

# Black holes, TeV-scale gravity and the LHC

Elizabeth Winstanley

Consortium for Fundamental Physics  
School of Mathematics and Statistics, University of Sheffield



The  
University  
Of  
Sheffield.

Thanks to STFC and EU-COST network MP0905 for financial support

# Outline

- 1 Black holes in classical and semi-classical gravity
  - Black holes in classical gravity
  - Black holes in semi-classical gravity
- 2 Large extra dimensions
- 3 Mini black hole production and decay
  - Balding phase
  - Spin-down and Schwarzschild phases
- 4 Experimental searches
- 5 Quantum black holes

# Black holes in classical and semi-classical gravity



# Black holes in four-dimensional general relativity

The simplest black hole: vacuum, static and spherically symmetric

## Schwarzschild black hole

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + r^2 d\Omega_2^2$$

Metric becomes singular at

- **Curvature** singularity  $r \rightarrow 0$
- **Co-ordinate** singularity  
 $r = r_H = 2M$
- **Schwarzschild radius**  $r_H = 2M$
- Black hole has **mass**  $M$



# Four-dimensional rotating black holes

Vacuum, axisymmetric

## Kerr black hole

$$ds^2 = -\frac{\Delta}{\Sigma} [dt - a \sin^2 \theta d\phi]^2 + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2 + \frac{\sin^2 \theta}{\Sigma} [(r^2 + a^2) d\phi - a dt]^2$$

$$\Delta = r^2 - 2Mr + a^2$$

$$\Sigma = r^2 + a^2 \cos^2 \theta$$

- Event horizon at  $\Delta = 0$

$$r_H = M + \sqrt{M^2 - a^2}$$

- Event horizon rotates with angular velocity

$$\Omega_H = \frac{a}{r_H^2 + a^2}$$

- Black hole has **mass**  $M$  and **angular momentum**  $J = aM$

## Higher-dimensional black holes $d = 4 + n$

Myers-Perry black hole [ Myers and Perry, *Annals Phys.* **172**, 304 (1986) ]

$$\begin{aligned}
 ds^2 = & \left(1 - \frac{\mu}{\Sigma r^{n-1}}\right) dt^2 + \frac{2a\mu \sin^2 \theta}{\Sigma r^{n-1}} dt d\varphi - \frac{\Sigma}{\Delta_n} dr^2 - \Sigma d\theta^2 \\
 & - \left(r^2 + a^2 + \frac{a^2 \mu \sin^2 \theta}{\Sigma r^{n-1}}\right) \sin^2 \theta d\varphi^2 - r^2 \cos^2 \theta d\Omega_n^2
 \end{aligned}$$

where

$$\Delta_n = r^2 + a^2 - \frac{\mu}{r^{n-1}}, \quad \Sigma = r^2 + a^2 \cos^2 \theta$$

Black hole mass  $M$  and angular momentum  $J$ :

$$M = \frac{(n+2) A_{n+2} \mu}{16\pi G_{4+n}}, \quad J = \frac{2aM}{n+2}$$

# Black holes in semi-classical gravity

## Quantum field theory in curved space-time

- Keep geometry fixed and classical (but not necessarily static)
- Quantum fields propagating on this background
- Semi-classical approximation to quantum gravity
- Semi-classical Einstein equations

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi\langle T_{\mu\nu}\rangle$$

- $\langle T_{\mu\nu}\rangle$  - renormalized expectation value of the quantum stress-energy tensor

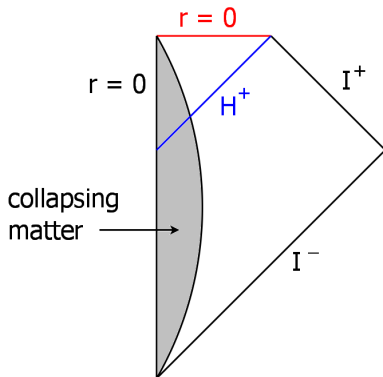
# “Black holes ain’t so black”

## Hawking radiation [ Hawking 1975 ]

- Black hole formed by gravitational collapse
- Quantum state empty of particles at  $I^-$
- A static observer at  $I^+$  sees an outgoing flux of thermal particles

## Hawking temperature

$$T = \frac{\kappa}{2\pi}$$





# Black hole evaporation

## Black hole temperature of a Schwarzschild black hole

$$T_{BH} = \frac{\kappa}{2\pi} = \frac{1}{8\pi M}$$

As the black hole radiates

- The black hole **shrinks** and its mass **decreases**
- The temperature  $T_{BH}$  increases
- Black holes have **negative specific heat**
- Rate of energy loss

$$\frac{dM}{dt} \propto \frac{1}{M^2}$$

- Lifetime of a Schwarzschild black hole of initial mass  $M_0$

$$t_{BH} \propto M_0^3$$

# Quantum fields on black hole space-times

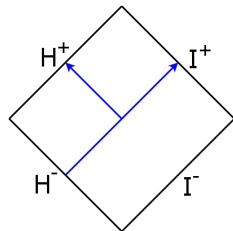
## Quantum field modes

- “Master” equation for fields of spin 0,  $\frac{1}{2}$ , 1 and 2 on Kerr  
[ Teukolsky, *Phys. Rev. Lett.* **29** 1114 (1972); *Astrophys. J.* **185** 635 (1973) ]
- Expand field  $\Psi$  in terms of modes of frequency  $\omega$ :

$$\Psi = \sum_{\omega l m} R_{s\omega l m}(r) S_{s\omega l m}(\theta) e^{-i\omega t} e^{im\varphi}$$

## “Up” modes

$$R_{s\omega l m} = \begin{cases} e^{i\tilde{\omega}r_*} + A_{\omega l m}^{\text{up}} e^{-i\tilde{\omega}r_*}, & r_* \rightarrow -\infty \\ B_{\omega l m}^{\text{up}} e^{i\omega r_*}, & r_* \rightarrow \infty \end{cases}$$



# Quantum fields on black hole space-times

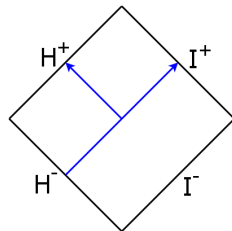
## Quantum field modes

- “Master” equation for fields of spin 0,  $\frac{1}{2}$ , 1 and 2 on Kerr [ Teukolsky, *Phys. Rev. Lett.* **29** 1114 (1972); *Astrophys. J.* **185** 635 (1973) ]
- Expand field  $\Psi$  in terms of modes of frequency  $\omega$ :

$$\Psi = \sum_{\omega lm} R_{s\omega lm}(r) S_{s\omega lm}(\theta) e^{-i\omega t} e^{im\varphi}$$

