

Introduction
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Calculation
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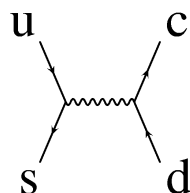
Comparison with PYTHIA
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Track-Based Observables
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Conclusions
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Jet Charge and Track-Based Observables at the LHC

Wouter Waalewijn
UCSD



SCET 2013

With: David Krohn, Tongyan Lin and Matthew Schwartz
arXiv: 1209.2421, 1209.3019

Hsi-Ming Chang, Massimiliano Procura, Jesse Thaler
arXiv: 1303.xxxx

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Outline

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② Calculation

③ Comparison with PYTHIA

④ Track-Based Observables

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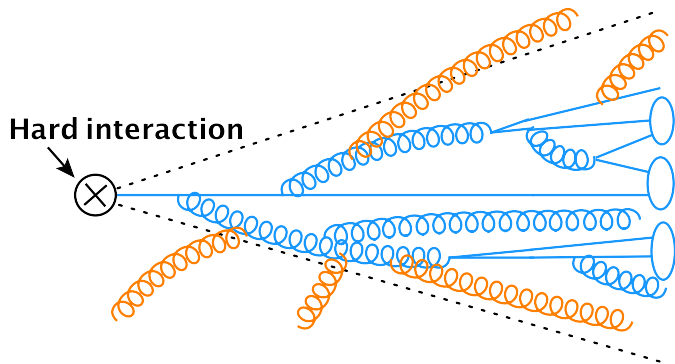
Track-Based Observables
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Introduction

Jet Charge

- ▶ Motivation: separate q and \bar{q} jets, flavor tag

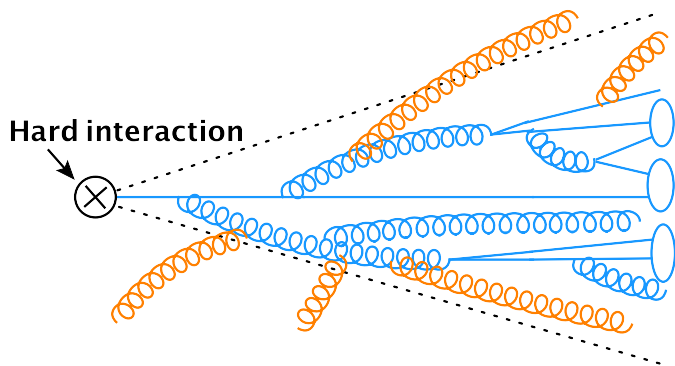


$$Q_\kappa = \sum_{h \in \text{jet}} \left(\frac{p_T^h}{p_T^{\text{jet}}} \right)^\kappa Q_h$$

[Feynman, Field (1977)]

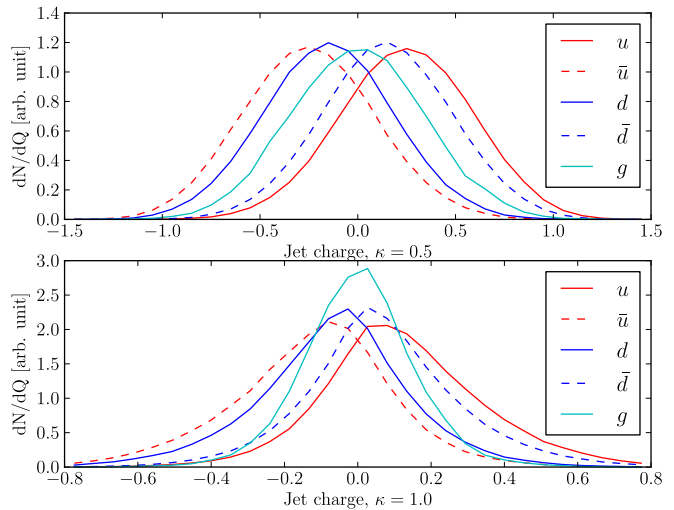
Jet Charge

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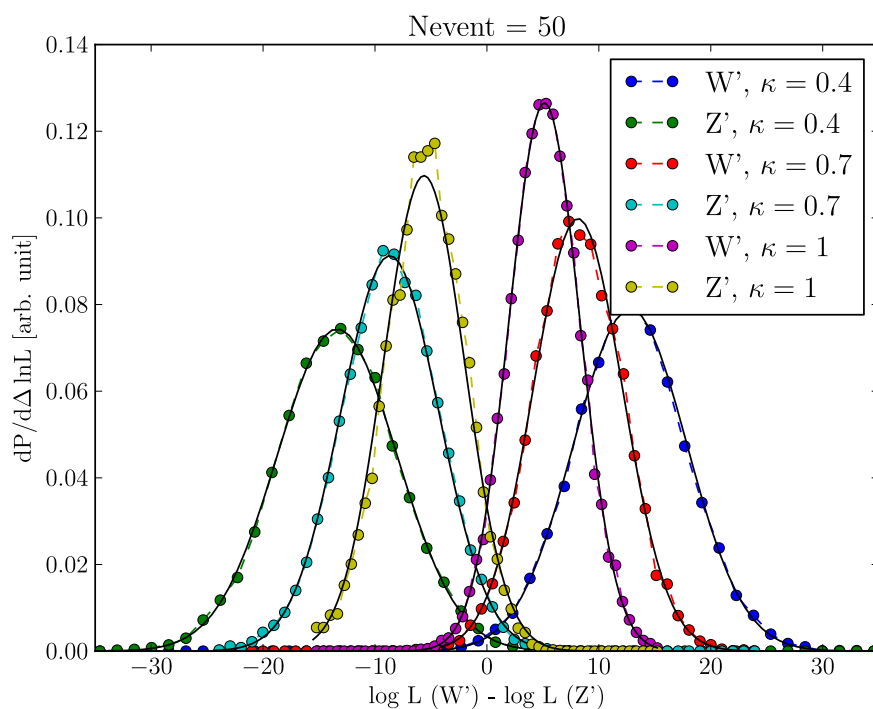
$$Q_\kappa = \sum_{h \in \text{jet}} \left(\frac{p_T^h}{p_T^{\text{jet}}} \right)^\kappa Q_h$$

[Feynman, Field (1977)]



- ▶ κ too small: measurement sensitive to soft hadrons → contamination
- ▶ κ too large: need more statistics to separate flavors

LHC Application: W' vs. Z'



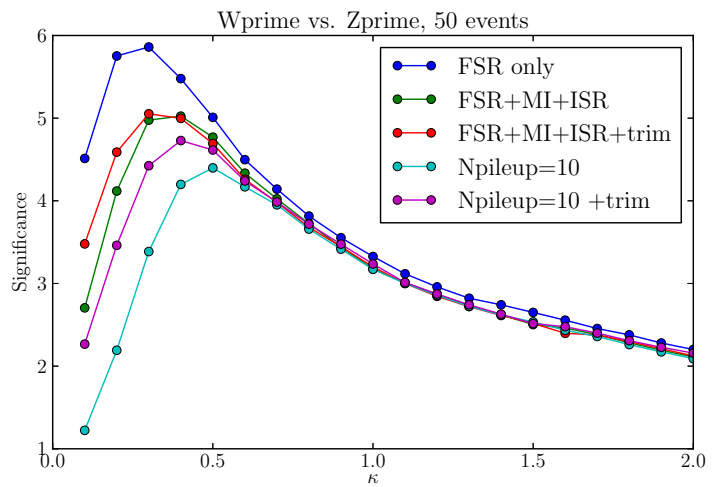
- ▶ Hadronically decaying W' and Z' with mass 1 TeV
- ▶ Likelihood discriminant based on charge of both jets
- ▶ With 50 events $\sim 4\sigma$ separation

