



J/ $\psi$  production in p-Pb collisions with



ALICE

at the LHC

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\*on behalf of the ALICE collaboration

Strangeness in Quark Matter 2013, Birmingham, 21-27.07.2013

# Outline

✧ Physics motivation

✧ Analysis

✧ Results

❖ Forward to Backward ratio  $R_{FB}^{J/\psi}$  integrated and vs  $p_T$ , vs  $y_{cms}$

❖ Nuclear modification factor  $R_{pPb}^{J/\psi}$  integrated and vs  $y$

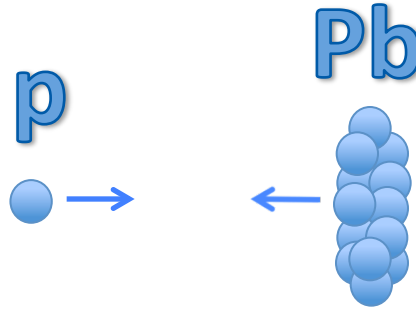
✧ Summary and outlook



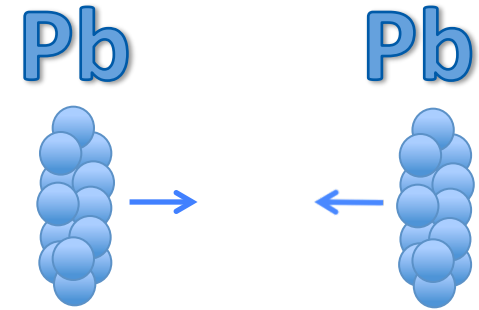
# Physics motivation



Elementary collision  
No nuclear matter effects



Cold nuclear matter effects –  
without Quark-Gluon Plasma (QGP)



Cold nuclear matter effects –  
without QGP  
+  
Hot nuclear matter effects –  
related to QGP formation

[\\*See the talk of  
Lizardo Valencia Palomo](#)

- To disentangle hot and cold nuclear matter (CNM) p-Pb measurements are needed as an intermediate step between Pb-Pb and benchmark pp collisions.



# Cold nuclear matter effects

In p-Pb different kinds of nuclear matter effects can be considered:

## ① Initial-state

✓ gluon shadowing[1] (or saturation[2]): at high energies gluons start shadowing each other (or recombining).

➤ At LHC energies large shadowing is expected.

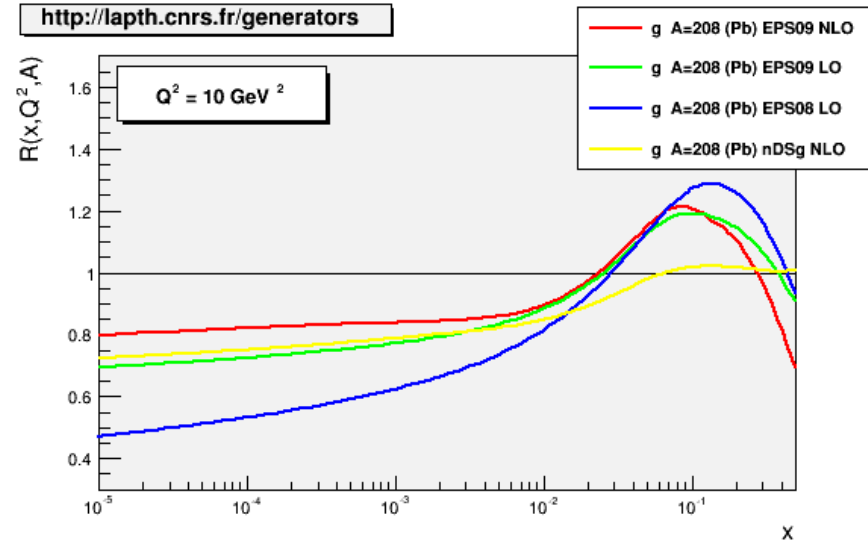
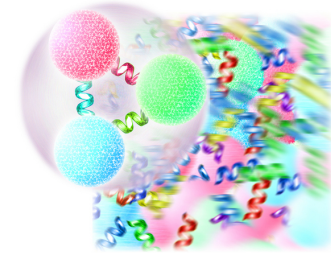
## ② Coherent energy loss [4]: gluon radiates a soft gluon.

➤ The amount of medium-induced gluon radiation defines the strength of the  $J/\psi$  suppression.

## ③ Final-state

✓ nuclear absorption:  $J/\psi$  pre-resonant state destruction by colliding nucleons.

➤ At the LHC at mid- and forward rapidity in p-Pb the  $c\bar{c}$  pair spends a very short time within cold nuclear matter, due to the large Lorentz gamma of the colliding nuclei. Consequently, nuclear absorption is then expected to be negligible [5].



[1] K. Eskola et al., JHEP 0904:065 (2009)

[2] D. E. Kharzeev et al., arXiv:1205.1554 (2012);

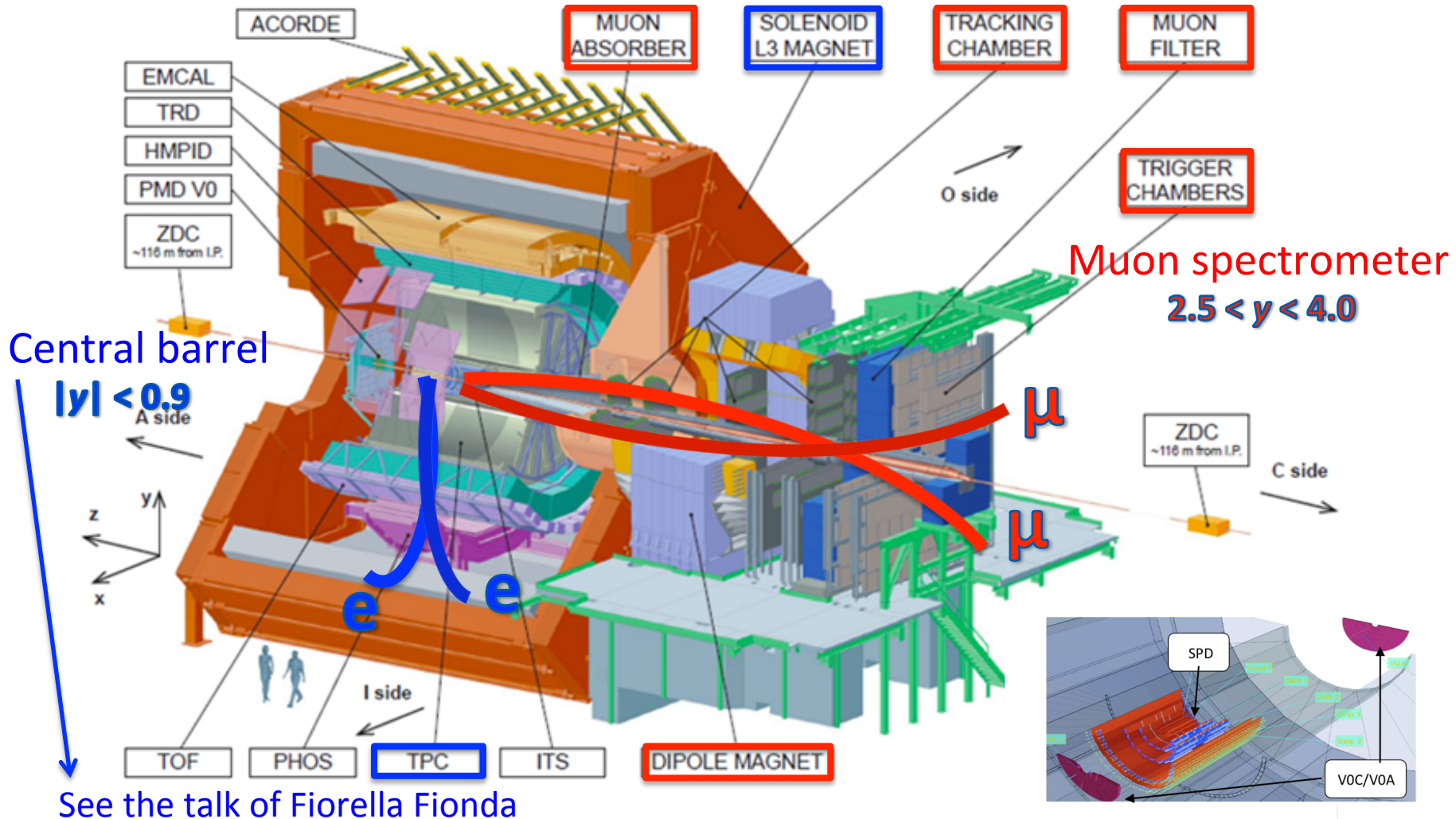
F. Dominguez et al. arXiv:1109.1250 (2012)

[3] R. Vogt Phys.Rev. C81 (2010) 044903

[4] F. Arleo, S. Peigne, arXiv1204.4609 (2012)

[5] Lourenco et al., JHEP 0902:014, 2009

# ALICE detector



# Event selection and analysis cuts



- Event selection

- ✓ MB trigger: Coincidence of the two sides of VZERO:  $2.8 < \eta < 5.1$ ,  $-3.7 < \eta < -1.7$
- ✓ MB trigger efficiency  $\sim 99\%$  for NSD events
- ✓ Rejection of beam-gas and electromagnetic interactions
- ✓ SPD used for vertex determination

- Dimuon trigger

- ✓ Coincidence of minimum bias (MB) interaction with two opposite sign muon tracks detected in the trigger chambers of the Muon spectrometer

- The following cuts (standard for  $J/\psi$  analysis) were also applied:

- ✓ Muon trigger matching
- ✓  $-4 < \eta_{\mu} < -2.5$
- ✓  $17.6 \text{ cm} < R_{\text{abs}} < 89.5 \text{ cm}$ , where  $R_{\text{abs}}$  – track radial position at the absorber end
- ✓ Unlike sign dimuon
- ✓  $2.5 < y_{\mu\mu}^{\text{lab}} < 4$

# Main observables ( $R_{pPb}$ , $R_{PbPb}$ )

## ➤ Nuclear modification factor $R_{pPb}$ and $R_{PbPb}$

$$R_{pPb}^{J/\psi} = \frac{Y_{pPb}}{\langle T_{pPb} \rangle \sigma_{pp}^{J/\psi \rightarrow \mu^+ \mu^-}},$$

$$Y_{pPb} = \frac{N_{J/\psi \rightarrow \mu^+ \mu^-}}{(A \times \varepsilon) N_{MB}}$$

