

Strange particle production in Monte Carlo generators

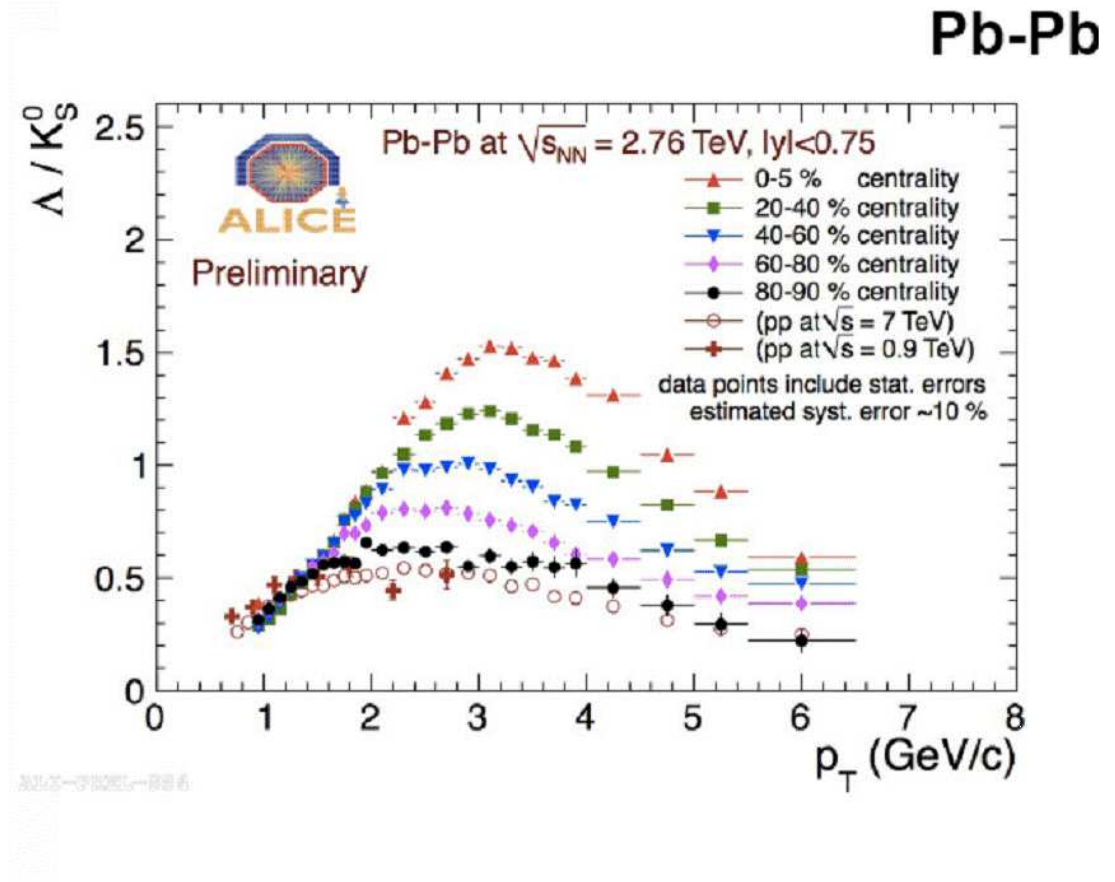
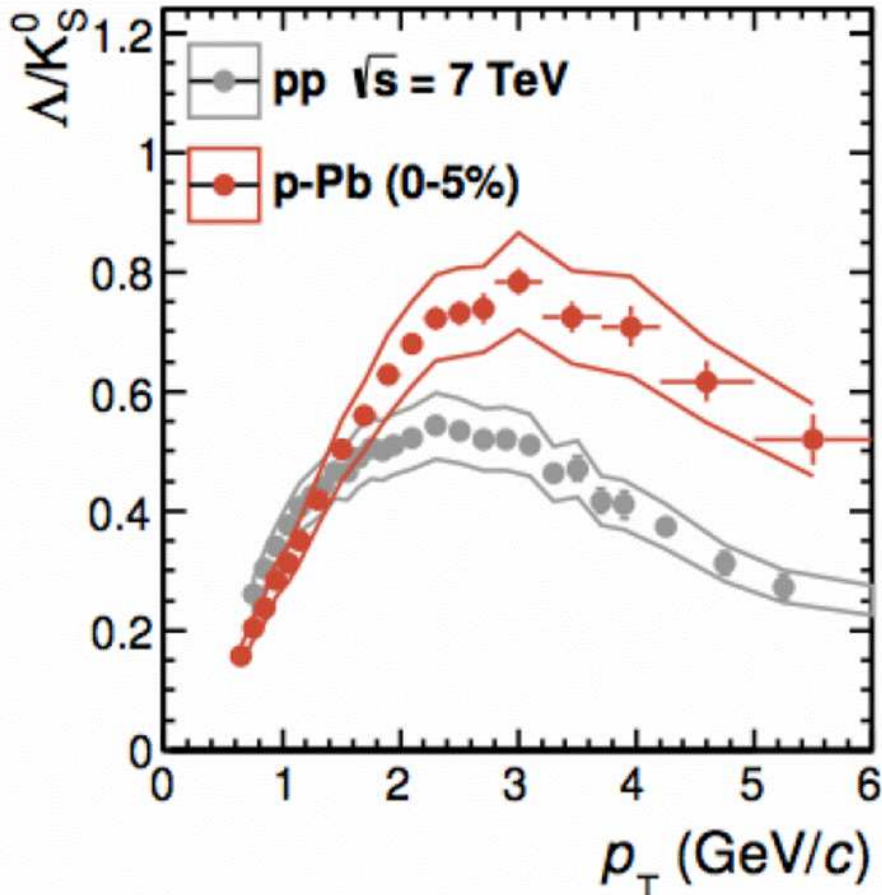
in pp and pPb collisions at the LHC

**pp: many simulation results
from various generators**

**pPb: much less simulated so far,
amazing experimental results,
only very recently also
on identified particles**

pPb at 5TeV

From R. Preghenella, ALICE, talk Trento workshop 2013

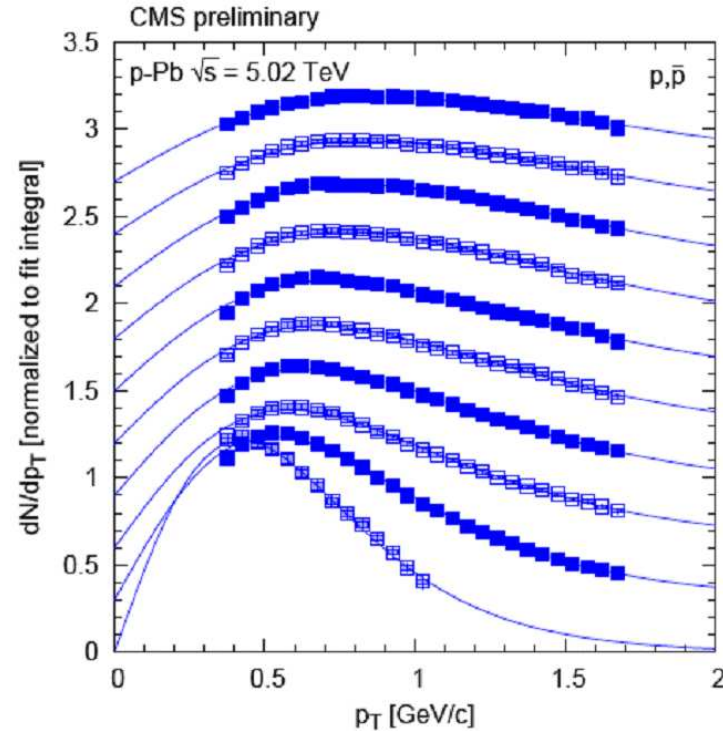
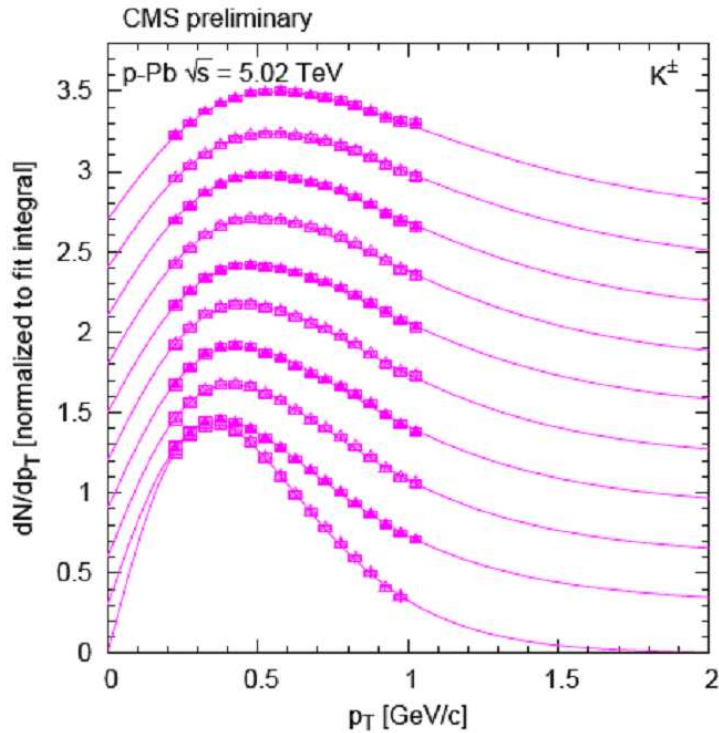


Lambda/Kshort: similar behavior as in PbPb

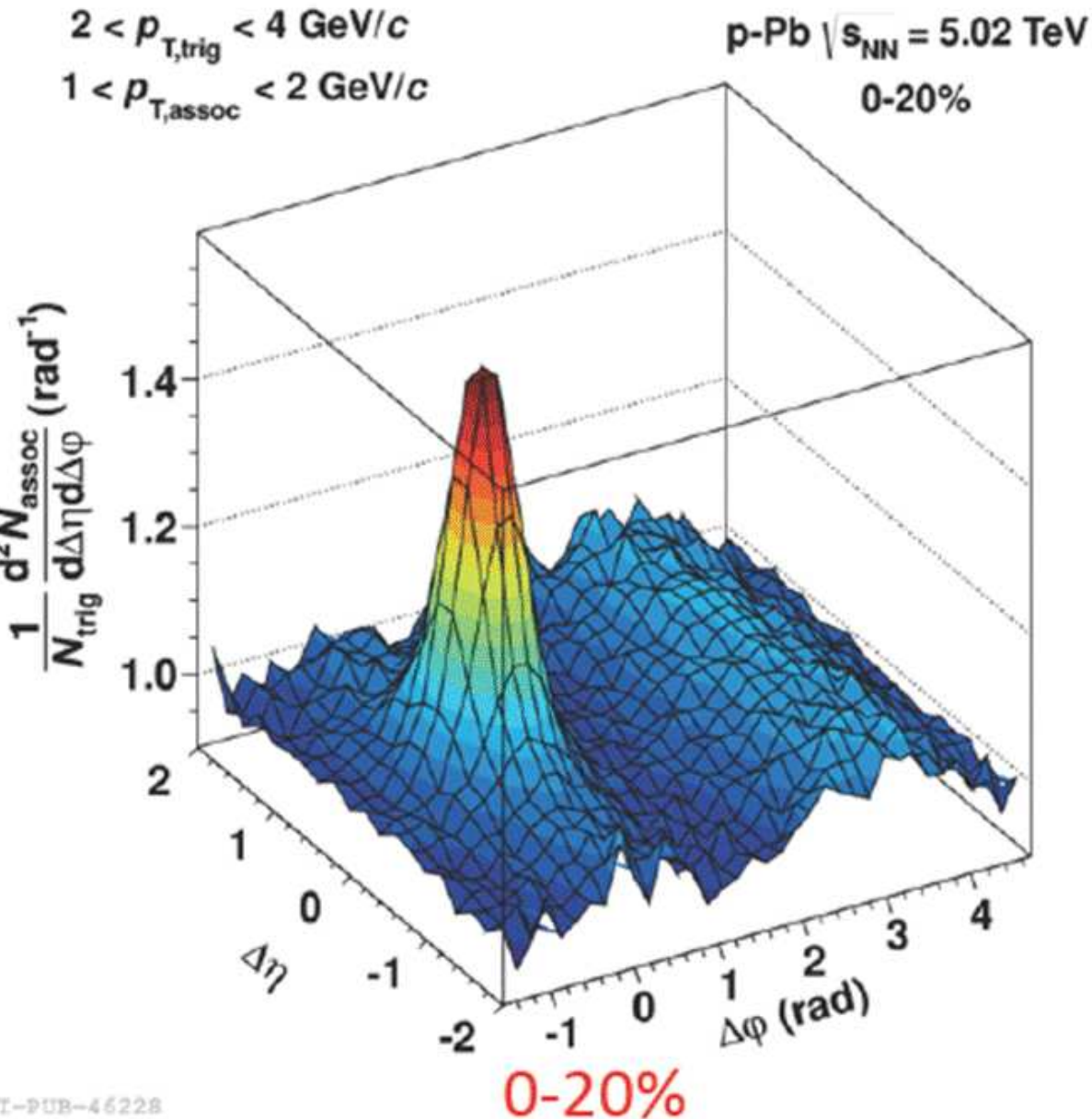
=> flow?

