



LUNDS
UNIVERSITET

Monte-Carlos of jet quenching: an overview (II)

Konrad Tywoniuk

Jet Modification in the RHIC and LHC era
(QM12 Satellite Workshop)

20-24 August 2012, Wayne State University

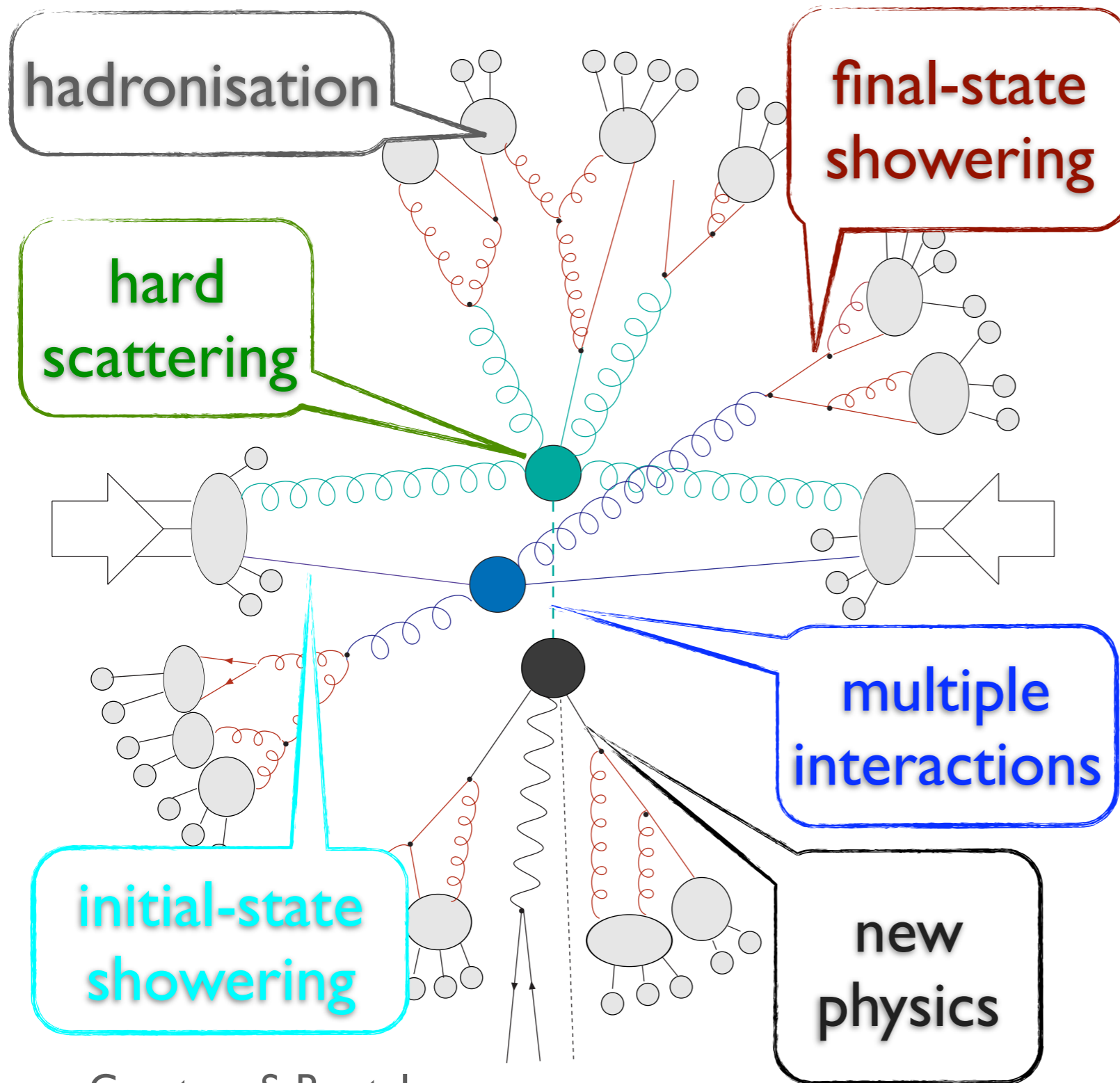
Outline

- classical vs. quantum
 - when is a probabilistic interpretation viable?
- key effects & concepts
- perspectives & challenges

[disclaimer: devil is often in the details...]



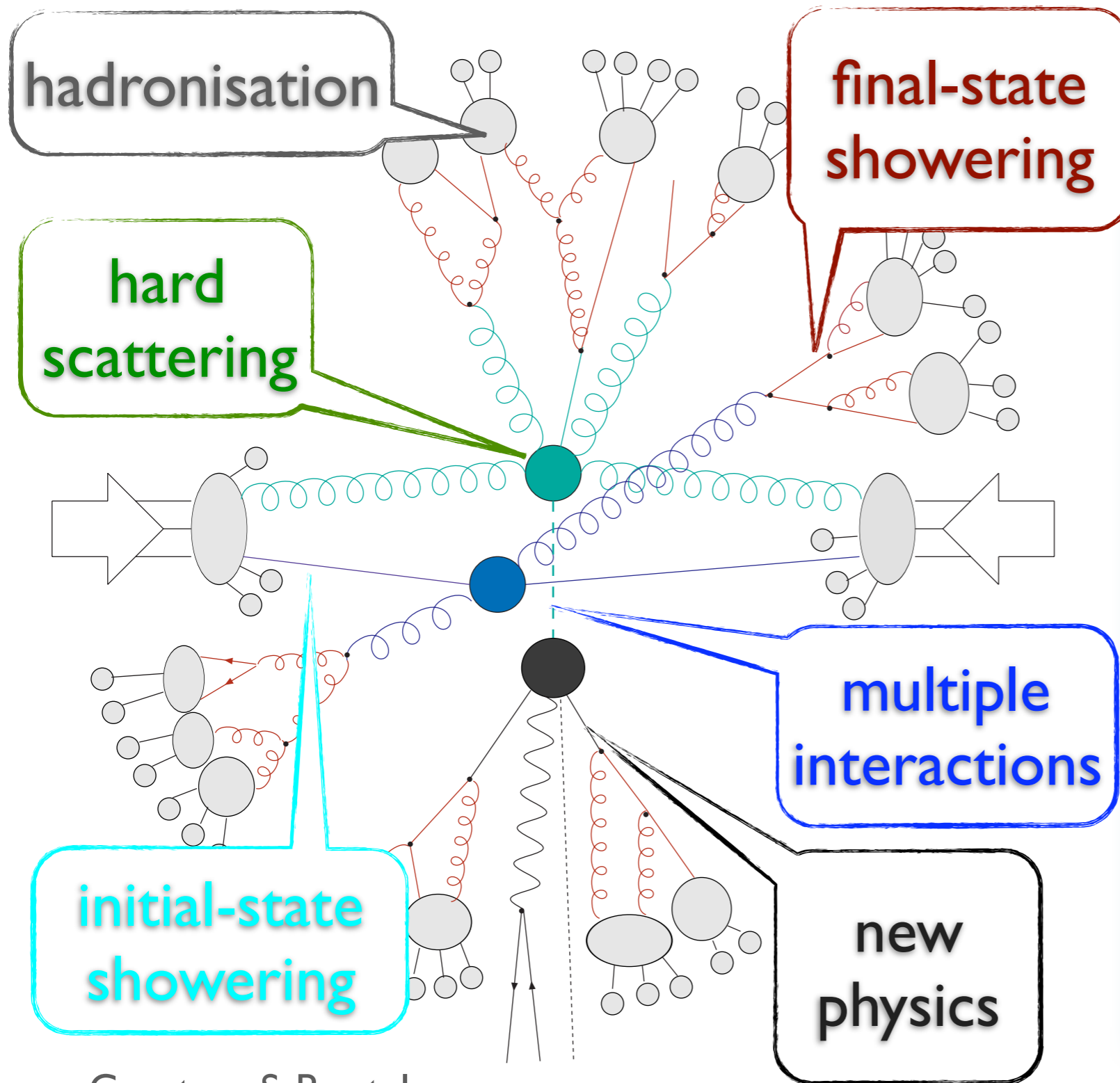
typical pp event [à la PYTHIA, HERWIG, SHERPA]



Courtesy: S. Prestel



typical pp event [à la PYTHIA, HERWIG, SHERPA]



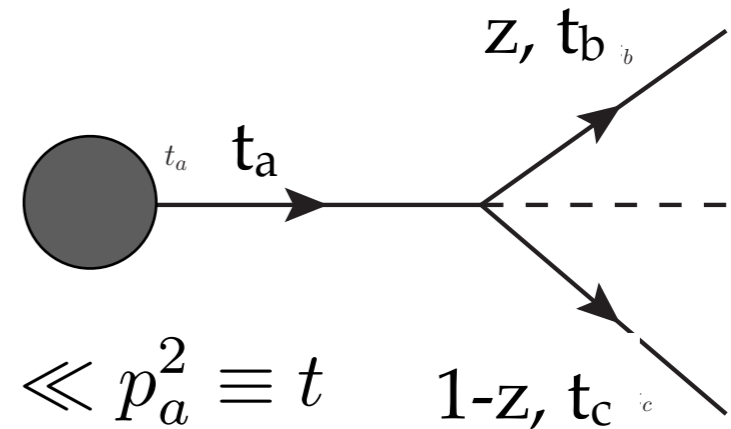
Issues in HIC:

- factorization?
- soft physics
- back-reaction
- showering
- final-state
- initial-state
- hadronisation
- other...

Courtesy: S. Prestel



MC shower: primer



- factorization of phase-space & splitting

$$p_b^2, p_c^2 \ll p_a^2 \equiv t \quad 1-z, t_c$$

$$d\sigma_{n+1} = d\sigma_n \frac{dt}{t} \frac{d\phi}{2\pi} dz \frac{\alpha_s}{2\pi} P(z)$$

$$t = 2E_b E_c (1 - \cos \theta) = z(1-z)E_a^2 \theta^2$$

- collinear & soft singularities
- coherence effects (soft)

not discernible at this level yet...



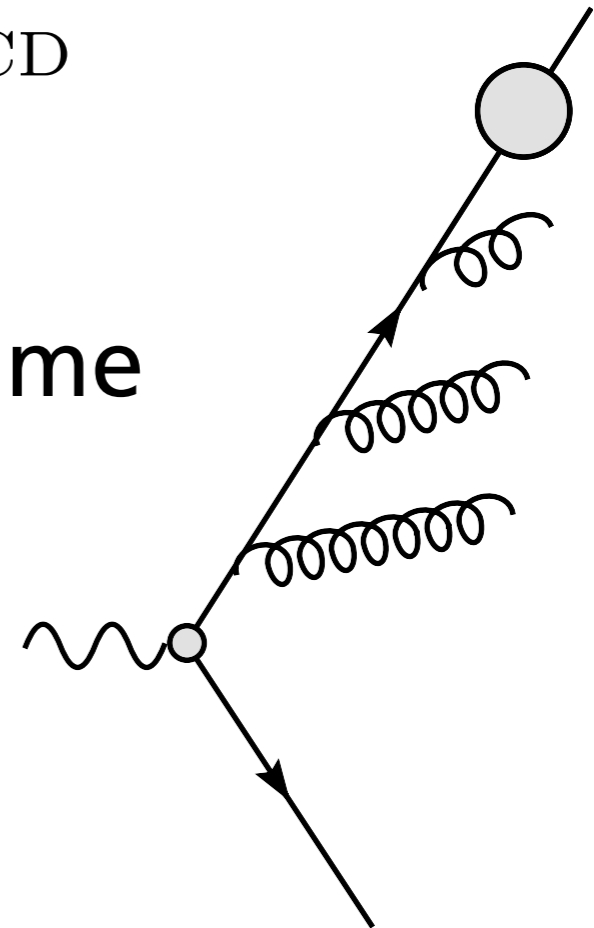
DGLAP evolution

$$t_0 \gg t_1 \gg t_2 \gg \dots \gg \Lambda_{\text{QCD}}^2$$

- strong ordering in virtuality
- strong ordering in formation time

$$t_f = \frac{2\omega}{k_{\perp}^2}$$

- allows for a probabilistic interpretation!



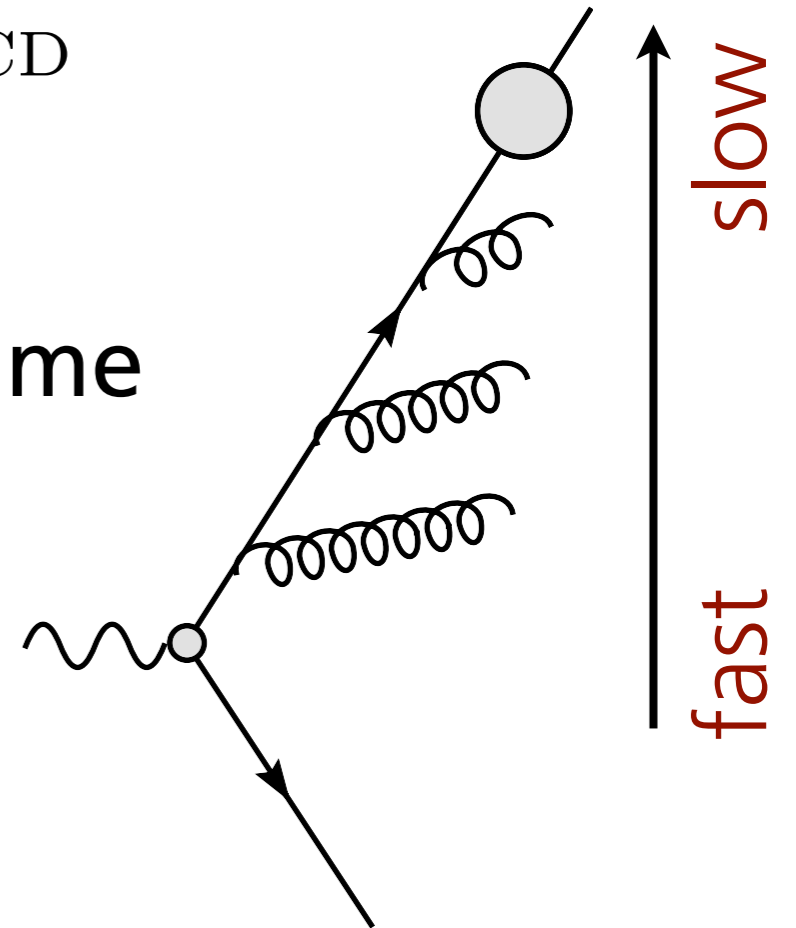
DGLAP evolution

$$t_0 \gg t_1 \gg t_2 \gg \dots \gg \Lambda_{\text{QCD}}^2$$

- strong ordering in virtuality
- strong ordering in formation time

$$t_f = \frac{2\omega}{k_{\perp}^2}$$

- allows for a probabilistic interpretation!



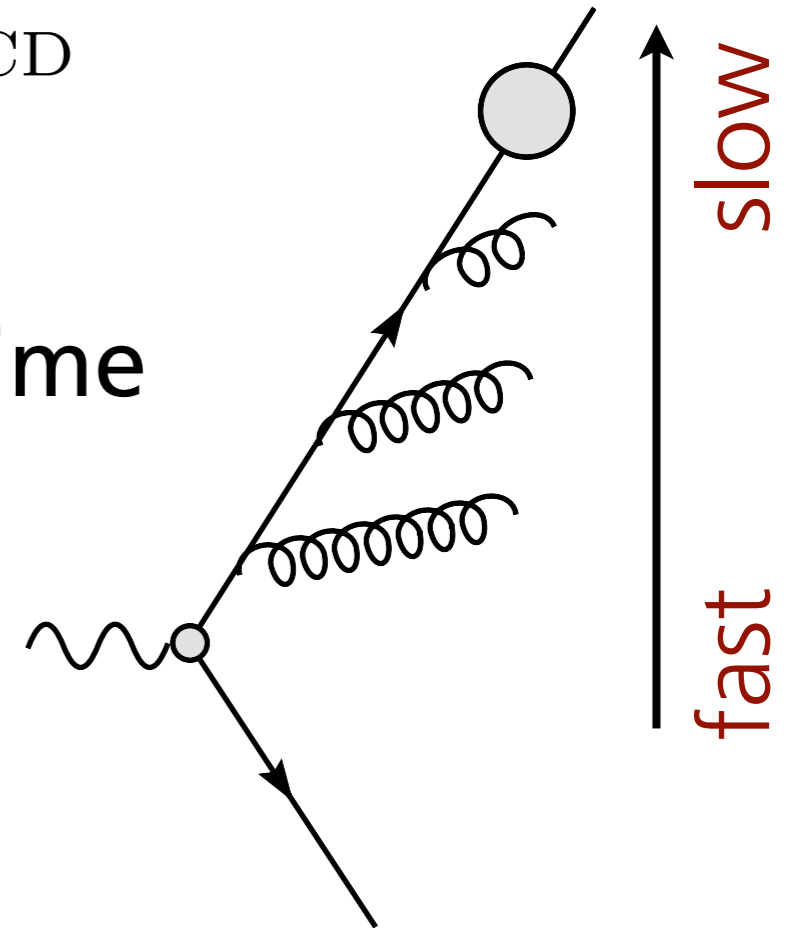
DGLAP evolution

$$t_0 \gg t_1 \gg t_2 \gg \dots \gg \Lambda_{\text{QCD}}^2$$

- strong ordering in virtuality
- strong ordering in formation time

$$t_f = \frac{2\omega}{k_{\perp}^2}$$

- allows for a probabilistic interpretation!