$R_{pPb}$  and  $R_{PbPb}$  are computed in the range  $2.5 < y_{lab} < 4$

$T_{pPb} = 0.0983 \pm 0.0034 \text{ mb}^{-1}$  – nuclear overlap function

See talk of Andreas Morsch

## ➤ Shift in $y_{cms}$ and rapidity coverage

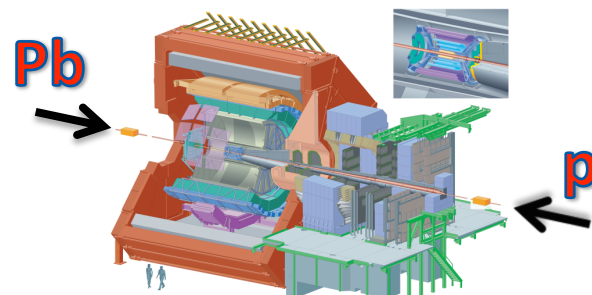
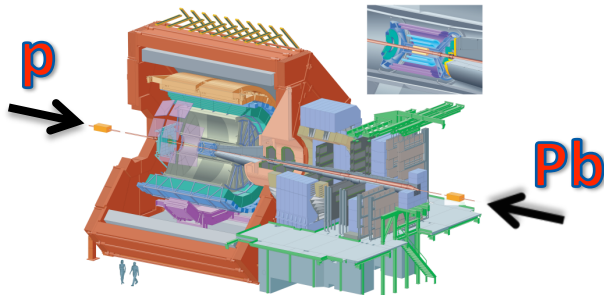
LHC beam asymmetry ( $E_{Pb} = 1.58 \cdot A \text{ TeV}$ ,  $E_p = 4 \text{ TeV}$ )  $\Rightarrow |\Delta y|_{cms} = 0.5 \text{ Log}(Z_{Pb} A_p / Z_p A_{Pb}) = 0.465$

**p-Pb:  $2.03 < y_{cms} < 3.53$**

$8.1 \cdot 10^{-5} > x_{Bjorken} > 1.8 \cdot 10^{-5}$

**Pb-p:  $-4.46 < y_{cms} < -2.96$**

$5.3 \cdot 10^{-2} > x_{Bjorken} > 1.2 \cdot 10^{-2}$



# Main observables ( $R_{FB}$ )

## ➤ Forward to Backward ratio $R_{FB}$

$$R_{FB}^{J/\psi} = \frac{R_{pPb}}{R_{Pbp}}$$

✧  $R_{FB}$  is computed in the  $y_{cms}$  range common to both p-Pb and Pb-p:  $2.96 < y_{cms} < 3.53$  which corresponds to the following range in lab.system:

$$\text{p-Pb: } 3.43 < y_{lab} < 4 \\ 3.2 \cdot 10^{-5} > x_{Bjorken} > 1.8 \cdot 10^{-5}$$

$$\text{Pb-p: } -3.07 < y_{lab} < -2.5 \\ 2.1 \cdot 10^{-2} > x_{Bjorken} > 1.2 \cdot 10^{-2}$$

In that case  $T_{ppb}$  and the pp cross-section cancel out in the ratio:

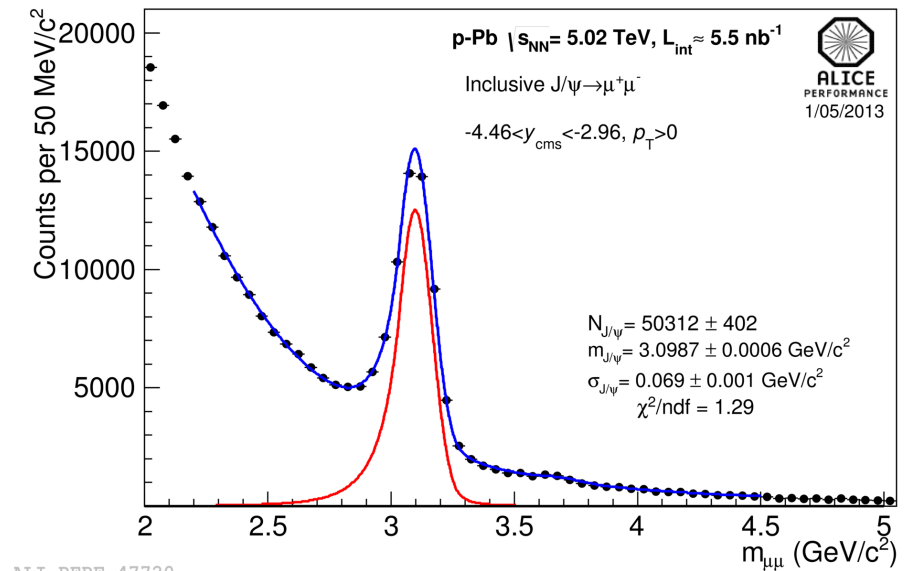
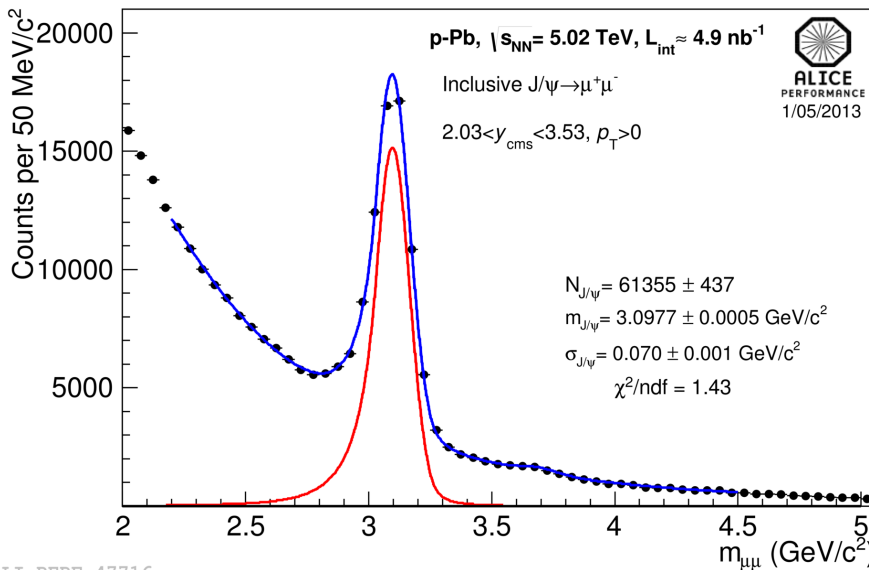
$$R_{FB}^{J/\psi} = \frac{Y_{pPb}^{Forward}}{Y_{pPb}^{Backward}} = \frac{N_{J/\psi \rightarrow \mu^+ \mu^-}^{Forward}}{(Acc \times \epsilon)^{Forward} N_{MB}^{Forward}} \times \frac{(Acc \times \epsilon)^{Backward} N_{MB}^{Backward}}{N_{J/\psi \rightarrow \mu^+ \mu^-}^{Backward}}$$



# Signal extraction

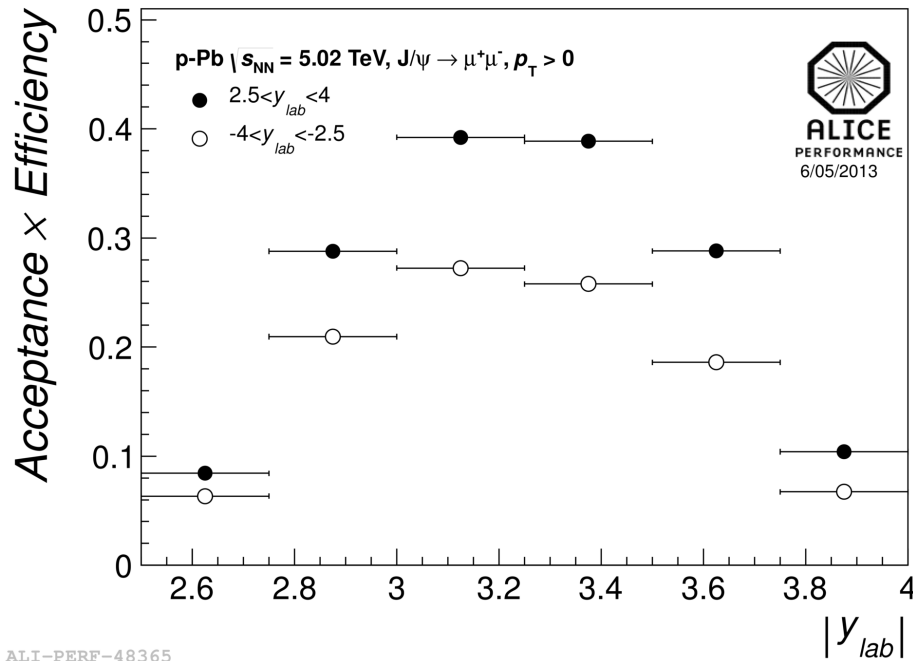
Signal extraction (and its syst. unc.) is based on fits of dimuon inv.mass distribution by varying:

- ① **Signal shape:** Extended Crystal Ball (CB2) or other pseudo-Gaussian functions (tails tuned on the corresponding Monte Carlo (MC))
- ② **Background shape:** Variable Width Gaussian (VWG) or Pol2\*Exp (or Pol4\*Exp)
- ③ **Fitting range**

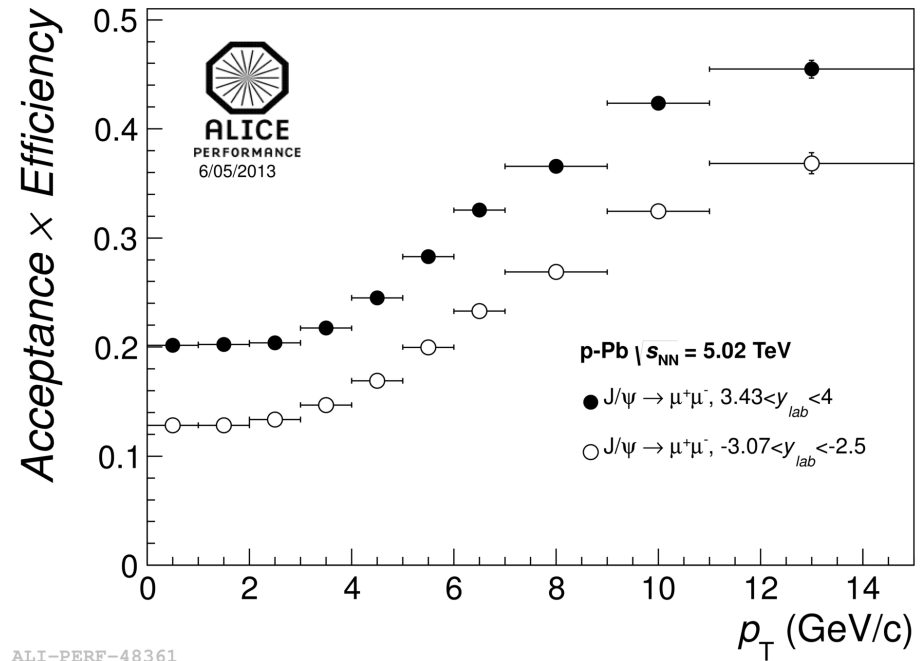


These plots are examples of the fit with **CB2+VWG**.

