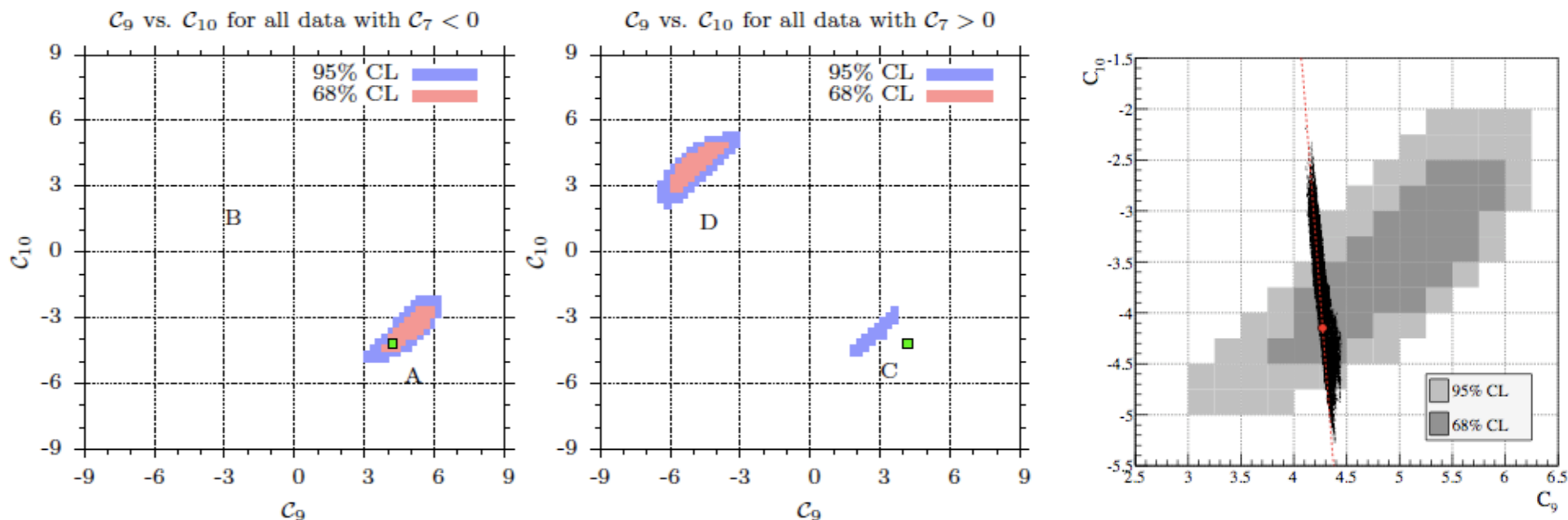


Figure 3: Future scenario of the model-independent bounds on real-valued C_9, C_{10} from $b \rightarrow s\mu^+\mu^-$ decays for $C_7 \simeq C_7^{\text{SM}} < 0$ (left-handed plot) and $C_7 \simeq -C_7^{\text{SM}} > 0$ (right-handed plot). The grey vertical bands denote the constraints arising if an A_{FB} zero at low q^2 could be established. There remain two allowed disconnected regions. schematic from 1106.1547 [hep-ph]. **Orange:** AFB at lo reco $A_{\text{FB}} \propto C_9 C_{10}$, **pink:** $Br \propto C_9^2 + C_{10}^2$, **grey** $q_0^2 \propto -C_7/C_9$. **Green point:** SM

Precision tests from global fits C_7, C_9, C_{10}

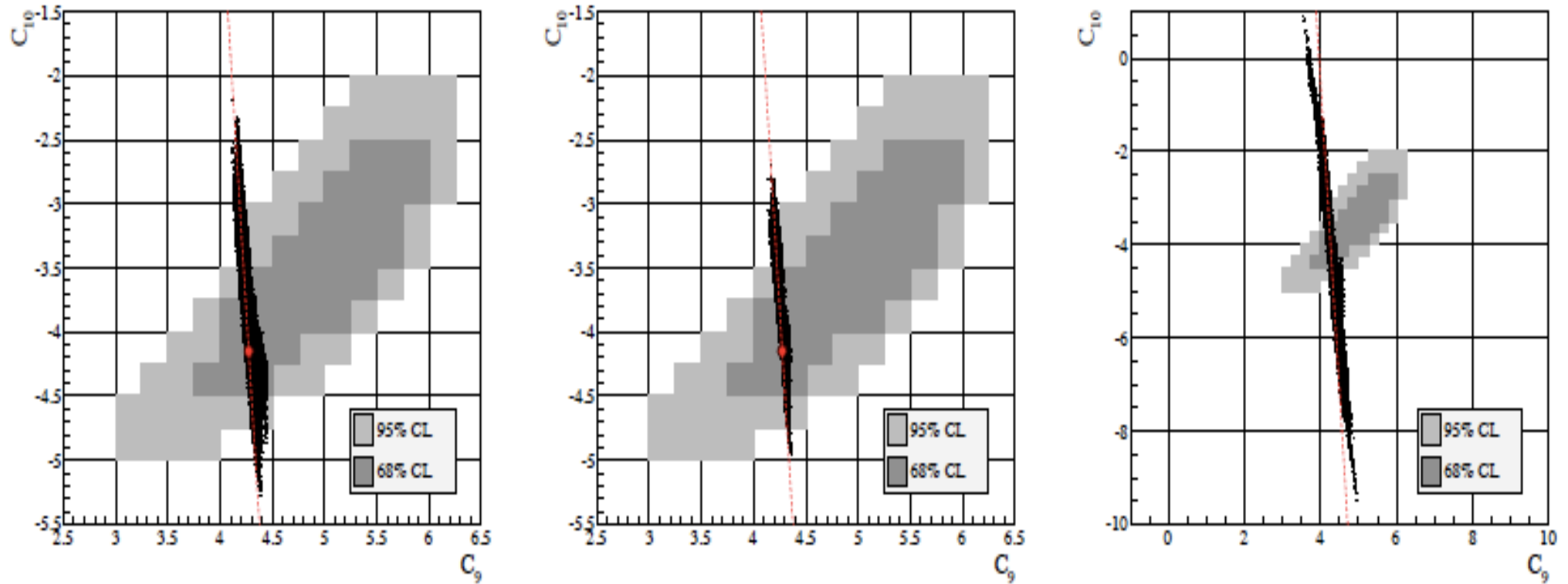


left: global analysis [1111.2558 \[hep-ph\]](#), also [Altmannshofer et al '11](#) ; solution C is ruled out by AFB-zero. right: red dot: SM; grey areas: allowed by $b \rightarrow s$ data; black points: SUSY model with squark flavor mixing [1205.1500 \[hep-ph\]](#)

flavor suppression with NP at $\Lambda_{NP} = 1 \text{ TeV}$: $|\tilde{c}_{10}| < 5 \cdot 10^{-4} (4 \cdot 10^{-3})$

limit on scale iff no suppression $\tilde{c}_{10} = 1$: $\Lambda_{NP} > 44 \text{ TeV}$ (16 TeV)

Flavor Interplay with Higgs Physics (in SUSY)



left: from 1205.1500 [hep-ph] with $m_{h^0} > 114.4$ GeV, right and mid: courtesy of Stefan Schacht, Talk given at FLASY'12, Dortmund; right: no Higgs constraint, mid: $120 < m_{h^0} < 130$ GeV; Higgs mass calculated with FeynHiggs 2.9.0-beta