## “Up” modes

$$R_{s\omega lm} = \begin{cases} e^{i\tilde{\omega}r_*} + A_{\omega lm}^{\text{up}} e^{-i\tilde{\omega}r_*}, & r_* \rightarrow -\infty \\ B_{\omega lm}^{\text{up}} e^{i\omega r_*}, & r_* \rightarrow \infty \end{cases}$$



## Computing Hawking radiation

Differential emission rates, integrated over all angles:

$$\frac{d^2}{dt d\omega} \begin{pmatrix} N \\ E \\ J \end{pmatrix} = \frac{1}{4\pi} \sum_{\text{modes}} \frac{|\mathcal{A}_{s\omega lm}|^2}{e^{\tilde{\omega}/T_H} \mp 1} \begin{pmatrix} 1 \\ \omega \\ m \end{pmatrix}$$

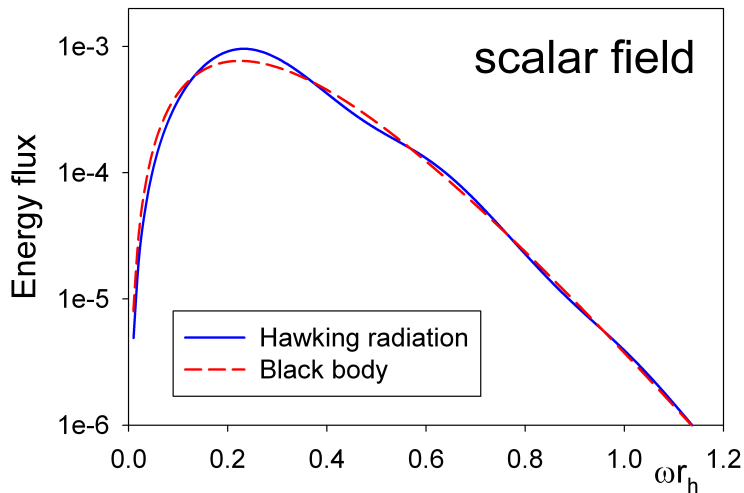
where  $\tilde{\omega} = \omega - m\Omega_H$

### Grey-body factor $|\mathcal{A}_{s\omega lm}|^2$

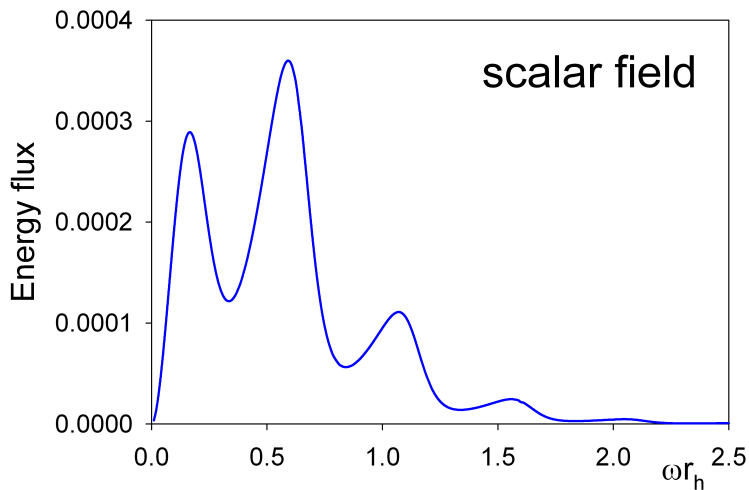
- Emitted radiation is not precisely thermal
- Interaction of emitted quanta with gravitational potential around the black hole
- For an outgoing wave from the event horizon of the black hole:

$$|\mathcal{A}_{s\omega lm}|^2 = 1 - |A_{\omega lm}^{\text{up}}|^2 = \frac{\mathcal{F}_{\text{infinity}}}{\mathcal{F}_{\text{horizon}}}$$

# Hawking emission from a Schwarzschild black hole



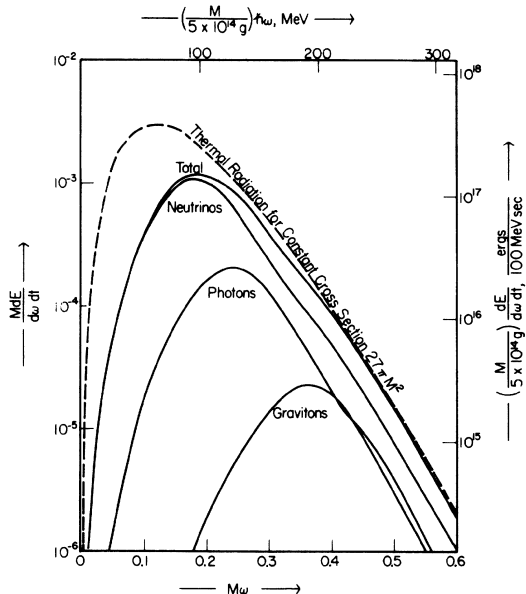
# Hawking emission from a Kerr black hole



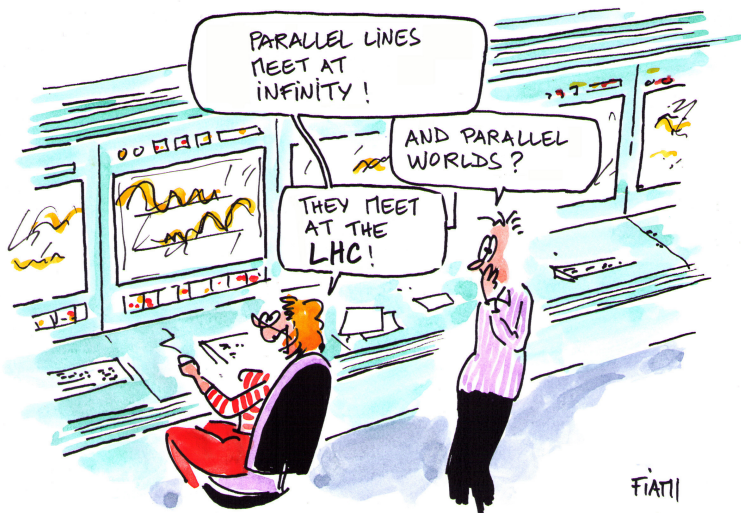
# Properties of Hawking radiation from 4D black holes