pPb at 5TeV

CMS, arXiv:1307.3442



**Strong variation of shape with multiplicity
for kaon and even more for proton p_T spectra
(flow like)**



pPb at 5TeV

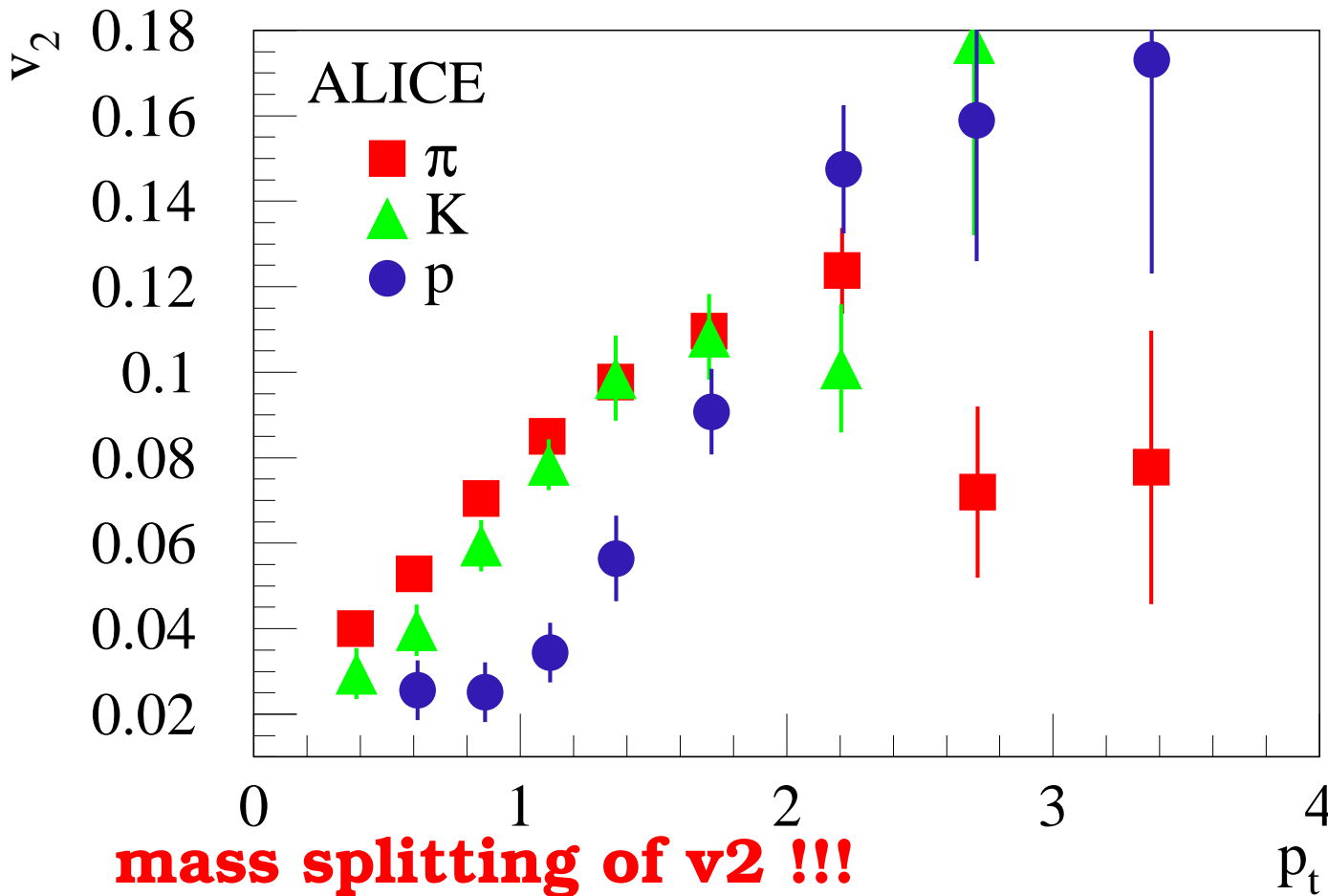
ALICE, Phys.Lett.
B719 (2013) 29,
arXiv:1212.2001

**ridge
& significant v_2**

as in in PbPb ...

pPb at 5TeV

ALICE, arXiv:1307.3237



one of the highlights in heavy ion collisions ...

pPb TeV looks very interesting!

**But even in pp some of these
“flow features” have been observed...**

=> closer look at pp first, then pPb.

**What do the MC's predict, concerning
identified (strange) particle production?**

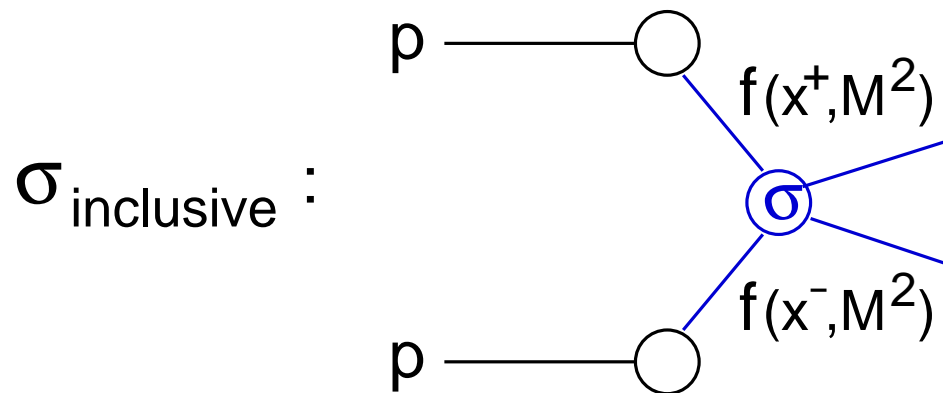
Models based on factorization

Pythia6
Pythia8
Herwig++
Sherpa

Phys.Rept. 504
(2011) 145-233

see also
mcplots.cern.ch/

Starting point:
Factorization formula
+ ISR + FSR



Multiple scattering
in a second step

Models based on Gribov Regge Theory

Epos3, Epos LHC
Qgsjet, Sybill, Phojet

EPOS3: K. Werner et al.,
arXiv:1307.4379

EPOS LHC: T. Pierog et al,
arXiv:1306.5413

QGSJET: S. Ostapchenko, Phys. Rev.
D74 (2006) 014026

SYBILL: E.-J. Ahn et al, Phys. Rev. D 80
(2009) 094003

PHOJET: R. Engel and J. Ranft, Phys.
Rev. D54 (1996) 4244

[Simulations of SYBILL, QGSJET, EPOS
LHC by T. Pierog](#)

Starting point:
Gribov Regge
multiple scattering

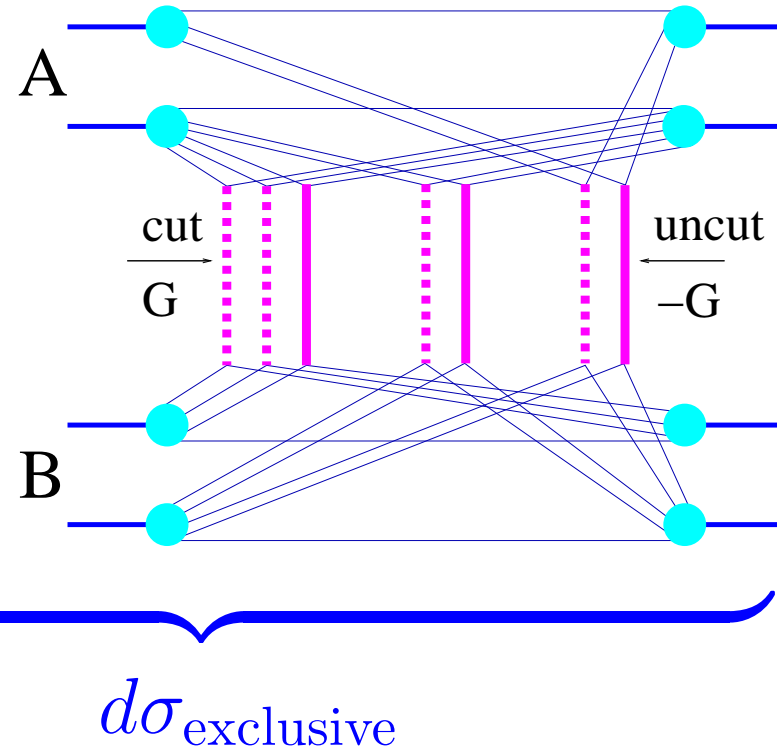
pQCD added

Treats exclusive
and
inclusive quantities

EPOS: Marriage pQCD+GRT+energy sharing

(Drescher, Hladik, Ostapchenko, Pierog, and Werner, Phys. Rept. 350, 2001)

$$\sigma^{\text{tot}} = \sum_{\text{cut P}} \int \sum_{\text{uncut P}} \int$$



$$\text{cut Pom} : G = \frac{1}{2\hat{s}} 2\text{Im} \{ \mathcal{FT} \{ T \} \} (\hat{s}, b), \quad T = i\hat{s} \sigma_{\text{hard}}(\hat{s}) \exp(R_{\text{hard}}^2 t)$$

Nonlinear effects considered via saturation scale $Q_s \propto N_{\text{part}} \hat{s}^\lambda$

$$\begin{aligned}
 \sigma^{\text{tot}} = & \int d^2b \int \prod_{i=1}^A d^2b_i^A dz_i^A \rho_A(\sqrt{(b_i^A)^2 + (z_i^A)^2}) \\
 & \prod_{j=1}^B d^2b_j^B dz_j^B \rho_B(\sqrt{(b_j^B)^2 + (z_j^B)^2}) \\
 & \sum_{m_1 l_1} \dots \sum_{m_{AB} l_{AB}} (1 - \delta_{0 \Sigma m_k}) \int \prod_{k=1}^{AB} \left(\prod_{\mu=1}^{m_k} dx_{k,\mu}^+ dx_{k,\mu}^- \prod_{\lambda=1}^{l_k} d\tilde{x}_{k,\lambda}^+ d\tilde{x}_{k,\lambda}^- \right) \left\{ \right. \\
 & \prod_{k=1}^{AB} \left(\frac{1}{m_k!} \frac{1}{l_k!} \prod_{\mu=1}^{m_k} G(x_{k,\mu}^+, x_{k,\mu}^-, s, |\vec{b} + \vec{b}_{\pi(k)}^A - \vec{b}_{\tau(k)}^B|) \right. \\
 & \left. \left. \prod_{\lambda=1}^{l_k} -G(\tilde{x}_{k,\lambda}^+, \tilde{x}_{k,\lambda}^-, s, |\vec{b} + \vec{b}_{\pi(k)}^A - \vec{b}_{\tau(k)}^B|) \right) \right\} \\
 & \prod_{i=1}^A \left(1 - \sum_{\pi(k)=i} x_{k,\mu}^+ - \sum_{\pi(k)=i} \tilde{x}_{k,\lambda}^+ \right)^\alpha \prod_{j=1}^B \left(1 - \sum_{\tau(k)=j} x_{k,\mu}^- - \sum_{\tau(k)=j} \tilde{x}_{k,\lambda}^- \right)^\alpha \left. \right\}
 \end{aligned}$$

Other models

AMPT

Z.-W. Lin, C. M. Ko, B.-A. Li, B. Zhang and S. Pal, Phys. Rev. C 72, 064901 (2005).

J. L. Albacete, N. Armesto, R. Baier, et al., Int. J. Mod. Phys. E 22, 1330007 (2013).

Simulations provided by [Ziwei Lin](#)