LHC Challenges

Contamination:

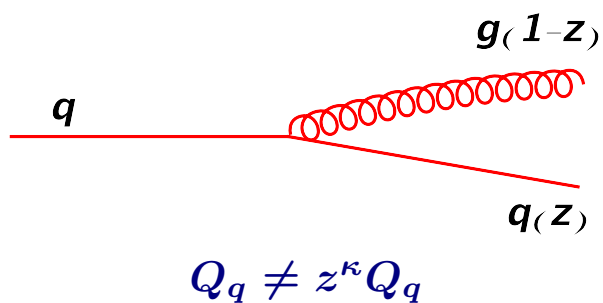
- ▶ Initial-state radiation (ISR)
- ▶ Multi-parton interactions (MI)
- ▶ Pile-up

Soft effects → issue for **small κ**



Jet Charge Not Infrared Safe

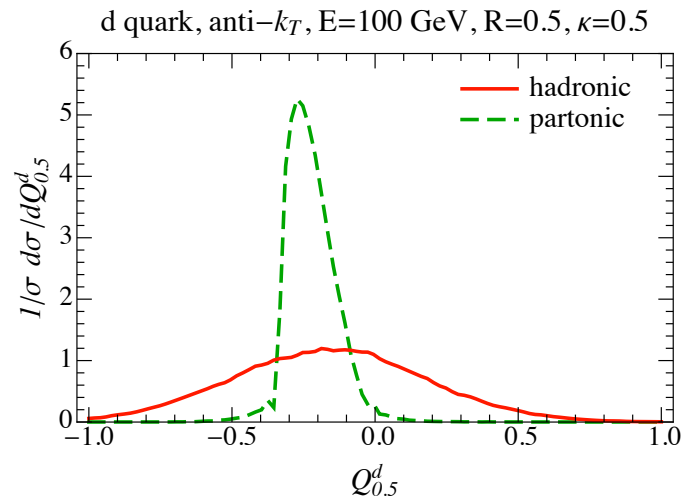
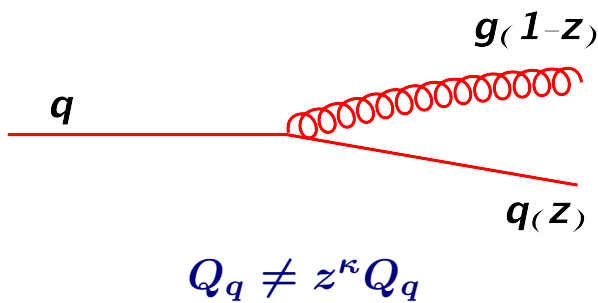
- ▶ Example: consider $q \rightarrow qg$ in collinear limit



- ▶ Jet charge only defined for hadrons

Jet Charge Not Infrared Safe

- ▶ Example: consider $q \rightarrow qg$ in collinear limit



- ▶ Jet charge only defined for hadrons
- ▶ Importance of hadronization observed in PYTHIA

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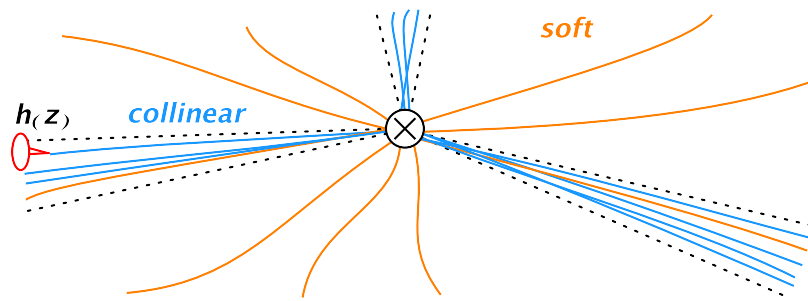
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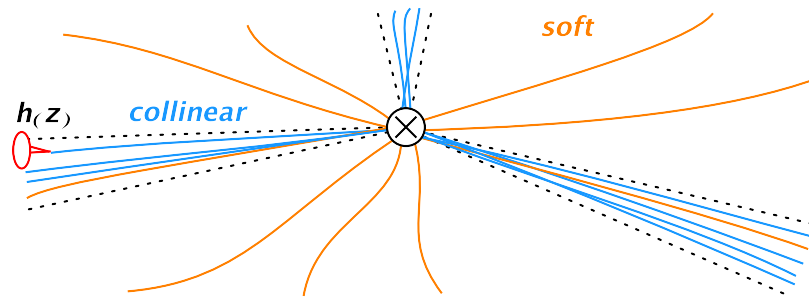
Average Jet Charge



Jet with energy E and radius R :

$$\langle Q_{\kappa}^q \rangle = \underbrace{\sum_h}_{\text{hadron } h \text{ with } z} \int dz \underbrace{z^{\kappa} Q_h}_{\text{charge}} \underbrace{\frac{1}{\sigma_{\text{jet}}} \frac{d\sigma_{h \in \text{jet}}}{dz}}_{\text{weight}}$$

Average Jet Charge



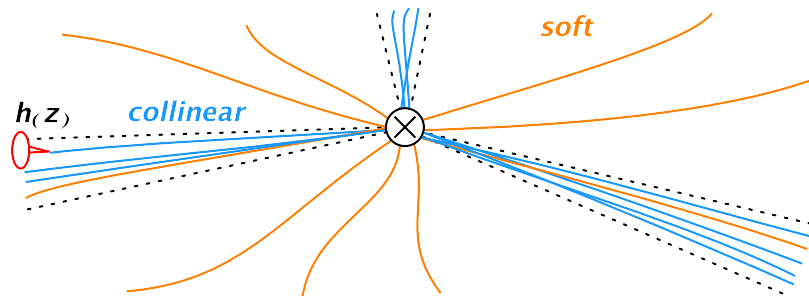
Jet with energy E and radius R :

[Using Ellis, Hornig, Lee, Vermilion and Walsh (2010)]

$$\langle Q_{\kappa}^q \rangle = \sum_h \int dz z^{\kappa} Q_h \frac{1}{\sigma_{\text{jet}}} \frac{d\sigma_{h \in \text{jet}}}{dz} = \sum_h \int dz z^{\kappa} Q_h \frac{\mathcal{G}_q^h(E R, z, \mu)}{J_q(E R, \mu)}$$

- ▶ Dependence on hard collision, soft radiation and other jets suppressed

Average Jet Charge



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- ▶ Dependence on hard collision, soft radiation and other jets suppressed
- ▶ Dependence on perturbative ER calculable:

$$\mathcal{G}_i^h(E R, z, \mu) = \sum_j \int_z^1 \frac{dx}{x} \mathcal{J}_{ij}(E R, x, \mu) D_j^h\left(\frac{z}{x}, \mu\right)$$