Distribution of gluons with mom fraction x and virtuality Q^2

$$t \frac{\partial f_i(x, t)}{\partial t} = \sum_j \int_x^1 \frac{dz}{z} \frac{\alpha_s}{2\pi} P_{ij}(z) f_j \left(\frac{x}{z}, t \right)$$

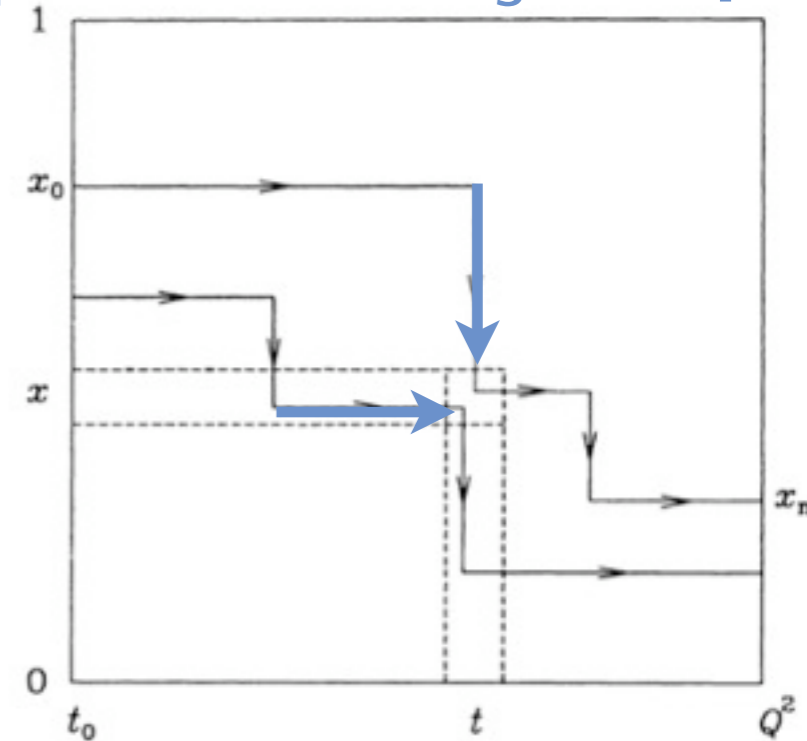


Stochastic process

[space-like branching here...]

$$\mathcal{P}(a \rightarrow b(z) + c(1-z) \text{ when } t \rightarrow t + \delta t) = \frac{\delta t}{t} \frac{\alpha_s}{2\pi} P_{ba}(z)$$

tap:
$$\delta f_{in}(x, t) = \frac{\delta t}{t} \int_x^1 dx' dz \frac{\alpha_s}{2\pi} P(z) f(x', t) \delta(x - zx')$$



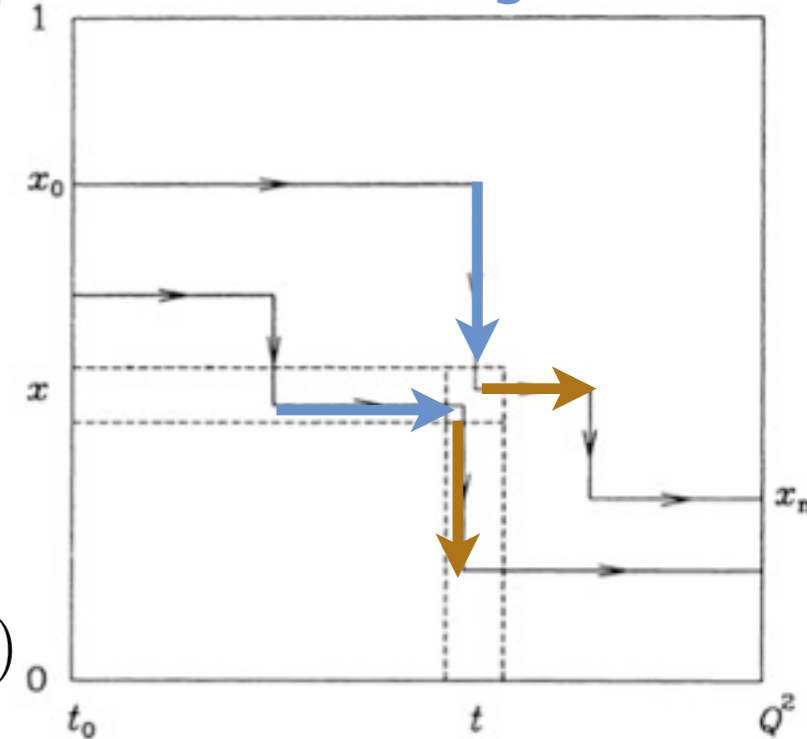
Stochastic process

[space-like branching here...]

$$\mathcal{P}(a \rightarrow b(z) + c(1-z) \text{ when } t \rightarrow t + \delta t) = \frac{\delta t}{t} \frac{\alpha_s}{2\pi} P_{ba}(z)$$

tap:
$$\delta f_{in}(x, t) = \frac{\delta t}{t} \int_x^1 dx' dz \frac{\alpha_s}{2\pi} P(z) f(x', t) \delta(x - zx')$$

sink:
$$\delta f_{out}(x, t) = \frac{\delta t}{t} f(x, t) \int_0^x dx' dz \frac{\alpha_s}{2\pi} P(z) \delta(x' - zx)$$



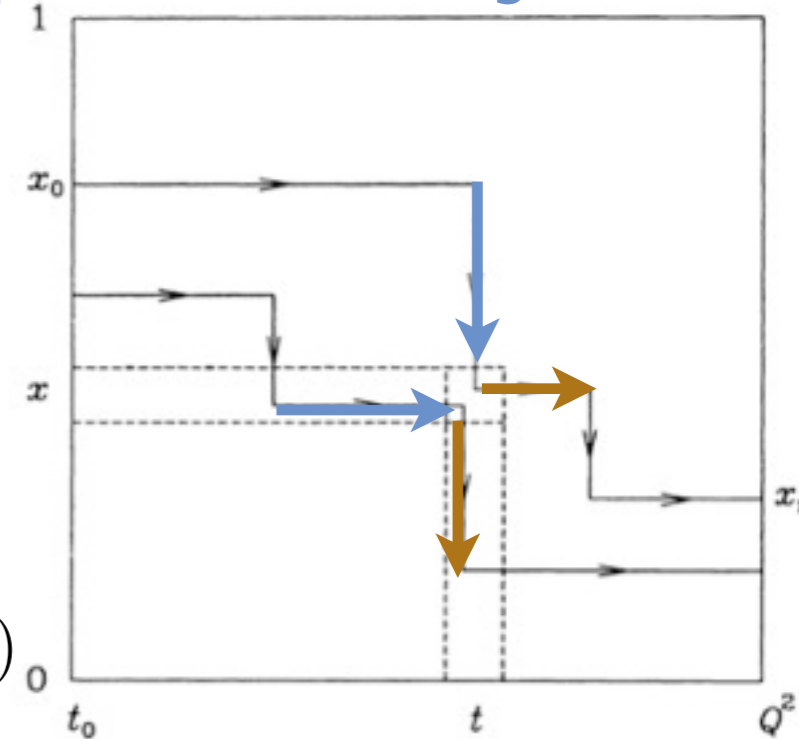
Stochastic process

[space-like branching here...]

$$\mathcal{P}(a \rightarrow b(z) + c(1-z) \text{ when } t \rightarrow t + \delta t) = \frac{\delta t}{t} \frac{\alpha_s}{2\pi} P_{ba}(z)$$

tap:
$$\delta f_{in}(x, t) = \frac{\delta t}{t} \int_x^1 dx' dz \frac{\alpha_s}{2\pi} P(z) f(x', t) \delta(x - zx')$$

sink:
$$\delta f_{out}(x, t) = \frac{\delta t}{t} f(x, t) \int_0^x dx' dz \frac{\alpha_s}{2\pi} P(z) \delta(x' - zx)$$



$$\delta f(x, t) = \delta f_{in}(x, t) - \delta f_{out}(x, t) = \frac{\delta t}{t} \int_0^1 dz \frac{\alpha_s}{2\pi} P(z) \left(\frac{1}{z} f\left(\frac{x}{z}, t\right) - f(x, t) \right)$$

virtual terms

- evolution “time” = virtuality



Sudakov form factor

$$\Delta(t_1, t_2) = \exp \left[- \int_{t_1}^{t_2} \frac{dt}{t} \int_{z_{\min}(t)}^{1-z_{\min}(t)} dz \frac{\alpha_s}{2\pi} P(z) \right] \quad \text{probability of no emission in } [t_1, t_2] \text{ interval}$$



Sudakov form factor

$$\Delta(t_1, t_2) = \exp \left[- \int_{t_1}^{t_2} \frac{dt}{t} \int_{z_{\min}(t)}^{1-z_{\min}(t)} dz \frac{\alpha_s}{2\pi} P(z) \right] \quad \text{probability of no emission in } [t_1, t_2] \text{ interval}$$

Integral equation:

$$f(x, t) = \Delta(t_0, t) f(x, t_0) + \int_{t_0}^t \frac{dt'}{t'} \Delta(t', t) \int \frac{dz}{z} \frac{\alpha_s}{2\pi} P(z) f\left(\frac{x}{z}, t'\right)$$



Sudakov form factor

$$\Delta(t_1, t_2) = \exp \left[- \int_{t_1}^{t_2} \frac{dt}{t} \int_{z_{\min}(t)}^{1-z_{\min}(t)} dz \frac{\alpha_s}{2\pi} P(z) \right] \quad \text{probability of no emission in } [t_1, t_2] \text{ interval}$$

Integral equation:

$$f(x, t) = \Delta(t_0, t) f(x, t_0) + \int_{t_0}^t \frac{dt'}{t'} \Delta(t', t) \int \frac{dz}{z} \frac{\alpha_s}{2\pi} P(z) f\left(\frac{x}{z}, t'\right)$$

Veto algorithm:

dicing down a step in the ladder using two random numbers

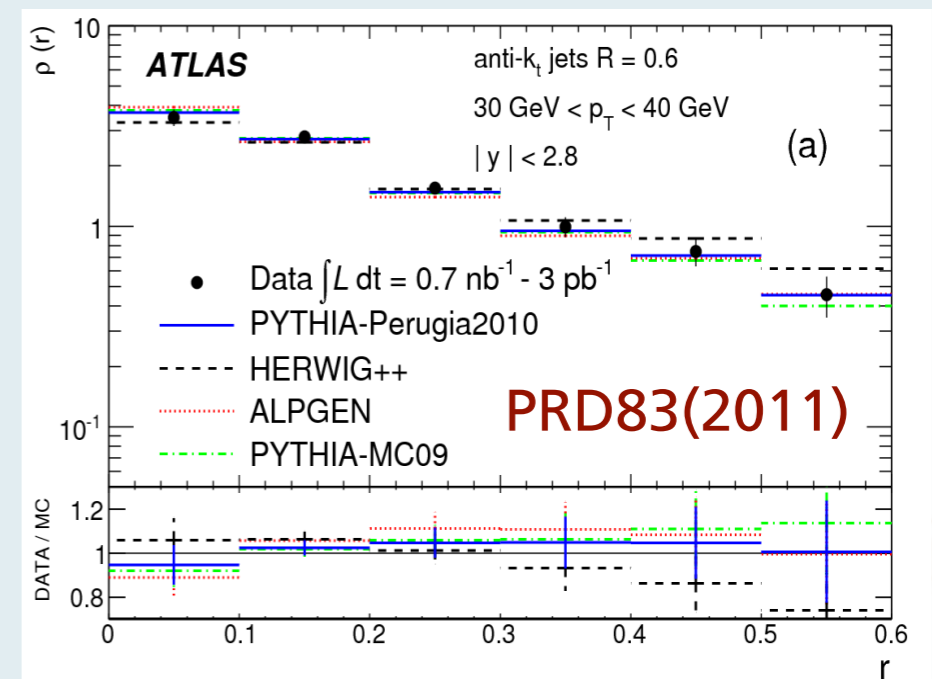
$$\Delta(t, t') = \mathcal{R}_1 \int_{z_{\min}(t')}^{z'} dz \frac{\alpha_s}{2\pi} P(z) = \mathcal{R}_2 \int_{z_{\min}(t')}^{1-z_{\min}(t')} dz \frac{\alpha_s}{2\pi} P(z)$$



Coherence effects

- soft gluon radiation implicate large angles
- another type of evolution equation!
- in MC: accounted for in an “average” sense
 - angular ordering
 - good enough for inclusive & collinear observables
 - inter-jet activity

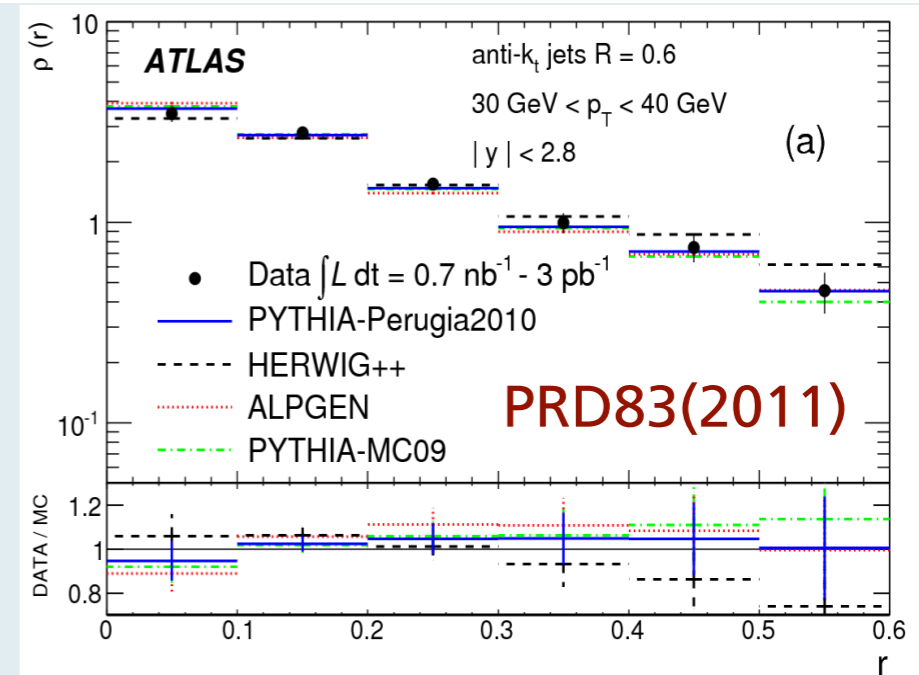
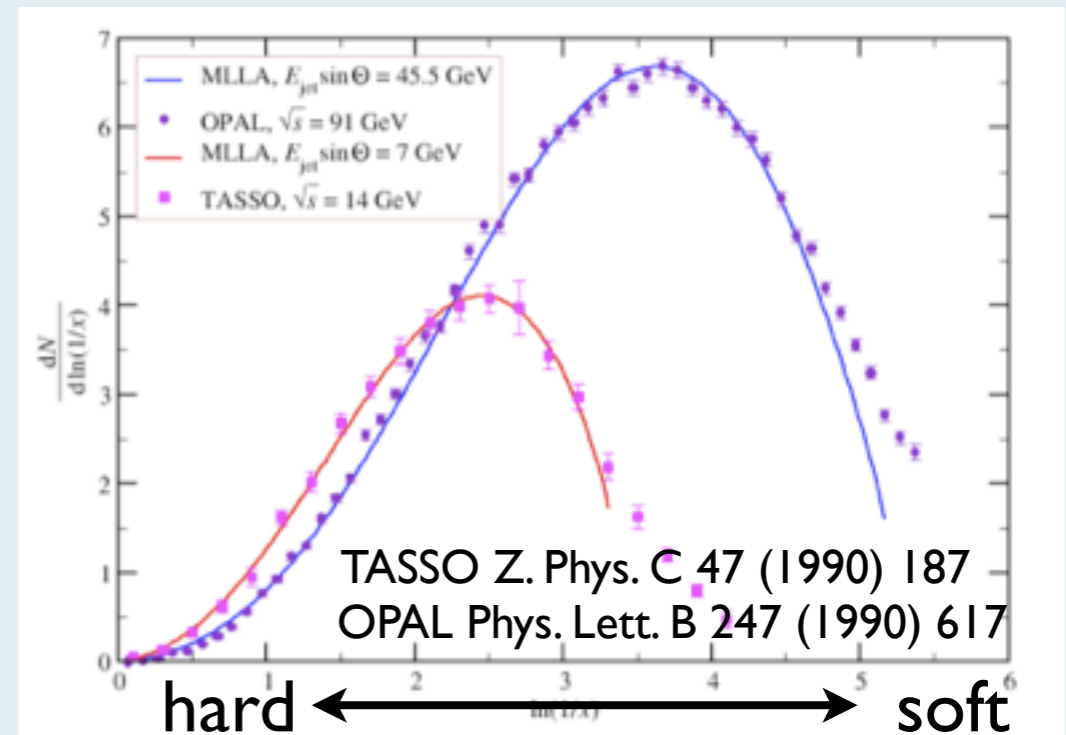
[see Yacine’s talk, Mon]



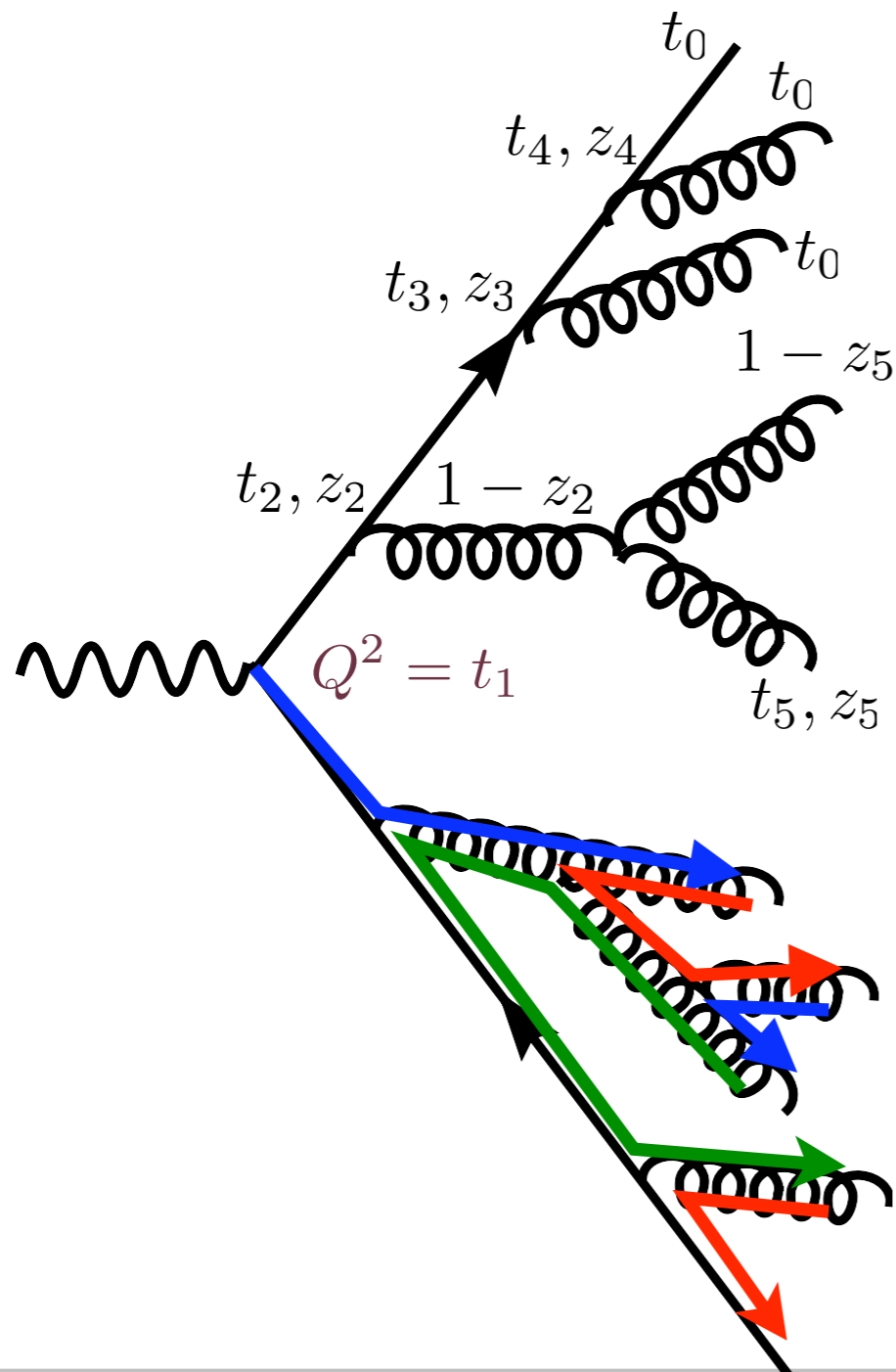
Coherence effects

- soft gluon radiation implicate large angles
- another type of evolution equation!
- in MC: accounted for in an “average” sense
 - angular ordering
 - good enough for inclusive & collinear observables
 - inter-jet activity

[see Yacine’s talk, Mon]