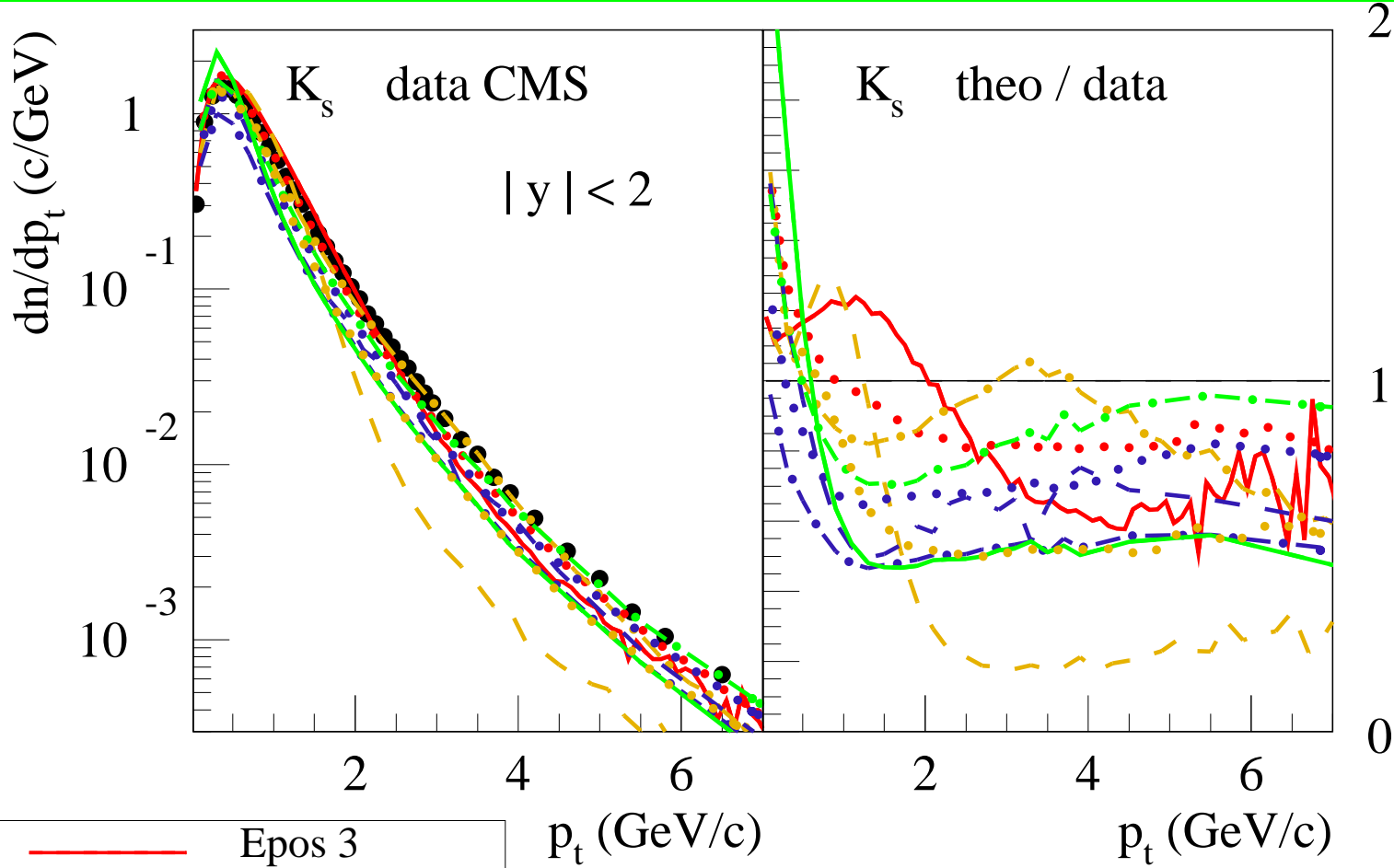
Model not optimized for LHC

Initial condition:
minijets from HIJING 1.0

Minijet partons rescatter in
the partonic stage

Hadronization via Lund
string fragmentation

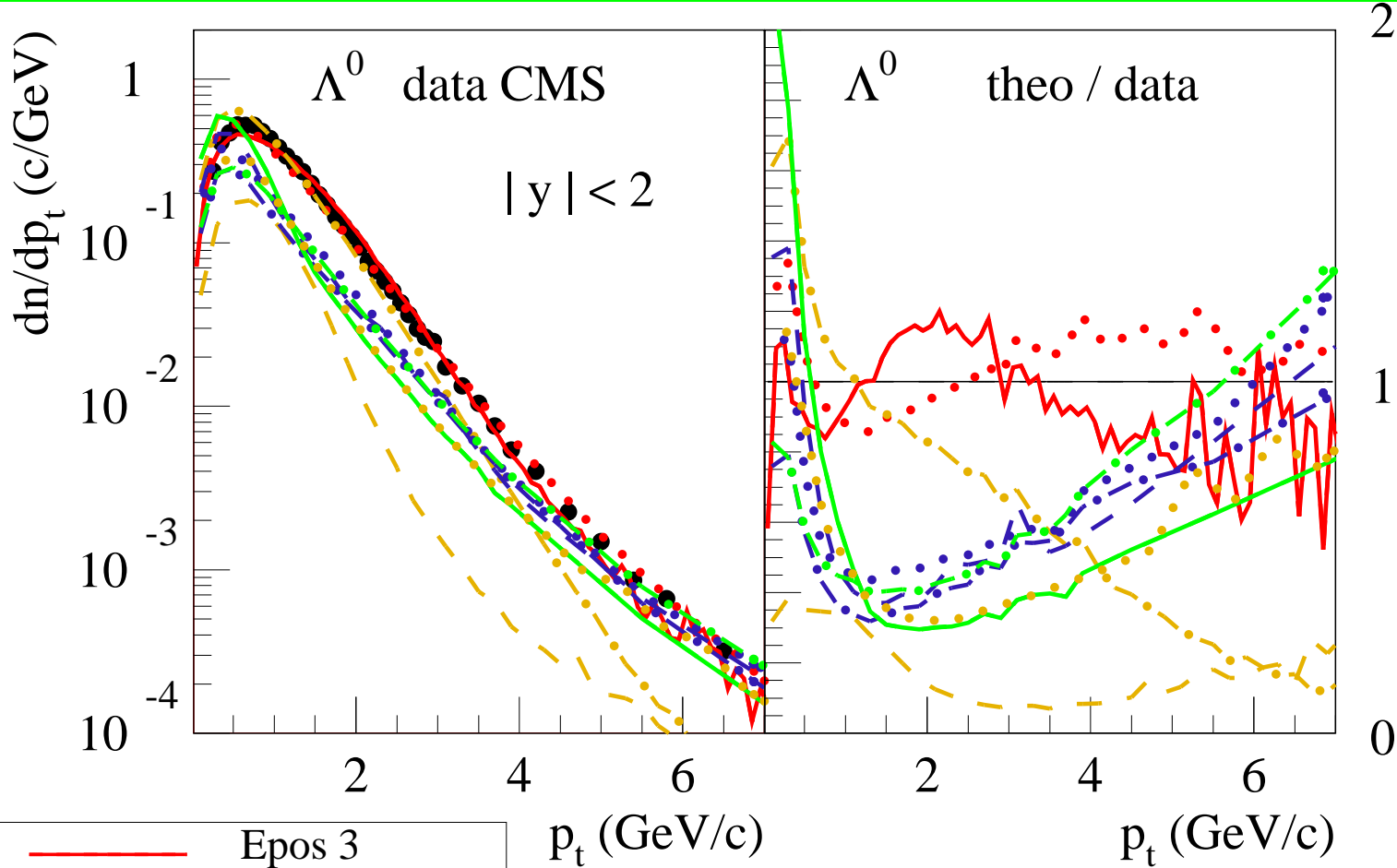
Final state hadronic cas-
cade



- Epos 3
- Epos LHC
- .-.- Qgsjet
- - - Sybill
- Ampt
- - - Phojet
- .-.- Pythia6
- Pythia8
- Herwig++
- .-.- Sherpa

Kaons

All models underpredict the data, 20-50%



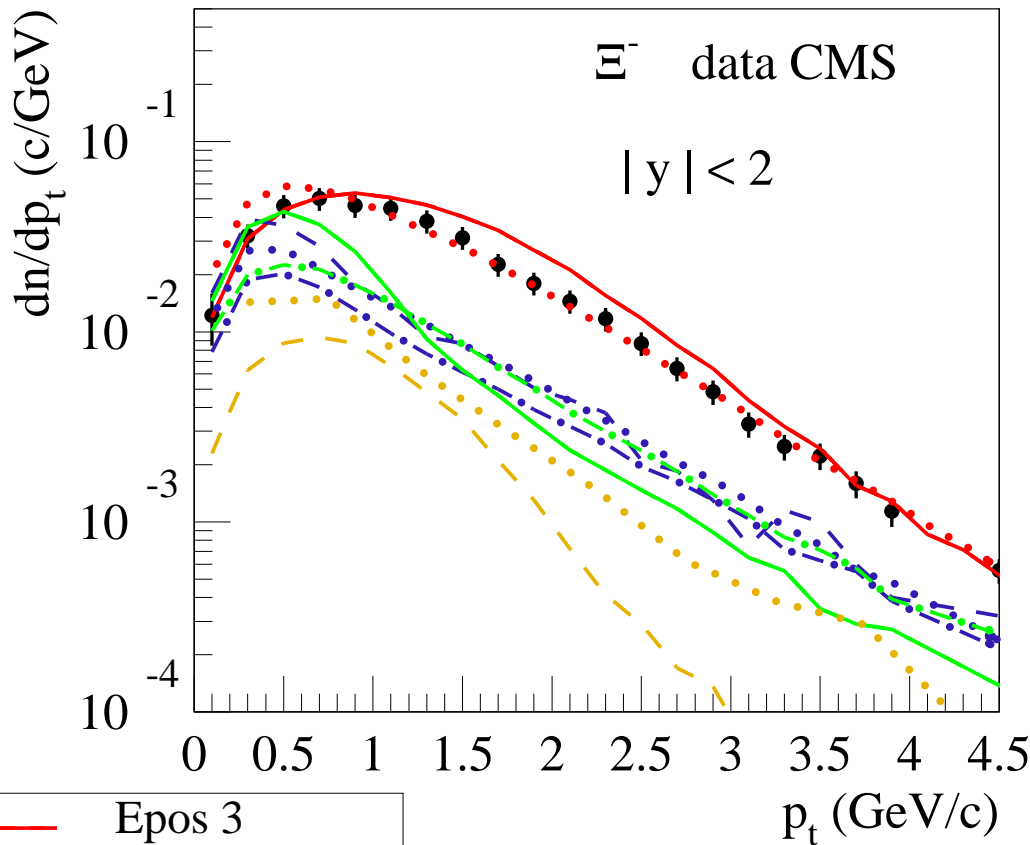
- Epos 3
- Epos LHC
- .-.-.- Qgsjet
- Sybill
- Ampt
- .-.-.- Phojet
- Pythia6
- Pythia8
- Herwig++
- .-.-.- Sherpa

Lambdas

Sybill, Qgsjet drop out

Factorization models: profound "hole"

EPOS's ~20%



- Epos 3
- Epos LHC
- Qgsjet
- - - Sybill
- Ampt
- - - Phojet
- Pythia6
- Pythia8
- Herwig++
- Sherpa

Xis

Sybill, Qgsjet (no Xis) drop out

**Factorization models:
considerably below the data**

EPOS's relatively close

EPOS3 and EPOS LHC:

collective flow, statistical hadronization

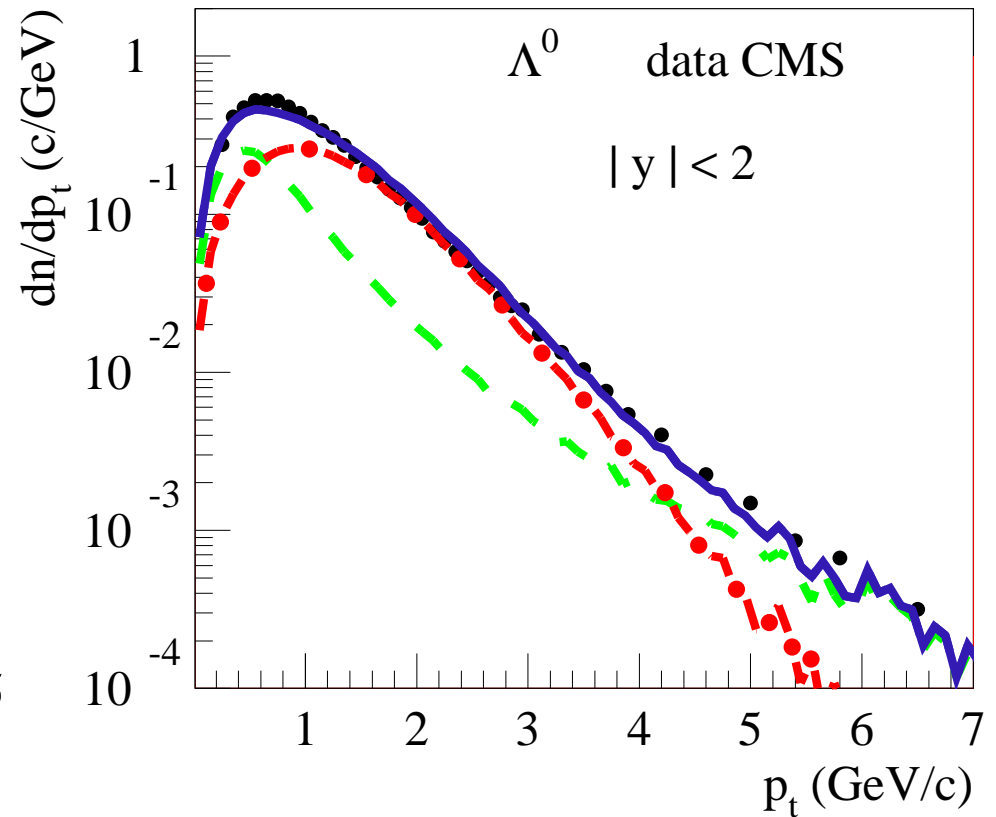
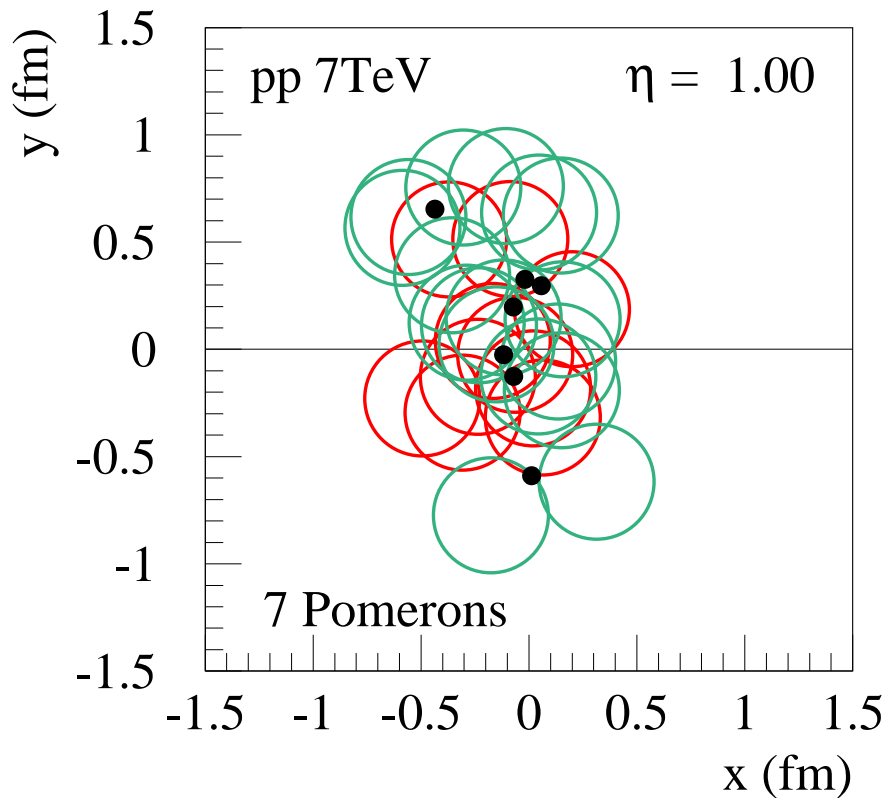
- **EPOS3: Gribov Regge (GT) approach provides the initial condition for subsequent hydro evolution & hadronization in pp, pA, AA**

EbE viscous 3D+1 dim hydro => Y. Karpenko

- **EPOS LHC (= LHC tune of EPOS1.99, T. Pierog): same GR, but uses parameterized flow**

EPOS3: Pomeron => parton ladder => flux tubes (kinky strings)

String segments with high p_t escape (corona),
the others form the “core” = initial condition for hydro



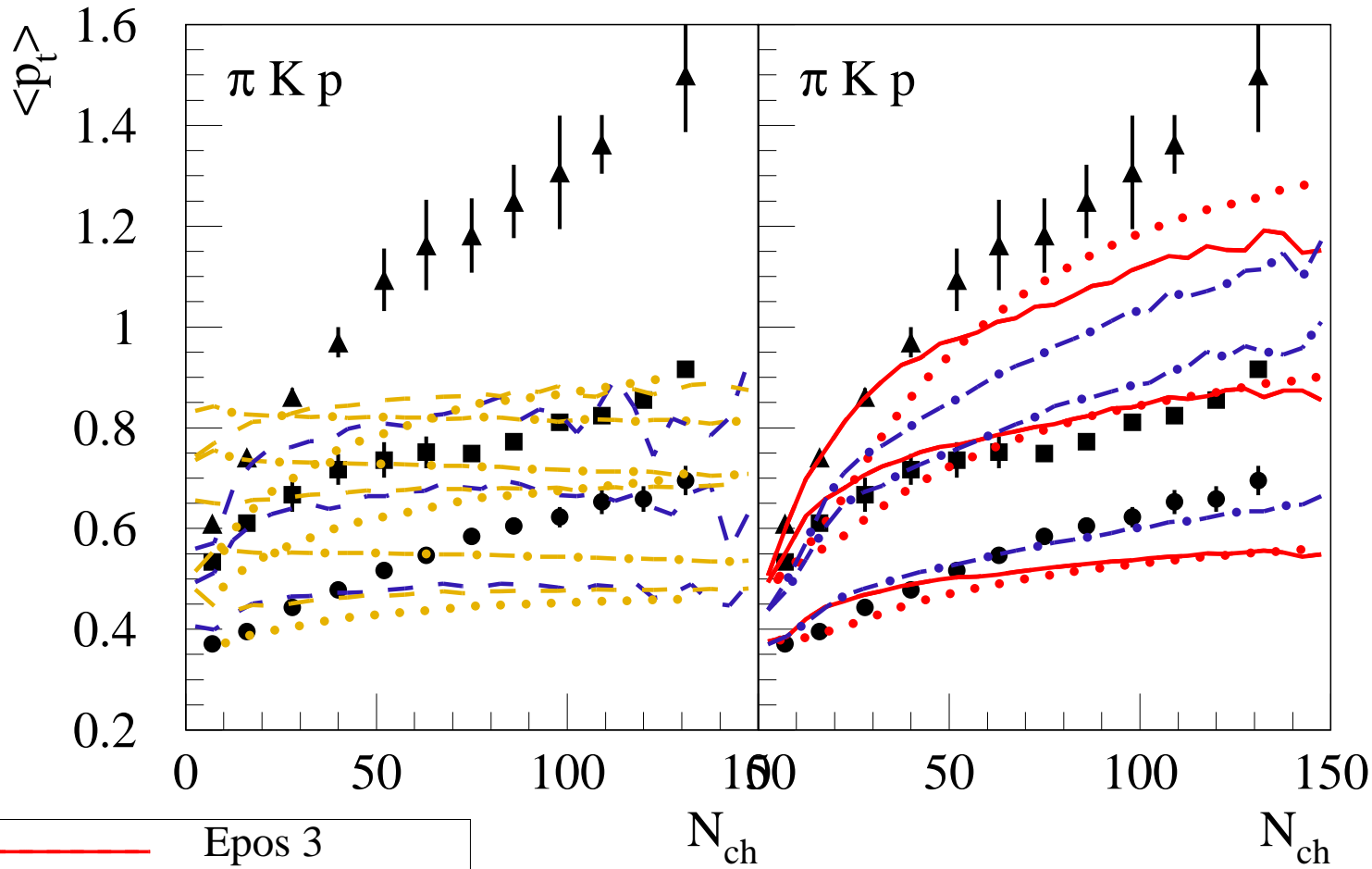
The hydrodynamic equations (Israel-Stewart formulation) in arbitrary coordinate system (implemented/solved by Yuri Karpenko), always $\eta/S = 0.08$, $\zeta/S = 0$

$$\partial_{;\nu} T^{\mu\nu} = \partial_{\nu} T^{\mu\nu} + \Gamma_{\nu\lambda}^{\mu} T^{\nu\lambda} + \Gamma_{\nu\lambda}^{\nu} T^{\mu\lambda} = 0$$

$$\gamma (\partial_t + v_i \partial_i) \pi^{\mu\nu} = -\frac{\pi^{\mu\nu} - \pi_{\text{NS}}^{\mu\nu}}{\tau_{\pi}} + I_{\pi}^{\mu\nu}$$

$$\gamma (\partial_t + v_i \partial_i) \Pi = -\frac{\Pi - \Pi_{\text{NS}}}{\tau_{\Pi}} + I_{\Pi}$$