[Procura, Stewart (2010), Jain, Procura, WW (2011)]

- ▶ μ -dependence cancels between \mathcal{G}_q^h and J_q
- ▶ Evaluate at jet scale $\mu \sim ER$ to avoid large logs

Interpretation

- ▶ At leading order:

$$\langle Q_{\kappa}^q \rangle = 1 \times \sum_h Q_h \tilde{D}_q^h(\kappa, \mu = ER)$$

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$$\mu \frac{d}{d\mu} \tilde{D}_i^h(\kappa, \mu) = \sum_j \frac{\alpha_s(\mu)}{\pi} \tilde{P}_{ji}(\kappa) \tilde{D}_j^h(\kappa, \mu),$$

- ▶ Mixing into gluons will vanish, since $D_g^{h^+} - D_g^{h^-} = 0$

Interpretation

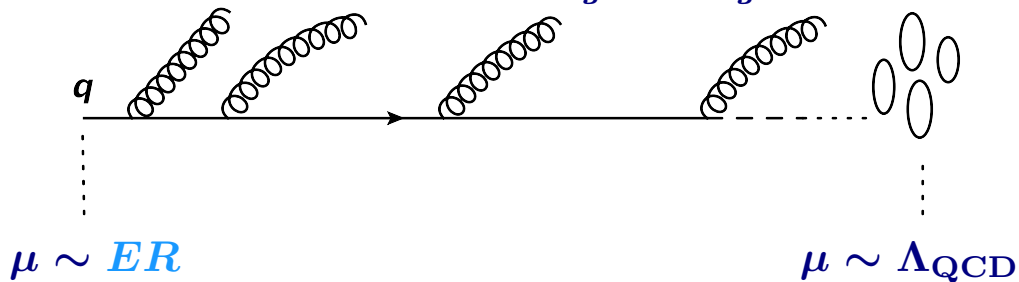
- ▶ At leading order:

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- ▶ Showering starts at **jet scale** not hard scale

Interpretation

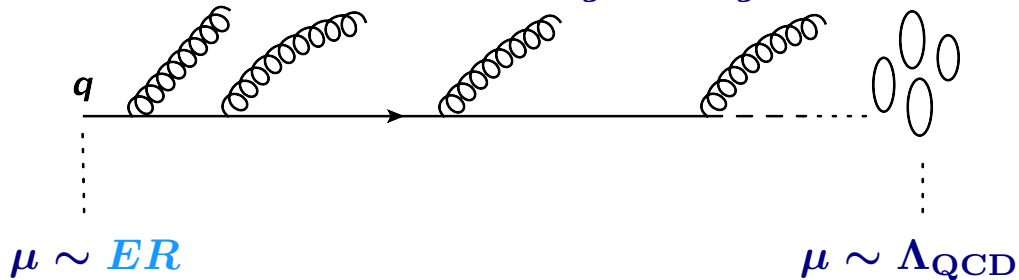
- ▶ At next-to-leading order:

$$\langle Q_{\kappa}^q \rangle = \tilde{\mathcal{J}}_{qq}(ER, \kappa, \mu = ER) \sum_h Q_h \tilde{D}_q^h(\kappa, \mu = ER)$$

- ▶ ER dependence from moment-space evolution:

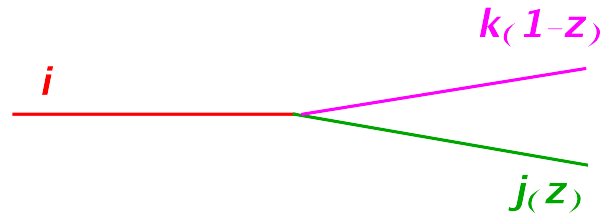
$$\mu \frac{d}{d\mu} \tilde{D}_i^h(\kappa, \mu) = \sum_j \frac{\alpha_s(\mu)}{\pi} \tilde{P}_{ji}(\kappa) \tilde{D}_j^h(\kappa, \mu),$$

- ▶ Mixing into gluons will vanish, since $D_g^{h^+} - D_g^{h^-} = 0$



- ▶ Showering starts at **jet scale** not hard scale
- ▶ Perturbative splittings at beginning of shower

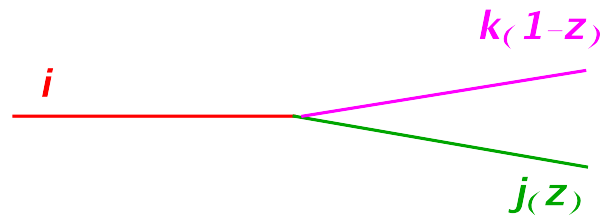
Charge Distribution



Nonperturbative charge distribution $D_i(Q_\kappa, \mu)$

$$\mu \frac{d}{d\mu} D_i(Q_\kappa, \mu) = \sum_j \int dz \frac{\alpha_s}{2\pi} P_{ji}(z) \int dQ_\kappa^a D_j(Q_\kappa^a, \mu) \int dQ_\kappa^b D_k(Q_\kappa^b, \mu) \times \delta[Q_\kappa - z^\kappa Q_\kappa^a - (1-z)^\kappa Q_\kappa^b]$$

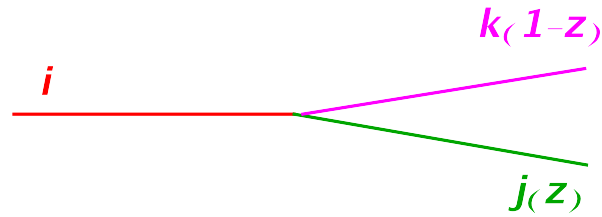
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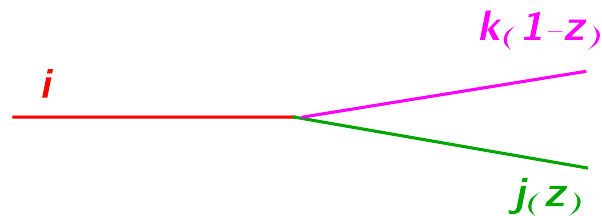
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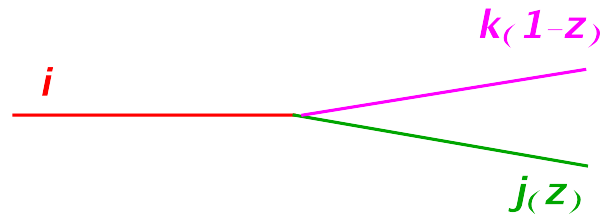
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$$\begin{aligned}
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 &\times \underbrace{\delta[Q_\kappa - z^\kappa Q_\kappa^a - (1-z)^\kappa Q_\kappa^b]}_{\text{Calculate charge from branches}}
 \end{aligned}$$

