

C-parameter distribution at N³LL and a determination of α_s

... a story of 3 π 's and 6's

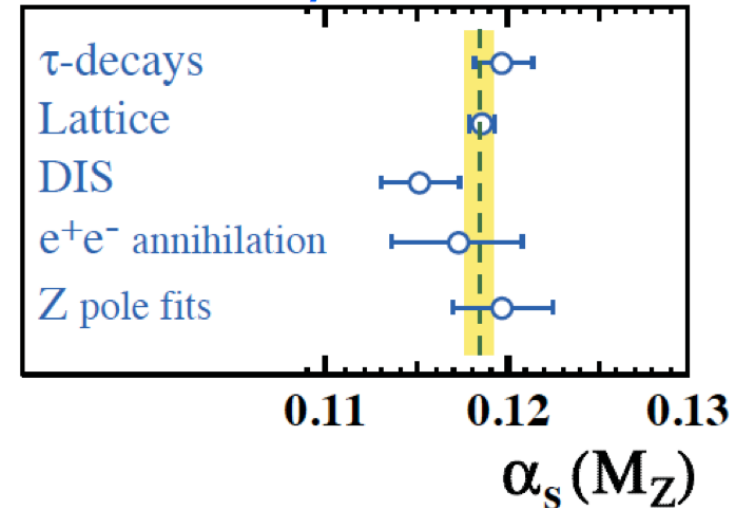
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work done with Andre Hoang, Vicent Mateu and Iain Stewart

Outline

- Introduction
- C-parameter factorization
 - Singular with resummation
 - Nonsingular cross section
 - Nonperturbative corrections
 - Profile Functions
- Analysis and Results
- Conclusions

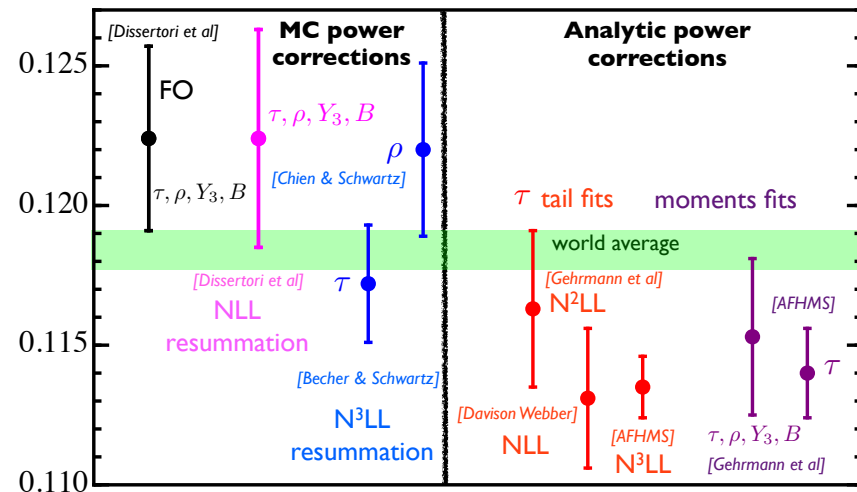
World average completely dominated by lattice result



α_s determination: compendium

Only consider analysis with 3-loop input

$\alpha_s(m_Z)$ determination from event shape fits



Definition of C-parameter

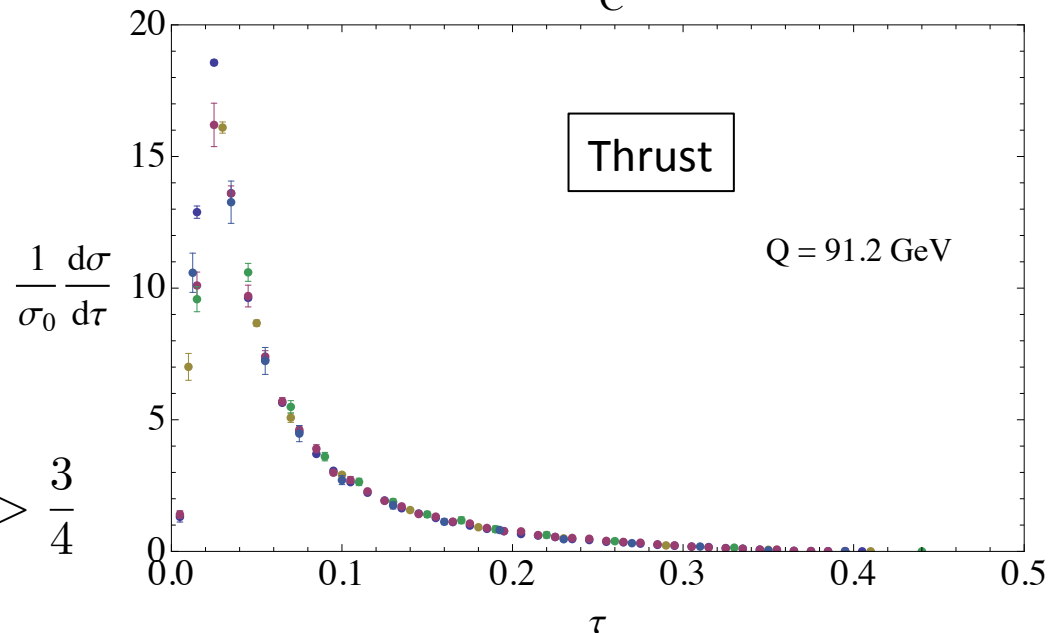
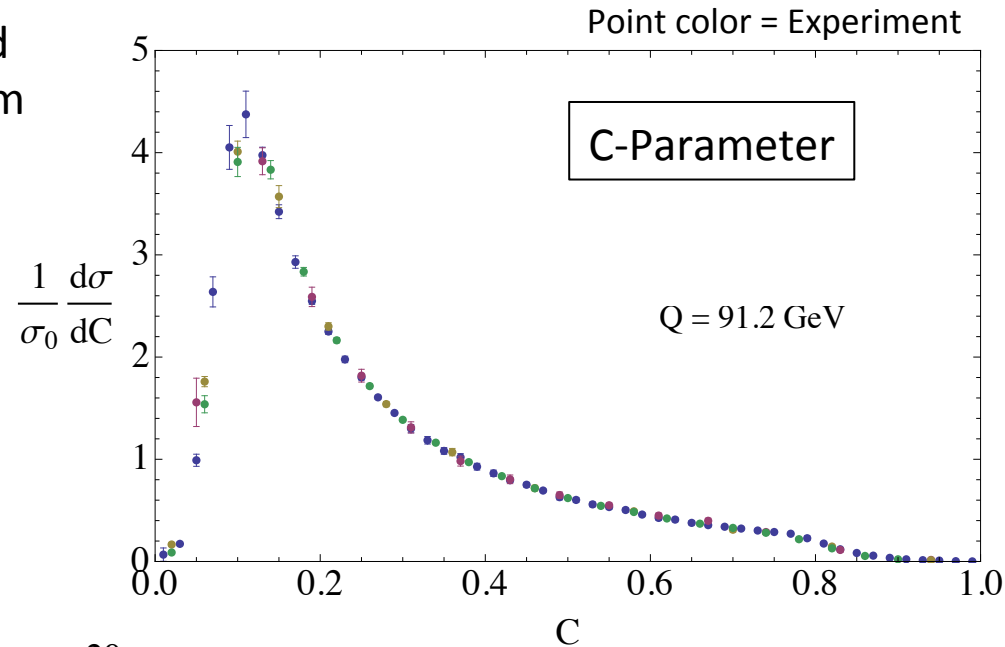
$$\theta_{\alpha\beta} = \frac{\sum_i \mathbf{p}_i^\alpha \mathbf{p}_i^\beta / |\mathbf{p}_i|}{\sum_j |\mathbf{p}_j|} \quad \text{Linearized momentum tensor}$$

- Eigenvalues of θ : $\lambda_1, \lambda_2, \lambda_3$
- Definition of C-parameter:

$$C = 3(\lambda_1\lambda_2 + \lambda_2\lambda_3 + \lambda_3\lambda_1)$$

$$C = \frac{3 \sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{2 (\sum_i |\vec{p}_i|)^2}$$

- Dijet limit: $C \rightarrow 0$
- Multijet region: $C \rightarrow 1$
- $C > \frac{3}{4}$ only for $\mathcal{O}(\alpha_s^2)$
 - Only non planar events have $C > \frac{3}{4}$



Definition of C-parameter

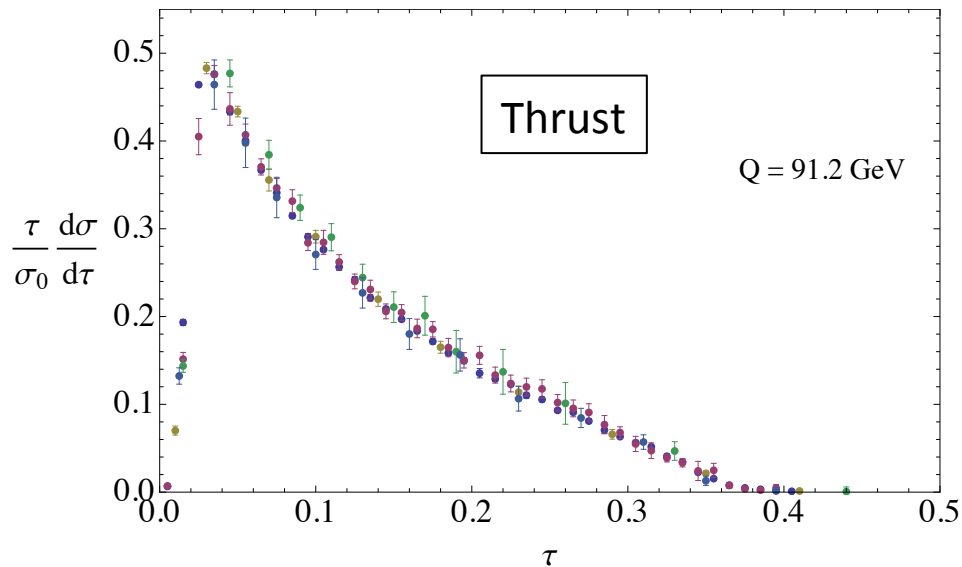
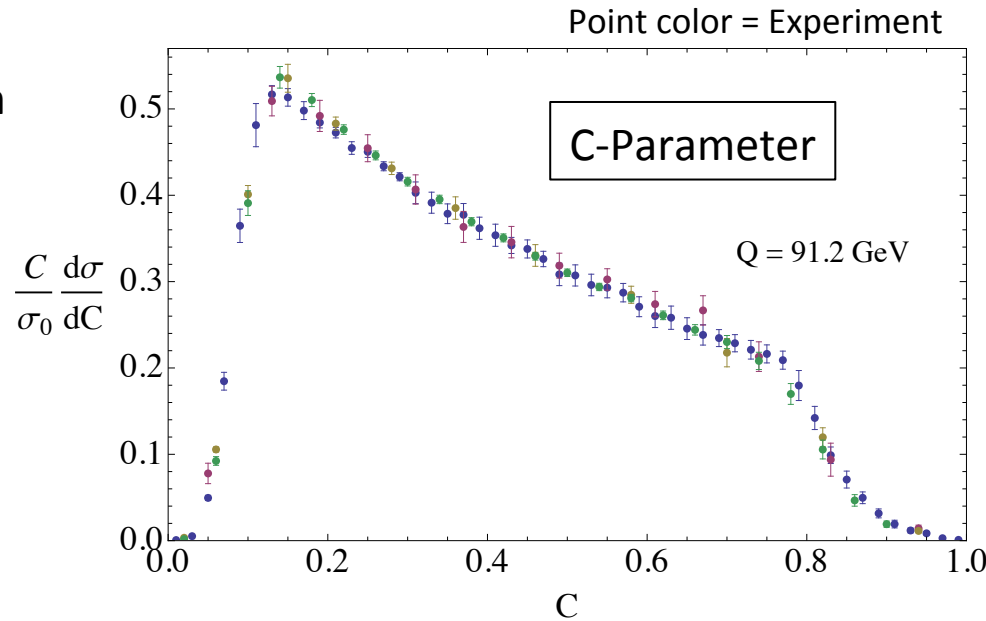
$$\theta_{\alpha\beta} = \frac{\sum_i \mathbf{p}_i^\alpha \mathbf{p}_i^\beta / |\mathbf{p}_i|}{\sum_j |\mathbf{p}_j|} \quad \begin{array}{l} \text{Linearized} \\ \text{momentum} \\ \text{tensor} \end{array}$$