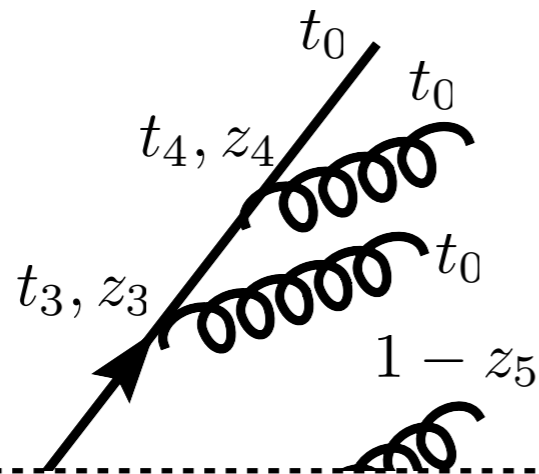
Parton-showers in vacuum



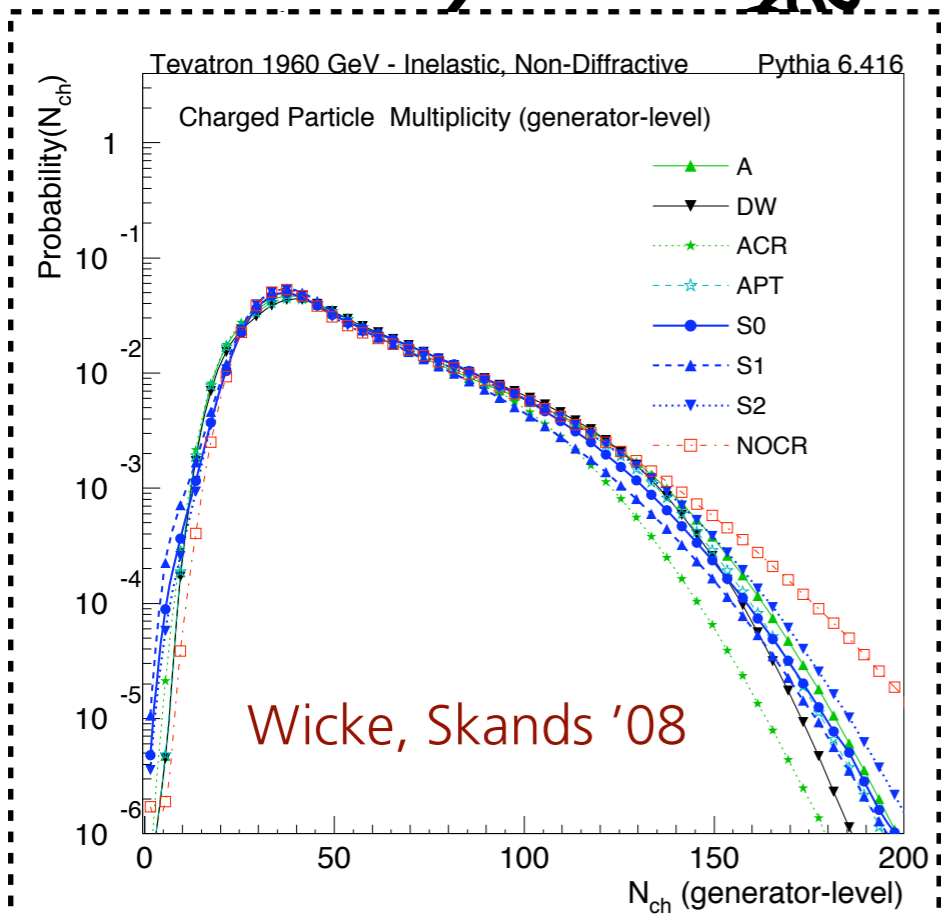
- no reference to space-time
- large N_c -limit (planar)
- only closest dipoles can radiate
- multi-jet and matching
- non-perturbative effects
 - color reconnections(!)
 - hadronization at Q_0
 - Lund string model
 - cluster hadronization



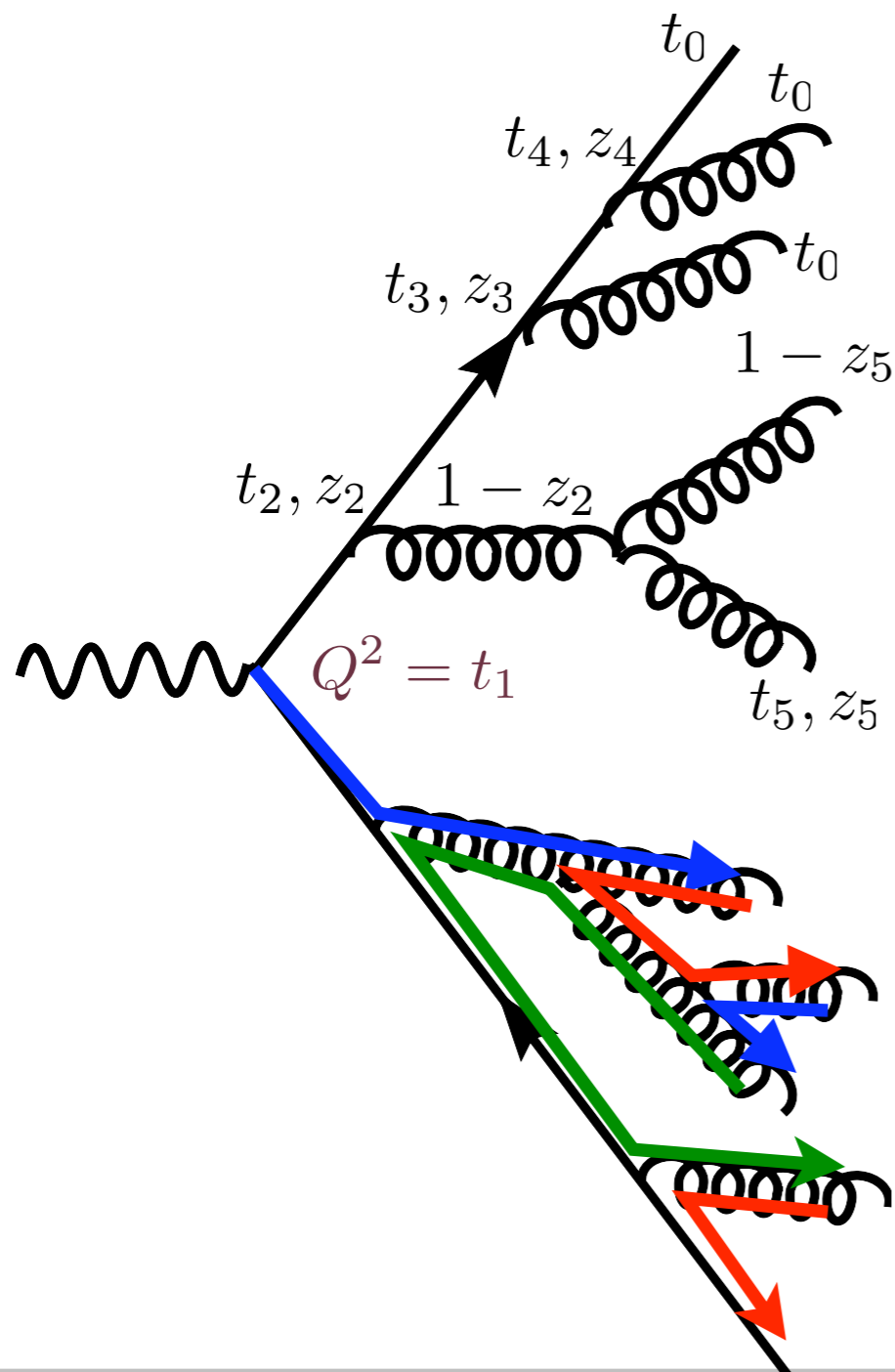
Parton-showers in vacuum



- no reference to space-time
- large N_c -limit (planar)
- only closest dipoles can radiate
- multi-jet and matching
- non-perturbative effects
- color reconnections(!)
- hadronization at Q_0
 - Lund string model
 - cluster hadronization



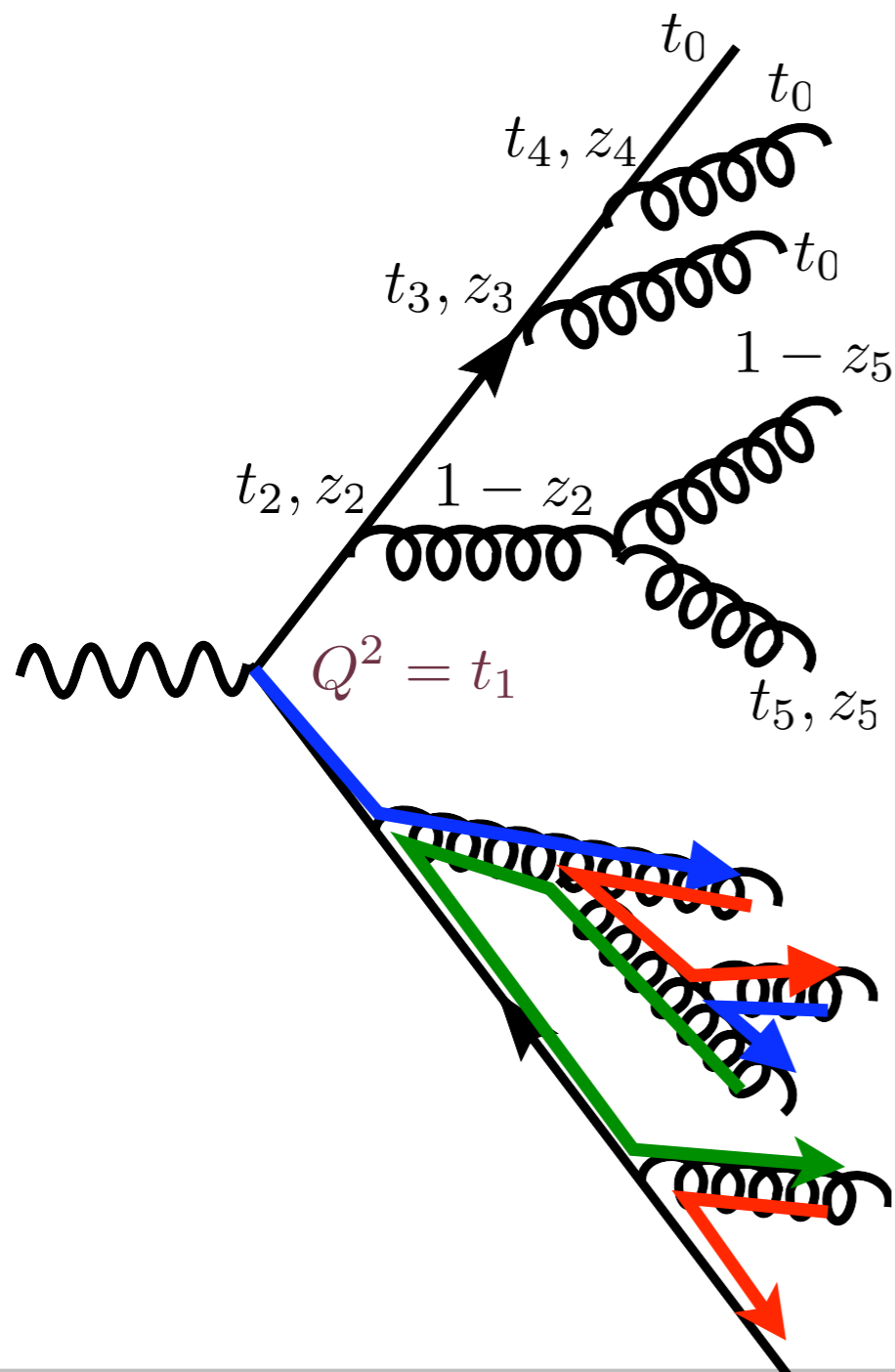
Parton-showers in vacuum



- probabilistic interpretation ensured!
- PYTHIA
 - virtuality (k_{\perp}) ordered (veto on angular ordering)
- HERWIG
 - angular ordered



Parton-showers in vacuum



- probabilistic interpretation ensured!
- PYTHIA
 - virtuality (k_{\perp}) ordered (veto on angular ordering)
- HERWIG
 - angular ordered

Issues in HIC:

- two types of radiation
- dispersion of momentum
- reference to space-time
- ...“(un)kown unkowns”



Models: an overview

[not comprehensive..! well of transport formulations!]

	vacuum rad.	med-induced rad.	elastic e.-loss	remarks
HIJING	✓	✓	(?)	full generator
PYQUEN	✓	✓		rough BDMPS
YaJEM	✓	✓	✓	mod kinematics
JEWEL1.0	✓	✓	✓	mod splitting functions + kinematics
JEWEL-LPM		✓		'exact' induced radiation
MARTINI	✓	✓	✓	rate equations
Q-PYTHIA	✓	✓		vacuum baseline
Q-HERWIG	✓	✓		vacuum baseline

Wang, Gyulassy PRD 44 (1991) 3501, Comput.Phys.Commun. 83 (1994) 307
Lokhtin, Snigirev EPJC 45 (2006) 211, Renk, PRC 78 (2008) 034908,
Ingelman, Rathsman, Stachel, Wiedemann, Zapp EPJC 60 (2009) 617,
Stachel, Wiedemann, Zapp JHEP 1107 (2011) 118,
Schenke, Gale, Jeon PRC 80 (2009) 054913
Armesto, Cunqueiro, Salgado EPJC 61 (2009) 775,
Armesto, Corcella, Cunqueiro, Salgado JHEP 0911 (2009) 122



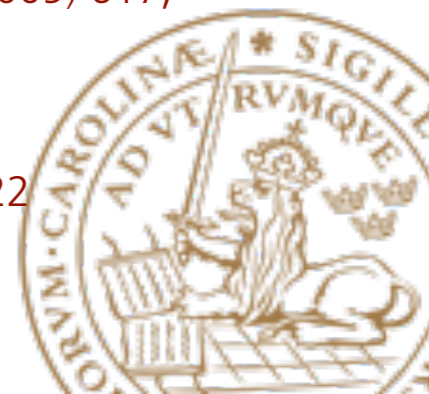
Models: an overview

[not comprehensive..! well of transport formulations!]

	vacuum rad.	med-induced rad.	elastic e.-loss	remarks
HIJING	✓	✓	(?)	full generator
PYQUEN	✓	✓		rough BDMPS
YaJEM	✓	✓	✓	mod kinematics
JEWEL1.0	✓	✓	✓	mod splitting functions + kinematics
JEWEL-LPM		✓		'exact' induced radiation
MARTINI	✓	✓	✓	rate equations
Q-PYTHIA	✓	✓		vacuum baseline
Q-HERWIG	✓	✓		vacuum baseline

only one Comput. Phys. Commun!

Wang, Gyulassy PRD 44 (1991) 3501, Comput.Phys.Commun. 83 (1994) 307
 Lokhtin, Snigirev EPJC 45 (2006) 211, Renk, PRC 78 (2008) 034908,
 Ingelman, Rathsman, Stachel, Wiedemann, Zapp EPJC 60 (2009) 617,
 Stachel, Wiedemann, Zapp JHEP 1107 (2011) 118,
 Schenke, Gale, Jeon PRC 80 (2009) 054913
 Armesto, Cunqueiro, Salgado EPJC 61 (2009) 775,
 Armesto, Corcella, Cunqueiro, Salgado JHEP 0911 (2009) 122



Standard features

Radiative processes

- $2 \rightarrow 3$ induced radiation (Gunion-Bertsch)
- medium-modified splitting functions
- absorptive reactions

Elastic processes

- transverse momentum broadening
- energy transfer, drag effects
- randomization of color
- back-reaction



Standard features

Radiative processes

- $2 \rightarrow 3$ induced radiation (Gunion-Bertsch)
- medium-modified splitting functions
- absorptive reactions

Elastic processes

- transverse momentum broadening
- energy transfer, drag effects
- randomization of color
- back-reaction

What are the typical timescales?

quantum \Leftrightarrow classical
[pQCD] [Boltzman eq., ...]



Standard features

Radiative processes

- $2 \rightarrow 3$ induced radiation (Gunion-Bertsch)
- medium-modified splitting functions
- absorptive reactions



Elastic processes

- transverse momentum broadening
- energy transfer, drag effects
- randomization of color
- back-reaction

What are the typical timescales?

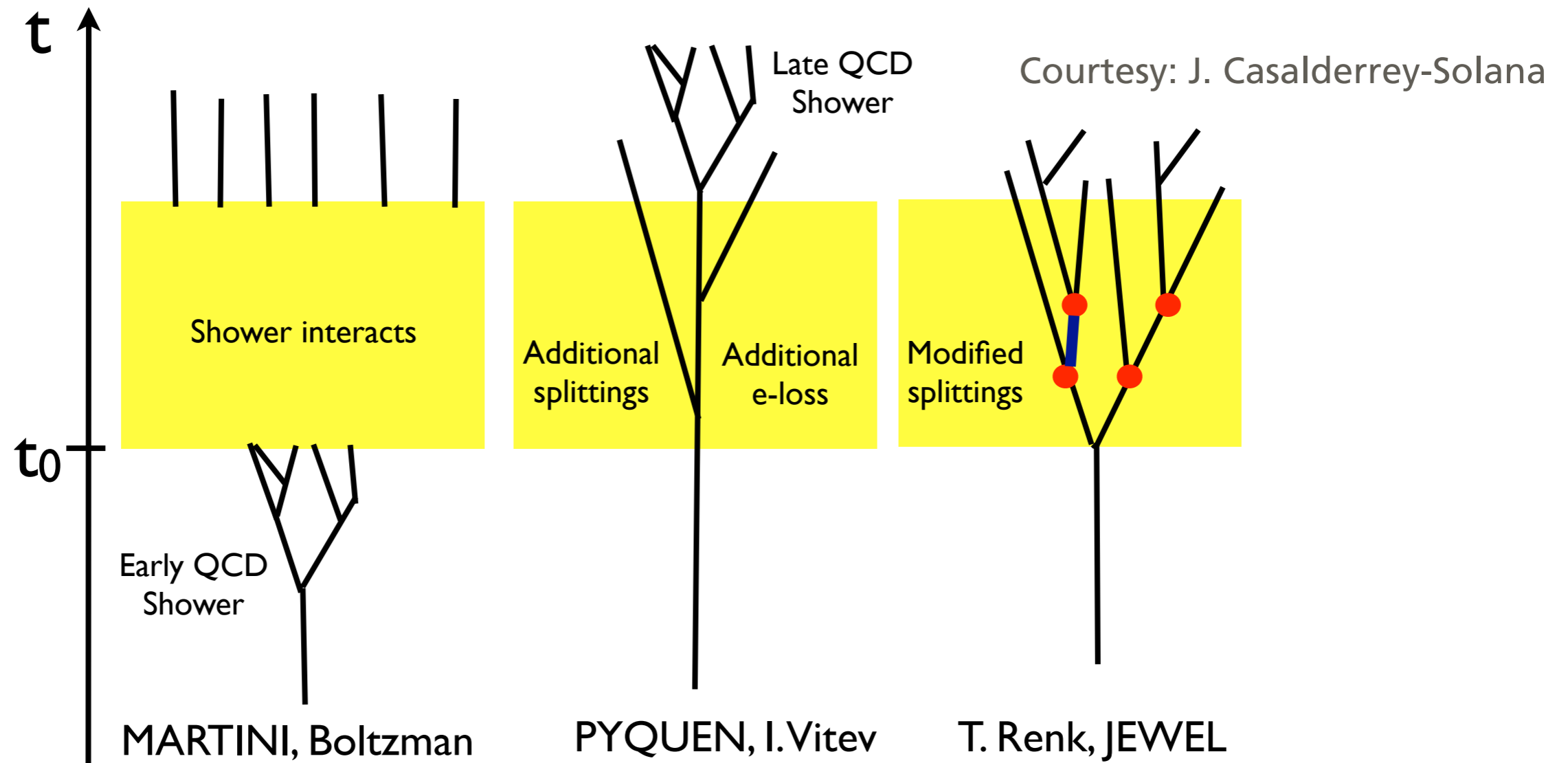
quantum
[pQCD]



classical
[Boltzman eq., ...]