- For a Schwarzschild black hole, the emission decreases significantly as the spin of the particle increases
- Kerr black hole sheds its angular momentum very rapidly

[ Figure taken from Page, *Physical Review D* **13** 198 (1976) ]

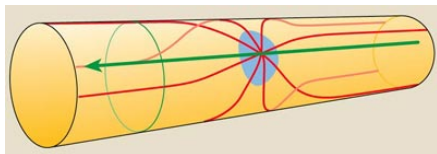


# Large extra dimensions

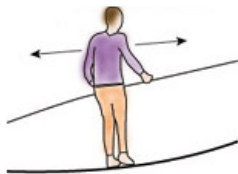




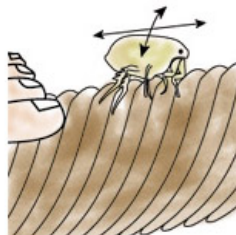
# Theories with extra dimensions



**Kaluza-Klein theory [ 1921/26 ]** Additional, compact space-like dimension smaller than any observable length-scale

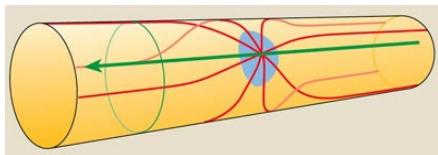


**An acrobat can only move in one dimension along a rope..**



**...but a flea can move in two dimensions.**

## Theories with extra dimensions



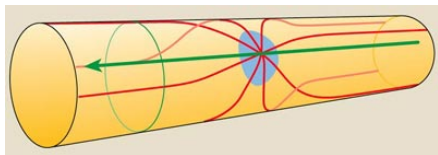
**Kaluza-Klein theory [ 1921/26 ]** Additional, compact space-like dimension smaller than any observable length-scale

**Superstring theory [ 1970's-80's ]** Unifies all forces in 10 space-time dimensions, 6 of which are compactified and roughly of the size of the Planck length  $L_P$

### Planck length

$$L_P = \sqrt{\frac{\hbar G}{c^3}} \sim 10^{-35} \text{ m}$$

# Theories with extra dimensions

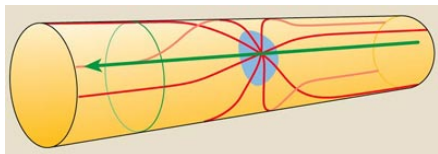


**Kaluza-Klein theory [ 1921/26 ]** Additional, compact space-like dimension smaller than any observable length-scale

**Superstring theory [ 1970's-80's ]** Unifies all forces in 10 space-time dimensions, 6 of which are compactified and roughly of the size of the Planck length  $L_P$

**M-theory [ mid 1990's ]** Five different string theories are different limits of underlying 11-dimensional theory

# Theories with extra dimensions



**Kaluza-Klein theory [ 1921/26 ]** Additional, compact space-like dimension smaller than any observable length-scale

**Superstring theory [ 1970's-80's ]** Unifies all forces in 10 space-time dimensions, 6 of which are compactified and roughly of the size of the Planck length  $L_P$

**M-theory [ mid 1990's ]** Five different string theories are different limits of underlying 11-dimensional theory

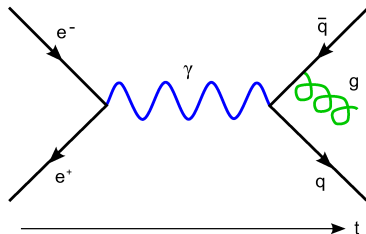
**Horava-Witten theory [ 1996 ]** Size of 11th dimension can be very much larger than  $L_P$

# The hierarchy problem

## Two very different energy scales in fundamental physics

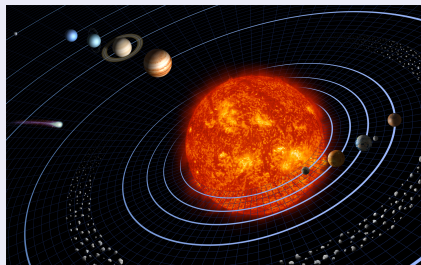
### Electroweak scale

Higgs mass  $\sim 100$  GeV



### Planck mass

$$M_P = \sqrt{\frac{\hbar c}{G}} \sim 10^{19} \text{ GeV}$$



# The hierarchy problem

Two very different energy scales in fundamental physics

## Electroweak scale

Higgs mass  $\sim 100$  GeV



## Planck mass

$$M_P = \sqrt{\frac{\hbar c}{G}} \sim 10^{19} \text{ GeV}$$



## The (A)ADD scenario

[ Antoniadis, Arkani-Hamed, Dimopoulos and Dvali, hep-ph/9803315 ;  
hep-ph/9804398 ]

Large volume for extra compact dimensions lowers the fundamental scale of quantum gravity to  $M_*$

$$M_P^2 \sim R^n M_*^{2+n}$$

$n$  - number of extra dimensions

$R$  - size of extra dimensions

If  $M_* \sim 1$  TeV

$$R \sim 10^{\frac{30}{n}-19} \text{ m}$$

For  $n = 5$ ,  $R \sim 10^{-13}$  m

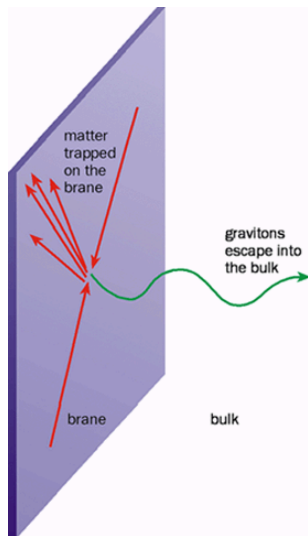


# The (A)ADD scenario

[ Antoniadis, Arkani-Hamed, Dimopoulos and Dvali, hep-ph/9803315 ;  
hep-ph/9804398 ]

To avoid a contradiction with Standard Model physics:

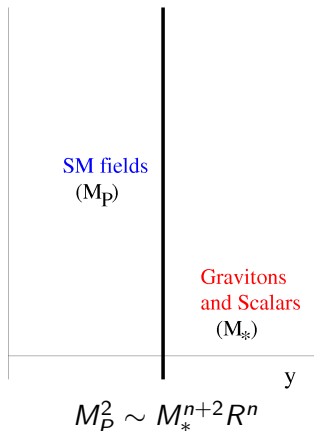
- 4D **brane** where all Standard Model particles live
- Effective scale for gravity is  $M_P$  on the brane
- $(4 + n)$  D **bulk**
- Only gravitons propagate in the bulk
- Higher-dimensional fundamental scale for gravity is  $M_*$