- Eigenvalues of θ : λ_1 , λ_2 , λ_3
- Definition of C-parameter:
$$C = 3(\lambda_1\lambda_2 + \lambda_2\lambda_3 + \lambda_3\lambda_1)$$

$$C = \frac{3}{2} \frac{\sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$$

Contrast with thrust

- Thrust peaks at a lower value
- Thrust peak is much higher
- Thrust falls off faster in tail
- No minimization
- No thrust axis



Regions of Distribution

Peak

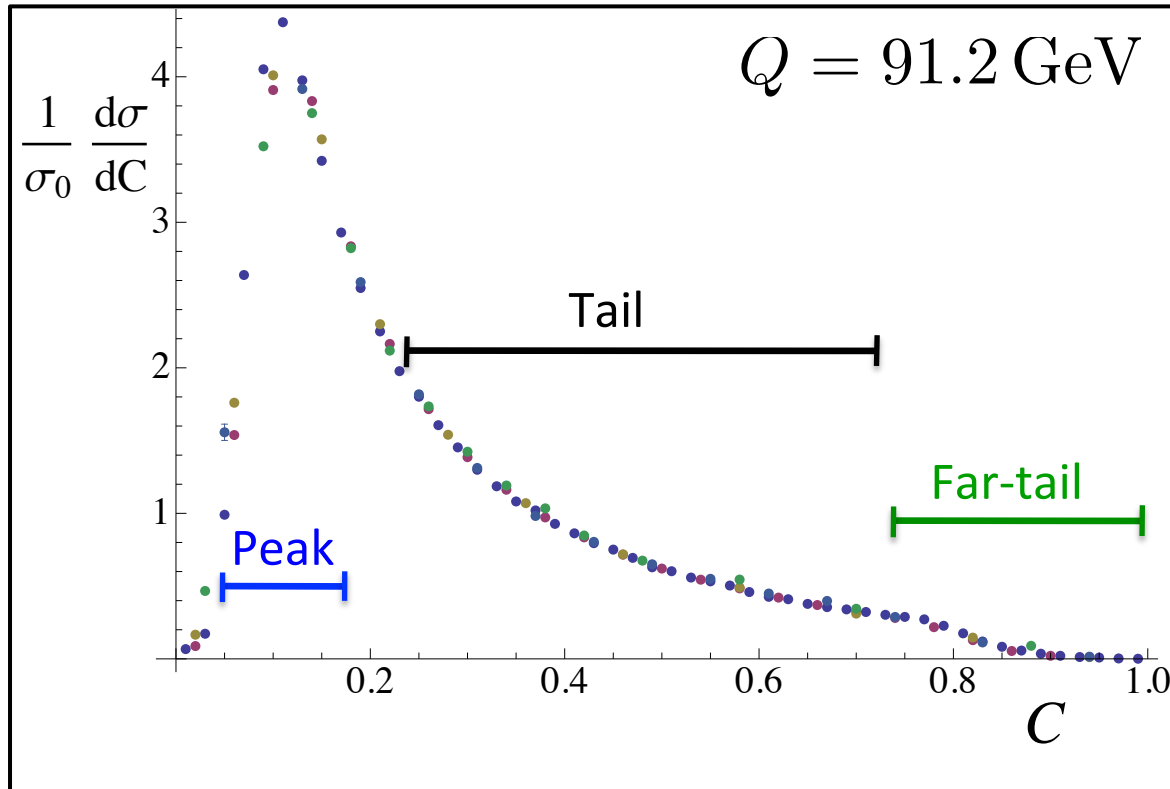
$$C \sim 3\pi \frac{\Lambda_{\text{QCD}}}{Q}$$

Tail

$$3\pi \frac{\Lambda_{\text{QCD}}}{Q} \ll C \lesssim \frac{3}{4}$$

Far-tail

$$\frac{3}{4} \lesssim C \leq 1$$

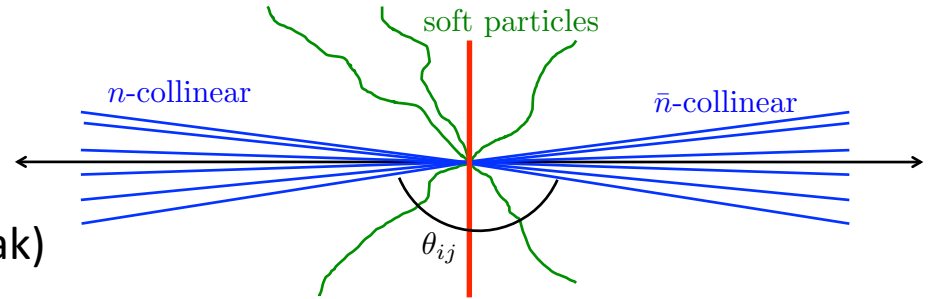


C-parameter Kinematics

$$C = \frac{3}{2} \frac{\sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$$

- In dijet limit C is small, expand in

$$\lambda \approx \sqrt{C} \quad (\text{tail}) \quad \lambda \approx \sqrt{\frac{\Lambda_{\text{QCD}}}{Q}} \quad (\text{peak})$$



- Split sum into pieces for various types of partons, **keep only leading order terms**

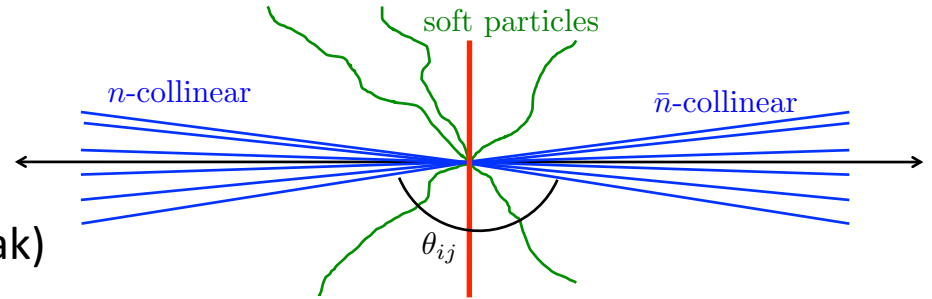
$$C_{\text{dijet}} = C_{ss} + C_{ns} + C_{\bar{n}s} + C_{nn} + C_{\bar{n}\bar{n}} + C_{n\bar{n}}$$

C-parameter Kinematics

$$C = \frac{3}{2} \frac{\sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$$

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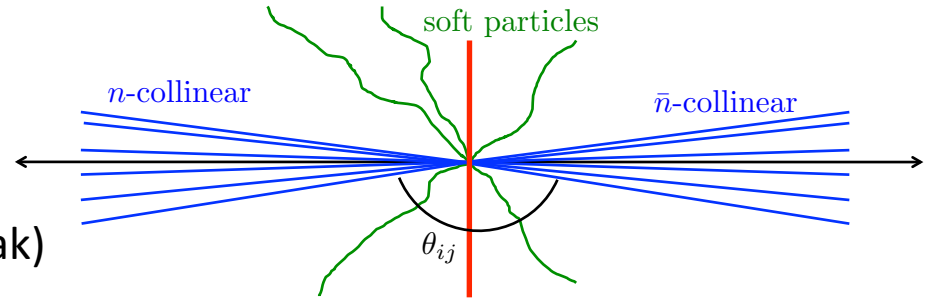
$\mathcal{O}(\lambda^4)$

C-parameter Kinematics

$$C = \frac{3}{2} \frac{\sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$$

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$$C_{\text{dijet}} = \cancel{C_{ss}} + C_{ns} + C_{\bar{n}s} + C_{nn} + C_{\bar{n}\bar{n}} + C_{n\bar{n}}$$

Using $\sin \theta_{sn} \approx \sin \theta_s$:

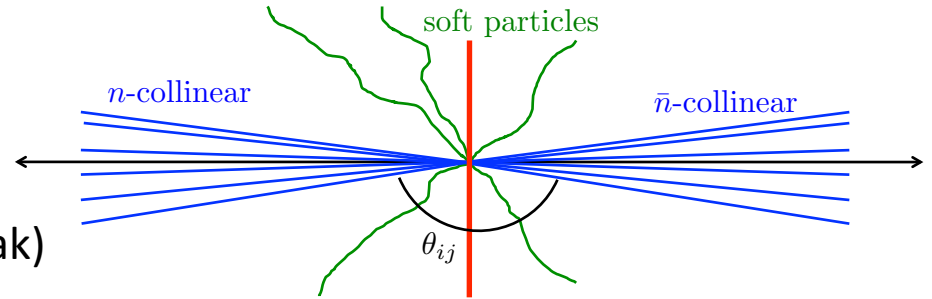
$$C_{ns} = C_{\bar{n}s} = \frac{3}{Q} \sum_{i \in \text{soft}} \frac{p_i^+ p_i^-}{p_i^+ + p_i^-}$$

C-parameter Kinematics

$$C = \frac{3}{2} \frac{\sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$$

- In dijet limit C is small, expand in

$$\lambda \approx \sqrt{C} \quad (\text{tail}) \quad \lambda \approx \sqrt{\frac{\Lambda_{\text{QCD}}}{Q}} \quad (\text{peak})$$



- Split sum into pieces for various types of partons, **keep only leading order terms**

$$C_{\text{dijet}} = \cancel{C_{ss}} + C_{ns} + C_{\bar{n}s} + C_{nn} + C_{\bar{n}\bar{n}} + C_{n\bar{n}}$$

Using momentum conservation and $\cos \theta_{nn} \approx 1$:

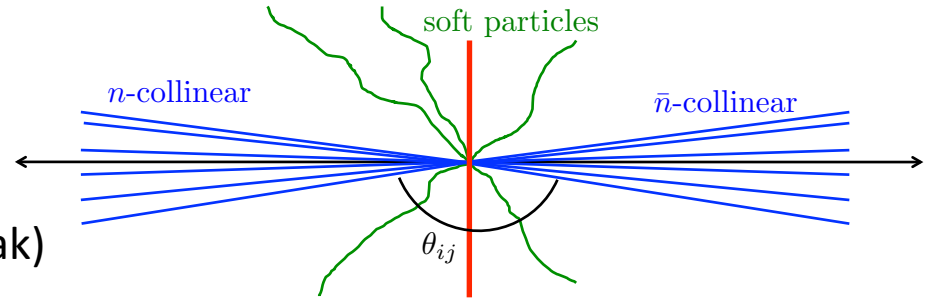
$$C_{nn} = \frac{3}{Q} \sum_{i \in n} p_i^+, \quad C_{\bar{n}\bar{n}} = \frac{3}{Q} \sum_{i \in \bar{n}} p_i^-$$

C-parameter Kinematics

$$C = \frac{3}{2} \frac{\sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$$

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$$C_{\text{dijet}} = \cancel{C_{ss}} + C_{ns} + C_{\bar{n}s} + C_{nn} + C_{\bar{n}\bar{n}} + C_{n\bar{n}}$$

Using energy conservation, momentum conservation and $\cos \theta_{n\bar{n}} \approx -1$:

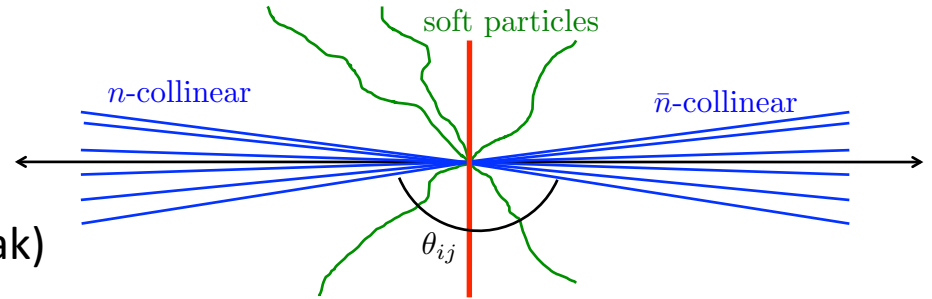
$$C_{n\bar{n}} = \frac{3}{Q} \sum_{i \in n} p_i^+ + \frac{3}{Q} \sum_{i \in \bar{n}} p_i^- = C_{nn} + C_{\bar{n}\bar{n}}$$

C-parameter Kinematics

$$C = \frac{3}{2} \frac{\sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$$

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$$C_s = \frac{6}{Q} \sum_{i \in \text{soft}} \frac{p_i^+ p_i^-}{p_i^+ + p_i^-}, \quad C_n = \frac{6}{Q} \sum_{i \in n} p_i^+, \quad C_{\bar{n}} = \frac{6}{Q} \sum_{i \in \bar{n}} p_i^-$$

