Electroweak boson production in heavy-ion collisions with CMS.

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Abstract. The Compact Muon Solenoid (CMS) is fully equipped to measure leptonic decays of electroweak probes in the high multiplicity environment of nucleus-nucleus collisions. Electroweak boson production is an important benchmark process at hadron colliders. Precise measurements of W and Z production in heavy-ion collisions can help to constrain nuclear PDFs as well as serve as a standard candle of the initial state in PbPb collisions at the LHC energies. The inclusive and differential measurements of the Z boson yield in the muon and electron decay channel will be presented, establishing that no modification is observed with respect to pp collisions, scaled by the number of incoherent nucleon-nucleon collisions.

1. Introduction

W and Z bosons were first observed by the UA1 and UA2 experiments at CERN thirty years ago in proton-antiproton collisions at $\sqrt{s} = 540$ GeV [1]. Since then a succession of electron-positron, proton-proton (pp) and proton-antiproton collider experiments have been used to measure their properties, such as mass, decay width and production cross sections at a variety of center-ofmass energies. Now for the first time, due to the large center-of-mass energy and luminosities offered at the LHC, the Z and W boson production can be studied in nucleus-nucleus collisions. Based on the first lead-lead (PbPb) collisions corresponding to 7.2 μ b⁻¹ integrated luminosity, the CMS collaboration reported first results on the Z $\rightarrow \mu\mu$ [2] and W $\rightarrow \mu\nu$ [3] processes, showing electroweak bosons are essentially unmodified by the hot and dense medium created in heavy-ion collisions, often referred to as the quark-gluon plasma (QGP).

Z-boson production is reported in this proceedings with the second set of PbPb collisions that occurred in late 2011 and accounts for 150 μ b⁻¹ integrated luminosity. In this document, the second set of pp data at the same energy in the center-of-mass, recorded in 2013, is used as a reference in order to calculate the nuclear modification factor, R_{AA} . The increases in statistics of both the pp and PbPb samples allow more detailed study of the transverse momentum (p_T) , rapidity (y) and collision centrality dependencies of the Z-boson yields. The dimuon channel is analyzed with an improved reconstruction algorithm with respect to Ref. [7], and electrons are used for the first time to reconstruct the Z decays in the heavy-ion environment in CMS [8].

Z bosons, once produced, decay within the medium with a typical lifetime of 0.1 fm/c. Since leptons pass freely through the strongly-interacting medium, dileptons from Z bosons can serve as a reference to the processes expected to be heavily modified in the QGP, such as quarkonium production, or the production of an opposite-side jet in Z+jet processes [4]. However, in heavy-ion collisions, Z-boson production can be slightly affected by initial-state effects. This modification is expected to be about 3% from the different mix of protons and neutrons in AA collisions (isospin effect) [5]. Energy loss and multiple scattering of the initial partons could produce modifications of the order of 3% [6]. In addition, the parton distribution functions can be depleted in nuclei, which could modify the Z-boson yield by at most 20% [5].

2. Lepton reconstruction and Signal Extraction

2.1. Muon reconstruction in CMS

With respect to the previous analysis reported in Ref. [7], a new muon offline reconstruction algorithm is used for the PbPb data, with increased efficiency due to an improvement in the track reconstruction. The single-muon reconstruction efficiency is raised from $\simeq 85\%$ to $\simeq 98\%$, for muons from Z decays, reaching the level of the algorithm used for pp collisions. Background muons from cosmic rays and heavy-quark semileptonic decays are rejected by requiring a transverse (longitudinal) distance of closest approach of less than 0.2 (5.0) mm from the measured vertex. These selection criteria further reject about 2% of real muons for both pp and PbPb collisions.

2.2. Electron reconstruction in CMS

The electron reconstruction uses information from the pixel and strip tracker, and the electromagnetic calorimeter (ECAL). Electrons produced in pp or PbPb collisions traversing the silicon layers of the tracker can emit bremsstrahlung photons and deposit energy in the ECAL with a significant spread in the azimuthal direction. An ECAL clustering algorithm and in particular the building of so-called *superclusters*, was used for collecting and estimating the proper energy of photons in the heavy-ion environment as in Ref. [9]. Pixels facing these superclusters are looked for, in order to initiate the building of trajectories in the inner tracker. A dedicated algorithm which takes into account the bremsstrahlung emission is used [10].

The reconstruction algorithm results in a reconstruction efficiency of 85% for electrons from Z decays. Further selection criteria reduce this efficiency by about 10%. For the pp sample, the standard algorithms and identification criteria presented in Ref. [11] are used, resulting in a reconstruction efficiency of 98% and a further reduction of 5% due to selection criteria.