Charge Distribution



Nonperturbative charge distribution $D_i(Q_\kappa, \mu)$

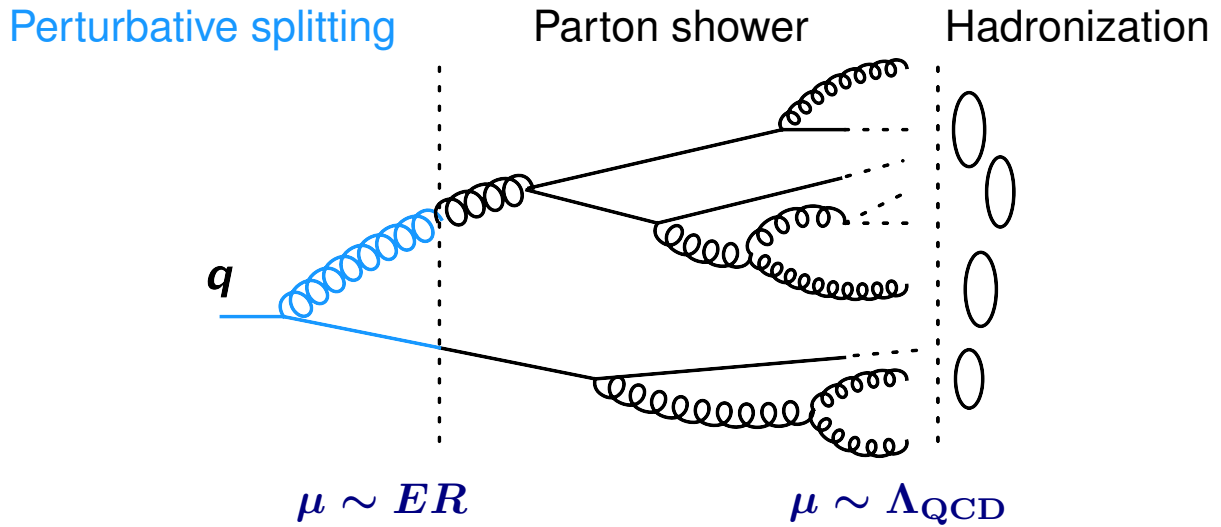
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- Preserves norm:

$$\frac{d}{d\mu} \int dQ_\kappa D_i(Q_\kappa, \mu) = 0$$

- First two moments reproduce average and width
- NLO perturbative splittings similar form: $P_{ji} \rightarrow \mathcal{J}_{ij}$

Parton Shower



- ▶ DGLAP evolution following **all branches** instead of single parton
- ▶ This parton shower does not yield full final-state kinematics
- ▶ Perturbative splitting generates perturbative scales (e.g. mass)

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Comparison with PYTHIA

Comparison with Fragmentation Functions

$$\langle Q_{\kappa}^q \rangle = \tilde{\mathcal{J}}_{qq}(ER, \kappa, \mu = ER) \sum_h Q_h \tilde{D}_q^h(\kappa, \mu = ER)$$

Average jet charge for $E = 100$ GeV and $R = 0.5$:

κ	u -quark			d -quark		
	PYTHIA	DSS	AKK08	PYTHIA	DSS	AKK08
0.5	0.271	0.237	0.221	-0.162	-0.184	-0.062
1	0.144	0.122	0.134	-0.078	-0.088	-0.046
2	0.055	0.046	0.064	-0.027	-0.030	-0.027

[DSS = de Florian, Sassot, Stratmann (2007), AKK08 = Albino, Kniehl, Kramer (2008)]

- ▶ PYTHIA consistent with fragmentation functions

Comparison with Fragmentation Functions

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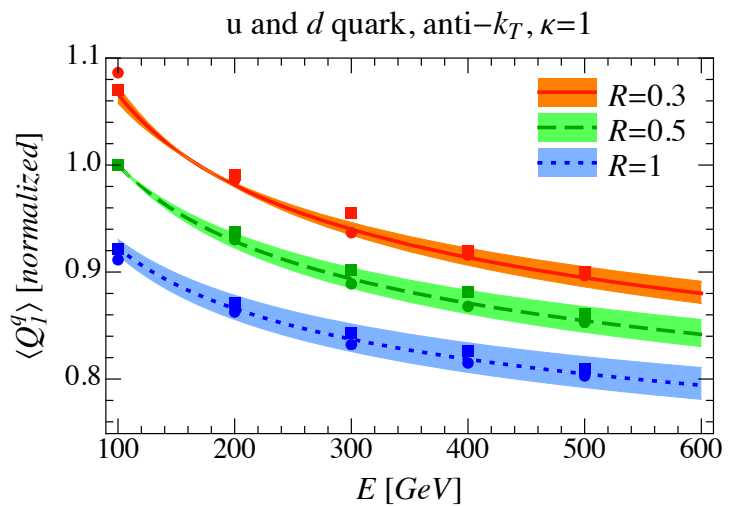
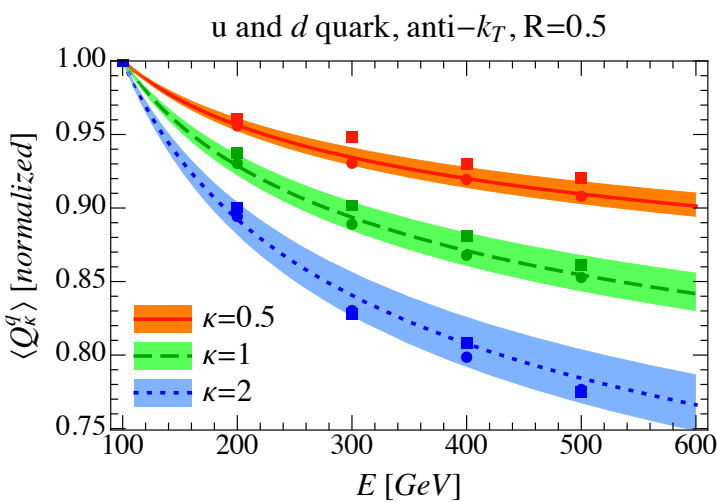
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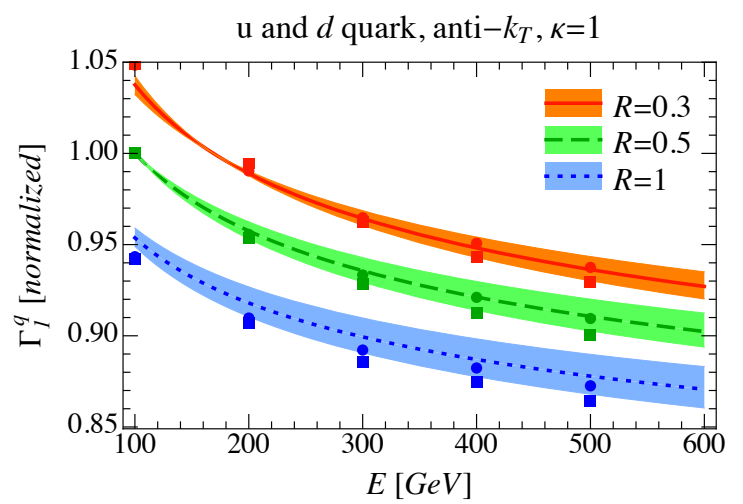
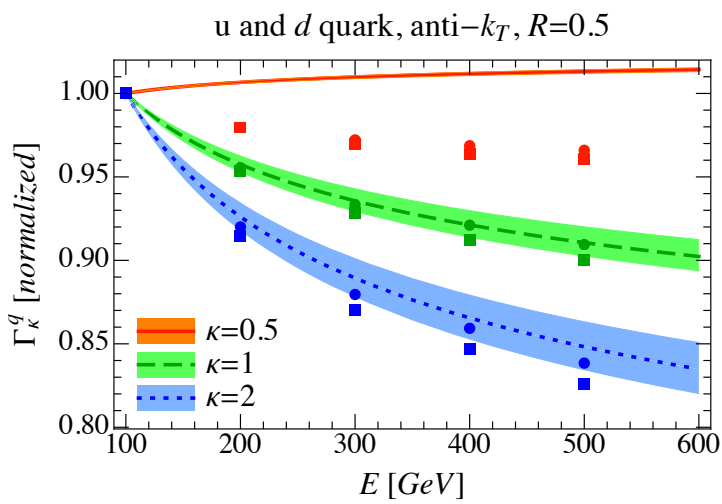
- ▶ PYTHIA consistent with fragmentation functions
- ▶ Large uncertainties because we need $D_q^{h^+} - D_q^{h^-} = D_q^{h^+} - D_{\bar{q}}^{h^+}$
Most fragmentation data is e^+e^- giving $D_q^{h^+} + D_{\bar{q}}^{h^+}$