Differences in evolution

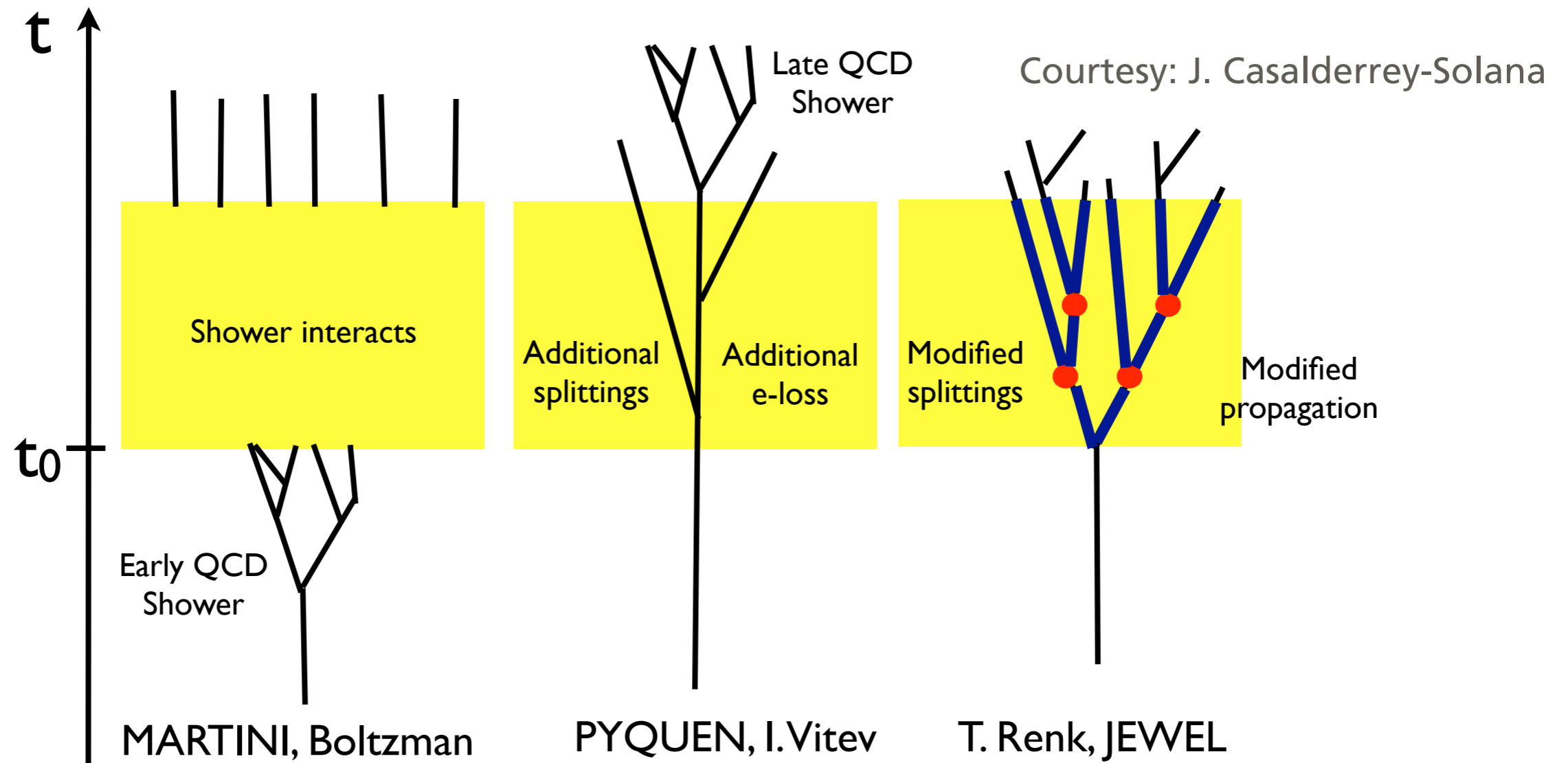


Is it reasonable to assume a separation of these processes?

We need guidance from theory!



Differences in evolution



Is it reasonable to assume a separation of these processes?

We need guidance from theory!



YaJEM

[lifetime of parton $i+1$]

$$\langle \tau_{i+1} \rangle = \frac{E_{i+1}}{Q_{i+1}^2} - \frac{E_{i+1}}{Q_i^2}$$

Renk, PRC 78 (2008) 034908

[emergence of parton $i+1$]

$$\tau_{i+1}^0 = \sum_{j=1}^i \tau_j^0$$

[smearing of lifetime]

$$P(\tau_i) = \exp \left(- \tau_i / \langle \tau_i \rangle \right)$$

- modifies PYSHOW
- implements space-time via formation time estimate
- in between splittings
 - “drag”
 - “broadening”
- if $Q^2 \ll \Delta Q^2$: t_f is found iteratively



YaJEM

[lifetime of parton $i+1$]

$$\langle \tau_{i+1} \rangle = \frac{E_{i+1}}{Q_{i+1}^2} - \frac{E_{i+1}}{Q_i^2}$$

Renk, PRC 78 (2008) 034908

[emergence of parton $i+1$]

$$\tau_{i+1}^0 = \sum_{j=1}^i \tau_j^0$$

[smearing of lifetime]

$$P(\tau_i) = \exp \left(- \tau_i / \langle \tau_i \rangle \right)$$

$$\Delta Q_i^2 = \int_{\tau_i^0}^{\tau_i^0 + \tau_i} d\xi \hat{q}(\xi)$$
$$\Delta E_i = \int_{\tau_i^0}^{\tau_i^0 + \tau_i} d\xi D\rho(\xi)$$

- modifies PYSHOW
- implements space-time via formation time estimate
- in between splittings
 - “drag”
 - “broadening”
- if $Q^2 \ll \Delta Q^2$: t_f is found iteratively

$$\hat{q}(\xi) = 2K \varepsilon(\xi)^{3/4} \left[\cosh \rho(\xi) - \sinh \rho(\xi) \cos \psi \right]$$

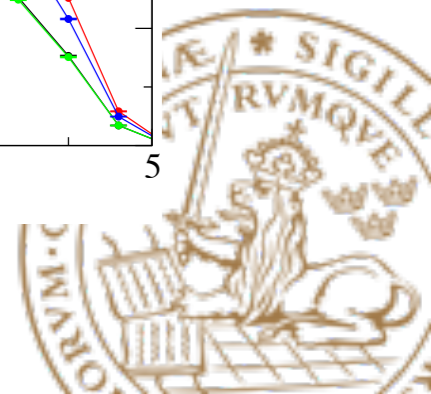
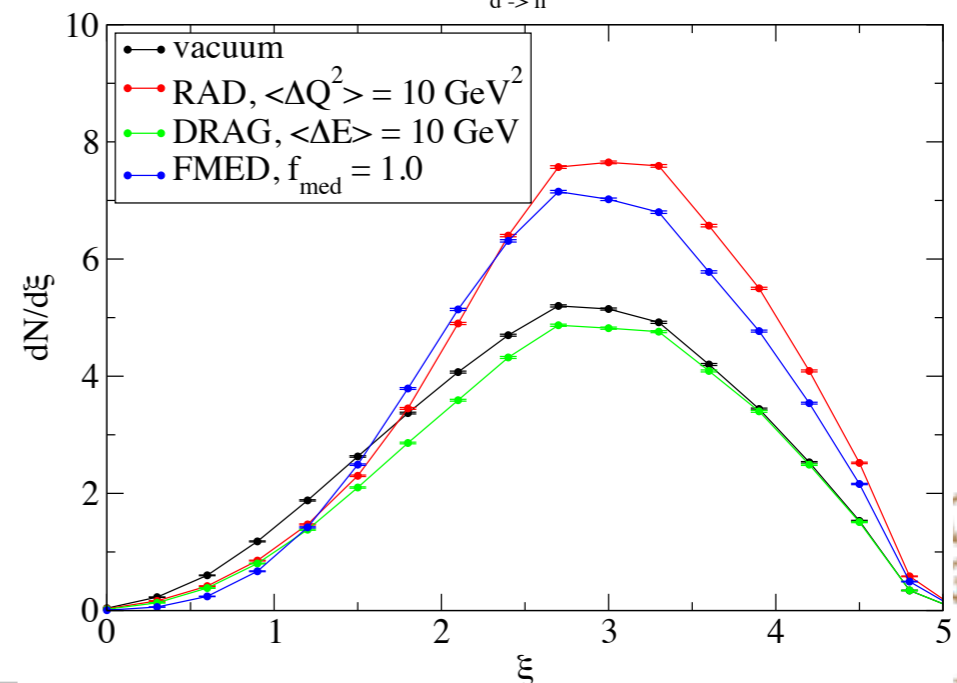
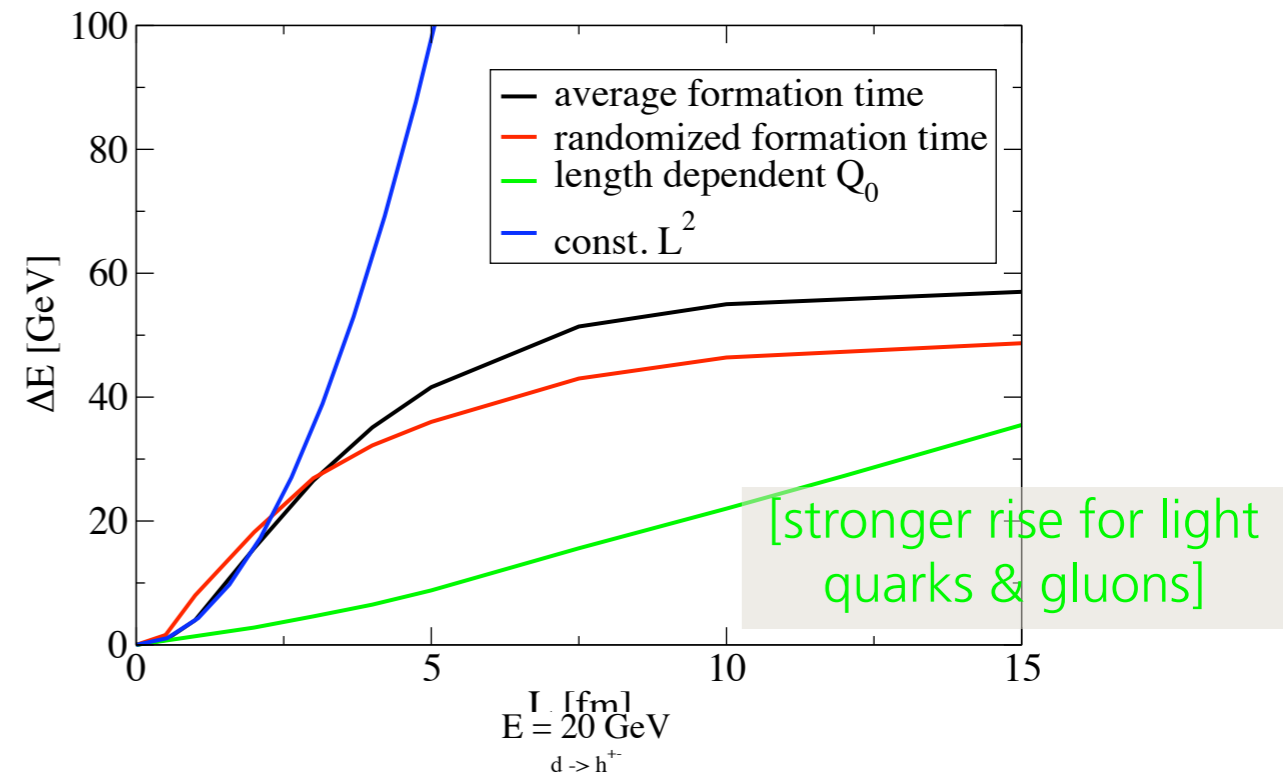


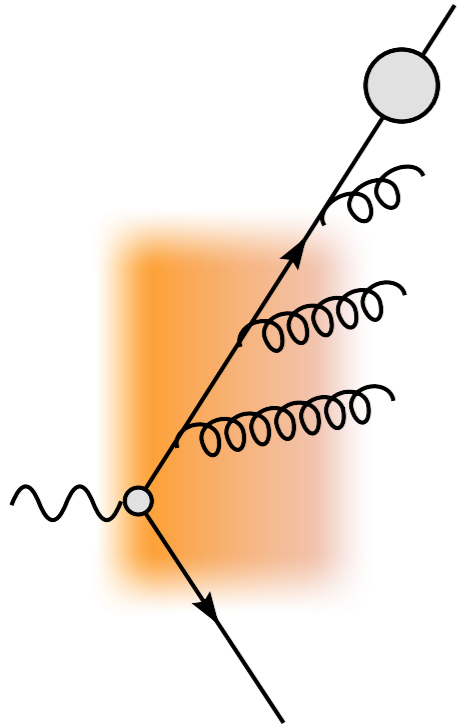
Path-dependence

- many modes
 - **RAD**: only radiative
 - **DRAG**: only drag
 - **FMED**: enhanced singularity in vacuum splitting function
- **ASW**: ?
- **YaJEM-D**: enhanced path-length dependence

[dynamical cut-off in medium]

$$Q_0 = \sqrt{\frac{E}{L}}$$



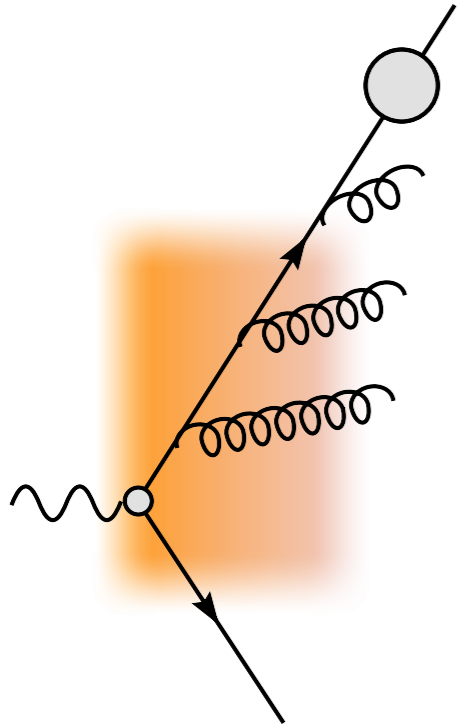


Q-PYTHIA

$$\frac{dN}{dz dk_{\perp}^2} = \frac{dN^{\text{med}}}{dz dk_{\perp}^2} + \frac{\alpha_s}{2\pi} \frac{P(z)}{k_{\perp}^2}$$

Armesto, Cunqueiro,
Salgado EPJC 61 (2009) 775





Q-PYTHIA

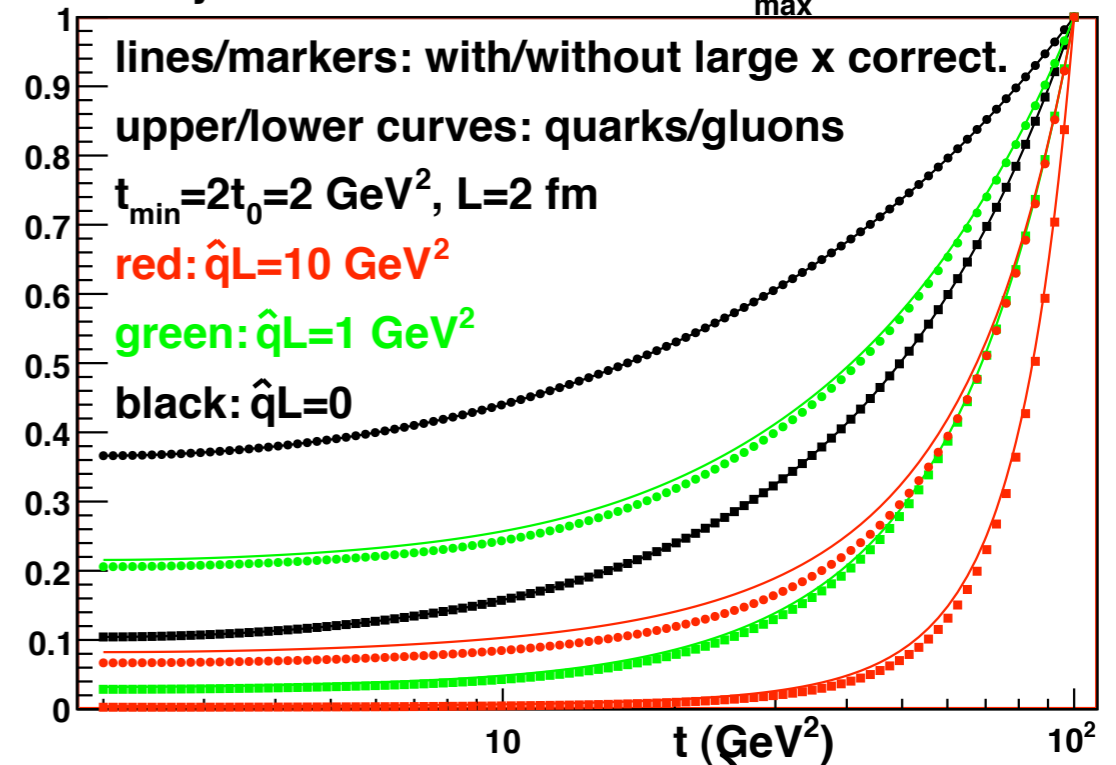
Armesto, Cunqueiro,
Salgado EPJC 61 (2009) 775

$$\frac{dN}{dz dk_{\perp}^2} = \frac{dN^{\text{med}}}{dz dk_{\perp}^2} + \frac{\alpha_s}{2\pi} \frac{P(z)}{k_{\perp}^2}$$

E=10 GeV

- define a medium-modified splitting function
- showering á la DGLAP as implemented in PYTHIA
 - also as Q-HERWIG
- similar in spirit to HT

Probability of no emission between $t_{\min}=E^2=10^2 \text{ GeV}^2$ and t_{\max}



$$\Delta(t_1, t_2) = \exp \left\{ - \int_{t_1}^{t_2} \frac{dt}{t} \int_{z_{\min}(t)}^{1-z_{\min}(t)} dz \frac{\alpha_s}{2\pi} [P(z) + \Delta P(z)] \right\}$$



Some motivation

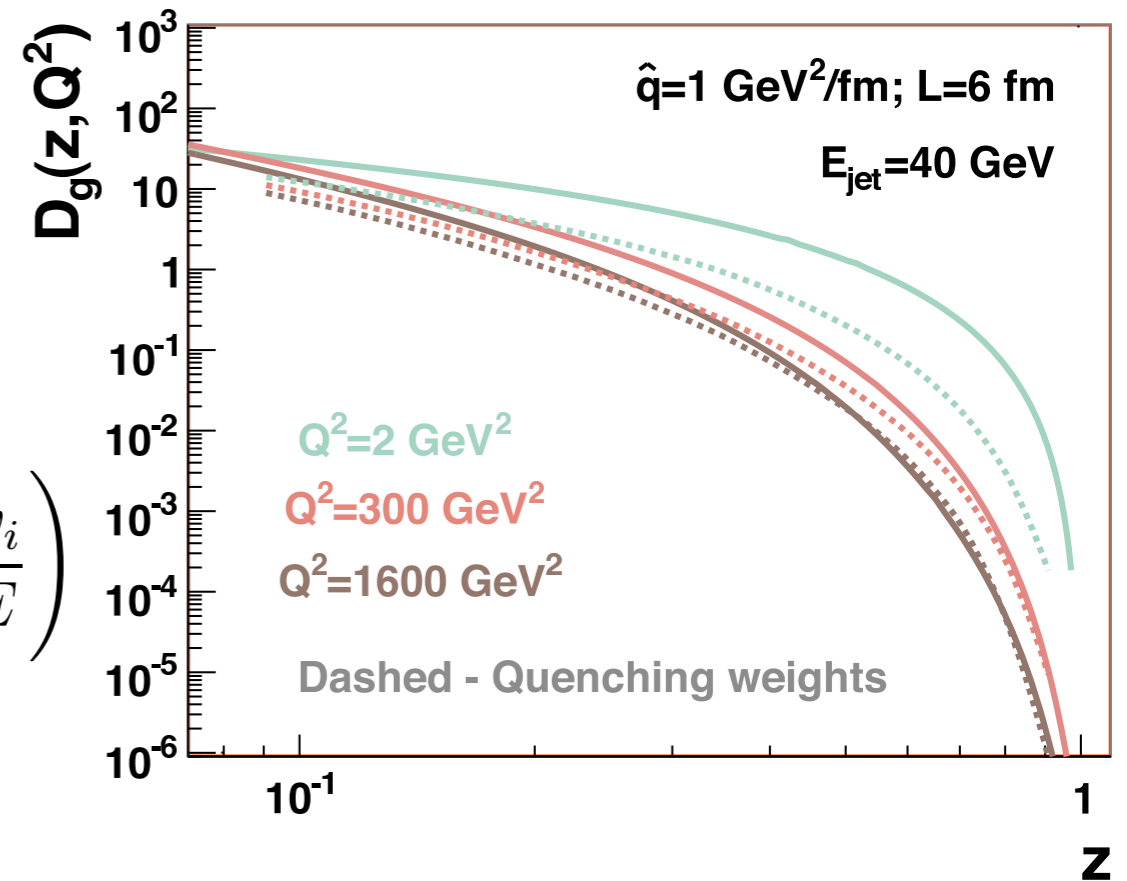
Armesto, Cunqueiro, Salgado, Xiang JHEP 0802 (2008) 048

- **neglecting ordering in virtuality: back to quenching weights!**

$$p(\epsilon) = p_0 \sum_{n=1}^{\infty} \prod_{i=1}^n \int d\omega_i \int d\mathbf{k}_{\perp i} \frac{dI^{\text{med}}}{d\omega_i d\mathbf{k}_{\perp i}} \delta \left(\epsilon - \sum_{j=1}^n \frac{\omega_j}{E} \right)$$

$$p_0 = \exp \left[- \int d\omega \int d\mathbf{k}_{\perp} \frac{dI^{\text{med}}}{d\omega d\mathbf{k}_{\perp}} \right]$$

$$D(x, t) \simeq p_0 D^{\text{vac}}(x, t) + \int \frac{d\epsilon}{1 - \epsilon} p(\epsilon) D^{\text{vac}} \left(\frac{x}{1 - \epsilon}, t \right)$$