- | | |
|--|---|
| <input type="checkbox"/> $T^{\mu\nu} = \epsilon u^{\mu} u^{\nu} - (p + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$, | <input type="checkbox"/> $\pi_{\text{NS}}^{\mu\nu} = \eta (\Delta^{\mu\lambda} \partial_{;\lambda} u^{\nu} + \Delta^{\nu\lambda} \partial_{;\lambda} u^{\mu}) - \frac{2}{3} \eta \Delta^{\mu\nu} \partial_{;\lambda} u^{\lambda}$ |
| <input type="checkbox"/> $\partial_{;\nu}$ denotes a covariant derivative, | <input type="checkbox"/> $\Pi_{\text{NS}} = -\zeta \partial_{;\lambda} u^{\lambda}$ |
| <input type="checkbox"/> $\Delta^{\mu\nu} = g^{\mu\nu} - u^{\mu} u^{\nu}$ is the projector orthogonal to u^{μ} , | <input type="checkbox"/> $I_{\pi}^{\mu\nu} = -\frac{4}{3} \pi^{\mu\nu} \partial_{;\gamma} u^{\gamma} - [u^{\nu} \pi^{\mu\beta} + u^{\mu} \pi^{\nu\beta}] u^{\lambda} \partial_{;\lambda} u_{\beta}$ |
| <input type="checkbox"/> $\pi^{\mu\nu}$ and Π are the shear stress tensor and bulk pressure, respectively. | <input type="checkbox"/> $I_{\Pi} = -\frac{4}{3} \Pi \partial_{;\gamma} u^{\gamma}$ |



Multipl. dependence of means p_t for π , K , p :

none for QGSJET, little for Sybill, Phojet

increase with N_{ch} in Pythia and EPOS has completely different origin

Models for pPb

Epos3
Epos LHC
Qgsjet
AMPT

“QCD generators”
Pythia, Herwig,...
do only pp

Piotr Bozek et al
(arXiv:1304.3044):
hydrodynamical evolution
based on Glauber initial con-
ditions
Results shown later at
this conference

CGC / satura-
tion / glasma
No particular
flavor depen-
dencies for
u,d,s

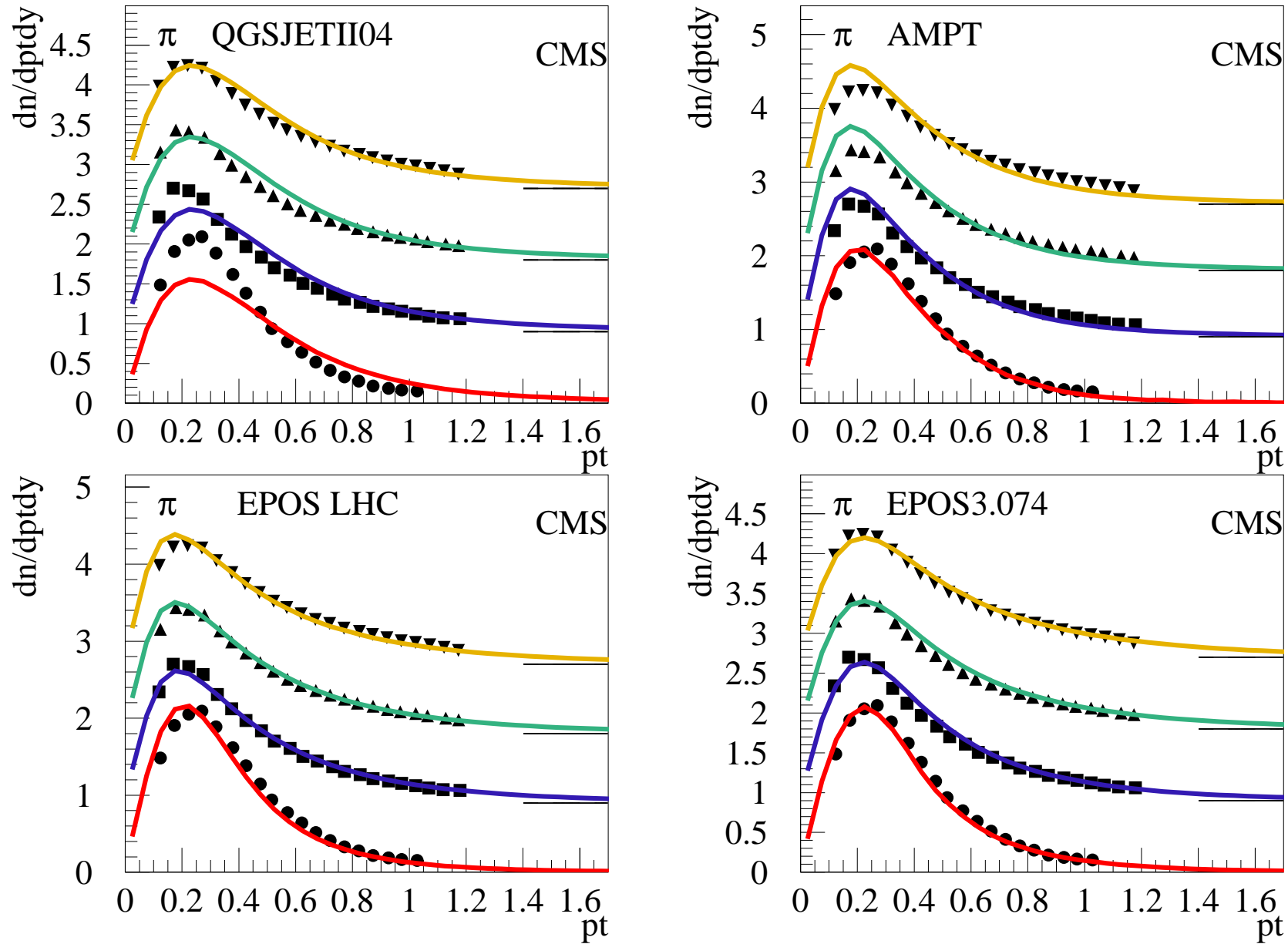
CMS: Multiplicity dependence of pion, kaon, proton pt spectra

CMS, arXiv:1307.3442

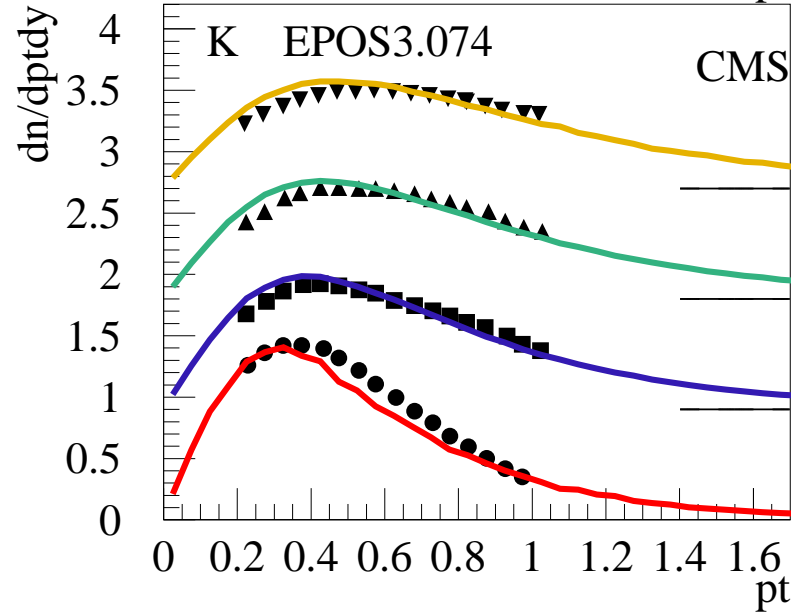
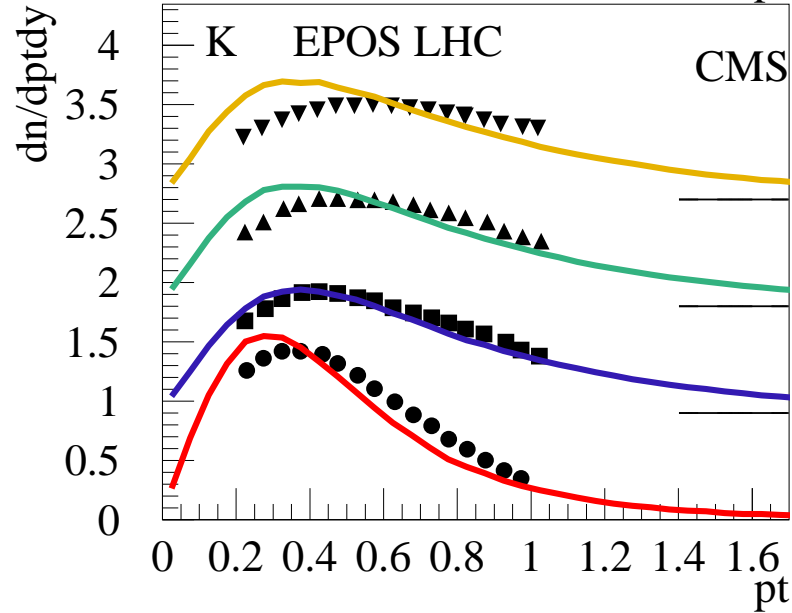
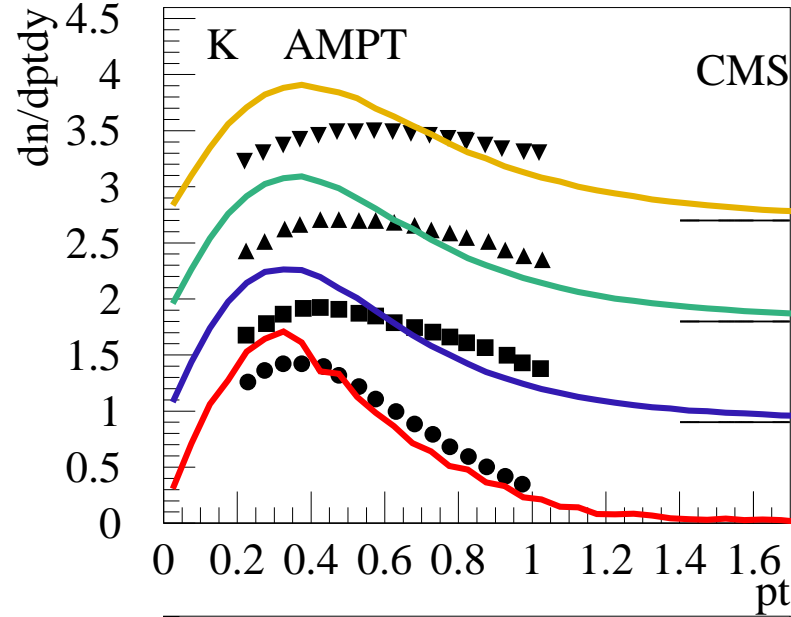
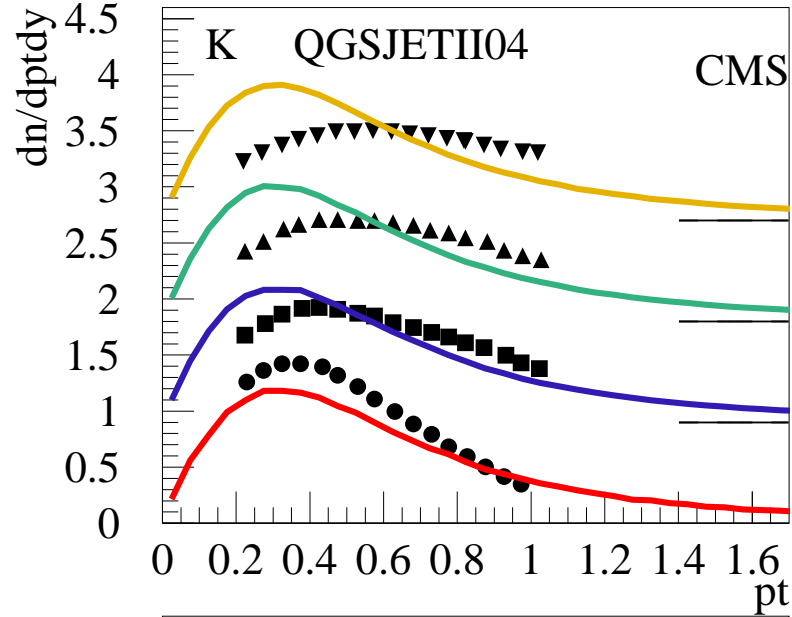
We plot 4 centrality classes:

$\langle N_{\text{tracks}} \rangle = 8, 84, 160, 235$ (in $|\eta| < 2.4$)