# Acceptance x Efficiency



ALI-PERF-48365



ALI-PERF-48361

➤ Average  $J/\psi$  acceptance x efficiency:

p-Pb:  $\sim 25\%$  in  $2.03 < y_{cms} < 3.53$

Pb-p:  $\sim 17\%$  in  $-4.46 < y_{cms} < -2.96$

- Difference in AccxEff between p-Pb and Pb-p is due to different efficiency of detector in two periods of data-taking

➤ Systematic uncertainties on acceptance inputs uncorrelated vs  $p_T$ ,  $y$  and collision system (different physics)

# Summary on the syst. uncertainties

Source of systematic uncertainty:	Systematic uncertainty
Signal extraction	1-4%
Nuclear thickness function $T_{\text{pPb}}$	3.4%
Acceptance inputs	1-3.5%
Tracking efficiency	4-6%
Trigger efficiency	3%
Matching efficiency	1%
Normalization dimuon-MB trigger	1%
<b>Total syst. uncertainty</b>	<b>7-12%</b>

\*(ranges correspond to values obtained in  $y$  or  $p_T$  bins)

# $d^2\sigma_{J/\psi}/dydp_T$

$$\sigma_{J/\psi \rightarrow \mu^+\mu^-}^{pPb} = \frac{N_{J/\psi \rightarrow \mu^+\mu^-}}{L_{\text{int}} \times \text{Acc} \times \varepsilon \times BR_{J/\psi \rightarrow \mu^+\mu^-}}$$

$$L_{\text{int}} = \frac{N_{MB}}{\sigma_{MB}}$$

✓  $\sigma_{MB}$  obtained using VdM scans:

**2.08 b ± 3.4%** for **p-Pb** period

**2.12 b ± 3.2%** for **Pb-p** period

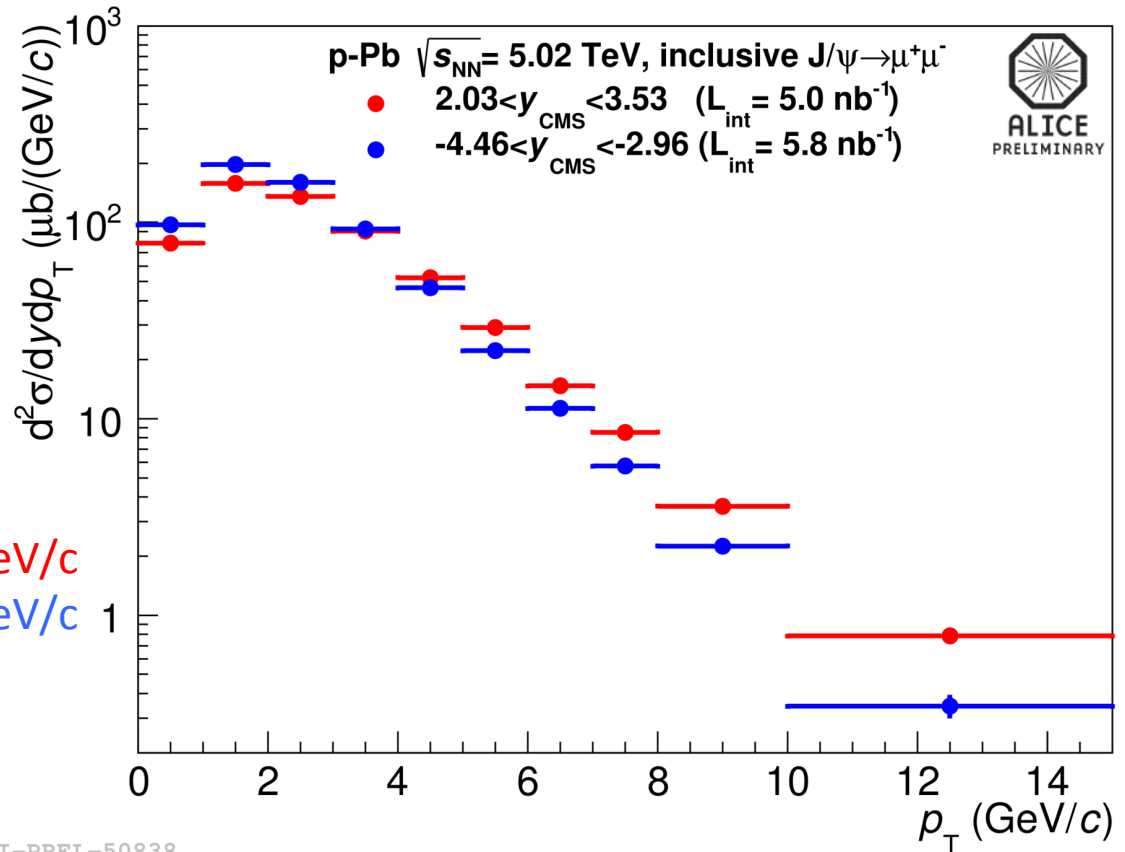
$L_{\text{int}}^{pPb} = 5.0 \text{ nb}^{-1}$ ;  $L_{\text{int}}^{Pbp} = 5.8 \text{ nb}^{-1}$

One can calculate  $\langle p_T \rangle |_{0-15 \text{ GeV}/c}$ :

➤  $\langle p_T \rangle = 2.77 \pm 0.01^{\text{stat.}} \pm 0.02^{\text{syst.}} \text{ GeV}/c$

➤  $\langle p_T \rangle = 2.47 \pm 0.01^{\text{stat.}} \pm 0.02^{\text{syst.}} \text{ GeV}/c$

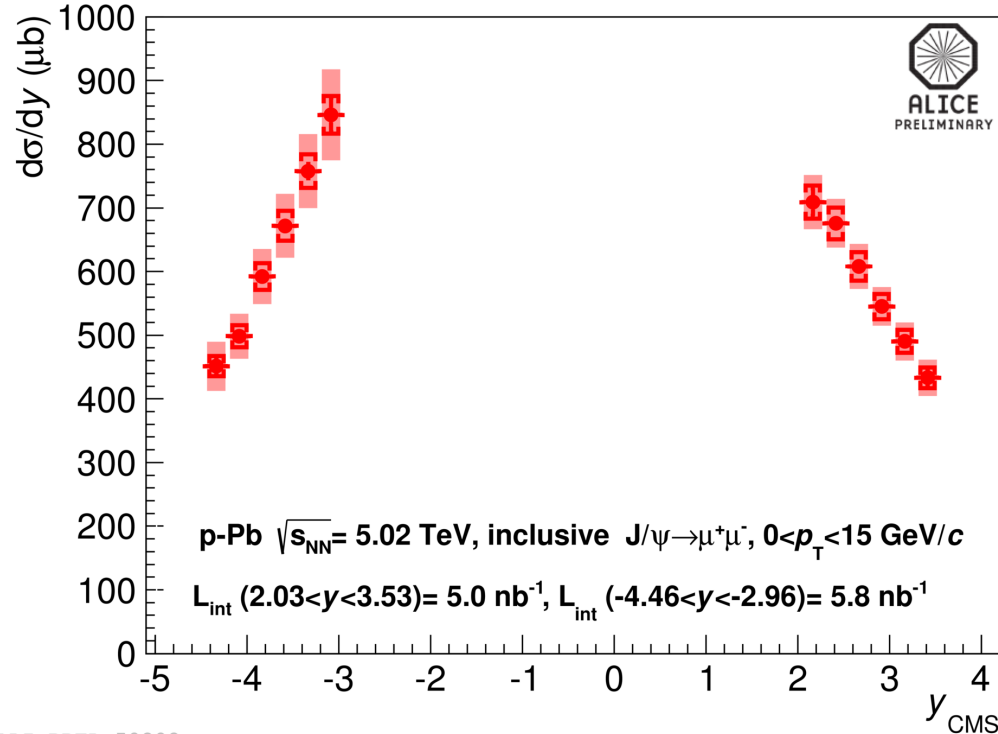
✓  $\sigma_{MB}$ : MB condition related to signal in VZERO



ALI-PREL-50838

I. Lakomov, SQM 2013, 21-27.07.2013

# $d\sigma_{J/\psi}/dy$



ALI-PREL-50808

$$\frac{d\sigma_{pPb}}{dy} = 588 \pm 4 \text{ (stat.)} \pm 38 \text{ (syst.) } \mu b$$

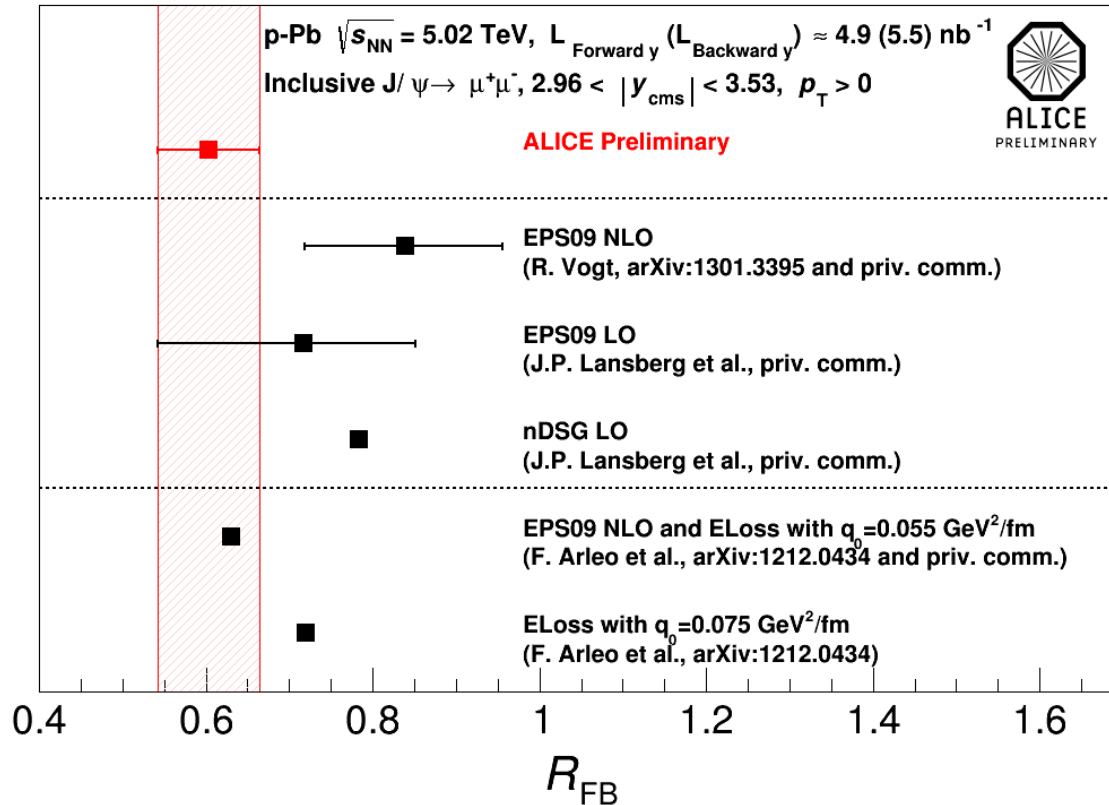
$$\frac{d\sigma_{Pbp}}{dy} = 644 \pm 5 \text{ (stat.)} \pm 51 \text{ (syst.) } \mu b$$

- ✓ **Correlated uncertainties** (brackets): luminosity, normalization factor, BR
- ❖ Luminosity is correlated within p-Pb or Pb-p, but not within the two systems
- ✓ **Uncorrelated uncertainties** (filled boxes): matching, trigger efficiency, tracking, acc. inputs, signal extraction
- ✓ **Statistical uncertainties** (line)

➤ Cross-sections are higher in the backward rapidity region (Pb-p).