Some motivation

Armesto, Cunqueiro, Salgado, Xiang JHEP 0802 (2008) 048

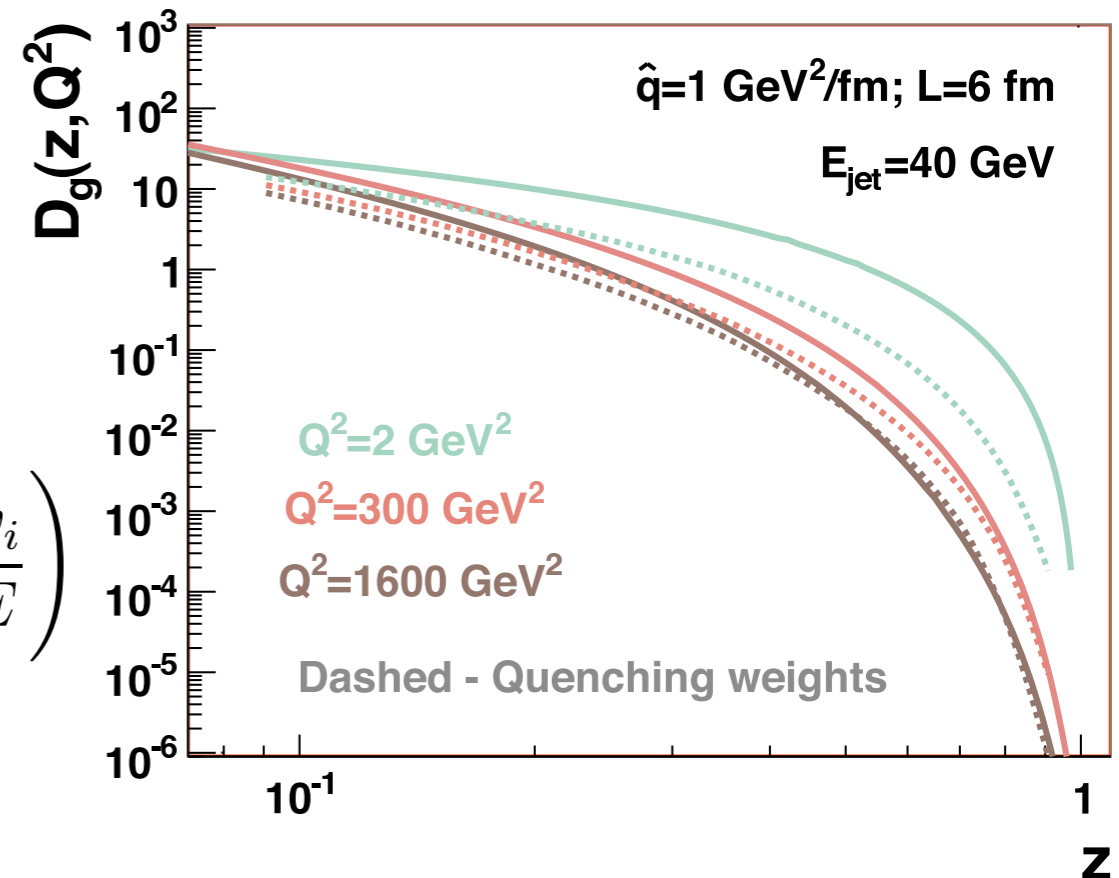
- **neglecting ordering in virtuality: back to quenching weights!**

$$p(\epsilon) = p_0 \sum_{n=1}^{\infty} \prod_{i=1}^n \int d\omega_i \int d\mathbf{k}_{\perp i} \frac{dI^{\text{med}}}{d\omega_i d\mathbf{k}_{\perp i}} \delta \left(\epsilon - \sum_{j=1}^n \frac{\omega_j}{E} \right)$$

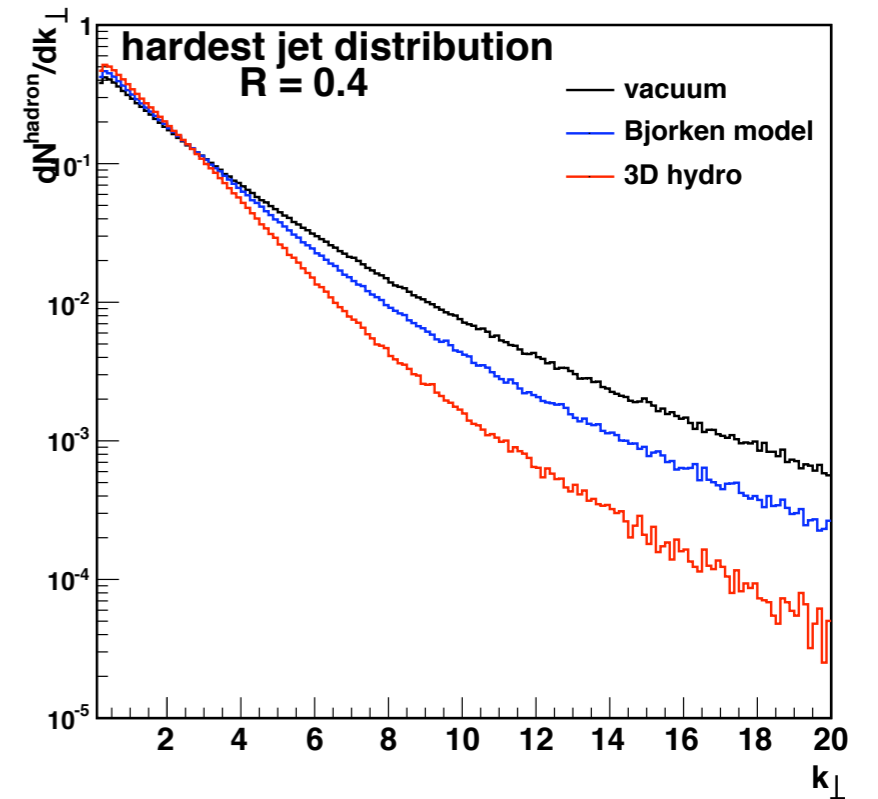
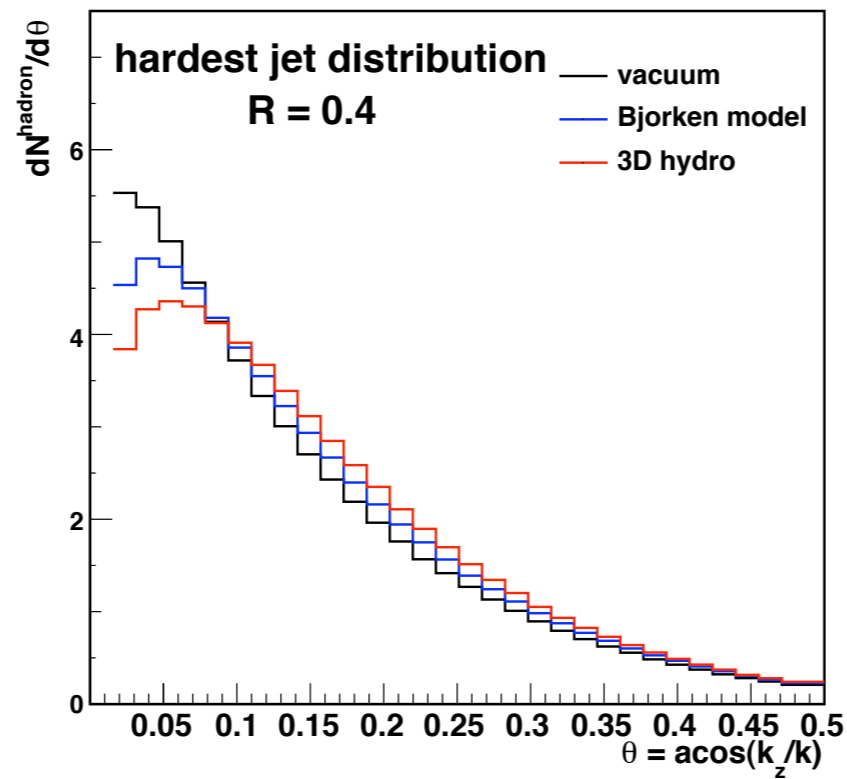
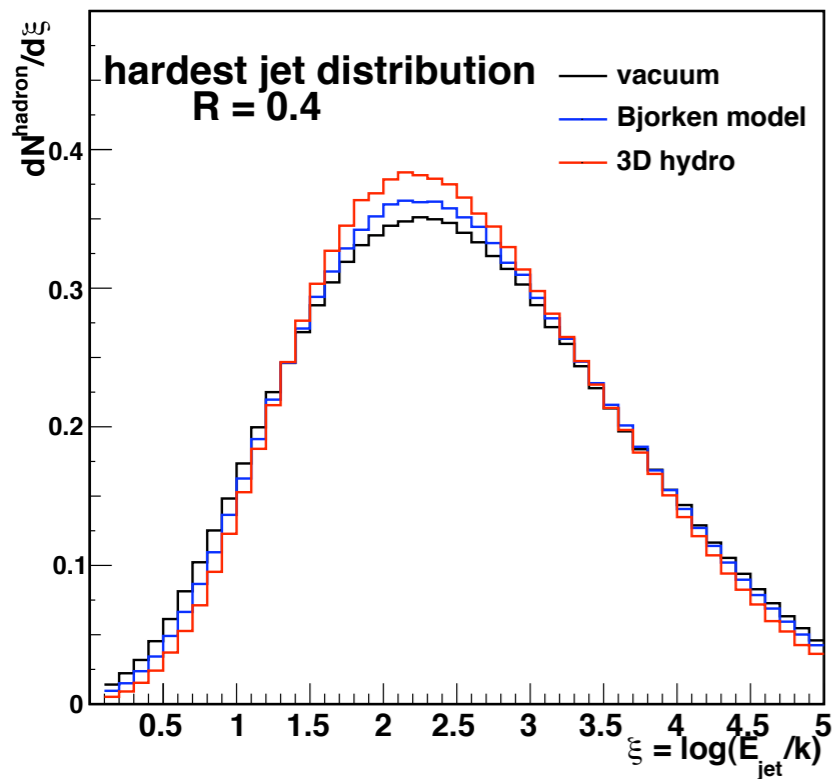
$$p_0 = \exp \left[- \int d\omega \int d\mathbf{k}_{\perp} \frac{dI^{\text{med}}}{d\omega d\mathbf{k}_{\perp}} \right]$$

$$D(x, t) \simeq p_0 D^{\text{vac}}(x, t) + \int \frac{d\epsilon}{1-\epsilon} p(\epsilon) D^{\text{vac}} \left(\frac{x}{1-\epsilon}, t \right)$$

Underlying: medium-rad small correction to virtuality-driven evolution



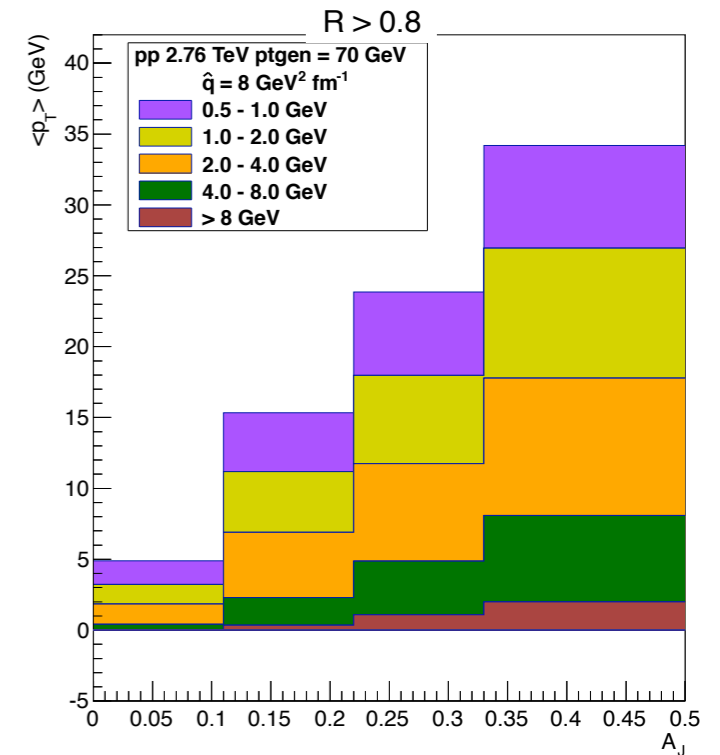
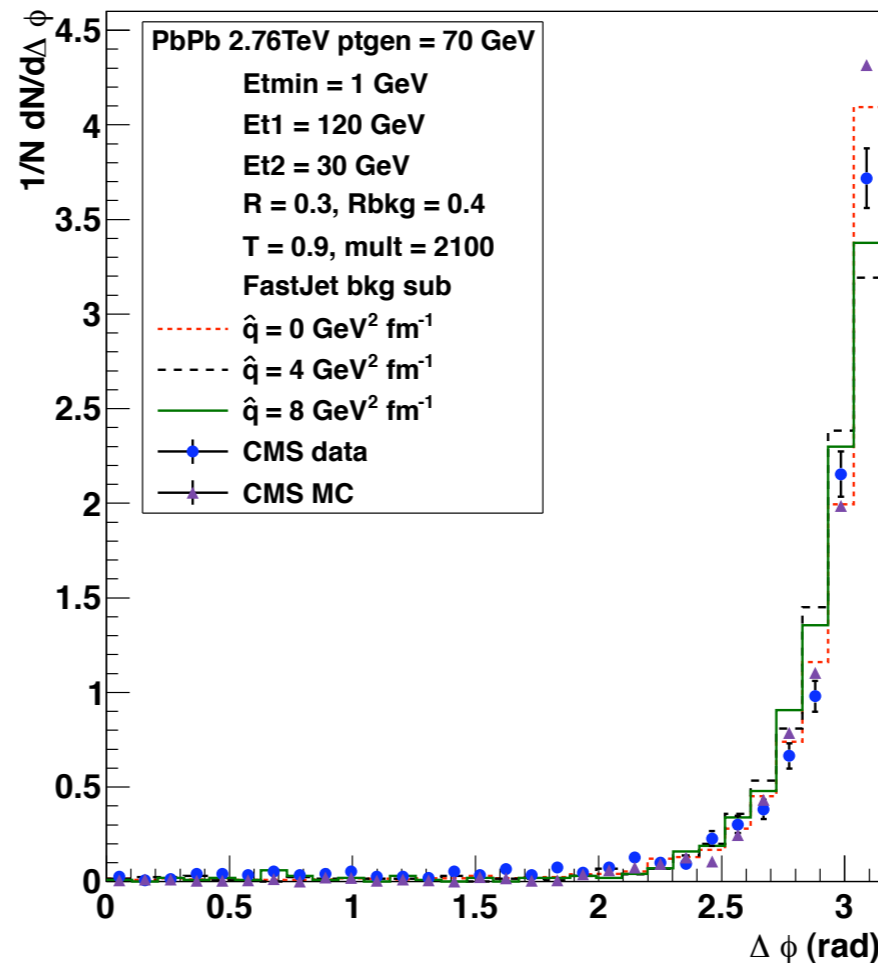
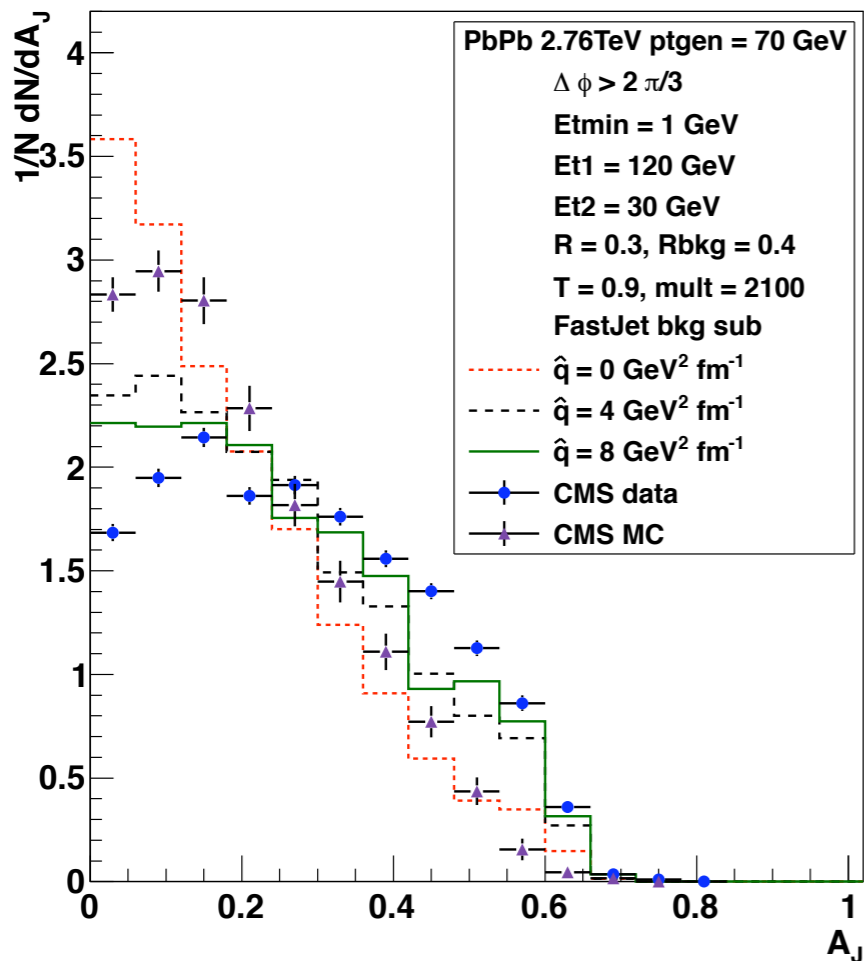
Selected results



- inter-jet distributions
- effects of expanding medium on jet quenching



Dijet asymmetry



- embedded in background (PSM)
- versatile tool

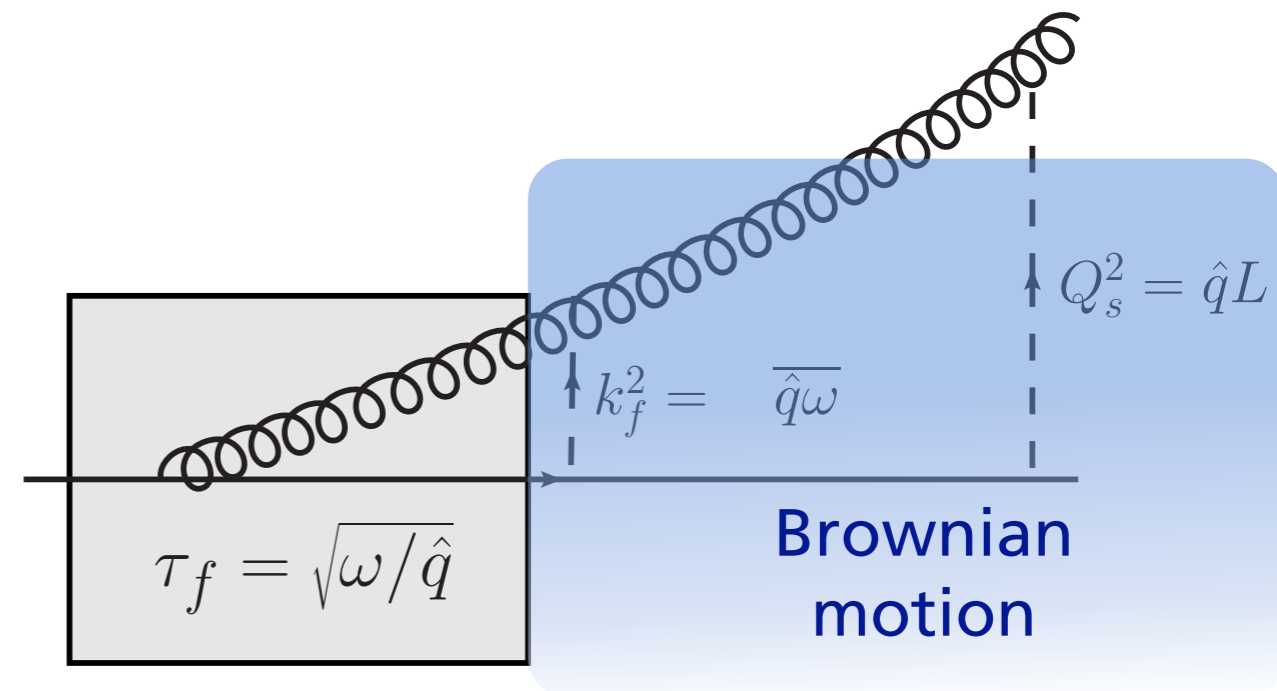
Apollinario, Armesto, Cunqueiro
 arXiv:1207.6587



