



Silicon detectors in HEP

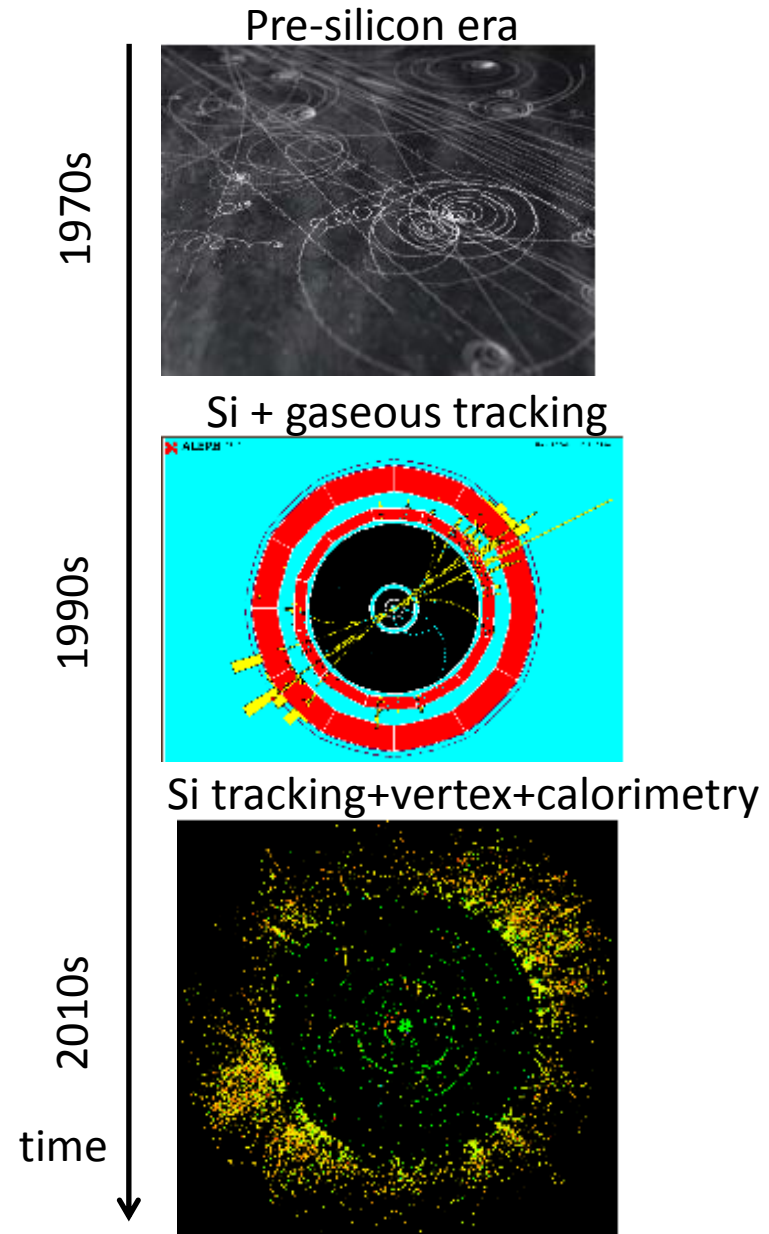
Enver Alagoz
ACCTR Meeting
18 August 2012

Introduction to silicon

- Semiconductor material
 - Used for radiation detection/imaging in
 - Particle and nuclear physics experiments
 - Space and ground-based telescopes
 - Medical applications
 - Industrial applications
 - ...
- Interacts with radiation
 - Radiation creates charge carriers
 - Electric field drifts carriers toward readout electronics
 - Signal integrated in front-end electronics
 - Digitize integrated signal
 - Readout and store digital signal

Particle reconstruction in time

- Why detectors change
 - Interested in rare events
 - in very high energies
 - with very small uncertainties
 - with very high background
- Need high precision
 - Very high granularity
 - low material budget
- High particle collision rates
 - High speed electronics
 - Good background tolerance
 - Long term reliability
 - Radiation hardness

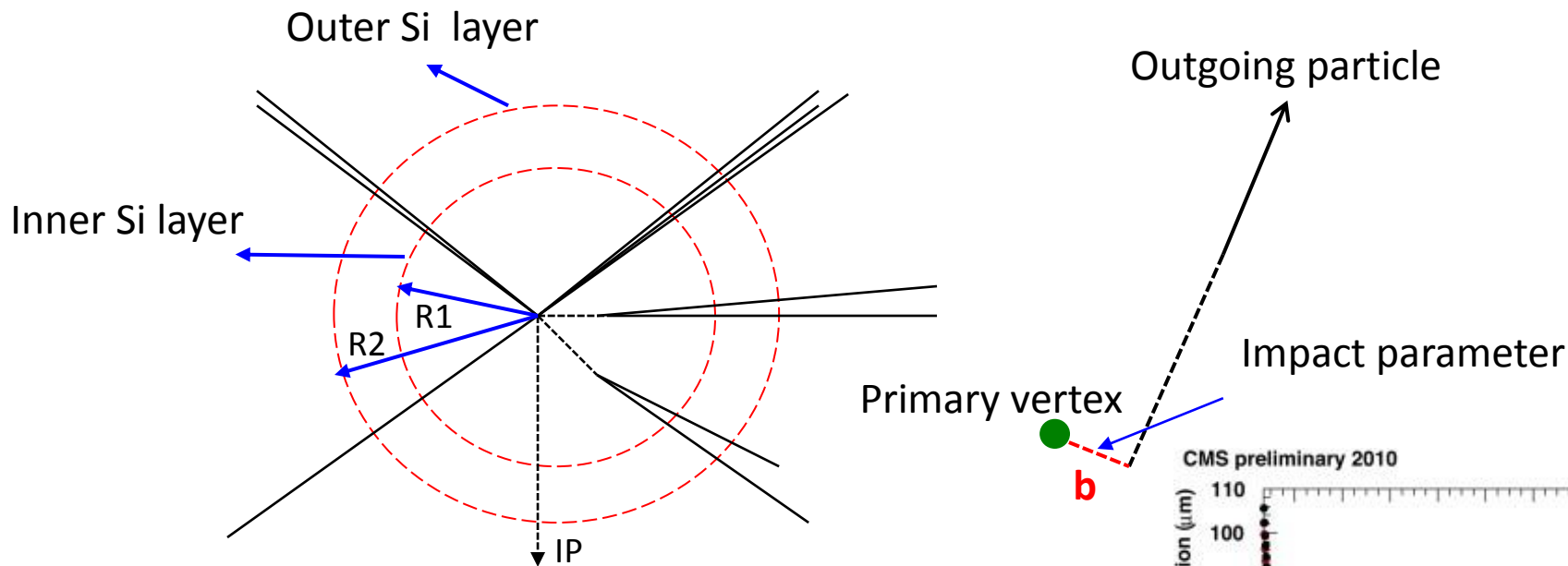


Gaseous vs silicon detectors

- In silicon detector
 - No multiplication → signal amplification needed
 - Optimum thickness to minimize multiple scattering

Property	Gas	Silicon	Importance
Density	Low	2.33 g/cm ³	Denser -> better spatial resolution
Charge	electrons and ions	electrons and holes	
Ionization energy	30eV per e- ion pair	3.6 eV per e- hole pair	Lower -> better energy resolution
Mobility	10 ns to 10 μs	few ns to 20 ns	Faster -> no dead time
Signal-to-noise	Low	High	Reliable signal

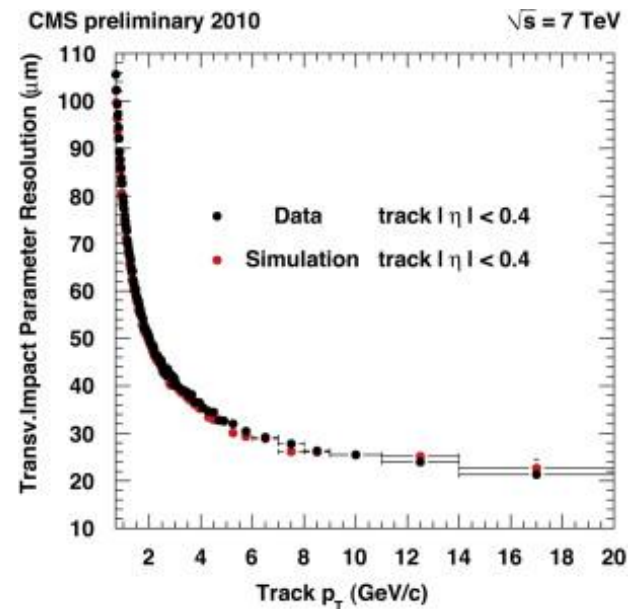
Vertex reconstruction



$$\sigma_b^2 \approx \frac{1}{(R_2 - R_1)^2} [(\sigma_1 R_2)^2 + (\sigma_2 R_1)^2]$$

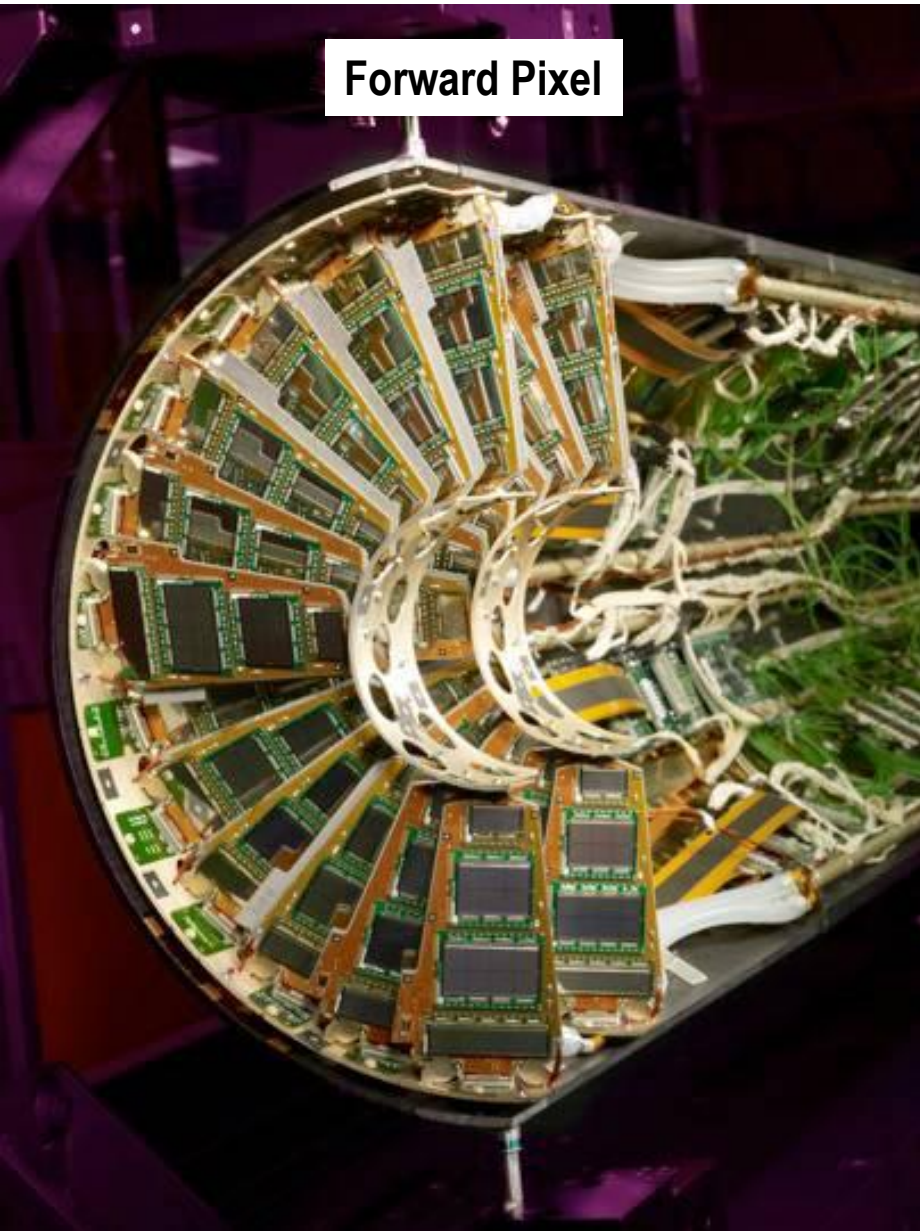
If both layers have same spatial resolution (σ)

$$\left(\frac{\sigma_b}{\sigma}\right)^2 \approx \left(\frac{1}{1 - R_1/R_2}\right)^2 + \left(\frac{1}{R_2/R_1 - 1}\right)^2$$

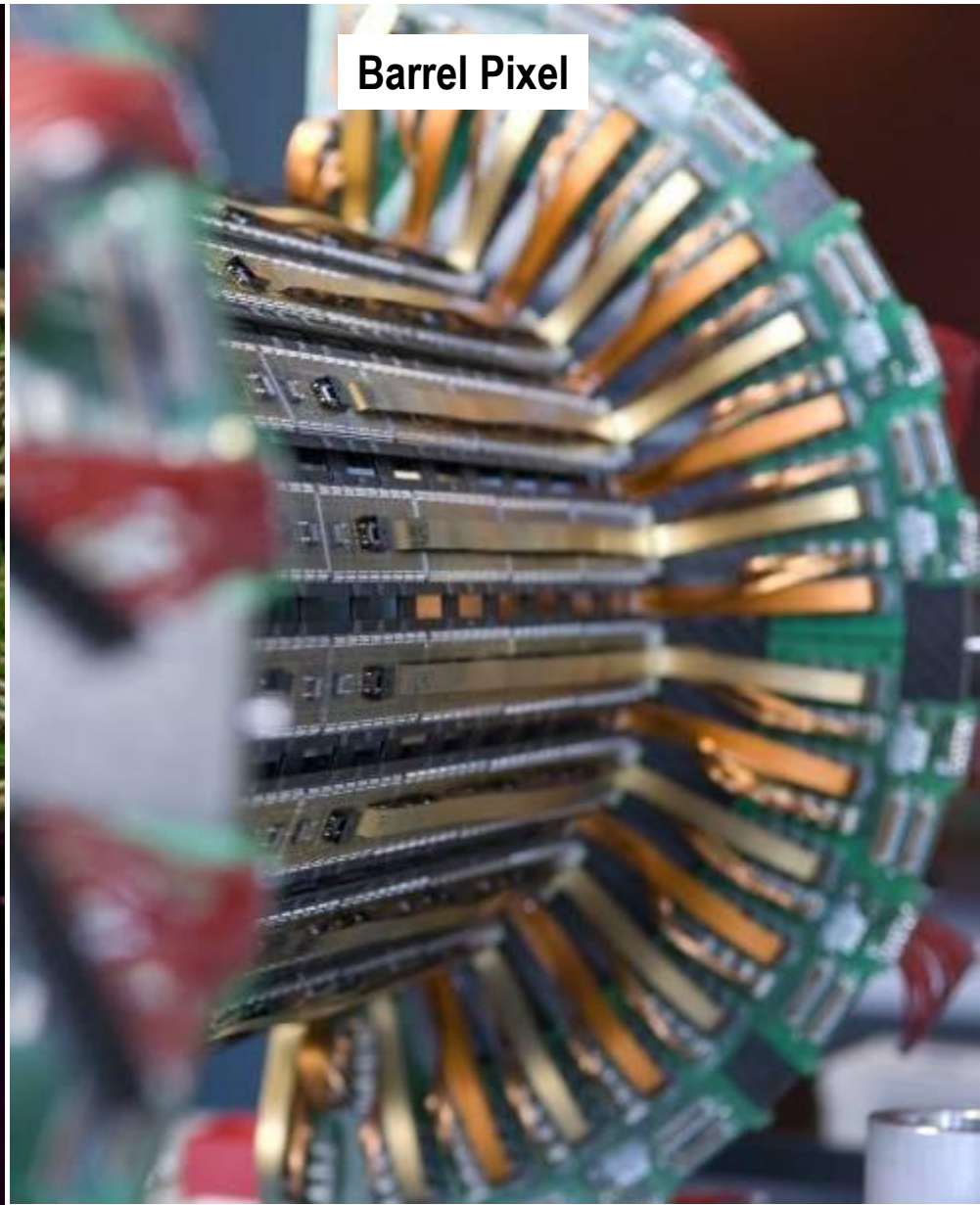


CMS pixel detector

Forward Pixel



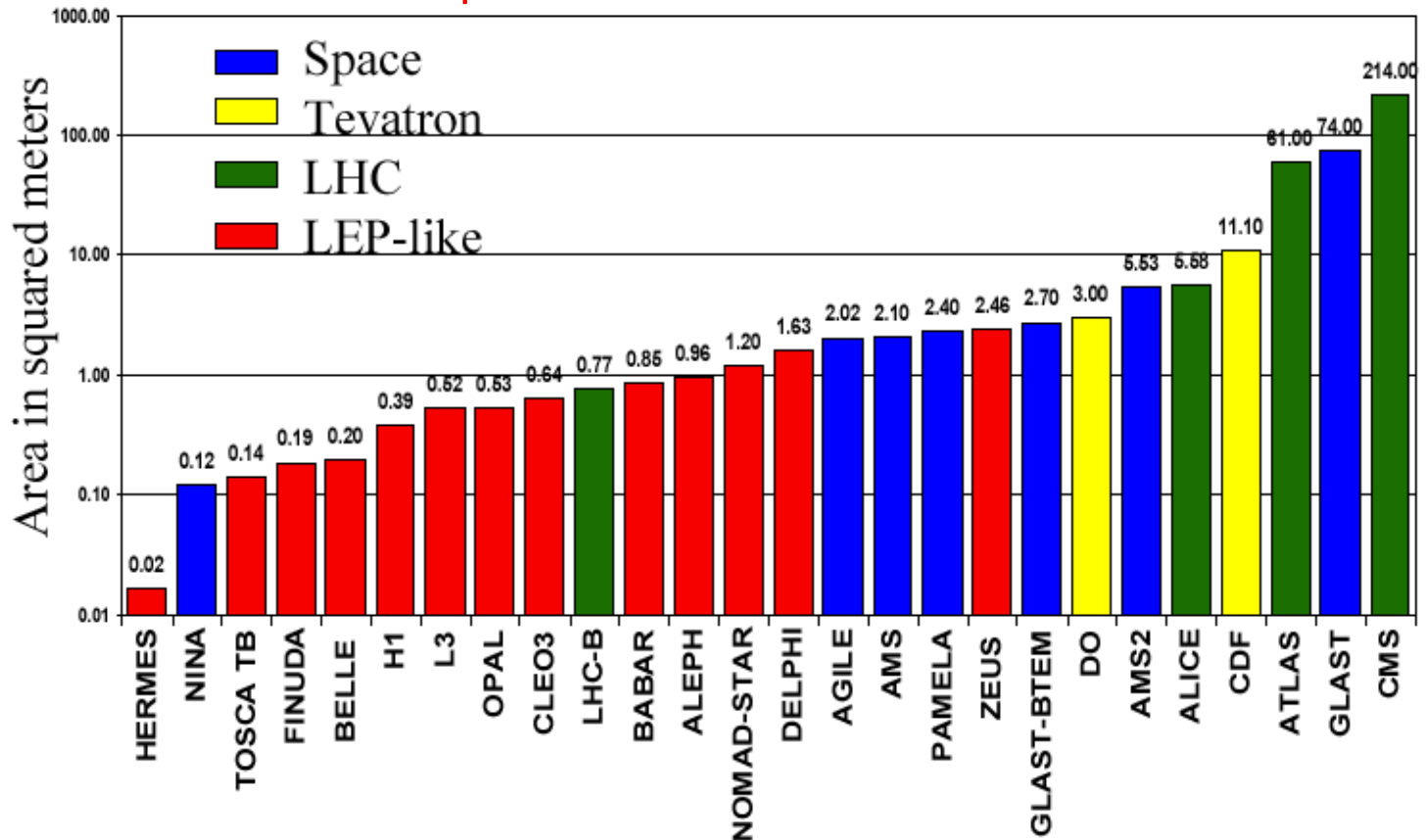
Barrel Pixel



Silicon detectors in HEP

- In HEP experiment
 - Vertex reconstruction
 - Fast detection and position resolution

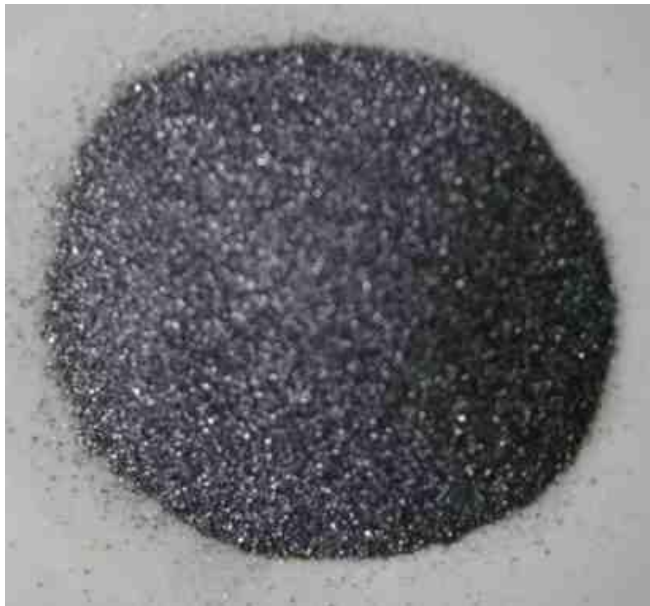
Experiments with silicon detector



E. Do Couto e Silva, Vertex 2000, Sept 10-15, National Lakeshore, MI, USA

Why silicon?