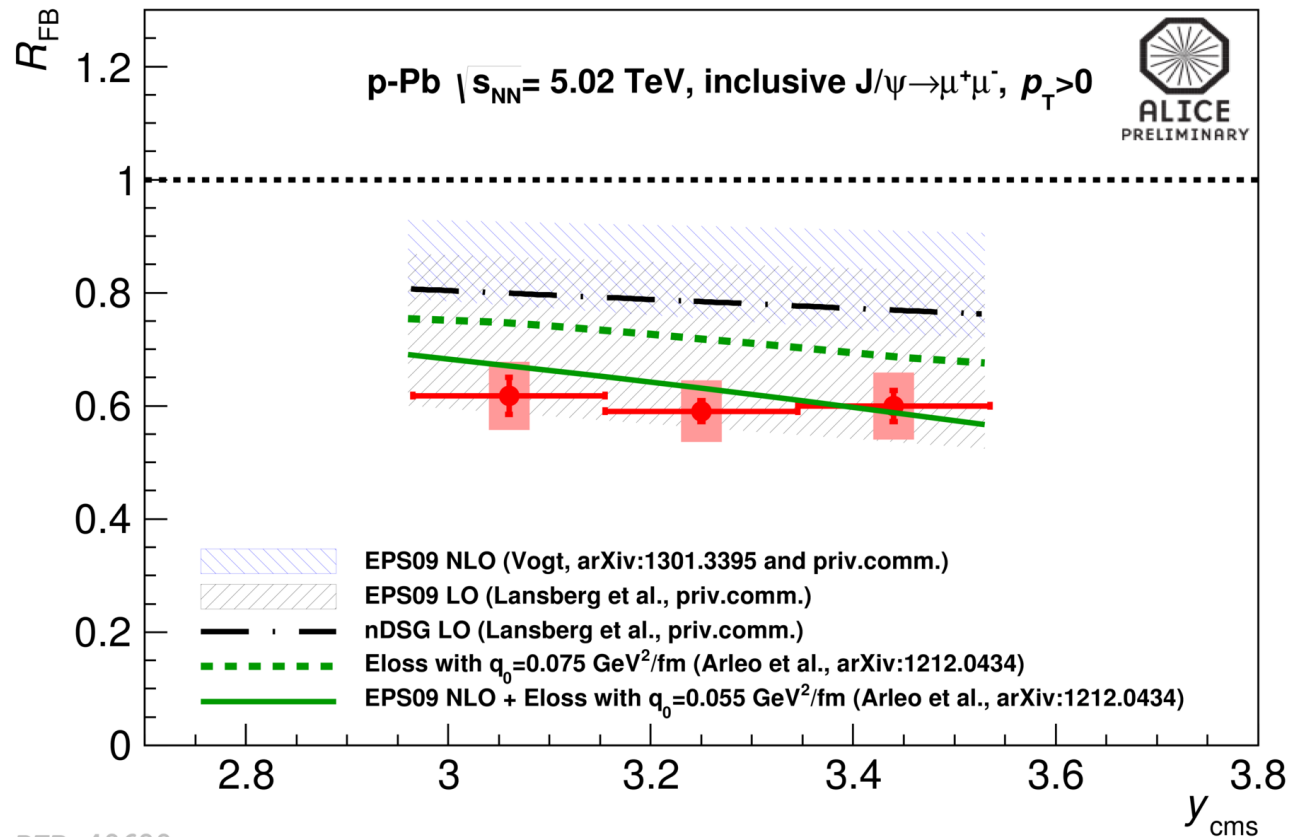
# Integrated $R_{FB}$

$$R_{FB} = 0.60 \pm 0.01 \text{ (stat.)} \pm 0.06 \text{ (syst.)}$$



- The uncertainty is small
- Pure shadowing slightly overestimates the data
- Model including energy loss contribution is rather good

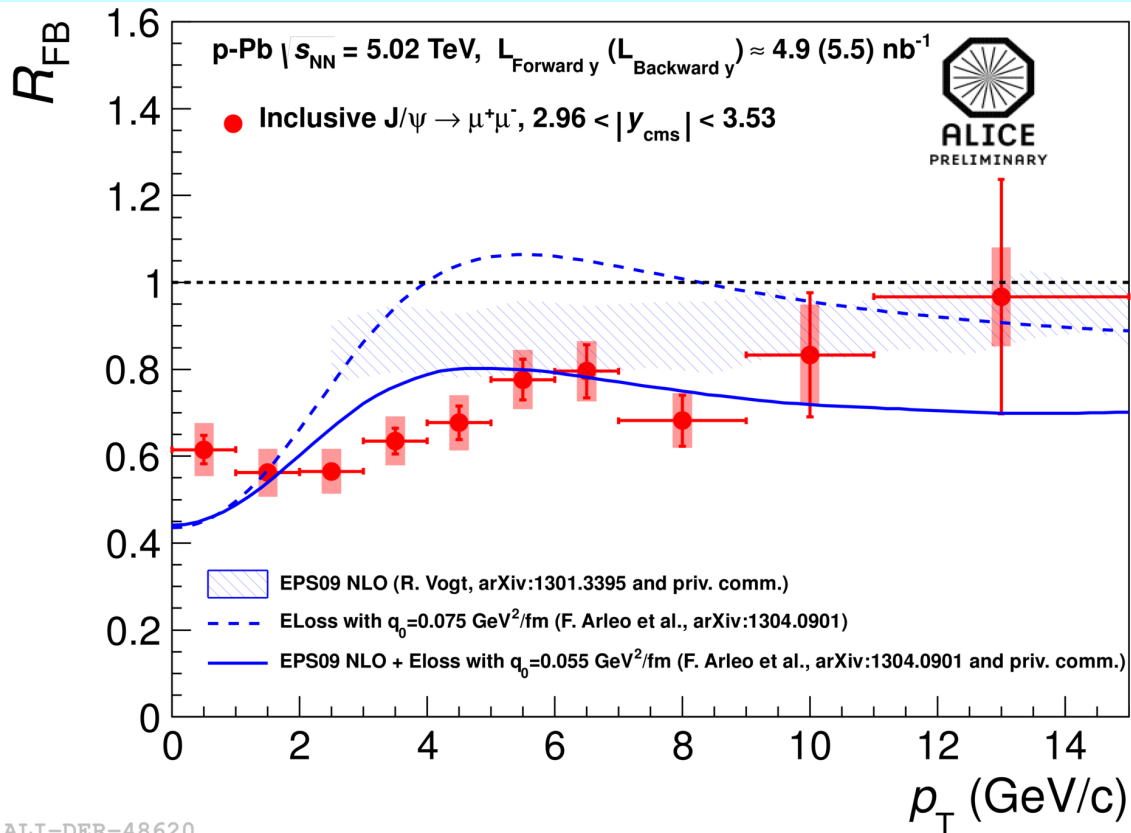
# $R_{FB}$ vs rapidity



ALI-DER-48628

- Comparison with theoretical models confirms previous observations done on the  $y$ -integrated results.
- Calculations including both shadowing and energy loss seems consistent with the data

# $R_{FB}$ vs $p_T$



ALI-DER-48620

- A sizeable  $p_T$ -dependence of  $R_{FB}$  is seen.
- Stronger suppression is found at low  $p_T$ .
- Theoretical models including energy loss show strong nuclear matter effects at low  $p_T$  in fair agreement with the data
- The observed  $p_T$ -dependence is smoother than expected in coherent energy loss models



# pp-reference

Phenomenological interpolation of the inclusive  $J/\psi$  x-section to pp collisions at  $\sqrt{s_{NN}}=5.02$  TeV from CDF, RHIC and LHC (2.76 and 7 TeV) based on the paper from arXiv:1103.2394v3.

## ① Energy dependence: pp cross-section at mid-rapidity

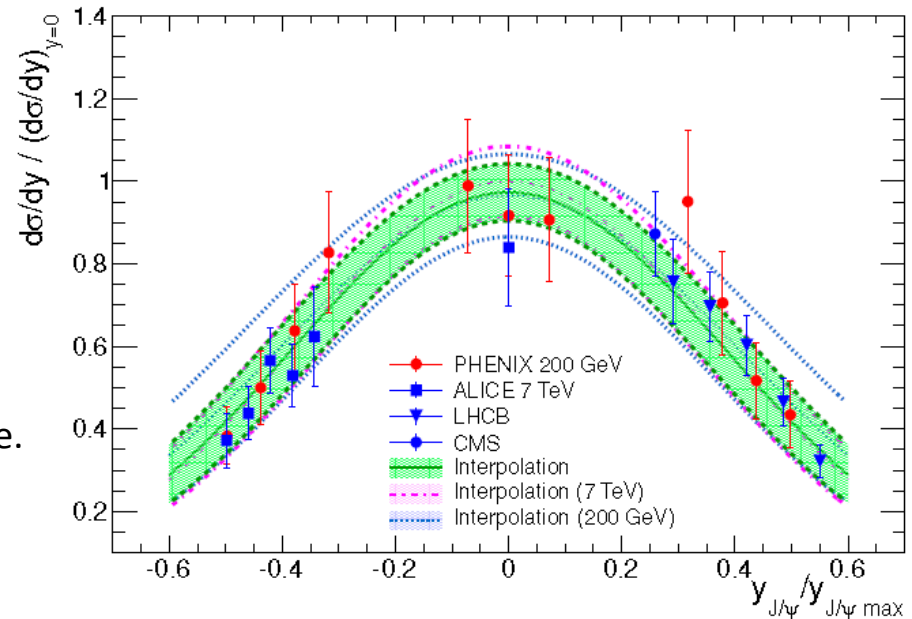
Calculations performed using a Monte Carlo toy.