A local implementation of LPM

Stachel, Wiedemann, Zapp PRL 103 (2009) 152302,
JHEP 1107 (2011) 118

- medium-induced radiation:
longitudinal coherence
- probabilistic picture interpolating
between known limits



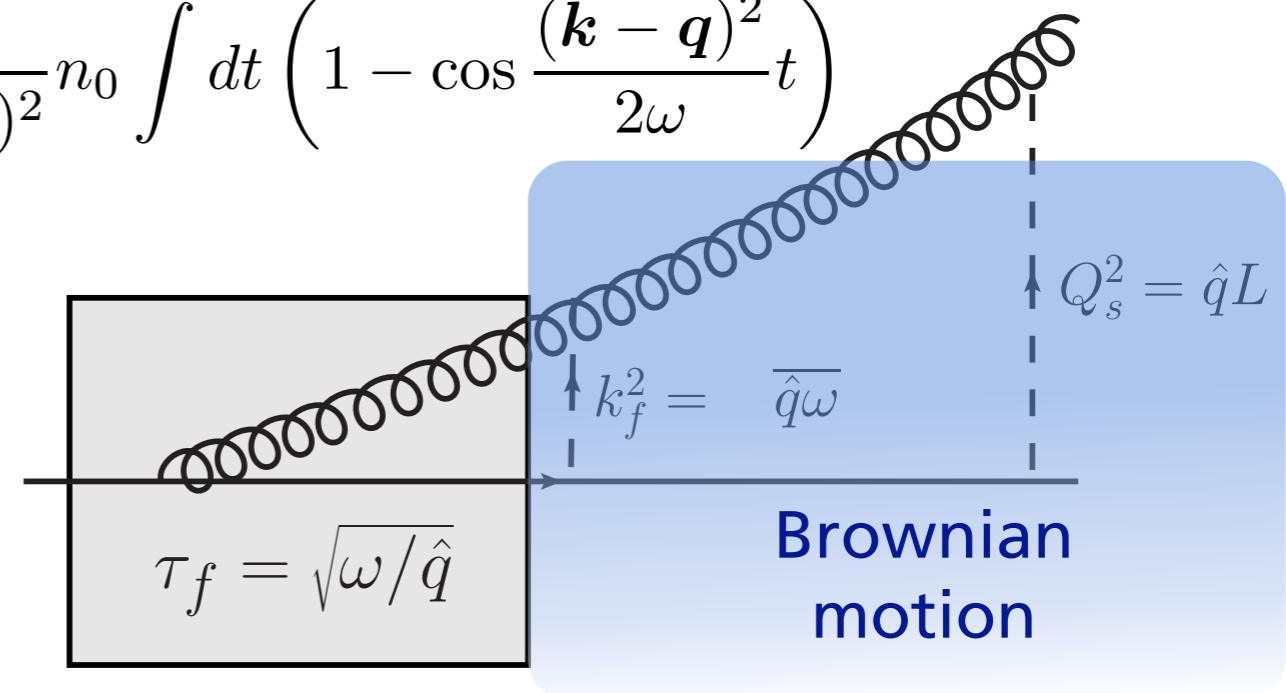
A local implementation of LPM

Stachel, Wiedemann, Zapp PRL 103 (2009) 152302,
JHEP 1107 (2011) 118

- medium-induced radiation:
longitudinal coherence
- probabilistic picture interpolating
between known limits

$$\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = \frac{\alpha_s C_R}{\pi^2} \frac{|A(\mathbf{q})|^2}{(2\pi)^2} \frac{2\mathbf{k} \cdot \mathbf{q}}{\mathbf{k}^2 (\mathbf{k} - \mathbf{q})^2} n_0 \int dt \left(1 - \cos \frac{(\mathbf{k} - \mathbf{q})^2 t}{2\omega} \right)$$

[at first order in medium opacity]



A local implementation of LPM

Stachel, Wiedemann, Zapp PRL 103 (2009) 152302,
JHEP 1107 (2011) 118

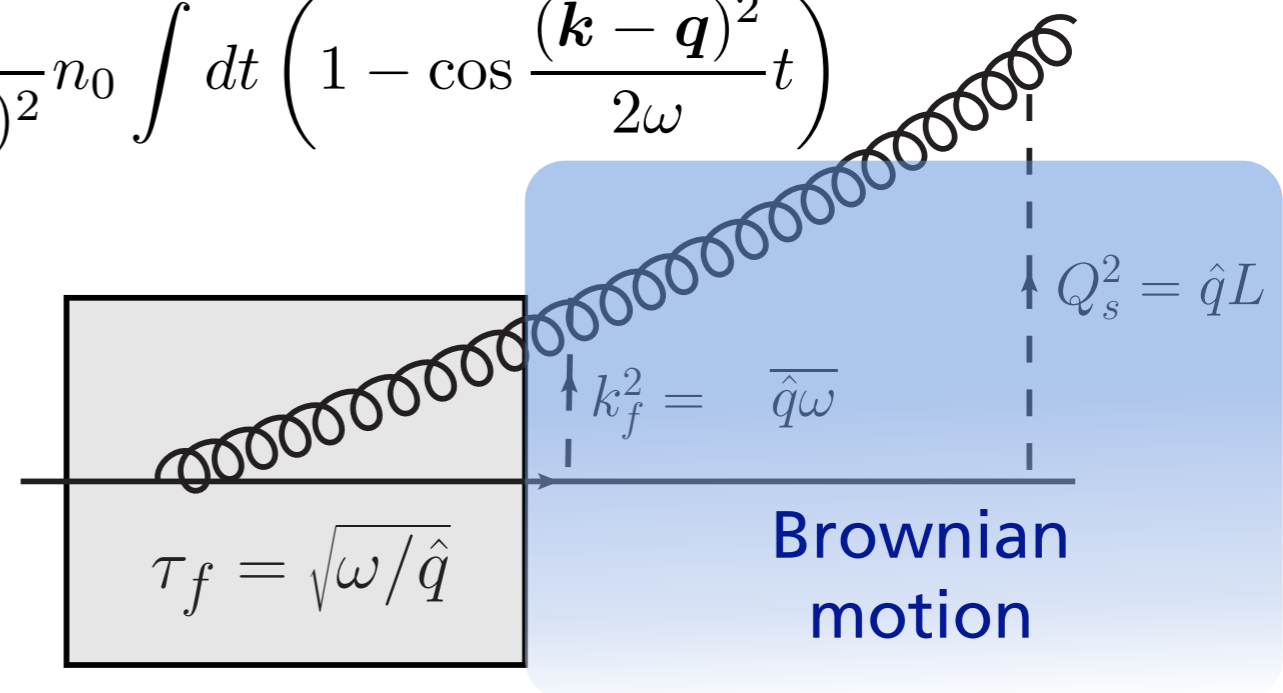
- medium-induced radiation:
longitudinal coherence
- probabilistic picture interpolating
between known limits

$$\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = \frac{\alpha_s C_R}{\pi^2} \frac{|A(\mathbf{q})|^2}{(2\pi)^2} \frac{2\mathbf{k} \cdot \mathbf{q}}{\mathbf{k}^2 (\mathbf{k} - \mathbf{q})^2} n_0 \int dt \left(1 - \cos \frac{(\mathbf{k} - \mathbf{q})^2 t}{2\omega} \right)$$

[at first order in medium opacity]

$$V_{\text{tot}} = \int \frac{d^2\mathbf{q}}{(2\pi)^2} |A(\mathbf{q})|^2 = \frac{\alpha_s C_A}{\pi} \int d^2\mathbf{q} \frac{m_D^2}{(\mathbf{q}^2 + m_D^2)^2}$$

[prob of scattering with the medium]



A local implementation of LPM

Stachel, Wiedemann, Zapp PRL 103 (2009) 152302,
JHEP 1107 (2011) 118

- medium-induced radiation:
longitudinal coherence
- probabilistic picture interpolating
between known limits

[formation time prior
to rescattering]

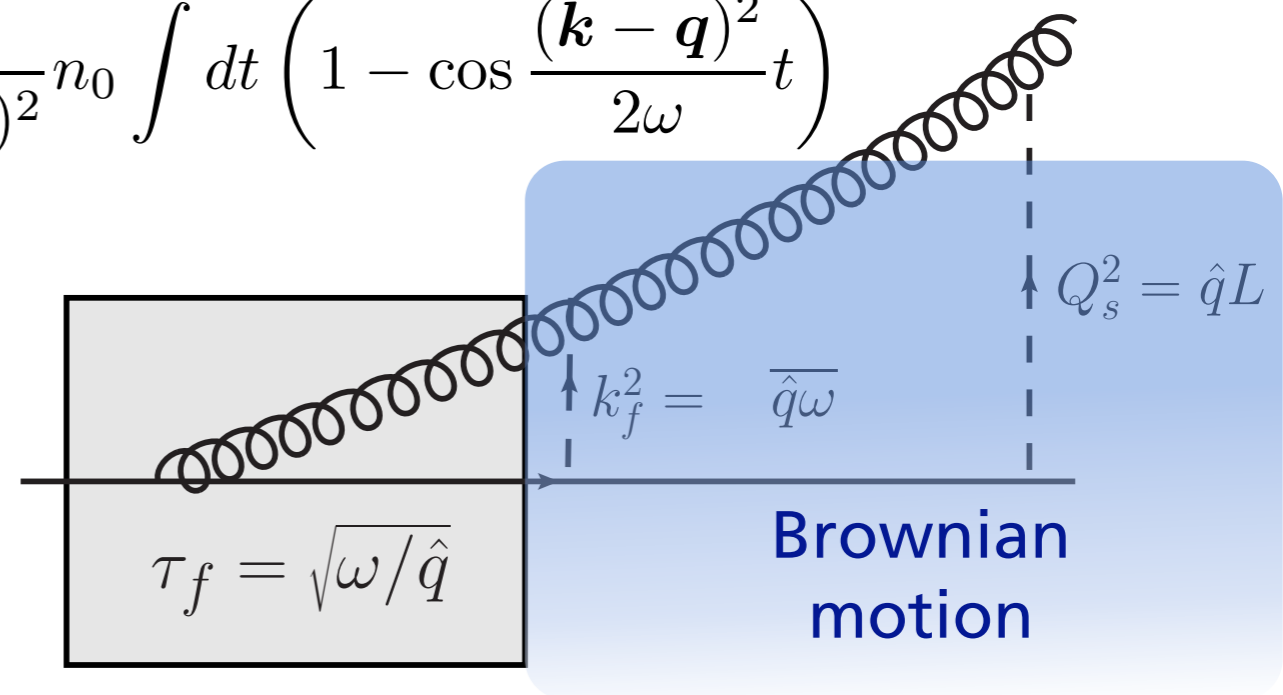
$$t_f^{(1)} = \frac{2\omega}{(\mathbf{k} - \mathbf{q})^2}$$

$$\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = \frac{\alpha_s C_R}{\pi^2} \frac{|A(\mathbf{q})|^2}{(2\pi)^2} \frac{2\mathbf{k} \cdot \mathbf{q}}{\mathbf{k}^2 (\mathbf{k} - \mathbf{q})^2} n_0 \int dt \left(1 - \cos \frac{(\mathbf{k} - \mathbf{q})^2 t}{2\omega} \right)$$

[at first order in medium opacity]

$$V_{\text{tot}} = \int \frac{d^2\mathbf{q}}{(2\pi)^2} |A(\mathbf{q})|^2 = \frac{\alpha_s C_A}{\pi} \int d^2\mathbf{q} \frac{m_D^2}{(\mathbf{q}^2 + m_D^2)^2}$$

[prob of scattering with the medium]



Two illustrative limits

Coherent limit: $L \ll t_f^{(1)}$ $\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = 0$



Two illustrative limits

Coherent limit: $L \ll t_f^{(1)}$ $\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = 0$

Incoherent limit: $L \gg t_f^{(1)}$

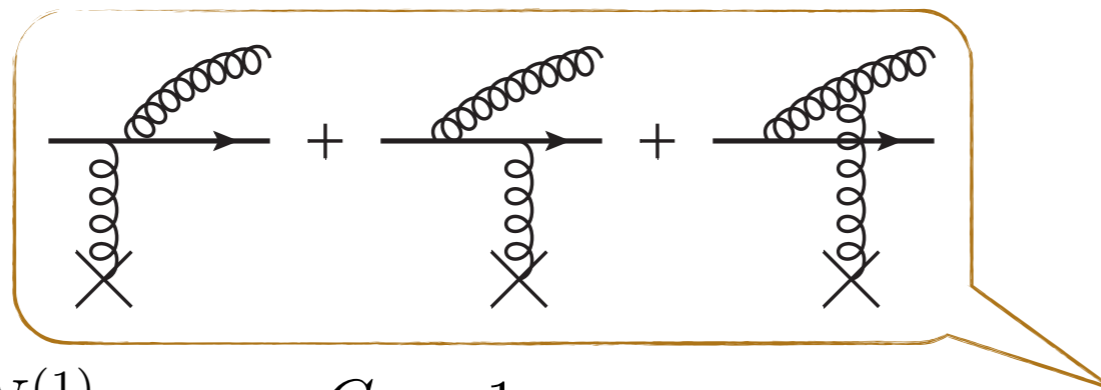
$$\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = \frac{\alpha_s C_R}{\pi^2} \frac{1}{(2\pi)^2} (n_0 L) \left\{ |A(\mathbf{q})|^2 [R(\mathbf{k}, \mathbf{q}) + H(\mathbf{k} - \mathbf{q})] - V_{\text{tot}} \bar{\delta}(\mathbf{q}) H(\mathbf{k}) \right\}$$



Two illustrative limits

Coherent limit: $L \ll t_f^{(1)}$ $\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = 0$

Incoherent limit: $L \gg t_f^{(1)}$



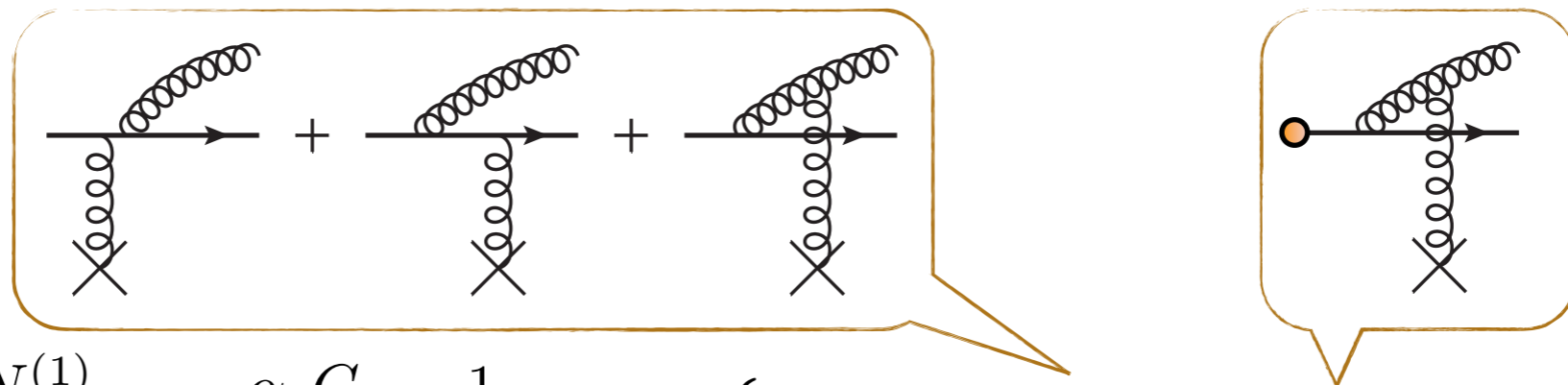
$$\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = \frac{\alpha_s C_R}{\pi^2} \frac{1}{(2\pi)^2} (n_0 L) \left\{ |A(\mathbf{q})|^2 [R(\mathbf{k}, \mathbf{q}) + H(\mathbf{k} - \mathbf{q})] - V_{\text{tot}} \bar{\delta}(\mathbf{q}) H(\mathbf{k}) \right\}$$



Two illustrative limits

Coherent limit: $L \ll t_f^{(1)}$ $\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = 0$

Incoherent limit: $L \gg t_f^{(1)}$



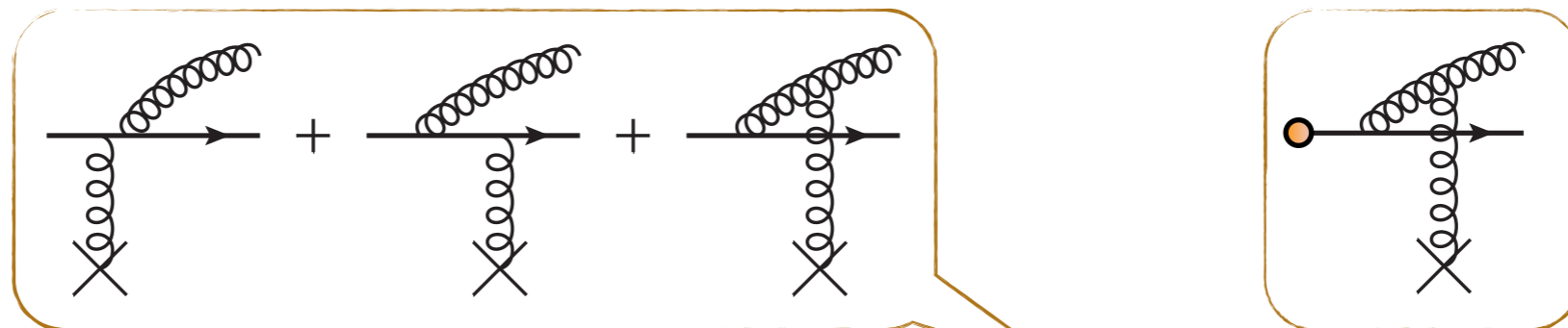
$$\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = \frac{\alpha_s C_R}{\pi^2} \frac{1}{(2\pi)^2} (n_0 L) \left\{ |A(\mathbf{q})|^2 [R(\mathbf{k}, \mathbf{q}) + H(\mathbf{k} - \mathbf{q})] - V_{\text{tot}} \bar{\delta}(\mathbf{q}) H(\mathbf{k}) \right\}$$



Two illustrative limits

Coherent limit: $L \ll t_f^{(1)}$ $\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = 0$

Incoherent limit: $L \gg t_f^{(1)}$



$$\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = \frac{\alpha_s C_R}{\pi^2} \frac{1}{(2\pi)^2} (n_0 L) \left\{ |A(\mathbf{q})|^2 [R(\mathbf{k}, \mathbf{q}) + H(\mathbf{k} - \mathbf{q})] - V_{\text{tot}} \bar{\delta}(\mathbf{q}) H(\mathbf{k}) \right\}$$

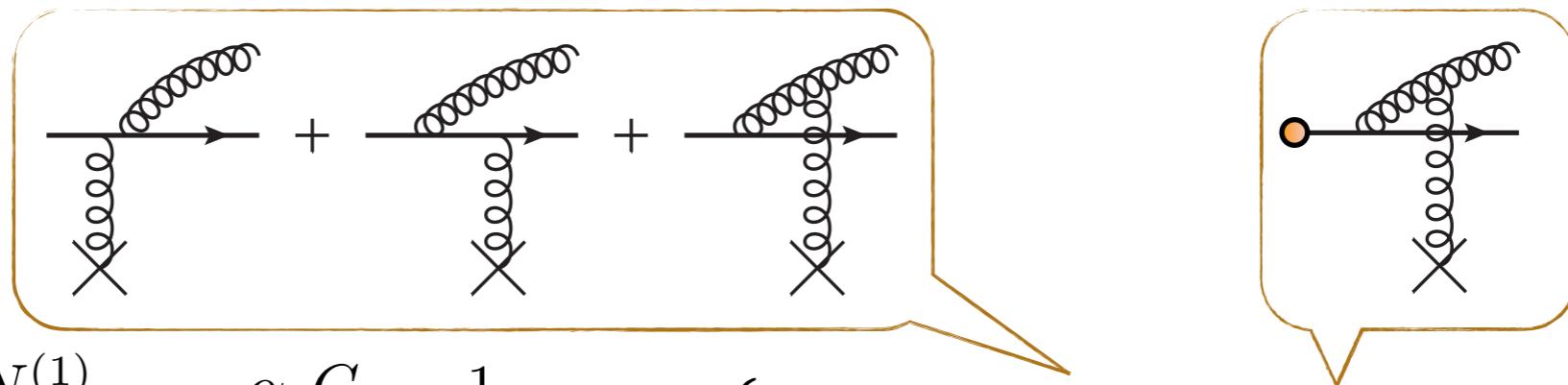
Conservation of probability!



