

Neutrino Physics

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Why are neutrinos so interesting ?

➤ **Cosmology:**

They played an important role during the Big Bang, they could explain the asymmetry among matter and anti-matter, they are the most abundant particles in the universe

➤ **Astrophysics:**

They are ruling the life and death of stars

➤ **Particle Physics:**

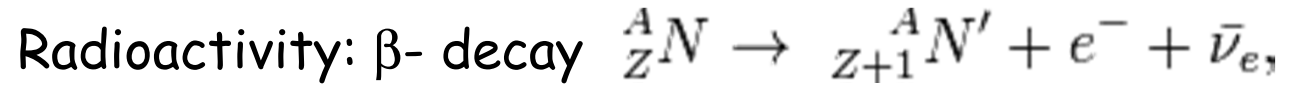
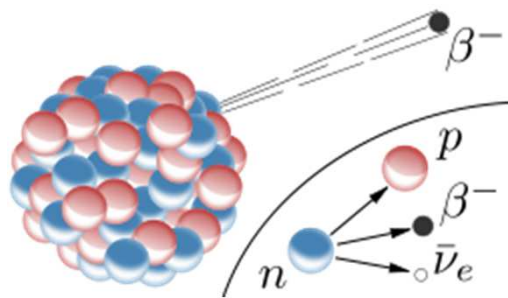
They are a window on physics beyond the Standard Model of particle physics: presently they represent the only experimental hint in that direction

Many properties of neutrinos were totally unexpected coming out as experimental results:

The history of neutrino physics is a real saga with an extraordinary richness of experimental techniques involved

There are still a lot of open questions in neutrino physics ...

The birth of the neutrino as a « **desperate remedy** » to solve apparent energy non-conservation in β decays (W. Pauli 1930)



Early 1900s: people thought they were dealing with a two body decay process:

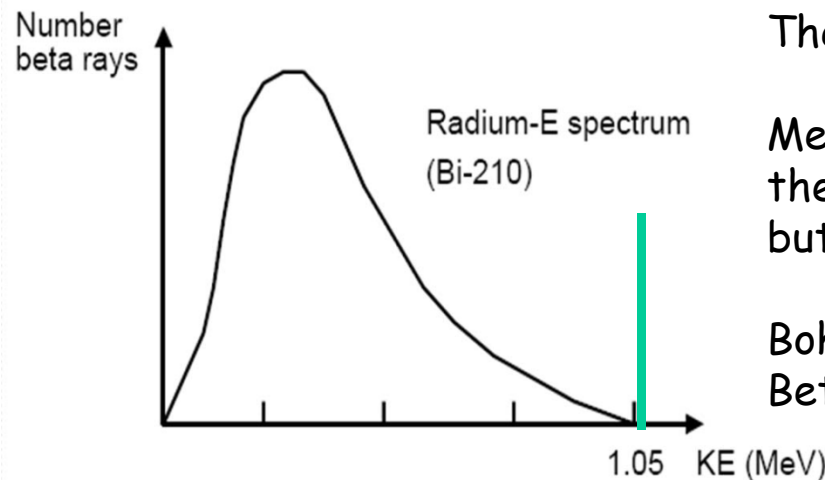
$$(A, Z) \rightarrow (A, Z + 1) + e^-$$

→ The energy spectrum of the electrons should be monochromatic:

$$p = \frac{1}{2M} \sqrt{[M^2 - (m_1 - m_2)^2][M^2 - (m_1 + m_2)^2]} \quad \text{with } M=M(A, Z), m_1=M(A, Z+1) \text{ and } m_2=m_e$$

First measurements of beta spectrum:

1911 Lise Meitner and Otto Hahn, 1914 Ellis and Chadwick



The beta spectrum is continuous

Meitner: electrons reinteract in the nuclei emitting gamma rays → but no gamma rays detected

Bohr: energy is not conserved in Beta decay !!!

Additional problem: the model of the nucleus (made of protons+electrons) and the spin of nuclei (Li and N) measured to be integer

Li nucleus: 6 protons + 3 electrons = 9 fermions

N nucleus: 14 protons + 7 electrons = 21 fermions



Original - Photostigma of 24. 0393
Abschrift/15.12.56 PM

Offener Brief an die Gruppe der Radioaktiven bei der
Gesellschafts-Tagung zu Tübingen.

Abschrift

Physikalisches Institut
der Eidg. Technischen Hochschule
Zürich

Zürich, 4. Dez. 1930
Uraniastrasse

Liebe Radioaktive Damen und Herren,

Wie der Ueberbringer dieser Zeilen, den ich halbvollst
anzuhören bitte, Ihnen des näheren auseinandersetzen wird, bin ich
angesichts der "falschen" Statistik der N- und Li-6 Kerne, sowie
des kontinuierlichen beta-Spektrums auf einen verzwweifelten Ausweg
verfallen um den "Wechselstz" (1) der Statistik und den Energienatz
zu retten. Nämlich die Möglichkeit, es könnten elektrisch neutrale
Teilchen, die ich Neutronen nennen will, in den Kernen existieren,
welche den Spin 1/2 haben und das Ausschliessungsprinzip befolgen und
sich von Lichtquanten ausserdem noch dadurch unterscheiden, dass sie
nicht mit Lichtgeschwindigkeit laufen. Die Masse der Neutronen
musste von derselben Grössenordnung wie die Elektronenmasse sein und
jedenfalls nicht grösser als 0,01 Protonenmasse. Das kontinuierliche
beta-Spektrum wäre dann verständlich unter der Annahme, dass beim
beta-Zerfall mit dem Elektron jeweils noch ein Neutron emittiert
wird, derart, dass die Summe der Energien von Neutron und Elektron
konstant ist.

1930: W. Pauli makes the hypothesis of an undetectable particle sharing the energy of beta decay with the emitted electron.

From Pauli's letter of the 4th of December 1930

Dear Radioactive Ladies and Gentlemen,

As the bearer of these lines, to whom I graciously ask you to listen, will explain to you in more detail, how because of the "wrong" statistics of the N and Li^6 nuclei and the continuous beta spectrum, I have hit upon a desperate remedy to save the "exchange theorem" of statistics and the law of conservation of energy. Namely, the possibility that there could exist in the nuclei electrically neutral particles, that I wish to call neutrons, which have spin $1/2$ and obey the exclusion principle and which further differ from light quanta in that they do not travel with the velocity of light. The mass of the neutrons should be of the same order of magnitude as the electron mass and in any event not larger than 0.01 proton masses. The continuous beta spectrum would then become understandable by the assumption that in beta decay a neutron is emitted in addition to the electron such that the sum of the energies of the neutron and the electron is constant... I agree that my remedy could seem incredible because one should have seen these neutrons much earlier if they really exist. But only the one who dare can win and the difficult situation, due to the continuous structure of the beta spectrum, is lighted by a remark of my honoured predecessor, Mr Debye, who told me recently in Bruxelles: "Oh, It's well better not to think about this at all, like new taxes". From now on, every solution to the issue must be discussed. Thus, dear radioactive people, look and judge. Unfortunately, I cannot appear in Tübingen personally since I am indispensable here in Zurich because of a ball on the night of 6/7 December. With my best regards to you, and also to Mr Back.

Your humble servant,

W. Pauli

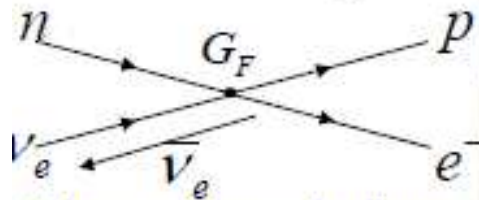
Today I have done something which no theoretical physicist should ever do in his life: I have predicted something which shall never be detected experimentally



1932 The neutron (as we know it today) was discovered, by J. Chadwick, two years after Pauli's proposal

- Solves nuclear spin problem: $A = Z(\text{protons}) + N(\text{neutrons})$
- But the mass of the neutron is similar to the proton mass
- cannot be the Pauli's particle

Fermi, 1933: coherent theory of beta decay



“Abstract speculations too far from physical reality to be of any interest to the readers”

Nature, rejecting the paper!

E. Fermi, *La Ricerca Scientifica* 4 (II), (1933), 491-495; and *Z.Physik*, 88 (1934) 161



Pauli thought his proposal of the "neutron" was too speculative, he did not publish it in a scientific journal until 1934, by which time Fermi had already developed his theory of beta decay incorporating the neutrino.

Fermi 4-fermion contact interaction, Lagrangian of interaction (in analogy with electrodynamics):

$$\mathcal{L}(x) = -\frac{G_F}{\sqrt{2}} \left[\bar{\phi}_p(x) \gamma^\mu \phi_n(x) \right] \left[\bar{\phi}_e(x) \gamma^\mu \phi_\nu(x) \right]$$

$G_F =$ Fermi coupling constant = $(1.16637 \pm 0.000001) 10^{-5} \text{ GeV}^{-2}$

In 1934, at a seminar Fermi was asked whether the neutral particle emitted in the nuclear beta-decay was the same as Chadwick's neutron.

Fermi clarified that he was talking about a different particle which he referred to as **neutrino** ("little neutral one").

Inverse beta decay process as a tool for neutrino detection:

$$\bar{\nu} + N(n, p) \rightarrow e^+ + N(n+1, p-1)$$

where n equals the number of neutrons and p equals the number of protons. If the nucleus happens to be that of hydrogen (a single proton), then the interaction produces a neutron and a positron:

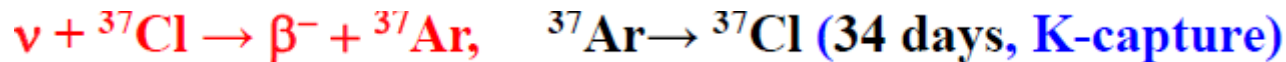
$$\bar{\nu} + p \rightarrow n + e^+$$

Bethe & Peierls, computing its cross section, 1934:

Few MeV neutrinos resulted to have an interaction length of about one light year of lead: “..this meant that one obviously would never be able to see a neutrino.”



1946 Pontecorvo proposes the following reaction for neutrino detection

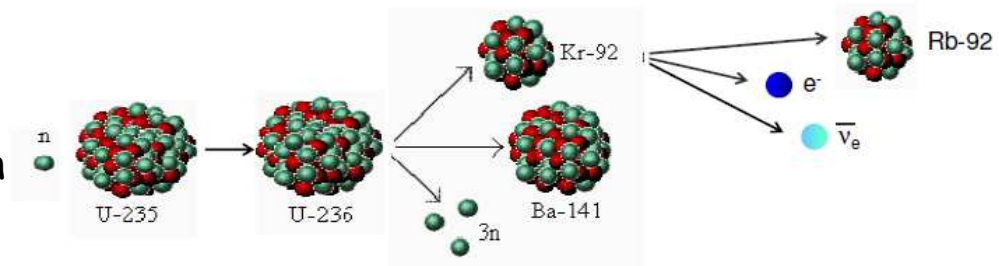


Davis exploits the Pontecorvo reaction $\nu + {}^{37}\text{Cl} \rightarrow \beta^- + {}^{37}\text{Ar}$:

- 1955-58, antineutrinos from reactor (Brookhaven, Savannah River)

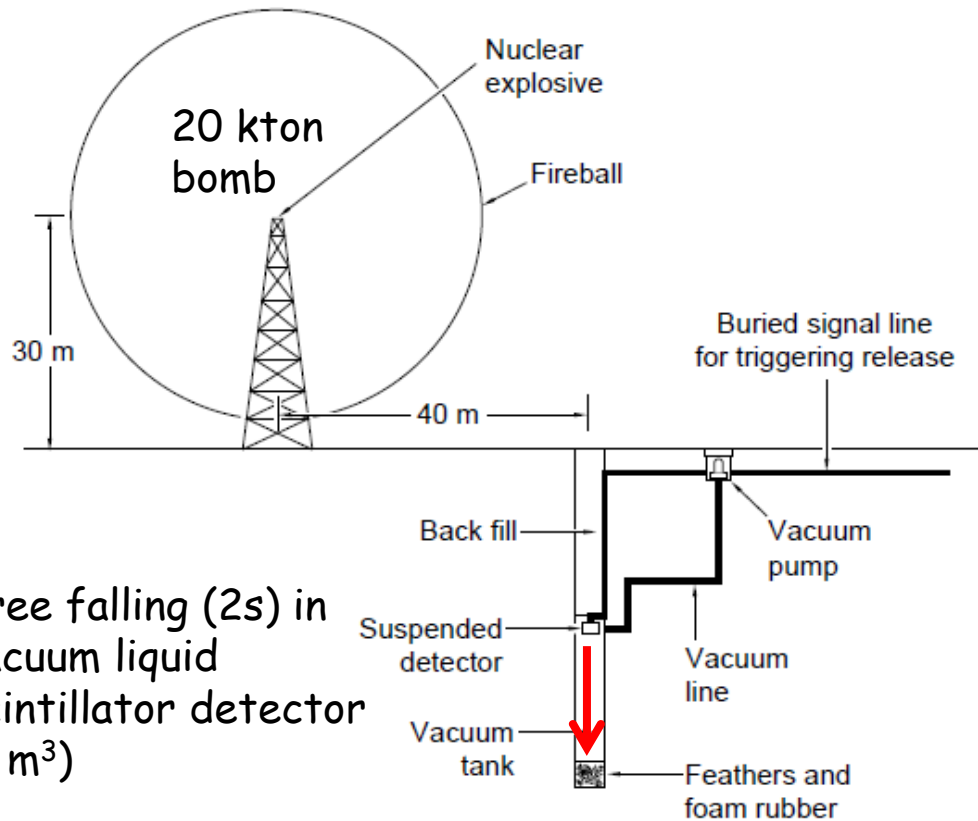
no signal => lepton number

- 1968, solar neutrinos detection



Neutrinos from decays of nuclear fission
Products →

How to detect neutrinos: producing them in a nuclear explosion



Free falling (2s) in vacuum liquid scintillator detector (1 m³)