Parametrization with a power-law shape.

$$\left. \frac{d\sigma^{pp}}{dy} \right|_{y=0} = 362 \pm 6(stat.)^{+55(syst.)}_{-37(syst.)} nb$$

## ② Rapidity dependence

Based on a universal, energy independent gaussian shape.



## ③ Systematic uncertainties

Evaluated within  $2.5\sigma$  in order to include most of the uncertainties from FONLL and CEM LO interpolation.

$$d\sigma^{pp} / dy (2.03 < y_{cms} < 3.53) = 231^{+41(syst.)}_{-32(syst.)} nb$$

$$d\sigma^{pp} / dy (-4.46 < y_{cms} < -2.96) = 159^{+40(syst.)}_{-27(syst.)} nb$$

# Summary on the systematics

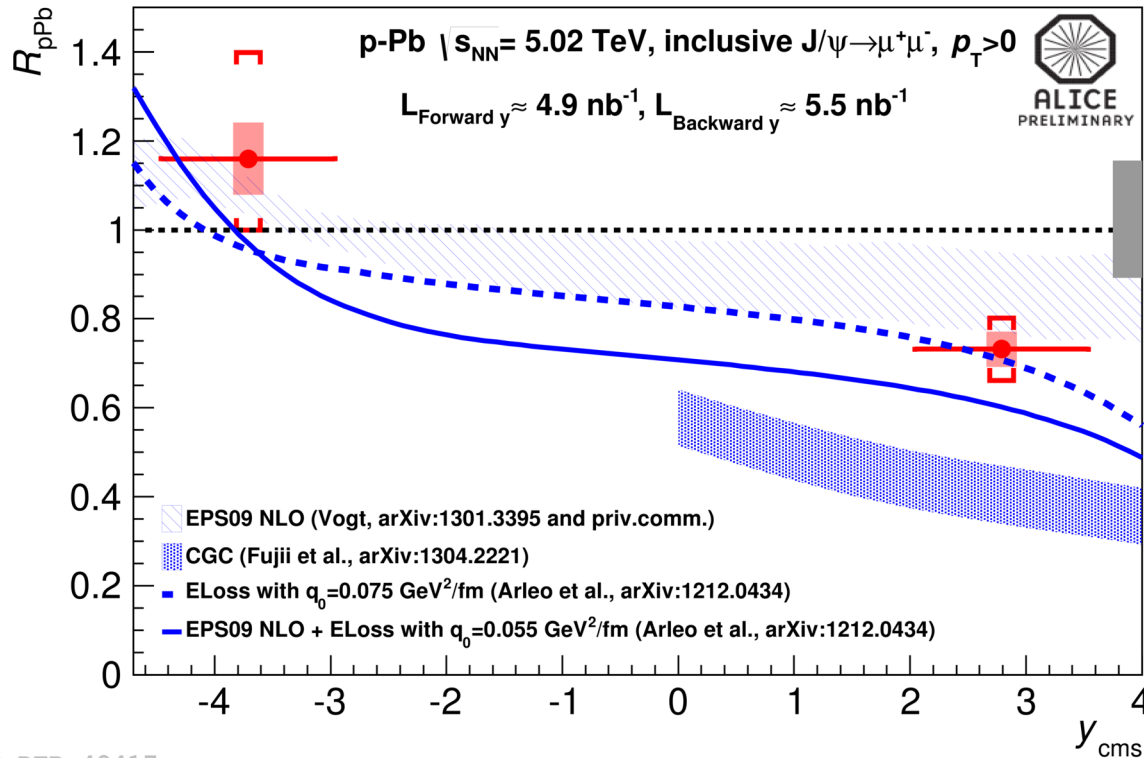
Source of systematic uncertainty:	Systematic Uncertainty
Signal extraction	1-4%
Nuclear thickness function $T_{\text{pPb}}$	3.4%
Acceptance inputs	1-3.5%
Tracking efficiency	4-6%
Trigger efficiency	3%
Matching efficiency	1%
Normalization dimuon-MB trigger	1%
pp reference @ $y=0$ , $\sqrt{s} = 5.02$ TeV	10-15%
y-dependence of pp interpolation @ $\sqrt{s}_{NN} = 5.02$ TeV	10-20%
<b>Total syst. uncertainty (excluding pp interpol.)</b>	<b>7-12%</b>

\*(ranges correspond to values obtained in  $y$  or  $p_T$  bins)

# $R_{pPb}$ and $R_{PbPb}$ integrated

$$R_{pPb} (2.03 < y_{cms} < 3.53) = 0.732 \pm 0.005(\text{stat}) \pm 0.059(\text{syst}) + 0.131(\text{syst. ref}) - 0.101(\text{syst.ref})$$

$$R_{pPb} (-4.46 < y_{cms} < -2.96) = 1.160 \pm 0.010(\text{stat}) \pm 0.096(\text{syst}) + 0.296(\text{syst. ref}) - 0.198(\text{syst.ref})$$



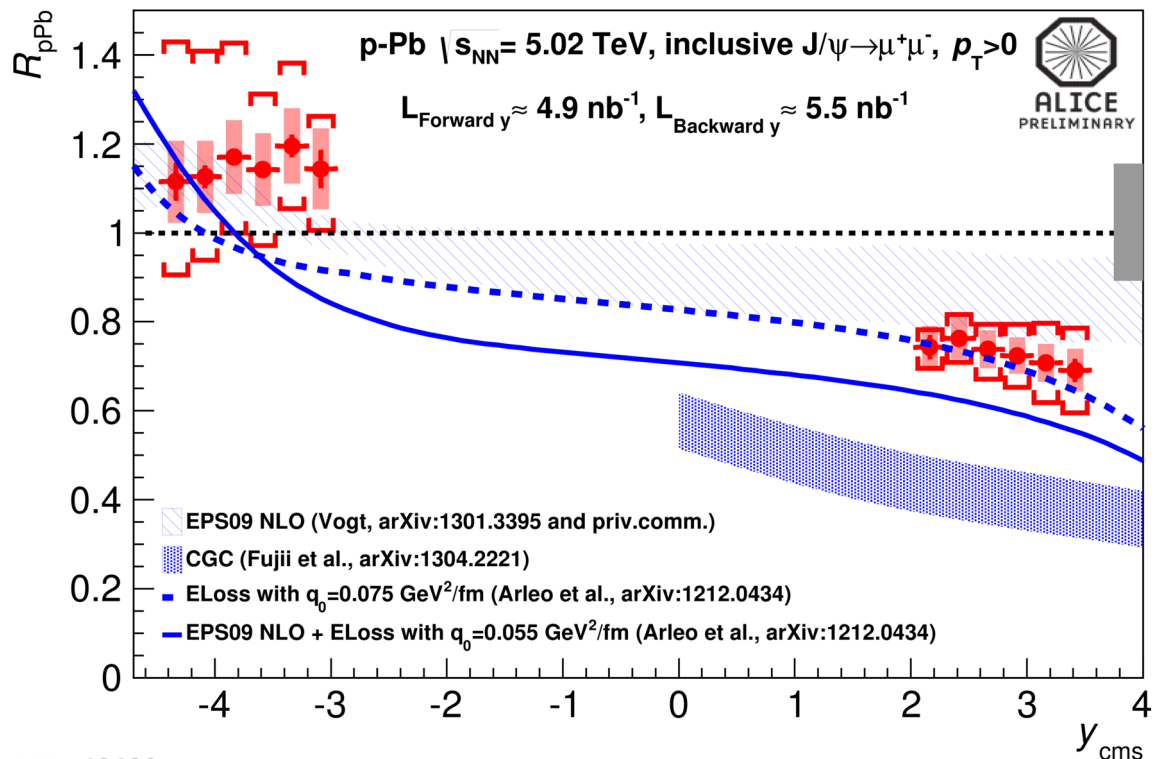
Error bars:

- ✓ boxes around the points: uncorrelated
- ✓ [ ]: partially correlated
- ✓ grey box around unity: fully correlated

ALI-DER-48417

- Large uncertainty (correlated and uncorrelated) from pp interpolation
- At forward rapidity, data in-between shadowing and energy loss models
- Color Glass Condensate (CGC) model underestimates the data

# $R_{pPb}$ and $R_{Pbp}$ vs rapidity



Error bars:

- ✓ boxes around the points: uncorrelated
- ✓ [ ]: partially correlated
- ✓ grey box around unity: fully correlated

ALI-DER-48480

- At backward rapidity, models including coherent parton energy loss show a slightly steeper pattern than the one observed in data
- Results dominated by a large uncertainty from pp interpolation

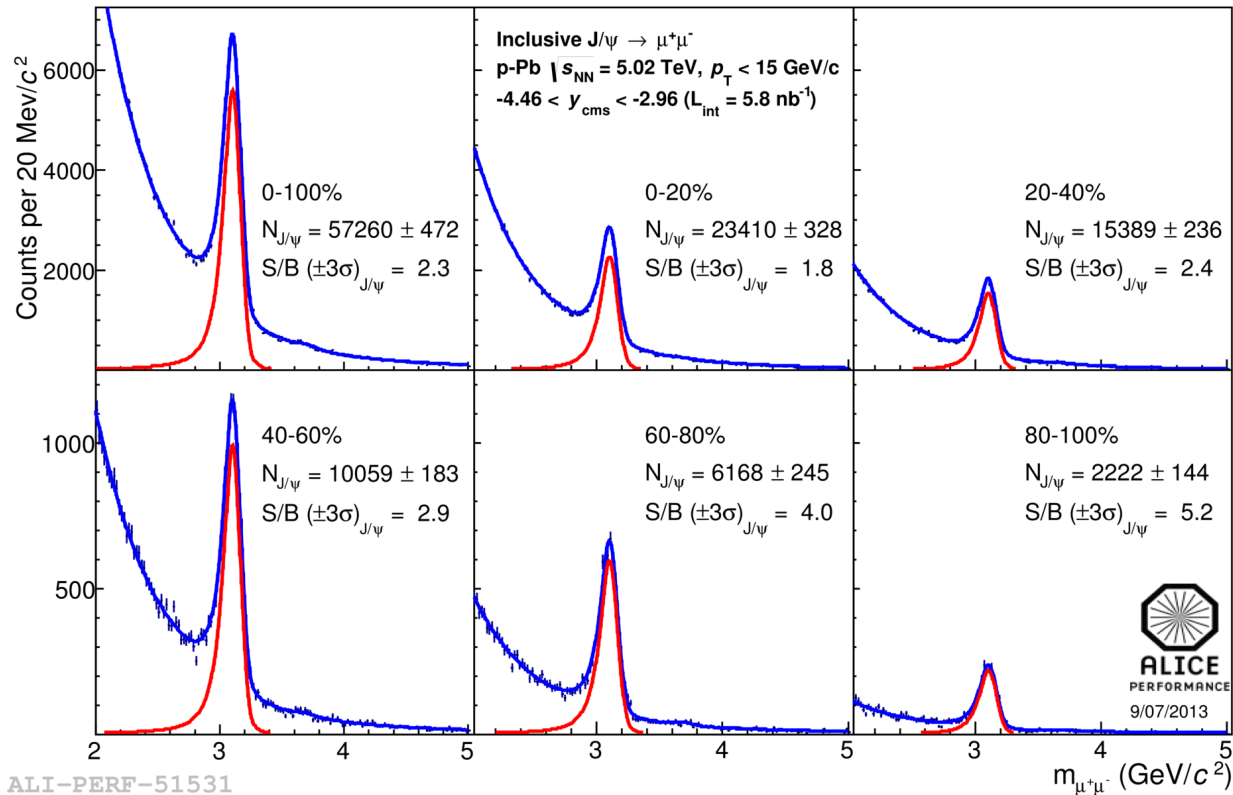
# Summary...