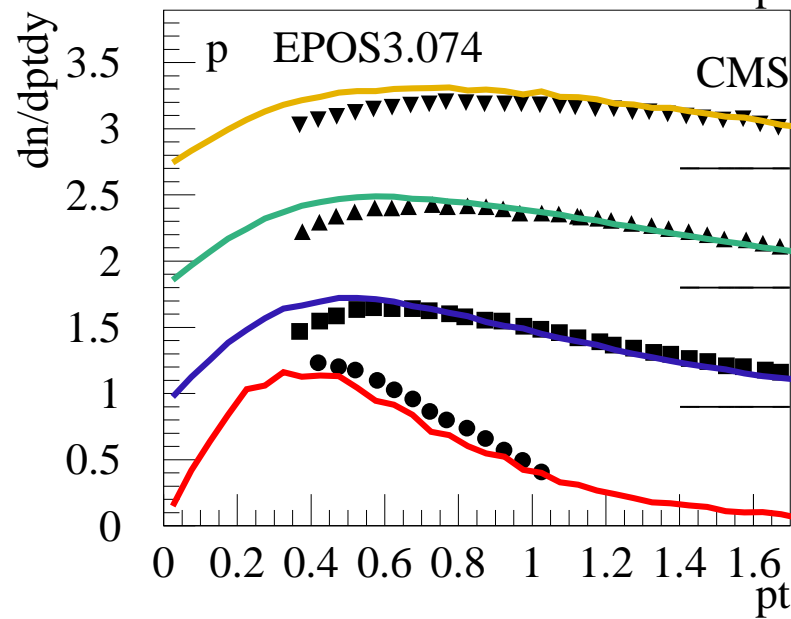
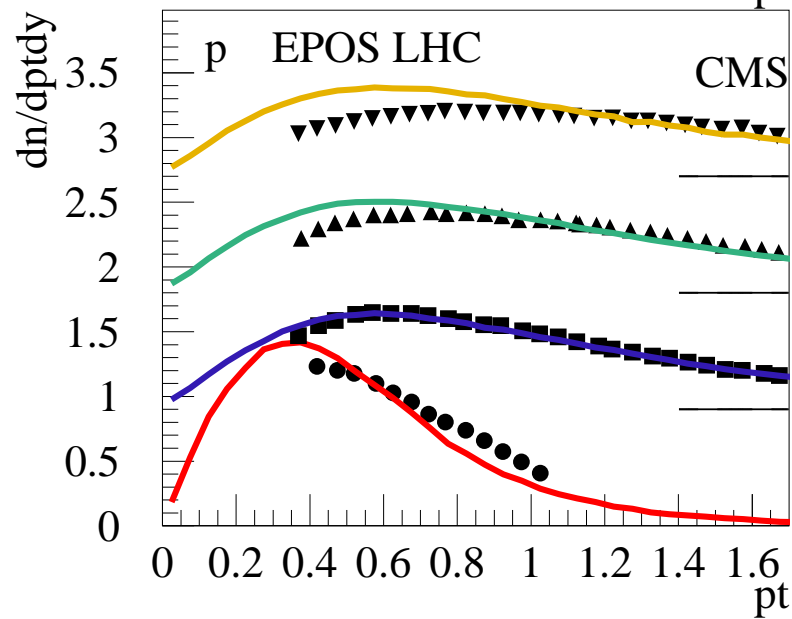
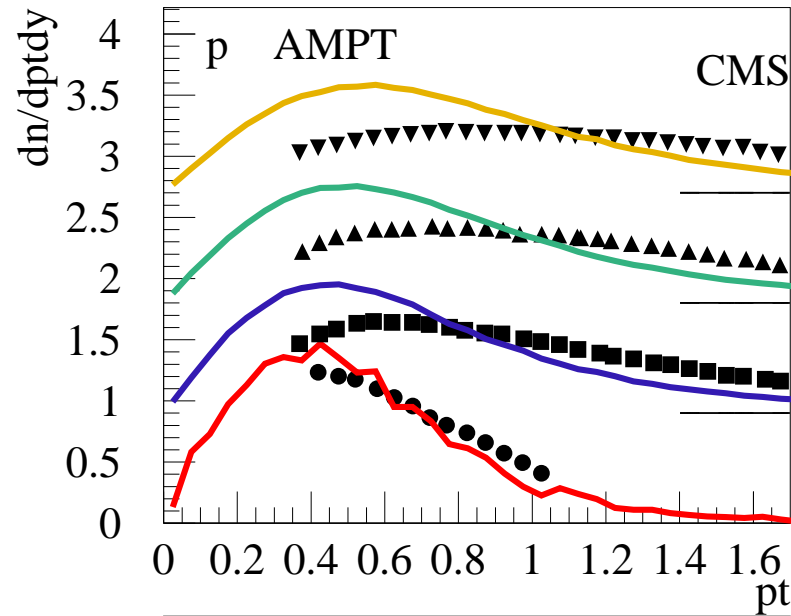
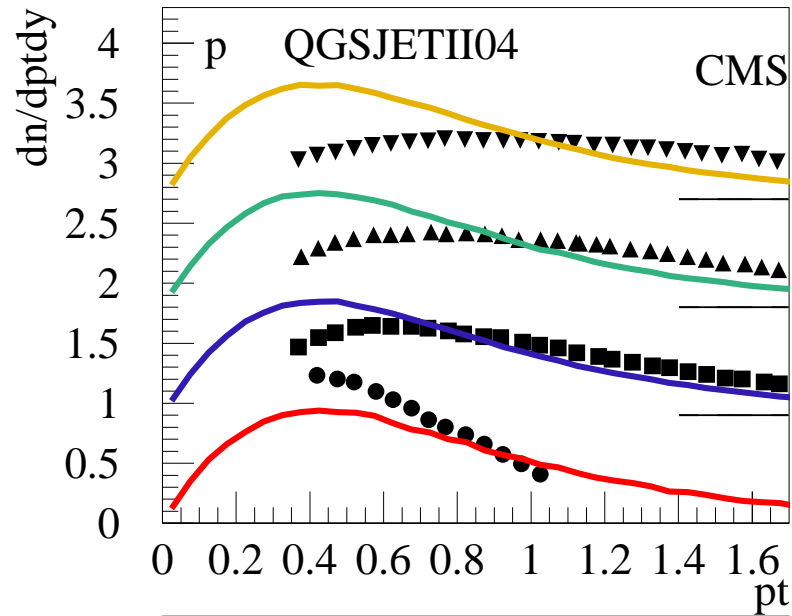
Multiplicity = centrality measure



Little change with $\langle N_{\text{tracks}} \rangle$ for pions



Kaon spectra change with $\langle N_{tracks} \rangle$

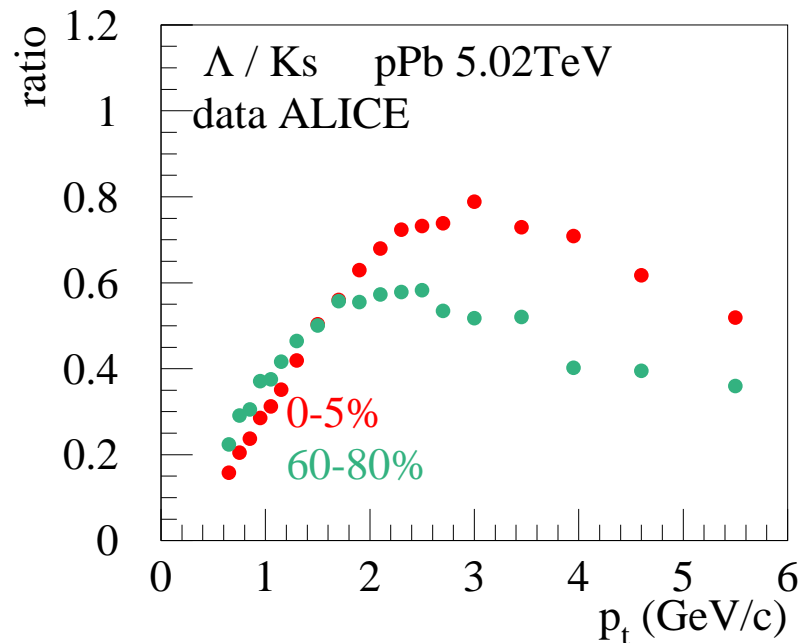
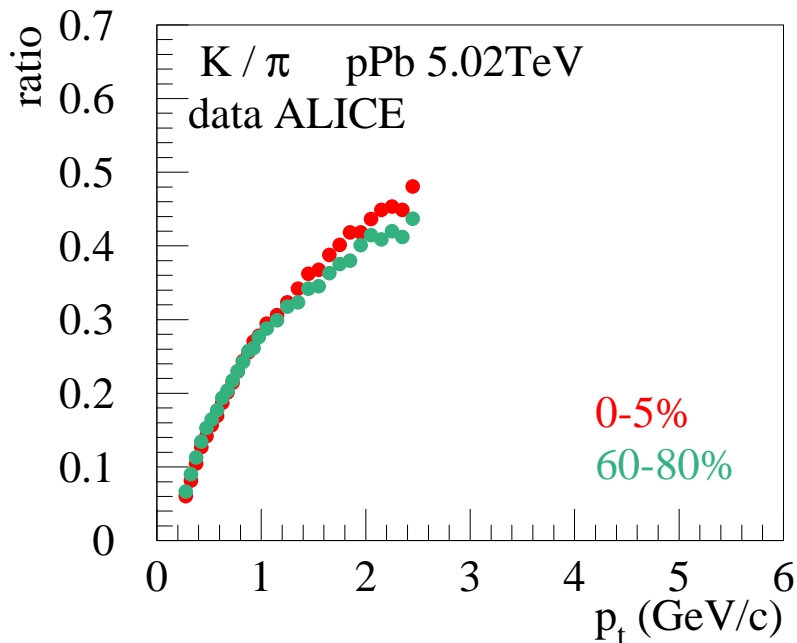


Strong variation of proton spectra => flow helps

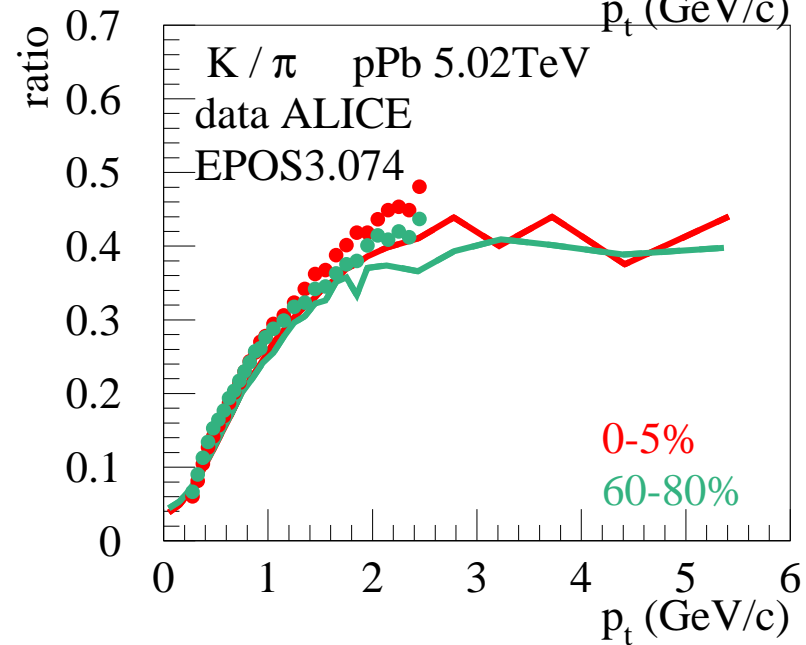
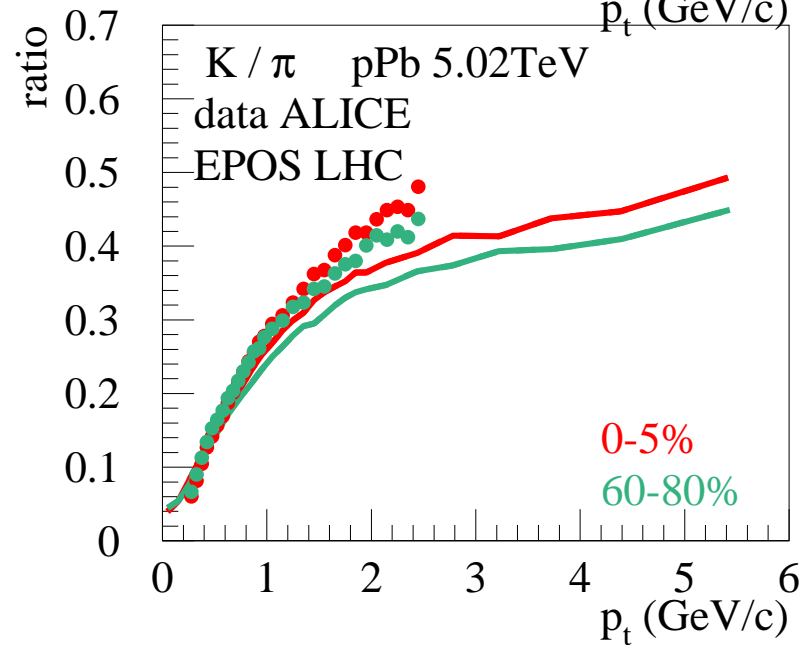
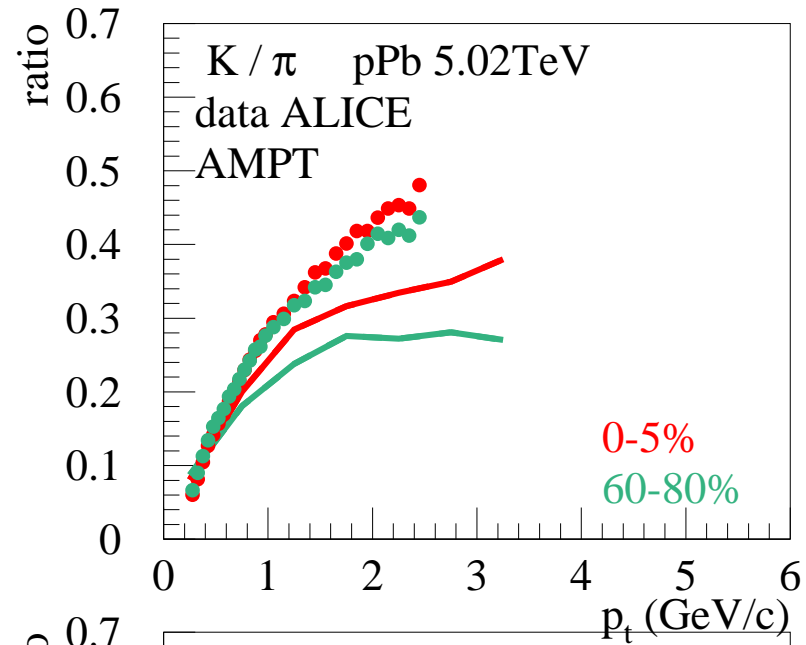
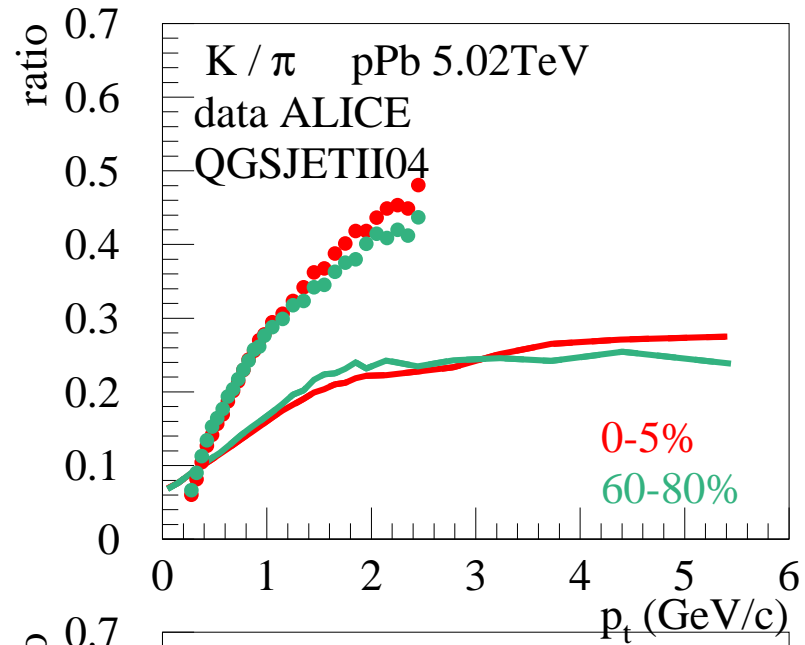
ALICE: compare p_t spectra for identified particles in different multiplicity classes: 0-5%,...,60-80%