Two illustrative limits

Coherent limit: $L \ll t_f^{(1)}$ $\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = 0$

Incoherent limit: $L \gg t_f^{(1)}$



$$\omega \frac{dN^{(1)}}{d\omega d^2\mathbf{k} d^2\mathbf{q}} = \frac{\alpha_s C_R}{\pi^2} \frac{1}{(2\pi)^2} (n_0 L) \left\{ |A(\mathbf{q})|^2 [R(\mathbf{k}, \mathbf{q}) + H(\mathbf{k} - \mathbf{q})] - V_{\text{tot}} \bar{\delta}(\mathbf{q}) H(\mathbf{k}) \right\}$$

Conservation of probability!

Resummed form factor: $S_{\text{el}} = \exp \left[- n_0 L V_{\text{tot}} \right]$



MC implementation

- neglecting final-state rescattering of vacuum radiation
 - equivalent to radiation off asymptotic charge
 - perfectly ok for $\omega \, dN/d\omega$
 - rescattering of induced radiation included in $R(k + \sum q_{rescat}, \sum q_{induced})$
- veto radiation that hasn't been formed inside the medium (qualitative guidance)
- dynamical determination of t_f
- numerical simplification: $R(\mathbf{k}, \mathbf{q}) \sim \delta(\mathbf{k} - \mathbf{q})$



Space-time propagation

- define mean path lengths for rescattering and radiation
- emitter scatters only inelastically, “emitee” scatters only elastically

$$\lambda_{\text{el}} \equiv \frac{1}{n_0 V_{\text{tot}}} \quad \lambda_{\text{inel}} \equiv \frac{1}{n_0 \sigma_{\text{inel}}}$$

BDMPS-Z limit: $\lambda_{\text{el}} \ll \lambda_{\text{inel}}$



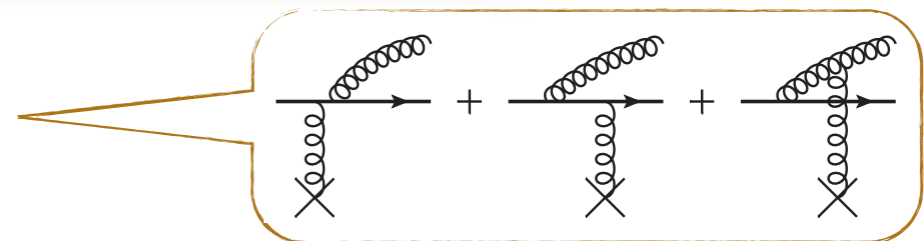
Space-time propagation

- define mean path lengths for rescattering and radiation
- emitter scatters only inelastically, “emitee” scatters only elastically

$$\lambda_{\text{el}} \equiv \frac{1}{n_0 V_{\text{tot}}} \quad \lambda_{\text{inel}} \equiv \frac{1}{n_0 \sigma_{\text{inel}}}$$

BDMPS-Z limit: $\lambda_{\text{el}} \ll \lambda_{\text{inel}}$

$$\sigma_{\text{inel}} \equiv \frac{\alpha_s C_R}{\pi^2} \int d\omega d^2 \mathbf{k} \int \frac{d^2 \mathbf{q}}{(2\pi)^2} |A(\mathbf{q})|^2 \frac{1}{\omega} R(\mathbf{k}, \mathbf{q})$$



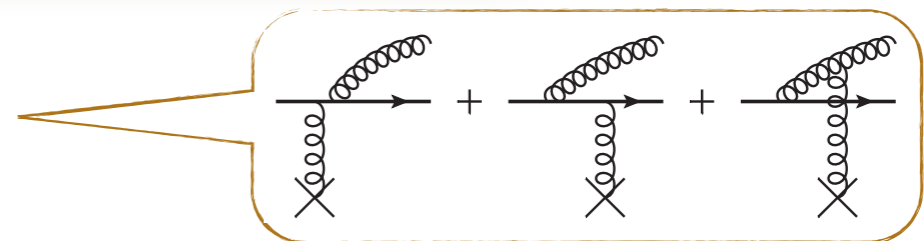
Space-time propagation

- define mean path lengths for rescattering and radiation
- emitter scatters only inelastically, “emitee” scatters only elastically

$$\lambda_{\text{el}} \equiv \frac{1}{n_0 V_{\text{tot}}} \quad \lambda_{\text{inel}} \equiv \frac{1}{n_0 \sigma_{\text{inel}}}$$

BDMPS-Z limit: $\lambda_{\text{el}} \ll \lambda_{\text{inel}}$

$$\sigma_{\text{inel}} \equiv \frac{\alpha_s C_R}{\pi^2} \int d\omega d^2 \mathbf{k} \int \frac{d^2 \mathbf{q}}{(2\pi)^2} |A(\mathbf{q})|^2 \frac{1}{\omega} R(\mathbf{k}, \mathbf{q})$$



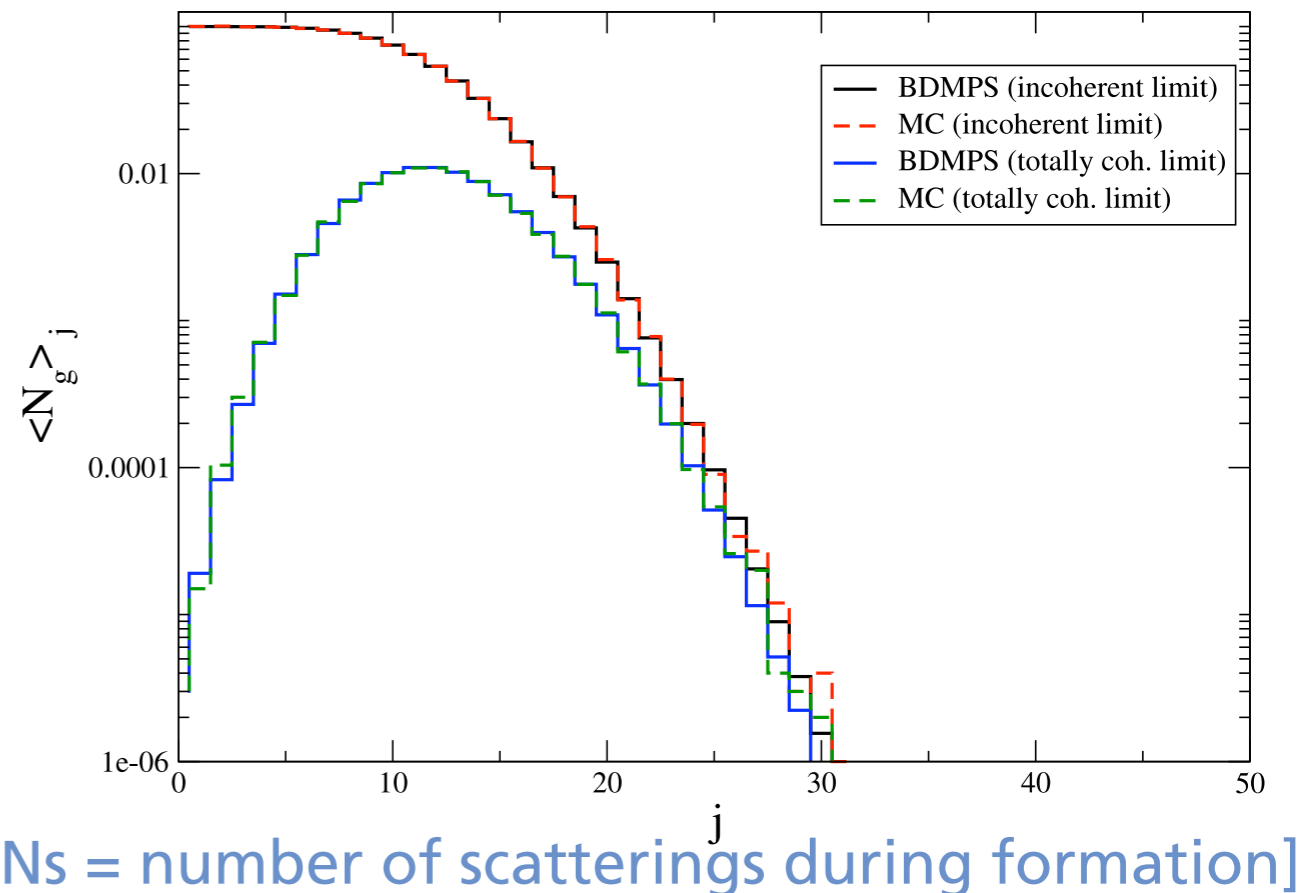
$$S_{\text{no}}^{(\text{in})\text{el}}(L) = \exp(-L/\lambda_{(\text{in})\text{el}})$$



Subtleties due to coherence

- t_f establishes a “zone” for emission
- $t_f \gg L$: elastic re-scatterings can take place along the **whole** (not remaining) path
- re-weight by $(N_s)^{-1}$

$\lambda_{\text{inel}} = 1.0 \text{ fm}, \quad \lambda_{\text{el}} = 0.1 \text{ fm}, \quad L = 1.3 \text{ fm}$



$$\langle N_g^{\text{incoh}} \rangle (N_s) = N_s \langle N_g^{\text{coh}} \rangle (N_s)$$



MC algorithm



MC algorithm

- determine position of (next) inelastic scattering: $1 - S^{\text{inel}}_{n_0}(L - \xi)$



MC algorithm

- determine position of (next) inelastic scattering: $1 - S^{\text{inel}}_{n_0}(L - \xi)$
- trial gluon is produced (remaining trial path length is $L_g = L!$)



MC algorithm

- determine position of (next) inelastic scattering: $1-S^{\text{inel}}_{no}(L-\xi)$
- trial gluon is produced (remaining trial path length is $L_g=L!$)
 - ω, k_{\perp} generated according to ω^{-1} and $|A(k_{\perp})|^2$ - gives trial t_f



MC algorithm

- determine position of (next) inelastic scattering: $1-S^{\text{inel}}_{\text{no}}(L-\xi)$
- trial gluon is produced (remaining trial path length is $L_g=L!$)
 - ω, k_{\perp} generated according to ω^{-1} and $|A(k_{\perp})|^2$ - gives trial t_f
 - coherent re-scattering with prob $1-S^{\text{el}}_{\text{no}}(\min(t_f, L_g))$ at distance ΔL according to $S^{\text{el}}_{\text{no}}(\Delta L)/\lambda_{\text{el}}$



MC algorithm

- determine position of (next) inelastic scattering: $1-S^{\text{inel}}_{\text{no}}(L-\xi)$
- trial gluon is produced (remaining trial path length is $L_g=L!$)
 - ω, k_{\perp} generated according to ω^{-1} and $|A(k_{\perp})|^2$ - gives trial t_f
 - coherent re-scattering with prob $1-S^{\text{el}}_{\text{no}}(\min(t_f, L_g))$ at distance ΔL according to $S^{\text{el}}_{\text{no}}(\Delta L)/\lambda_{\text{el}}$
 - update remaining path length L_g to $L_g - \Delta L$



MC algorithm

- determine position of (next) inelastic scattering: $1-S^{\text{inel}}_{\text{no}}(L-\xi)$
- trial gluon is produced (remaining trial path length is $L_g=L!$)
 - ω, k_{\perp} generated according to ω^{-1} and $|A(k_{\perp})|^2$ - gives trial t_f
 - coherent re-scattering with prob $1-S^{\text{el}}_{\text{no}}(\min(t_f, L_g))$ at distance ΔL according to $S^{\text{el}}_{\text{no}}(\Delta L)/\lambda_{\text{el}}$
 - update remaining path length L_g to $L_g - \Delta L$
 - update t_f to $t_f + \Delta L k_{\perp}^2/2\omega$



MC algorithm

- determine position of (next) inelastic scattering: $1-S^{\text{inel}}_{\text{no}}(L-\xi)$
- trial gluon is produced (remaining trial path length is $L_g=L!$)
 - ω, k_{\perp} generated according to ω^{-1} and $|A(k_{\perp})|^2$ - gives trial t_f
 - coherent re-scattering with prob $1-S^{\text{el}}_{\text{no}}(\min(t_f, L_g))$ at distance ΔL according to $S^{\text{el}}_{\text{no}}(\Delta L)/\lambda_{\text{el}}$
 - update remaining path length L_g to $L_g - \Delta L$
 - update t_f to $t_f + \Delta L k_{\perp}^2/2\omega$
 - find momentum transfer according to $|A(q_{\perp})|^2$ and update k_{\perp}



MC algorithm

- determine position of (next) inelastic scattering: $1-S^{\text{inel}}_{\text{no}}(L-\xi)$
- trial gluon is produced (remaining trial path length is $L_g=L!$)
 - ω, k_{\perp} generated according to ω^{-1} and $|A(k_{\perp})|^2$ - gives trial t_f
 - coherent re-scattering with prob $1-S^{\text{el}}_{\text{no}}(\min(t_f, L_g))$ at distance ΔL according to $S^{\text{el}}_{\text{no}}(\Delta L)/\lambda_{\text{el}}$
 - update remaining path length L_g to $L_g - \Delta L$
 - update t_f to $t_f + \Delta L k_{\perp}^2/2\omega$
 - find momentum transfer according to $|A(q_{\perp})|^2$ and update k_{\perp}
 - ... continue until no further re-scattering is found



MC algorithm

- determine position of (next) inelastic scattering: $1-S^{\text{inel}}_{\text{no}}(L-\xi)$
- trial gluon is produced (remaining trial path length is $L_g=L!$)
 - ω, k_{\perp} generated according to ω^{-1} and $|A(k_{\perp})|^2$ - gives trial t_f
 - coherent re-scattering with prob $1-S^{\text{el}}_{\text{no}}(\min(t_f, L_g))$ at distance ΔL according to $S^{\text{el}}_{\text{no}}(\Delta L)/\lambda_{\text{el}}$
 - update remaining path length L_g to $L_g - \Delta L$
 - update t_f to $t_f + \Delta L k_{\perp}^2/2\omega$
 - find momentum transfer according to $|A(q_{\perp})|^2$ and update k_{\perp}
 - ... continue until no further re-scattering is found
 - **accept gluon with prob N_s^{-1}** and determine a random production point in a box with size t_f around ξ



MC algorithm

- determine position of (next) inelastic scattering: $1-S^{\text{inel}}_{\text{no}}(L-\xi)$
- trial gluon is produced (remaining trial path length is $L_g=L!$)
 - ω, k_{\perp} generated according to ω^{-1} and $|A(k_{\perp})|^2$ - gives trial t_f
 - coherent re-scattering with prob $1-S^{\text{el}}_{\text{no}}(\min(t_f, L_g))$ at distance ΔL according to $S^{\text{el}}_{\text{no}}(\Delta L)/\lambda_{\text{el}}$
 - update remaining path length L_g to $L_g - \Delta L$
 - update t_f to $t_f + \Delta L k_{\perp}^2/2\omega$
 - find momentum transfer according to $|A(q_{\perp})|^2$ and update k_{\perp}
 - ... continue until no further re-scattering is found
 - **accept gluon with prob N_s^{-1}** and determine a random production point in a box with size t_f around ξ
 - propagate further to find remaining elastic scatterings



Limit: BDMPS-Z spectrum

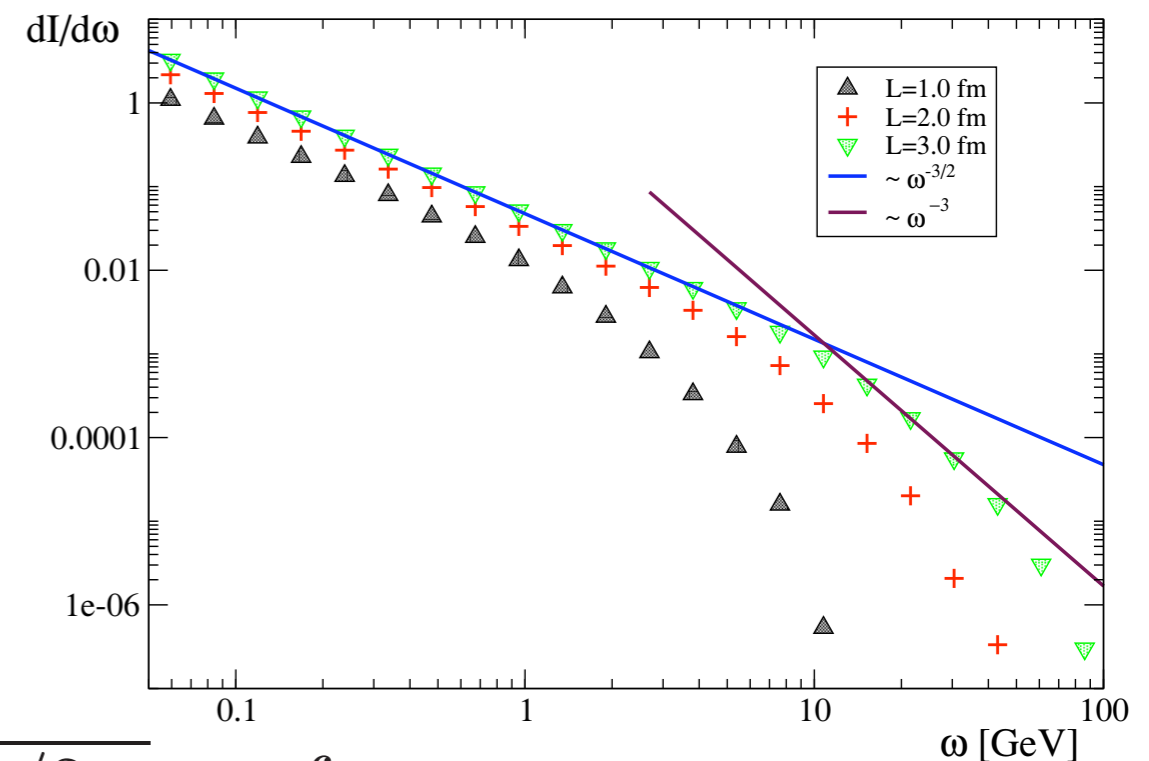
- “BDMPS-Z scenario”
 - only soft scattering
 - only projectile is allowed to radiate
- spectrum is nicely reproduced with applied approximations
- allows to go beyond!