« El Monstro »

Reines and Cowan 1951-1952

Approved after discussing with Fermi and Bethe who were convinced that this was the most promising (anti)neutrino source

- ✓ Intense
- ✓ Short flash (less environmental background)

but then abandoned in favour of the detection at a nuclear reactor:

Bomb: flux $\sim 10^4$ times larger than with a reactor

Background from neutrons and gammas similar to reactor

→ But a new idea on how to reduce the background and detect neutrinos over a long time scale with the low reactor flux

Figure 1. Detecting Neutrinos from a Nuclear Explosion

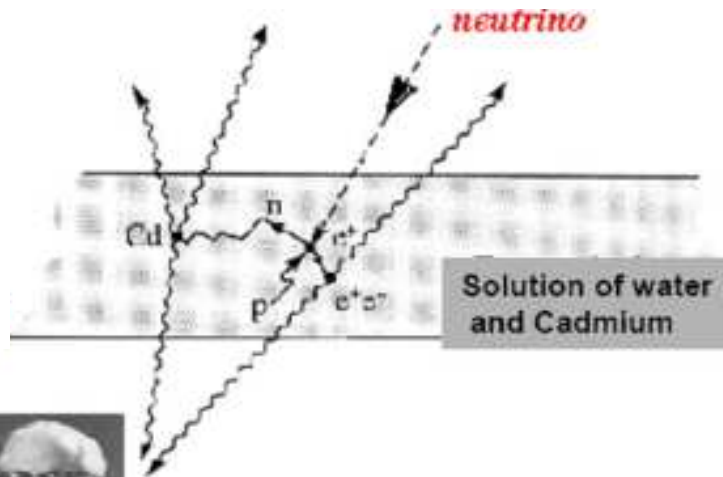
Antineutrinos from the fireball of a nuclear device would impinge on a liquid scintillation detector suspended in the hole dug below ground at a distance of about 40 meters from the 30-meter-high tower. In the original scheme of Reines and Cowan, the antineutrinos would induce inverse beta decay, and the detector would record the positrons produced in that process. This figure was redrawn courtesy of Smithsonian Institution.

Smithsonian Institution.

1956 (anti)neutrino detection at the Savannah River reactor

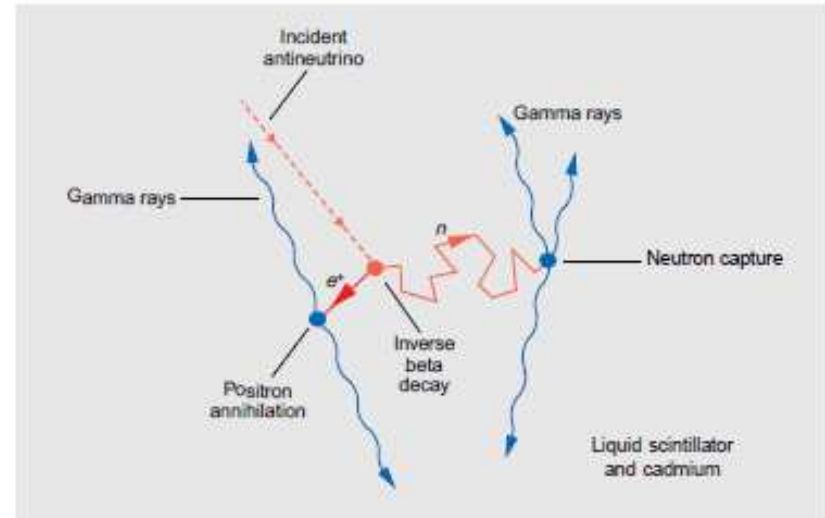
flux $\sim 10^{13}$ neutrino / (cm² s)

the idea: detect also the delayed neutron capture signal after the positron \rightarrow

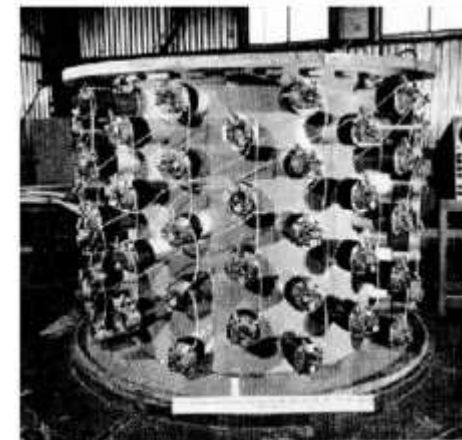


Reines: 1995 Nobel Prize

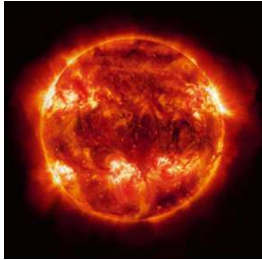
We are happy to inform you (Pauli) that we have definitely detected the neutrino



Detector 12 m underground and 11 m from reactor
 ~ 3 neutrinos detected/hour



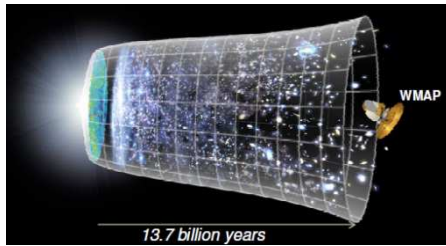
Neutrino sources:



Sun:
65 billions/s/cm²
on the earth surface
~ MeV



Nuclear reactors:
1 GW → 2^{E20} anti-nue/s
~ few MeV



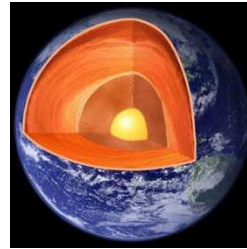
Big Bang
Relic neutrinos
330/cm³
1.95 K



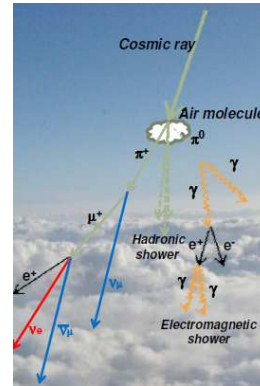
Particle accelerators
~few GeV



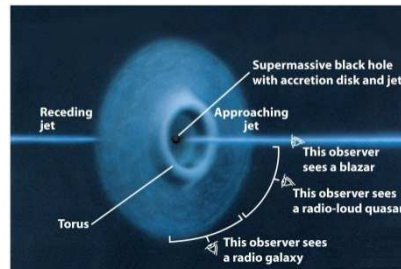
Supernova
explosion
99% of collapse
energy in neutrinos
10-30 MeV



Earth radioactivity
U, Th, K
→ Geoneutrinos
4^{E6} / (cm² s)
~ MeV



Cosmic rays
~ GeV
~ 1 / (cm² minute)



Extragalactic:
Active galactic nuclei
Gamma ray bursts
PeV

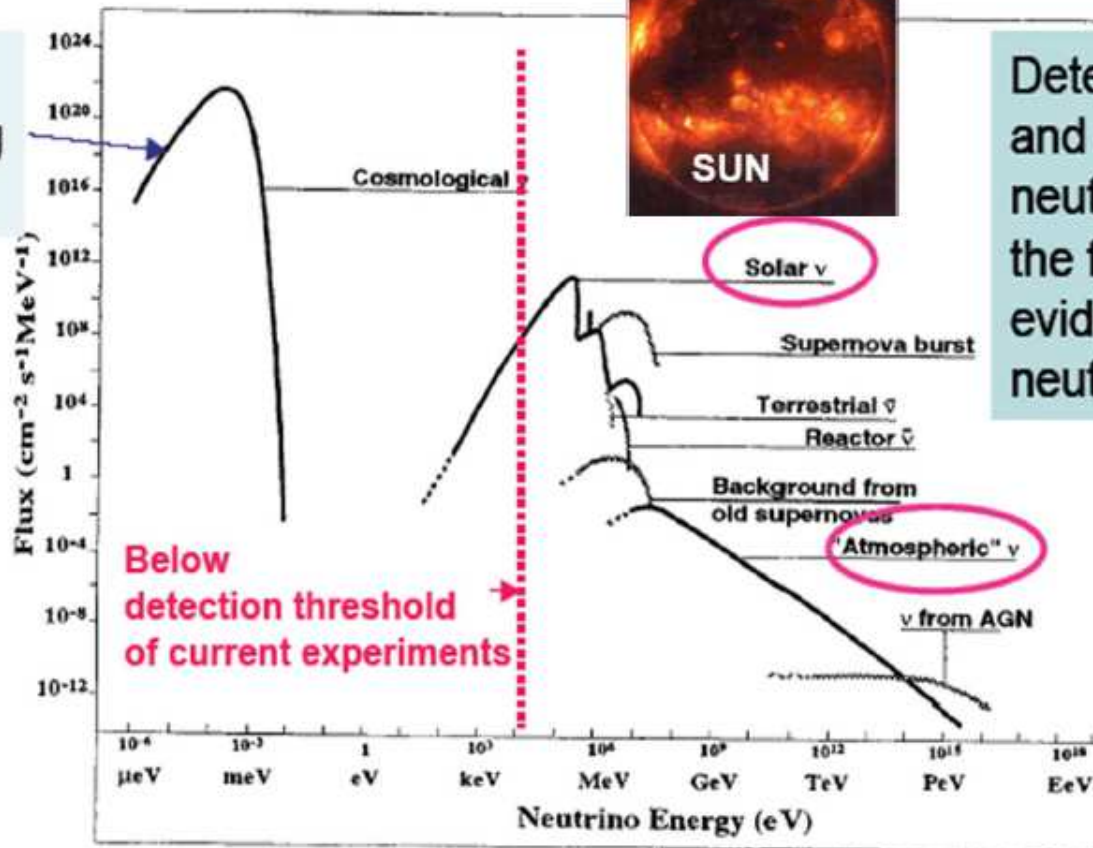
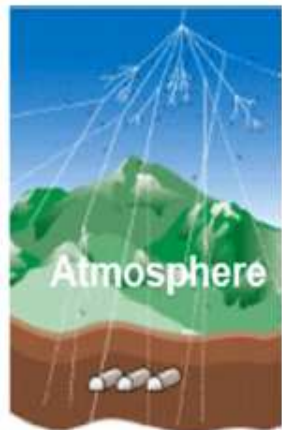


Human Body
20 mg of K 40
340 millions/day

Sources of neutrinos

The Sun is the most intense detected source with a flux on Earth of $6 \times 10^{10} \nu/\text{cm}^2\text{s}$

Abundant but challenging detection

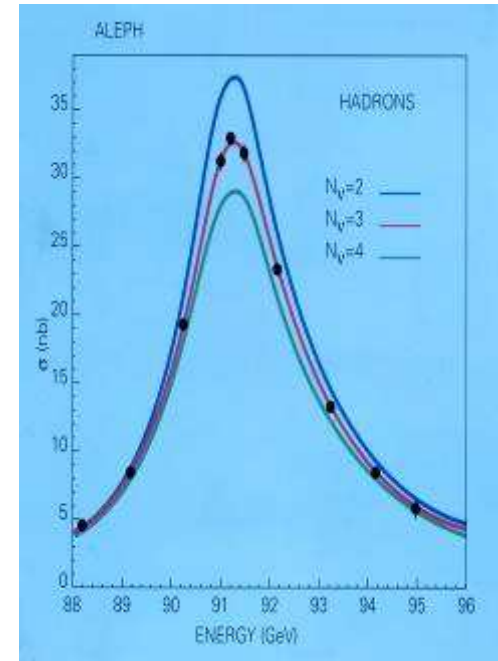


Detection of solar and atmospheric neutrino has provided the first compelling evidence of neutrino oscillations



Standard Model of Elementary particles

		Model of Elementary Particles				
		Three Generations of Matter (Fermions)			Force Carriers (Gauge Bosons)	
		I	II	III	Range	
Q u a r k s	Up	+2/3	Charm	+2/3	Photon	0
	stable	3	variable	3	stable	Electro-magnetism
	u	1.5 - 4 Mev	c	1.15 - 1.35 Gev	γ	Infinite
	Down	-1/3	Strange	-1/3	Gluon	0
	variable	3	variable	3	stable	Strong Interactions
	d	4 - 8 Mev	s	80 - 130 Mev	g	10^{-13} cm
L e p t o n s	Electron	0	Muon	0	Z zero	0
	Neutrino	stable	Neutrino	stable	10^{-25} s	Weak Interactions
	ν_e	< 3 ev	ν_μ	< 0.19 Mev	Z⁰	91.19 Gev
	Electron	-1	Muon	-1	W plus	±1
	stable	2×10^{-6} s	2x10 ⁻⁶ s	3x10 ⁻¹³ s	minus	10^{-25} s
	e	0.511 Mev	μ	105.6 Mev	W[±]	80.4 Gev
	Tau	-1	Tau	-1		
	τ	1.777 Gev				



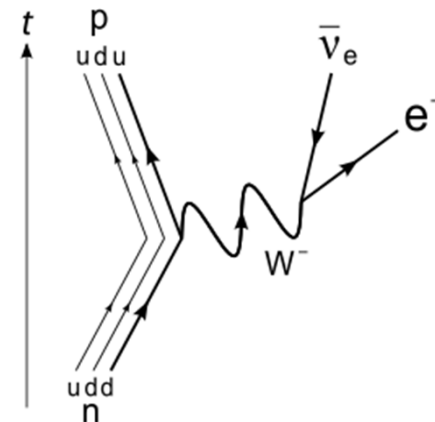
1989 LEP results:
only 3 neutrinos coupled to the Z⁰
(M_ν < M_Z/2)

Why 3 families ?
Why so different masses ?

