

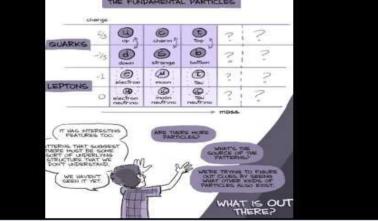
Unity of elementary particles and forces for the 3rd family^[1]

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Abstract:

- We propose a non-supersymmetric $SU(5) \times SU(3)'_C \times SU(2)'_L \times U(1)'_Y$ model in which only the 3rd family of fermions are unified in the SU(5) while the first two families belong to $SU(3)'_C \times SU(2)'_L \times U(1)'_Y$. We call it a top SU(5) model.
- The model remedies the non-unification of the three Standard Model couplings in non-supersymmetric SU(5).
- It also provides a mechanism of **baryon number violation** which is needed for the baryon asymmetry of the universe and not present in the Standard Model.
- Current experimental constraints on the leptoquark and diquark gauge bosons, mediating such baryon and lepton violating interactions, allow their masses to be at the TeV scale.
- These can be searched for as a (b tau) or (t t) resonance at the Large Hadron Collider. **Our Model and Formalism:**
- First two families of the SM fermions are charged under SM and singlet under the SU(5), while the third family is charged under SU(5) and singlet under SM.
- The SM gauge couplings are given by: $\frac{1}{g_j^2} = \frac{1}{g_5^2} + \frac{1}{(g_j)}, \frac{1}{g_Y^2} = \frac{1}{(g_5^Y)^2} + \frac{1}{(g_Y^Y)^2}$, with j = 1; 2; 3.
- Thus no unification of the SM couplings.
- $SU(5) \times SU(3)'_C \times SU(2)'_L \times U(1)'_Y$ is broken to the SM gauge symmetry $SU(3)_C \times SU(2)_L \times U(1)_Y$ via Higgs mechanism.
- Particle Content of the model :





Particles	Quantum Numbers	Γ
Q_i	$({f 1};{f 3},{f 2},{f 1/6})$	
U_i^c	$({f 1};{f ar 3},{f 1},-{f 2}/{f 3})$	
D_i^c	$({f 1};{f ar 3},{f 1},{f 1}/{f 3})$	
F_3	$({f 10};{f 1},{f 1},{f 0})$	
H	$({f 1};{f 1},{f 2},-{f 1}/{f 2})$	
U_T	$({f 5};{f ar 3},{f 1},{f 1}/{f 3})$	
XT	$({f 1};{f ar 3},{f 1},{f 1}/{f 3})$	
Xf	$({f 5};{f 1},{f 1},{f 0})$	
XD	$({f 1};{f 3},{f 1},-{f 1}/{f 3})$	
XL	(1; 1, 2, -1/2)	
Table: The complete particle conten		
under $SU(5) \times SU(3)'_C \times SU(2)'_L \times$		
	el. Here, $i = 1, 2$ and $k =$	
20 (0) mode	in the state of th	

The Higgs potential breaking gauge symmetry is given by:

$$\mathbf{V} = -\mathbf{m}_{T}^{2} \left| U_{T}^{2} \right| - \mathbf{m}_{D}^{2} \left| U_{D}^{2} \right| + \lambda_{T} \left| U_{T}^{2} \right|^{2} + \lambda_{D} \left| U_{D}^{2} \right|^{2} + \lambda_{TD} \left| U_{T}^{2} \right| \left| U_{D}^{2} \right| + \left[A_{T} \Phi U_{T} X T^{\dagger} + A_{D} \Phi U_{D} H^{\dagger} + \frac{y_{TD}}{M_{*}} U_{T}^{3} U_{D}^{2} + H.C \right]$$

Where

$$\Theta, \qquad \langle \boldsymbol{U}_D \rangle = \boldsymbol{v}_D \begin{pmatrix} \boldsymbol{0}_{3\times 2} \\ \boldsymbol{I}_{2\times 2} \end{pmatrix}, \langle \boldsymbol{U}_D \rangle = \boldsymbol{v}_D \begin{pmatrix} \boldsymbol{0}_{3\times 2} \\ \boldsymbol{I}_{2\times 2} \end{pmatrix}$$