[see Fabio’s talk, Wed]

$$|A(\mathbf{q})|^2 \rightarrow |A(\mathbf{q})|^2 \Theta(2m_D - \mathbf{q})$$

$$\lambda_{\text{inel}}=0.1 \text{ fm}, \lambda_{\text{el}}=0.01 \text{ fm}, \mu=0.2 \text{ GeV}, \omega_{\text{max}}=100 \text{ GeV}$$

Multiple soft scattering limit

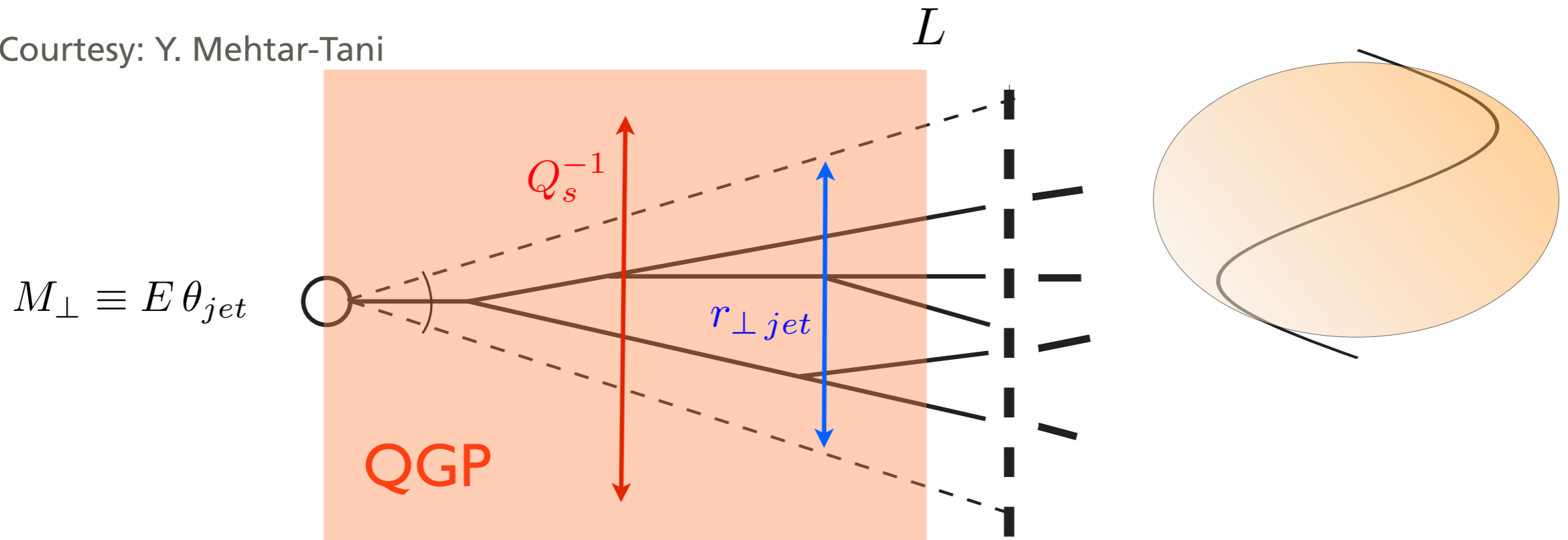


$$\omega \frac{dI}{d\omega} \simeq \frac{2\alpha_s C_R}{\pi} \begin{cases} \sqrt{\omega_c/2\omega} & \text{for } \omega \ll \omega_c \\ \frac{1}{12} \left(\frac{\omega_c}{\omega}\right)^2 & \text{for } \omega \gg \omega_c \end{cases}$$



Vacuum-medium interface

Courtesy: Y. Mehtar-Tani

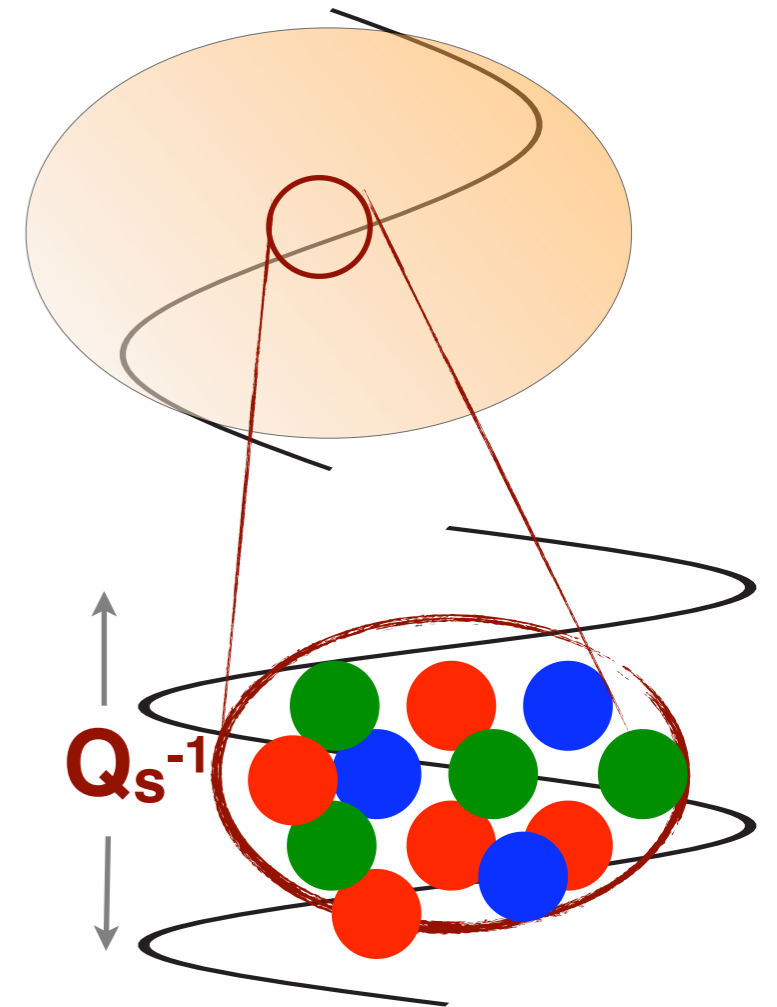
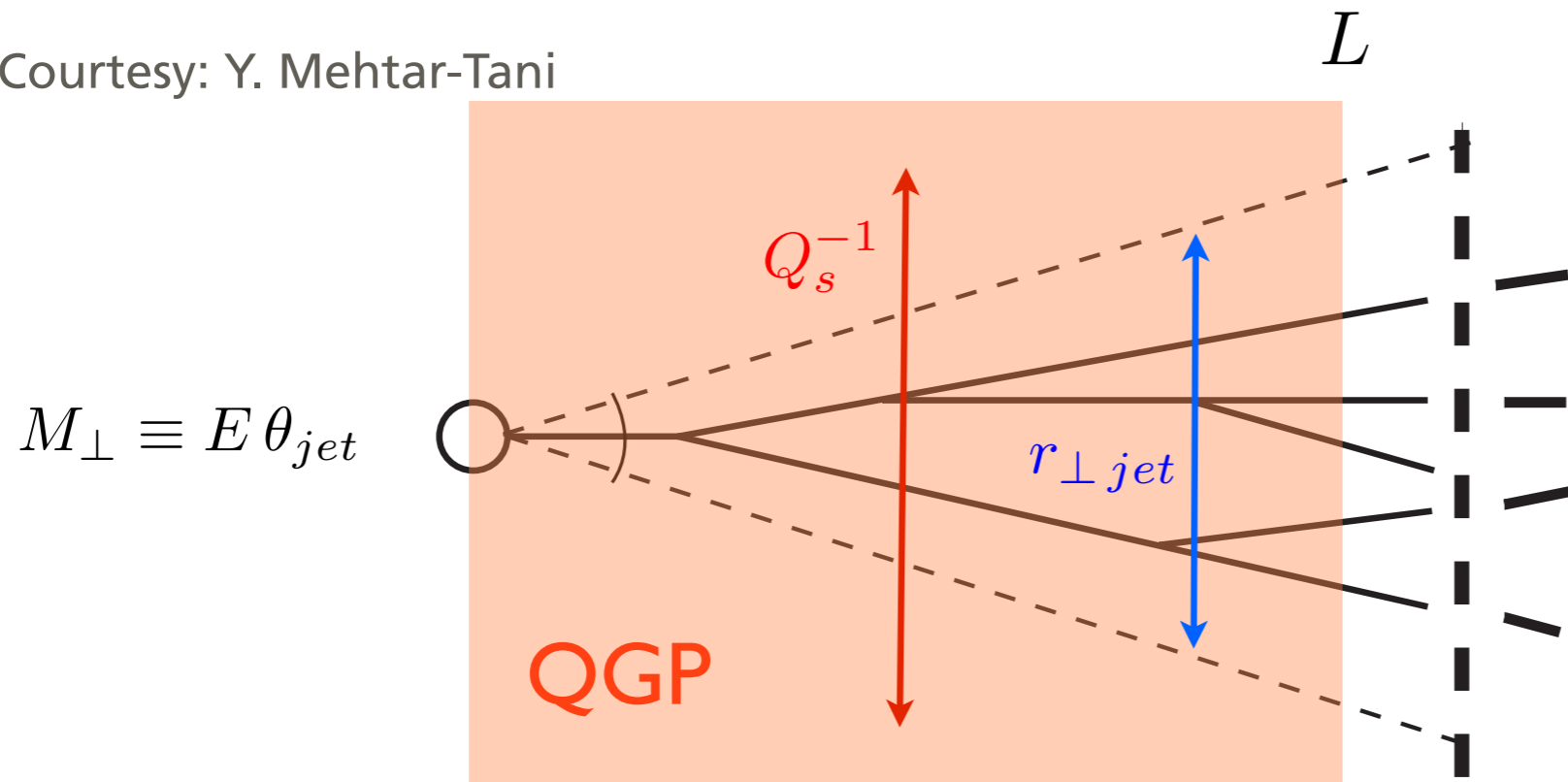


- **bremstrahlung** is treated quite differently in the various codes
- **separation of scales** at LHC!



Vacuum-medium interface

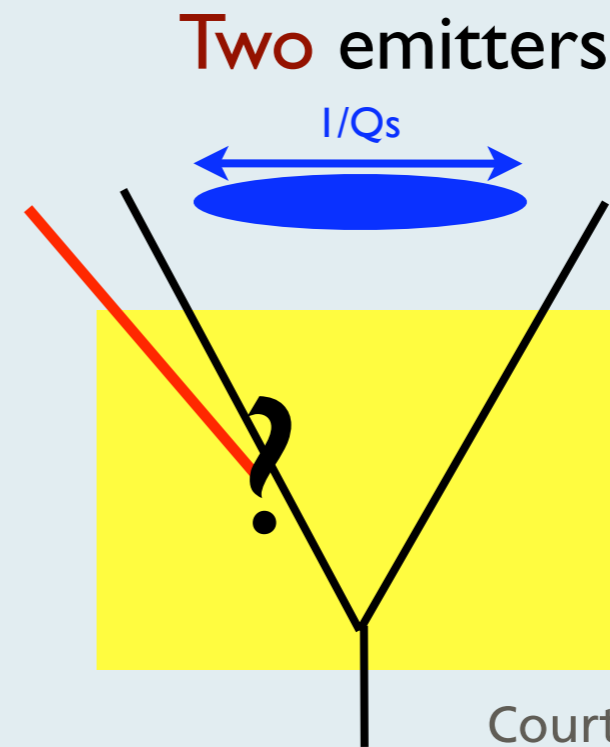
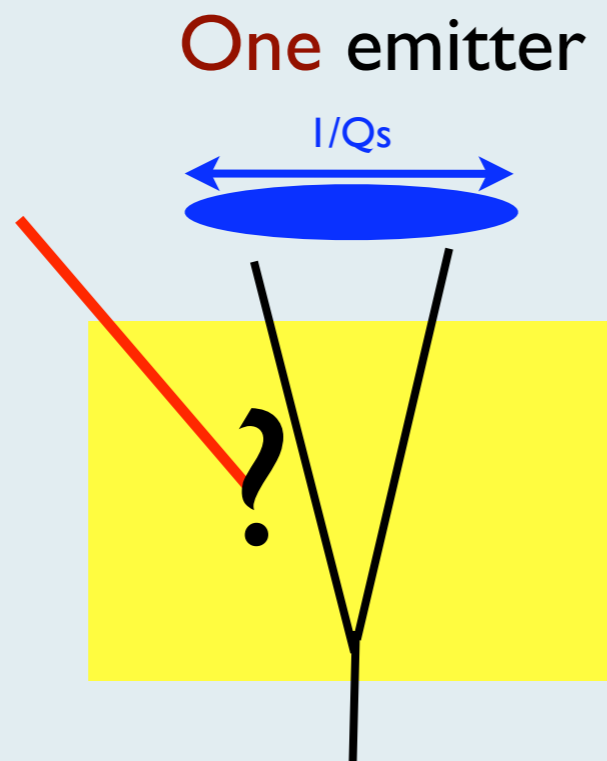
Courtesy: Y. Mehtar-Tani



- **bremstrahlung** is treated quite differently in the various codes
- **separation of scales** at LHC!



Guidance from the antenna



Courtesy: J. Casalderrey-Solana

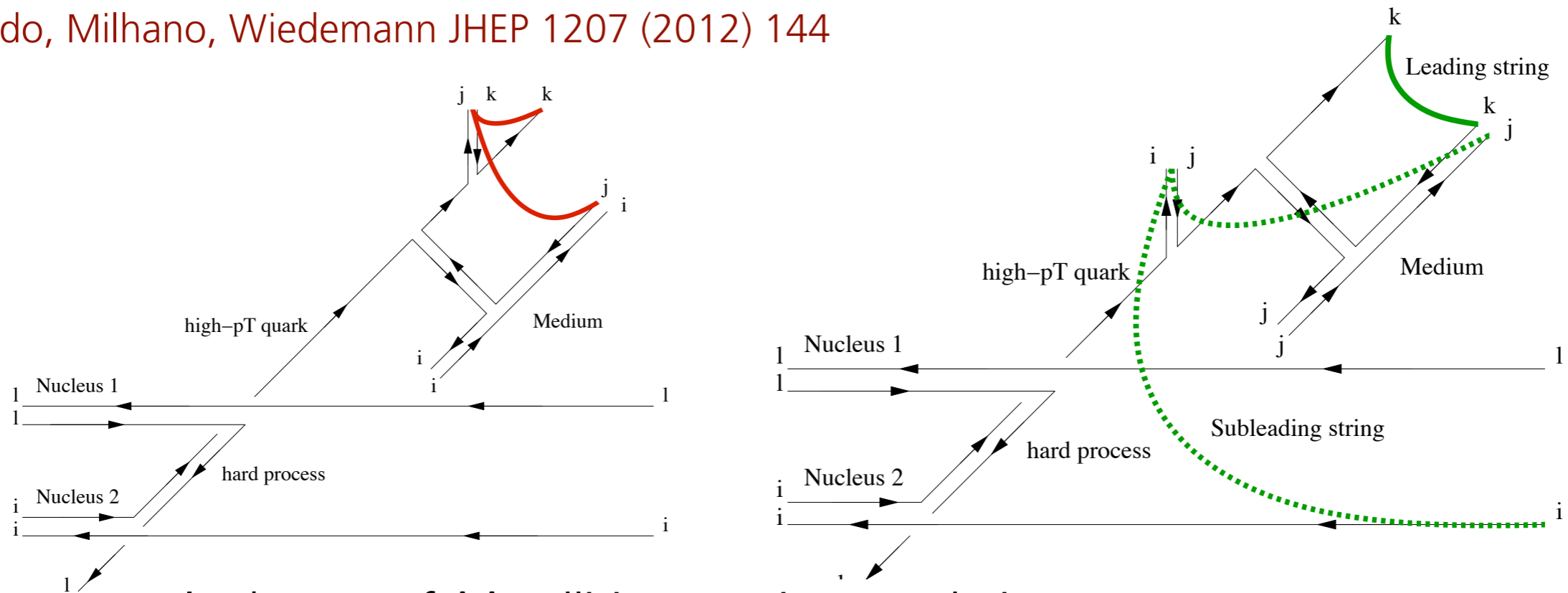
- Shower transverse size $< 1/Q_s \Rightarrow$ radiation as a single parton
- Shower transverse size $> 1/Q_s \Rightarrow$ radiation as a independent partons

Genuine pQCD effect: color transparency
[tunneling]



An important point...

Beraudo, Milhano, Wiedemann JHEP 1207 (2012) 144



In the case of AA collisions a naive convolution

Parton Energy loss \otimes Vacuum Fragmentation

without accounting for the modified color-flow would result into a too hard hadron spectrum: fitting the experimental amount of quenching would require an **overestimate of the energy loss at the partonic level**;

Andrea Beraudo, Hard Probes 2012



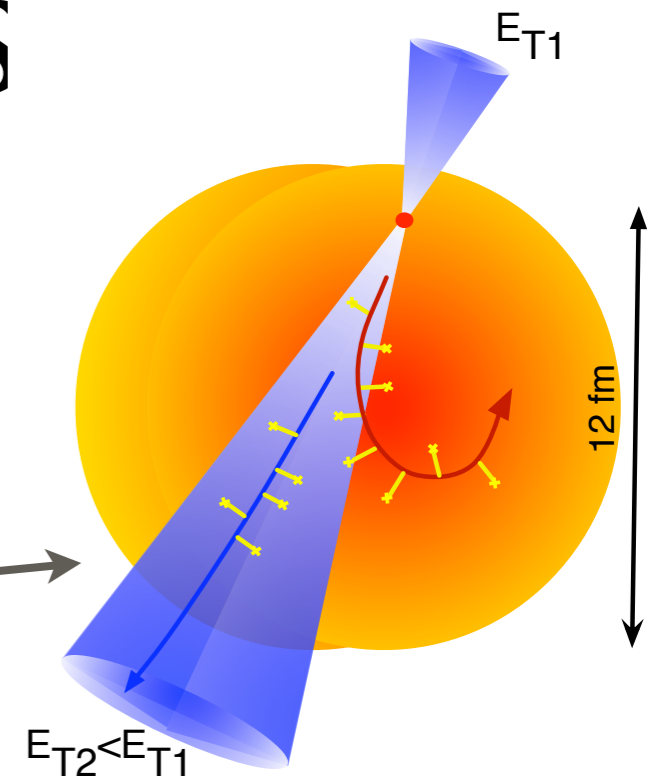
MC prospects

- versatile tools vs. detailed descriptions
- exact kinematics
 - conservation of energy-momentum
- track all particles
- implement 'advanced' space-time picture
 - evolution of energy-density, flow fields
- must be checked against well-controlled limits!
- extensions beyond theory
 - recoil/back-reaction: source terms



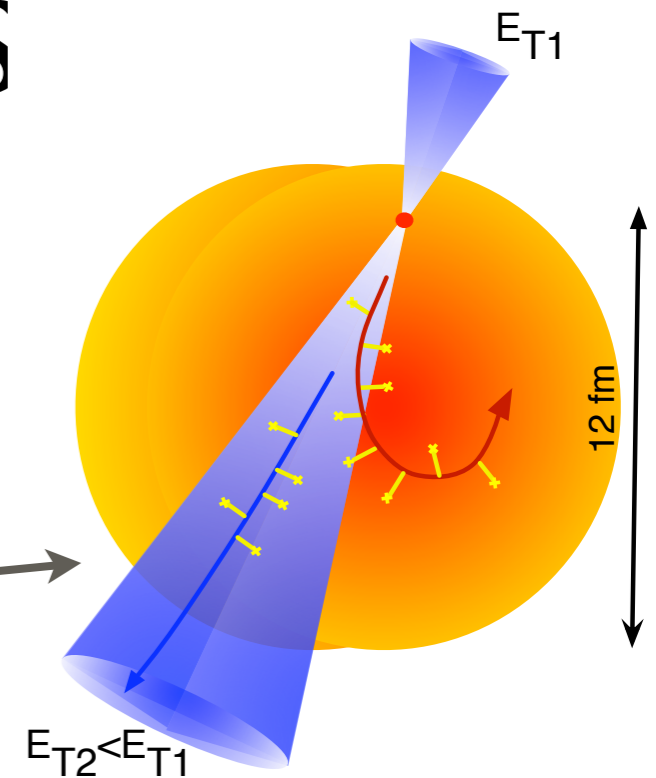
MC prospects

- versatile tools vs. detailed descriptions
- exact kinematics
- conservation of energy-momentum
- track all particles
- implement 'advanced' space-time picture
- evolution of energy-density, flow fields
- must be checked against well-controlled limits!
- extensions beyond theory
- recoil/back-reaction: source terms



MC prospects

- versatile tools vs. detailed descriptions
- exact kinematics
- conservation of energy-momentum
- track all particles
- implement 'advanced' space-time picture
- evolution of energy-density, flow fields
- must be checked against well-controlled limits!
- extensions beyond theory
- recoil/back-reaction: source terms



the “truth” is out there...



Outlook

- generic features: energy loss & softening
- how robust are the experimental signals to:
 - collisional (drag), radiative, collimation (broadening), NLO, non-perturbative (hadronization).....
 - how to get a handle?
- do we need to rethink approach to the problem?