- 2nd most abundant element in Earth (28% by mass)
- Easy to process and purify to 0.001 ppb
- Natural SiO₂ as insulation during fabrication
- Found in form of silicate minerals and SO₂ (silica – sand)
- Discovered in 1824 by Swedish chemist J.J. Berzelius
- Commercially utilized since 1824



Silicon powder

Silicon crystal



Spectral lines of Silicon

Silicon - $_{14}\text{Si}^{28}$

PERIODIC TABLE OF THE ELEMENTS

<http://www.ktf-split.hr/periodni/en/>

PERIOD	GROUP IUPAC		GROUP CAS										GROUP IUPAC					
	1 IA	2 IIA	3 IIIB	4 IVB	5 VB	6 VIB	7 VIIB	8 VIII B	9 VIII B	10 VIII B	11 IB	12 IIB	13 IIIA	14 IVA	15 VA	16 VIA	17 VIIA	18 VIIIA
1	1 1.0079 H HYDROGEN																	2 4.0026 He HELIUM
2	3 6.941 Li LITHIUM	4 9.0122 Be BERYLLIUM			5 10.811 B BORON								6 12.011 C CARBON	7 14.007 N NITROGEN	8 15.999 O OXYGEN	9 18.998 F FLUORINE	10 20.180 Ne NEON	
3	11 22.990 Na SODIUM	12 24.305 Mg MAGNESIUM										13 26.982 Al ALUMINIUM	14 28.086 Si SILICON	15 30.974 P PHOSPHORUS	16 32.065 S SULPHUR	17 35.453 Cl CHLORINE	18 39.948 Ar ARGON	
4	19 39.098 K POTASSIUM	20 40.078 Ca CALCIUM	21 44.956 Sc SCANDIUM	22 47.867 Ti TITANIUM	23 50.942 V VANADIUM	24 51.996 Cr CHROMIUM	25 54.938 Mn MANGANESE	26 55.845 Fe IRON	27 58.933 Co COBALT	28 58.693 Ni NICKEL	29 63.546 Cu COPPER	30 65.39 Zn ZINC	31 69.723 Ga GALLIUM	32 72.64 Ge GERMANIUM	33 74.922 As ARSENIC	34 78.96 Se SELENIUM	35 79.904 Br BROMINE	36 83.80 Kr KRYPTON
5	37 85.468 Rb RUBIDIUM	38 87.62 Sr STRONTIUM	39 88.906 Y YTTORIUM	40 91.224 Zr ZIRCONIUM	41 92.906 Nb NIOBIUM	42 95.94 Mo MOLYBDENUM	43 (98) Tc TECHNETIUM	44 101.07 Ru RUTHENIUM	45 102.91 Rh RHODIUM	46 106.42 Pd PALLADIUM	47 107.87 Ag SILVER	48 112.41 Cd CADMIUM	49 114.82 In INDIUM	50 118.71 Sn TIN	51 121.76 Sb ANTIMONY	52 127.60 Te TELLURIUM	53 126.90 I IODINE	54 131.29 Xe XENON
6	55 132.91 Cs CAESIUM	56 137.33 Ba BARIUM	57-71 La-Lu Lanthanide	72 178.49 Hf HAFNIUM	73 180.95 Ta TANTALUM	74 183.84 W TUNGSTEN	75 186.21 Re RHENIUM	76 190.23 Os OSMIUM	77 192.22 Ir IRIDIUM	78 195.08 Pt PLATINUM	79 196.97 Au GOLD	80 200.59 Hg MERCURY	81 204.38 Tl THALLIUM	82 207.2 Pb LEAD	83 208.98 Bi BISMUTH	84 (209) Po POLONIUM	85 (210) At ASTATINE	86 (222) Rn RADON
7	87 (223) Fr FRANCIUM	88 (226) Ra RADIUM	89-103 Ac-Lr Actinide	104 (261) Rf RUTHERFORDIUM	105 (262) Db DUBNIUM	106 (266) Sg SEABORGIUM	107 (264) Bh BOHRNIUM	108 (277) Hs HASSIUM	109 (268) Mt MEITNERIUM	110 (281) Uun UNUNNIUM	111 (272) Uuu UNUNUNIUM	112 (285) Uub UNUNBIUM		114 (289) Uuq UNUNQUADIUM				

LANTHANIDE

Copyright © 1998-2003 EniG. (eni@ktf-split.hr)

(1) Pure Appl. Chem., 73, No. 4, 667-683 (2001)

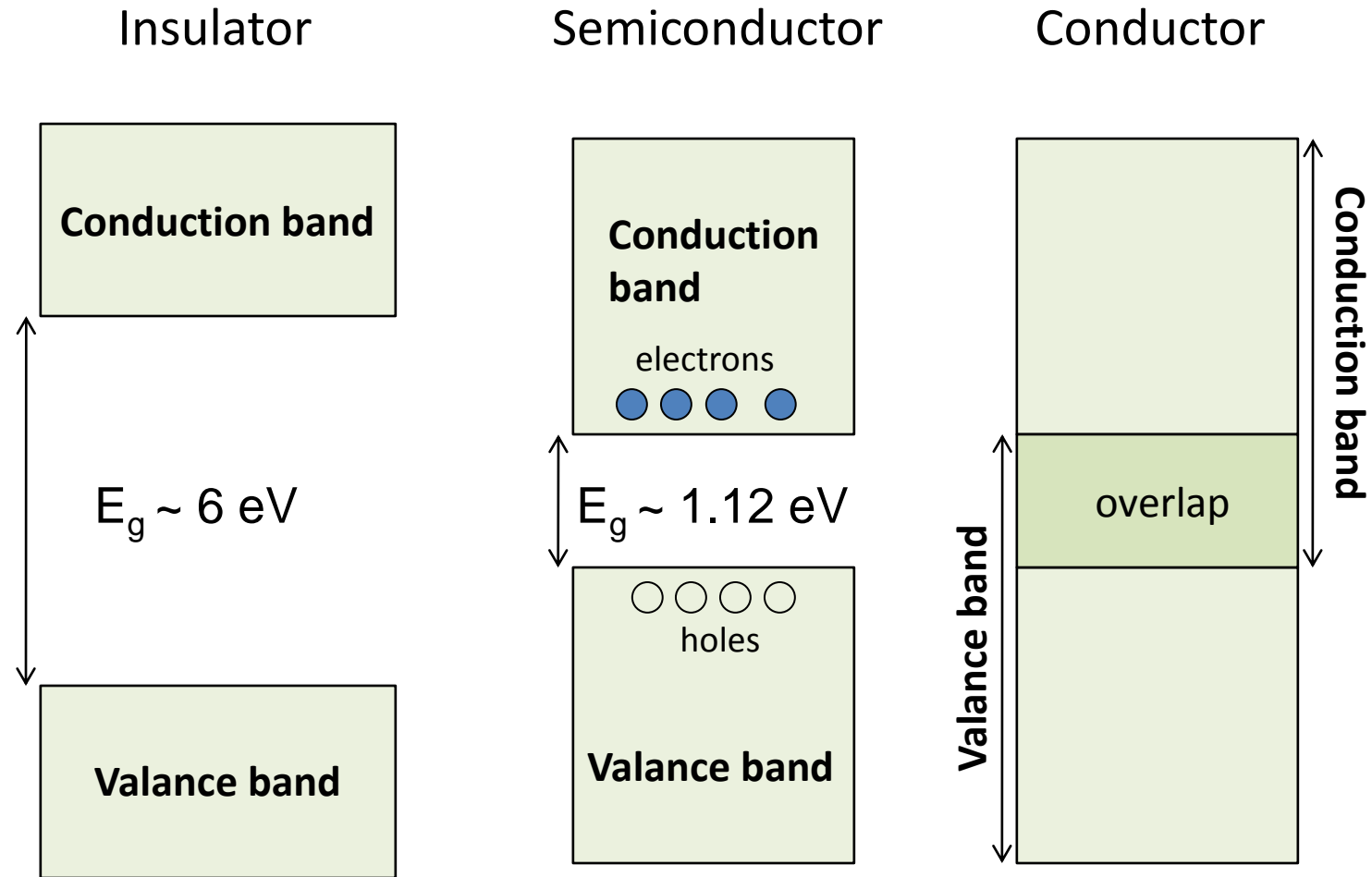
Relative atomic mass is shown with five significant figures. For elements having no stable nuclides, the value enclosed in brackets indicates the mass number of the longest-lived isotope of the element.

However three such elements (Th, Pa, and U) do have a characteristic terrestrial isotopic composition, and for these an atomic weight is tabulated.

ACTINIDE

57 138.91 La LANTHANUM	58 140.12 Ce CERIUM	59 140.91 Pr PRASEODYMIUM	60 144.24 Nd NEODYMIUM	61 (145) Pm PROMETHIUM	62 150.36 Sm SAMARIUM	63 151.96 Eu EUROPIUM	64 157.25 Gd GADOLINIUM	65 158.93 Tb TERBIUM	66 162.50 Dy DYSPROSIUM	67 164.93 Ho HOLMIUM	68 167.26 Er ERBIUM	69 168.93 Tm THULIUM	70 173.04 Yb YTTERIUM	71 174.97 Lu LUTETIUM
89 (227) Ac ACTINIUM	90 232.04 Th THORIUM	91 231.04 Pa PROTACTINIUM	92 238.03 U URANIUM	93 (237) Np NEPTUNIUM	94 (244) Pu PLUTONIUM	95 (243) Am AMERICIUM	96 (247) Cm CURIUM	97 (247) Bk BERKELIUM	98 (251) Cf CALIFORNIUM	99 (252) Es EINSTEINIUM	100 (257) Fm FERMIUM	101 (258) Md MENDELEVIUM	102 (259) No NOBELIUM	103 (262) Lr LAWRENCIUM

Basic properties – Band gap

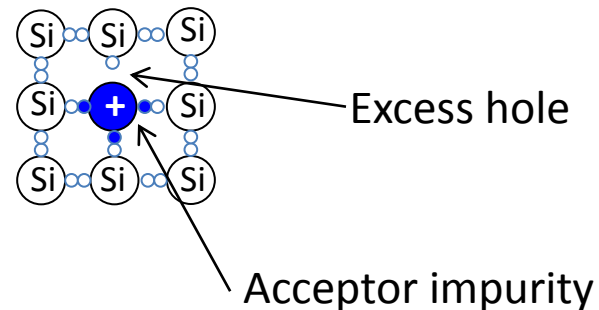
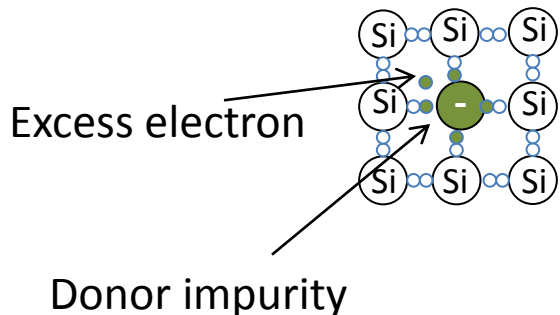
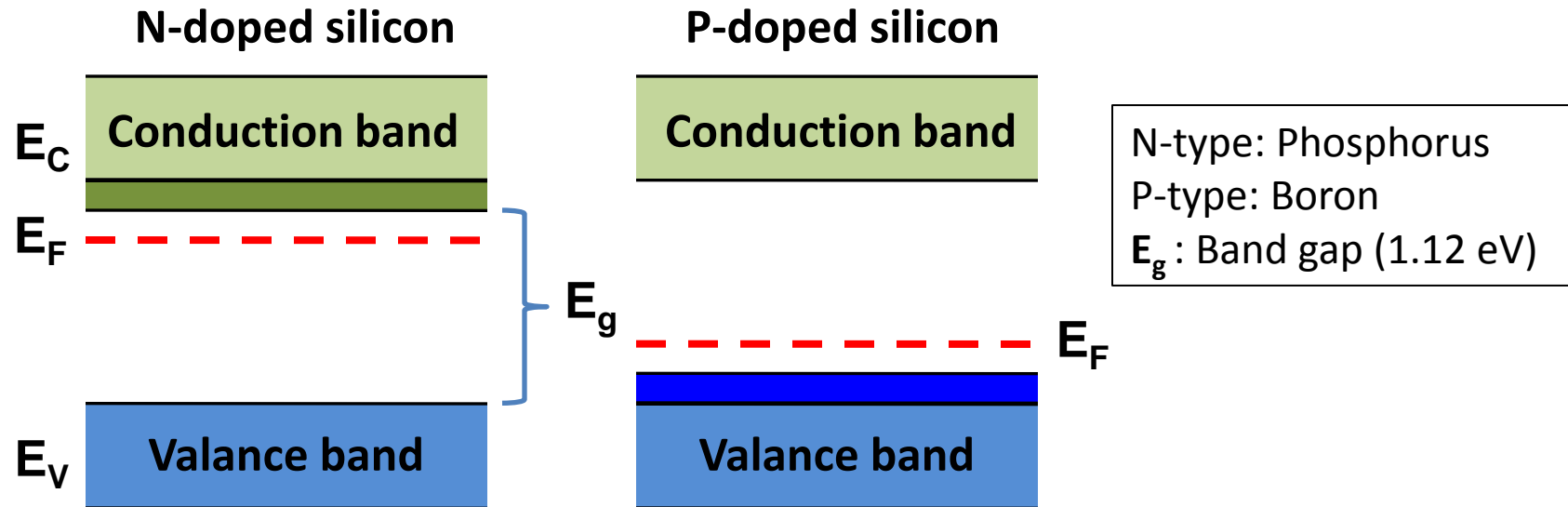


Basic properties

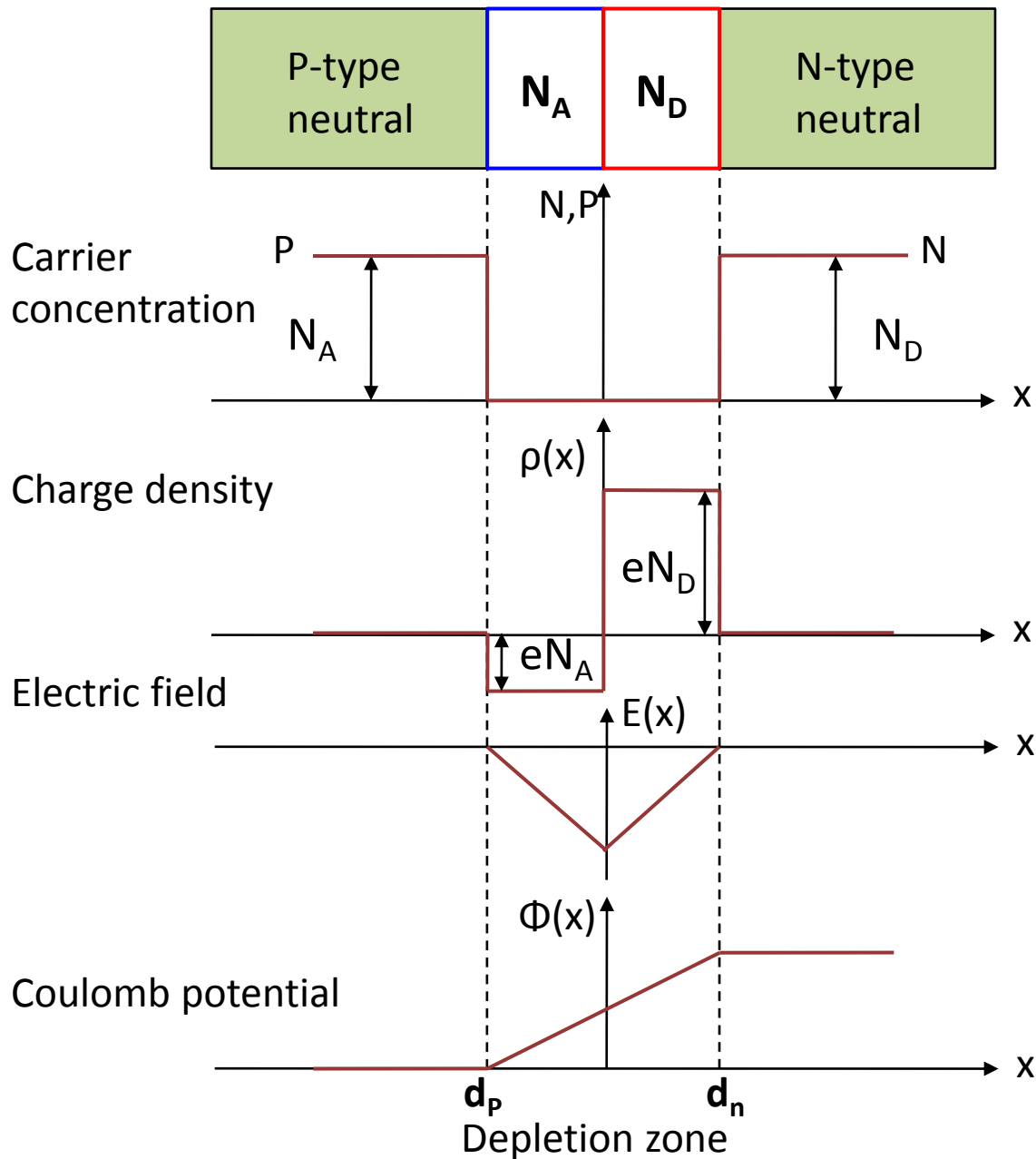
- Very pure material
- Charge carriers are created by
 - thermal, optical, and other excitations or ionization
- Four valance electrons (covalent bonds)
- Silicon (Si) and Germanium (Ge) are common
- Doped with extrinsic semiconductors
 - N-type: excess electrons (e- donor), i.e. P, As etc.
 - P-type: excess holes (e- acceptor), i.e. Al, B, Ga, etc.