(in $2.8 < \eta_{\text{lab}} < 5.1$) From R. Preghenella, ALICE, talk Trento workshop 2013

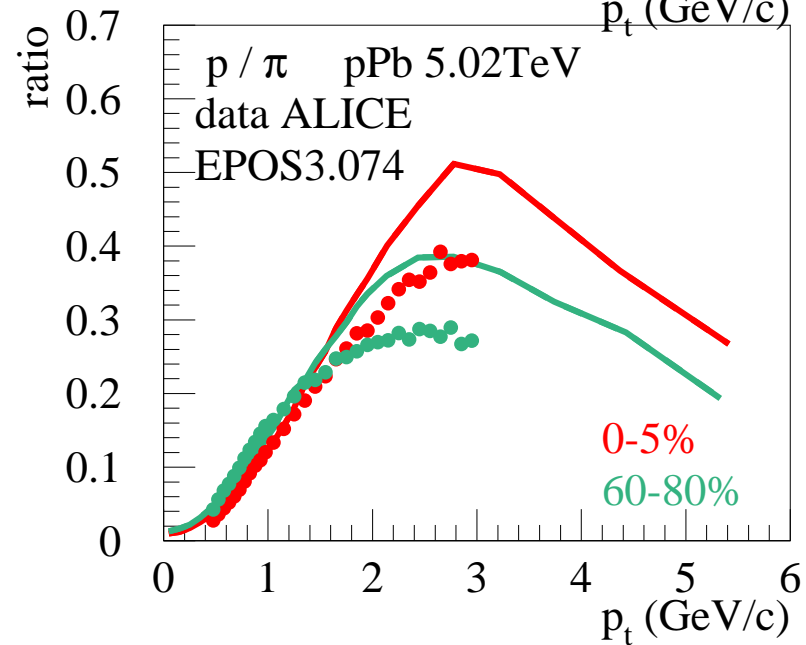
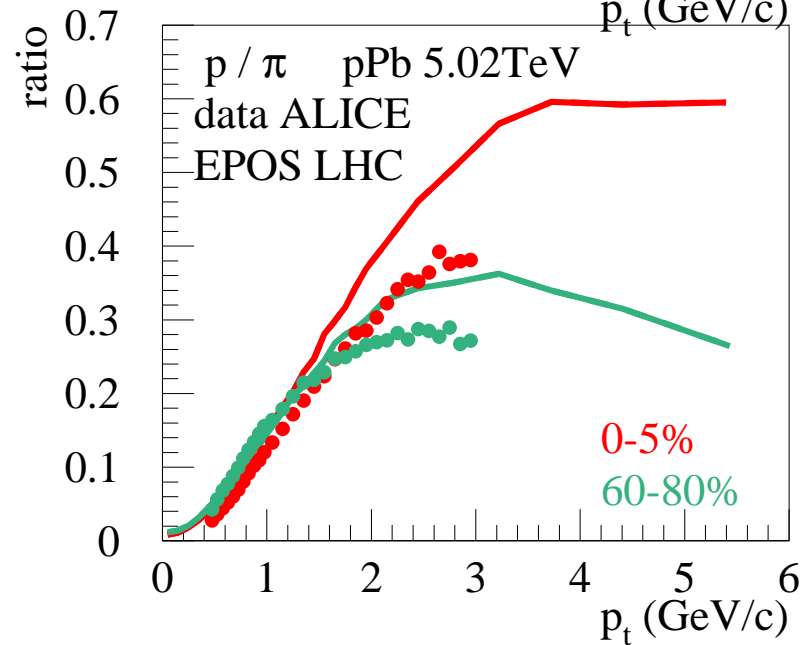
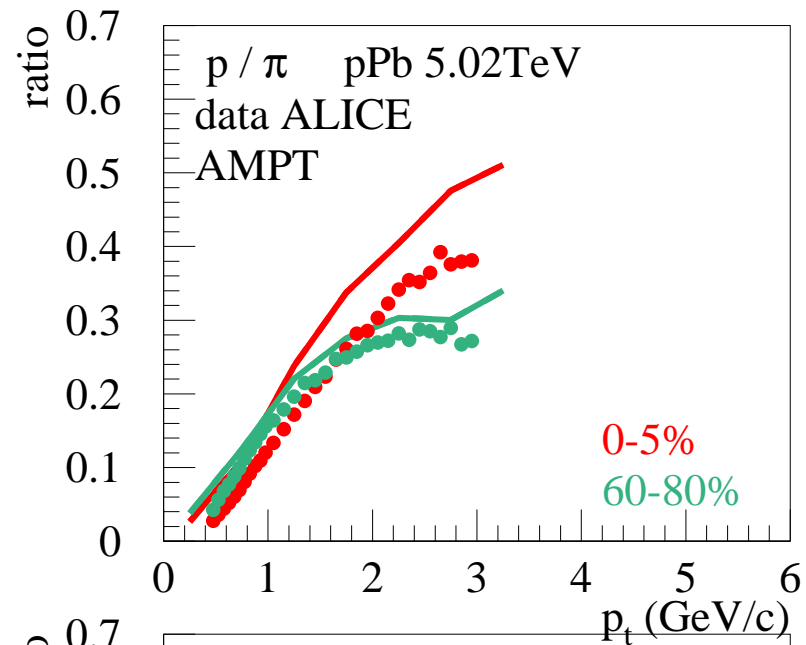
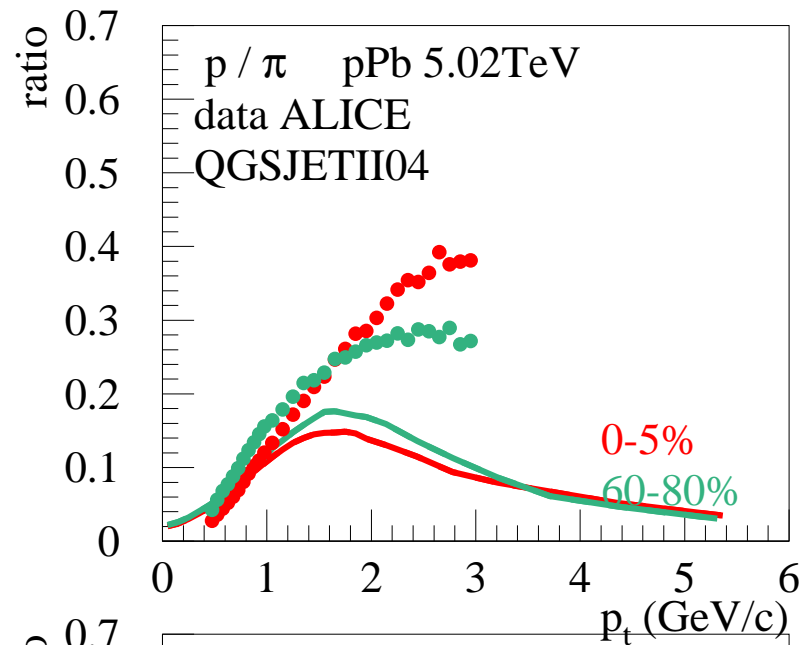
Useful : ratios (K/pi, p/pi...)



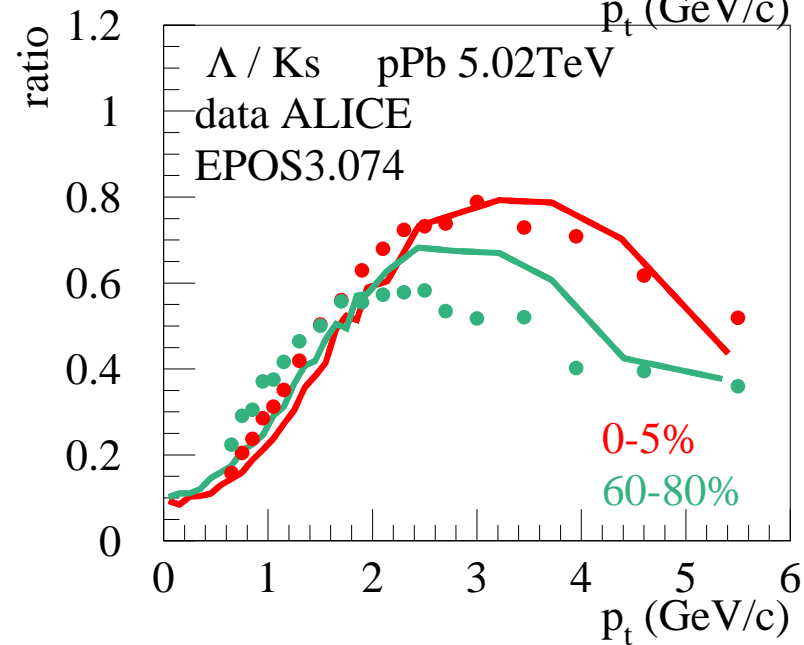
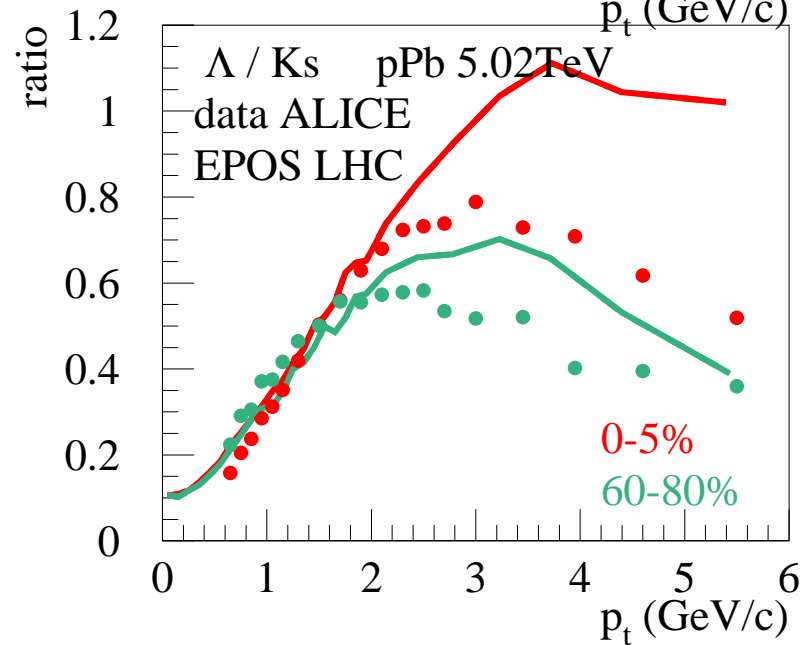
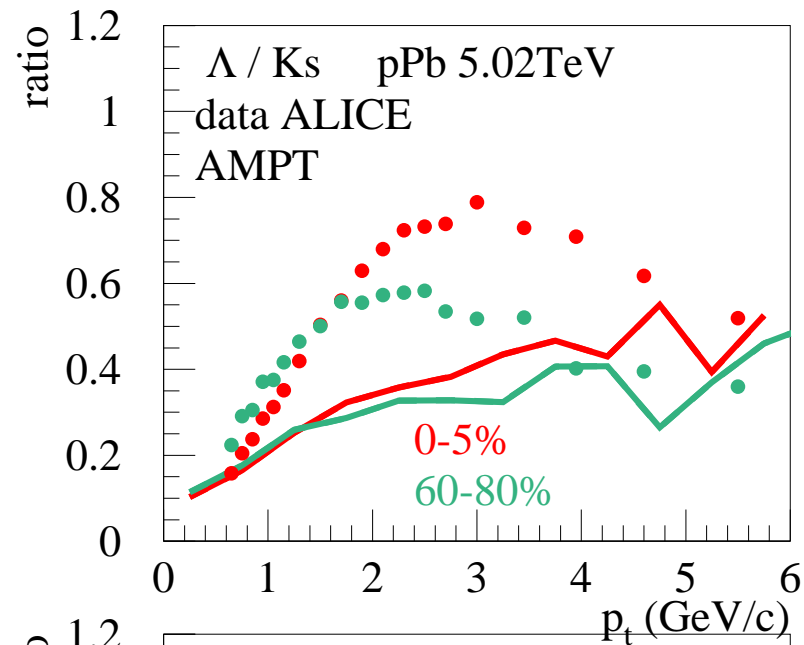
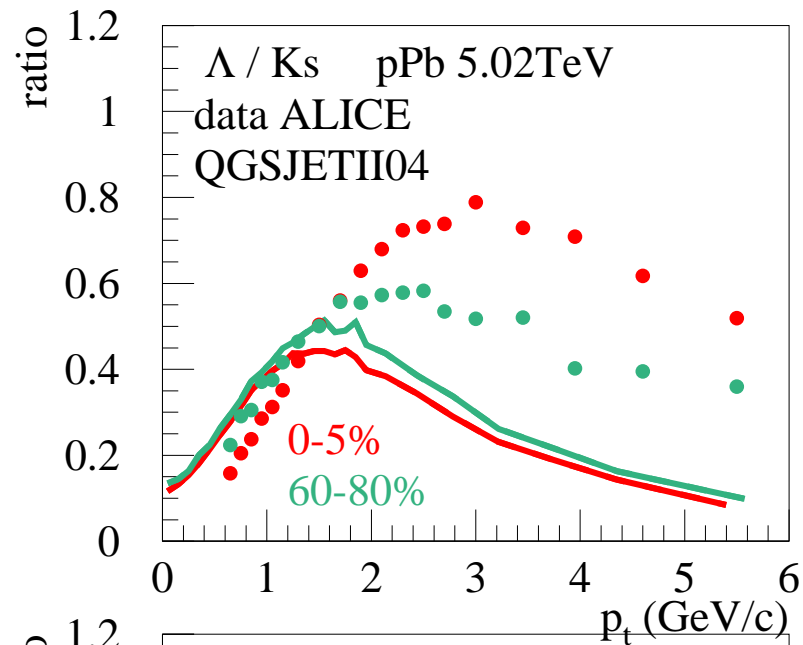
Significant variation of lambda/K – like in PbPb