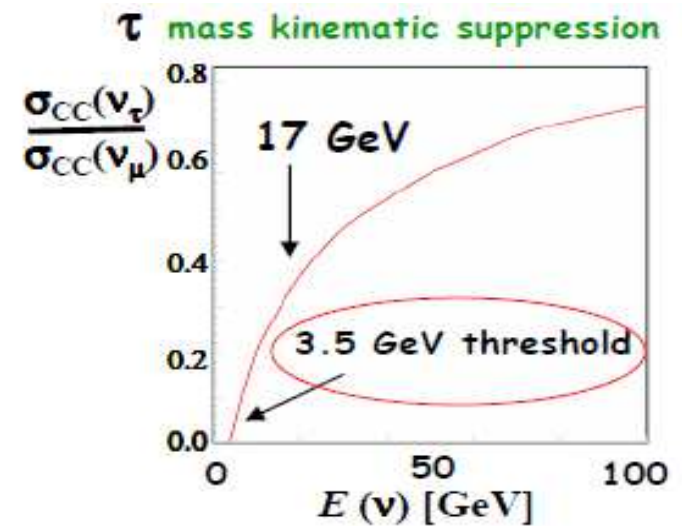
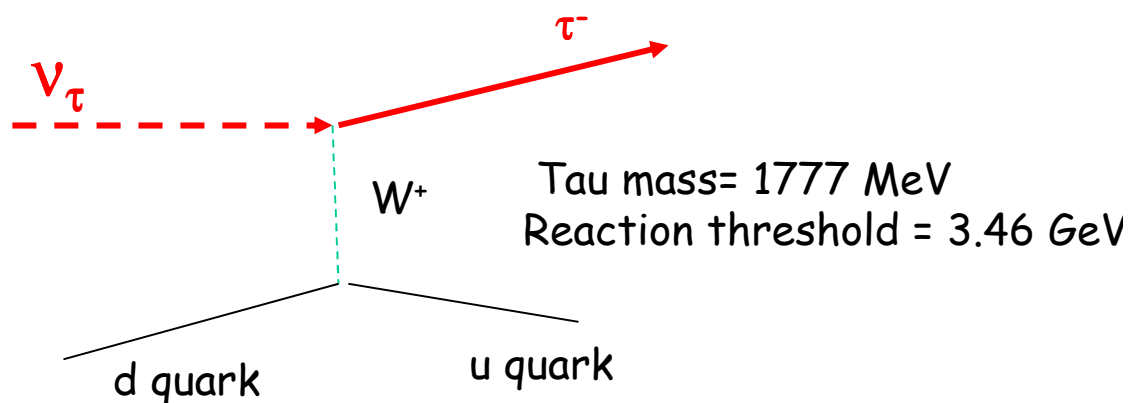
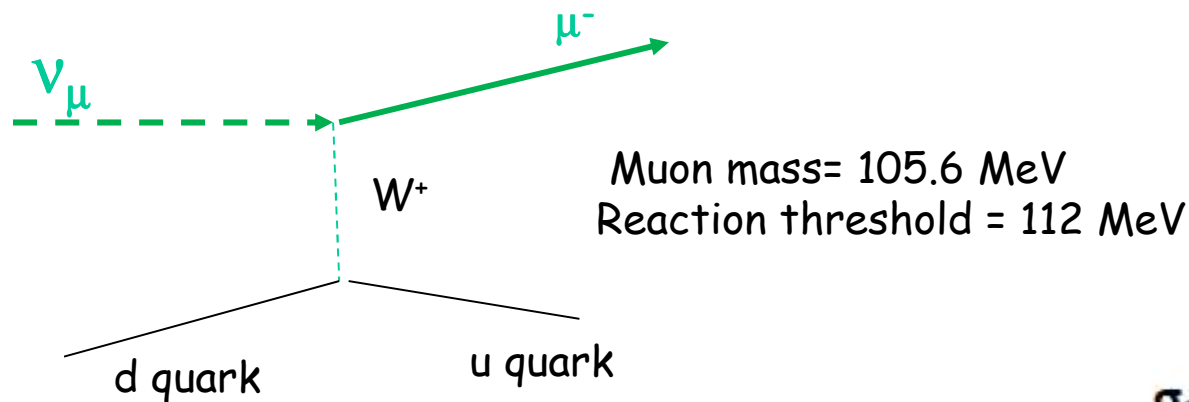
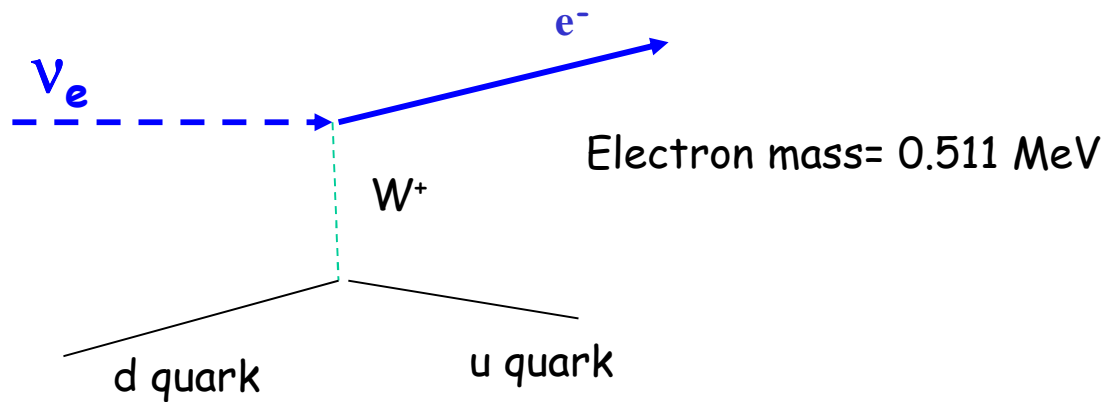
In SM massless neutrinos:
Neutrino: helicity -1 (+1 not existing)
Antineutrino: +1 (-1 not existing)



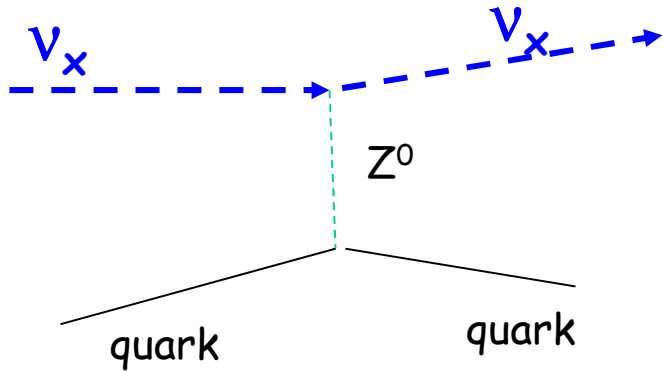
Modern description of the Beta decay



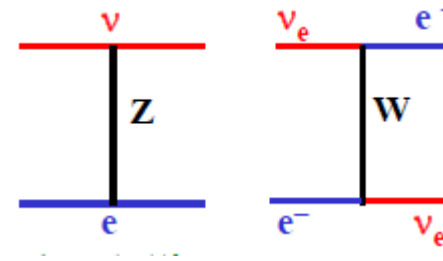
How can we detect different neutrinos: Charged Current reactions



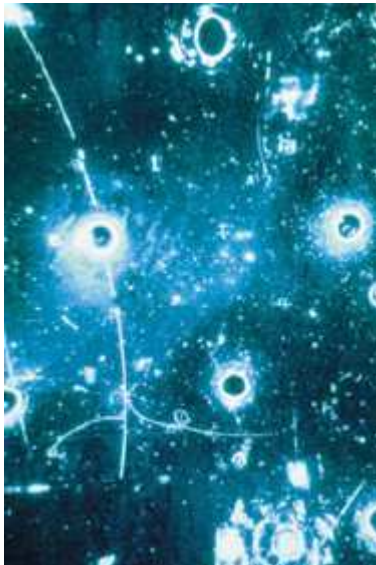
Neutral current reactions (Z exchange), do not distinguish neutrinos, no threshold



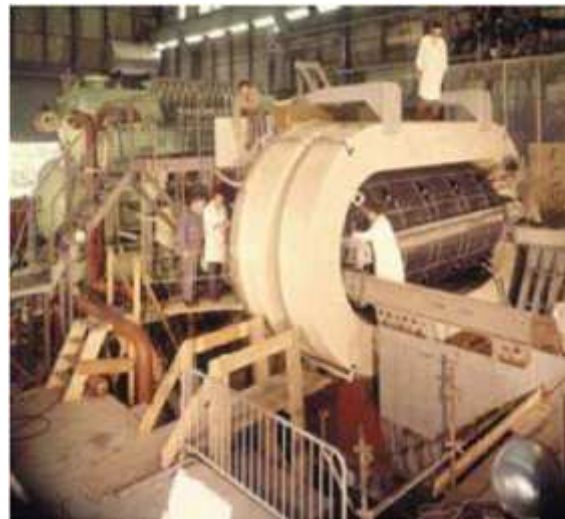
Elastic scattering neutrino-electron



Discovery of neutral currents 1973 (10 years before the discovery of the Z)



Bubble chamber experiment Gargamelle



1962 Discovery of the muonic neutrino with the first neutrino beam produced with an accelerator (pion decays)

Nobel 1988

Lederman, Schwarz, Steinberger



1959 **Pontecorvo** raised the question whether ν from β -decay processes is identical with ν from pion decay (Sov. Phys. JETP 10 (1960) 1236)

1960 **Pontecorvo and Schwartz** (PRL 4 (1960) 306) suggested to study neutrino reactions with high energy muons coming from proton accelerator ($\pi \rightarrow \mu + \nu_\mu$ $K \rightarrow \mu + \nu_\mu$)

$$\nu_\mu + N \rightarrow \mu^- + X$$

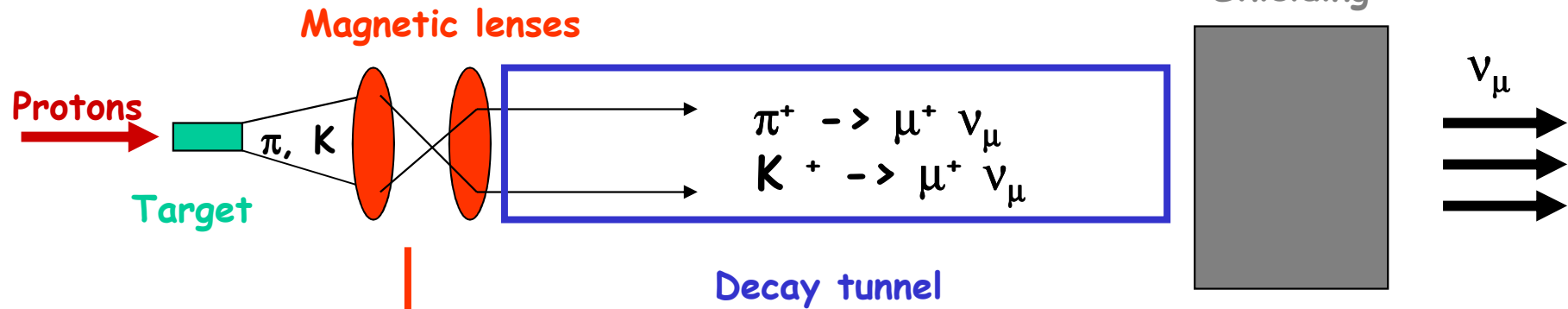


The two neutrinos experiment:

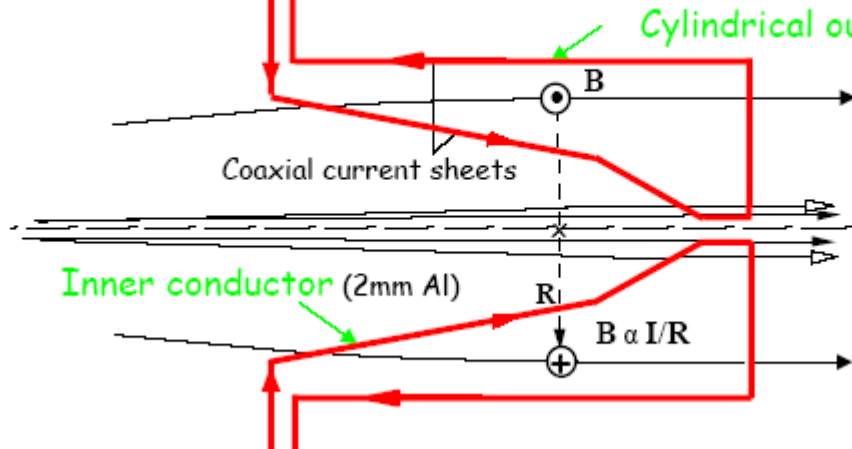
Muonic neutrino different than electronic neutrino

→ Conservation of leptonic number

Typical high energy Wide Band neutrino beam



$I = O(100 \text{ KA})$



Horns: sign selection,
focalization: flux $\times 10$

Contaminations:

$\bar{\nu}_\mu$ (wrong sign parents)	$O(5\%)$
ν_e (K_{e3} decays, μ decays)	$O(1\%)$
ν_τ (D_s decays)	$O(10^{-6})$

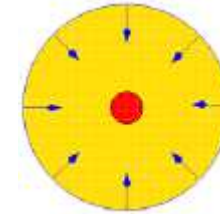


Note that the π/K abundances and spectra at the target are not easy to predict: to reduce systematics perform ad hoc hadron-production experiments (Spy, Harp etc ...)

Solar Neutrinos

Sun birth for the gravitational collapse of a primordial gas cloud ($\sim 75\% \text{H}_2$, $\sim 25\% \text{He}$)
→ Increase of density and temperature in the core → nuclear fusion reactions

→ Hydrostatic equilibrium between pressure from fusion reactions and gravitational attraction



Final result of chain of fusion reactions: $4\text{p} \rightarrow \text{He}^4 + 2\text{e}^+ + 2\nu_e$

Average energy emitted under the form of **electromagnetic radiation**:

$$Q = (4M_p - M_{\text{He}^4} + 2m_e)c^2 - \langle E(2\nu_e) \rangle \approx 26.1 \text{ MeV}$$

↳ (from $2\text{e}^+ + 2\text{e}^- \rightarrow 4\gamma$)

$$\langle E(2\nu_e) \rangle \approx 0.59 \text{ MeV}$$

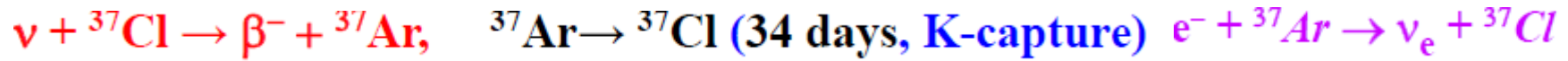
2.2% of the total

Solar luminosity: $L_\odot = 3.846 \times 10^{26} \text{ W} = 2.401 \times 10^{39} \text{ MeV/s}$

Rate of neutrino emission: $dN(\nu_e)/dt = 2 L_\odot / Q \approx 1.84 \times 10^{38} \text{ s}^{-1}$

Flux of neutrinos on earth: $\Phi(\nu_e) \approx 6.4 \times 10^{10} \text{ cm}^{-2} \text{ s}^{-1}$

First detection of solar neutrinos 1968: Homestake mine experiment R. Davis
 Depth equivalent to 4100 m of water



$E(\text{neutrino}) > 0.814 \text{ MeV}$

Tank with 390 m^3 of C_2Cl_4
 ${}^{37}\text{Cl} \sim 24\%$ of natural Cl

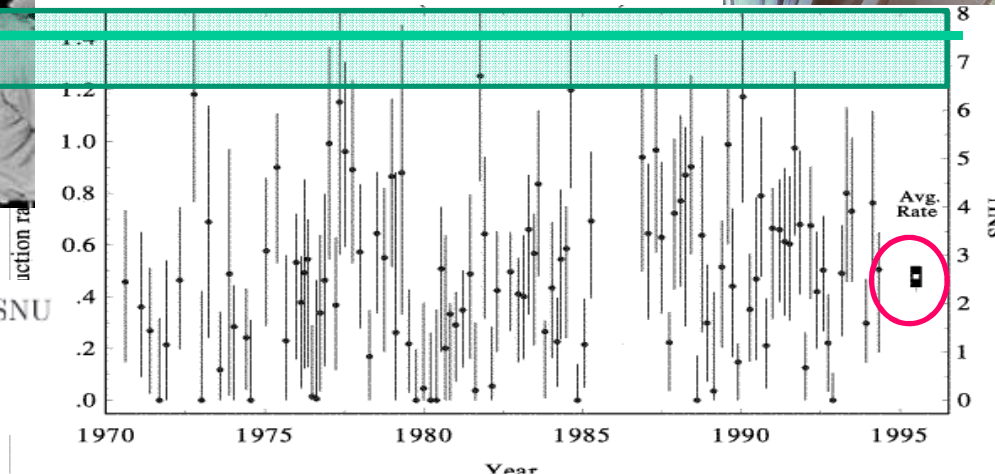
$\sim 1.5 \text{ Ar atoms/day}$ produced by solar neutrinos
 Extracted every 3 months with a flux of N

Final state ${}^{37}\text{Cl}$ excited emitting Augier electrons e/o x rays

Results compared to the neutrino flux predicted by the Standard Solar Model (J. Bahcall)



$\rightarrow 1/3$ of expected rate
 Solar neutrinos deficit



$$R({}^{37}\text{Cl}) = 2.56 \pm 0.16 \pm 0.16 \text{ SNU}$$

$$R_{\text{SSM}} = 7.6_{-1.1}^{+1.3} \text{ SNU}$$

$$R_{\text{Données/SSM}} = 0.33 \pm 0.03$$

Interpretations:

I [J.N. Bahcall] want to tell you an illustrative story about neutrino research ... One of the miners came over to our bench, said : "Hello, Dr. Davis. How is it going ? You don't look too happy." And, Ray replied : "Well, I don't know ...I am capturing in my tank many fewer of those neutrinos than this young man says I should be capturing." The miner [...] finally said : "Never mind, Dr. Davis, it has been a very cloudy summer here in South Dakota. "

More seriously debated for long ... long time:

The trivial ones:

- The Homestake experiment which is quite delicate has a bias in the neutrino detection
- The Standard Solar Model is not correct



The fascinating one:

Pontecorvo: the Davis experiment and the SSM are both correct it is new physics: neutrinos change their nature during their trip to the earth

→ Neutrino oscillations

Electronic neutrinos from the sun become muonic neutrinos

The energy of the muonic neutrinos is too low for the charged current reaction → neutrino disappearance

But neutrinos must be massive particles

ОБЪЕДИНЕННЫЙ ИНСТИТУТ ЯДЕРНЫХ ИССЛЕДОВАНИЙ

JOINT INSTITUTE FOR NUCLEAR RESEARCH

Москва, Главный почтамт П/Я 7В.

Head Post Office, P. O. Box 77, Moscow, USSR

№ 994/31

April 6, 19 72

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Prof. J.N.Bahcall

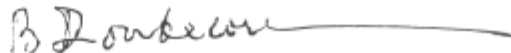
The Institute for Advanced Study
School of Natural Science
Princeton, New Jersey 08540, USA

Dear Prof. Bahcall,


Thank you very much for your letter and the abstract of the new Davis investigation the numerical results of which I did not know. It starts to be really interesting! It would be nice if all this will end with something unexpected from the point of view of particle physics. Unfortunately, it will not be easy to demonstrate this, even if nature works that way.

I will attend the Balaton meeting on neutrinos and looking forward to see you there.

Yours sincerely,



B. Pontecorvo



Pontecorvo was
predictive:
It took 30 years for the
demonstration!

Neutrino oscillations

Neutrino mixing (Pontecorvo 1958; Maki, Nakagawa, Sakata 1962):

3 neutrinos framework, neutrinos are massive particles and they mix similarly to quarks; the flavour eigenstates ν_e, ν_μ, ν_τ are not mass eigenstates but linear superpositions of the mass eigenstates ν_1, ν_2, ν_3 with eigenvalues m_1, m_2, m_3 :

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i} |\nu_i\rangle$$

$\alpha = e, \mu, \tau$ (flavor index)

$i = 1, 2, 3$ (mass index)

$U_{\alpha i}$ = unitary mixing matrix

Today favorite parametrization of U : in terms of 3 mixing angles $\theta_{12}, \theta_{23}, \theta_{13}$ and one Dirac-like CP phase δ (two extra phases in case of Majorana neutrinos):

$$U \equiv U_{23}U_{13}U_{12} \equiv \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Atmospheric ν oscillations

Solar ν oscillations

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta} & c_{13}c_{23} \end{pmatrix}$$

where: $s_{ij} \equiv \sin \theta_{ij}$, $c_{ij} \equiv \cos \theta_{ij}$.

Considering the **time evolution** of a flavour eigenstate ν_α produced at $t = 0$:

$$|\nu(t)\rangle = e^{i\mathbf{p}\cdot\mathbf{r}} \sum_k U_{\alpha k} e^{-iE_k t} |\nu_k\rangle \quad E_k = \sqrt{p^2 + m_k^2}$$

The phases: $e^{-iE_k t}$ will be different if $m_j \neq m_k$

Projecting $\nu(t)$ on the flavor basis one can obtain the probability of finding other flavours:

→ Appearance of the flavour $\nu_\beta \neq \nu_\alpha$ for $t > 0$

Simplified case: two neutrinos mixing

Only one mixing angle θ is needed

$$|\nu_\alpha\rangle = \cos\theta |\nu_1\rangle + \sin\theta |\nu_2\rangle$$

$$|\nu_\beta\rangle = -\sin\theta |\nu_1\rangle + \cos\theta |\nu_2\rangle$$

If $\nu = \nu_\alpha$ at ($t = 0$):

$$|\nu(t)\rangle = e^{i(\mathbf{p}\cdot\mathbf{r} - E_1 t)} \left[\cos\theta |\nu_1\rangle + e^{-i(E_2 - E_1)t} \sin\theta |\nu_2\rangle \right]$$

Probability of detecting ν_β at the instant t if $\nu(0) = \nu_\alpha$:

$$\mathcal{P}_{\alpha\beta}(t) = \left| \langle \nu_\beta | \nu(t) \rangle \right|^2 = \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2 t}{4E}\right) \quad \hbar = c = 1$$

$$\Delta m^2 \equiv m_2^2 - m_1^2$$

Oscillatory behaviour of $\mathcal{P}_{\alpha\beta}$ with time ruled by two parameters:

θ is related to the amplitude of the oscillation
 Δm^2 is related to the wavelength

For $m \ll p$, and assuming propagation in vacuum:

$$E_2 - E_1 \approx \frac{m_2^2 - m_1^2}{2p} \approx \frac{m_2^2 - m_1^2}{2E} \equiv \frac{\Delta m^2}{2E}$$

$$E = \sqrt{p^2 + m^2} \approx p + \frac{m^2}{2p}$$

In more empirical units:

$$\mathcal{P}_{\alpha\beta}(L) = \sin^2(2\theta) \sin^2\left(1.267 \Delta m^2 \frac{L}{E}\right)$$