ALICE has measured **inclusive  $J/\psi$  production in p-Pb run in backward and forward rapidity regions at  $\sqrt{s_{NN}} = 5.02$  TeV**. Many interesting results are obtained:

- Measured **strong  $p_T$  dependence of  $R_{FB}$  with a decrease at low  $p_T$**  is in a fair agreement with models including coherent energy loss contribution.
- **$R_{pPb}$  and  $R_{Pbp}$  show an increase of suppression towards forward rapidity** in agreement with energy loss model and/or shadowing model EPS09 NLO.
- pure nuclear shadowing and/or energy loss seem to reasonably describe the data, indicating that **final state absorption may indeed be negligible at LHC energies**

# ...and outlook

- Many other interesting results are under study:  $R_{pPb}$  vs centrality,  $\Psi(2S)$  yield...

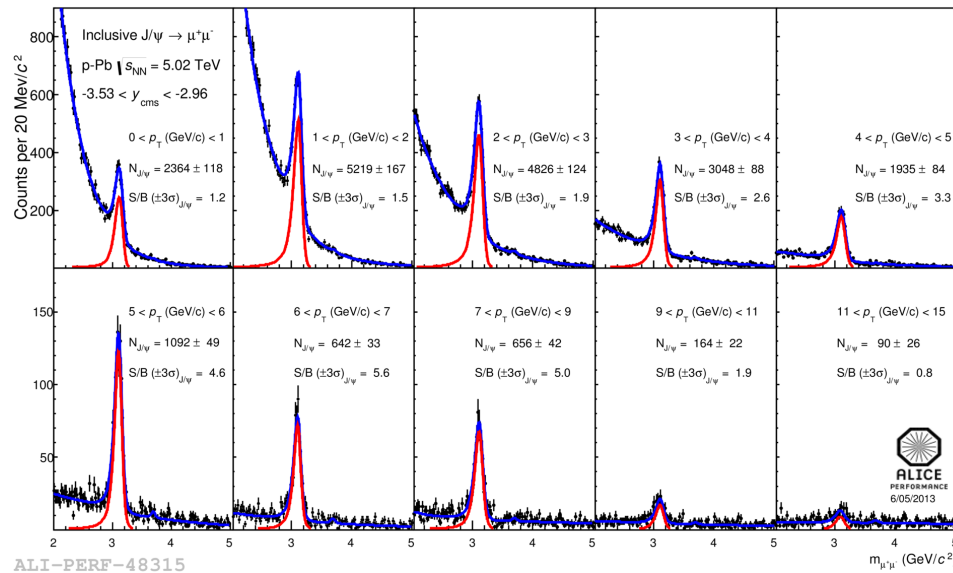
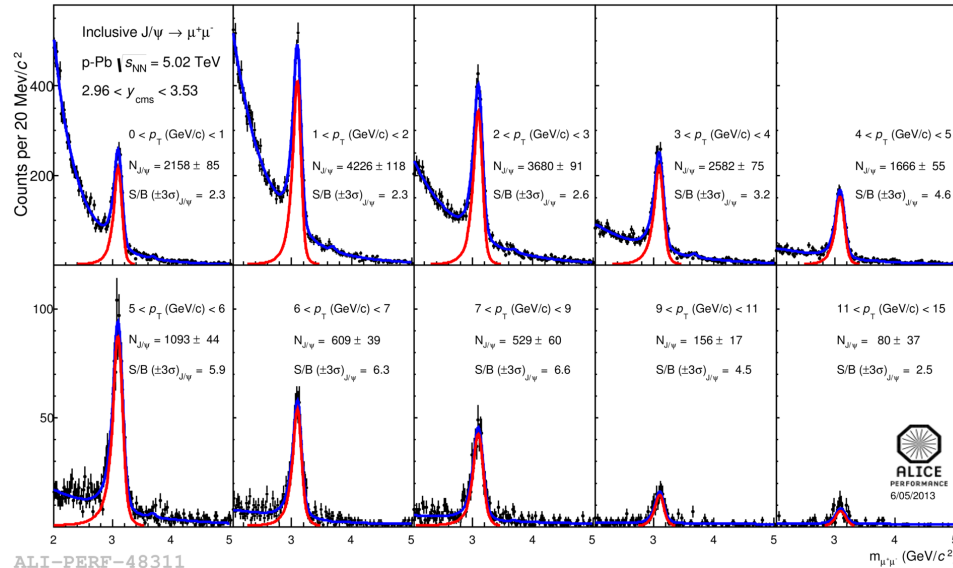


Stay tuned...

➤ Thank you for your attention!

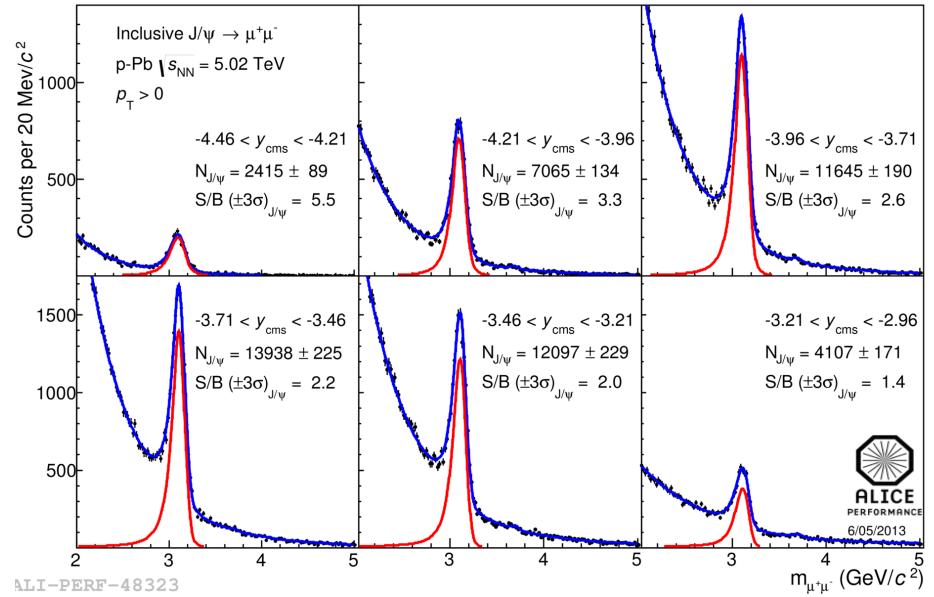
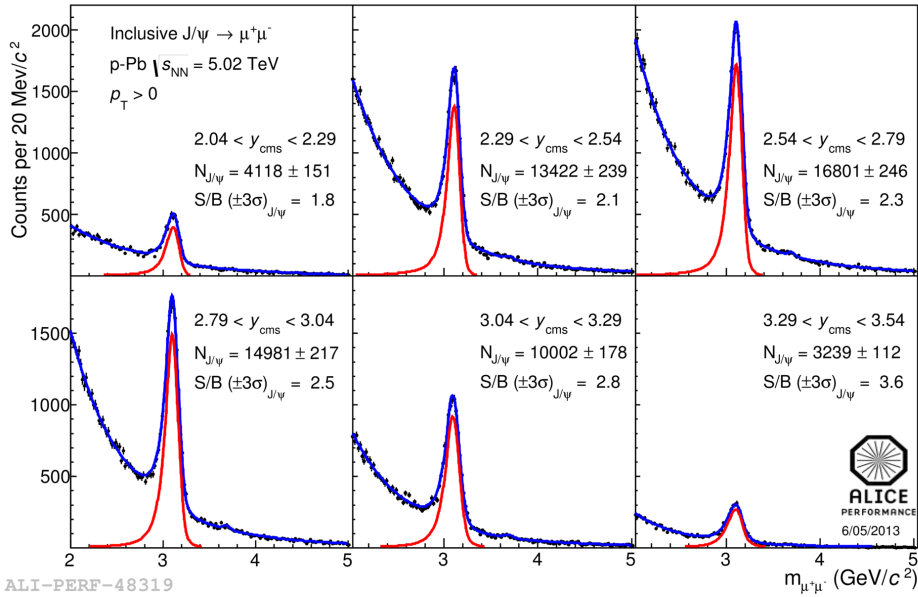
# Backup slides