No multiplicity dependence (not trivial to get the peripheral right)

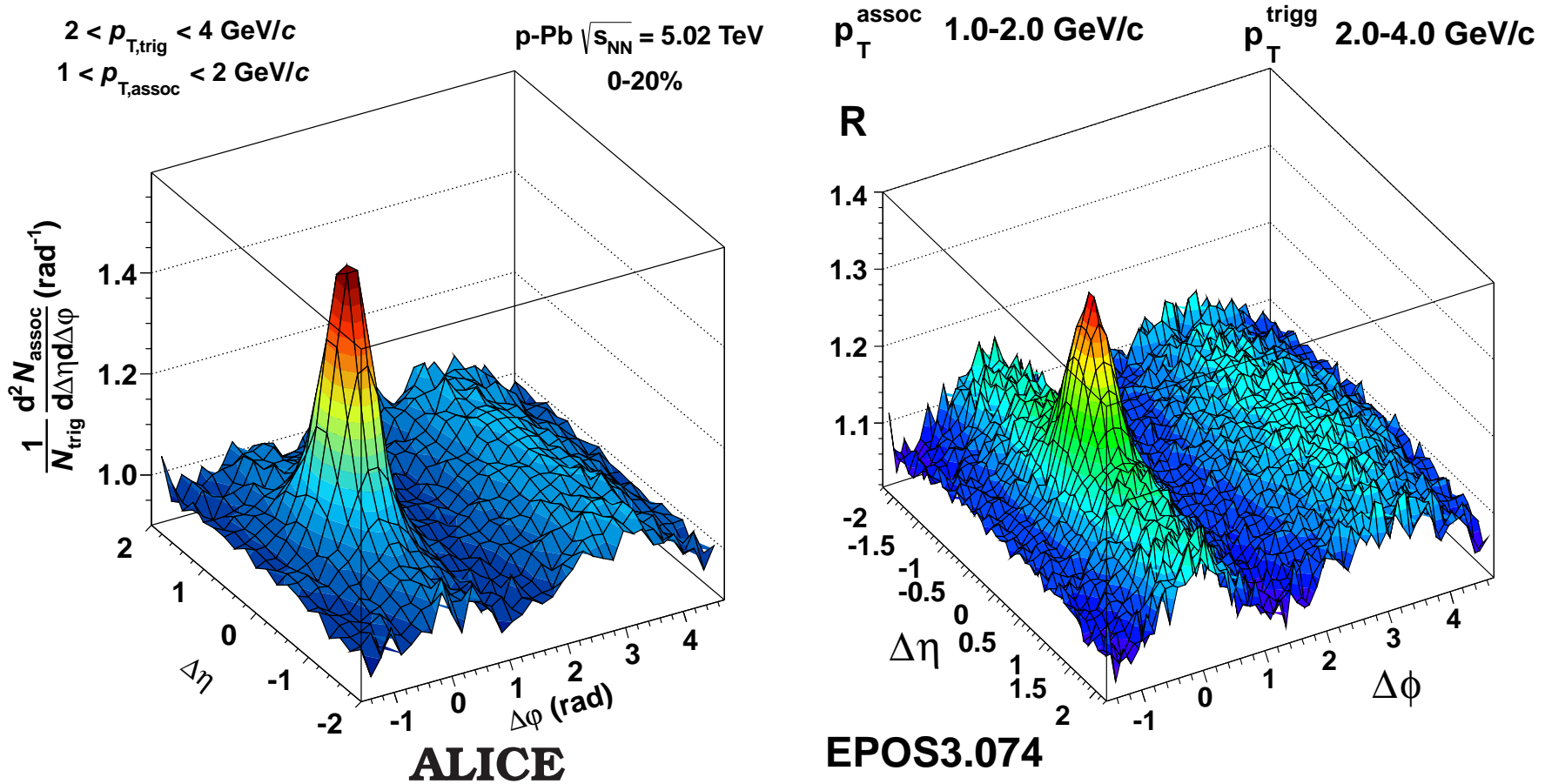


Significant multiplicity dependence



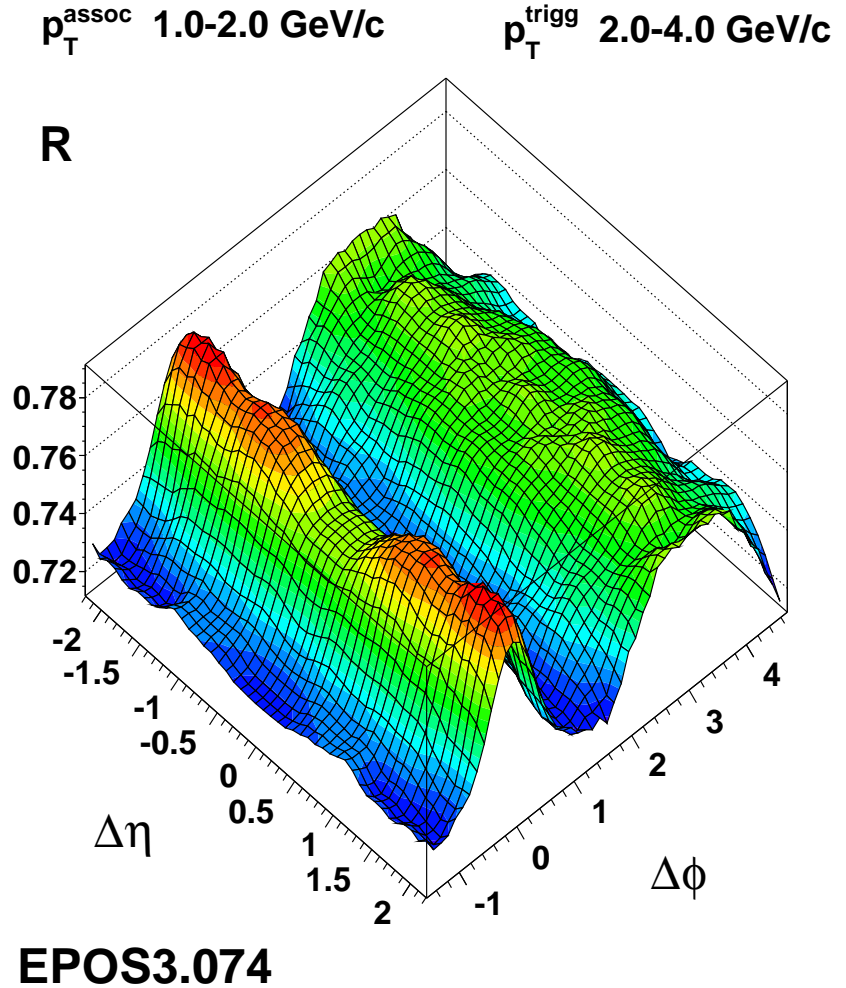
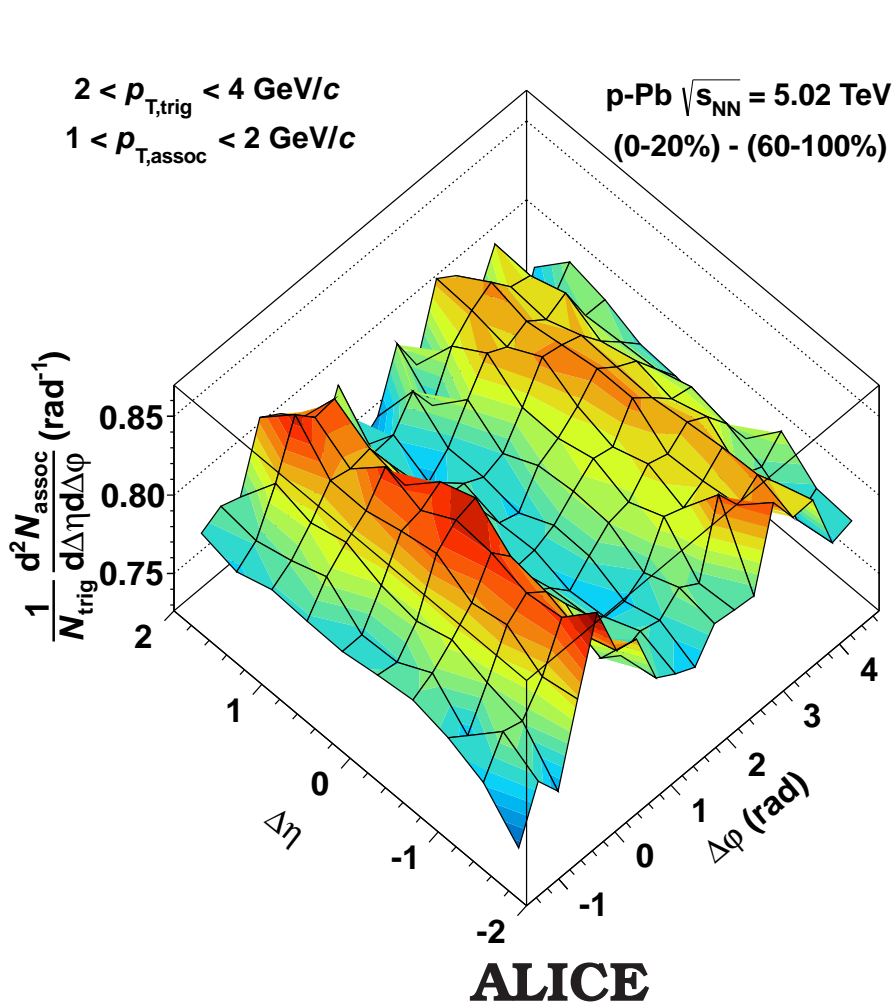
Significant multiplicity dependence => flow peak

“Ridges” in pA (EPOS3, arXiv:1307.4379)

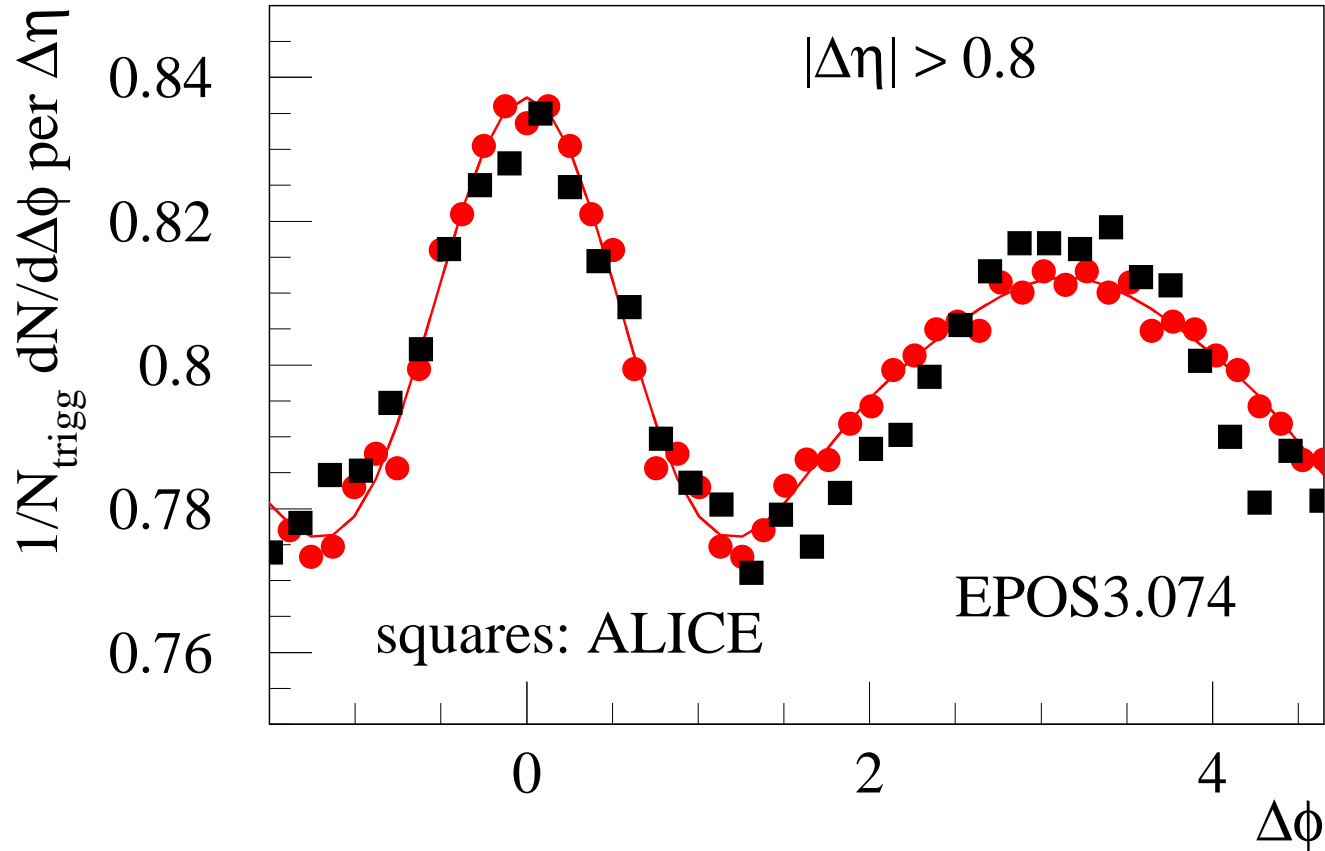


ALICE, Phys.Lett. B719 (2013) 29, arXiv:1212.2001

Central - peripheral (to get rid of jets)