Silicon PN-junction

- Silicon detectors are in basic form of PN-junction

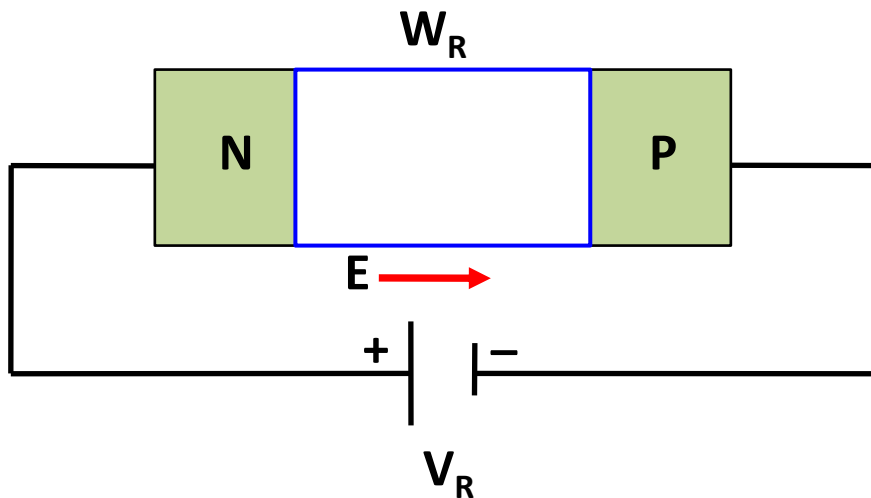
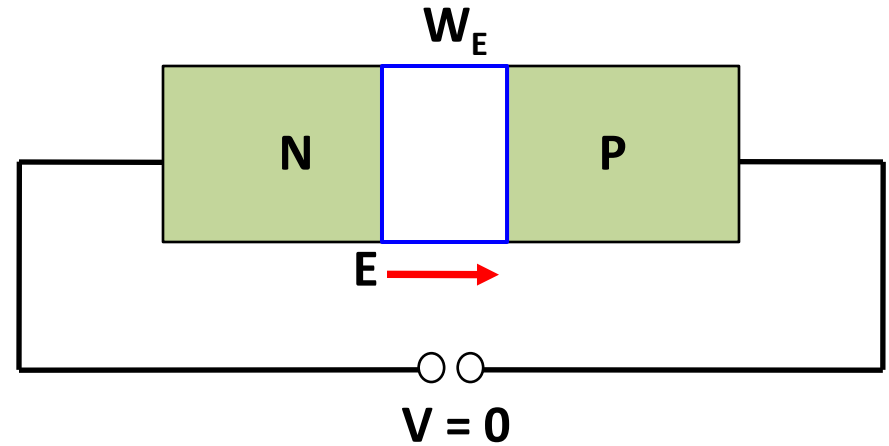
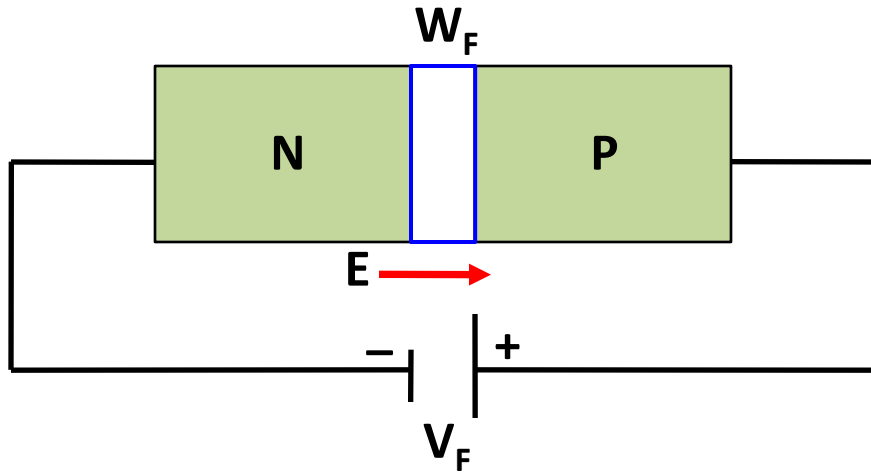


Silicon pn-junction formation



N_A : # of acceptor impurities
 N_D : # of donor impurities

PN-junction bias schemes



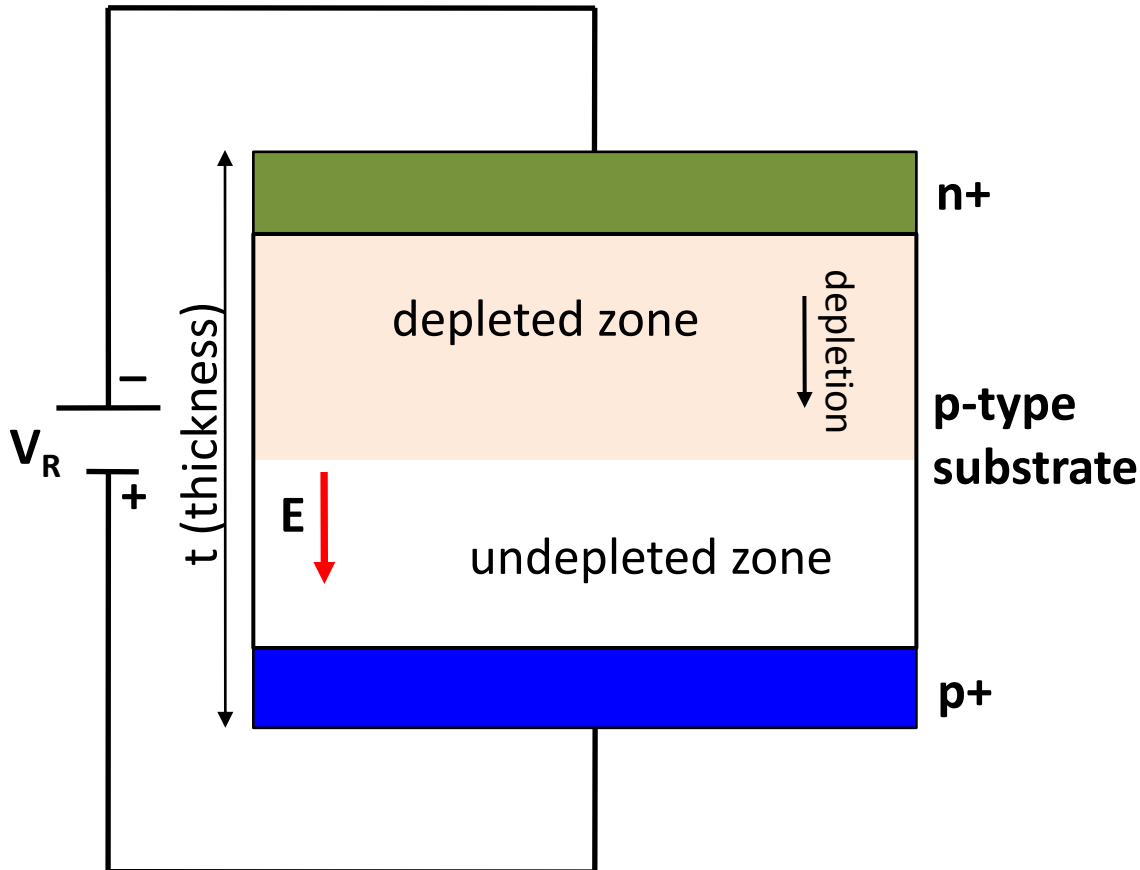
W_F : depletion width at forward bias
 W_R : depletion width at reverse bias
 W_E : depletion width in equilibrium

$$W_F < W_E < W_R$$

Reverse bias scheme is used for silicon detectors in HEP

A real life case

- Larger active area provided by n(p)-type substrate



Depletion with

$$W = \sqrt{2\epsilon\rho\mu V_R}$$

Resistivity

$$\rho = \frac{1}{q\mu N_{\text{eff}}}$$

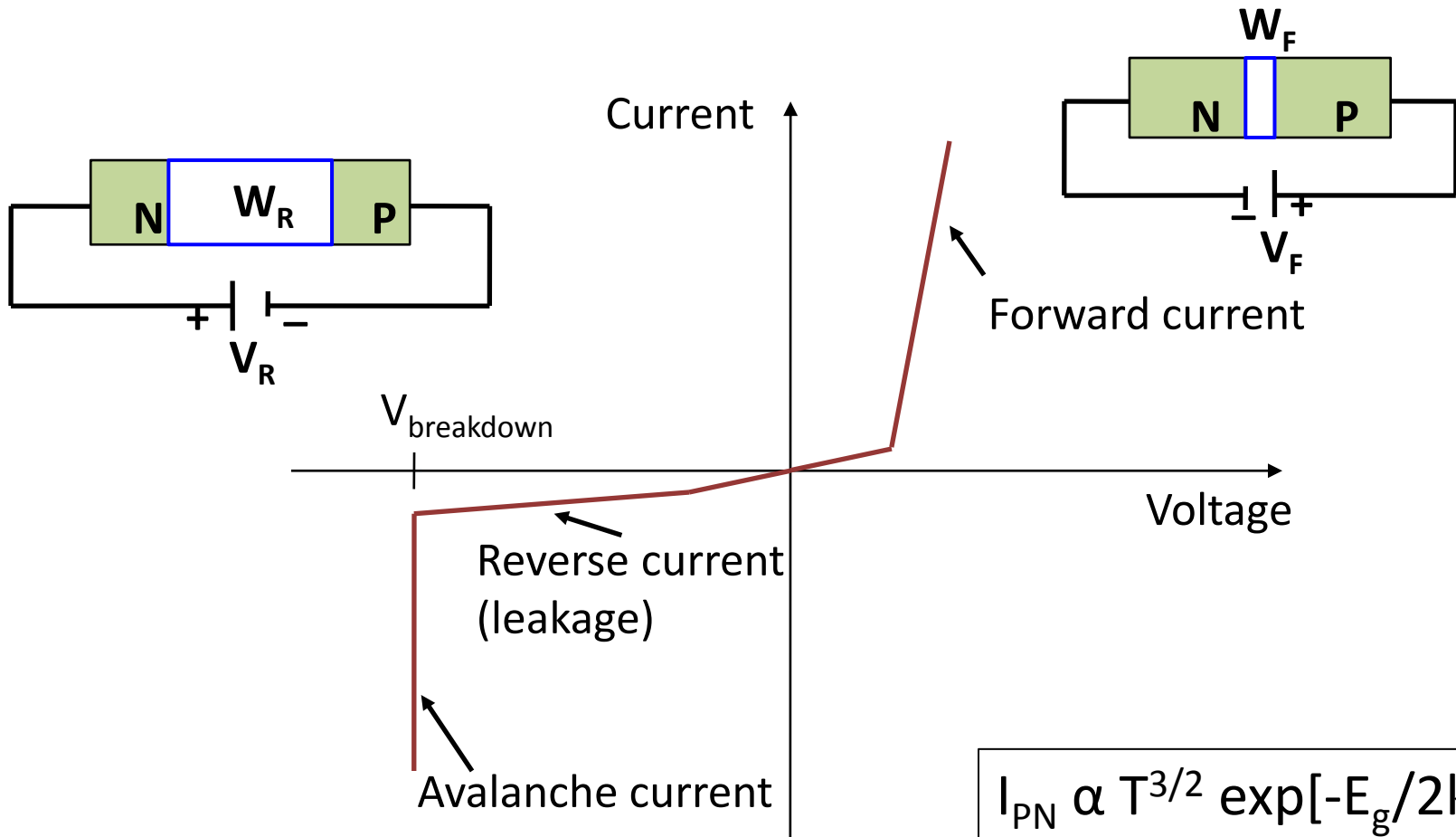
Depletion voltage

$$V_d = \frac{t^2}{2\epsilon\rho\mu}$$

Effective doping concentration

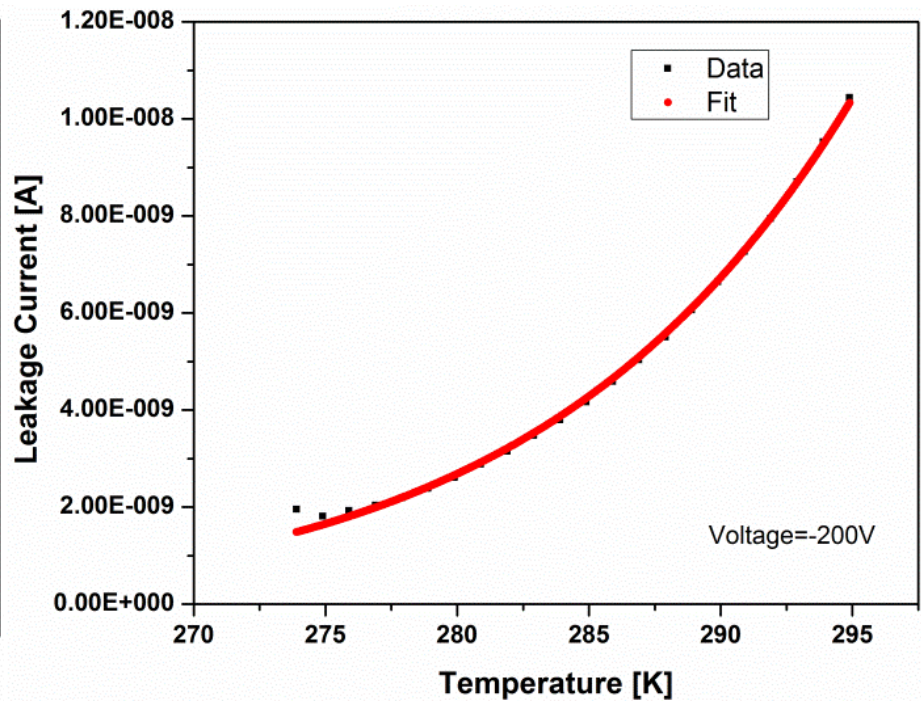
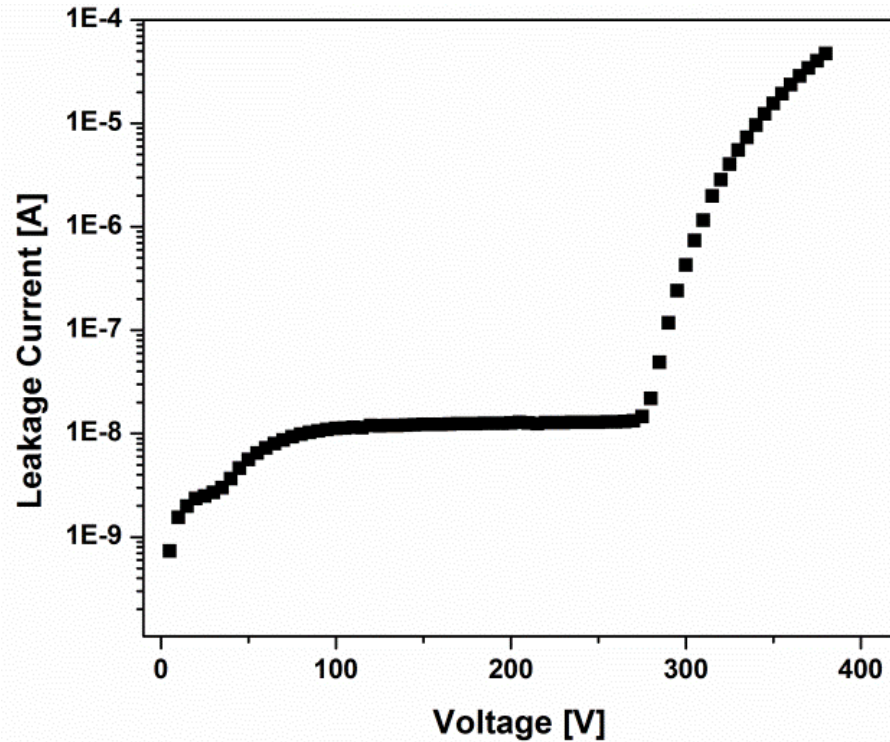
$$N_{\text{eff}} = |N_D - N_A|$$

PN-junction current



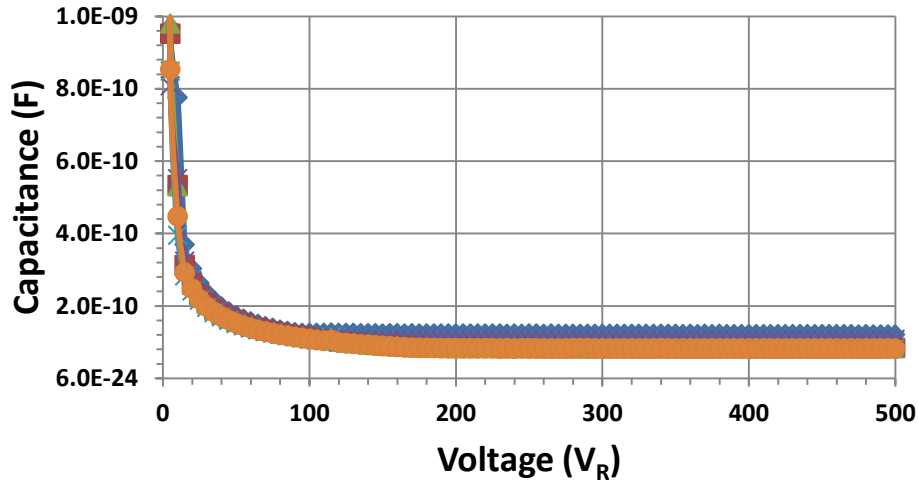
Leakage current

$$\text{Current} \propto T^{3/2} \exp[-E_g/2k_B T]$$

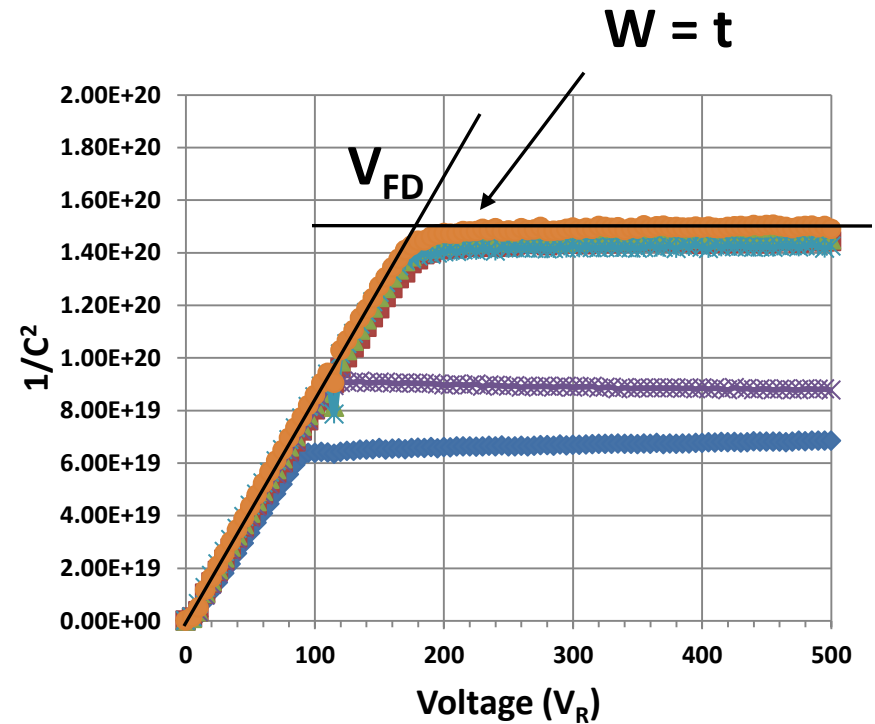


Breakdown voltage

- Extracted from capacitance measurements



$$C = A \sqrt{\frac{\epsilon}{2\rho\mu V_R}}$$



Charge collection

- Carriers are drifted under the electric field

$$v_{e,h}(x) = \mu_{e,h}E(x) \quad \mu_{e^-} = 1500 \text{ cm}^2/\text{Vs} \quad \mu_h = 450 \text{ cm}^2/\text{Vs}$$

