

Superconformal Operator Product Expansion and General Gauge Mediation

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with Kenneth Intriligator and Andreas Stergiou

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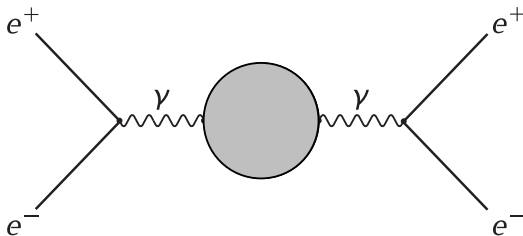
(S)OPE UV/IR applications

- Non-CFTs \Rightarrow Conformally covariant OPEs
- Softly broken symmetries as spontaneously broken symmetries
 - Symmetry breaking seen as IR effect (via field or spurion vev)
 - Symmetry restored in UV theory \Rightarrow (S)OPE selection rules

\Rightarrow Strongly-coupled IR physics described by weakly-coupled UV physics through (S)OPE

- Example: QCD
 - Not conformal (non-trivial RG flow)
 - IR physics \Rightarrow Theory with chiral symmetry breaking ($\langle \bar{Q}Q \rangle \neq 0$) and confinement ($\langle G_{\mu\nu}^A G^{A\mu\nu} \rangle \neq 0$)
 - UV physics \Rightarrow Asymptotically free CFT \Rightarrow QCD sum rules [Shifman, Vainshtein, Zakharov \(1979\)](#)

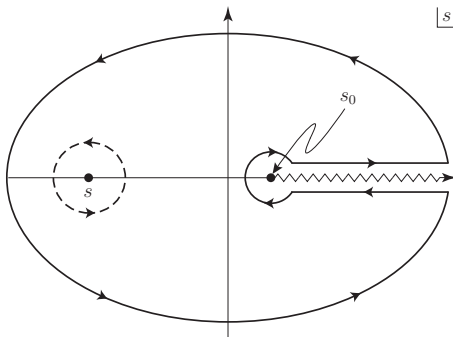
Probe hadron sector from lepton sector through gauge interactions
(UV/IR physics separated using OPE)



$$\Pi_{\text{hadron},\mu\nu}(p) = ie^2 \int d^4x e^{-ip \cdot x} \langle j_\mu(x) j_\nu(0) \rangle$$

$$\mathcal{O}_i(x) \mathcal{O}_j(0) = \sum_k c_{ij}^k(x) \mathcal{O}_k(0)$$

OPE, analyticity and optical theorem \Rightarrow Relations between total cross section and OPE coefficients



$$\Pi_{\text{hadron}}(s) = \frac{1}{2\pi i} \int_{s_0}^{\infty} ds' \frac{\text{Disc } \Pi_{\text{hadron}}(s')}{s' - s}$$

$$\sigma_{e^+e^- \rightarrow \text{hadrons}}(s) = \frac{2\pi i \alpha}{s} \text{Disc } \Pi_{\text{hadron}}(s)$$

Operator product expansion: Review

Operator product expansion

$$\begin{aligned}\mathcal{O}_i(x)\mathcal{O}_j(0) &= \sum_k \frac{c_{ij}^k}{x^{\Delta_i+\Delta_j-\Delta_k}} \mathcal{O}_k(0) \\ &= \sum_{\text{primary } k} \frac{c_{ij}^k}{x^{\Delta_i+\Delta_j-\Delta_k}} F_{\Delta_i\Delta_j}^{\Delta_k}(x, P) \mathcal{O}_k(0)\end{aligned}$$

- Short distance physics expressed in terms of local operators
- Wilson coefficients \Rightarrow UV physics
- Vacuum expectation values \Rightarrow IR physics
- OPE constrained by conformal symmetry \Rightarrow Wilson coefficients of descendants determined by Wilson coefficients of primaries [Ferrara, Gatto, Grillo \(1971\)](#)

Superconformal operator product expansion

$$\mathcal{O}_i(x)\mathcal{O}_j(0) \stackrel{?}{=} \sum_{\text{superprimary } k} \frac{c_{ij}^k}{x^{\Delta_i+\Delta_j-\Delta_k}} F_{ij}^k(x, P, Q, \bar{Q}) \mathcal{O}_k(0)$$

$$\mathcal{T}_\mu(z) = \textcolor{red}{j}_\mu^R(\textcolor{red}{x}) + \theta^\alpha \mathcal{S}_{\alpha\mu}(x) + \bar{\theta}^{\dot{\alpha}} \bar{\mathcal{S}}_{\dot{\alpha}\mu}(x) + 2\theta^\nu \bar{\theta} T_{\nu\mu}(x) + \dots$$

$$\mathcal{J}(z) = \textcolor{red}{J}(\textcolor{red}{x}) + i\theta j(x) - i\bar{\theta} \bar{j}(x) - \theta\sigma^\mu \bar{\theta} j_\mu(x) + \dots$$