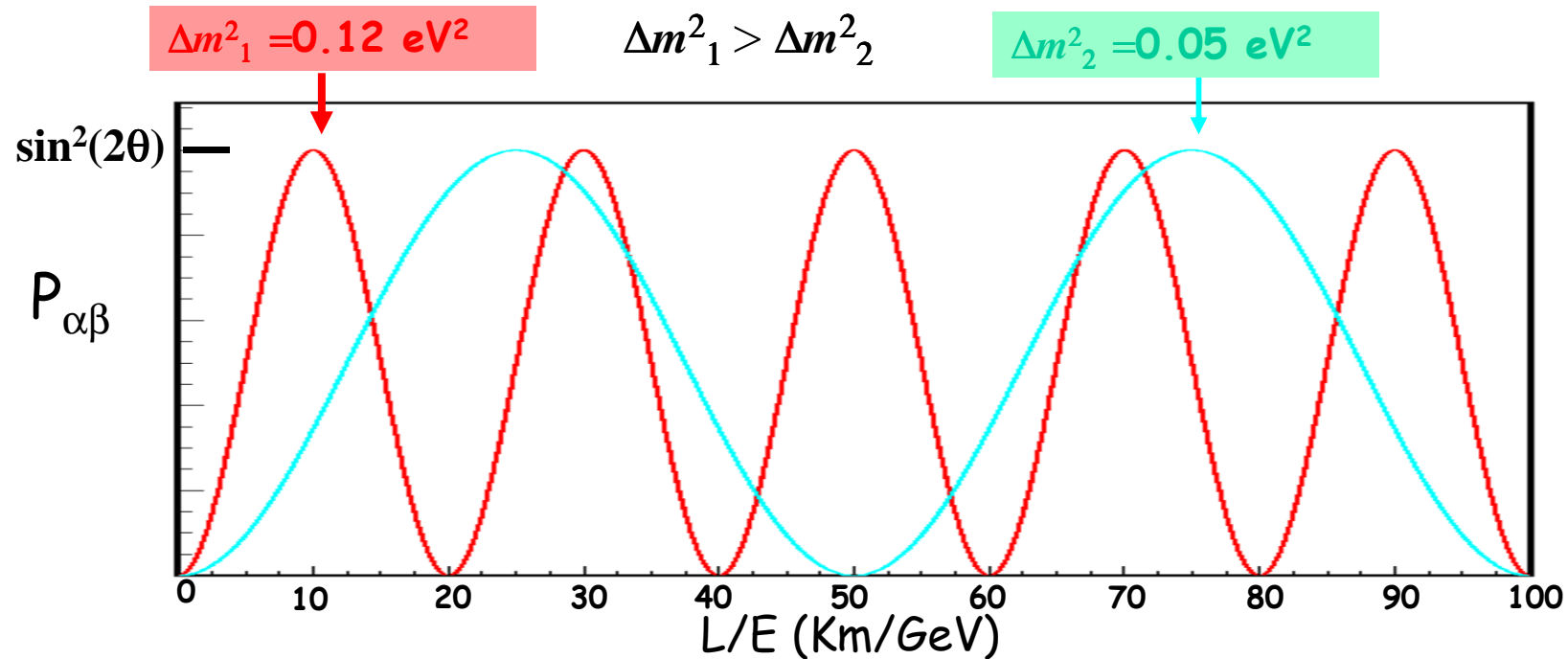
Δm^2 [eV²]

$L=ct$ [km] (distance among the neutrino source and the detector)

E [GeV] (neutrino energy)

Given Δm^2 the experimental quantity for the study of neutrino oscillations is the ratio L/E [km/GeV]: first oscillation maximum at $L/E \sim 1.24 / \Delta m^2$

$$\mathcal{P}_{\alpha\beta}(L) = \sin^2(2\theta) \sin^2\left(\pi \frac{L}{\lambda}\right) \quad \lambda = 2.48 \frac{E}{\Delta m^2} \quad \text{Oscillation wavelength}$$



For very large Δm^2 the oscillations become very fast and average over the dimensions of the source and of the detector:

→ $\langle \sin^2(1.267 \Delta m^2 L/E) \rangle = 1/2$ $P = (1/2) \sin^2(2\theta)$

The baseline is related to the L/E ratio of the experimental setup:

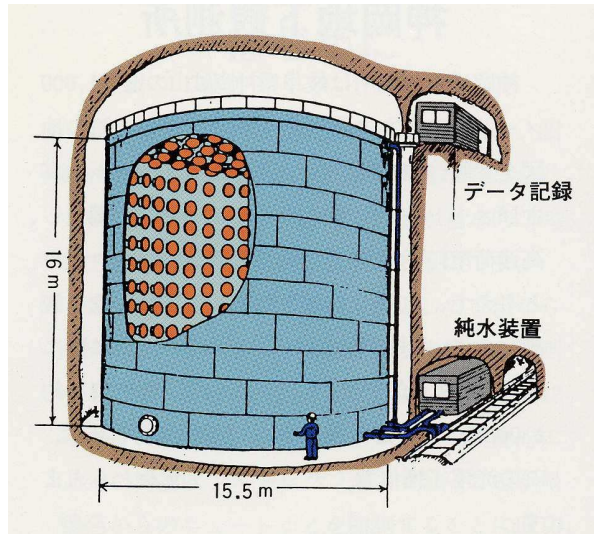
Short Baseline experiments: sensitive to large Δm^2 ($> 1 \text{ eV}^2$)

Long Baseline experiments: sensitive at least to Δm^2 of interest for the atmospheric neutrino anomaly ($< 10^{-2} \text{ eV}^2$), $L/E > 100 \text{ Km/GeV}$

Reactors: $L > 0.3 \text{ Km}$, $E \sim 3 \text{ MeV}$

Accelerators: $L > 100 \text{ Km}$, $E \sim 1 \text{ GeV}$

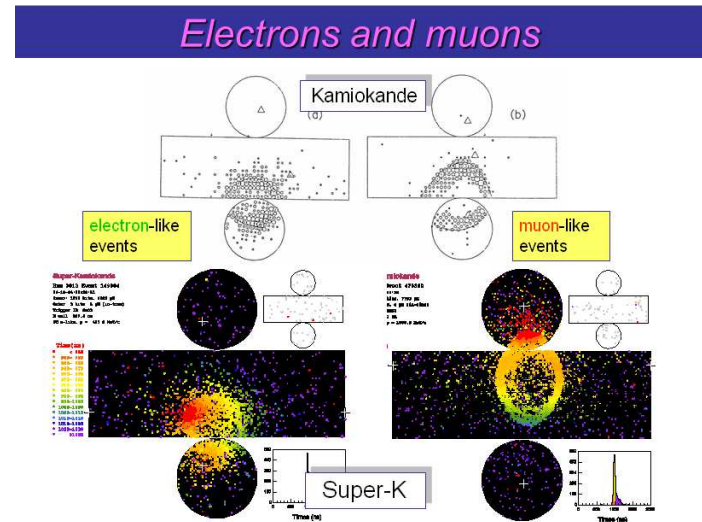
Water Cerenkov experiment (Kamiokande 1987-1994)



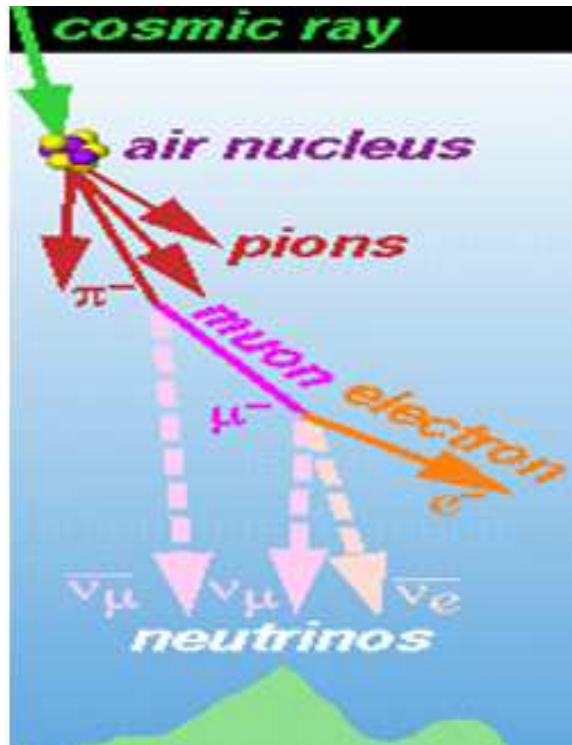
Particle detection by emission of Cerenkov light in water (680 tons) → (electrons, muons)

Built for proton decay search
 Neutrinos produced by cosmic rays in the atmosphere are a background for cosmic rays
 → Studying this background people realize that it is different than expectations

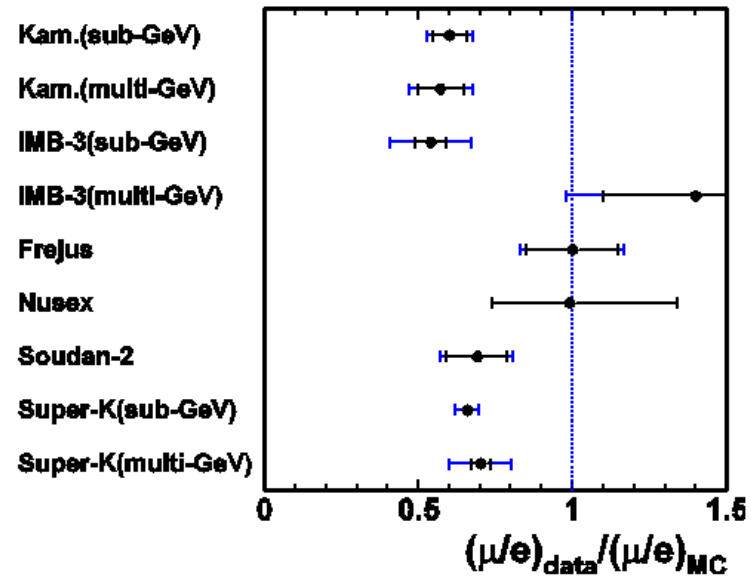
→ Can look at solar neutrinos (high threshold > 5 MeV) by elastic scattering on electrons (emitted electron at 5 MeV stops in ~ 2 cm in water)
 → Deficit of solar neutrinos $\sim 50\%$



Atmospheric neutrinos anomaly



Let's write the atmospheric ν_μ deficit by $(\mu/e)_{\text{data}}/(\mu/e)_{\text{MC}}$



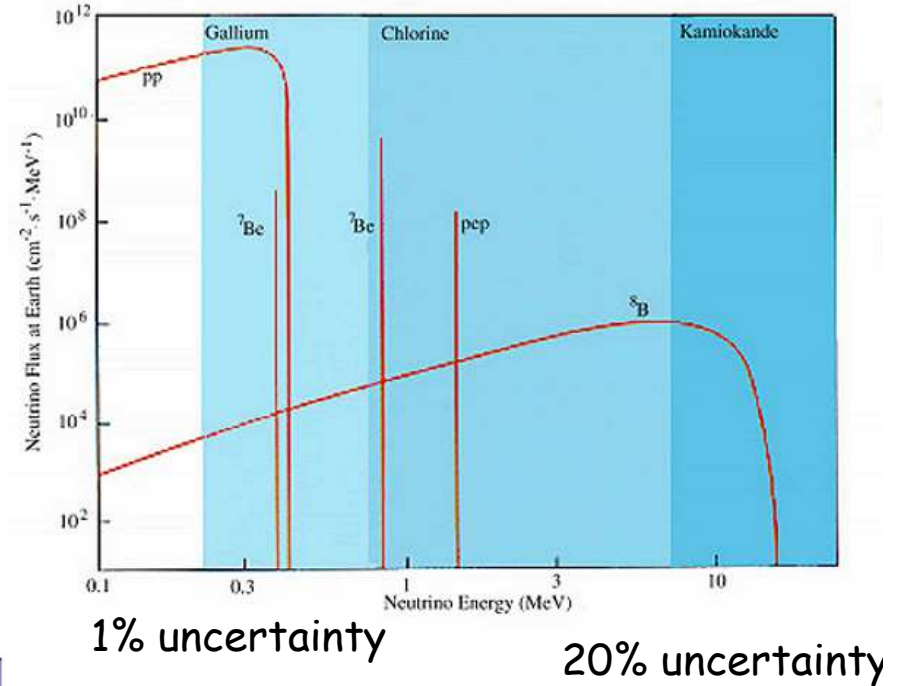
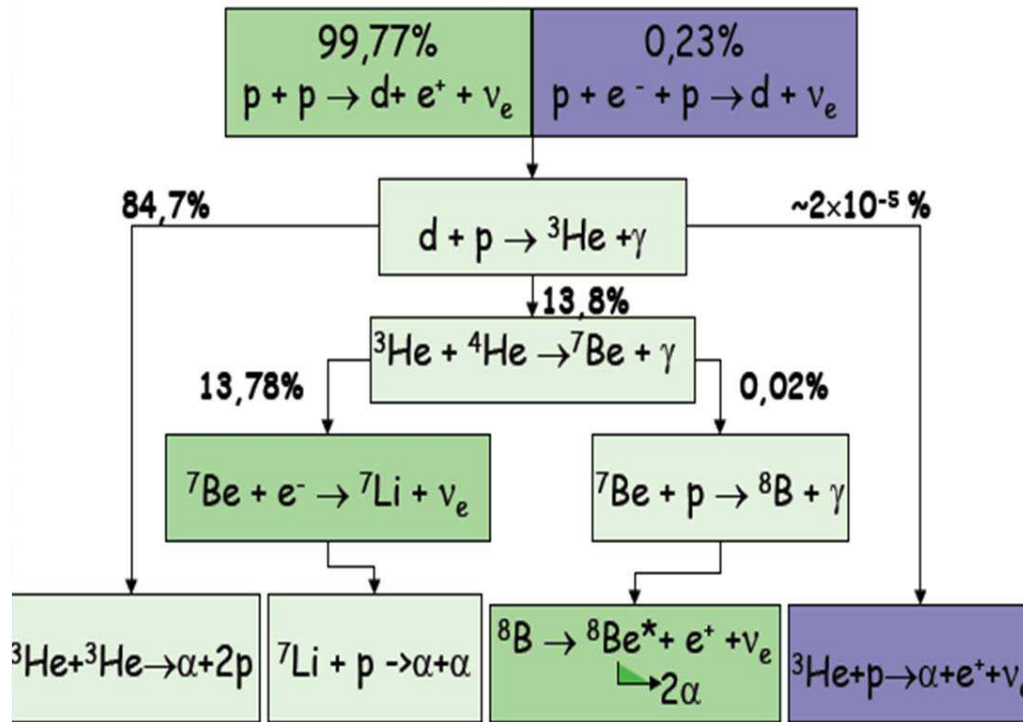
Unclear situation among different experiments