- Total drift time for a carrier created at depth x

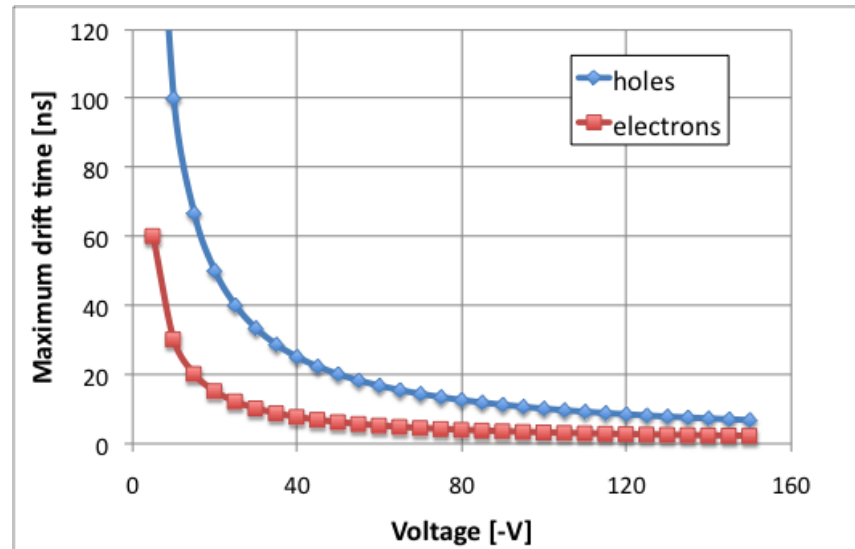
$$t(x) = \frac{d^2}{2\mu V_{FD}} \ln \left(\frac{V_B - V_{FD}}{V_B - V_{FD} + 2V_{FD}(1 - x/d)} \right)$$

V_B = applied voltage
 V_{FD} = full depletion voltage
 d = detector thickness

- Maximum drift time ($V_B \gg V_{FD}$)

$$t_{\max} = t(x = 0) = \frac{d^2}{2\mu V_B}$$

For $d = 300 \mu\text{m} \rightarrow$

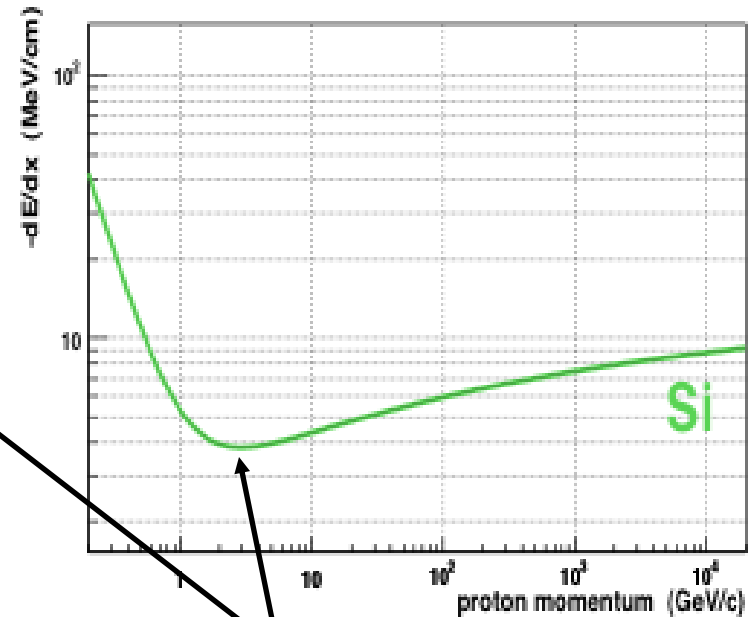
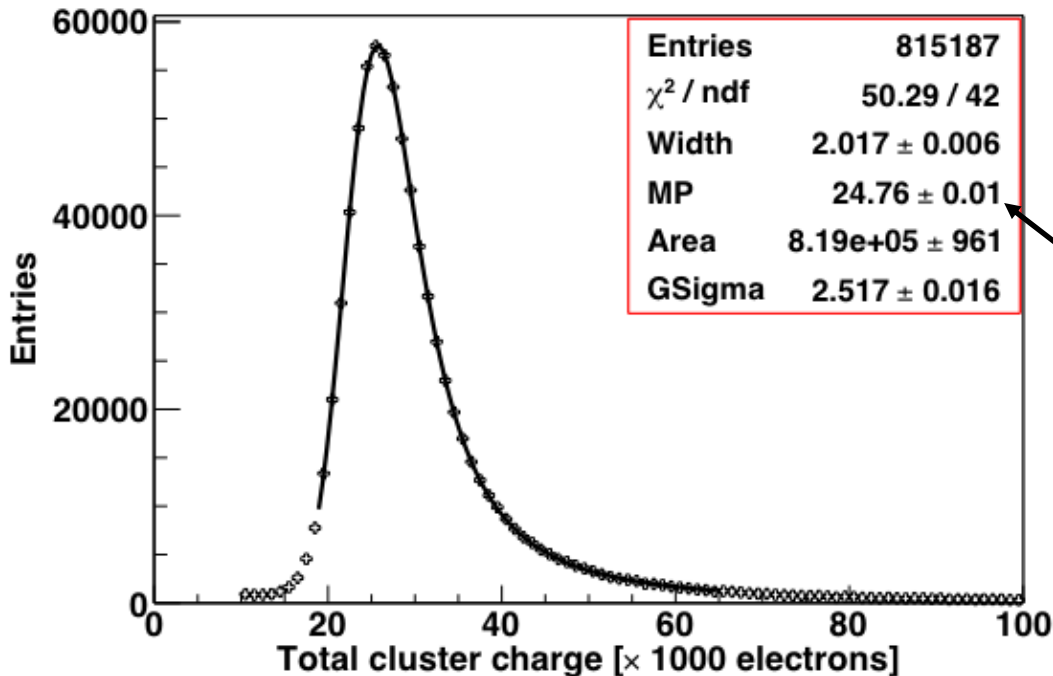


Radiation detection

Ionization energy loss:

$$-\frac{dE}{dx} [\text{MeV/cm}] = 2\pi N_A r_e^2 m_e c^2 \rho \frac{Z}{A} \frac{z^2}{\beta^2} \left[\ln \left(\frac{2m_e c^2 \beta^2 \gamma^2 E_{\text{kin}}^{\text{max}}}{I^2} - 2\beta^2 - \delta \right) \right]$$

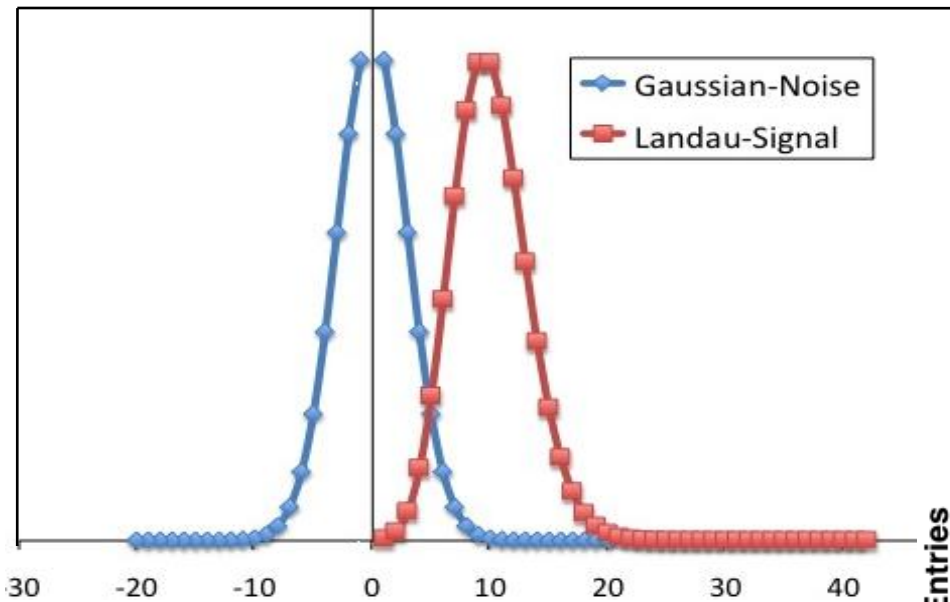
- Ionization loss follows Landau statistics
- Average energy loss is 390 eV/ μm of Si
- Average charge is 108 e-h/ μm of Si
- Most probable charge 80 e-h/ μm of Si
 - 24 ke in 300 μm thick Si



Most ionizing particle = Landau MP

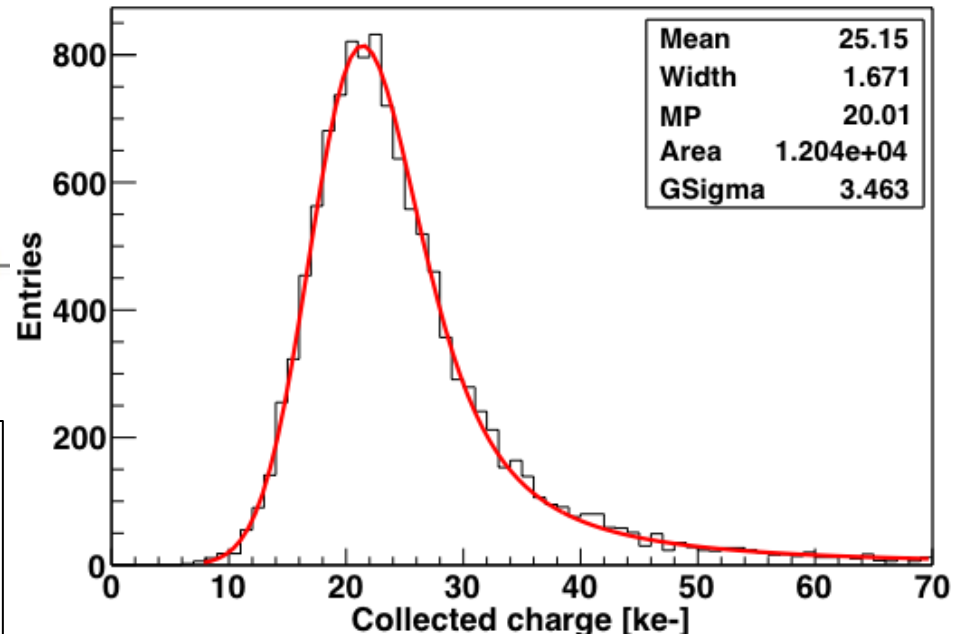
Signal-to-Noise (S/N)

- Signal is Landau, noise is Gaussian



$$S/N = \frac{\text{Most probable signal}}{\text{Noise}}$$

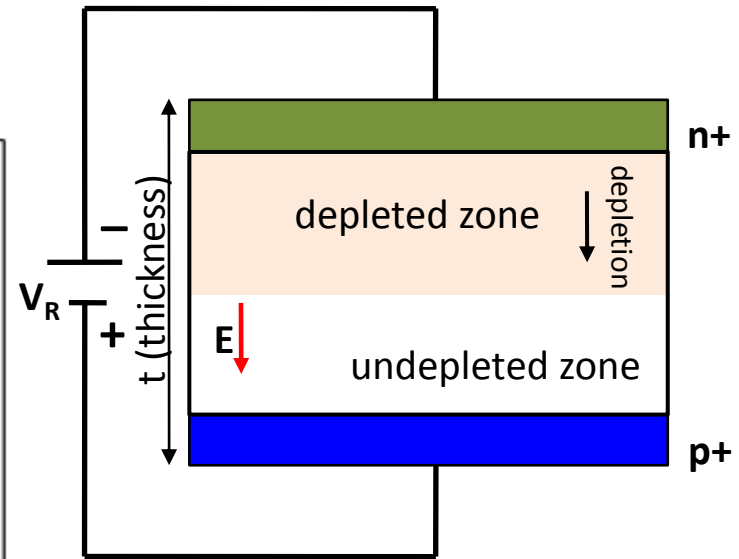
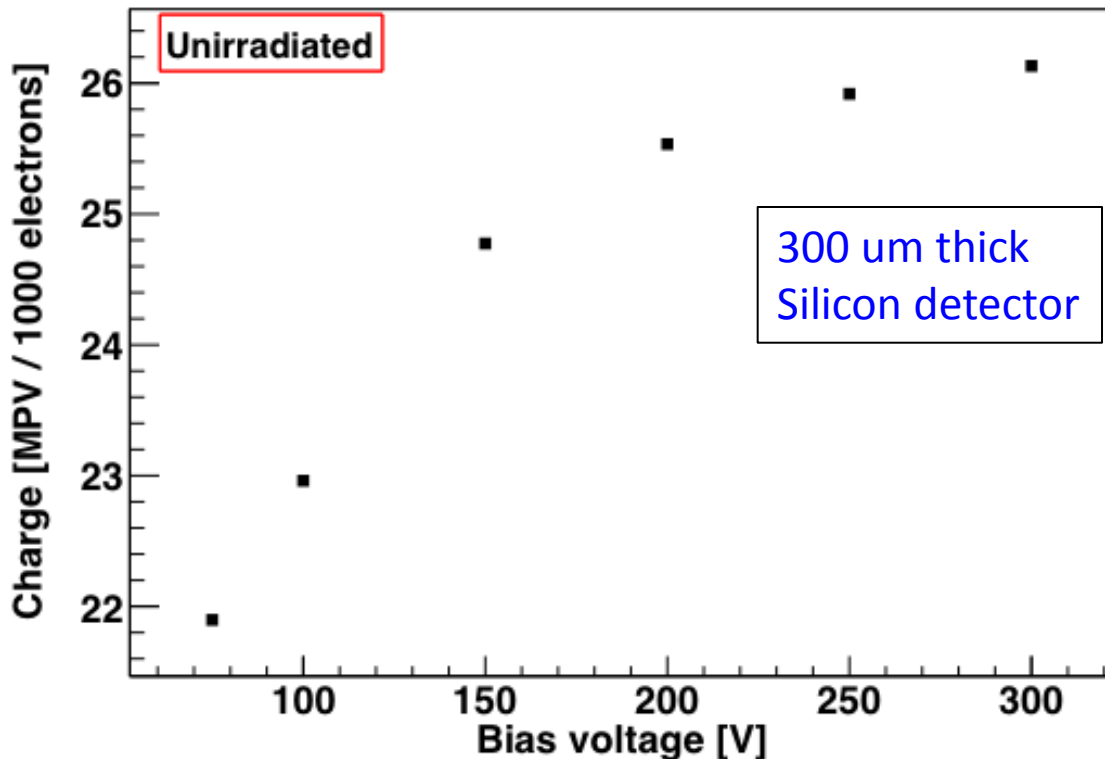
Landau+Gaussian fit



- Ionization energy loss follows Landau statistics
- Random electronic noise centered at zero is Gaussian