Average Jet Charge



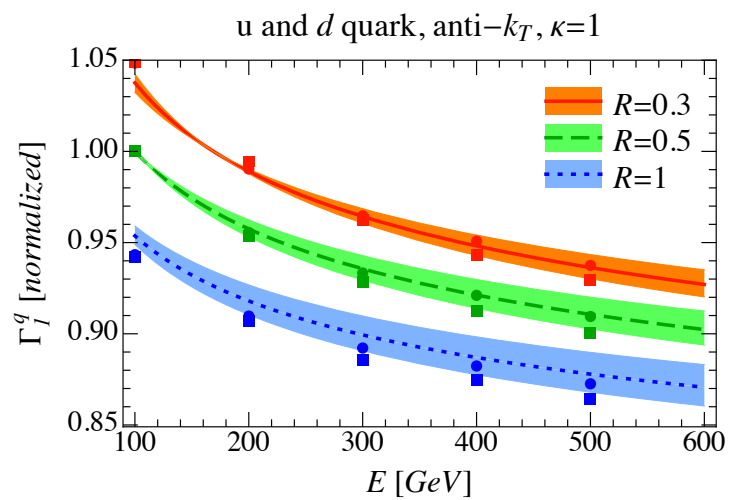
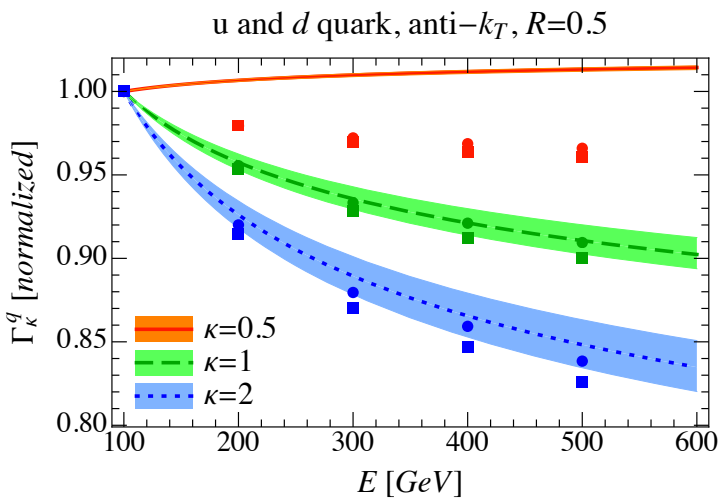
- ▶ PYTHIA results for d = squares, u = circles
- ▶ Normalizing removes dependence on nonperturbative input and flavor
- ▶ Uncertainty bands from varying μ by factors of 2
- ▶ Good agreement

Width of Quark Jet Charge



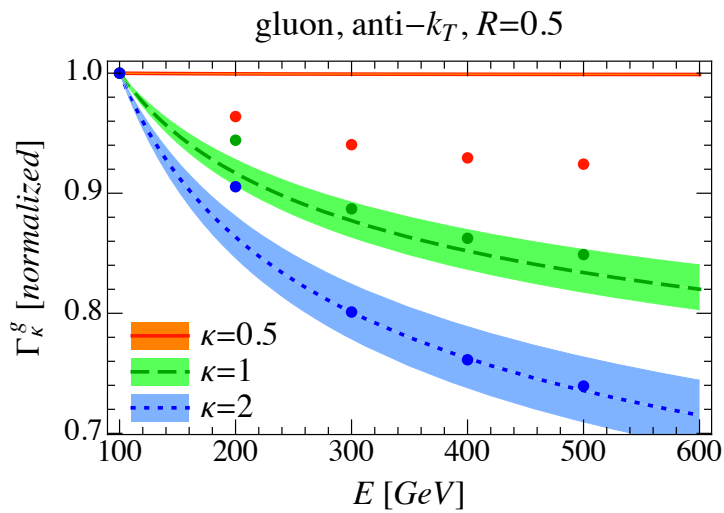
► Good agreement for $\kappa \gtrsim 1$

Width of Quark Jet Charge



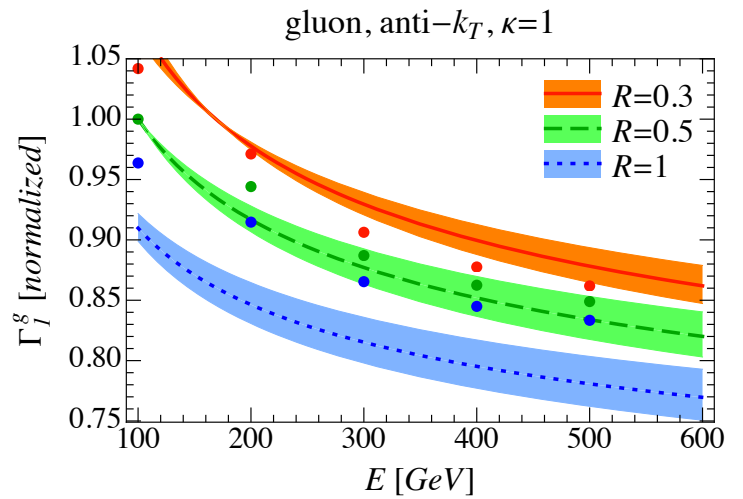
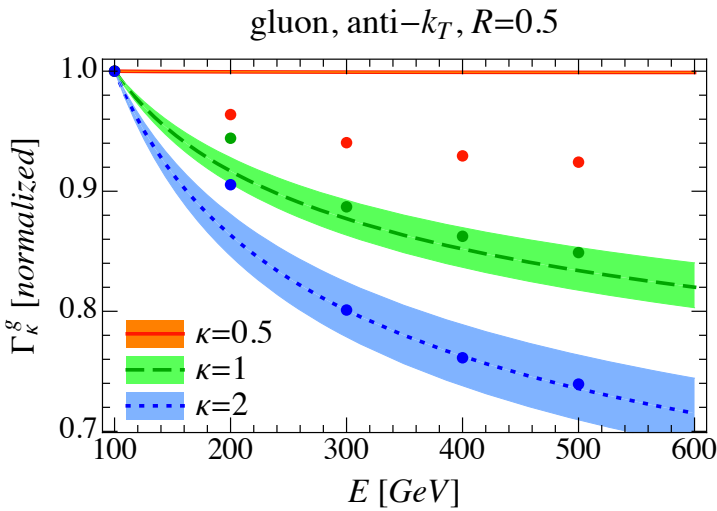
- ▶ Good agreement for $\kappa \gtrsim 1$
- ▶ Disagreement for $\kappa = 0.5$:
 - ▶ Rise caused by mixing with gluons
 - ▶ Large power corrections? (scale like $\sim \lambda^{2\kappa}$, absent for average)
 - ▶ Issue with PYTHIA? (rises when hadronization off)