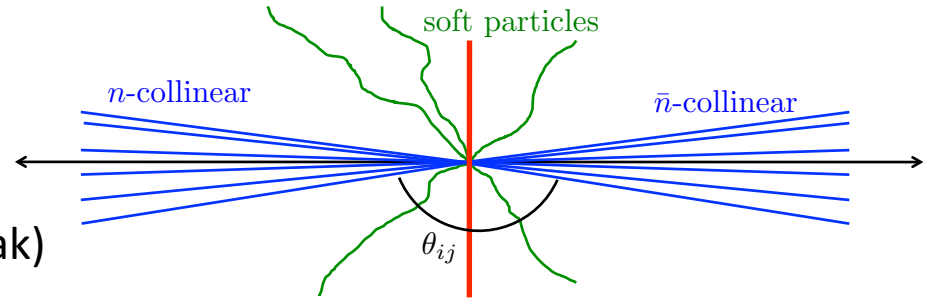
$$\Rightarrow C_{\text{dijet}} = C_s + C_n + C_{\bar{n}} = \frac{3}{Q} \sum_i \frac{p_i^\perp}{\cosh \eta_i}$$

C-parameter Kinematics

$$C = \frac{3}{2} \frac{\sum_{ij} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$$

- In dijet limit C is small, expand in

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$$\Rightarrow C_{\text{dijet}} = C_s + C_n + C_{\bar{n}} = \frac{3}{Q} \sum_i \frac{p_i^\perp}{\cosh \eta_i}$$

For the purposes of logs in factorization theorem (dijet)

$$(\text{C-parameter}) \quad \tilde{C} \equiv \frac{C}{6} \approx \tau \quad (\text{Thrust})$$

Factorization

$$\frac{1}{\sigma_0} \frac{d\sigma}{dC} = \frac{Q}{6} H(Q, \mu) \int ds J_\tau(s, \mu) S_{\tilde{C}} \left(\frac{QC}{6} - \frac{s}{Q}, \mu \right)$$

Catani, Weber:
NLL

Korchemsky,
Sterman, Tafat:
NLL+Shape
Function

Bauer, Fleming,
Lee, Sterman:
Factorization

- Both **Jet** and **Hard** functions are the same as in thrust
- Soft function differs, given by:

$$S_{\tilde{C}}(\ell, \mu) = \langle 0 | \overline{Y}_{\bar{n}}^\dagger Y_n^\dagger \delta \left(\ell - \frac{Q\hat{C}}{6} \right) Y_n \overline{Y}_{\bar{n}} | 0 \rangle$$

- Anomalous dimension for soft function same as thrust!

$$\frac{d\sigma^{\text{NLL}}}{dC} = \frac{1}{6} \frac{d\sigma^{\text{NLL}}}{d\tau} \Bigg|_{\tau = \frac{C}{6}}$$

- Relations between large logs to all orders in α_s
(which are nontrivial for NNLL and beyond)

Soft Function

$$S_{\tilde{C}}(\ell, \mu) = \langle 0 | \bar{Y}_{\bar{n}}^\dagger Y_n^\dagger \delta \left(\ell - \frac{Q\hat{C}}{6} \right) Y_n \bar{Y}_{\bar{n}} | 0 \rangle$$

- In Fourier space

$$S_{\tilde{C}}(y, \mu) = e^{K(y, \mu)} e^{S_0(\alpha_s(\mu_y))}$$

$$\mu_y^{-1} = ie^{\gamma_E} y$$

- $K(y, \mu)$ contains logs and is the same as thrust

- S_0 is C-parameter dependent piece

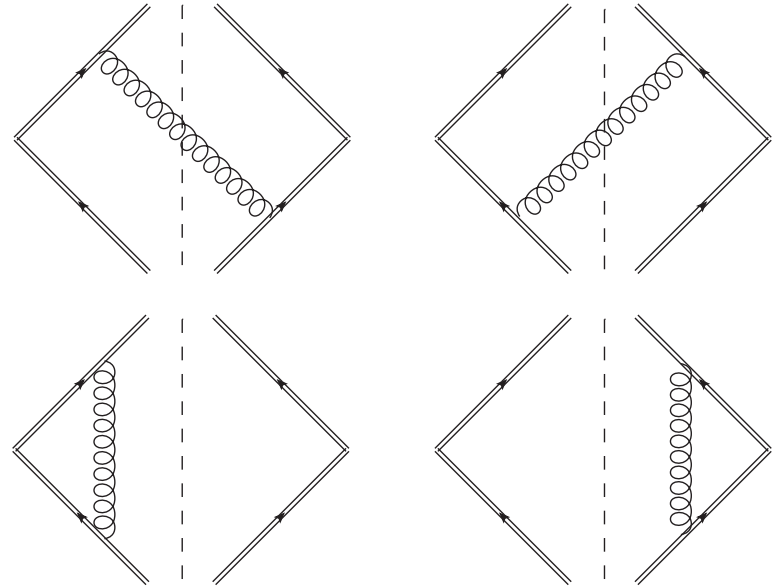
$$S_0(\alpha_s) = \frac{\alpha_s}{4\pi} s_1 + 2 \left(\frac{\alpha_s}{4\pi} \right)^2 s_2 + \dots$$

$$s_1 = -\frac{C_F \pi^2}{3}$$

$$s_1^T = -C_F \pi^2$$

Thrust Result

$$s_2 = -43 \pm 1 \quad (\text{From later fit to EVENT2})$$



Diagrams to calculate one loop soft function

Full Cross Section

From renormalon free definition of first moment

$$\frac{d\sigma}{dC} = \int dk \left(\frac{d\hat{\sigma}_s}{dC} + \frac{d\hat{\sigma}_{ns}}{dC} \right) \left[C - \frac{6k}{Q} \right] F_{\tilde{C}} \left(k - \frac{3\pi}{6} \bar{\Delta}(R, \mu_S) \right)$$

- **Singular** piece calculated from factorization theorem

$$\frac{d\hat{\sigma}_s}{dC} \sim \frac{\alpha_s^k \ln^\ell \left(\frac{C}{6} \right)}{C} \quad \text{At N}^3\text{LL}$$

- **Nonsingular** piece calculated numerically to $\mathcal{O}(\alpha_s^3)$

$$\frac{d\hat{\sigma}_{ns}}{dC} = \frac{\alpha_s}{2\pi} f_1 + \left(\frac{\alpha_s}{2\pi} \right)^2 f_2 + \left(\frac{\alpha_s}{2\pi} \right)^3 f_3 + \dots$$

- **Shape function** gives nonperturbative corrections

In tail region

$$\frac{1}{6} F \left(\frac{\ell}{6} \right) = \delta(\ell) - 3\pi \Omega_1^{gC} \delta'(\ell) + \dots$$

[Adopting a notation that makes is easier to compare with previous talk by Vicent]

Resummation

$$\frac{1}{\sigma_0} \frac{d\sigma_s}{dC} = \frac{Q}{6} H(Q, \mu_H) U_H(\mu_H, \mu_J) \int ds' ds J_\tau(s', \mu_J) U_S(s - s', \mu_J, \mu_S) S_{\tilde{C}} \left(\frac{QC}{6} - \frac{s}{Q}, \mu_S \right)$$

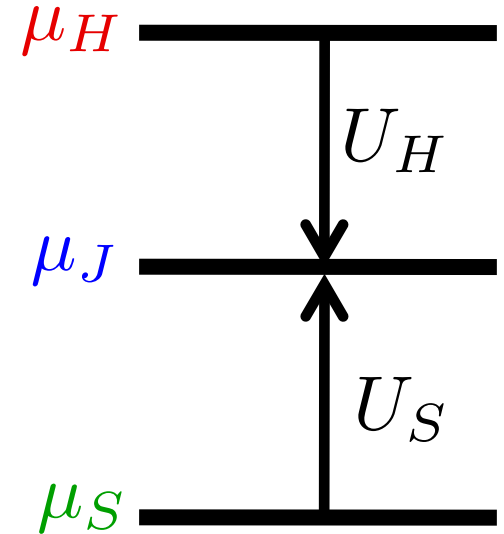
- $H(Q, \mu_H)$, $J_\tau(s, \mu_J)$, $S_{\tilde{C}}(\ell, \mu_S)$ will have no more large logs
- Running U_H and U_S resum the large logs
- To get N³LL need 3-loop anomalous dimension and 4-loop Γ_{cusp}

Padé Approximation \pm 200%

- N³LL + $\mathcal{O}(\alpha_s^3)$ needs 3-loop matching for

$$H(Q, \mu_H), J_\tau(s, \mu_J), S_{\tilde{C}}(\ell, \mu_S)$$

- j_3 , s_3 varied as part of uncertainty

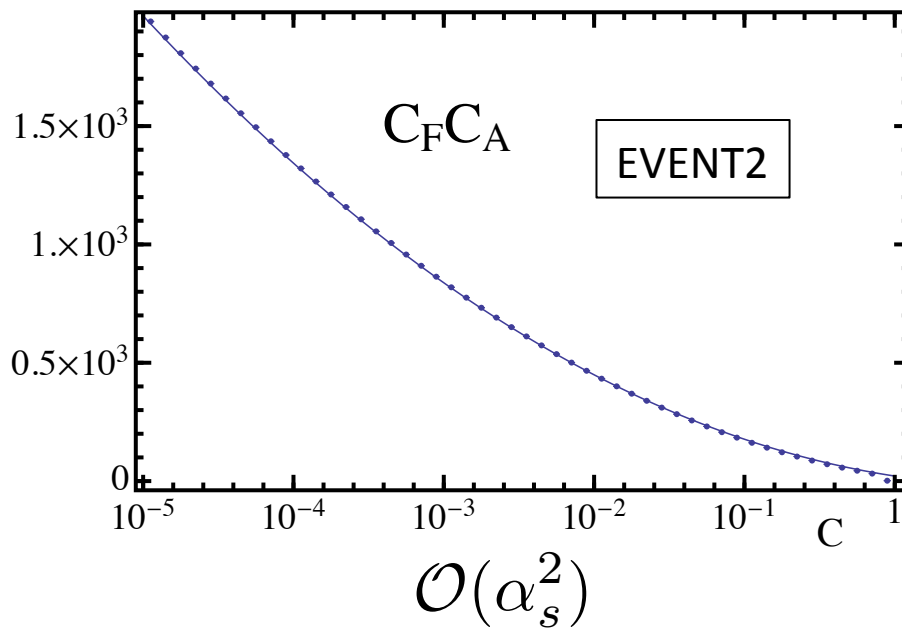


Nonsingular Distributions

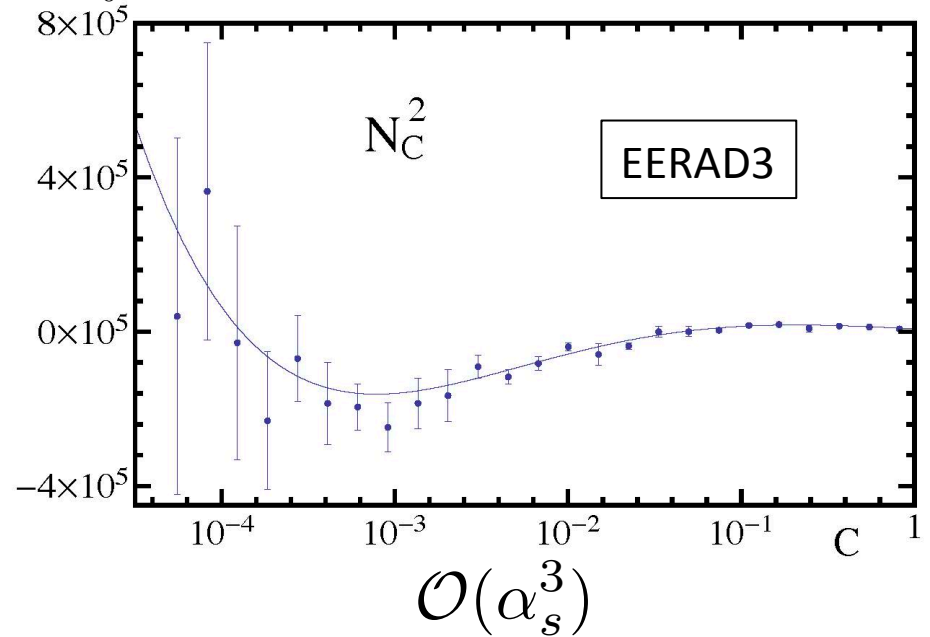
$$\sigma_{\text{nonsingular}} = \sigma_{\text{full}}^{\text{Fixed Order}} - \sigma_{\text{singular}}^{\text{Fixed Order}}$$