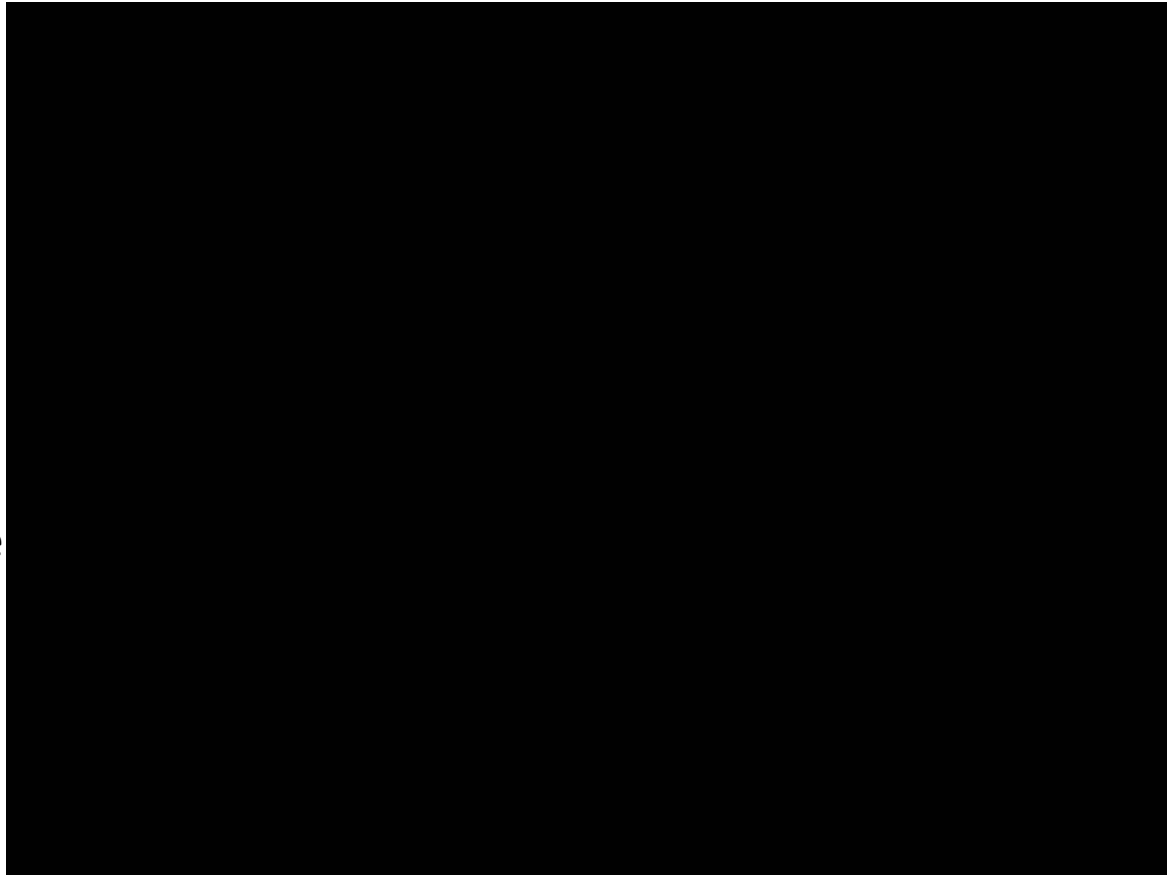
Charge vs applied bias

- Collected charge increases until full depletion is reached

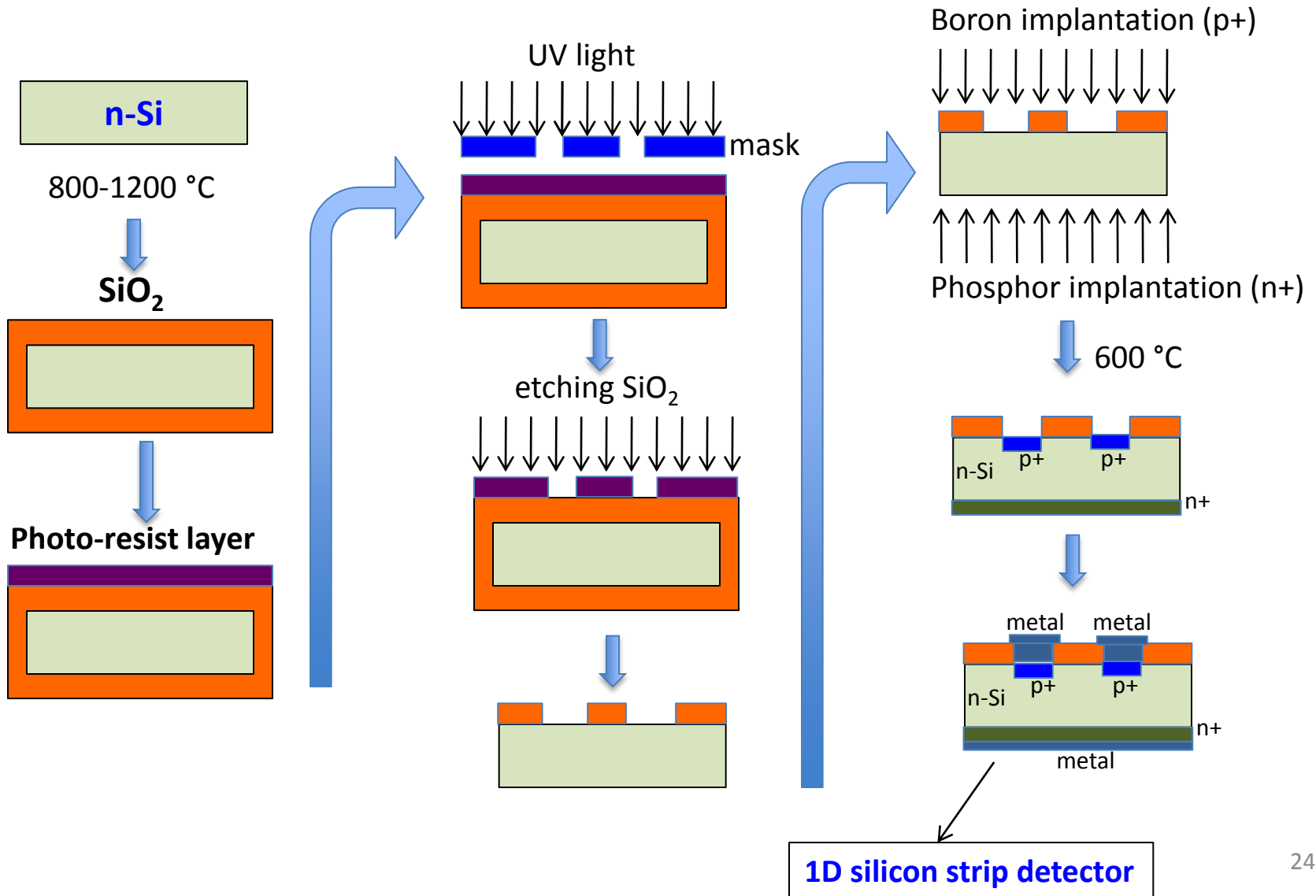


Silicon wafer fabrication

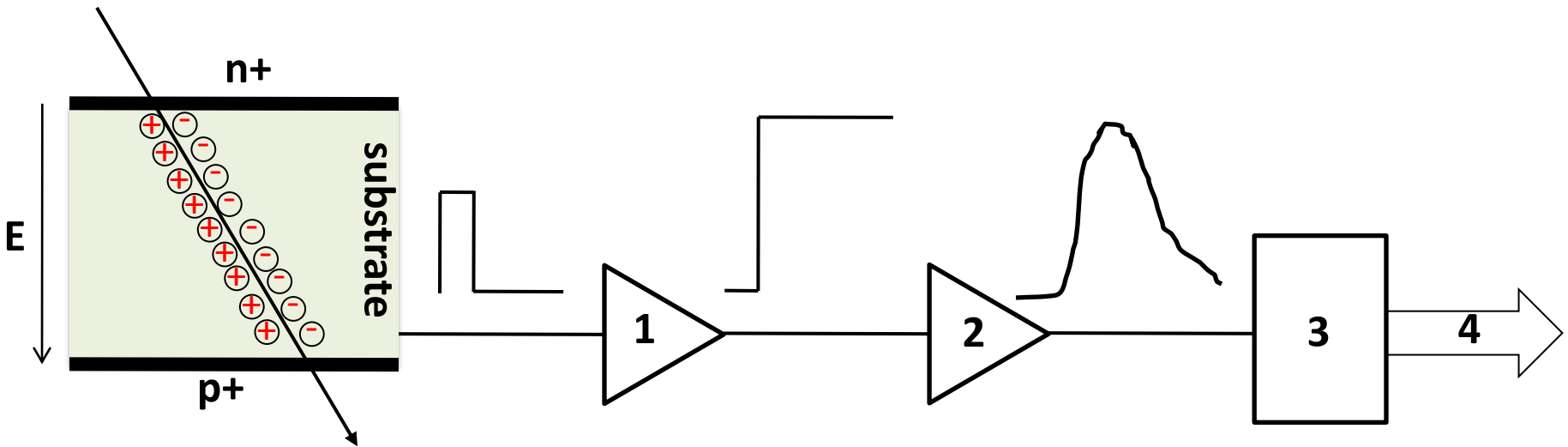
- Hyper-pure polysilicon chunks
- Melt and add impurity to make n-type (P) or p-type (B)
- Pour into a mold to make a polysilicon cylinder
- Melt onto a mono-crystal silicon seed by means of RF power
- Seed and melted polysilicon rotations are opposite to grow a round shape mono-crystal **ingot**
- Employ a grindwheel to form the **ingot** into a desirable diameters
- Cut ingot into wafers by means multi-wire-sawing
- Lapping to flatten wafers
- Chemical etching to smoothen wafers
- Edge rounding for wafer robustness



Device processing



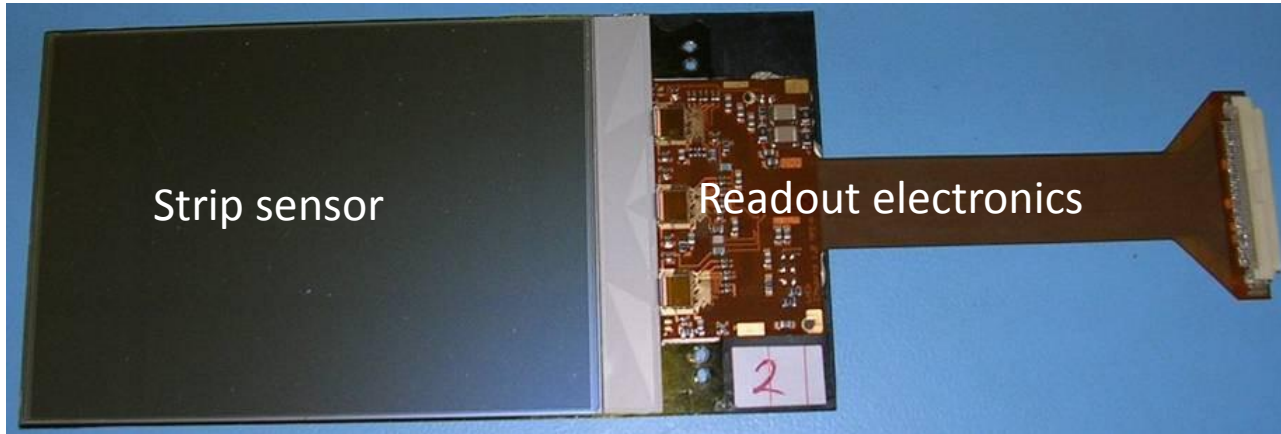
Principle of operation



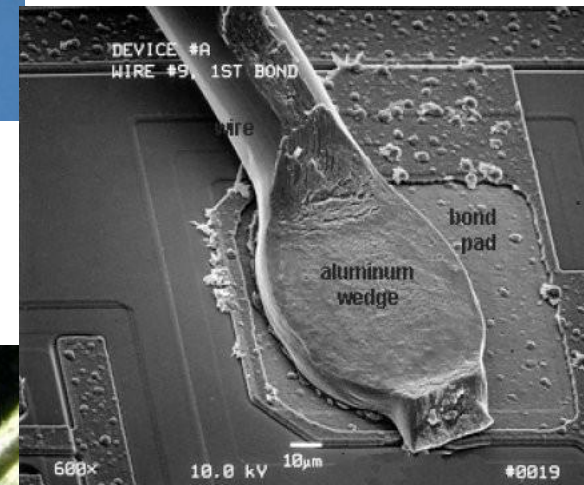
1. **Preamplifier**: amplifying small sensor signal
2. **Pulse shaper**: improving signal-to-noise ratio by filtering signal and attenuating electronic noise
3. **Analog-to-Digital Converter (ADC)**
4. **Buffer**: Store data

Front-end electronics add noise to signal (smearing). Low noise readout is essential

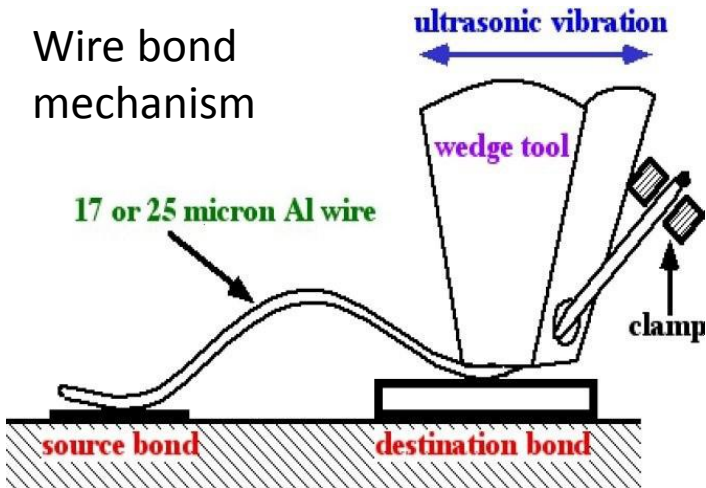
A silicon strip detector



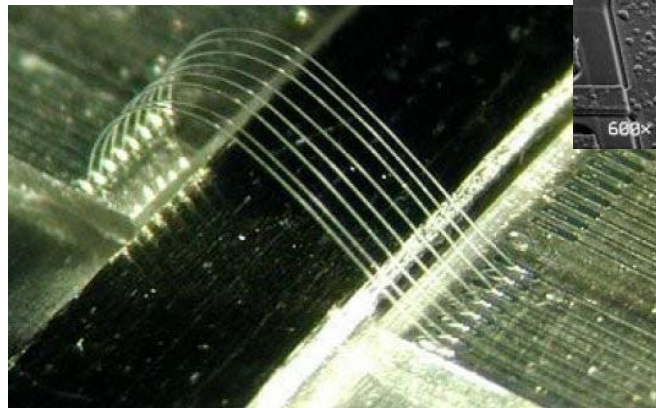
Wire bond (SEM image)



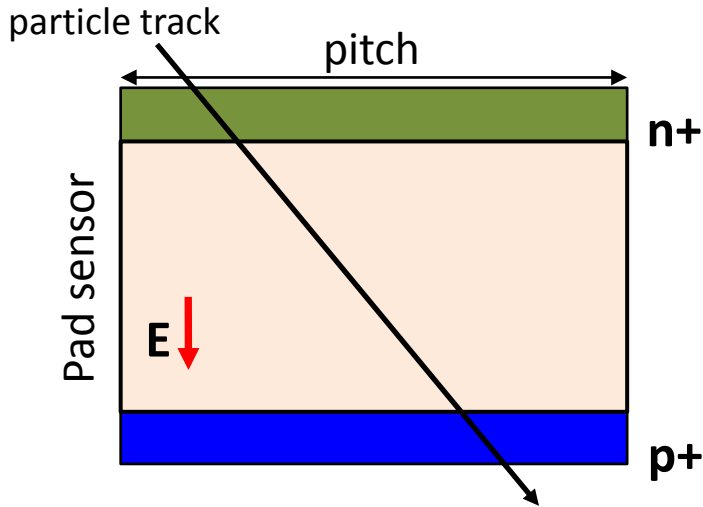
Wire bond mechanism



Wire bonds



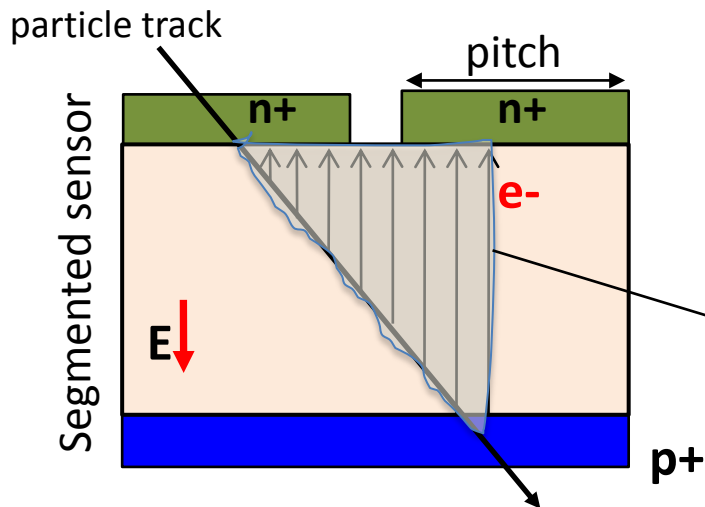
Position resolution



$$\sigma_{\text{hit}}^2 = \int_{-\text{pitch}/2}^{\text{pitch}/2} \frac{x^2}{\text{pitch}} dx = \frac{\text{pitch}^2}{12}$$

$$\sigma_{\text{hit}} = \frac{\text{pitch}}{\sqrt{12}}$$

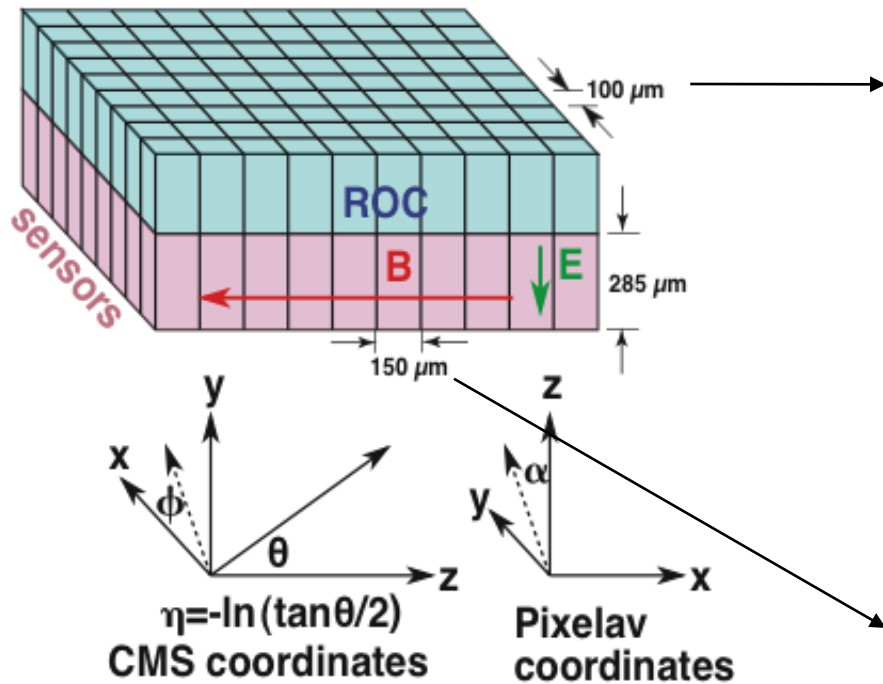
- Charge interpolation improves spatial resolution



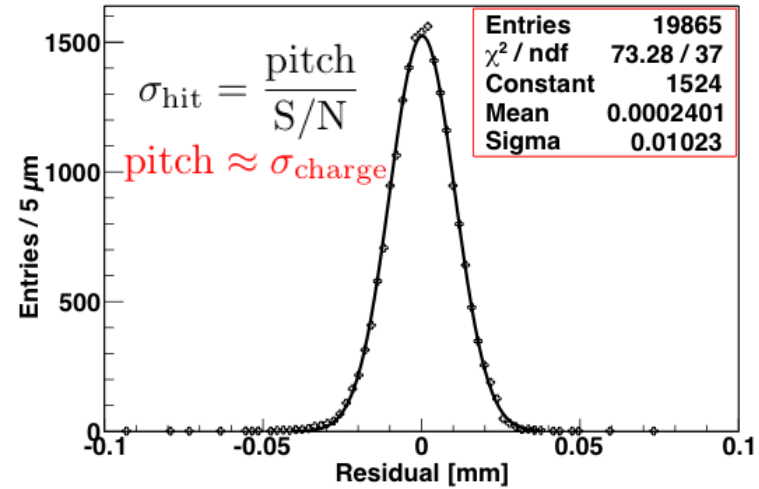
$$\sigma_{\text{hit}} = \frac{\text{pitch}}{S/N}$$