Width of Gluon Jet Charge



- ▶ Good agreement for $\kappa \gtrsim 1$
- ▶ Disagreement for $\kappa = 0.5$ (same possible explanations)

Width of Gluon Jet Charge



- ▶ Good agreement for $\kappa \gtrsim 1$
- ▶ Disagreement for $\kappa = 0.5$ (same possible explanations)
- ▶ PYTHIA has smaller R -dependence than we predicted

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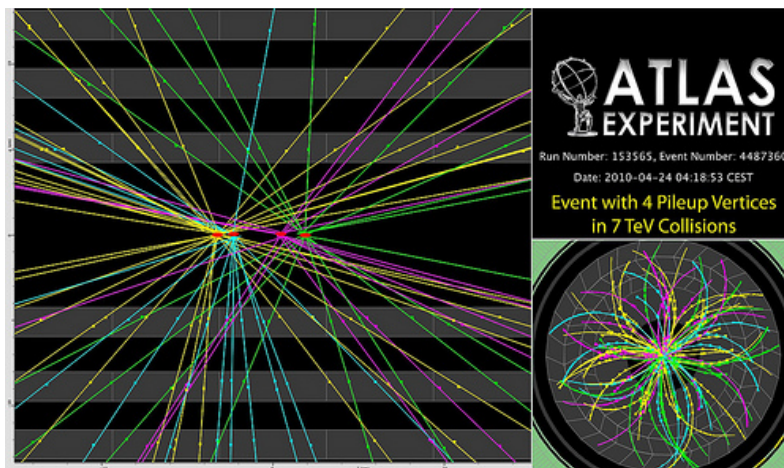
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Track-Based Observables

Tracks Rule!



Advantages:

- ▶ Better angular resolution
- ▶ Much better pointing → less sensitive to pile-up

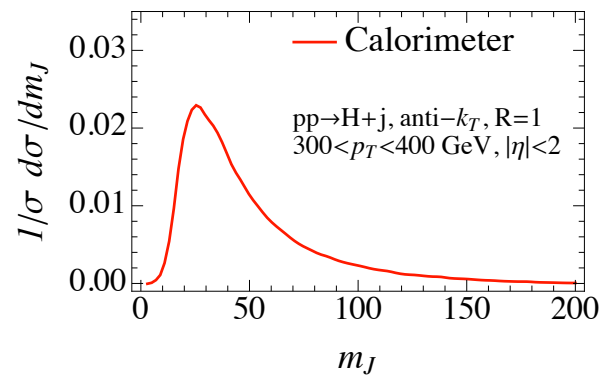
Disadvantages:

- ▶ Smearing of (resonance) peaks
- ▶ Not infrared safe

Track-Based Observables

- ▶ IR safe observable e :

$$\frac{d\sigma}{de} = \sum_N \int d\Pi_N \frac{d\sigma_N}{d\Pi_N} \delta[e - \hat{e}(\{p_i^\mu\})]$$



Track-Based Observables

- ▶ IR safe observable e :

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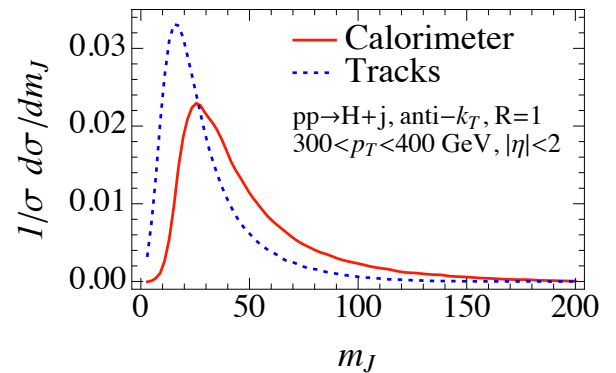
- ▶ Track-based version \bar{e} :

$$\frac{d\sigma}{d\bar{e}} = \sum_N \int d\Pi_N \frac{d\bar{\sigma}_N}{d\Pi_N} \underbrace{\int \prod_{i=1}^N dx_i T_i(x_i)}_{\text{hadronization}} \delta[\bar{e} - \hat{e}(\{\mathbf{x}_i p_i^\mu\})]$$

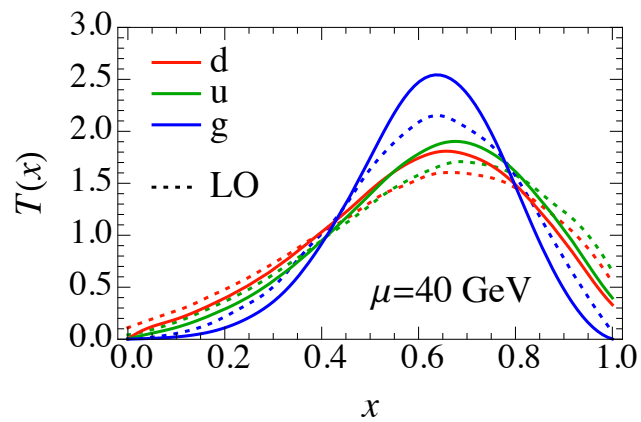
- ▶ Partons hadronize independently at leading power
- ▶ Momentum of charged particles

$$\bar{p}_i^\mu = \mathbf{x}_i p_i^\mu + \mathcal{O}(\Lambda_{\text{QCD}})$$

- ▶ Distribution in \mathbf{x}_i described by track function $T_i(\mathbf{x}_i)$

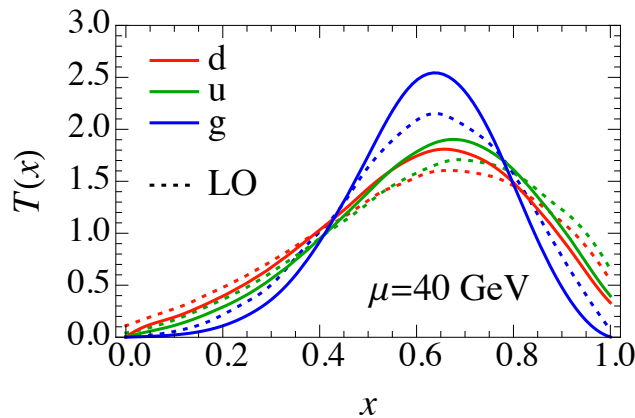


Track Function



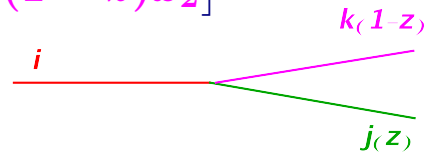
- ▶ Extracted using $x = \text{track}/\text{total jet energy}$ at $\mu = ER$

