

# Multiplicity fluctuations of identified hadrons in p+p interactions at SPS energies

Based on Ph.D. Thesis of Maja Maćkowiak-Pawłowska

Andrzej Wilczek  
for the NA61 Collaboration

Institute of Physics, University of Silesia, Katowice

SQM 2013, Birmingham  
23 VII 2013

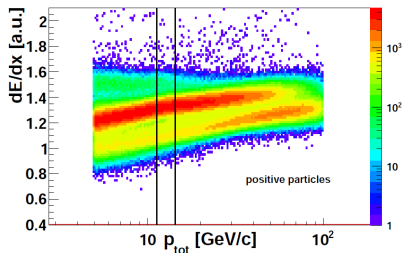
# Data sets

- The results are based on the NA61/SHINE data collected in 2009:
  - 1  $p(31 \text{ GeV}/c)+p$  ( $\sqrt{s} = 7.6 \text{ GeV}$ ):  $3.5 \times 10^6$  events
  - 2  $p(40 \text{ GeV}/c)+p$  ( $\sqrt{s} = 8.7 \text{ GeV}$ ):  $5.8 \times 10^6$  events
  - 3  $p(80 \text{ GeV}/c)+p$  ( $\sqrt{s} = 12.3 \text{ GeV}$ ):  $5.0 \times 10^6$  events
  - 4  $p(158 \text{ GeV}/c)+p$  ( $\sqrt{s} = 17.3 \text{ GeV}$ ):  $4.0 \times 10^6$  events
- event and track cuts were chosen to select only inelastic interactions with particles produced in strong and EM processes within the NA61/SHINE acceptance.
- analysis focuses on fluctuations of  $\pi = \pi^+ + \pi^-$ ,  $K = K^+ + K^-$  and  $p + \bar{p}$  as well as positively charged hadrons ( $p, K^+, \pi^+$ ) by getting first and second (pure and mixed) moments of identified particle multiplicity distributions.
- second moments of identified particle multiplicity distributions are corrected for the misidentification effect using the identity method.<sup>1</sup>
- Presented results of NA61/SHINE include statistical errors and first estimates of systematic uncertainties (work to finalize systematic uncertainties is in progress e.g. feed down and detector effects).

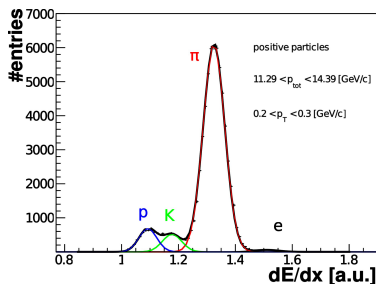
<sup>1</sup>PRC83:054907, PRC84:024902, PRC86:044906

# Particle identification

Particle identification for chemical fluctuation analysis is based on  $dE/dx$  measurements in relativistic rise region.



Inclusive  $dE/dx$  spectra  
are sliced in  $p_{tot}$ ,  $p_T$



sum of Gaussian functions is fitted  
in each phase-space bin.

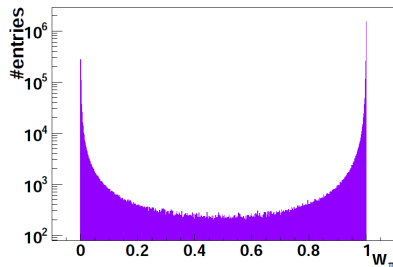
# Identity method

The identity method allows to obtain second and third moments (pure and mixed) of identified particle multiplicity distributions corrected for misidentification effect.

The particle identity is calculated as:

$$w_i = \frac{\rho_i(dE/dx)}{\rho(dE/dx)},$$

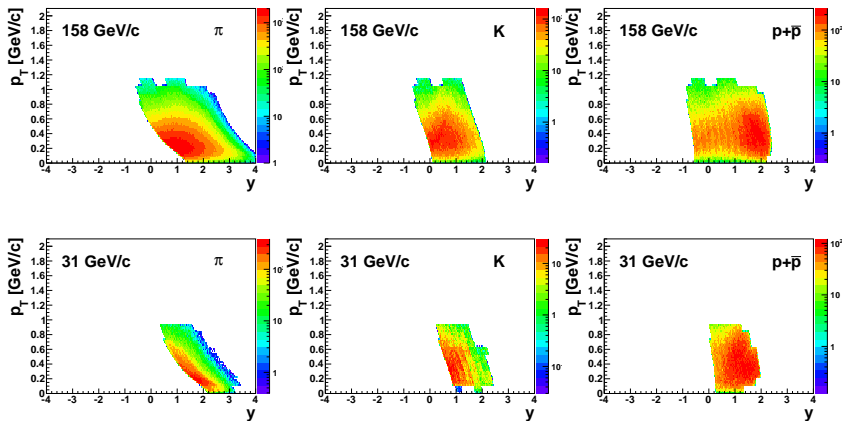
where  $\rho_i$  - function fitted to the  $i^{\text{th}}$  particle type and  $\rho$  - function fitted to the total  $dE/dx$  distribution in a given phase-space bin.



example of  $w_\pi$  distribution for p+p at 12.3 GeV

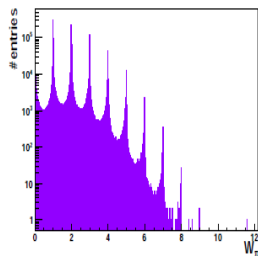
# Identified hadron acceptance

Single particle identity allows to obtain identified hadron acceptance by taking each particle with its corresponding  $w$  as weight.

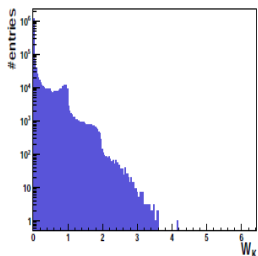


# Identity method - event identity measure

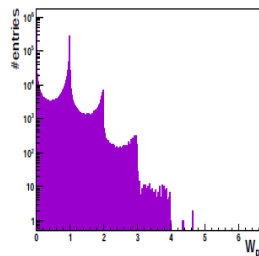
Event quantity  $W_i$  defined as:  $W_i = \sum w_j$  where summation runs over all particles in an event



example of event variable  $W_\pi$ ,



$W_K$ ,



$W_{p+\bar{p}}$  for p+p at 12.3 GeV

Once, detector response ( $\rho_i$ ) and  $W$  distributions are known the identity method is used to obtain moments of identified particle multiplicity distributions.

# Fluctuation quantities

In the Wounded Nucleon Model or the GCE one may define two families of quantities:

## Intensive quantities

A ratio of two extensive quantities ( $N_W$ ) is an intensive measure e.g.:

$$\omega_i = \frac{\langle N_i^2 \rangle - \langle N_i \rangle^2}{\langle N_i \rangle}$$

- independent of  $N_W$
- depends on fluctuations of  $N_W$  - not good for comparison of p+p with wide A+A centrality bins
- $\omega = 1$  for Poisson distribution

In WNM  $\omega_i = \frac{\text{Var}(n_i)}{\langle n_i \rangle} + \langle n_i \rangle \frac{\text{Var}(N_W)}{\langle N_W \rangle}$ , where  $n_i$  - particles produced from single wounded nucleon

## Strongly intensive quantities

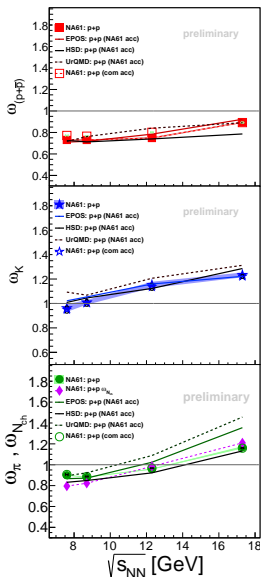
Special combinations of extensive quantities can be strongly intensive measures e.g.:

$$\Phi_{ij} = \frac{\sqrt{\langle N_i \rangle \langle N_j \rangle}}{\langle N_i + N_j \rangle} \times \left( \sqrt{\Sigma^{ij}} - 1 \right),$$

- independent of  $N_W$
- independent of fluctuations of  $N_W$
- $\Phi_{ij} = 0$  for independent particle emission

where  $\Sigma^{ij} = [\langle N_i \rangle \omega_j + \langle N_j \rangle \omega_i - 2(\langle N_i N_j \rangle - \langle N_i \rangle \langle N_j \rangle)] / \langle N_i + N_j \rangle$