First compare the EVENT2 and EERAD3 results with singular cross section very close to zero

$$\frac{1}{\sigma_0} \frac{d\sigma^{(2)}}{dC}$$



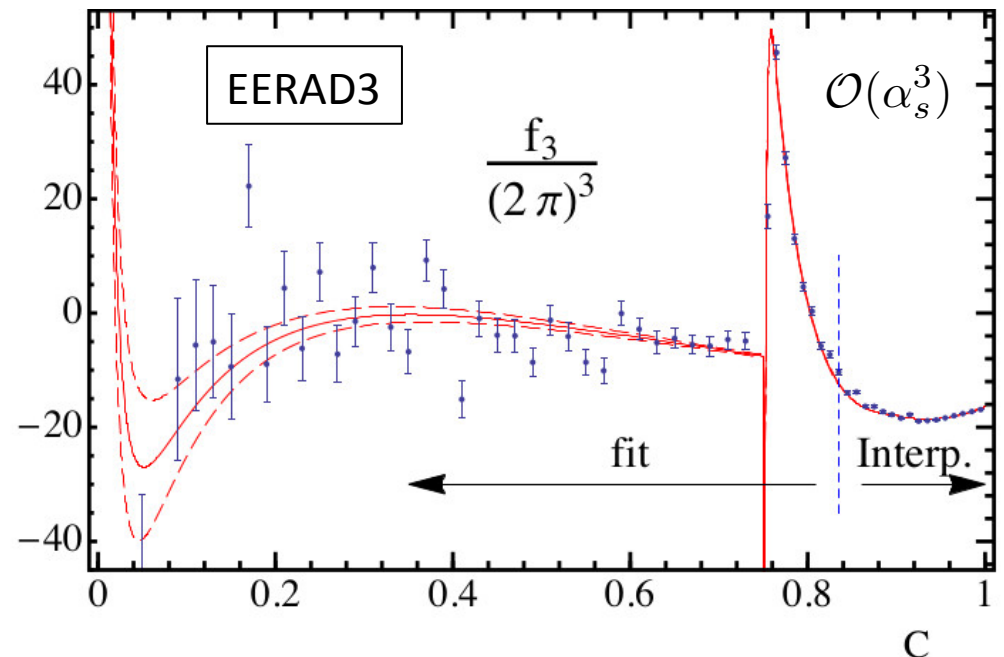
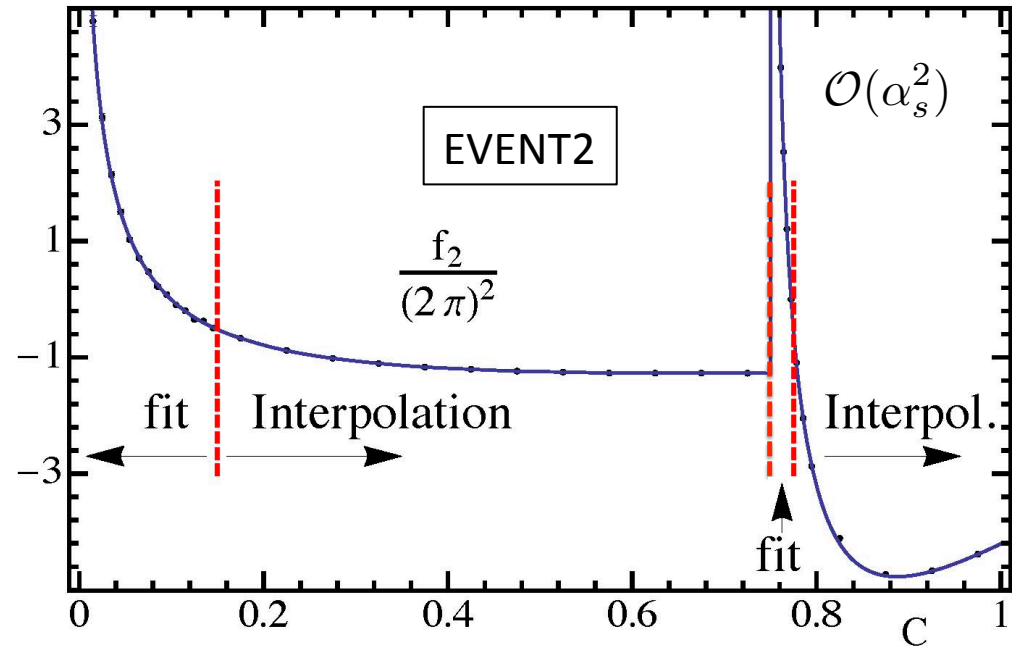
$$\frac{1}{\sigma_0} \frac{d\sigma^{(3)}}{dC}$$



Nonsingular Distributions

$$\sigma_{\text{ns}} = \sigma_{\text{full}}^{\text{FO}} - \sigma_{\text{s}}^{\text{FO}}$$

- Exact at one loop
- Subtract singular from EVENT2 (3×10^{11} events) and EERAD3 data (Gehrmann)
- Use norm to fit for $s_2 = -43 \pm 1$
- Fit in low region
- Interpolate in tail
- Fit around **shoulder** at $C = 0.75$
- More on this later

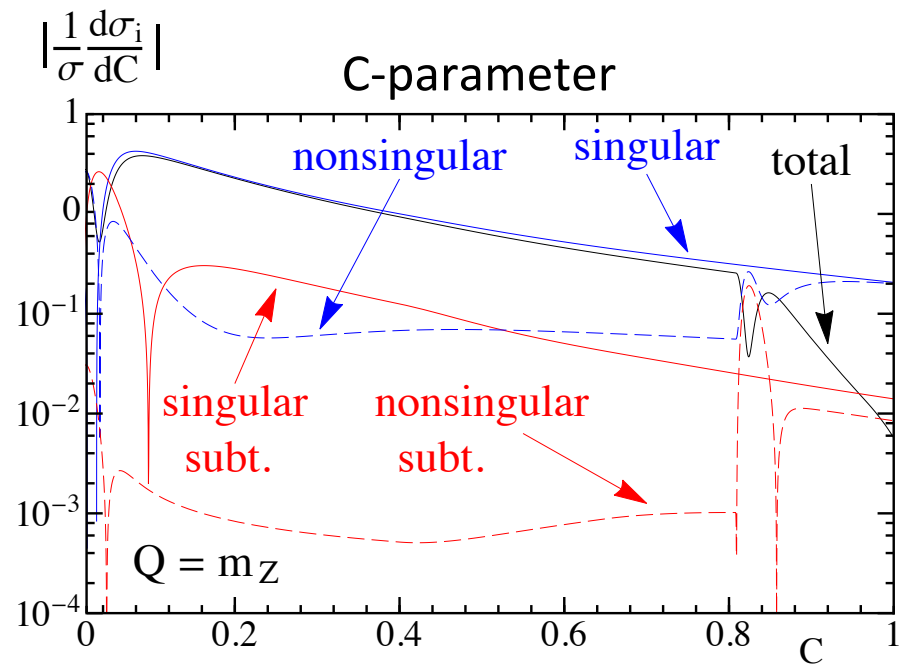
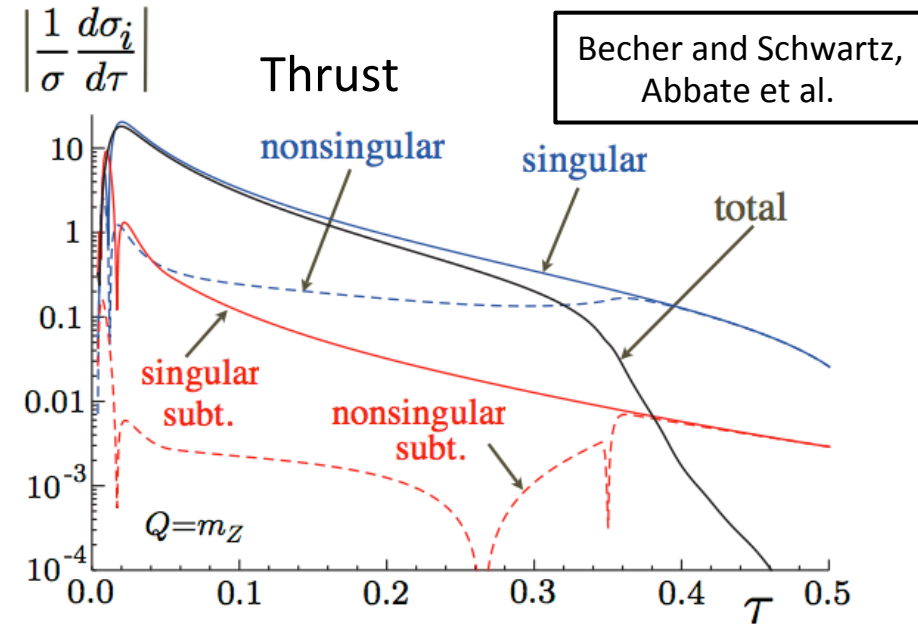


Nonsingular Distributions

$$\sigma_{\text{ns}} = \sigma_{\text{full}}^{\text{FO}} - \sigma_{\text{s}}^{\text{FO}}$$

Contrast with thrust

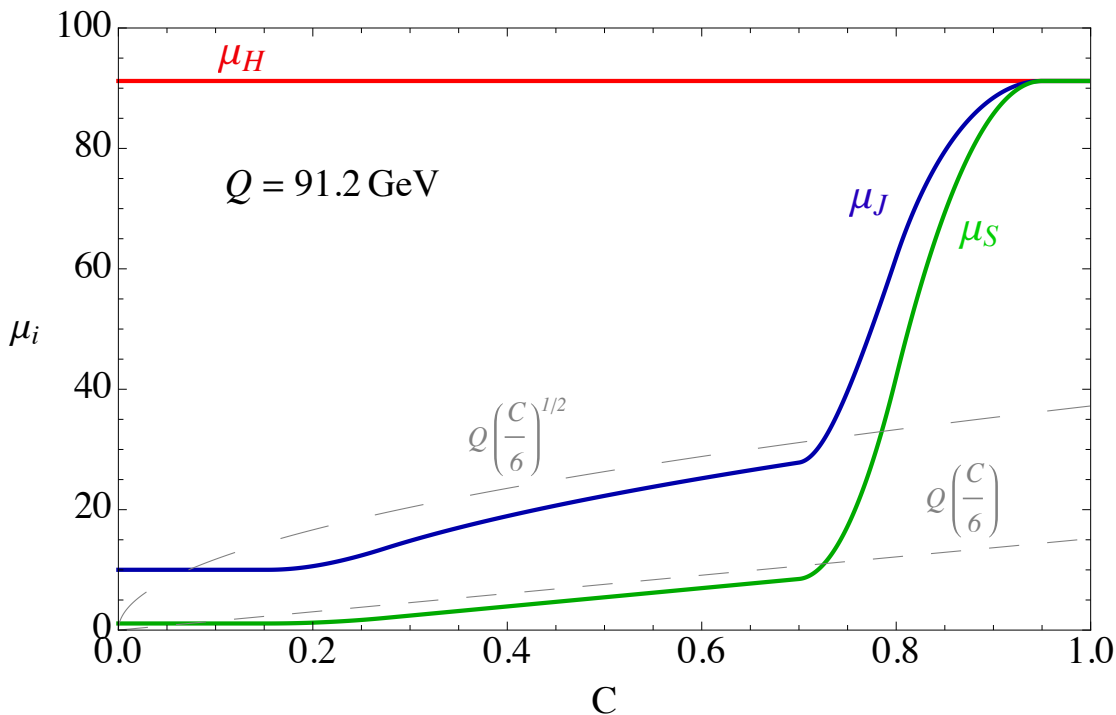
- Nonsingular is smaller up to $C = 0.8$
- Singular dominant region is larger in C-parameter than thrust



Profile Functions

Used by many authors

Pick scales to avoid large logs



$$\mu_H \sim Q$$

$$\mu_J \sim \sqrt{\mu_S \mu_H}$$

- Peak: $\mu_S \sim \Lambda_{\text{QCD}}$

- Tail: $\mu_S \sim \frac{QC}{6}$

Much gentler slope than for thrust

- Far-tail: $\mu_S \sim Q$

Shape Function

- Complete basis

Ligeti, Stewart,
Tackmann

$$F(k, \lambda, \{c_i\}) = \frac{1}{\lambda} \left[\sum_{n=0}^N c_n f_n \left(\frac{k}{\lambda} \right) \right]^2$$