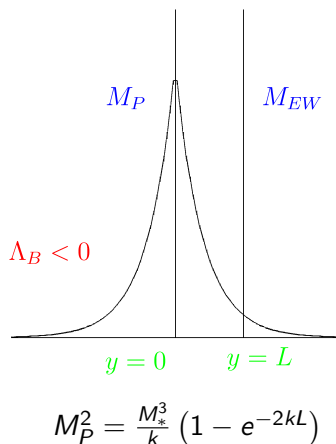


## Brane worlds

ADD model



RS I model



[ Figures taken from Kanti, arXiv:0802.2218 [hep-th] ]

# Consequences of the ADD model

## KK-modes

- A field in the higher-dimensional theory is seen by an observer on the brane as an infinite tower of massive modes with  $m \sim 1/R$
- Produced at energies  $E > 1/R$ , modifying processes on the brane
- Modification of short-range gravitational potential

$$V(r) = -\frac{G_N M_1 M_2}{r} \left(1 + \alpha e^{-r/R}\right)$$

# Consequences of the ADD model

## KK-modes

- A field in the higher-dimensional theory is seen by an observer on the brane as an infinite tower of massive modes with  $m \sim 1/R$
- Produced at energies  $E > 1/R$ , modifying processes on the brane
- Modification of short-range gravitational potential

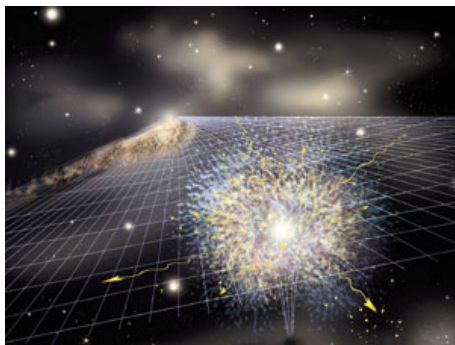
$$V(r) = -\frac{G_N M_1 M_2}{r} \left( 1 + \alpha e^{-r/R} \right)$$

## Probing quantum gravity at colliders

- Collider experiment with centre-of-mass energy  $\sqrt{s} > M_*$  will probe strong-gravity regime
- Creation of heavy extended gravitational objects

[ Banks and Fischler, [hep-th/9906038](#) ]

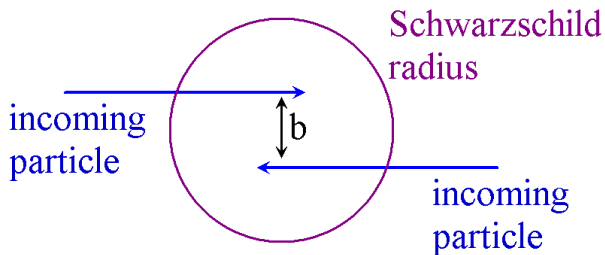
# Mini black hole production and decay



[ Image credit: Aurore Simonet]

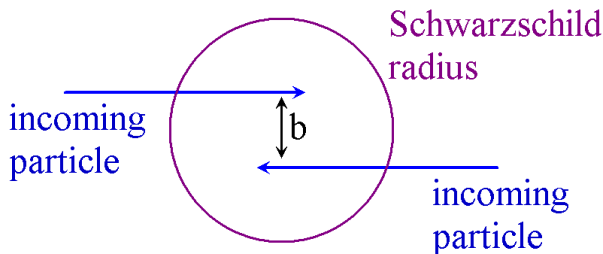
## Black hole production

Two particles with centre-of-mass energies greater than  $M_*$ , and impact parameter  $b$ :



## Black hole production

Two particles with centre-of-mass energies greater than  $M_*$ , and impact parameter  $b$ :

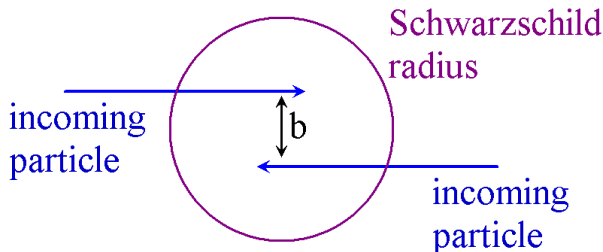


4D hoop conjecture [ Thorne, *Magic Without Magic* 231 (1972) ]

A black hole is formed when a mass  $M$  gets compacted into a region whose circumference in every direction is less than  $2\pi r_H(M)$

## Black hole production

Two particles with centre-of-mass energies greater than  $M_*$ , and impact parameter  $b$ :



Hyperhoop conjecture [ Ida and Nakao, [gr-qc/0204082](#) ]

A black hole is formed when a mass  $M$  gets compacted into a region whose  $D - 3$ -dimensional volume in every direction is less than  $\alpha_D M$

Modified hyperhoop conjecture required when extra dimensions are compactified [ Yoo et al, [arXiv:0906.0689](#) [gr-qc] ]

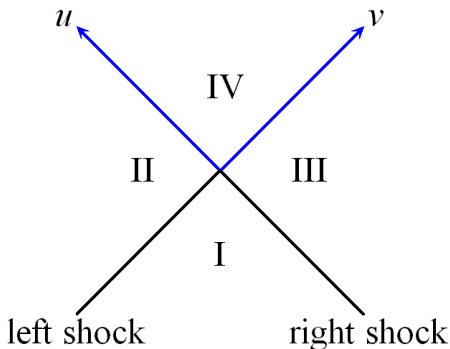
# Modelling the formation of mini black holes

- Classical process
- Model particles as infinitely-boosted black holes
- Search for apparent horizon
- Maximum impact parameter  $b_{max}$
- Parton-level production cross-section

$$\sigma \sim \pi b_{max}^2 \sim 3\pi r_h^2$$

[ Yoshino and Rychkov,  
hep-th/0503171 ]

Aichelberg-Sexl shock waves  
[ Aichelberg and Sexl, GRG 2 303  
(1971) ]





# Classical high energy collisions in numerical relativity

Model particles as colliding boson stars, fluid particles or black holes

Key questions [ Sperhake, arXiv:1301.3772 [gr-qc] ]

- Validity of the hoop conjecture
- Scattering threshold for black hole formation
- Mass and spin of the formed black holes
- Effects of internal structure of colliding objects

Results to date

- Four-dimensional collisions best studied
- Higher-dimensional work in early days  
[ Witek et al, arXiv:1006.3081 [gr-qc] ; arXiv:1011.0742 [gr-qc] ]  
  
[ Okawa et al, arXiv:1105.3331 [gr-qc] ]