We assume that VT and VD are in the TeV scale. The masses of the gauge bosons are,

$$\sum_{i=T,D} \left\langle (D_{\mu}U_{i})^{\dagger} D^{\mu}U_{i} \right\rangle = \frac{1}{2} v_{T}^{2} (g_{5}\hat{A}_{\mu}^{a3} - g_{3}^{\dagger}\tilde{A}_{\mu}^{a3})^{2} + \frac{1}{2} v_{D}^{2} (g_{5}\hat{A}_{\mu}^{a2} - g_{2}^{\dagger}\tilde{A}_{\mu}^{a2})^{2} + (\frac{v_{T}^{2}}{3} + \frac{v_{D}^{2}}{2}) (g_{5}^{Y}\hat{A}_{\mu}^{a1} - g_{Y}^{\dagger}\tilde{A}_{\mu}^{a1})^{2} + \frac{1}{2} g_{5}^{2} (g_{5}\hat{A}_{\mu}^{a2} - g_{2}^{\dagger}\tilde{A}_{\mu}^{a2})^{2} + (\frac{v_{T}^{2}}{3} + \frac{v_{D}^{2}}{2}) (g_{5}^{Y}\hat{A}_{\mu}^{a1} - g_{Y}^{\dagger}\tilde{A}_{\mu}^{a1})^{2} + \frac{1}{2} g_{5}^{2} (g_{5}\hat{A}_{\mu}^{a2} - g_{2}^{\dagger}\tilde{A}_{\mu}^{a2})^{2} + (\frac{v_{T}^{2}}{3} + \frac{v_{D}^{2}}{2}) (g_{5}^{Y}\hat{A}_{\mu}^{a1} - g_{Y}^{\dagger}\tilde{A}_{\mu}^{a1})^{2} + \frac{1}{2} g_{5}^{2} (g_{5}\hat{A}_{\mu}^{a2} - g_{2}^{\dagger}\tilde{A}_{\mu}^{a2})^{2} + (g_{5}\hat{A}_{\mu}^{a2}$$

Yukawa Coupling:

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$$-L = y_{ij}^{u}U_{i}^{c}Q_{j}\tilde{H} + y_{kj}^{v}N_{k}^{c}L_{j}\tilde{H} + y_{ij}^{d}D_{i}^{c}Q_{j}H + y_{ij}^{e}E_{i}^{c}L_{j}H + y_{33}^{u}F_{3}F_{3}\Phi^{\dagger} + y_{33}^{d,e}F_{3}\bar{f}_{3}\Phi + y_{k3}^{v}N_{k}^{c}\bar{f}_{3}\Phi^{\dagger} + m_{h}^{c}D_{h}^{c}Q_{j}H + y_{ij}^{e}E_{i}^{c}L_{j}H + y_{33}^{u}F_{3}F_{3}\Phi^{\dagger} + y_{33}^{d,e}F_{3}\bar{f}_{3}\Phi + y_{k3}^{v}N_{k}^{c}\bar{f}_{3}\Phi^{\dagger} + m_{h}^{c}D_{h}^{c}Q_{j}H + y_{ij}^{e}E_{i}^{c}L_{j}H + y_{33}^{u}F_{3}F_{3}\Phi^{\dagger} + y_{33}^{d,e}F_{3}\bar{f}_{3}\Phi + y_{k3}^{v}N_{k}^{c}\bar{f}_{3}\Phi^{\dagger} + m_{h}^{c}D_{h}^{c}Q_{j}H + y_{ij}^{e}E_{i}^{c}L_{j}H + y_{ij}^{u}F_{3}F_{3}\Phi^{\dagger} + y_{ij}^{d,e}F_{3}\bar{f}_{3}\Phi + y_{k3}^{v}N_{k}^{c}\bar{f}_{3}\Phi^{\dagger} + m_{h}^{c}D_$$

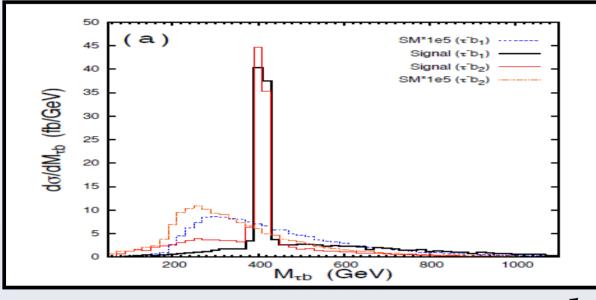
- Choose a basis in which up quark mass matrix is diagonal.. So no mixing between u, c, t.
- CKM mixing arises from purely down quark sector. The mixing between the first two families and the 3rd family is done by using dimension 5 operators generated by the fields in the model via renormalizable interaction.
- CKM Mixings: The dimension 5 interactions are generated at the renormalizable level by using the vector-like fermions (Xf, XD, XL) with masses ~1; 000TeV. This gives M* to be 1,000 TeV.

$$L = \frac{1}{M_*} (y_{i3}^d D_i^c F_3 \Phi U_T^{\dagger} + y_{i3}^e E_i^c \bar{f}_3 H U_D + y_{3i}^d \bar{f}_3 Q_i H U_T + y_{3i}^e F_3 L_i \Phi U_D^{\dagger}) + H.C$$

Correct CKM mixing is generated by using M* ~ 1,000 TeV.