2.3. Signal Extraction

Figures 1 and 2 show the dimuon and dielectron invariant mass spectra in the 60–120 GeV/c^2 mass range after applying selection criteria in pp and PbPb collisions. Only leptons with $p_T > 20$ GeV/c and $|\eta| < 2.4$ and 1.44 for the muons and the electrons, respectively, are considered. The dimuon rapidity is limited to |y| < 2, while the dielectron rapidity is kept in the |y| < 1.44 range. Overlaid histograms are from the Monte Carlo (MC) simulation.

In the $60-120 \text{ GeV}/c^2$ mass range of the PbPb sample, 1022 events are found with oppositecharge (OC, black solid circles) muons and not a single same charge (SC, black open squares) muon pair, while the pp sample has 830 OC and 1 SC pair. In the restricted dielectron rapidity range, 328 (388) OC electrons are measured in the PbPb (pp) data, with 27 (17) SC pairs resulting from higher fake and charge mis-identification rates. The SC pairs provide a measure of uncorrelated combinatorial background: negligible (1/1000 level) in the muon channel, and 8% (4)% in the electron channel. The number of Z candidates is taken as the OC–SC difference. Other sources of background are found to be negligible and covered by systematic uncertainties.

3. Results

Based on pp and PbPb Z-boson yields at the same energy in the center-of-mass, we compute the nuclear modification factors R_{AA} for both electron and muon channels as a function of the Z-boson p_T and rapidity, and event centrality, as the following: $R_{AA} = N_{AA}/T_{AA} \times \sigma_{pp}$, where



Figure 1. Dimuon invariant mass spectra for muons with $|\eta^{\mu}| < 2.4$ and $p_T^{\mu} > 20$ GeV/c, and dimuons with |y| < 2.0 in pp (left) and in PbPb (right) collisions. Full black circles are oppositecharge pairs, open black squares are same-charge pairs. Superimposed is the MC simulation from PYTHIA pp $\rightarrow Z \rightarrow \mu^+\mu^-$ embedded in HYDJET for the PbPb case.



Figure 2. Dielectron invariant mass spectra for electrons with $|\eta^e| < 1.44$ and $p_T^e > 20$ GeV/c, and dielectrons with |y| < 1.44 in pp (left) and in PbPb (right) collisions.

 $N_{\rm AA}$ are the Z boson yields per event estimated from PbPb collisions corrected for acceptance and efficiency measured in simulation, and $\sigma_{\rm pp}$ the differential cross sections estimated from pp collisions by dividing the corrected Z-boson yields by the integrated luminosity. The nuclear overlap function $T_{\rm AA}$, computed with a Glauber model [12] is proportional to the number of elementary nucleon-nucleon binary collisions $N_{\rm coll} = T_{\rm AA} \times \sigma_{\rm NN}$ where the latter is an assumed inelastic nucleon-nucleon cross section that cancels to first order in $T_{\rm AA}$.

The data are divided in different independent ranges: 6 in event centrality, 8 (5) in dimuon (dielectron) rapidity, and 7 in Z-boson p_T . Figure 3 shows the obtained AA as a function of the Z-boson p_T , rapidity and event centrality (translated to the average number of participants (N_{part}) , through the Glauber model mentioned above). The T_{AA} uncertainties, displayed as vertical green bands are common to all points when centrality-integrated, and to the dielectron and dimuon points when centrality-binned. The pp luminosity uncertainty, common to all points is displayed as a vertical grey band.

Within our statistical and systematic uncertainties (dominated by uncertainty in the efficiency from a data-driven Tag and Probe method) the nuclear modification factors in both the electron and the muon channels as a function of p_T , rapidity and event centrality are compatible with unity and no strong deviations are observed. The integrated R_{AA} values are $1.06 \pm 0.05(stat) \pm 0.11(syst)$ in the dimuon channel and $1.08 \pm 0.09(stat) \pm 0.14(syst)$ for dielectrons, and the differential values are shown in Figure 3.



Figure 3. Nuclear modification factor (R_{AA}) extracted from PbPb and pp collisions at $\sqrt{s_{NN}} = 2.76$ TeV for $Z \rightarrow e^+e^-$ and $Z \rightarrow \mu^+\mu^-$ as a function of Z-boson p_T (left), rapidity (center) and number of participants in the collision (right).

4. Conclusions

The Z-boson R_{AA} has been measured from PbPb and pp collisions at $\sqrt{s_{NN}} = 2.76$ TeV. This was done as a function of Z-boson transverse momentum, rapidity and event centrality in both the dimuon and dielectron decay channels with 150 μb^{-1} integrated luminosity. No significant deviation is observed, underlying that within uncertainties, no nuclear effects are observed.

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