# Classical high energy collisions in numerical relativity

Model particles as colliding boson stars, fluid particles or black holes

Key questions [ Sperhake, arXiv:1301.3772 [gr-qc] ]

- Validity of the hoop conjecture
- Scattering threshold for black hole formation
- Mass and spin of the formed black holes
- Effects of internal structure of colliding objects

## Results to date

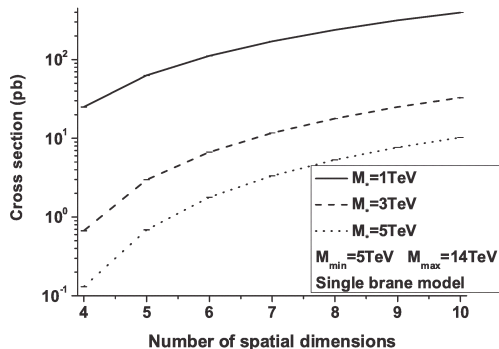
- Four-dimensional collisions best studied
- Higher-dimensional work in early days  
[ Witek et al, arXiv:1006.3081 [gr-qc] ; arXiv:1011.0742 [gr-qc] ]  
  
[ Okawa et al, arXiv:1105.3331 [gr-qc] ]

# Black hole production cross-sections

## Parton-level BH production cross-section

$$\sigma_{ij \rightarrow \text{BH production}} \propto \pi r_H^2 \sim \frac{1}{M_*^2} \left( \frac{E}{M_*} \right)^{\frac{2}{n+1}}$$

[ Giddings and Thomas, hep-ph/0106219 ;  
Dimopoulos and Landsberg, hep-ph/0106295 ]



For  $M_* = 1 \text{ TeV}$  and  
 $M_{BH} = 5 \text{ TeV}$  with  
 $n = 6$ ,  $\sigma \sim 10^2 \text{ pb}$  :  
about one black hole per  
second!

[ Figure taken from  
Dai et al,  
arXiv:0711.3012 ]

# Stages in the life of a mini black hole

[ Giddings and Thomas, hep-ph/0106219 ]

**Balding phase** Shedding of hair and asymmetries

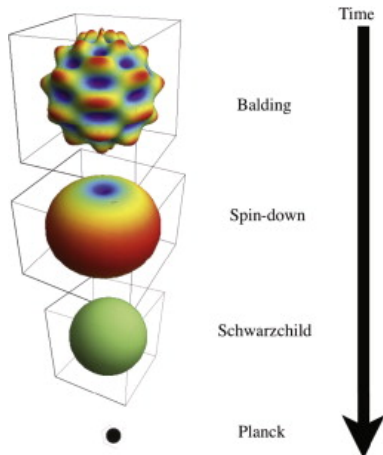
**Spin-down phase** Loss of angular momentum via Hawking radiation

**Schwarzschild phase** Loss of mass via Hawking radiation

**Planck phase**  $M_{BH} \sim M_*$

Image credit:

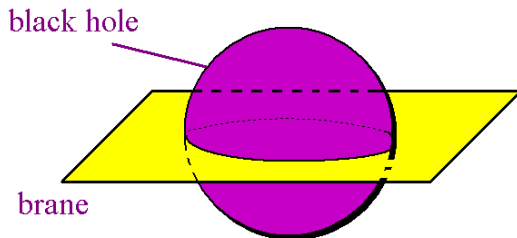
[ Park, arXiv:1203.4683 [hep-ph] ]



# Modelling mini black holes at the end of the balding stage

## Small black holes in ADD

- Metric of higher-dimensional black holes in general relativity is known [ Myers and Perry, *Annals Phys.* **172**, 304 (1986) ]
- Take a 'slice' through a higher-dimensional black hole to give a brane black hole



# Modelling mini black holes in ADD

## Slice of Myers-Perry black hole

$$ds^2 = \left(1 - \frac{\mu}{\Sigma r^{n-1}}\right) dt^2 + \frac{2a\mu \sin^2 \theta}{\Sigma r^{n-1}} dt d\varphi - \frac{\Sigma}{\Delta_n} dr^2 - \Sigma d\theta^2 \\ - \left(r^2 + a^2 + \frac{a^2 \mu \sin^2 \theta}{\Sigma r^{n-1}}\right) \sin^2 \theta d\varphi^2$$

where

$$\Delta_n = r^2 + a^2 - \frac{\mu}{r^{n-1}}, \quad \Sigma = r^2 + a^2 \cos^2 \theta$$

and  $n$  is the number of extra dimensions.

## Usual Kerr black hole

Set  $n = 0$  in the above metric

## Balding phase

Shedding of mass and angular momentum through gravitational radiation modeled as part of formation process

### Limits on amount of energy lost in gravitational radiation

- Colliding shock waves:  $\leq 30\%$  ( $n = 0$ ),  $\leq 40\%$  ( $n = 7$ )  
[ Yoshino and Rychkov, hep-th/0503171 ]
- Four-dimensional numerical relativity:  $\leq 50\%$  ( $n = 0$ )  
[ Sperhake et al, arXiv:1211.6114 [gr-qc] ]

### Angular momentum of formed black hole

- Angular momentum of black holes with  $n > 1$  potentially unbounded
- Limited by maximum impact parameter
- Colliding shock waves:  $j \sim 0.93$  ( $n = 1$ )  
[ Yoshino and Rychkov, hep-th/0503171 ]

## Balding phase

Shedding of mass and angular momentum through gravitational radiation modeled as part of formation process

### Limits on amount of energy lost in gravitational radiation

- Colliding shock waves:  $\leq 30\%$  ( $n = 0$ ),  $\leq 40\%$  ( $n = 7$ )  
[ Yoshino and Rychkov, hep-th/0503171 ]
- Four-dimensional numerical relativity:  $\leq 50\%$  ( $n = 0$ )  
[ Sperhake et al, arXiv:1211.6114 [gr-qc] ]

### Angular momentum of formed black hole

- Angular momentum of black holes with  $n > 1$  potentially unbounded
- Limited by maximum impact parameter
- Colliding shock waves:  $j \sim 0.93$  ( $n = 1$ )  
[ Yoshino and Rychkov, hep-th/0503171 ]



## Balding phase

Very little work done on shedding charges or gauge field hair

### QCD effects

Likely to be significant, but little work on this

[ Calmet et al, arXiv:0806.4605 [hep-ph] ]

[ Gingrich, arXiv:0912.0826 [hep-ph] ]