Projection

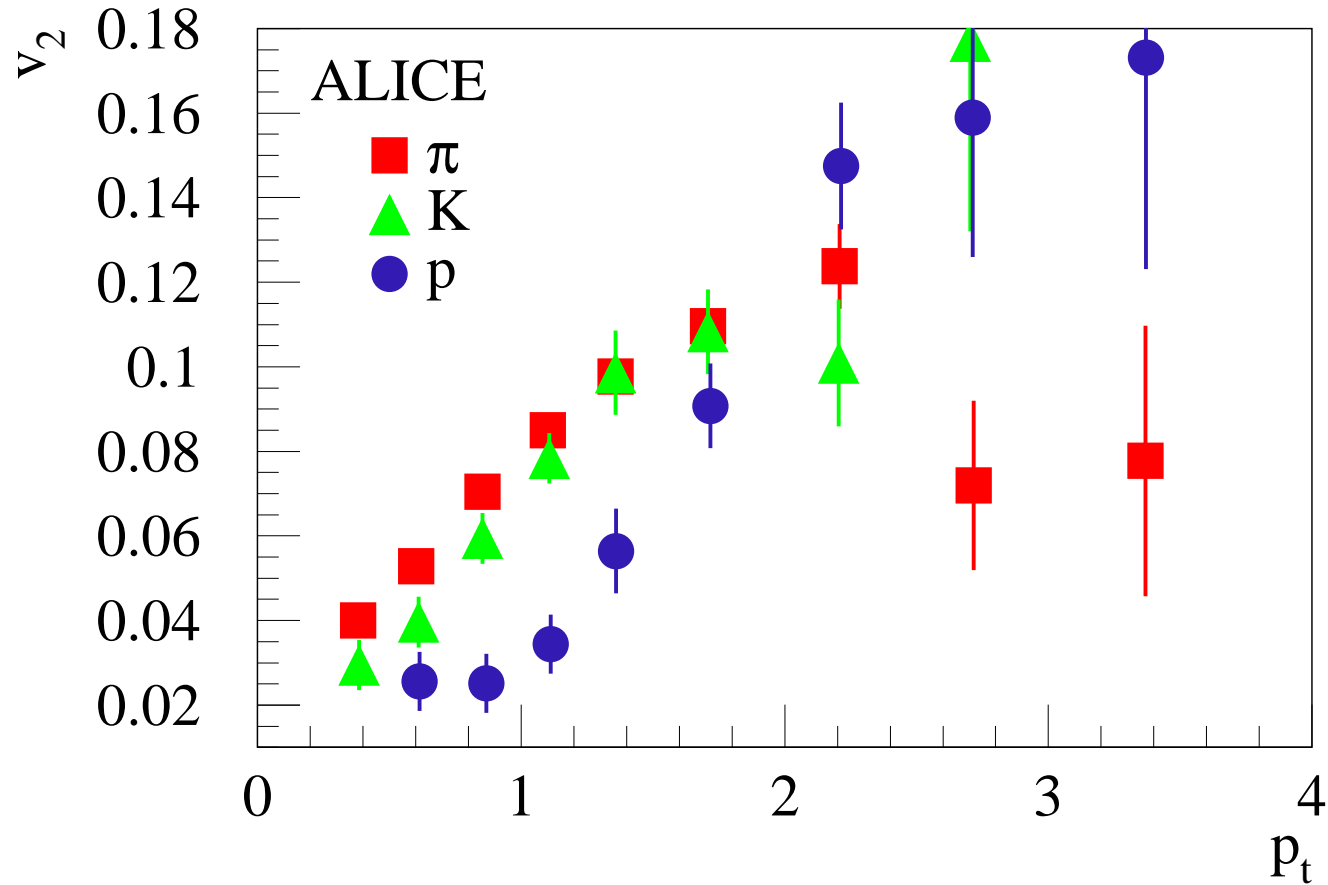


red line : $\sum 2a_i \cos(i\Delta\phi);$

$$\implies v_n = \sqrt{\frac{a_n}{b}}$$

Identified particle v_2

ALICE, arXiv:1307.3237



mass splitting, as in PbPb !!!

pPb in EPOS3

Pomerons (number and positions)
characterize geometry (P. number \propto multiplicity)

random

azimuthal

asymmetry

=>

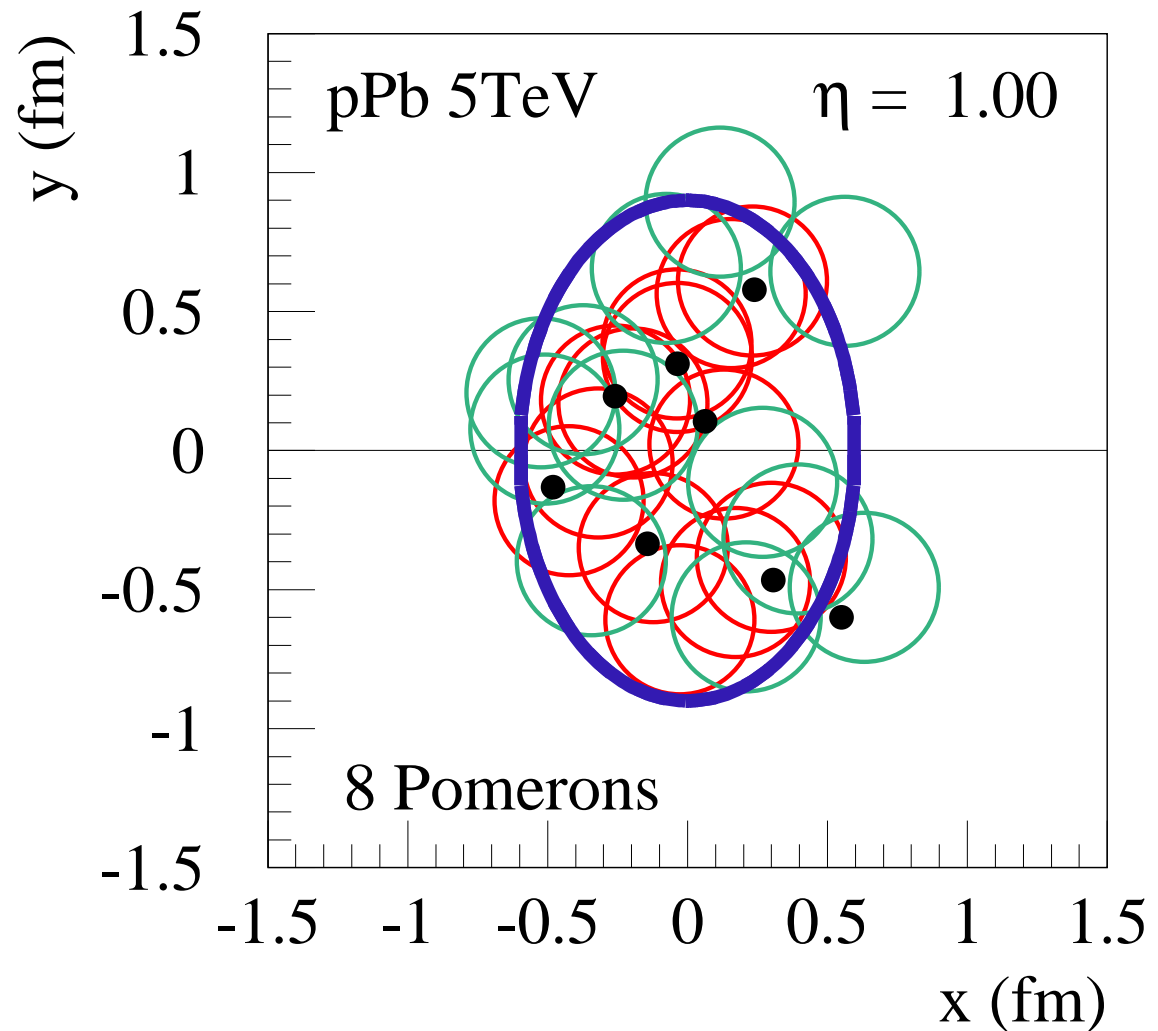
asymmetric flow

seen at higher

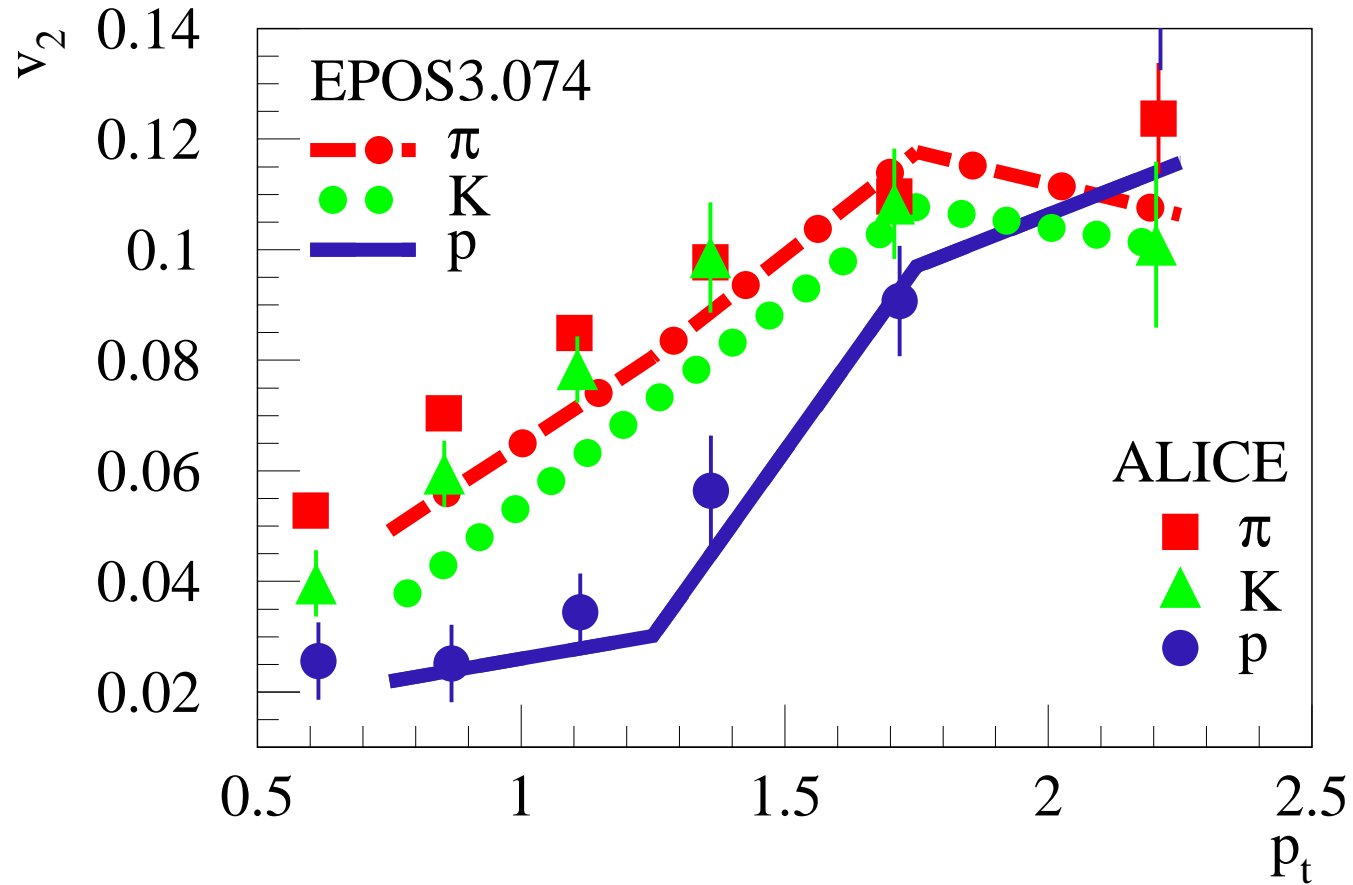
pt for heavier

ptls

Robust results

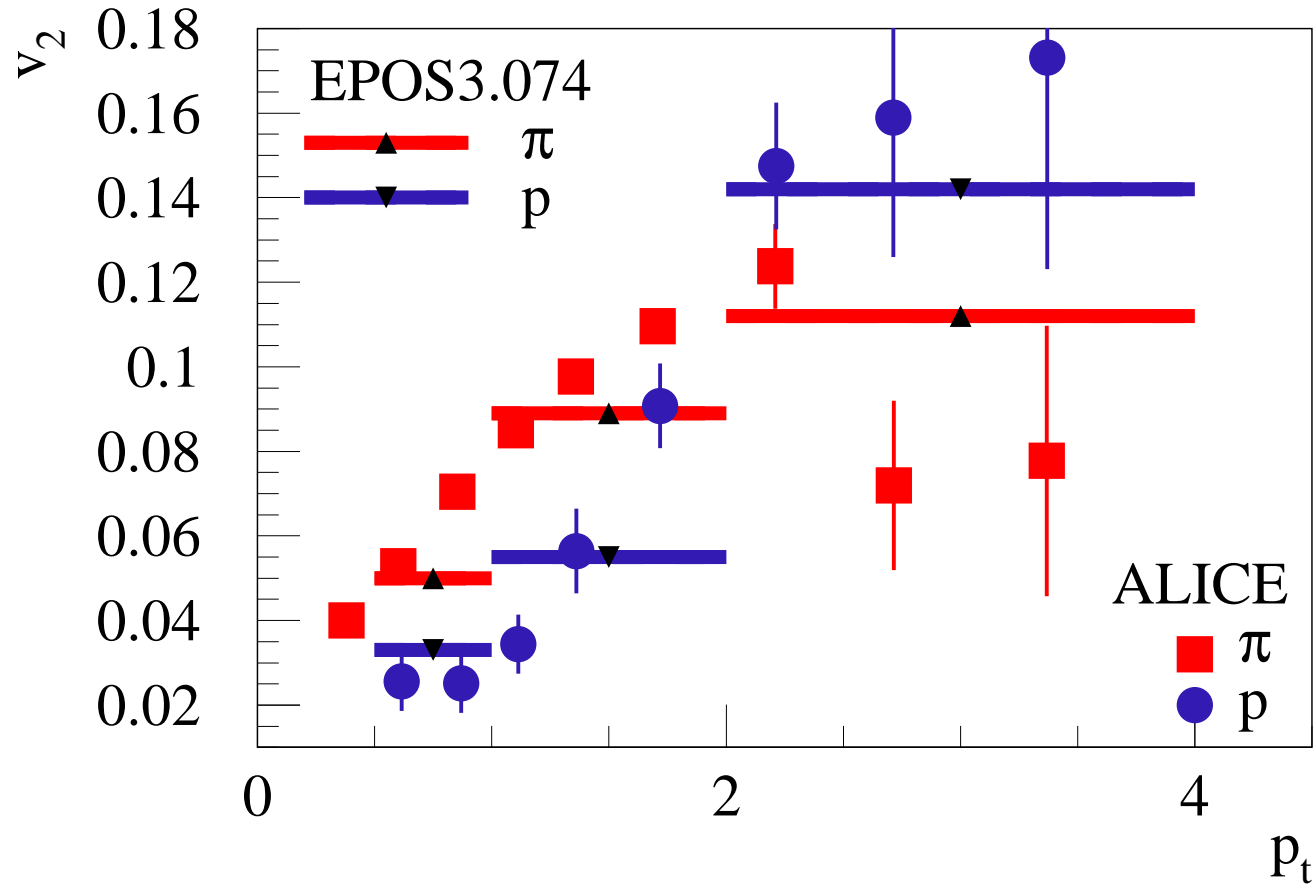


v_2 for π , K, p clearly differ



mass splitting, due to flow

different binning:



$v_2(\text{protons}) > v_2(\text{pions})$ beyond 2GeV

Summary

- **Traditional Monte Carlos have difficulties to describe strange particle production in pp.**
Hydrodynamical picture helps.
- **pPb looks very much like a hydrodynamically expanding system**
(more clean than PbPb, where hydro and minijets heavily interact, as well as the final hadrons among themselves)