# Signal extraction in $p_T$ bins



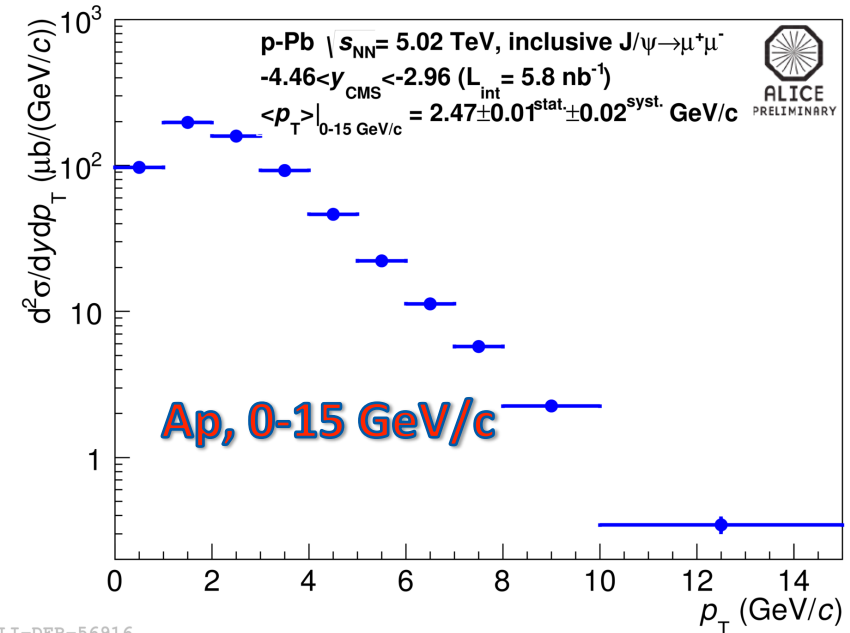
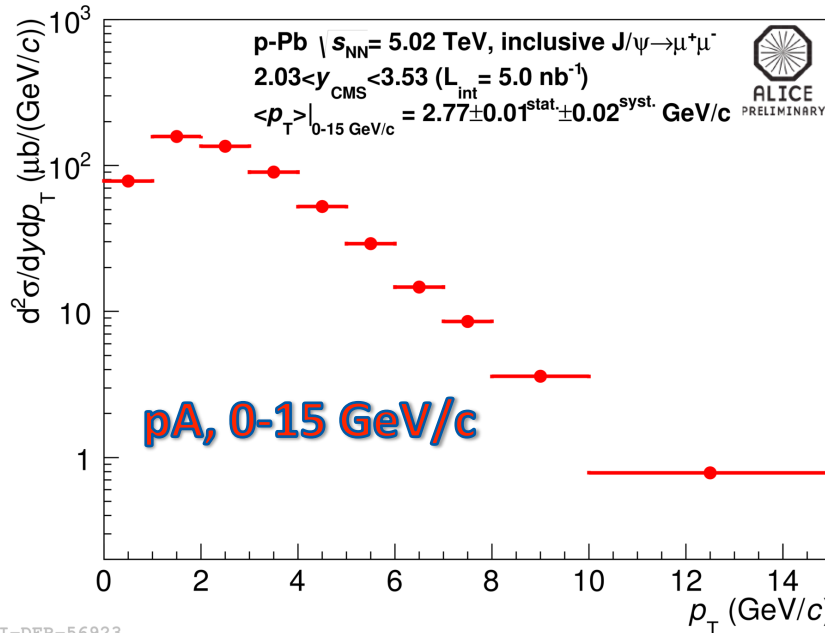


# Signal extraction in $y$ bins



# $d^2\sigma_{J/\psi}/dp_T dy$ and $\langle p_T \rangle$ in full $y$ -range

From the  $d^2\sigma/dydp_T$  distributions one can calculate the mean  $p_T$  in the full  $y$ -range



ALI-DER-56923

**pp, 0-8 GeV/c**

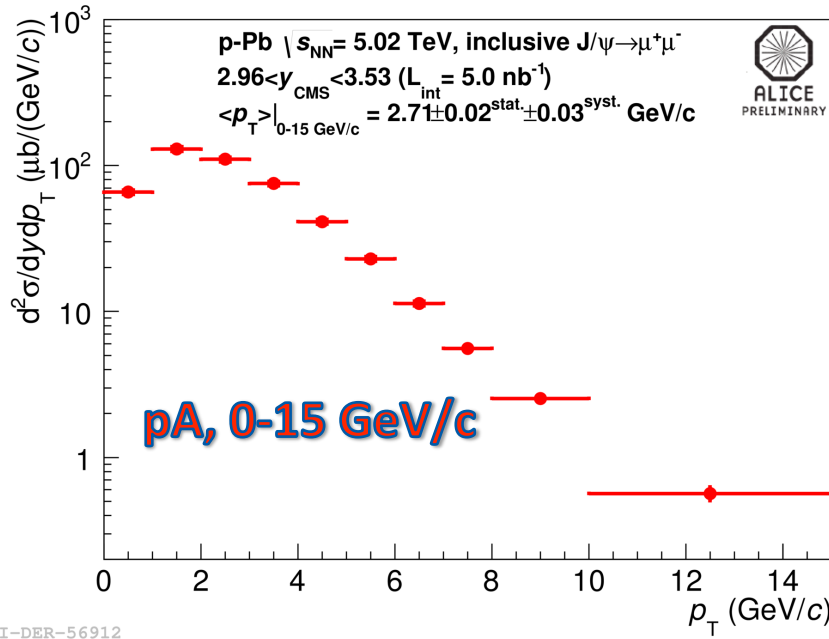
	$\langle p_t \rangle$ (GeV/c)	$\langle p_t^2 \rangle$ (GeV <sup>2</sup> /c <sup>2</sup> )
$\sqrt{s} = 2.76$ TeV $2.5 < y < 4$	$2.28 \pm 0.07 \pm 0.04$	$7.06 \pm 0.40$
$\sqrt{s} = 7$ TeV, $ y  < 0.9$	$2.72 \pm 0.21 \pm 0.28$	$10.02 \pm 1.4$
$\sqrt{s} = 7$ TeV, $2.5 < y < 4$	$2.44 \pm 0.09 \pm 0.06$	$8.32 \pm 0.50$

ALI-DER-56916

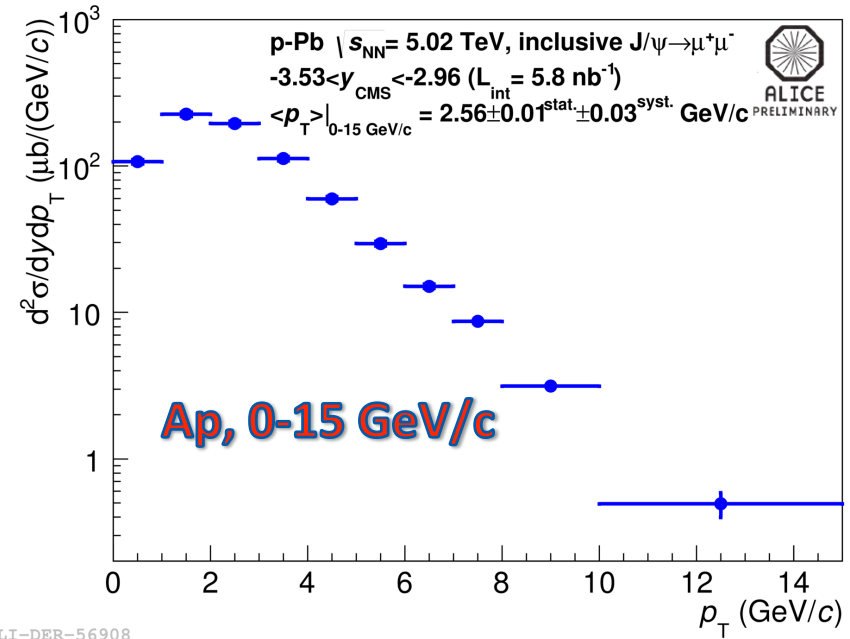
- $\langle p_T \rangle |_{0-15 \text{ GeV/c}} = 2.77 \pm 0.01^{\text{stat.}} \pm 0.02^{\text{syst.}}$  GeV/c in p-Pb
- $\langle p_T \rangle |_{0-15 \text{ GeV/c}} = 2.47 \pm 0.01^{\text{stat.}} \pm 0.02^{\text{syst.}}$  GeV/c in Pb-p
- $\langle p_T \rangle$  is higher in p-Pb than in pp at  $\sqrt{s}_{\text{NN}} = 7$  TeV
- $\langle p_T \rangle$  in Pb-p is close to the one in pp at  $\sqrt{s}_{\text{NN}} = 7$  TeV

# $d^2\sigma_{J/\psi}/dp_T dy$ and $\langle p_T \rangle$ in common $y$ -range

From the  $d^2\sigma/dydp_T$  distributions one can calculate the mean  $p_T$  in the common  $y$ -range



ALI-DER-56912

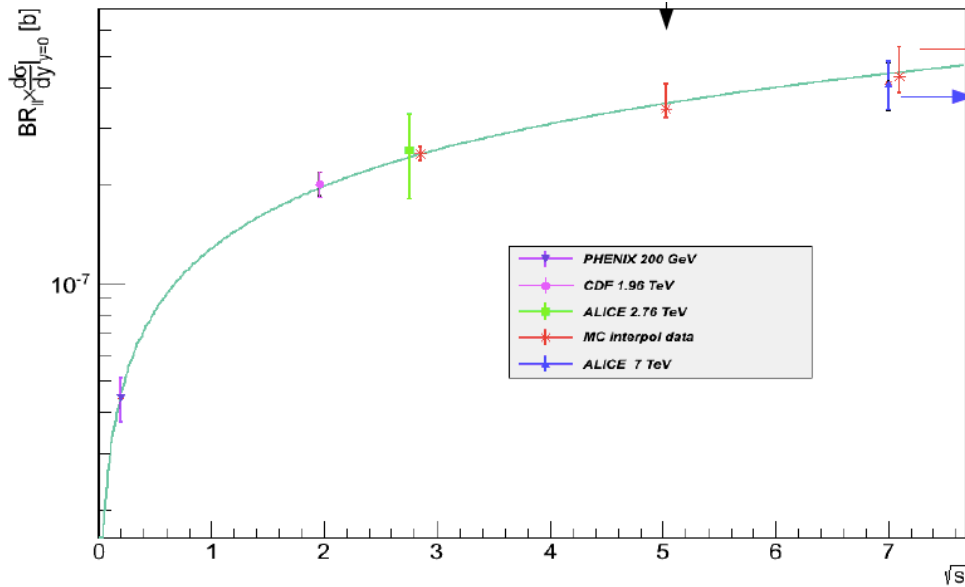


ALI-DER-56908

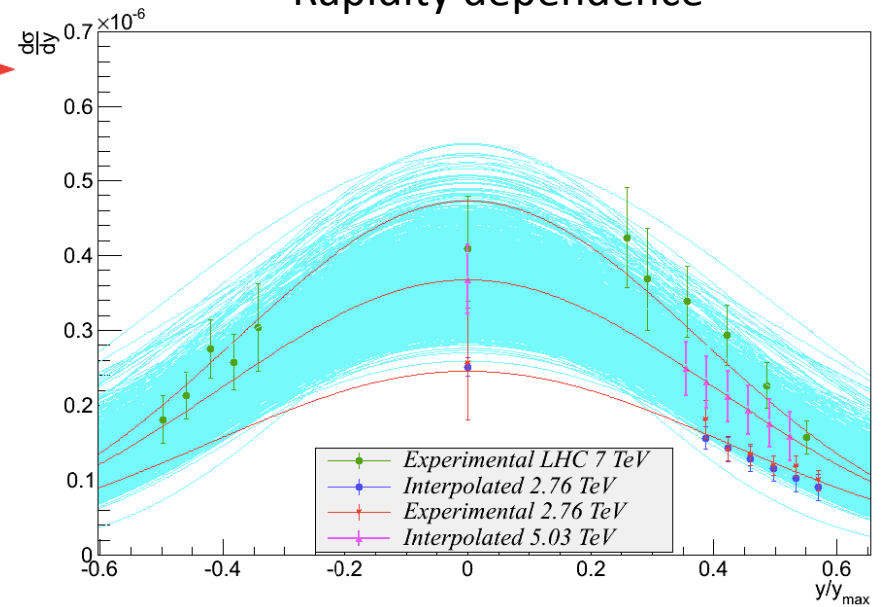
- $\langle p_T \rangle |_{0-15 \text{ GeV/c}} = 2.71 \pm 0.02^{\text{stat}} \pm 0.03^{\text{syst}} \text{ GeV/c}$  in p-Pb  
 -> compare to  $\langle p_T \rangle |_{0-15 \text{ GeV/c}} = 2.77 \pm 0.01^{\text{stat}} \pm 0.02^{\text{syst}} \text{ GeV/c}$  in the full  $y$ -range
- $\langle p_T \rangle |_{0-15 \text{ GeV/c}} = 2.56 \pm 0.01^{\text{stat}} \pm 0.03^{\text{syst}} \text{ GeV/c}$  in Pb-p  
 -> compare to  $\langle p_T \rangle |_{0-15 \text{ GeV/c}} = 2.47 \pm 0.01^{\text{stat}} \pm 0.02^{\text{syst}} \text{ GeV/c}$  in the full  $y$ -range