### Electromagnetic effects

- Effect of charge on formation process  
[ Zilhao et al, arXiv:1205.1063 [gr-qc] ]
- Classical Maxwell field on the brane only - modifies the “slice” of the Myers-Perry black hole
- Loss of black hole charge is not rapid in TeV gravity models

[ Sampaio, arXiv:0907.5107 [hep-th] ]

## Balding phase

Very little work done on shedding charges or gauge field hair

### QCD effects

Likely to be significant, but little work on this

[ Calmet et al, arXiv:0806.4605 [hep-ph] ]

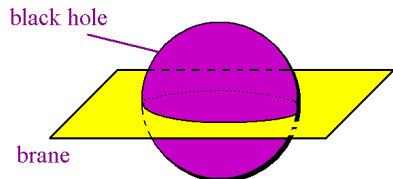
[ Gingrich, arXiv:0912.0826 [hep-ph] ]

### Electromagnetic effects

- Effect of charge on formation process  
[ Zilhao et al, arXiv:1205.1063 [gr-qc] ]
- Classical Maxwell field on the brane only - modifies the “slice” of the Myers-Perry black hole
- Loss of black hole charge is not rapid in TeV gravity models

[ Sampaio, arXiv:0907.5107 [hep-th] ]

# Hawking radiation on the brane and in the bulk



Hawking temperature

$$T_H = \frac{(n+1)r_h^2 + (n-1)a^2}{4\pi(r_h^2 + a^2)r_h}$$

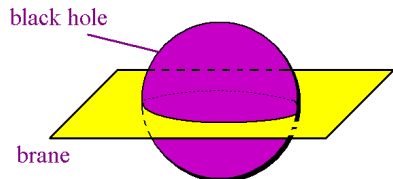
## Particles on the brane

- Standard model particles: fermions, gauge bosons, Higgs
- Also gravitons and scalars
- Live on the brane “slice” of the black hole geometry

## Particles in the bulk

- Gravitons and scalars
- Will be invisible
- Live on the higher-dimensional black hole geometry

# Hawking radiation on the brane and in the bulk



Hawking temperature

$$T_H = \frac{(n+1)r_h^2 + (n-1)a^2}{4\pi(r_h^2 + a^2)r_h}$$

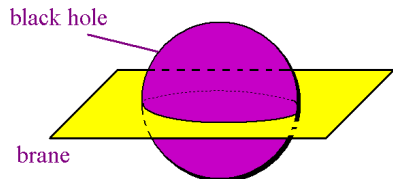
## Particles on the brane

- Standard model particles: fermions, gauge bosons, Higgs
- Also gravitons and scalars
- Live on the brane “slice” of the black hole geometry

## Particles in the bulk

- Gravitons and scalars
- Will be invisible
- Live on the higher-dimensional black hole geometry

# Hawking radiation on the brane and in the bulk



## Particles on the brane

- Standard model particles: fermions, gauge bosons, Higgs
- Also gravitons and scalars
- Live on the brane “slice” of the black hole geometry

## Hawking temperature

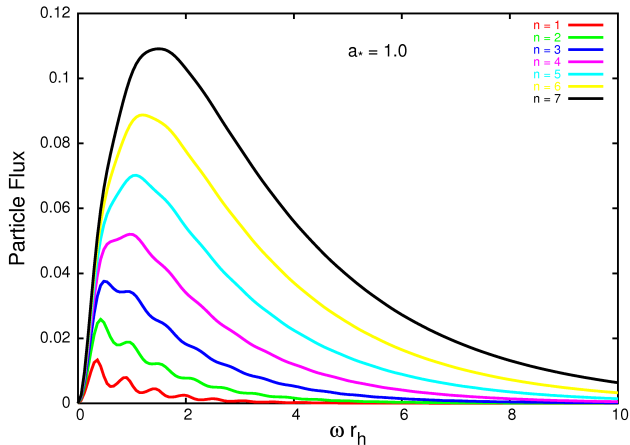
$$T_H = \frac{(n+1)r_h^2 + (n-1)a^2}{4\pi(r_h^2 + a^2)r_h}$$

## Particles in the bulk

- Gravitons and scalars
- Will be invisible
- Live on the higher-dimensional black hole geometry

## Emission spectra

Fermion emission spectra for a rotating black hole, integrated over all angles

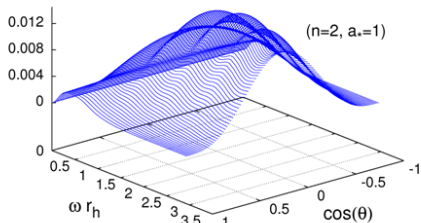


[ Figure taken from Casals et al, hep-th/0608193 ]

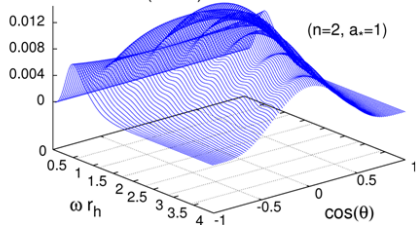
# Angular distribution of energy flux

Six-dimensional  
black hole  
 $n=2$

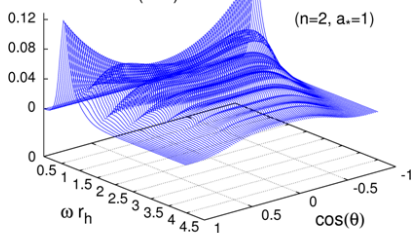
Power Flux ( $s=0$ )



Power Flux ( $s=1/2$ )



Power Flux ( $s=1$ )



# What we know about the Hawking radiation phases

## “Spin-down” phase

- Brane emission - scalars, fermions, gauge bosons done
- Bulk emission - scalars done
- Graviton emission - partial results only

## “Schwarzschild” phase

- Brane emission - scalars, fermions, gauge bosons done
- Bulk emission - scalars done
- Graviton emission - bulk and brane done

## “Black holes radiate mainly on the brane”

[ Emparan, Horowitz and Myers, hep-th/0003118 ]

Ratio of bulk/brane emission for massless scalars,  $n = 2$

$a_* = 0.0$	$a_* = 0.2$	$a_* = 0.4$	$a_* = 0.6$	$a_* = 0.8$	$a_* = 1.0$
19.9%	18.6%	15.3%	11.7%	9.0%	7.1%

[ Casals et al, arXiv:0801.4910 [hep-th] ]



# What we know about the Hawking radiation phases

## “Spin-down” phase

- Brane emission - scalars, fermions, gauge bosons done
- Bulk emission - scalars done
- Graviton emission - partial results only