- First moment is $\Omega_1^C = 3\pi \Omega_1^{gC}$

- Cannot expand in peak region

- Can expand in tail

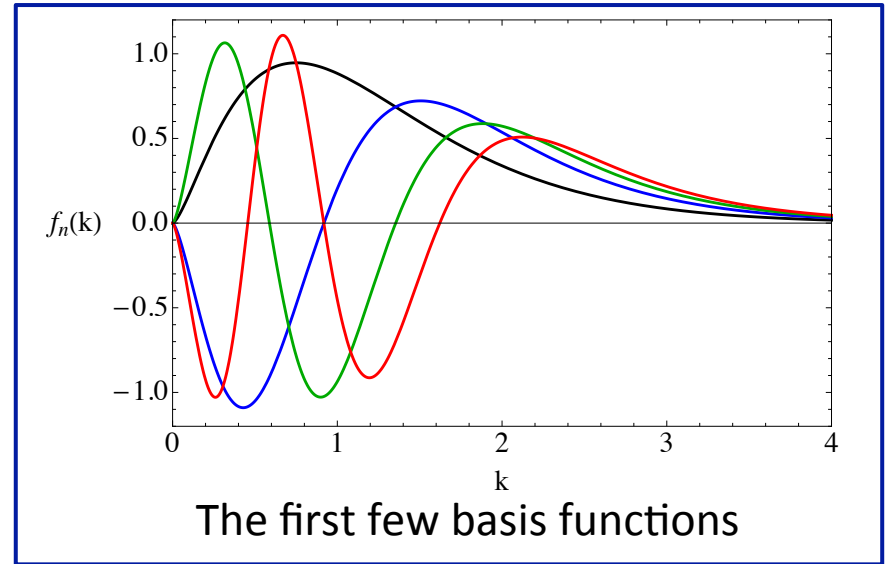
$$F_C(k, \lambda, \{c_i\}, \mu) = \delta(k) - 3\pi \bar{\Omega}_1^{gC}(\lambda, \{c_i\}, \mu) \delta'(k) + \dots$$

- Consider two renormalization schemes:

$\overline{\text{MS}}$ scheme: $\bar{\Omega}_1^{gC}$

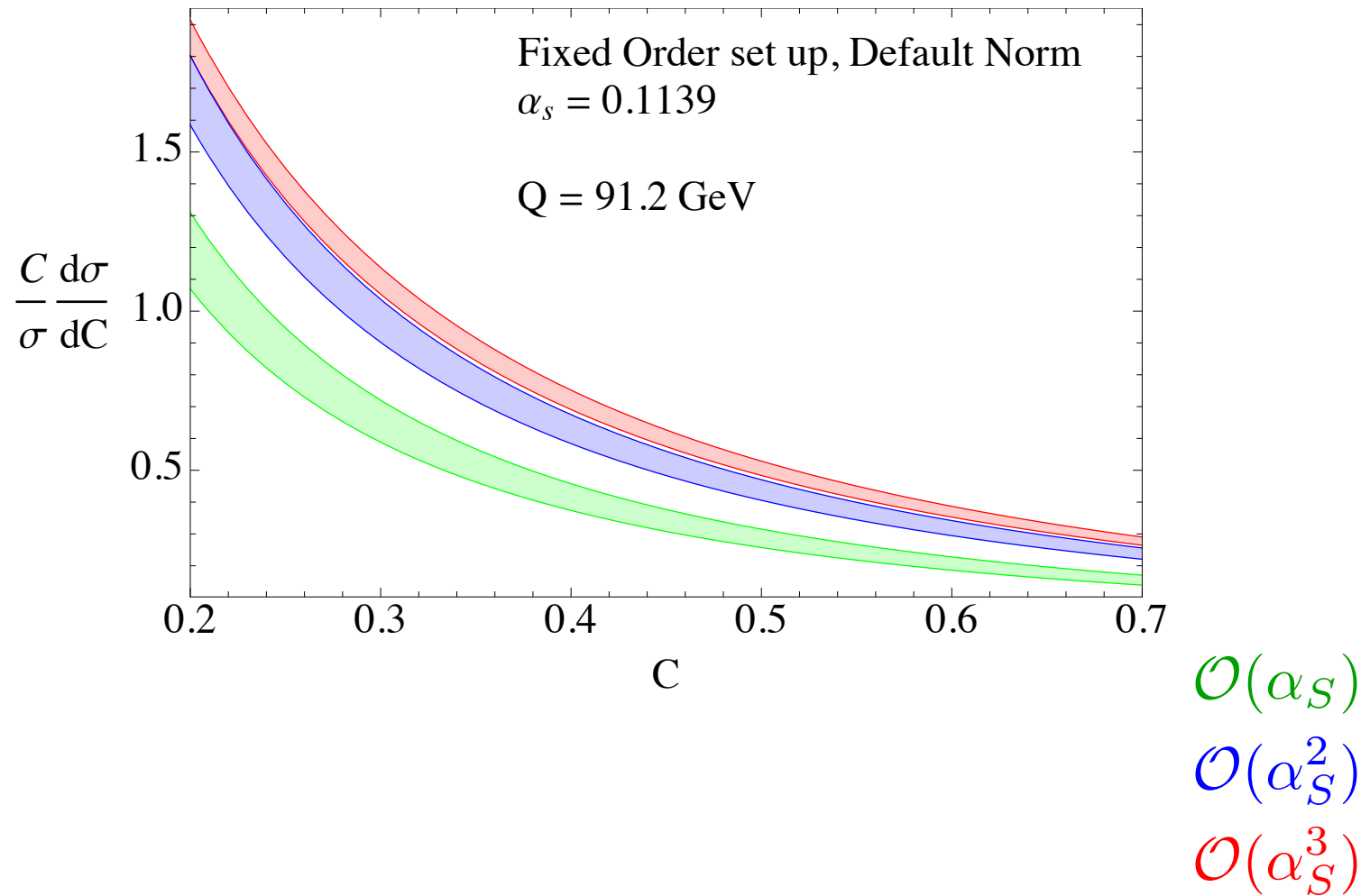
Renormalon free gap scheme: $\Omega_1^{gC}(R, \mu_S) = \bar{\Omega}_1^{gC} - \delta(R, \mu_S)$

$$\delta(R, \mu_S) = e^{\gamma_E} R \sum_k \alpha_s^k(\mu_S) \delta_k(R, \mu_S), \quad \frac{d\hat{\sigma}}{dC} = e^{-\left(\frac{3\pi\delta(R, \mu_S)}{Q} \frac{\partial}{\partial C}\right)} \frac{d\sigma}{dC}$$



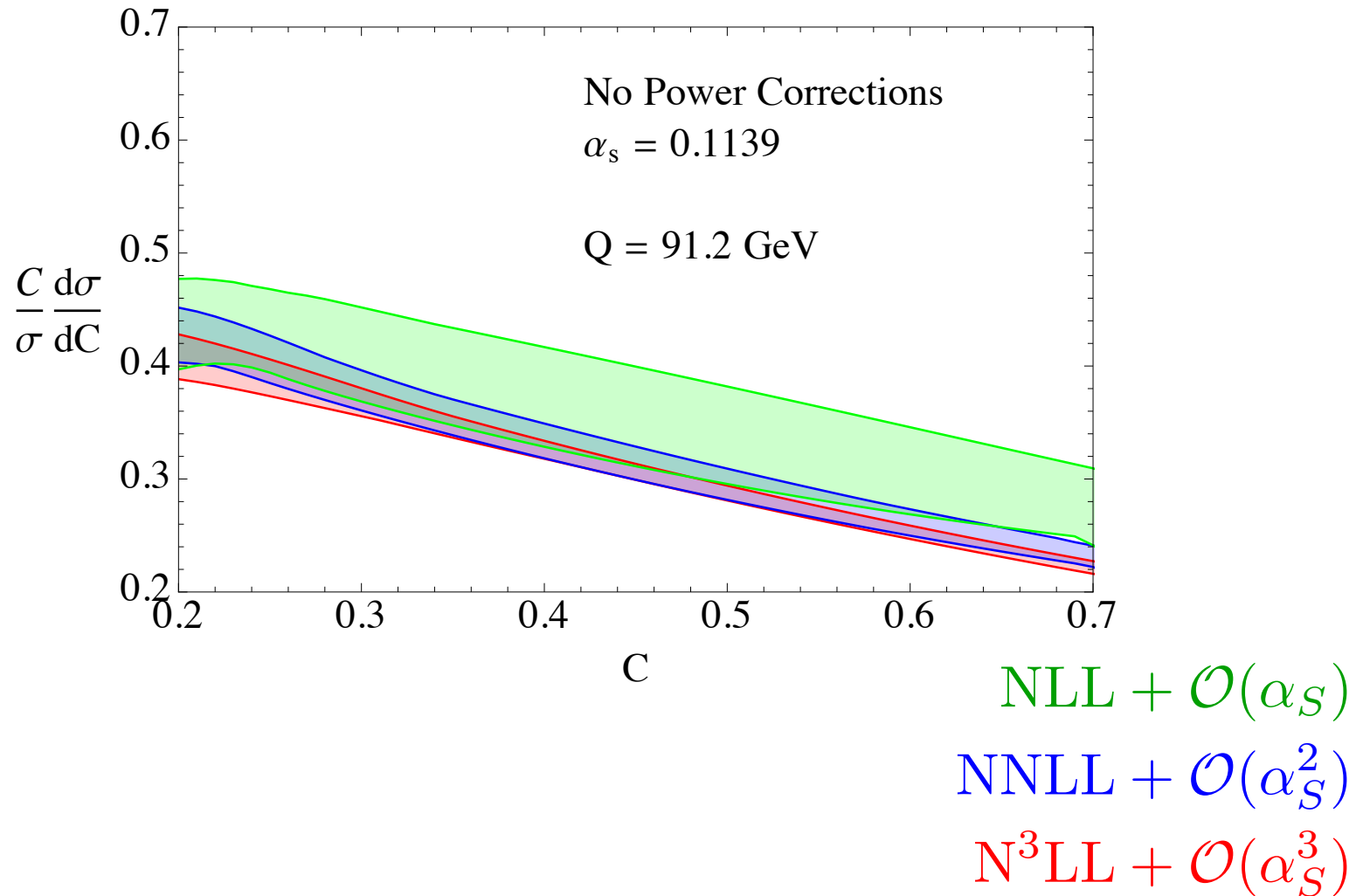
Full Distributions

Fixed Order



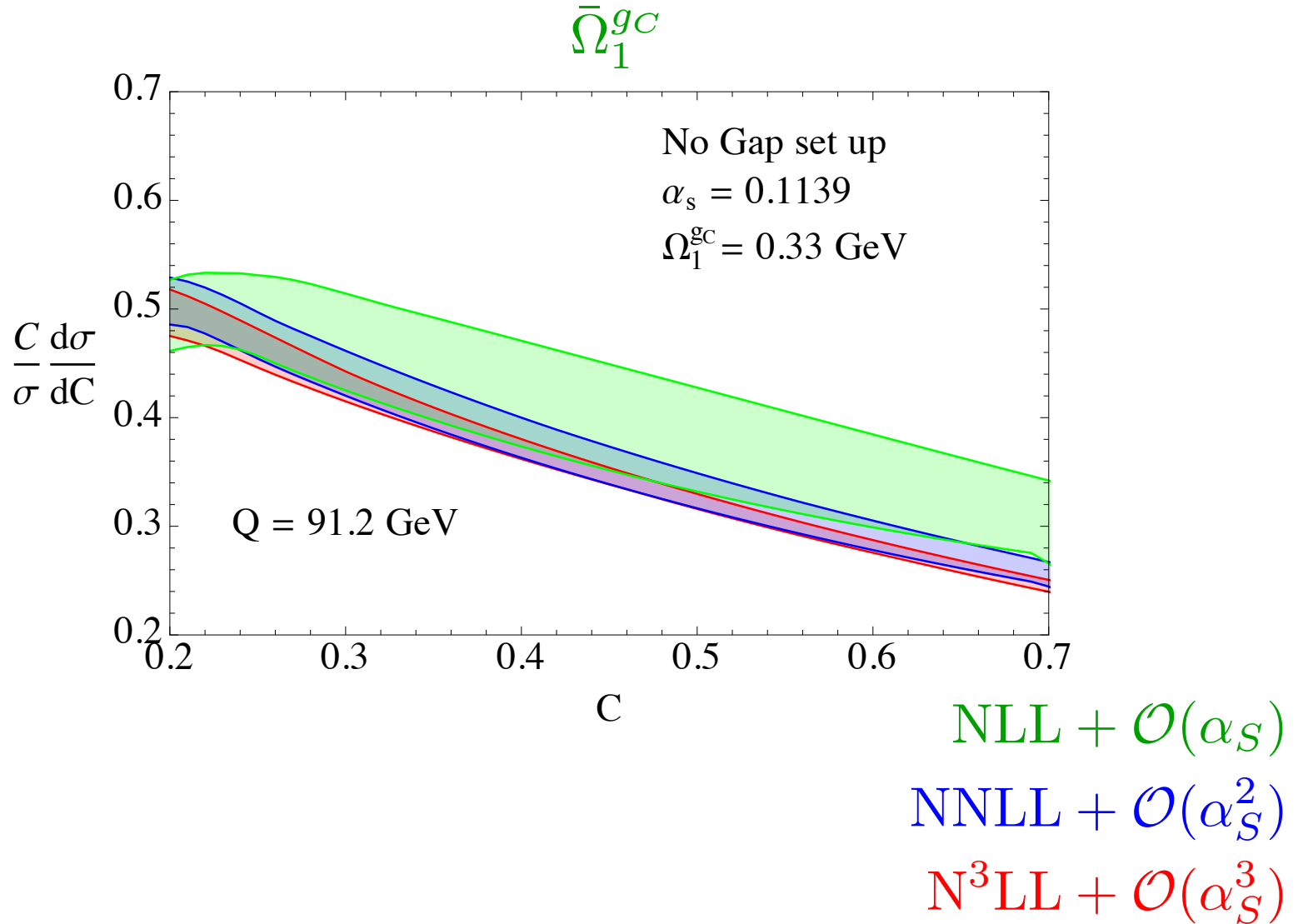
Full Distributions

Pure Resummation
(no nonperturbative shape function)