# Multiplicity fluctuations of identified hadrons

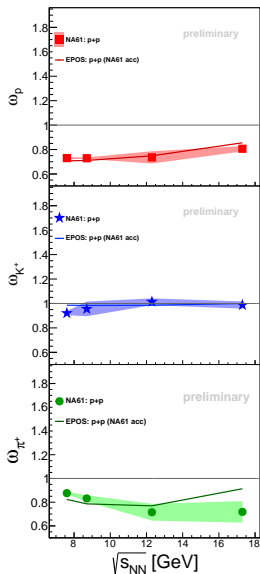


- Results are presented in NA61 acceptance<sup>a</sup> and in common acceptance of NA49 and NA61
- $\omega_{p+\bar{p}}$  is below 1 for all SPS energies possibly due to baryon number conservation ( $B = 2$ )
- $\omega_K$  is above 1 for all SPS energies possibly due to strangeness conservation ( $S = 0$ )
- $\omega_\pi$  increases with increasing energy
- HSD, EPOS, and UrQMD reproduce measured scaled variances

<sup>a</sup><https://edms.cern.ch/document/1237791/1>

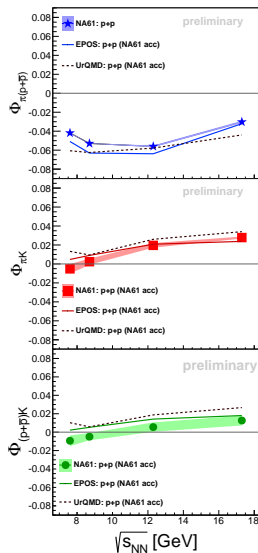


# Multiplicity fluctuations of positively charged hadrons



- $\omega_p$  similar to  $\omega_{p+\bar{p}}$ . Small fraction of antiprotons
- $\omega_{K^+} \approx 1$   
 $\omega_{K^+} < \omega_K$  which suggests that strangeness conservation contributes to  $\omega_K$
- suppressed fluctuations of  $\pi^+$ .  $\omega_{\pi^+} < \omega_{\pi^-}$  possibly due to charge conservation
- EPOS reproduces scaled variances of positively charged hadrons

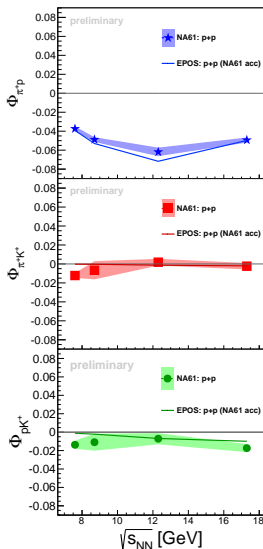
# Chemical fluctuations of charged hadrons



- fluctuations of  $\pi(p + \bar{p})$  are affected by conservation laws and resonance decays<sup>a</sup> -  $\Phi_{\pi(p+\bar{p})} < 0$
- small increase of  $\pi K$  fluctuations with increasing energy
- $\Phi_{(p+\bar{p})K} \approx 0$  indicates weak if any correlations between  $p + \bar{p}$  and  $K$  production
- UrQMD and EPOS closely reproduces the data

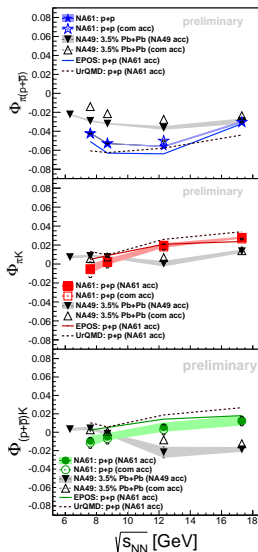
<sup>a</sup>PRC70:064903

# Chemical fluctuations of positively charged hadrons



- fluctuations of  $\pi^+ p$  agree with both charge fluctuations
- no energy dependence of  $\pi^+ K^+$  and  $p K^+$  fluctuations
- EPOS closely reproduces data

# Chemical fluctuations of charged hadrons in p+p (NA61) and Pb+Pb (NA49)



- limiting the analysis to the region of common NA49 and NA61 acceptance does not impact the results in any significant way
- qualitatively similar energy dependence of  $\pi(p + \bar{p})$  and  $\pi K$  fluctuations for p+p and central Pb+Pb collisions
- $(p + \bar{p})K$  fluctuations show different tendency in p+p (increase with energy) than in Pb+Pb (decrease)
- note that the results are still preliminary