Phenomenological Implications:

- Leptoquark gauge bosons X and Y can be pair produced at the LHC via QCD strong interactions. $gg \to X\overline{X}, Y\overline{Y}, \quad q\overline{q} \to X\overline{X}, Y\overline{Y}$
- X and Y subsequently decays to:
- $X \to \overline{b} \tau^+, tt; \overline{X} \to b \tau^-, \overline{t}\overline{t}$. $Y \to \overline{b} \nu_\tau, Y \to \overline{t} \tau^+, tb$.
- Consider signal for production : $\Rightarrow (b \tau^+)(b \tau^-)$
- Both b and tau can be tagged, so the resonance X in the b tau mode can be reconstructed in Fig.
- Dominant SM background : $pp \rightarrow 2b2\tau; 4b; 2j2b; 4j; t\bar{t}$, which can be easily eliminated using suitable cuts.



<u>X resonance in the $b\tau$ mode at 7 TeV LHC</u>

- Invariant mass distribution for the b channel for $M_x = 400$ GeV and $M_x = 600$ GeV.
- Cuts used: $p_T > 80 \text{GeV}; \eta < 2:5; \Delta R > 0:2$ Also used efficiency for **b** and τ tagging to be 0.5, mistag rate for light quark 1%, for charmed quark 10%.
- Background has been multiplied by 10⁵ in (a) and 10⁴ in (b).

Leptoquark Signal at LHC:

- To prove baryon and lepton violation, X need to be reconstructed also in the (t t) mode in addition to the $(b\tau)$ mode to show that it is a leptoquark as well as a diquark and hence baryon and lepton number violating.
- For heavier leptoquark, one can look for the final state $bb \tau \tau$ signal.
- For 7 TeV LHC, our model gives 5 sigma reach for mass as high as 750 GeV with 5fb⁻¹ luminosity.

Conclusions:

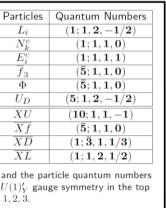
- Presented TeV scale model for quark lepton unification in 4 Dimensions.
- Has leptoquark gauge bosons X and Y coupling only to the 3rd family of fermions, and hence produces B and L violation only involving the 3rd family.
- Can be observed as $(b\tau)$ and (t t) resonance at the LHC at the 7 TeV LHC, with 5fb⁻¹ luminosity, the discovery reach is 750 GeV.

Further Works on this model:

- Motivated by our model, CMS collaboration has searched for third generation leptoquarks with integrated luminosity of 4.8 fb⁻¹ as a resonance in the $(b\tau)$ mode and has set a limit of exclusion of SU(5) vector leptoquarks with masses below 760 GeV at 95% CL.
- *Further theoretical work for X,Y on the other modes at 8TeV and 14TeV LHC are in progress.[3]

References:

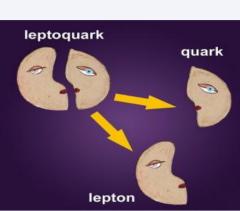
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- [2] CMS Physics Analysis Summary: CMS PAS EXO-12-002.
- [3] S.Chakdar, T.Li, S.Nandi and S.K.Rai (in preparation).
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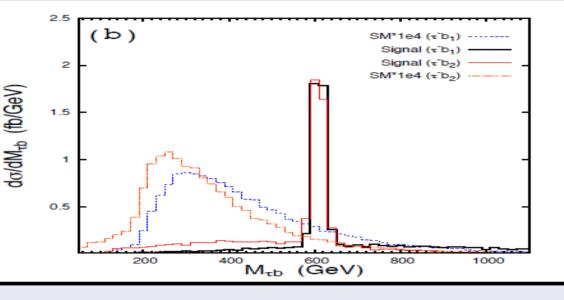


 $v_T^2 + v_D^2)(X_\mu \overline{X}_\mu + Y_\mu \overline{Y}_\mu).$

$P_{kl}^N N_k^c N_l^c + H.C$









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