$$\text{pitch} \approx \sigma_{\text{charge}}$$

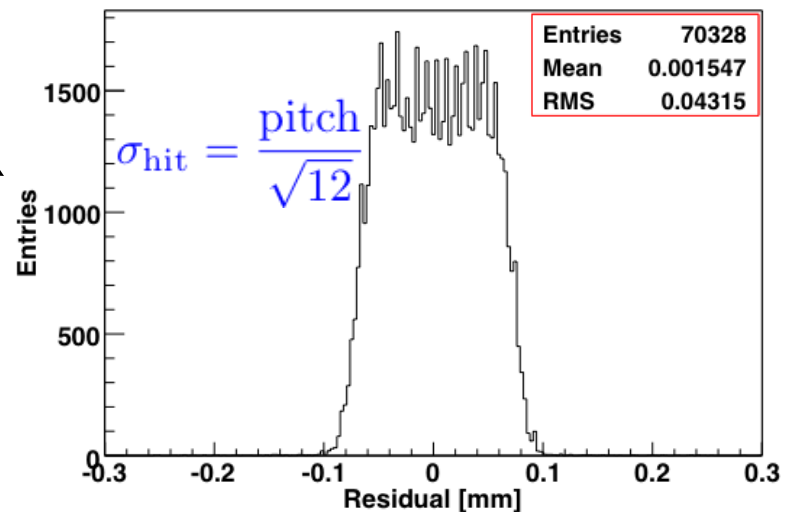
Position resolution



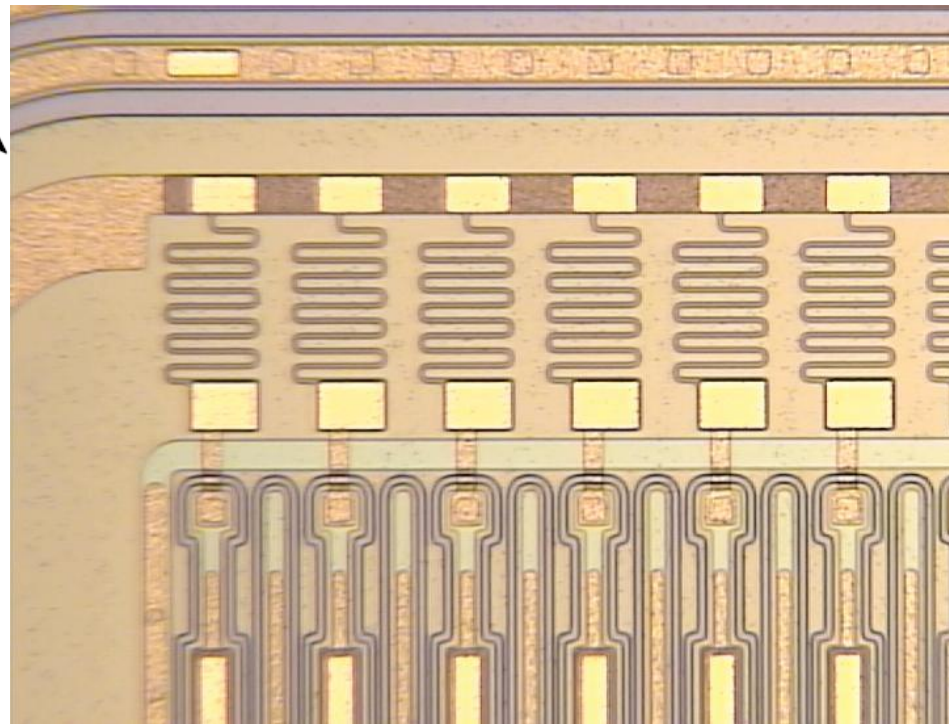
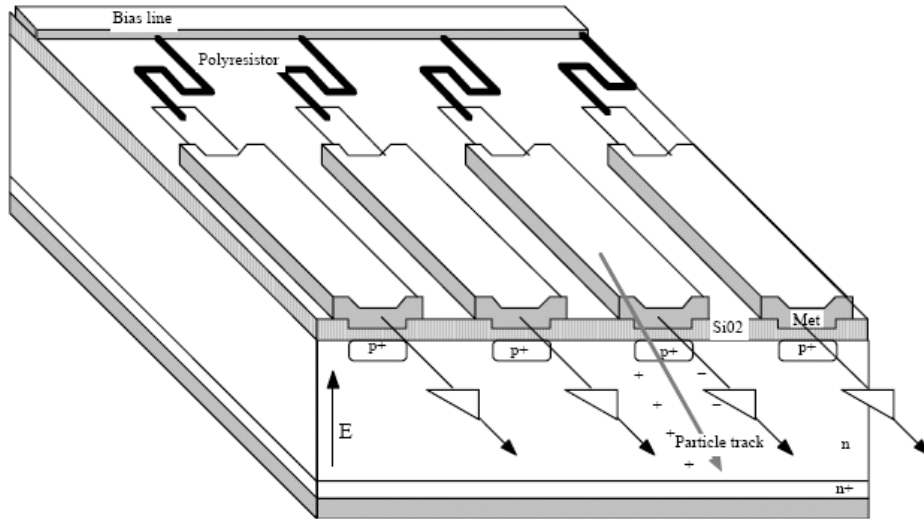
Charge interpolated spatial resolution



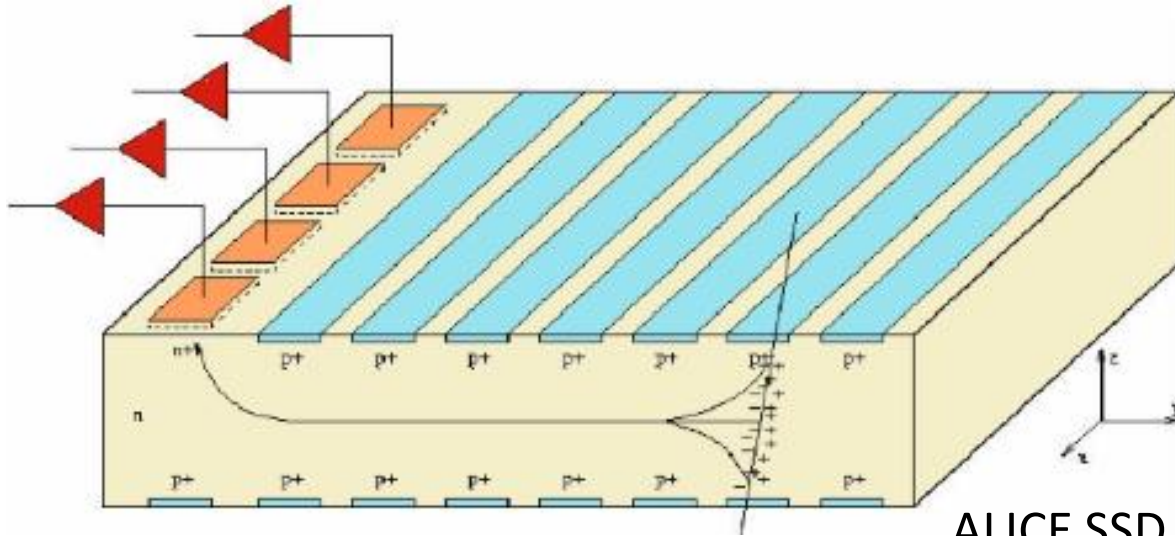
Binary spatial resolution



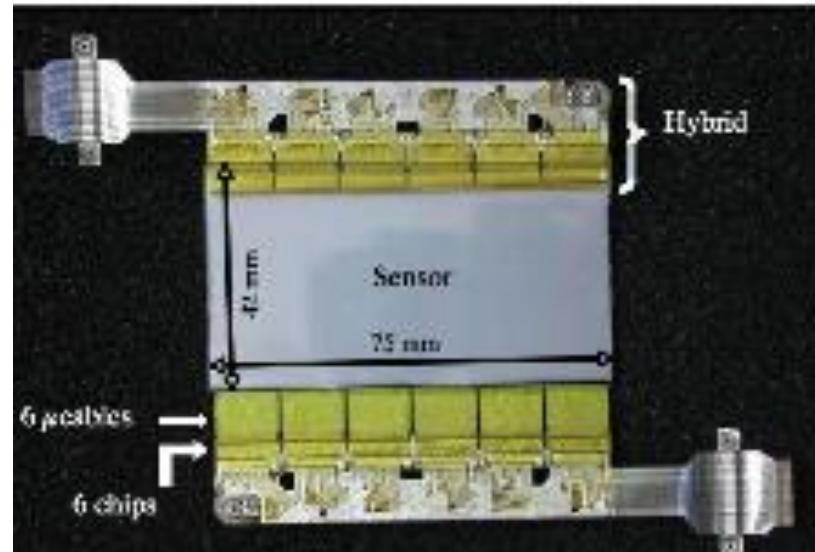
Microstrip detectors (1D)



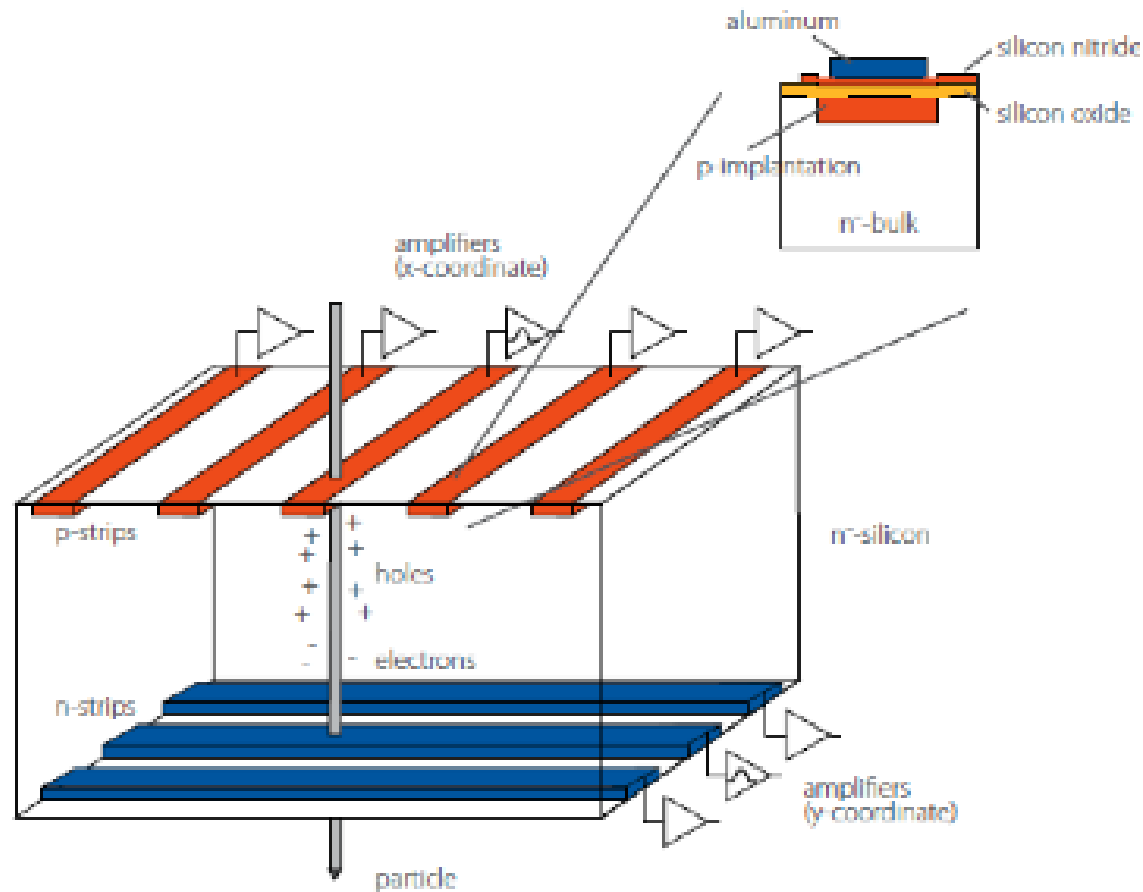
Silicon drift detector (SSD)



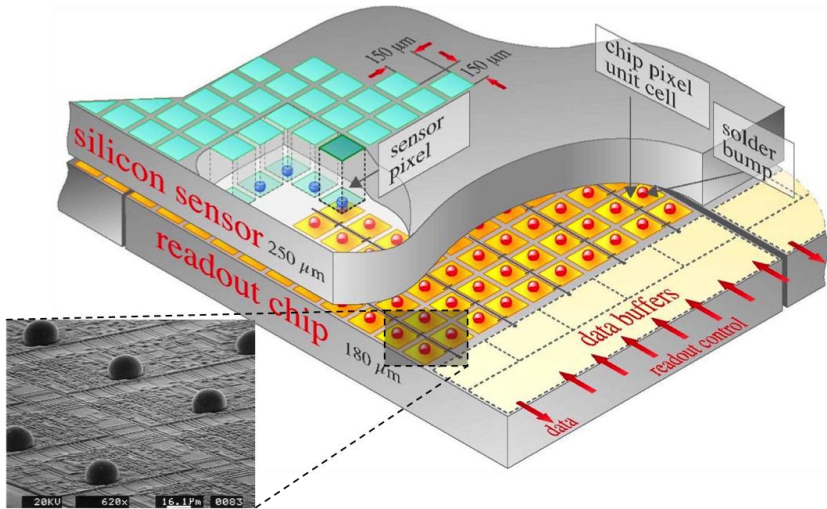
ALICE SSD module



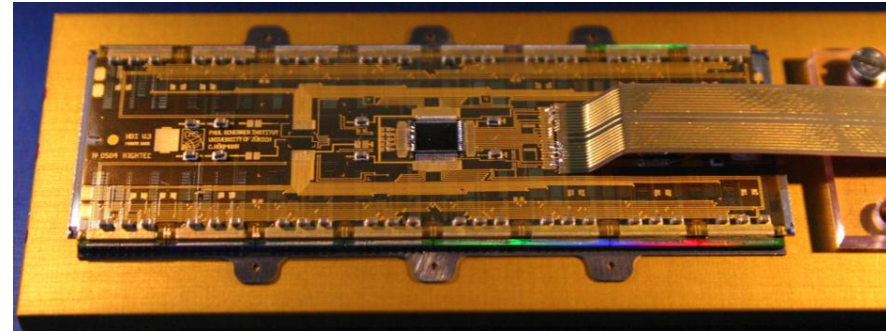
Double-side strip detector



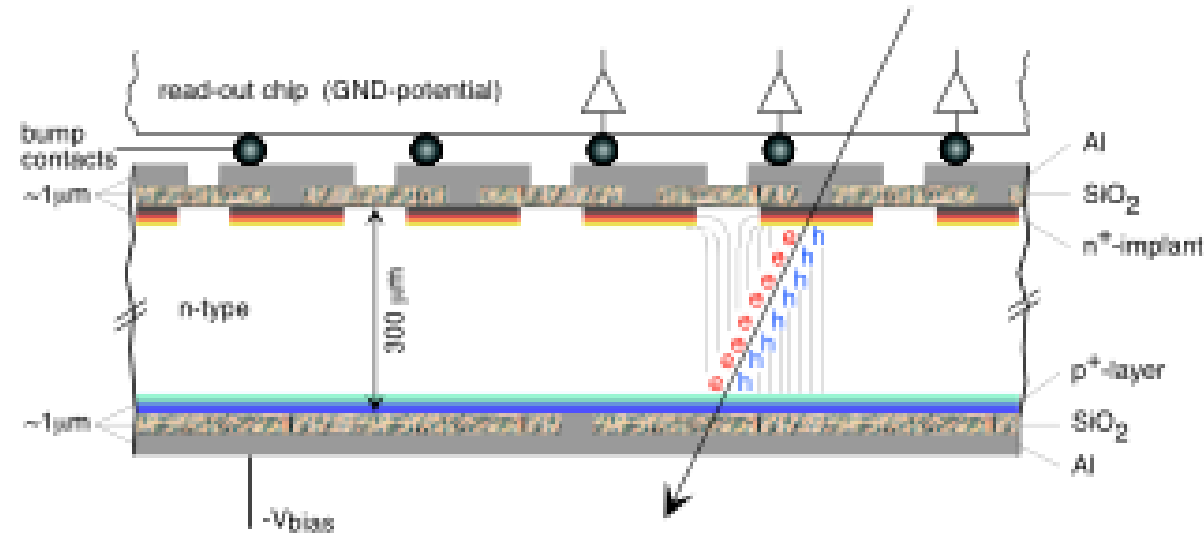
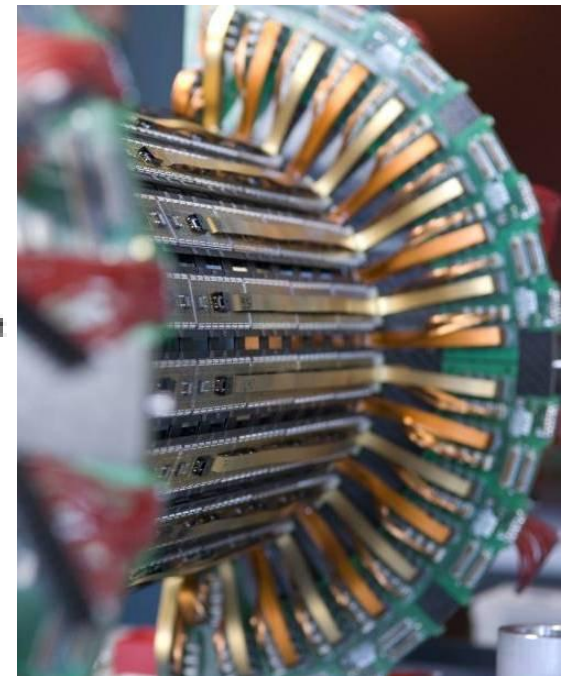
Pixel detector



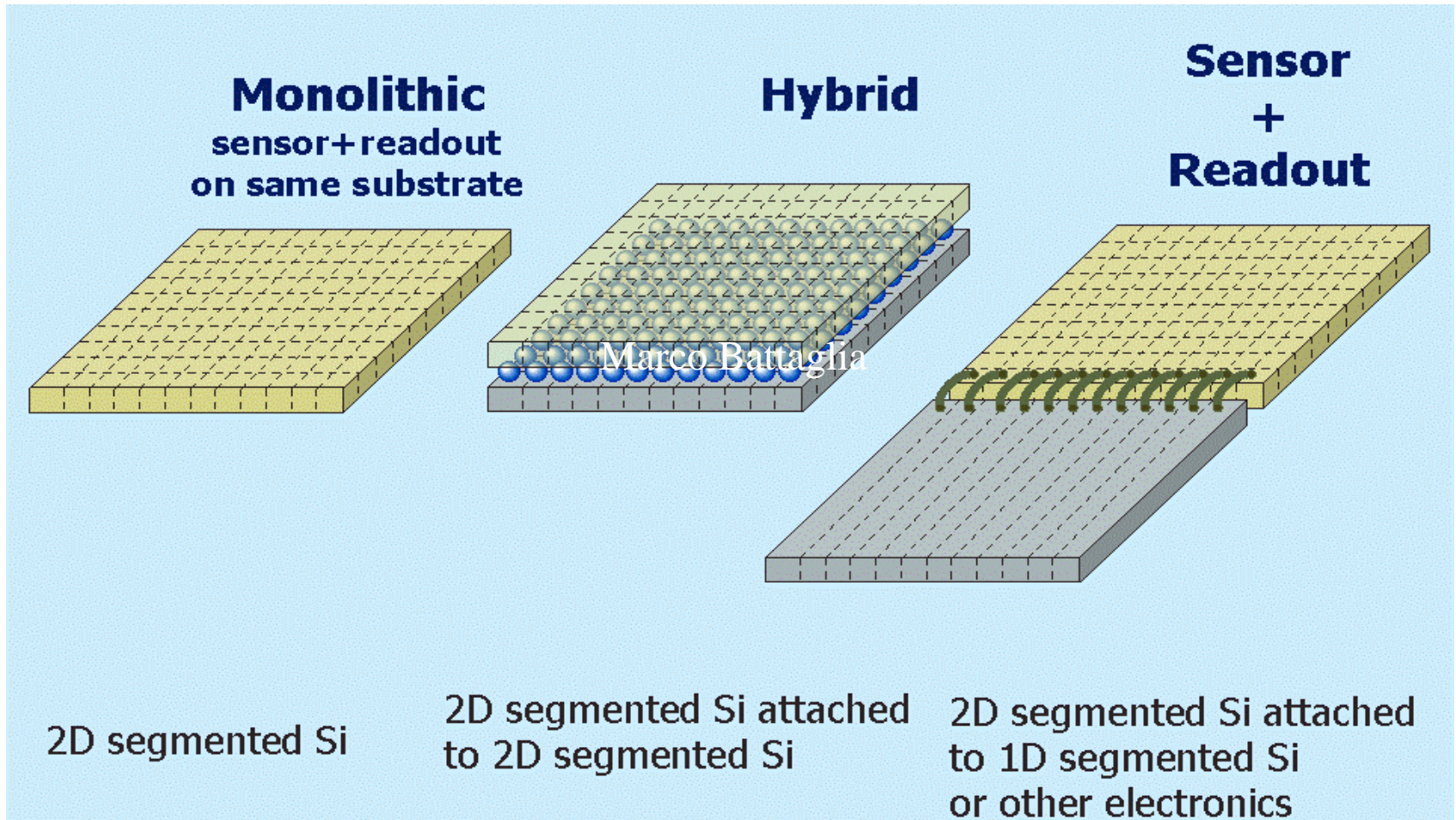
CMS BPIX module



CMS barrel pixel detector



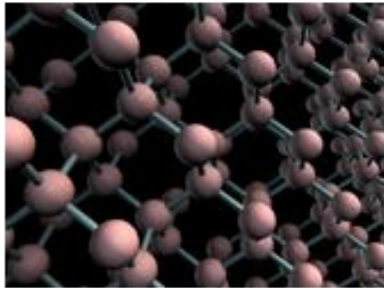
Detector topology



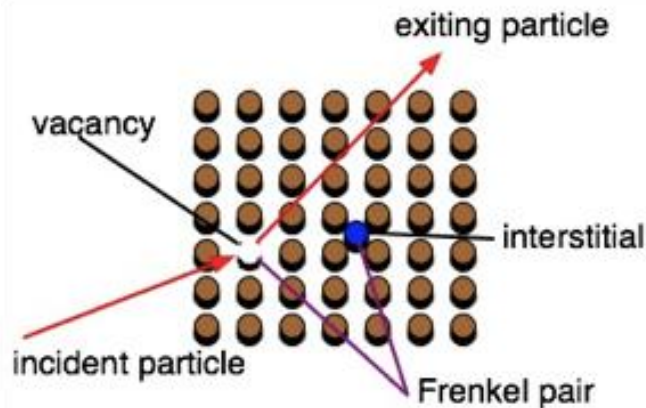
Marco Battaglia, EDIT 2012, Silicon Track, February 2012