Track Function



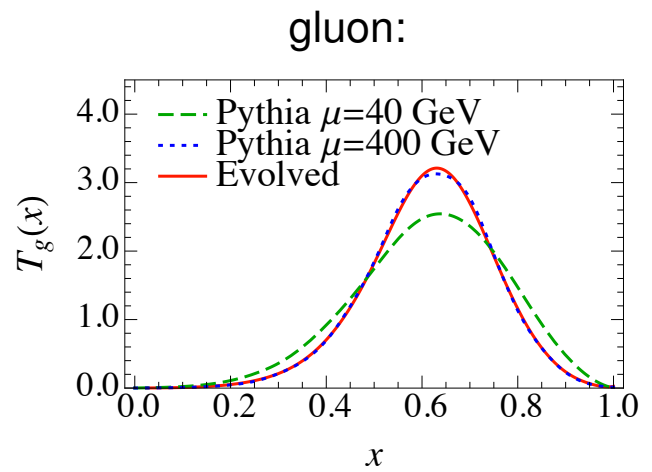
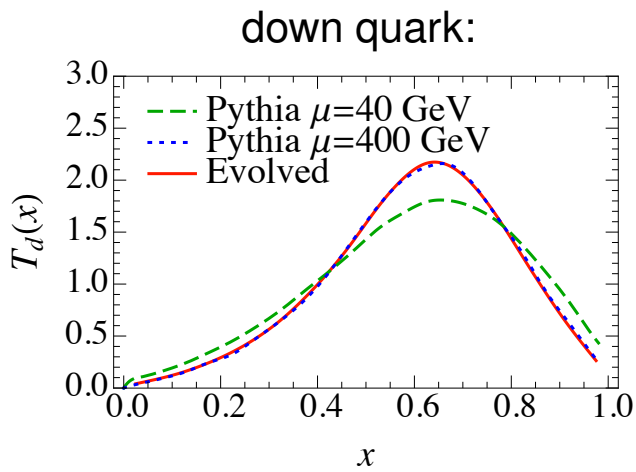
- ▶ Extracted using $x = \text{track/total jet energy}$ at $\mu = ER$
- ▶ NLO partonic track function:

$$T_{i,\text{bare}}^{(1)}(x) = \sum_j \int dz \left[\frac{\alpha_s(\mu)}{4\pi} \left(\frac{1}{\epsilon_{UV}} - \frac{1}{\epsilon_{IR}} \right) P_{ji}(z) \right] \int dx_1 T_j^{(0)}(x_1, \mu) \\ \times \int dx_2 T_k^{(0)}(x_2, \mu) \delta[x - zx_1 - (1-z)x_2]$$



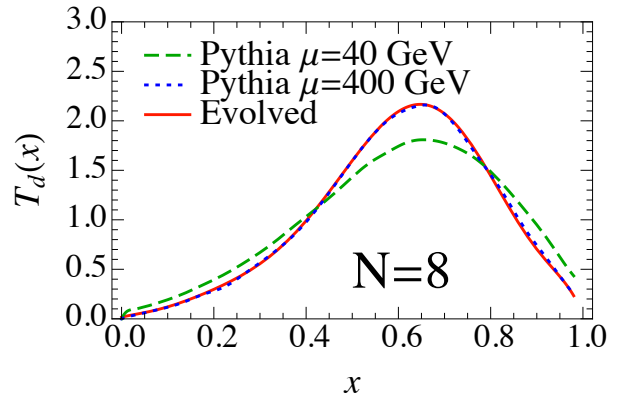
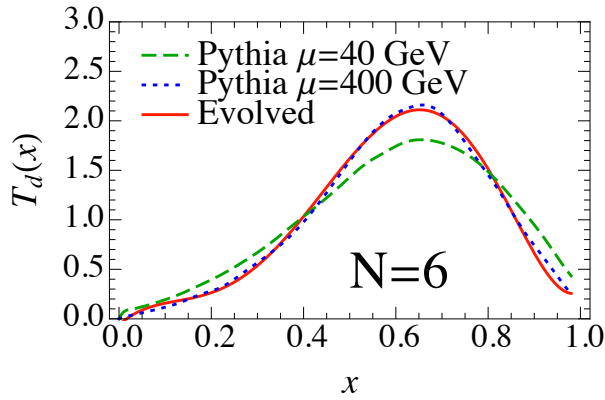
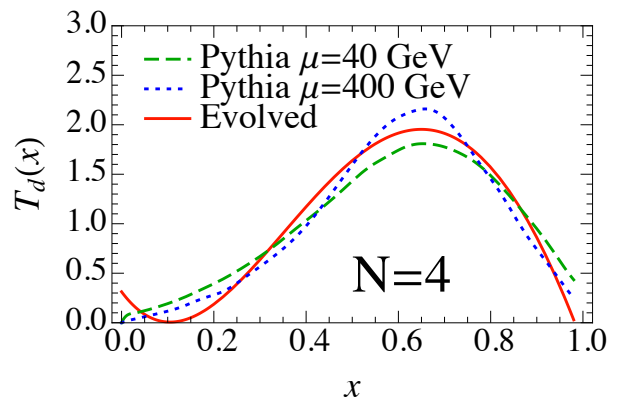
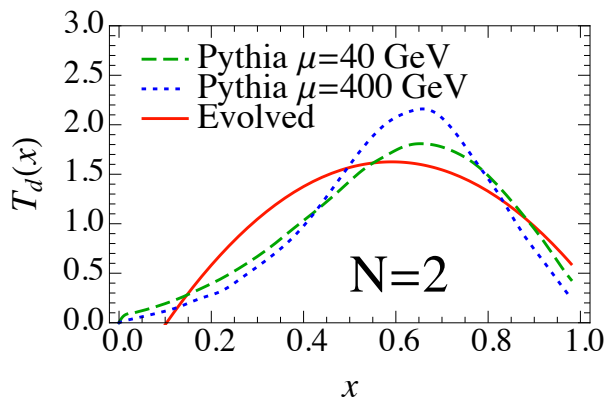
- ▶ IR divergence cancels in matching
- ▶ UV divergence gives parton shower evolution

Track Function Evolution



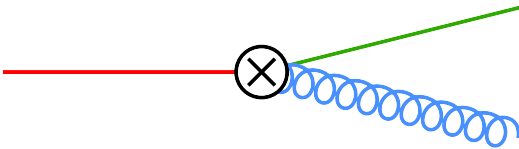
- ▶ Excellent agreement with PYTHIA
- ▶ Track functions for q and g similar at small μ

Comparison to Moment Space Evolution



- ▶ Evolving $N + 1$ moments and matching onto $\{1, x, \dots, x^N\}$

Energy Fraction w of Charged Particles in e^+e^-



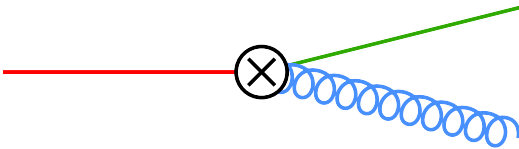
- ▶ Up to NLO:

$$\frac{d\sigma}{dw} = \int dy_1 dy_2 \frac{d\bar{\sigma}}{dy_1 dy_2} \int dx_1 T_q(x_1) \int dx_2 T_q(x_2) \int dx_3 T_g(x_3) \\ \times \delta\{w - [x_1 y_1 + x_2 y_2 + x_3(2 - y_1 - y_2)]/2\}$$