- OPE constrained by superconformal symmetry
 - Sum over superprimaries instead of primaries
- ⇒ Wilson coefficients of superdescendants NOT fully determined by Wilson coefficients of superprimaries (existence of superconformal 3-point invariants) ! Osborn (1998)

Current-current SOPE

Superconformal 3-point correlation functions for current superfields

$$\langle \mathcal{J}(z_1) \mathcal{J}(z_2) \mathcal{O}^{\mu_1 \dots \mu_\ell}(z_3) \rangle = \frac{1}{x_{\bar{1}3}^2 x_{\bar{3}1}^2 x_{\bar{2}3}^2 x_{\bar{3}2}^2} t_{\mathcal{J}\mathcal{J}\mathcal{O}}^{\mu_1 \dots \mu_\ell}(X_3, \Theta_3, \bar{\Theta}_3)$$

- $\mathcal{O}^{\mu_1 \dots \mu_\ell}$ real spin- ℓ superfield with vanishing R-charge
- $t_{\mathcal{J}\mathcal{J}\mathcal{O}}(X, \Theta, \bar{\Theta})$ fully determined by superconformal symmetry and current superconservation [JFF, Intriligator, Stergiou \(2011\)](#)

⇒ (Non-universal) 4-point $\langle J(x_1) J(x_2) J(x_3) J(x_4) \rangle$
superconformal blocks

General gauge mediation: Overview

- SUSY breaking hidden sector connected to visible sector through gauge interactions [Buican, Meade, Seiberg, Shih \(2008\)](#)
 - Decoupled hidden sector in $g_{\text{SM}} \rightarrow 0$ limit
 - Universal visible sector SUSY breaking effects introduced via loops

⇒ Current-current correlation functions (even without hidden sector Lagrangian)

$$\mathcal{J}(z) = J(x) + i\theta j(x) - i\bar{\theta}\bar{j}(x) - \theta\sigma^\mu\bar{\theta}j_\mu(x) + \dots$$

$$\langle J(x)J(0) \rangle, \langle j_\alpha(x)\bar{j}_{\dot{\alpha}}(0) \rangle, \langle j_\mu(x)j_\nu(0) \rangle, \langle j_\alpha(x)j_\beta(0) \rangle$$

Cross sections

OPE constraints, analyticity and optical theorem

$$\sigma_{\text{visible} \rightarrow D^* \rightarrow \text{hidden}}(s) = -\frac{(4\pi\alpha_{\text{SM}})^2}{s} \sum_k \text{Im}[\tilde{c}_{JJ}^k(s)] \langle \mathcal{O}_k(0) \rangle$$

$$\sigma_{\text{visible} \rightarrow \lambda_\alpha^* \rightarrow \text{hidden}}(s) = f_{1/2}(\text{Im}[\tilde{c}_{JJ}^k(s)])$$

$$\sigma_{\text{visible} \rightarrow A_\mu^* \rightarrow \text{hidden}}(s) = f_1(\text{Im}[\tilde{c}_{JJ}^k(s)])$$

- Consistent with direct computation in ordinary minimal gauge mediation [Martin \(1996\)](#)
- Good approximation with first few terms

Visible sector SUSY breaking masses

OPE constraints and analyticity

$$M_{\text{gaugino}} \approx \sum_k \frac{\alpha_{\text{SM}} \text{Im}[s^{d_k/2} \tilde{c}_{JJ}^k(s)]}{2^{d_k-1} d_k M^{d_k}} \langle Q^2(\mathcal{O}_k(0)) \rangle$$

$$m_{\text{sfermion}}^2 \approx 4\pi\alpha_{\text{SM}} Y \langle J(x) \rangle - \sum_k \frac{\alpha_{\text{SM}}^2 c_2 \text{Im}[s^{d_k/2} \tilde{c}_{JJ}^k(s)]}{2^{d_k+1} \pi d_k^2 M^{d_k}} \langle \bar{Q}^2 Q^2(\mathcal{O}_k(0)) \rangle$$

- Exactly equivalent to the usual $f(x)$ and $g(x)$ functions of ordinary minimal gauge mediation [Martin \(1996\)](#)
- Good approximation with first few terms

Other research interests

Scale and conformal invariance

- Non-conformal scale-invariant theories
 - RG flow recurrent behaviors
 - Consistent with weak and strong version of c -theorem
- Generalized c -theorem
 - Conformal recurrent behaviors

Dark matter

- Model building
- Constraints
 - Direct and indirect detection experiments
 - Collider experiments
 - Astrophysical observations (e.g. white dwarfs)