# Summary

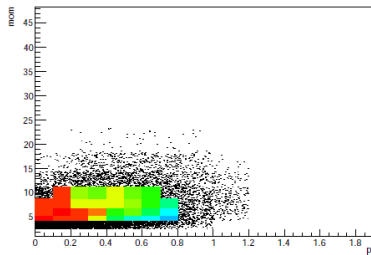
- multiplicity fluctuations of  $\pi$ ,  $K$ ,  $p$  and chemical fluctuations of  $\pi p$ ,  $\pi K$ ,  $pK$  as well as fluctuations of positively charged hadrons were measured by NA61/SHINE at  $\sqrt{s_{NN}} = 7.6 - 17.3$  GeV.
- models (EPOS, HSD, UrQMD) describe fluctuations in p+p interactions
- qualitatively similar energy dependence of  $\pi(p + \bar{p})$  and  $\pi K$  fluctuations for p+p and central Pb+Pb collisions
- when strongly intensive measure  $\Phi$  is used,  $(p + \bar{p})K$  fluctuations show different tendency in p+p (increase with energy) than in Pb+Pb (decrease)

BACKUP

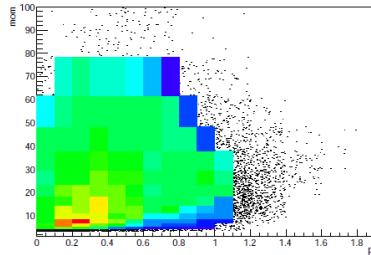
# Common acceptance of NA61 and NA49

In order to compare p+p (NA61) and Pb+Pb (NA49) results the common phase space region for the chemical fluctuation analysis was defined.

Due to significantly lower particle multiplicity in p+p interactions, the region where the track density is sufficient for energy loss analysis is limited in the case of NA61.



$$\sqrt{s_{NN}} = 7.6 \text{ GeV}$$



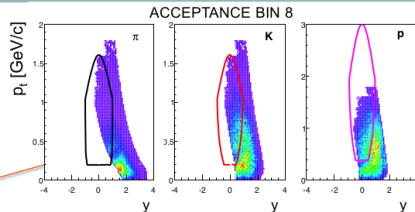
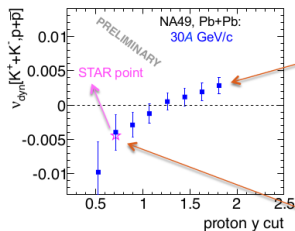
$$\sqrt{s_{NN}} = 17.3 \text{ GeV}$$

Coloured region marks common phase-space region used for comparison of p+p and Pb+Pb results (scattered points indicate region used for Pb+Pb analysis only). For details see <https://edms.cern.ch/document/1237791/1>

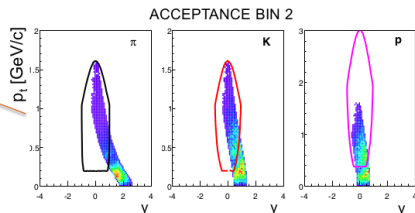
## Acceptance issues (STAR and NA61)

## Dependence on acceptance

NA49 central Pb+Pb data at 30A GeV/c



Lines indicate the corresponding STAR acceptance





## Other measures of fluctuations

$$\begin{array}{ll}
 \text{NA49:} & \sigma_{dyn} = \text{sgn}(\sigma_{data}^2 - \sigma_{mixed}^2) \sqrt{|\sigma_{data}^2 - \sigma_{mixed}^2|} & \sigma = \frac{\sqrt{\text{Var}(A/B)}}{\langle A/B \rangle} & \frac{A}{B} = \frac{K}{\pi}, \frac{p}{\pi}, \frac{K}{p} \\
 \text{STAR:} & \nu_{dyn} = \frac{\langle N_1^2 \rangle}{\langle N_1 \rangle^2} + \frac{\langle N_2^2 \rangle}{\langle N_2 \rangle^2} - 2 \frac{\langle N_1 N_2 \rangle}{\langle N_1 \rangle \langle N_2 \rangle} - \left( \frac{1}{\langle N_1 \rangle} + \frac{1}{\langle N_2 \rangle} \right) & \nu_{dyn} \approx \text{sgn}(\sigma_{dyn}) \sigma_{dyn}^2
 \end{array}$$