# Interpolation of $\sigma_{J/\psi}^{pp}$ at $\sqrt{s_{NN}}=5.02$ TeV

Energy dependence



Rapidity dependence



# Signal shapes functions

## CB extended

$$f(x; N, \bar{x}, \sigma, \alpha, n, \alpha', n') = N \cdot \begin{cases} \exp\left(-\frac{(x-\bar{x})^2}{2\sigma^2}\right), & \text{for } \frac{(x-\bar{x})}{\sigma} > -\alpha \\ A \cdot \left(B - \frac{x-\bar{x}}{\sigma}\right)^{-n}, & \text{for } \frac{(x-\bar{x})}{\sigma} \leq -\alpha \\ C \cdot \left(D + \frac{x-\bar{x}}{\sigma}\right)^{-n'}, & \text{for } \frac{(x-\bar{x})}{\sigma} \geq \alpha' \end{cases}$$

$$A = \left(\frac{n}{|\alpha|}\right)^n \cdot \exp\left(-\frac{|\alpha|^2}{2}\right)$$

$$B = \frac{n}{|\alpha|} - |\alpha|$$

$$C = \left(\frac{n'}{|\alpha'|}\right)^{n'} \cdot \exp\left(-\frac{|\alpha'|^2}{2}\right)$$

$$D = \frac{n'}{|\alpha'|} - |\alpha'|$$

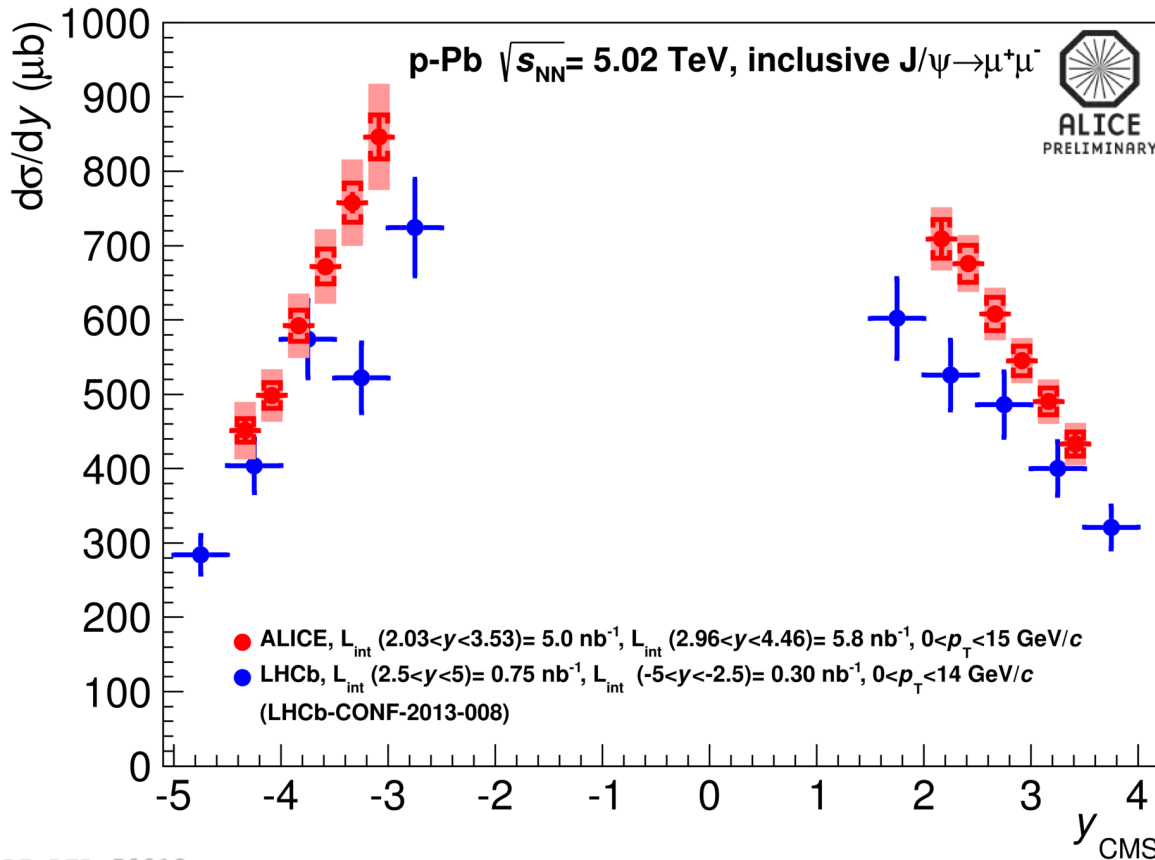
## NA60 function

$$f(x; N, \bar{x}, \sigma, x_1, x_2, p_1, \dots, p_6) = N \cdot \exp\left(-\frac{(x-\bar{x})^2}{2\sigma_{NA60}^2}\right)$$

$$\sigma_{NA60} = \begin{cases} \sigma \cdot \left(1 + p_1(x_1 - x)^{p_2 - p_3 \sqrt{x_1 - x}}\right), & \text{for } x < x_1 \\ \sigma, & \text{for } x_1 \leq x < x_2 \\ \sigma \cdot \left(1 + p_4(x - x_2)^{p_5 - p_6 \sqrt{x - x_2}}\right), & \text{for } x \geq x_2 \end{cases}$$

- ✓ Parameters  $N, \sigma, x$  are left free in the data
- ✓ All the other parameters are fixed on the tuned MC

# Comparison of ALICE results with LHCb - 1



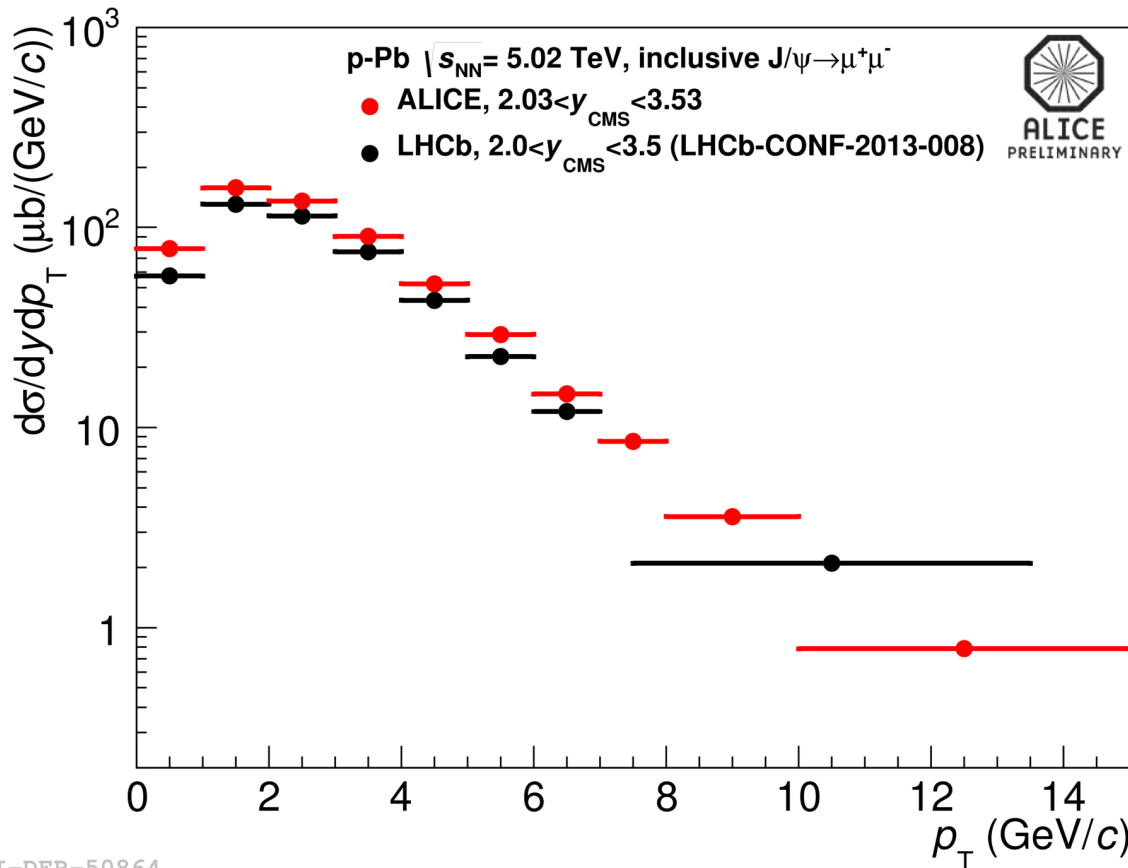
## ALICE uncertainties:

- ◇ **Statistical uncertainties** (line)
- ◇ **Systematic uncertainties:**
  - Corr. uncertainties** (brackets): luminosity, normalization factor, BR (Luminosity is correlated within p-Pb or Pb-p, but not within the two systems)
  - Uncorr. uncertainties** (filled boxes): matching, trigger, tracking, acc. inputs, signal extraction

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- Visible disagreement in results.
- Only half of statistics analyzed by LHCb.
- Difference in cross-sections seems to be the main source of the discrepancy.
- Work in progress on understanding the discrepancy between experiments.

# Comparison of ALICE results with LHCb - 2



## ALICE uncertainties:

- ◇ **Statistical uncertainties** (line)
- ◇ **Systematic uncertainties:**
  - Corr. uncertainties** (brackets): luminosity, normalization factor, BR (Luminosity is correlated within p-Pb or Pb-p, but not within the two systems)
  - Uncorr. uncertainties** (filled boxes): matching, trigger, tracking, acc. inputs, signal extraction

- Visible disagreement in results.
- Same conclusions as in the previous slide.
- Work in progress on understanding the discrepancy between experiments.