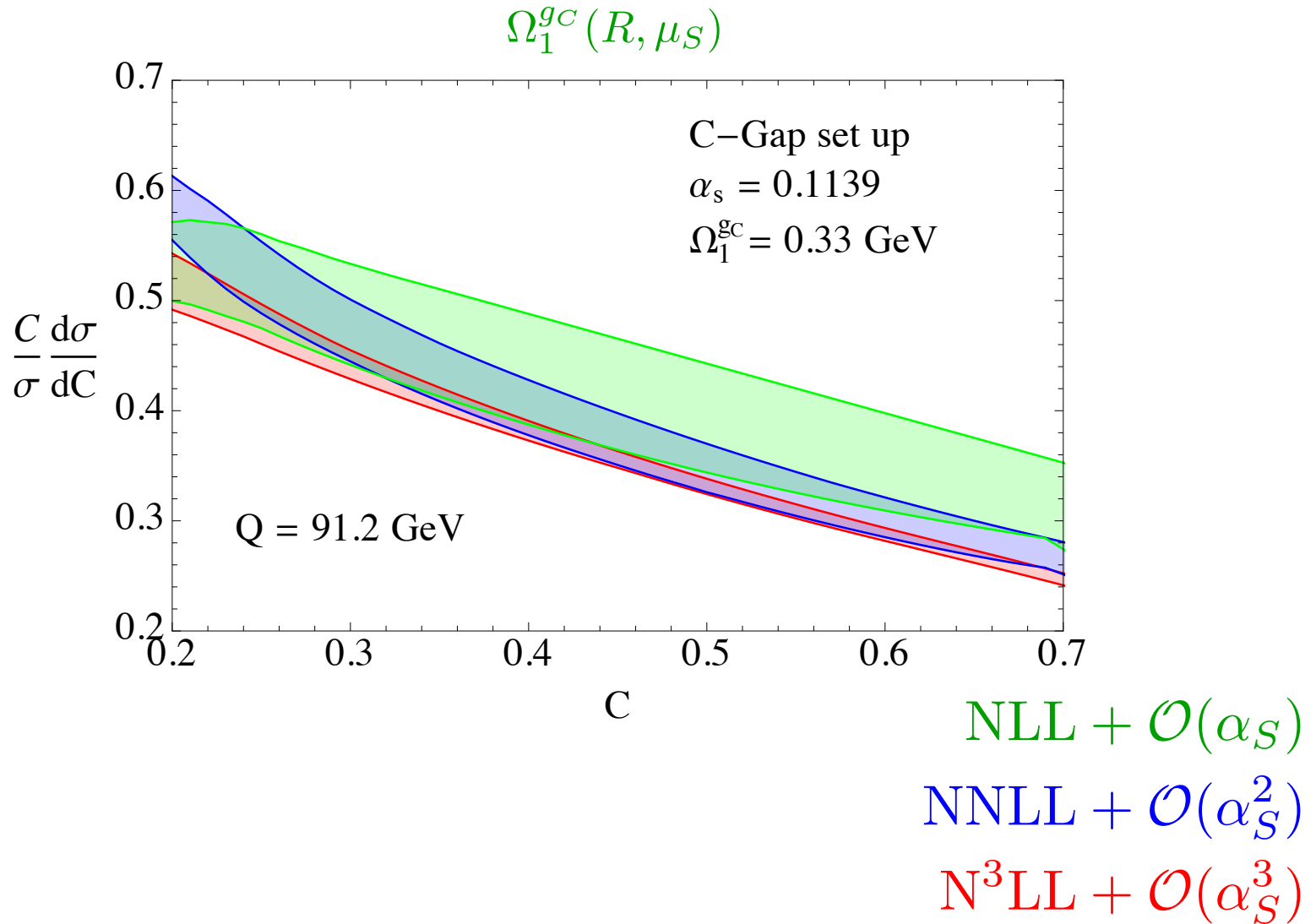
Full Distributions

Resummation and shape function



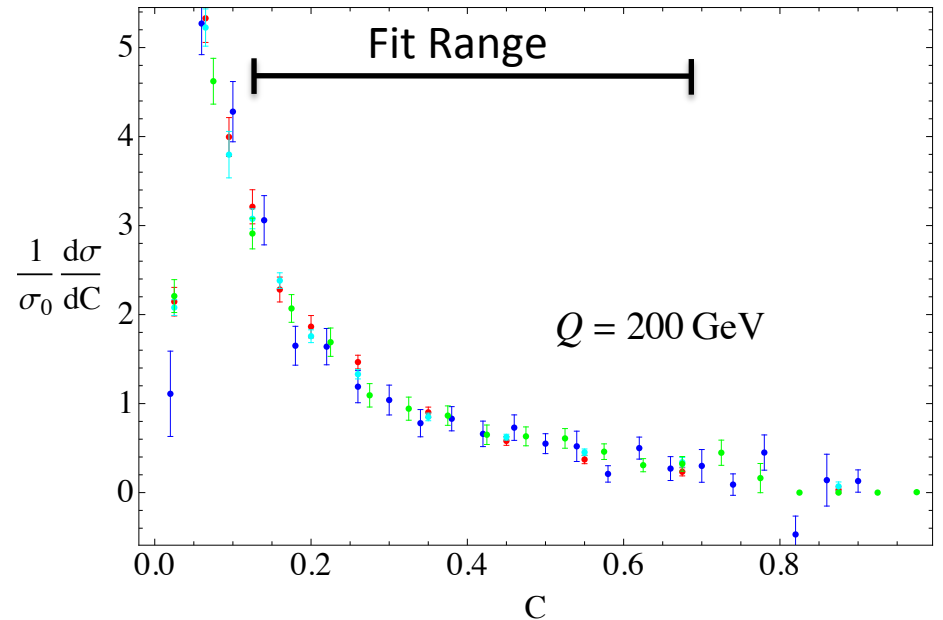
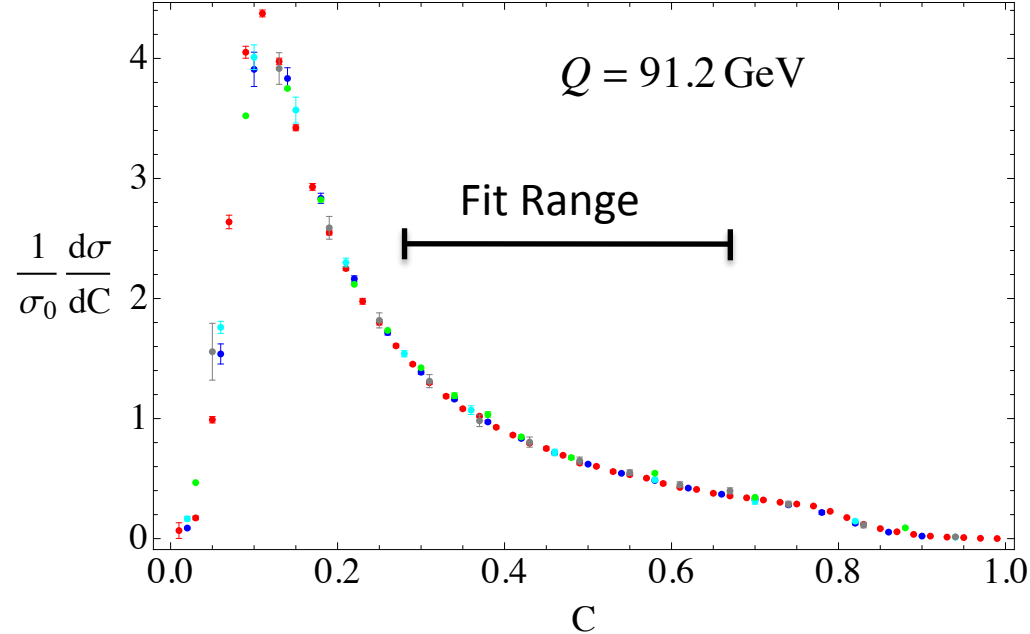
Full Distributions

Full result: Resummation, Shape function and renormalon free scheme



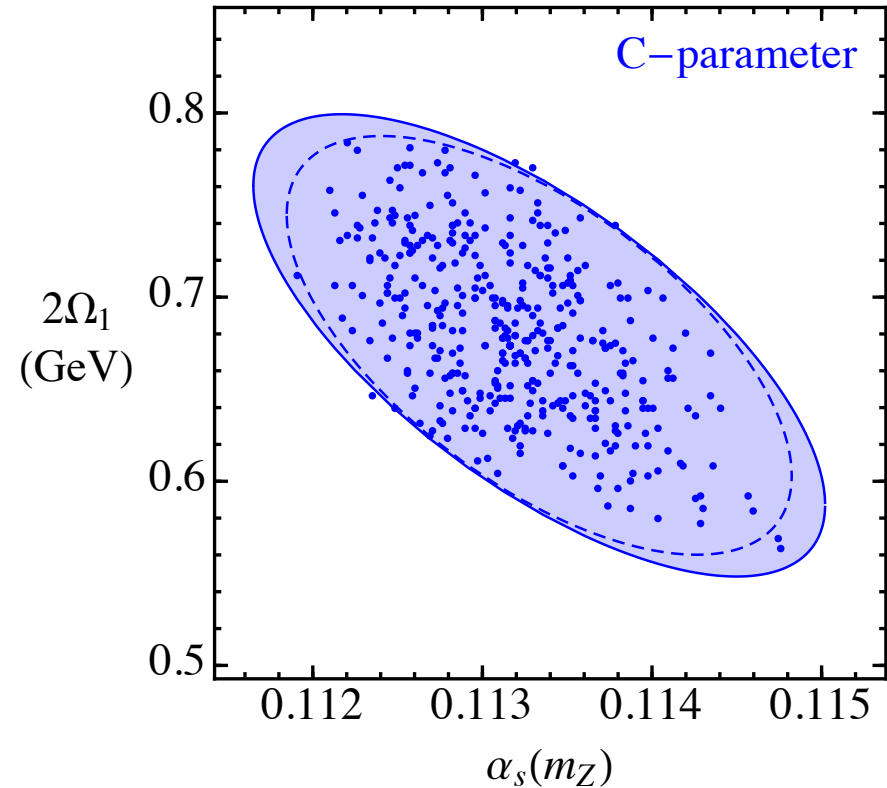
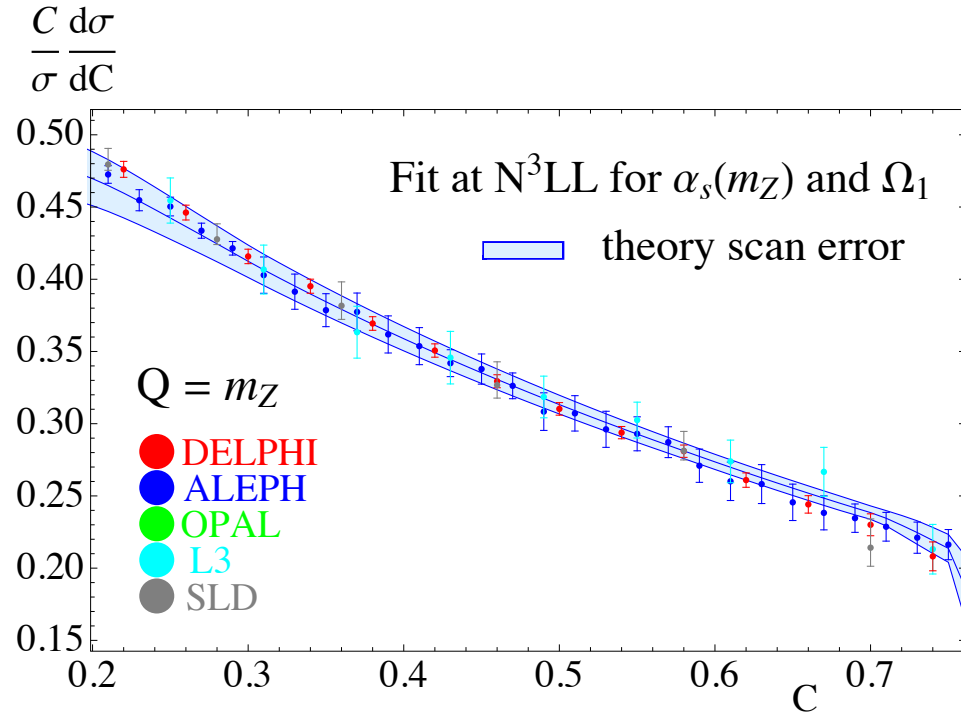
Data