## “Schwarzschild” phase

- Brane emission - scalars, fermions, gauge bosons done
- Bulk emission - scalars done
- Graviton emission - bulk and brane done

## “Black holes radiate mainly on the brane”

[ Emparan, Horowitz and Myers, hep-th/0003118 ]

Ratio of bulk/brane emission for massless scalars,  $n = 2$

$a_* = 0.0$	$a_* = 0.2$	$a_* = 0.4$	$a_* = 0.6$	$a_* = 0.8$	$a_* = 1.0$
19.9%	18.6%	15.3%	11.7%	9.0%	7.1%

[ Casals et al, arXiv:0801.4910 [hep-th] ]

# What we know about the Hawking radiation phases

## “Spin-down” phase

- Brane emission - scalars, fermions, gauge bosons done
- Bulk emission - scalars done
- Graviton emission - partial results only

## “Schwarzschild” phase

- Brane emission - scalars, fermions, gauge bosons done
- Bulk emission - scalars done
- Graviton emission - bulk and brane done

## “Black holes radiate mainly on the brane”

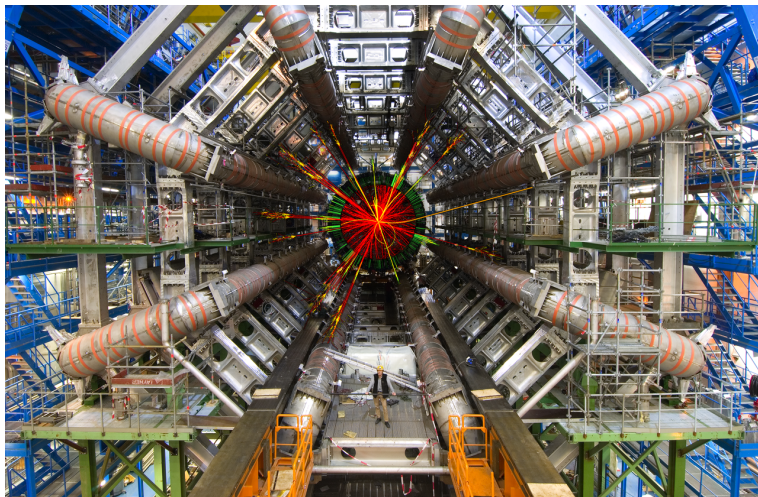
[ Emparan, Horowitz and Myers, hep-th/0003118 ]

Ratio of bulk/brane emission for massless scalars,  $n = 2$

$a_* = 0.0$	$a_* = 0.2$	$a_* = 0.4$	$a_* = 0.6$	$a_* = 0.8$	$a_* = 1.0$
19.9%	18.6%	15.3%	11.7%	9.0%	7.1%

[ Casals et al, arXiv:0801.4910 [hep-th] ]

# Experimental searches



[ Image credit: ATLAS experiment ©2012 CERN]

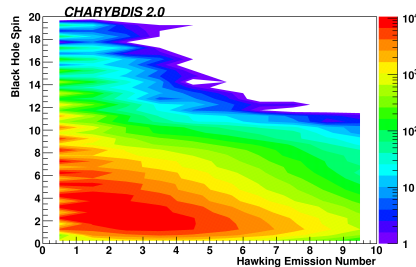
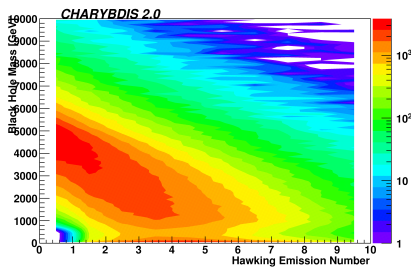
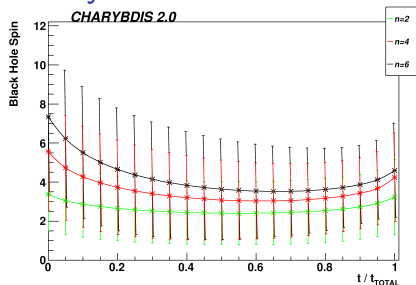
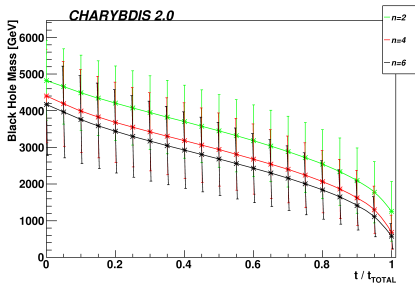
# Simulations of black hole events at the LHC

To discover black holes at the LHC, accurate event simulations are required

## Black hole event generators [ Gingrich, hep-ph/0610219 ]

- TrueNoir [ Landsberg, hep-ph/0607297 ]  
<http://hep.brown.edu/users/Greg/TrueNoir/>
- CATFISH [ Cavaglia et al, hep-ph/0609001 ]  
<http://www.phy.olemiss.edu/GR/catfish/introduction.html>
- BlackMax [ Dai et al, arXiv:0902.3577 [hep-ph] ]  
<http://projects.hepforge.org/blackmax/>
- CHARYBDIS2 [ Frost et al, arXiv:0904.0979 [hep-ph] ]  
<http://projects.hepforge.org/charybdis2/>
- QBH [ Gingrich, arXiv:0911.5370 [hep-ph] ]  
<http://projects.hepforge.org/qbh/>

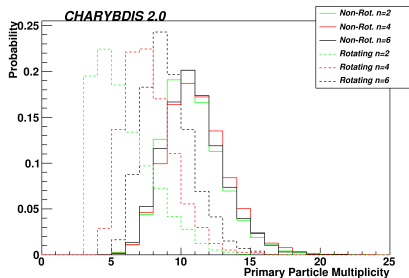
# Evolution of black holes simulated by CHARYBDIS2



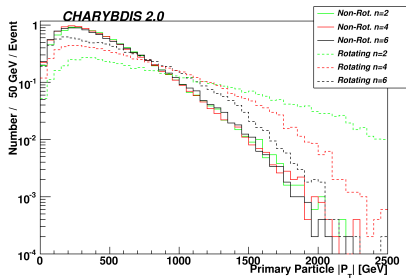
[ Figures taken from Frost et al, arXiv:0904.0979 [hep-ph] ]