- ▶ Hard scattering:

$$\frac{d\sigma^{(0)}}{dy_1 dy_2} = \sigma^{(0)} \delta(1 - y_1) \delta(1 - y_2)$$

Energy Fraction w of Charged Particles in e^+e^-



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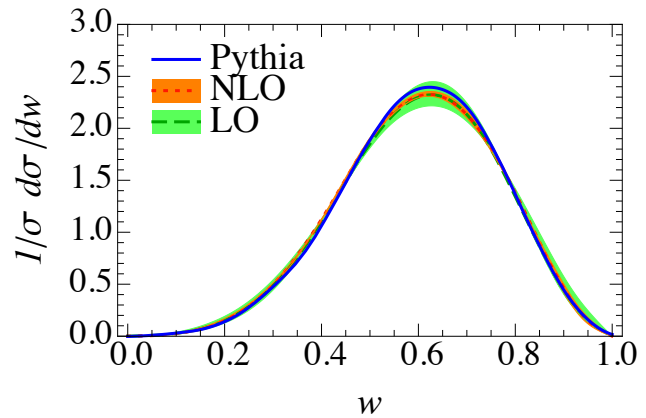
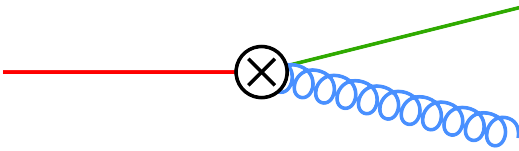
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- ▶ Hard scattering:

$$\frac{d\sigma^{(1)}}{d\mathbf{y}_1 d\mathbf{y}_2} = \sigma^{(0)} \frac{\alpha_s C_F}{2\pi} \left[-\frac{1}{\epsilon_{\text{IR}}} P_{qq}(y_1) \delta(1 - y_2) + \dots \right]$$

- ▶ IR divergence cancels against $T_q \rightarrow$ finite remainder is $\bar{\sigma}$

Energy Fraction w of Charged Particles in e^+e^-



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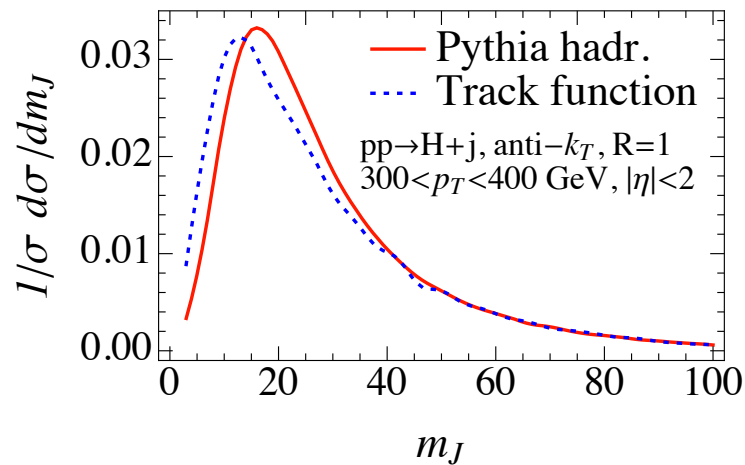
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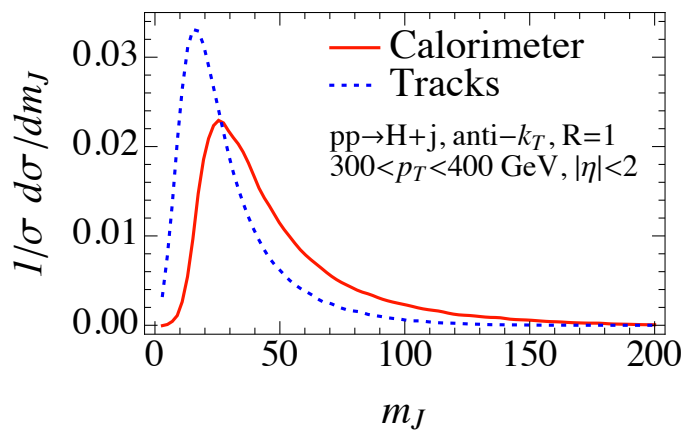
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Track Mass



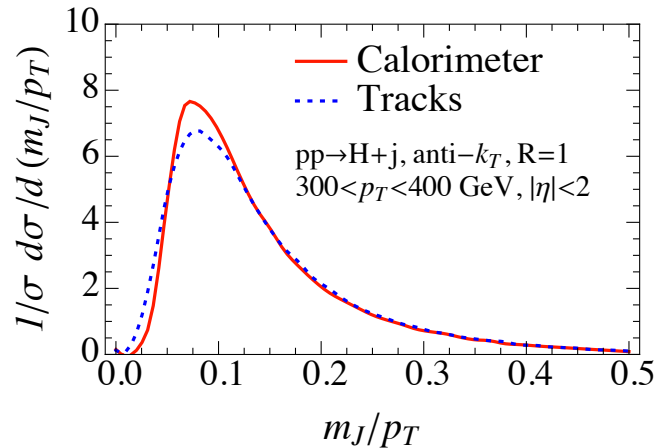
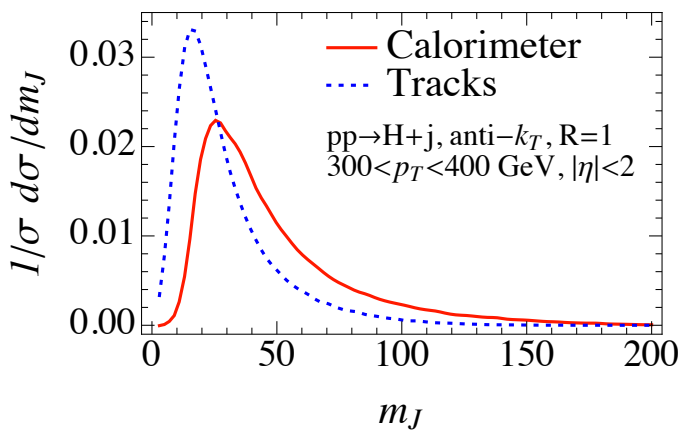
- ▶ Pythia with hadronization replaced by track functions:
Each emission in shower/resummation hadronizes independently
- ▶ Good agreement (difference is usual nonperturbative shift)
- ▶ Working on resummed track mass calculation

Ratio Observables



- ▶ Smearing from hadronization fluctuations (width of track function)

Ratio Observables



- ▶ Smearing from hadronization fluctuations (width of track function)
- ▶ Reduced smearing for ratio observables since fluctuations correlated
- ▶ E.g. ratios of energy correlation functions [Andrew's talk]
 N -subjettiness ratios [Thaler, Tilburg (2010)]

Conclusions



- ▶ Jet charge is old tool with new LHC applications
- ▶ Track-based observables less sensitive to pile-up, better angular resolution
- ▶ Partonic calculation not IR safe → hadronization crucial
- ▶ Computational framework:
 - ▶ Match on track function (charge distr.) with parton-shower evolution
 - ▶ Agrees with PYTHIA (except $\kappa = 0.5$)
 - ▶ Systematically improvable

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Thank You!