- Six experiments:
ALEPH, **DELPHI**, **OPAL**,
L3, SLD and JADE
- Center of mass energies:
 $Q = 35 \text{ GeV} - 207 \text{ GeV}$
- Default fit range:
$$\frac{25 \text{ GeV}}{Q} \leq C \leq 0.7$$
- 404 bins in default fit region



Fit in the tail

Fit for Ω_1^{gC} and $\alpha_s(m_Z)$ simultaneously



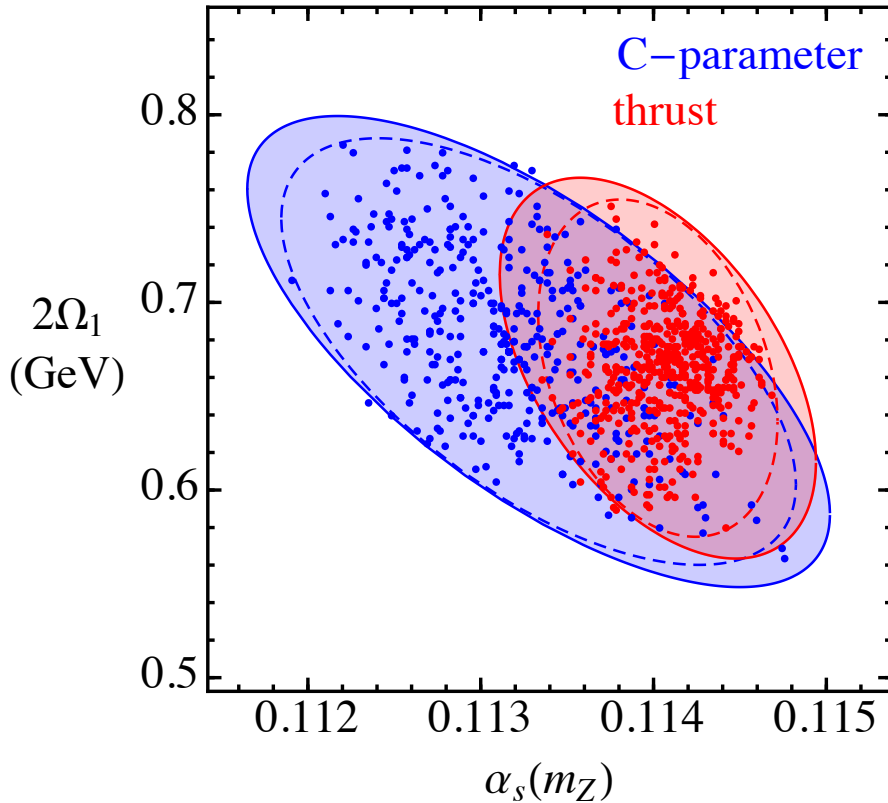
$$\alpha_s(m_Z) = 0.1133 \pm 0.0016$$

(Preliminary)

$$\Omega_1^{gC} = 0.34 \pm 0.06 \text{ GeV}$$

$$\frac{\chi^2}{\text{d.o.f.}} = 0.97$$

Comparison with Thrust Results



C-Parameter (pure QCD)

$$\alpha_s(m_z) = 0.1133 \pm 0.0016$$

$$\Omega_1^{gC} = 0.34 \pm 0.06 \text{ GeV}$$

Thrust (pure QCD, Abbate et al.)

$$\alpha_s(m_z) = 0.1140 \pm 0.0011$$

$$\Omega_1^{g\tau} = 0.33 \pm 0.05 \text{ GeV}$$

Even including hadron mass effects,

$$\Omega_1^{gC} = \Omega_1^{g\tau} + \mathcal{O}(10\%)$$

(Salam and Wicke;
Mateu, Stewart, Thaler)

Thrust (with QED and m_b)

$$\alpha_s(m_z) = 0.1135 \pm 0.0011$$

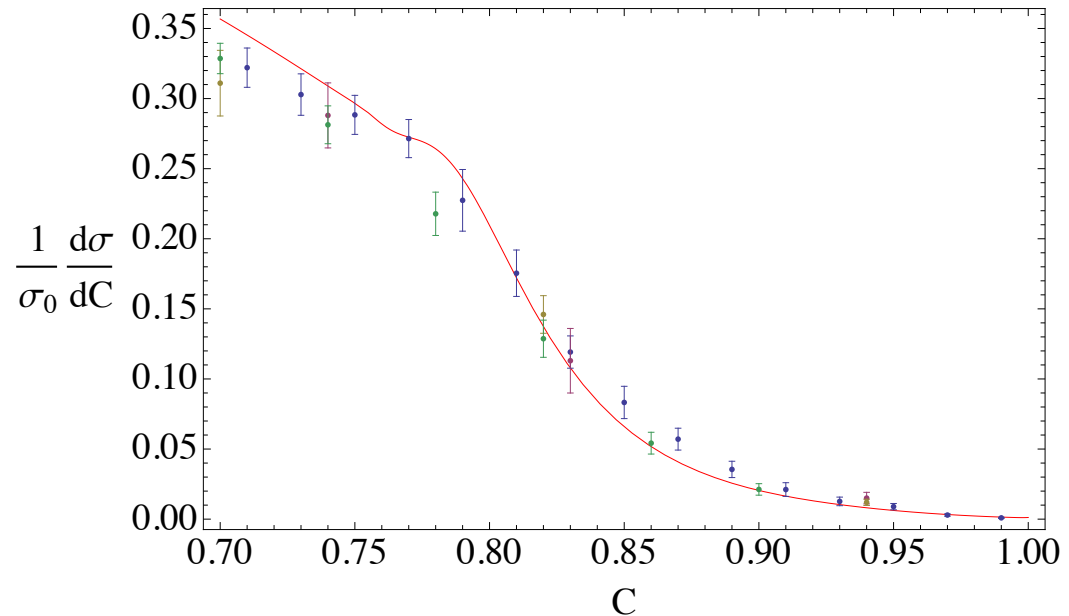
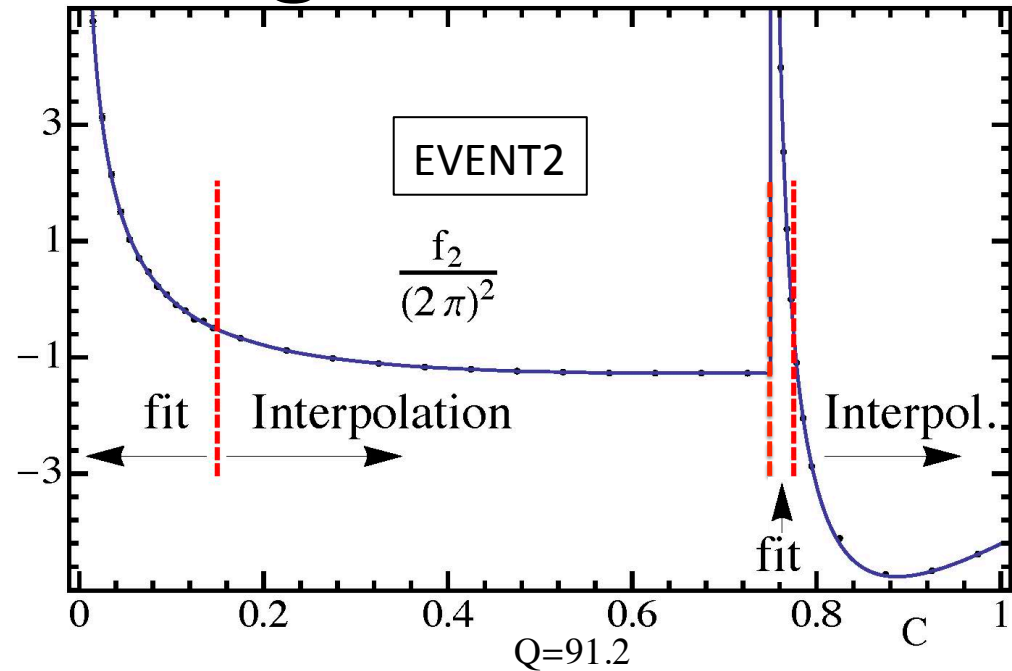
$$\Omega_1^{g\tau} = 0.32 \pm 0.05 \text{ GeV}$$

Shoulder Region

- Treat to order α_s^3
- Smear with F
- No subtractions
- This region is **not** in fit
- Proper treatment needs another resummation

$$\sim \alpha_s \left(\alpha_s \ln^2 \left(C - \frac{3}{4} \right) + \dots \right)$$

(NLL exists, Catani and Webber)



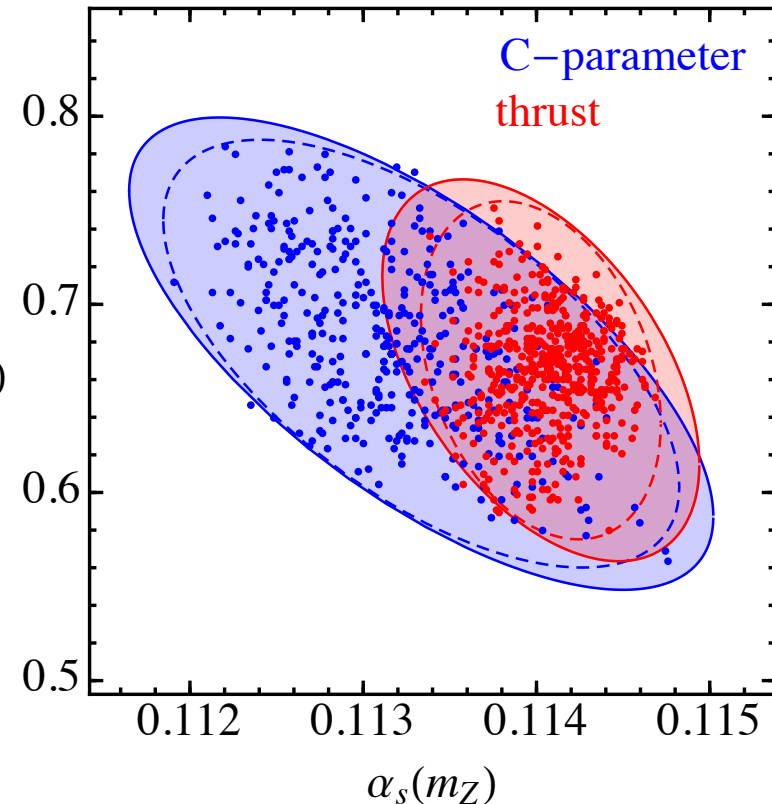
Conclusions

- Have used factorization for C-parameter to obtain
 - N³LL + $\mathcal{O}(\alpha_s^3)$ result
 - Field theory treatment of power corrections with renormalon subtractions
 - Valid treatment of all kinematic regions
- Global fit gives values