# Results from CHARYBDIS2

## Primary particle multiplicity

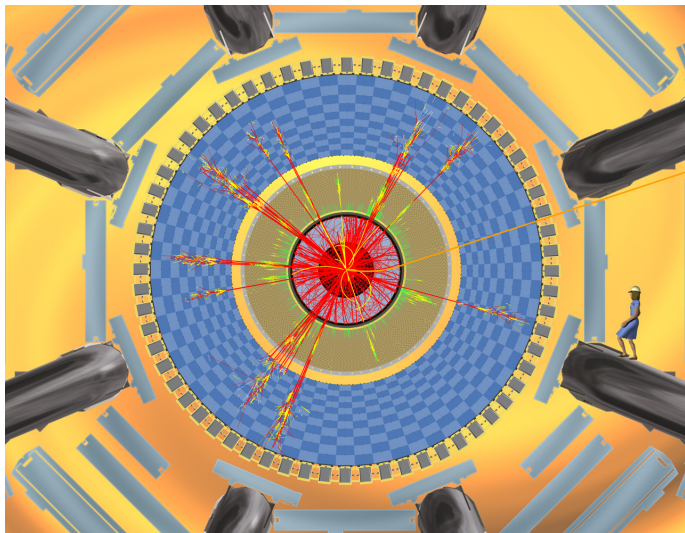


## Primary particle $P_T$



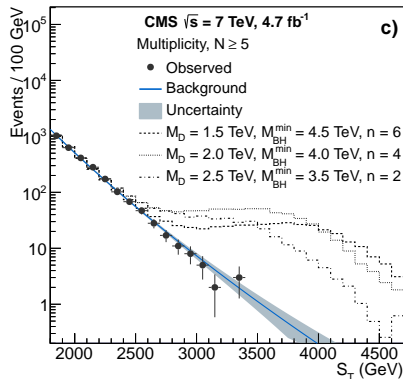
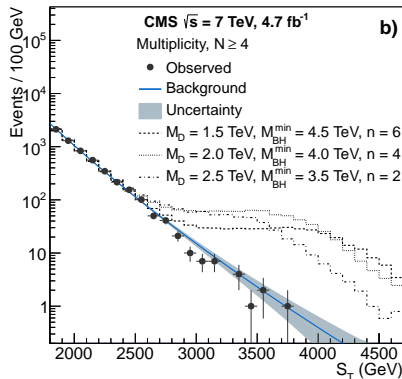
[ Figures taken from Frost et al, arXiv:0904.0979 [hep-ph] ]

# An example black hole event



[ Image credit: ATLAS experiment ©2012 CERN]

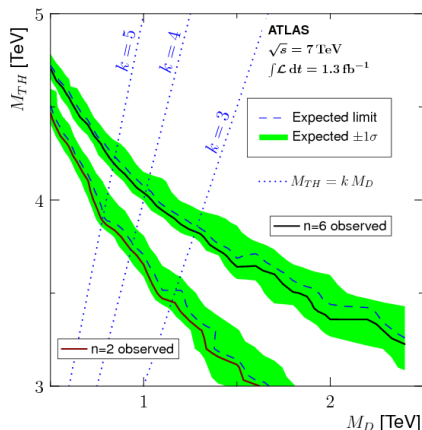
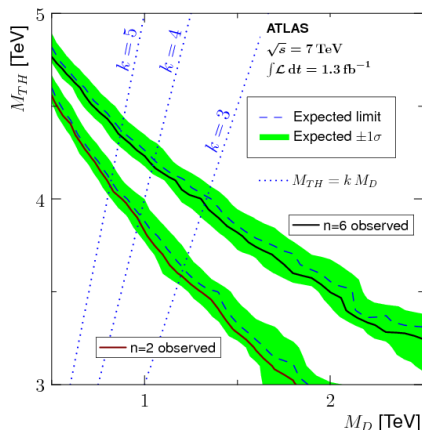
## CMS search for black hole events



[ CMS collaboration, arXiv:1202.6396 [hep-ex] ]

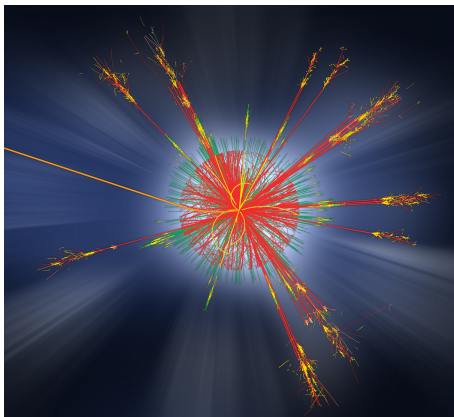


## ATLAS search for black hole events



[ ATLAS collaboration, arXiv:1111.0080 [hep-ex] ]

# Quantum black holes



[ Image credit: ATLAS experiment ©2012 CERN]

# Drawbacks of the “standard” black hole model

## Black hole formation

### Classical approximation

- Colliding particles described by general relativity
- Difficult to include quantum fields, particularly QCD

## Black hole evolution

### Semi-classical approximation

- Black hole geometry described by general relativity
- Quantum fields on this background
- Semi-classical approximation breaks down for  $M_{BH} \sim M_*$
- Details of quantum gravity unknown

## Beyond the semi-classical approximation

### Validity of semi-classical approximation

Compton wavelength of colliding particle of energy  $E/2$  must lie within the Schwarzschild radius:

$$4\pi/E < r_h(E)$$

Therefore  $E/M_* \gtrsim 10$  in order for semi-classical black holes to form  
 [ Meade and Randall, arXiv:0708.3017 [hep-ph] ]

### Planckian quantum black holes

Masses close to  $M_*$

- Do not have thermal decay
- Particle physics symmetries used to constrain decay products

[ Calmet et al, arXiv:0806.4605 [hep-ph] ]

[ Gingrich, arXiv:0912.0826 [hep-ph] ]

## Beyond the semi-classical approximation

### Validity of semi-classical approximation

Compton wavelength of colliding particle of energy  $E/2$  must lie within the Schwarzschild radius:

$$4\pi/E < r_h(E)$$

Therefore  $E/M_* \gtrsim 10$  in order for semi-classical black holes to form  
 [ Meade and Randall, arXiv:0708.3017 [hep-ph] ]

### Planckian quantum black holes

Masses close to  $M_*$

- Do not have thermal decay
- Particle physics symmetries used to constrain decay products

[ Calmet et al, arXiv:0806.4605 [hep-ph] ]

[ Gingrich, arXiv:0912.0826 [hep-ph] ]

# Conclusions

- Large-extra-dimension scenarios
- Scale of quantum gravity  $M_*$  could be as low as a few TeV
- Possibility of making microscopic black holes at the LHC
- Semi-classical model - decay by Hawking radiation
- Non-observation sets bound on  $M_*$