Radiation damage

- 1) nuclear particle (i.e. p,n etc.) collides with a lattice atom,
- 2) the lattice atom is kicked out from its position leaving behind a vacancy
- 3) the displaced atom can collide with others creating more interstitial positions (clusters)

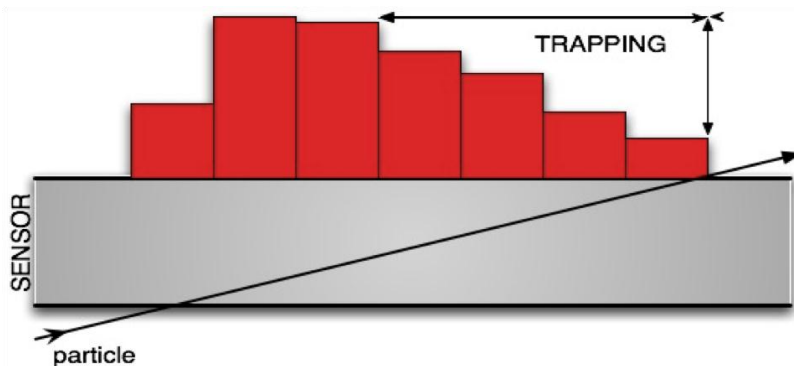
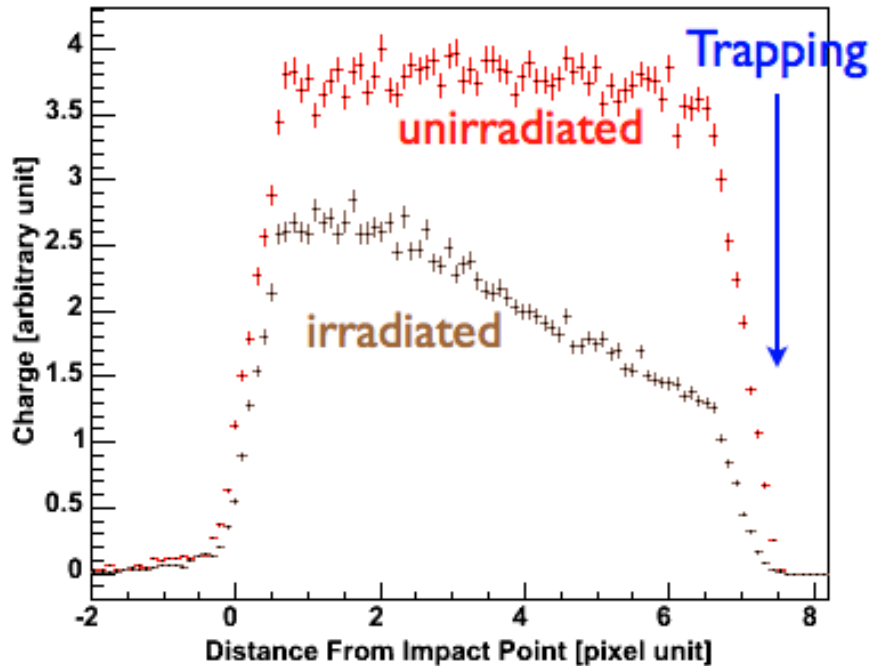


Si Atoms

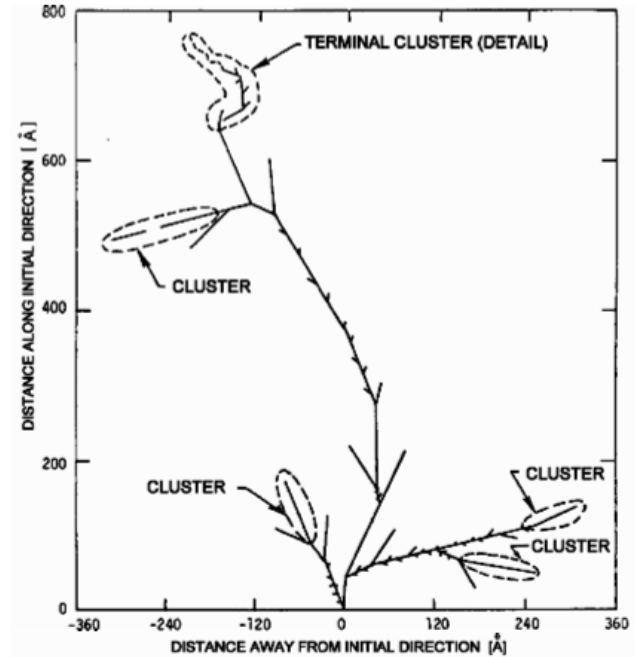


- an energetic nuclear particle knocks out an atom resulting in **Frenkel pairs** (defects)
- above 150 K 90 % of pairs recombine (thermal vibrations in lattice)

Radiation damage



Recoil after 1 MeV neutron collision



Radiation induced damages

- cause severe signal losses across sensor thickness

Signal loss can be covered partially

- by increasing high voltage

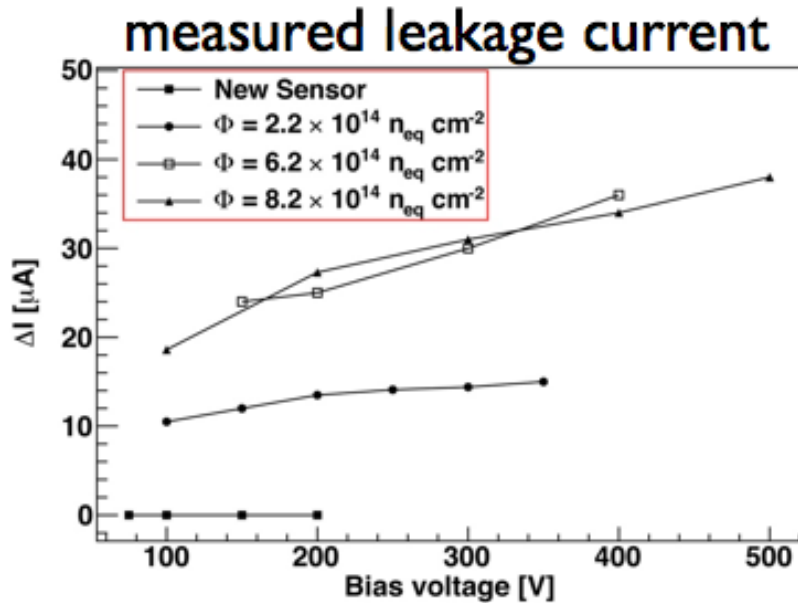
But high voltage

- degrades the position resolution

Radiation damage effects

- Leakage current increases $\Delta I = \alpha \Phi V$
 - ΔI change in current in volume V
 - α is damage constant
 - Φ is time-integrated radiation flux
- Depletion voltage increases
$$V_{FD} = q |N_{eff}| d^2 / 2\epsilon\epsilon_0$$
 - $N_{eff} = |N_{donor} - N_{acceptor}|$
 - $\epsilon\epsilon_0$ permittivity
 - d is sensor thickness
- Charge collection degrades
- Position resolution degrades

Radiation damage – leakage current



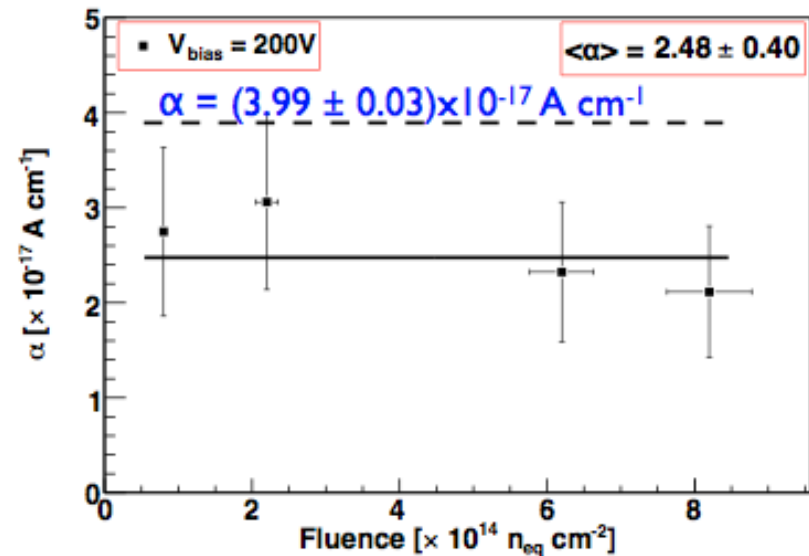
$$\alpha = \frac{\Delta I}{V \cdot \Phi}$$

α - damage rate

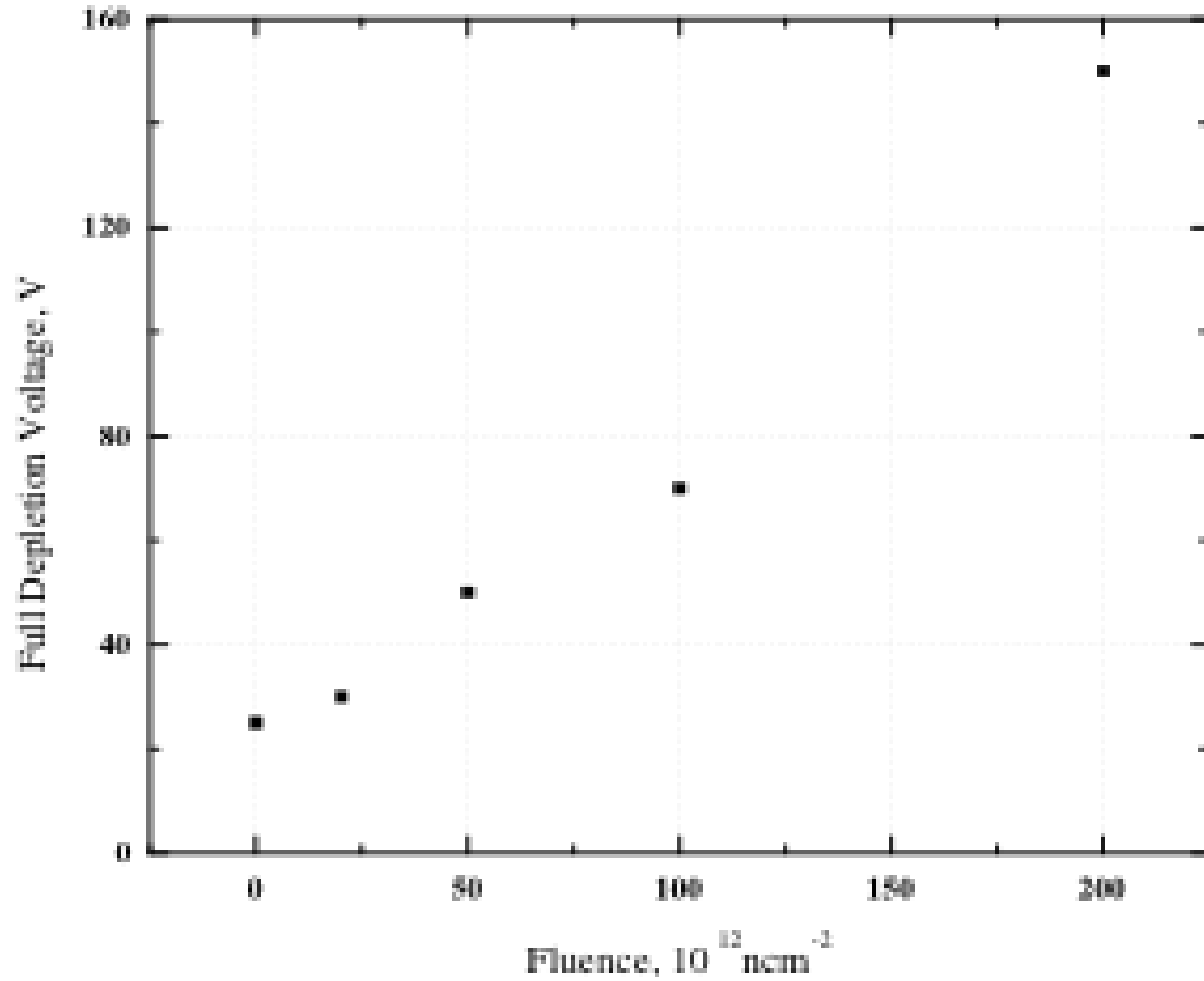
Φ - fluence

ΔI - leakage current

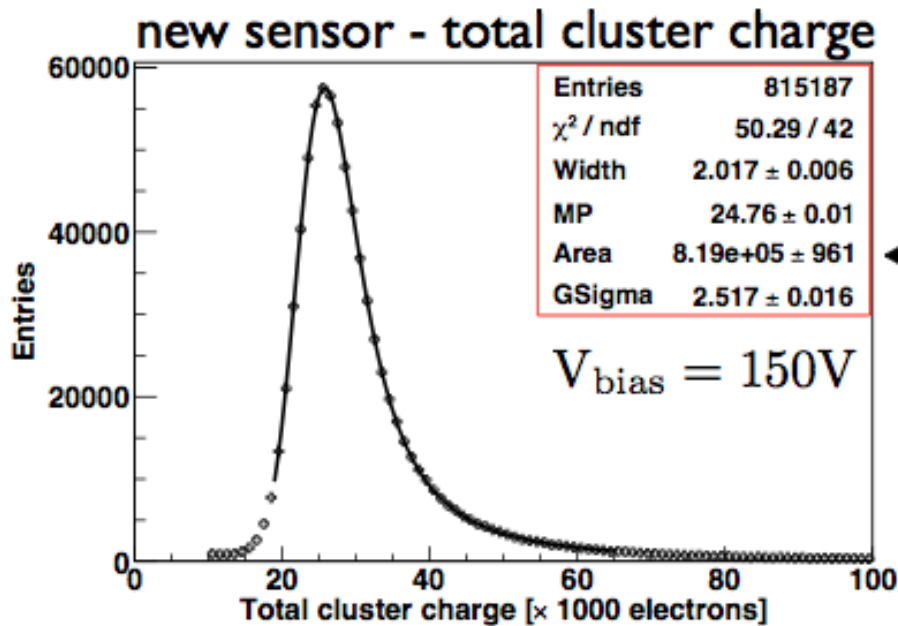
V - volume (0.018 cm^3)



Radiation damage - V_{FD}

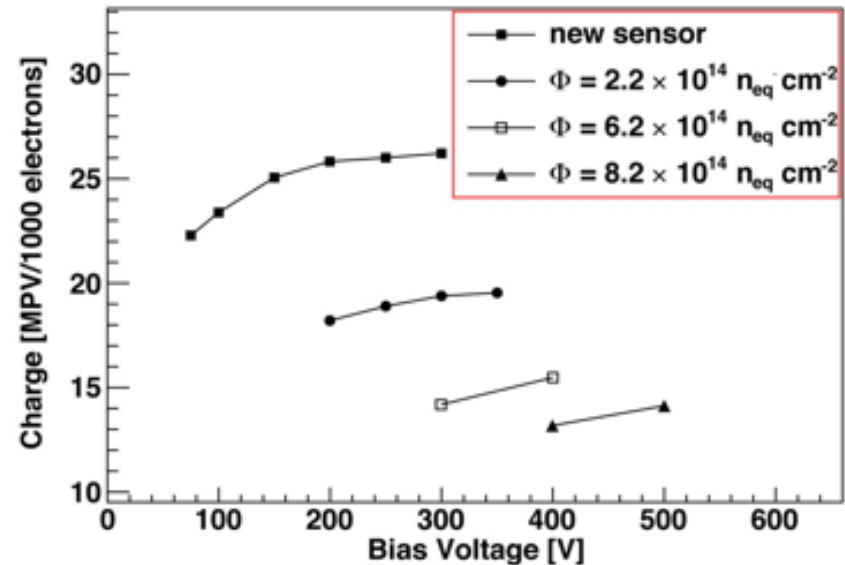


Radiation damage – charge collection



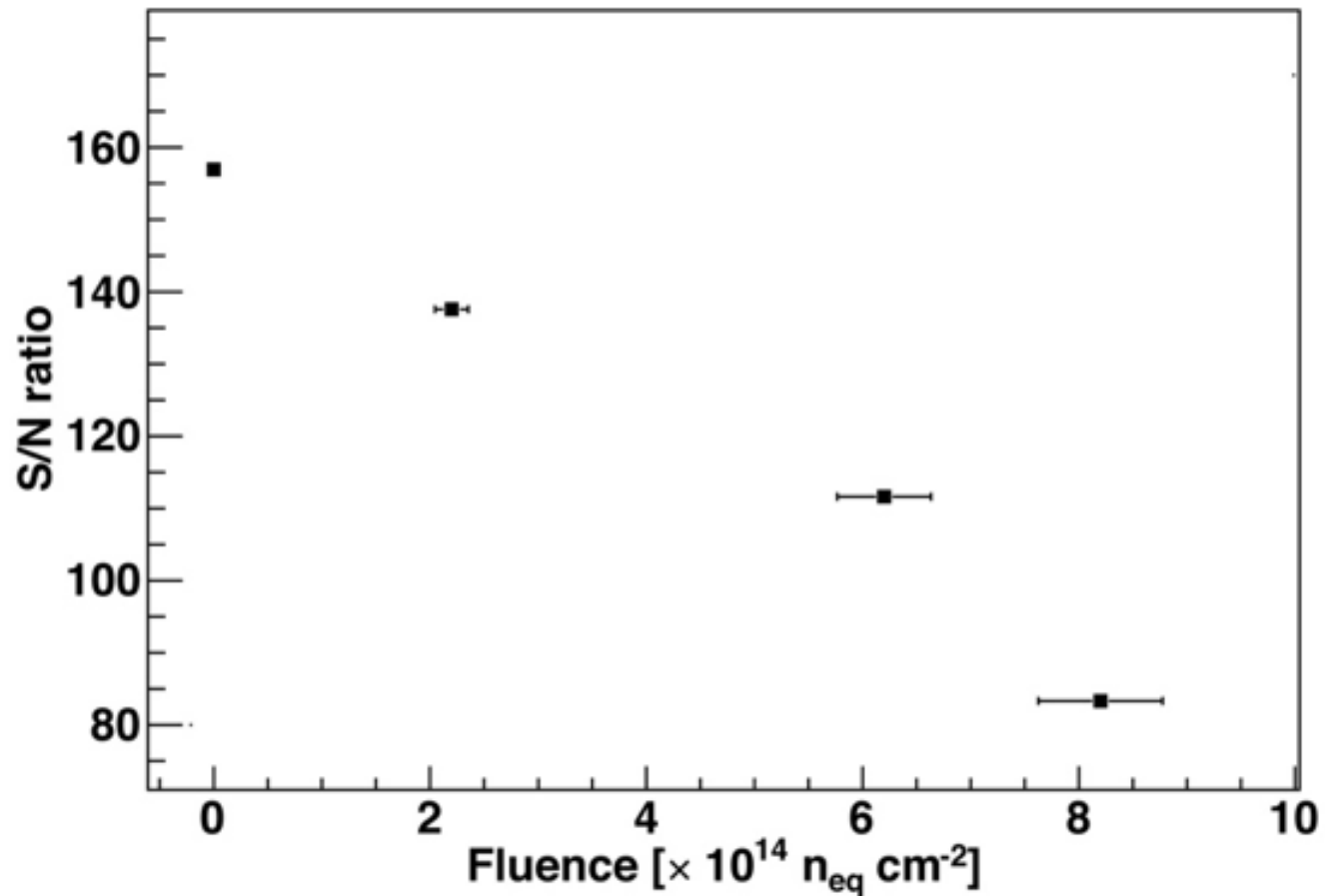
Gaussian convoluted with a Landau function fit to total cluster charge