$$\alpha_s(m_z) = 0.1133 \pm 0.0016$$

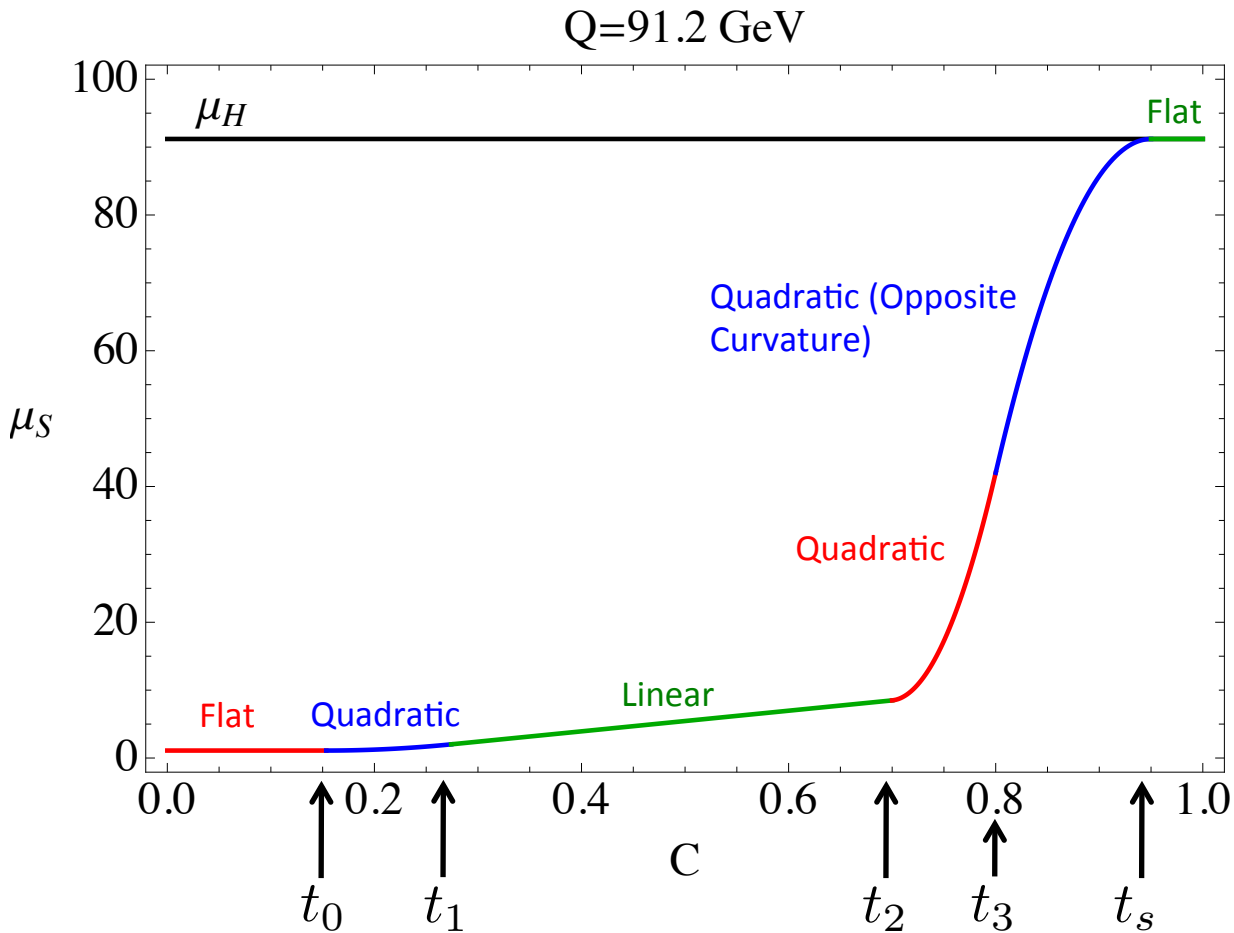
$$\Omega_1^{gC} = 0.34 \pm 0.06 \text{ GeV} \quad 2\Omega_1 \text{ (GeV)}$$

- Agrees with earlier thrust results



Backup Slides

Profile Functions

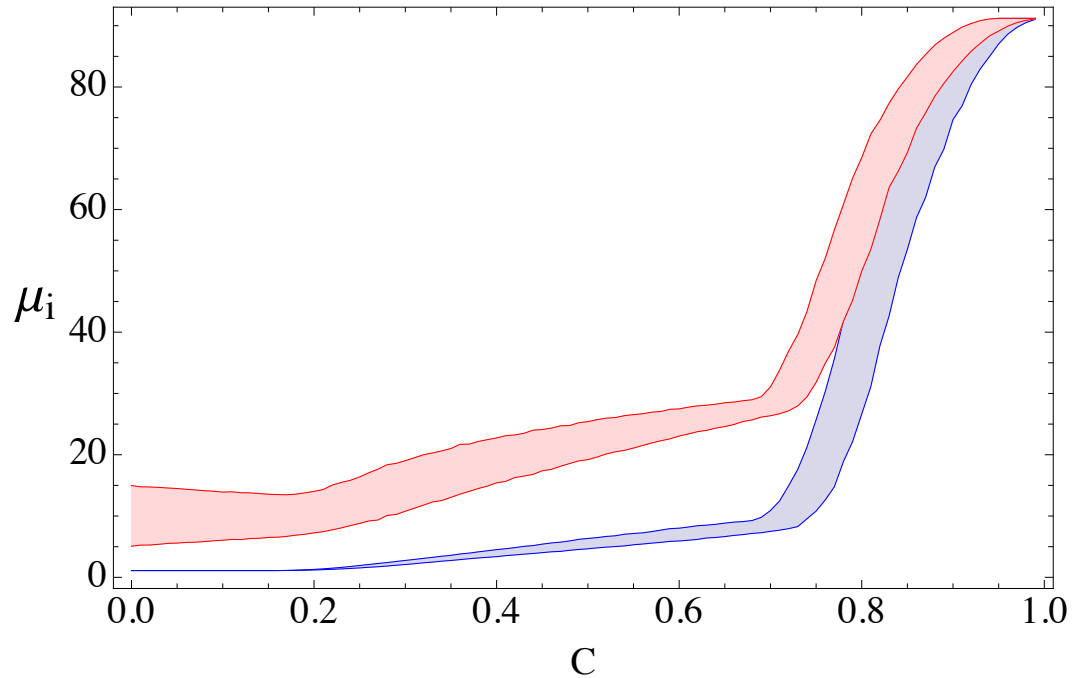


Parameter	Default Value
μ_0	1.1 GeV
r	$\frac{1}{6}$
t_0	$\frac{14 \text{ GeV}}{Q}$
t_1	$\frac{25 \text{ GeV}}{Q}$
t_2	0.7
t_3	0.8
t_s	0.95
$e_H = \frac{\mu_H}{Q}$	1

Parameters varied to
get the perturbative
uncertainty

Profile Variation for Fits

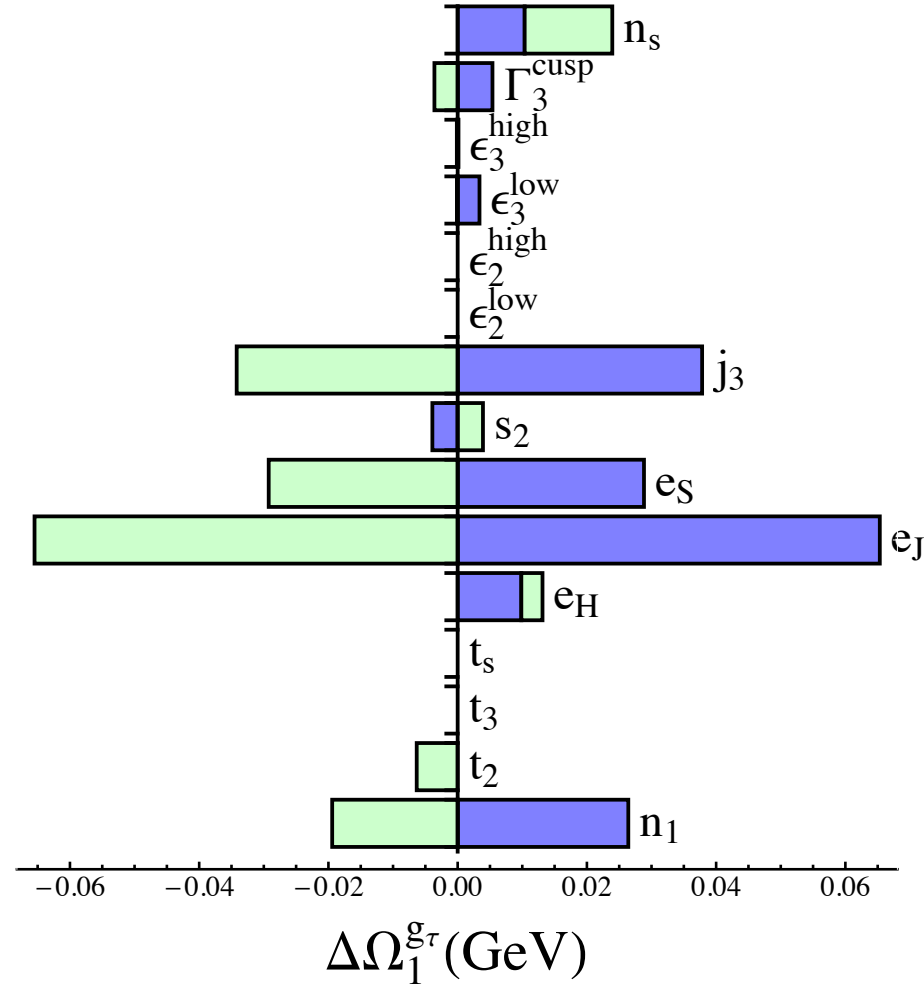
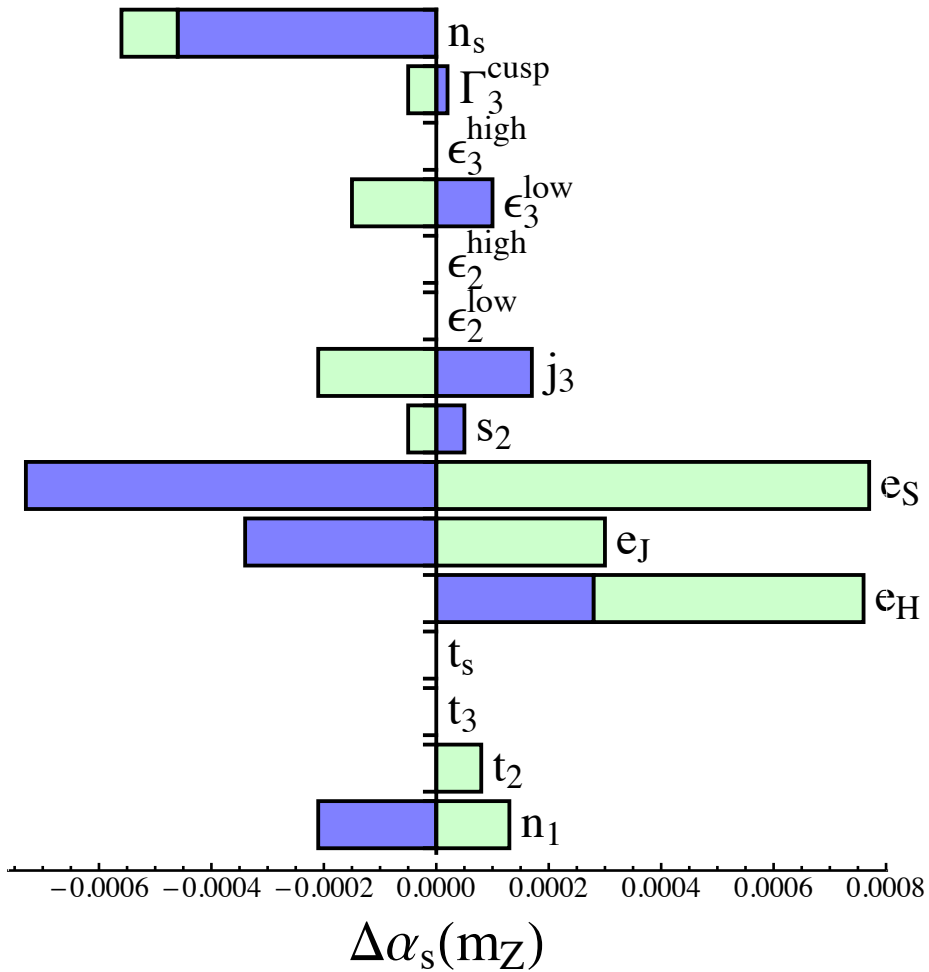
Variation of the soft (blue) and jet (red) scales at $Q=91.2$ GeV



Plus there is an additional up & down variation of all scales by factor of 2 & $\frac{1}{2}$

Theory Error

Vary parameters in profile and model to estimate theory error



Preliminary