- measurable Landau most probable values (MP) are shown
- charge collection decreases with irradiation
- higher bias voltage can recover some lost charge



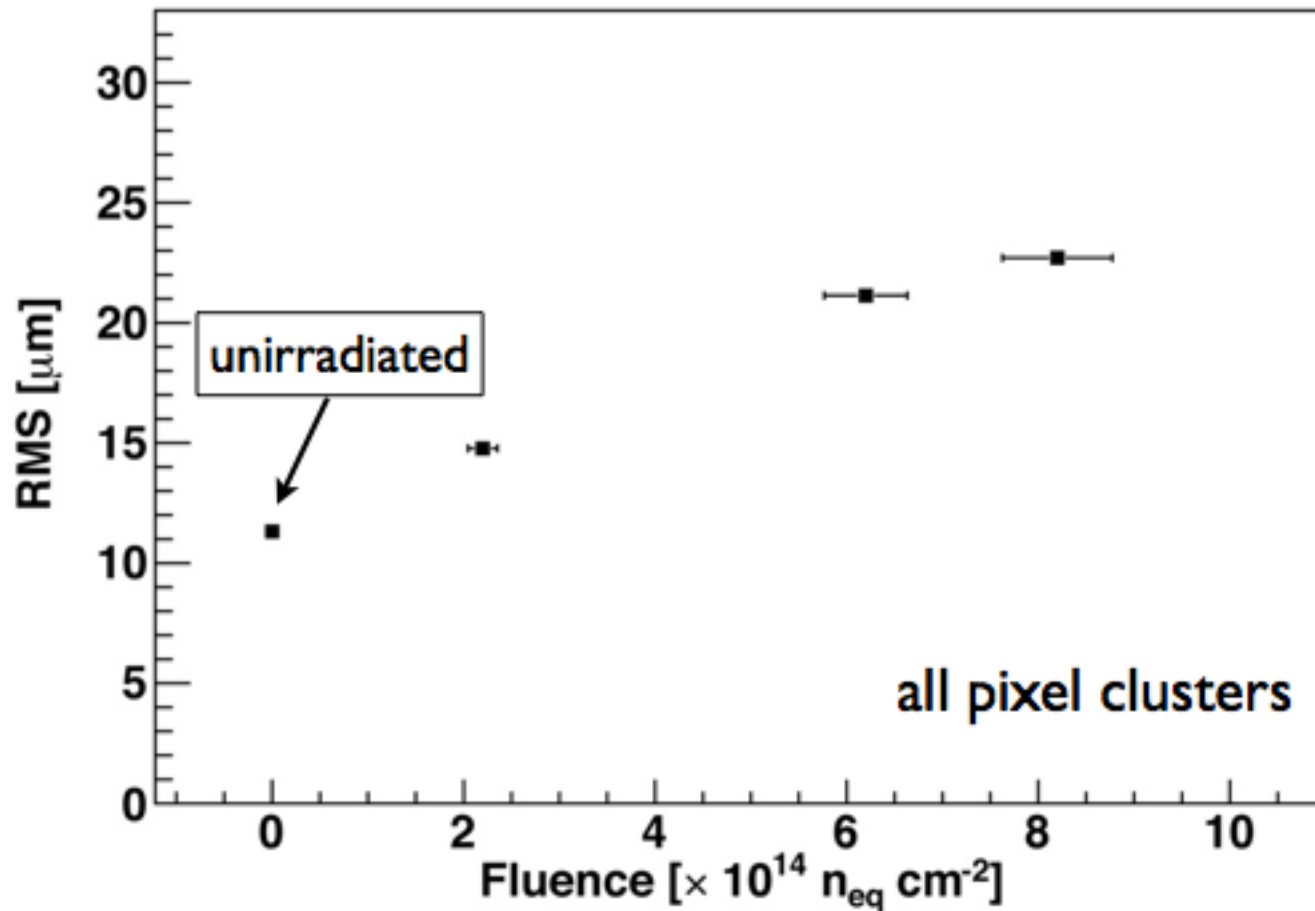
Radiation damage – S/N

- signal (S) is the MP value of the Landau fit
- noise (N) is from the noise test



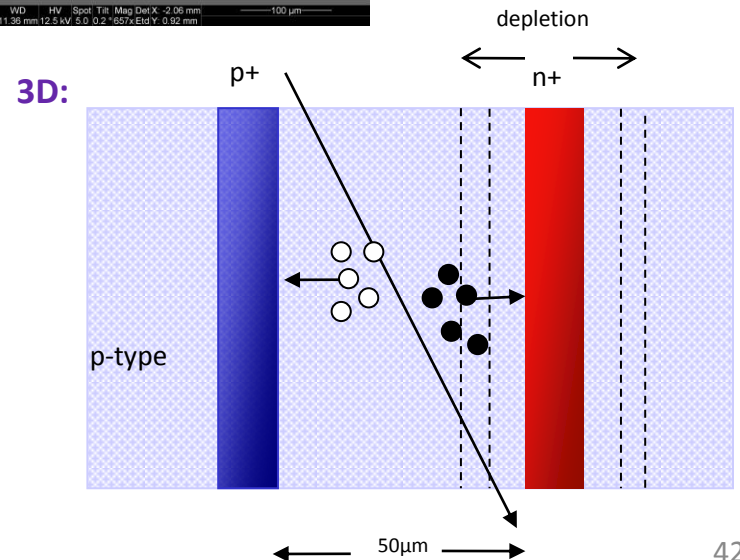
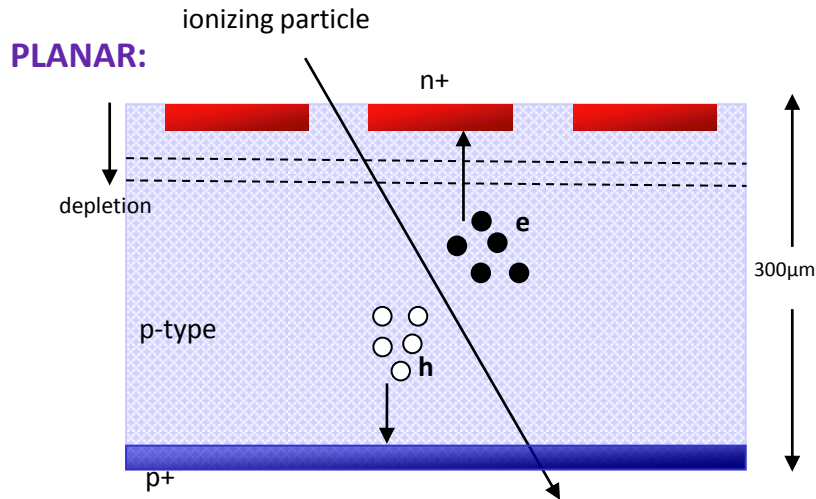
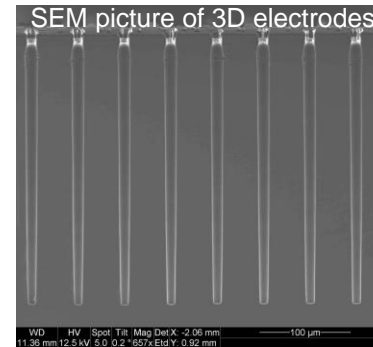
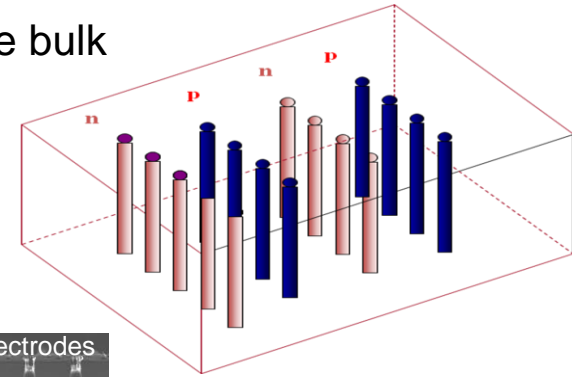
Radiation damage – spatial resolution

- overall resolution (RMS) decreases with irradiation



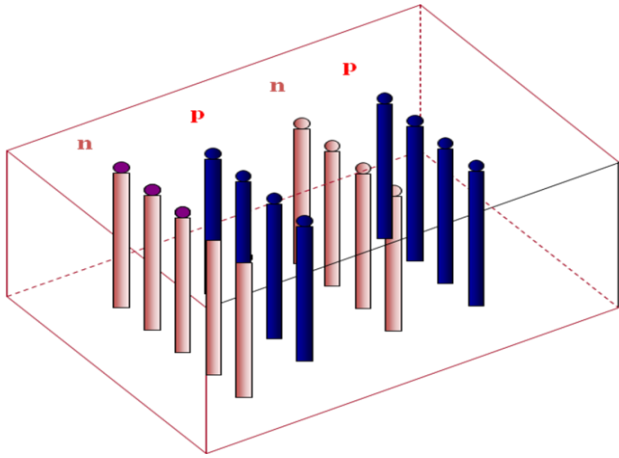
Radiation harder approaches – 3D

- p+ and n+ electrodes are arrays of columns that penetrate into the bulk
- Lateral depletion
- Charge collection is sideways
- Superior radiation hardness due to smaller electrode spacing:
 - smaller carrier drift distance
 - faster charge collection
 - less carrier trapping
 - lower depletion voltage
- Higher noise
- Complex, non-standard processing

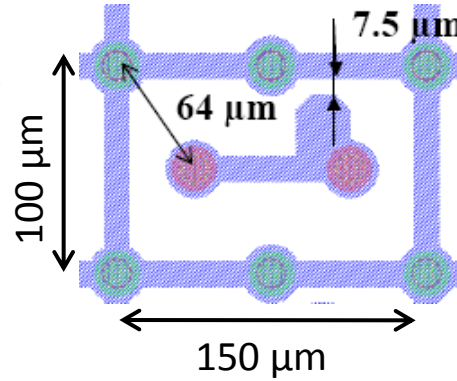


Radiation harder approaches – 3D

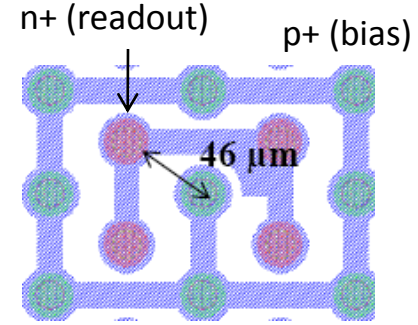
SINTEF 3D
(200 μm thick)



2E Configuration

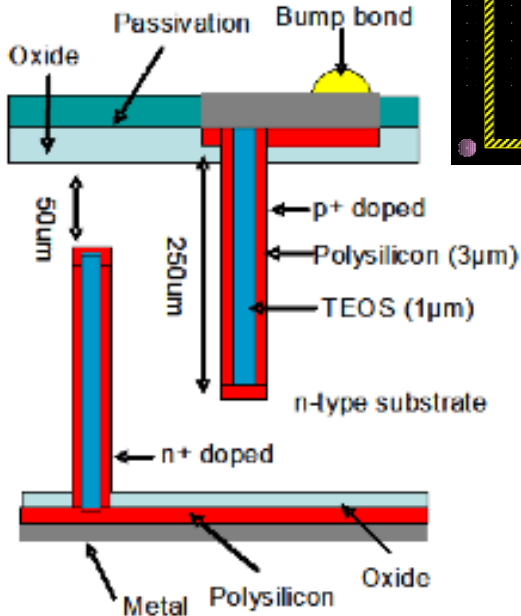


4E Configuration

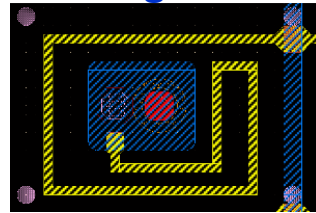


Single-side etching

CNM 3D (200 μm thick)

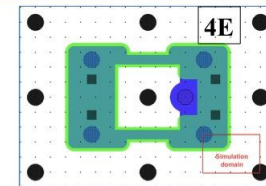
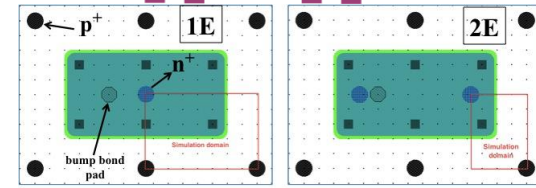
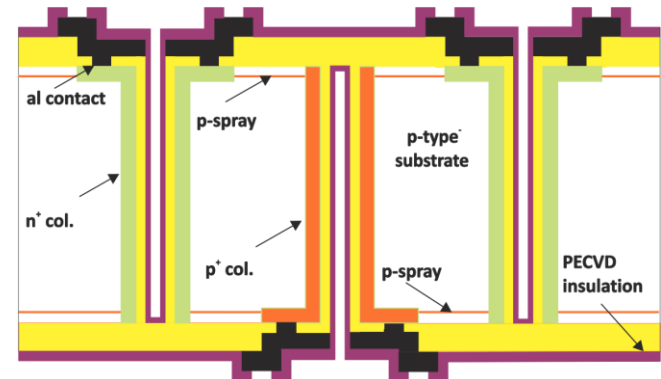


1E Configuration



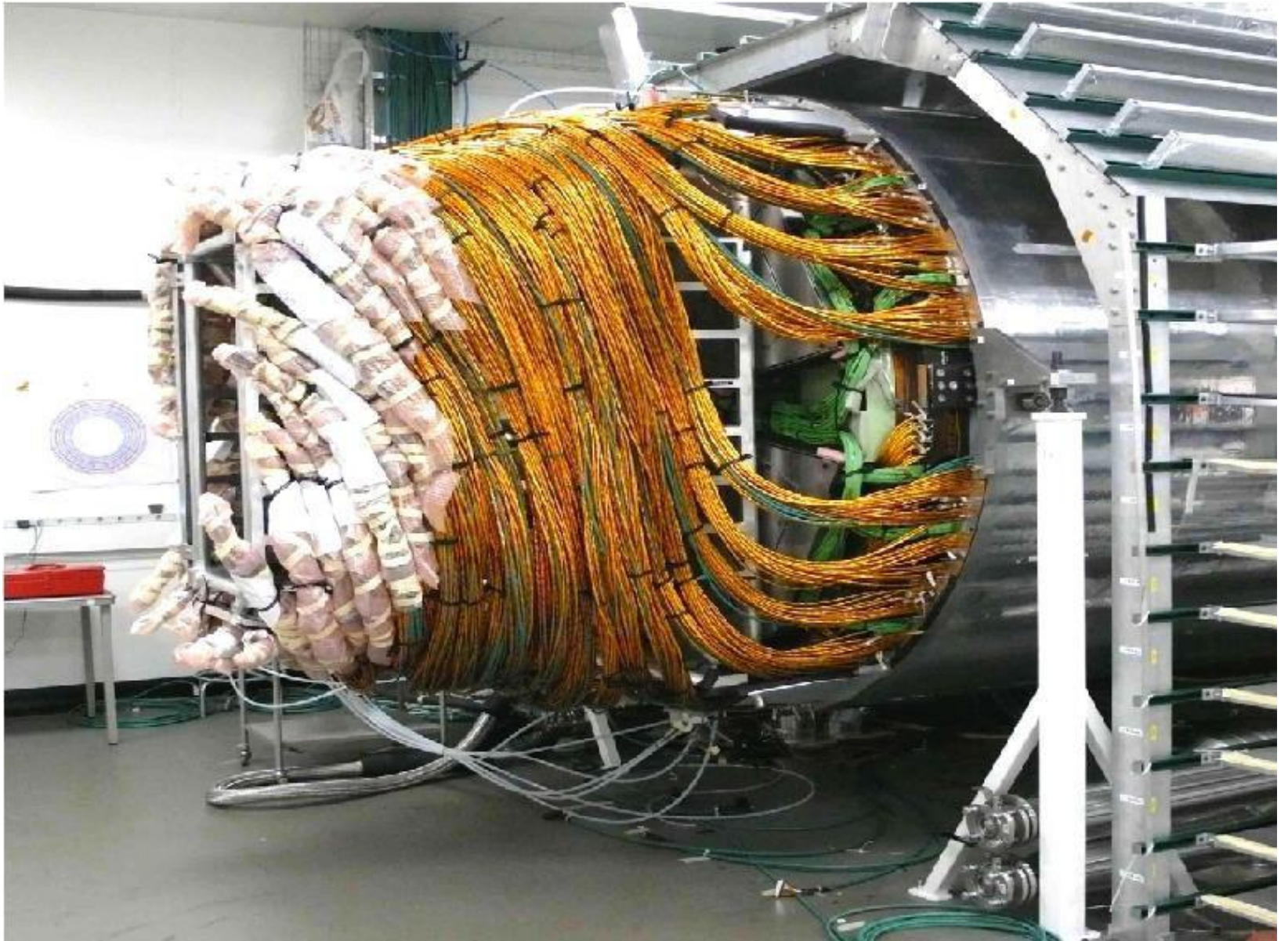
Double-side etching

FBK 3D (200 μm thick)



Double-side etching

COMPLEXITY – CMS Tracker



References

- Sze, Physics of semiconductor devices, 2nd Edition
- Helmuth Spieler, Semiconductor Detector Systems
- Olaf Steinkamp, Experimental Methods of Particle Physics, 2011
- Enver Alagoz, Simulation and beam test measurements of the CMS pixel detector, PhD Thesis, 2009
- Daniela Bortoletto, An introduction to semiconductor detectors, Vienna Conference, VCI 2004