Interpretation in terms of neutrino oscillations (both $\nu_\mu \rightarrow \nu_e$ and $\nu_\mu \rightarrow \nu_\tau$) with

$$\Delta m^2 \sim 10^{-2} \text{ eV}^2$$

Some hints of dependence on the zenith angle

The pp-chain



Gallex (1991-2002): radiochemical experiment with Gallium looking at low energy neutrinos (>0.233 MeV) from pp cycle

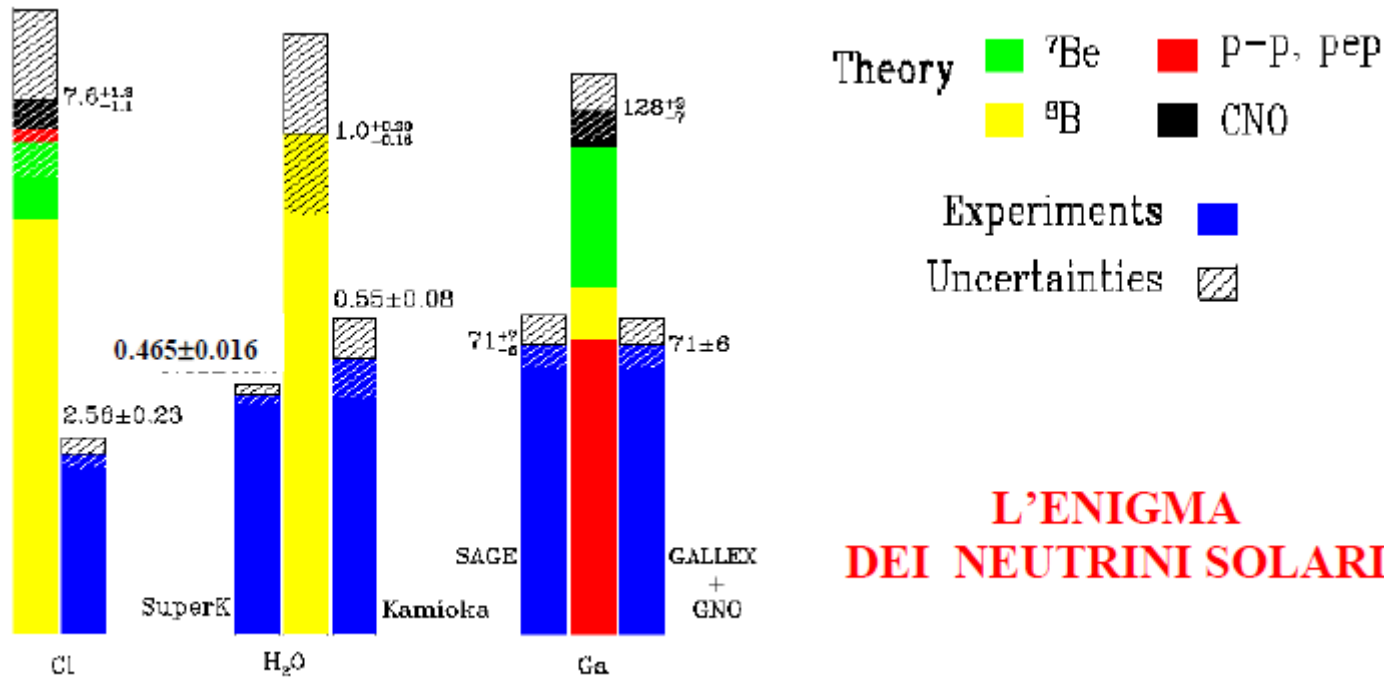
→ Confirms the deficit but:
Data/SSM = 0.56



In parallel many checks are performed also on the Standard Solar Model

The 3 experiments give different results (in particular Homestake) even considering neutrino oscillations as an explanation

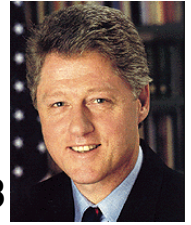
Is there an energy dependence of the solar neutrino deficit?



A more complex mechanism MSW which includes in the oscillations mechanism resonant effects of neutrino interactions with matter can explain the 3 results

Neutrino oscillation searches at the beginning of 90s

(long time ago in neutrino physics, not so much in everyday life ...)



U.S. new president in 1993

- The long standing (since 1968) problem of the solar neutrino deficit opened by the Homestake measurements (+ Kamiokande since 1986)
In 1992 first Gallex results confirm the deficit also for neutrinos from the pp cycle
- Atmospheric neutrino anomaly still quite weak

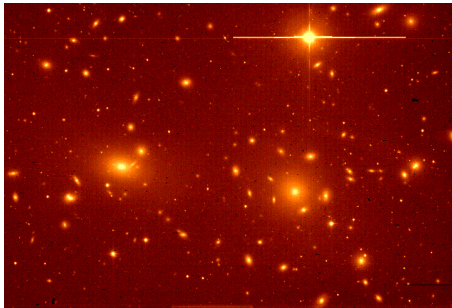
The controlled observation of neutrino oscillations with an accelerator neutrino beam would have been a great discovery, **where to search ?**

Prejudice towards **small mixing angles** and **large Δm^2**

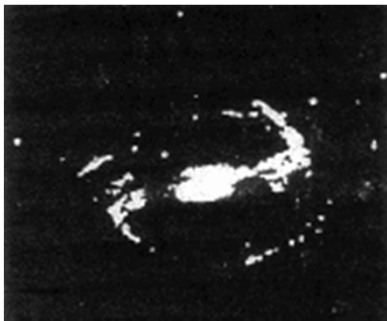
- ✓ Take the MSW solution of the solar neutrino deficit: $\Delta m_{\mu e}^2 \sim 10^{-5} \text{ eV}^2$
- ✓ Assume a strong hierarchy: $m_{\nu e} \ll m_{\nu \mu} \ll m_{\nu \tau} \rightarrow m_{\nu \mu} \sim 3 \times 10^{-3} \text{ eV}$
- ✓ Assume the See-Saw mechanism: $m(\nu_i) = m^2(f_i) / M$
 $M = \text{very large Majorana mass}$ $m(f_i) = \text{e.g. quark masses}$

Then: $m_{\nu \tau} \sim 30 \text{ eV}$ (Cosmological relevance)

Dark matter

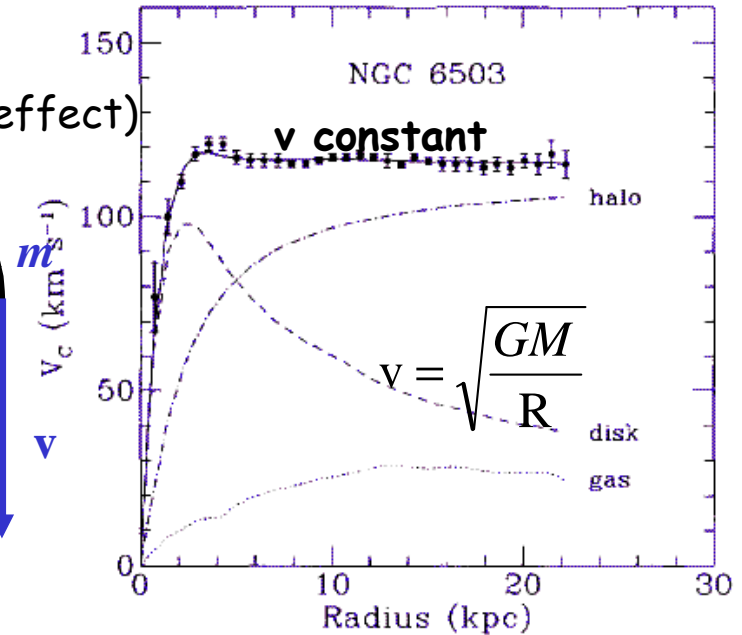
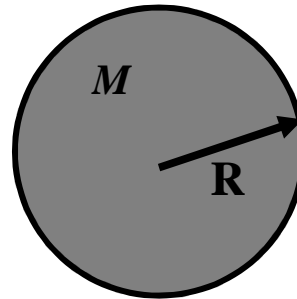


Coma cluster of galaxies, application of the virial theorem by F. Zwicky (1933)
 velocity dispersion, geometric size → **total mass (x400 luminous mass)**



Rotational velocity curves of galaxies (Hydrogen, doppler effect)

$$m \frac{v^2}{R} = \frac{GMm}{R^2}$$



Could it be due to the BIG BANG relic neutrinos ?
 112 ν / cm^3 per flavour

if $\sum m_i \approx 52 \text{ eV}$ then $\Omega_\nu = 1$

« ν are an important component of the dark matter » ~ a few 10 eV
 Harari PLB 1989

1992 first measurements from the COBE satellite $\Omega \sim 1$

J. Ellis PLB 292 1992 $\Omega_{\text{HDM}} = 0.3$, $\Omega_{\text{CDM}} = 0.7$

Recent cosmological results: $\sum m_\nu < 0.26 \text{ eV}$ 95% CL

With $m_{\nu_\tau} \sim 30 \text{ eV}$ cosmological neutrinos important component of dark matter $\Delta m^2_{\mu\tau} \sim O(100 \text{ eV}^2)$

Look for $\nu_\mu \rightarrow \nu_\tau$ with short baseline experiments at accelerators, high energy beam

CERN ν_τ appearance experiments:

Search for ν_τ appearance from oscillations in the CERN wide band neutrino beam (WANF)

Pioneers of the technique also for long baseline experiments, important samples of neutrino interactions well measured

NOMAD:

- Proposal 1991
- Detector 1995
- Data-taking 95-98 (1.35 M ν_μ CC)

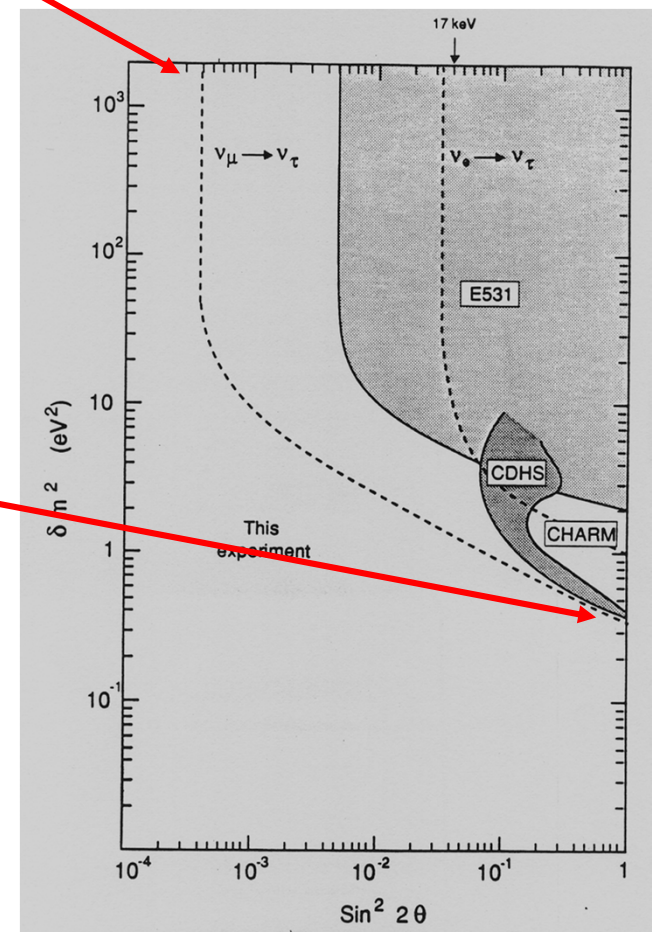
CHORUS:

Data-taking 1994-1997 (0.71 M ν_μ CC)

sensitive down to:

$P_{\mu\tau} \sim 1.5 \times 10^{-4}$ (90% CL) (x10) improvement

$\langle E_\nu \rangle = 24 \text{ GeV}$
 $\langle L \rangle = 600 \text{ m}$
sensitive to:
 $1 \text{ eV}^2 < \Delta m^2$

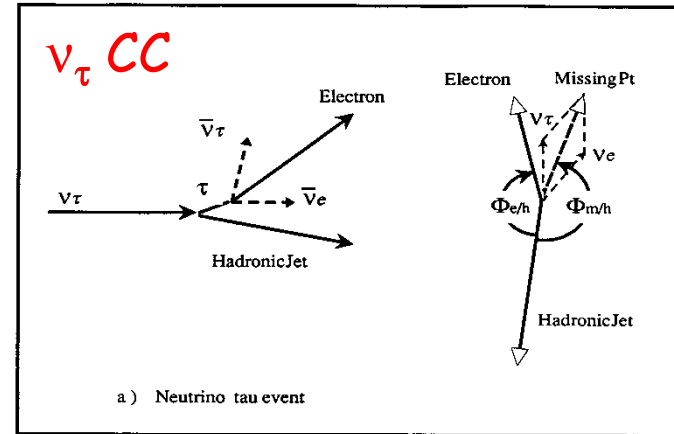
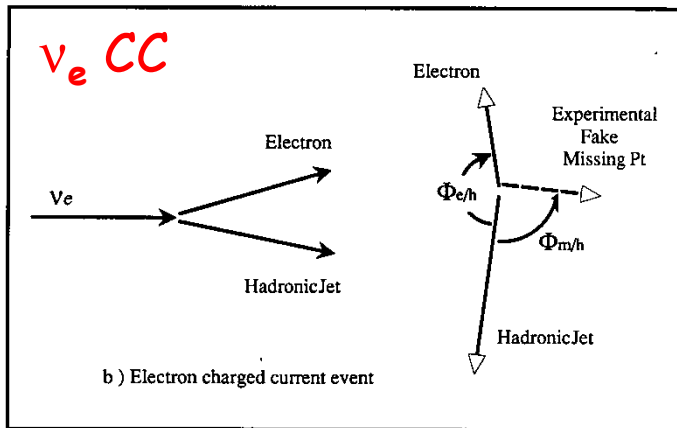


Use of kinematics to extract a ν_τ signal:

(First proposed by Albright and Shrock P.L.B. 1979)

NOMAD: fully reconstruct 1.7 M neutrino interactions, with good resolution, at single particles level:

→ Kinematics closure on the transverse plane



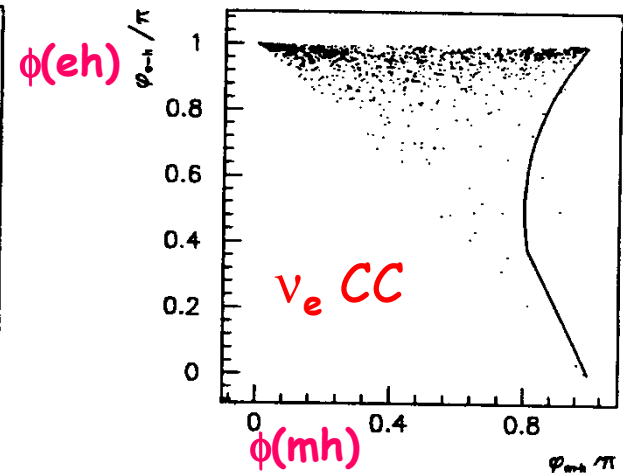
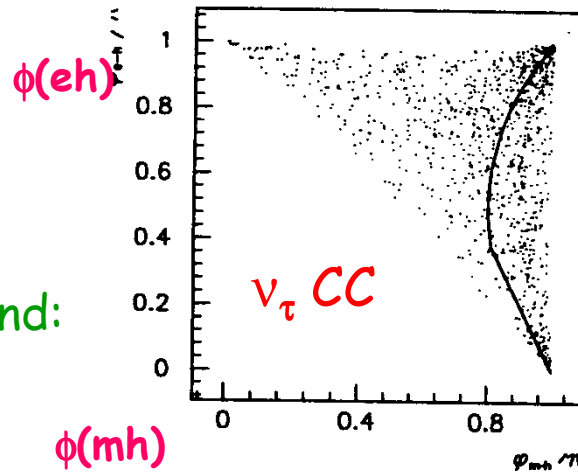
Find ν_τ down to $P_{\mu\tau} \sim 10^{-4}$ in a large background:

- 1.3 M ν_μ CC
- 0.4 M ν_μ NC
- 13 K ν_e CC

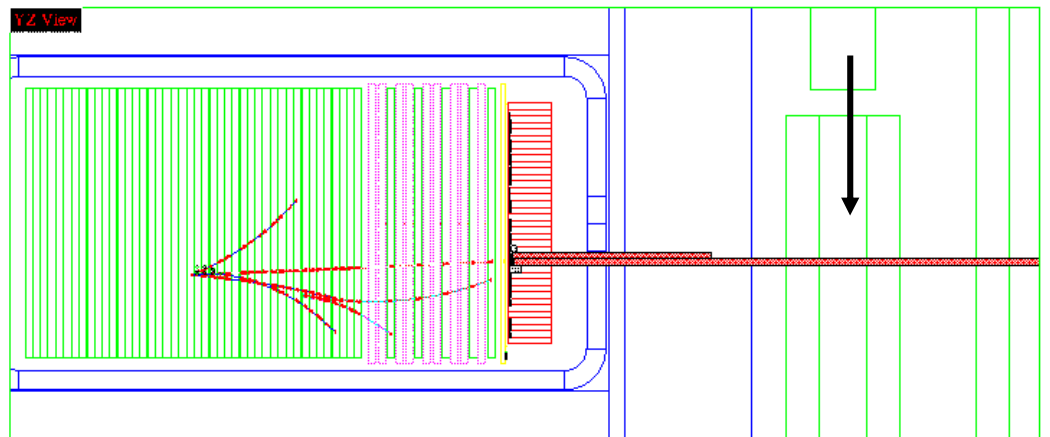
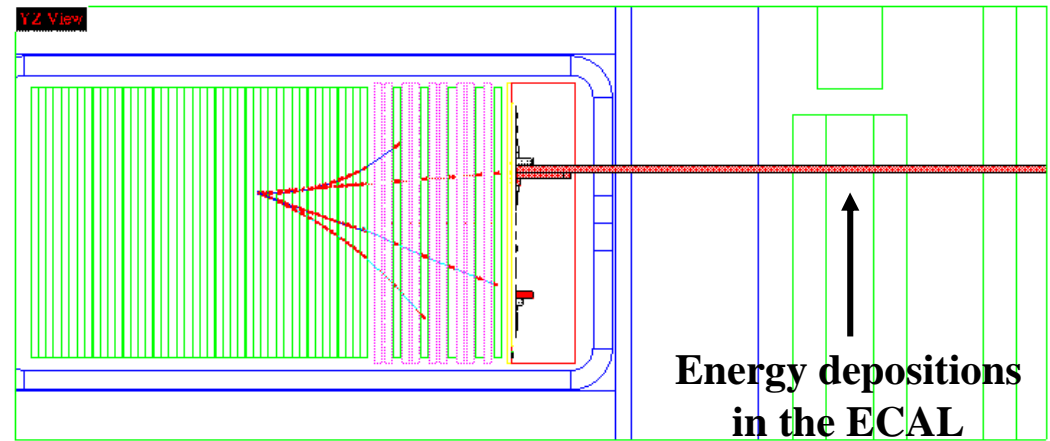
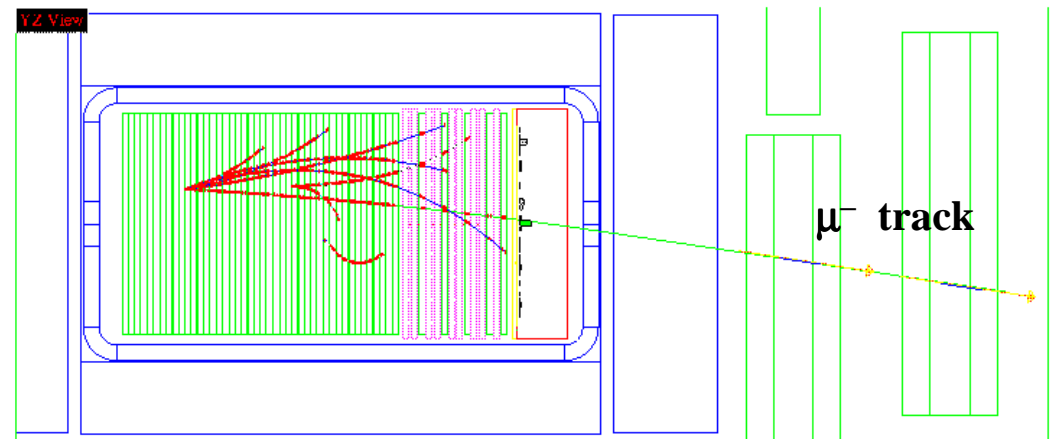
Exploit the small ν_e background:

$\tau \rightarrow e$ channel: electron id

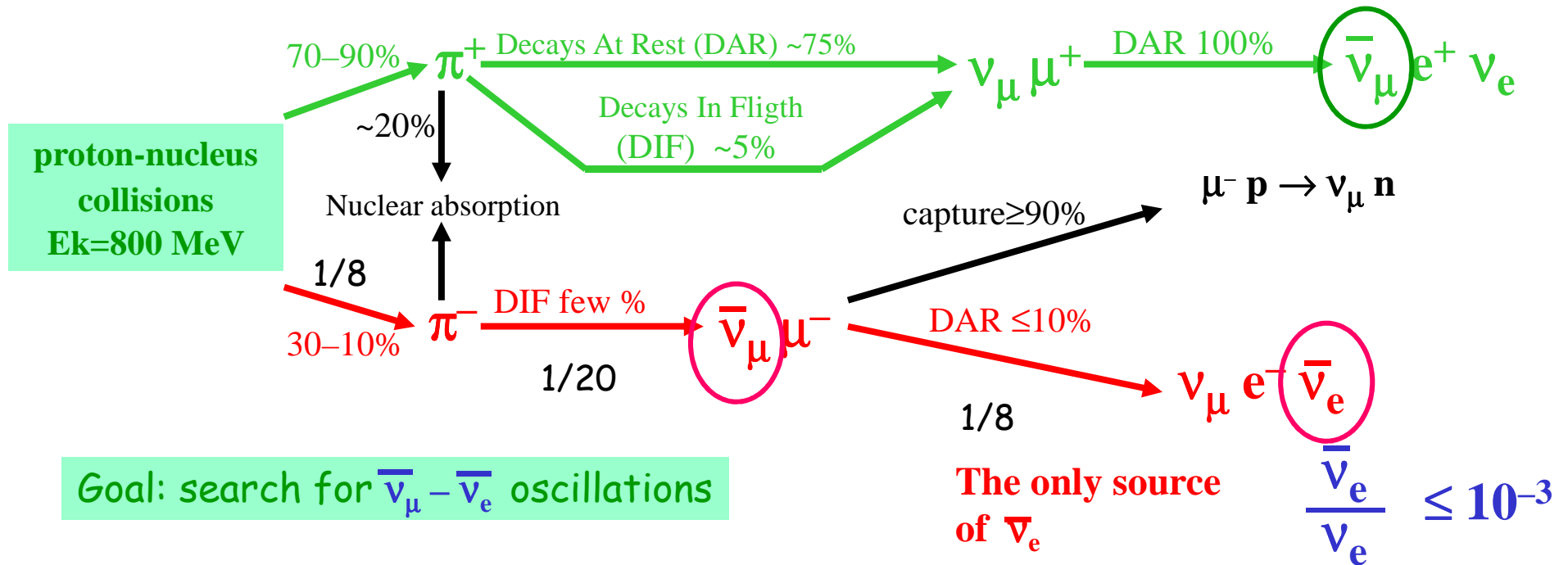
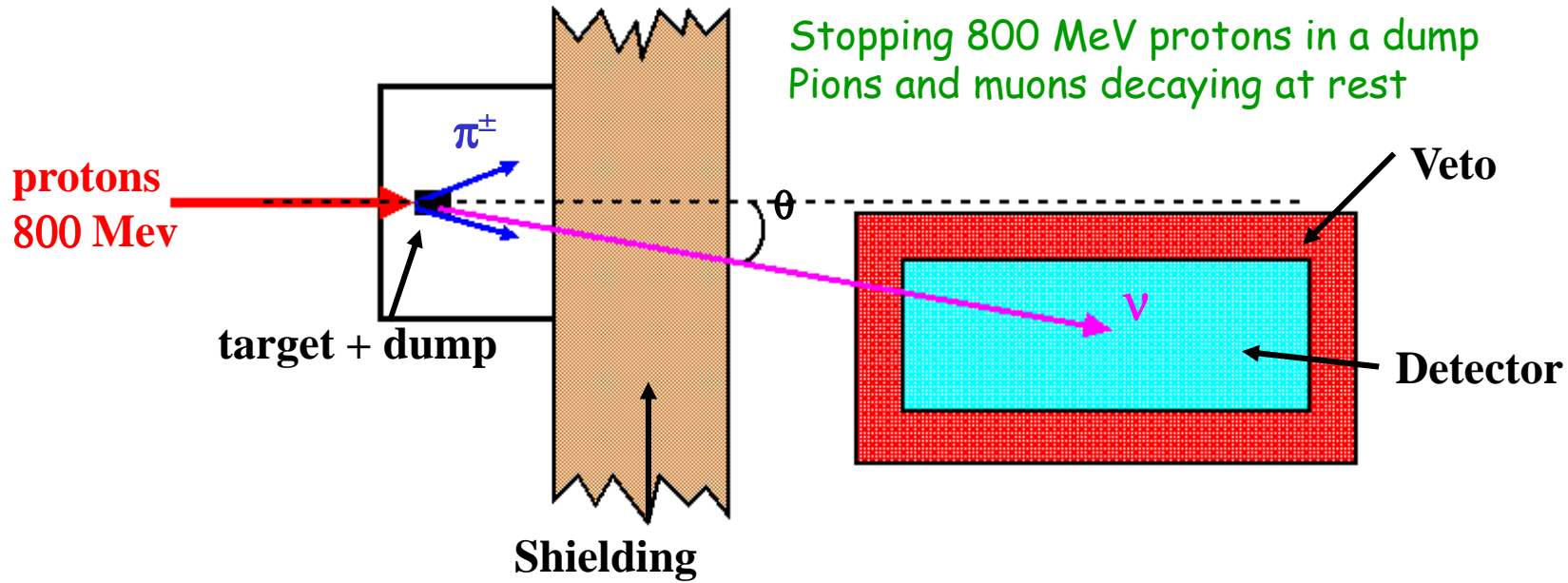
The $\phi - \phi$ plot:



Nomad typical events:



The LSND experiment (1993-2001)



LSND result: evidence for $\bar{\nu}_\mu - \bar{\nu}_e$ oscillations

Signal: Positrons with $20 < E < 200$ MeV correlated in space and in time with the γ rays of 2.2 MeV expected from the neutron capture:

$$N(\text{“beam-on”}) - N(\text{“beam-off”}) = 117.9 \pm 22.4 \text{ events}$$

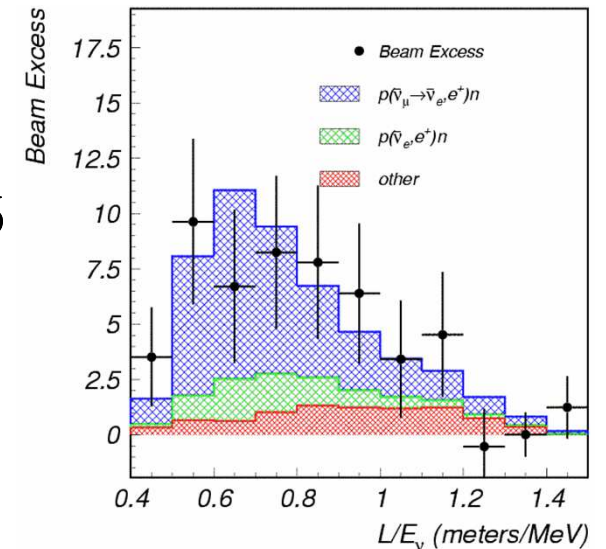
$$\text{Background due to } \mu^- \text{ DAR} = 19.5 \pm 3.9$$

$$\text{Background from } \pi \text{ DIF} + (\bar{\nu}_\mu + p \rightarrow \mu^+ + n) = 10.5 \pm 4.6$$

$$\text{Signal } \bar{\nu}_e = 87.9 \pm 22.4 \pm 6.0 \text{ events} \quad 3.8 \sigma \text{ effect}$$

(stat.) (syst.)

$$\mathcal{P}_{\text{osc}}(\bar{\nu}_\mu - \bar{\nu}_e) = (0.264 \pm 0.067 \pm 0.045) \times 10^{-2}$$



LSND has not been confirmed by dedicated experiment MINIBOONE (2001-2008)

However some small anomalies are still floating around in this domain feeding speculations and additional experimental activity

The oscillation signal $\bar{\nu}_\mu - \bar{\nu}_e$ in LSND complicates the global scenario:
 with 3 neutrinos only two independent Δm^2 are possible:

$$\Delta m_{12}^2 + \Delta m_{23}^2 + \Delta m_{31}^2 = 0$$

$$\Delta m_{\text{solar}}^2 + \Delta m_{\text{atm}}^2 \neq \Delta m_{\text{LSND}}^2$$

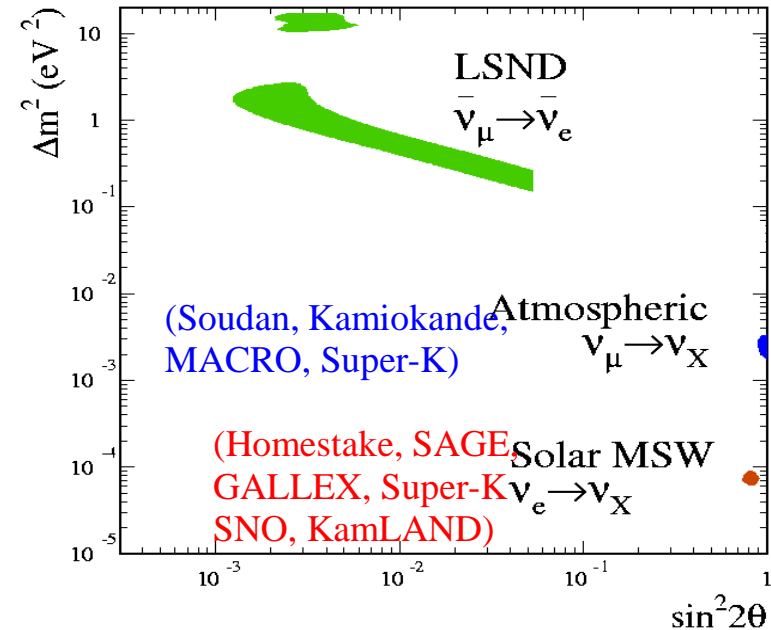
Oscillation signals:

- Solar: $\Delta m_{12}^2 \approx 7 \times 10^{-5} \text{ eV}^2$

- Atmospheric: $\Delta m_{23}^2 \approx 2.5 \times 10^{-3} \text{ eV}^2$

- LSND: $|\Delta m_{31}^2| = 0.2 - 2 \text{ eV}^2$

$$|\Delta m_{12}^2 + \Delta m_{23}^2 + \Delta m_{31}^2| = 0.2 - 2 \text{ eV}^2$$



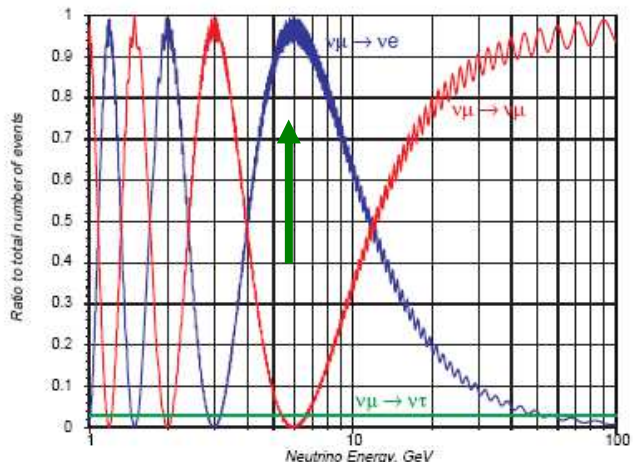
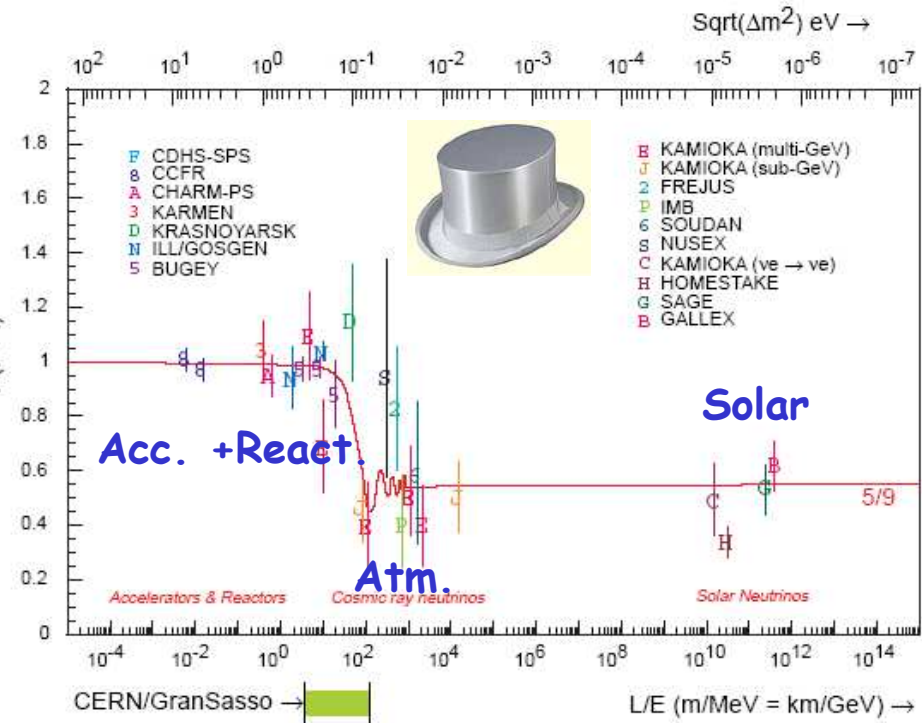
➤ At least 4 neutrinos are needed to reconcile all the results, from LEP it is known that the number of active light neutrinos is 3, so the other neutrinos must be sterile

➤ Even under this assumption the global fit of oscillation signals is poor: oscillations involving sterile neutrinos are disfavoured for the Atmospheric and Solar neutrinos, more sophisticated mechanisms like CPT violation must be advocated

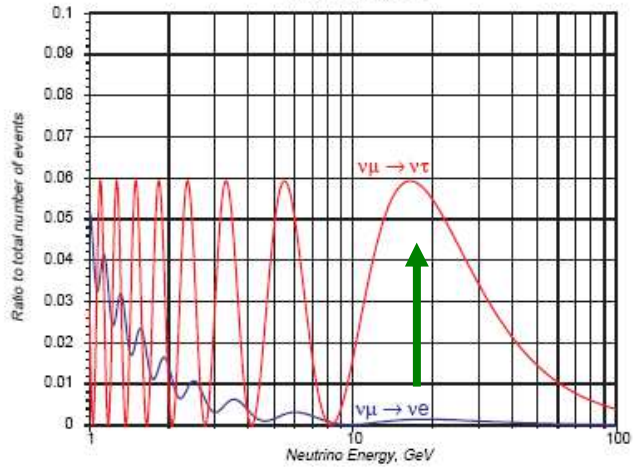
The Perkins plot (PLB 349 1995)

Interpretation of solar + atmospheric data in terms of just one $\nu_{\mu} \rightarrow \nu_e$ oscillation with $\Delta m^2 \sim 10^{-2} \text{ eV}^2$

The Acker-Pakvasa 3 flavours model hep-ph/9611423 included also LSND ($\Delta m^2 \sim 1 \text{ eV}^2$)



(C) Gran Sasso (735 km)



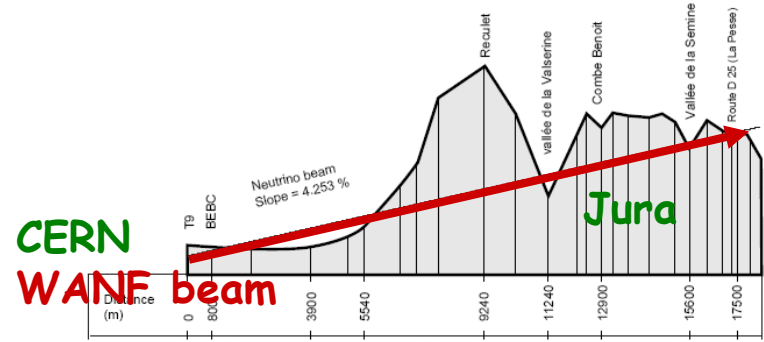
(B) Behind Jura (17 km)

Medium-baseline L/E ~ 1 Km/GeV

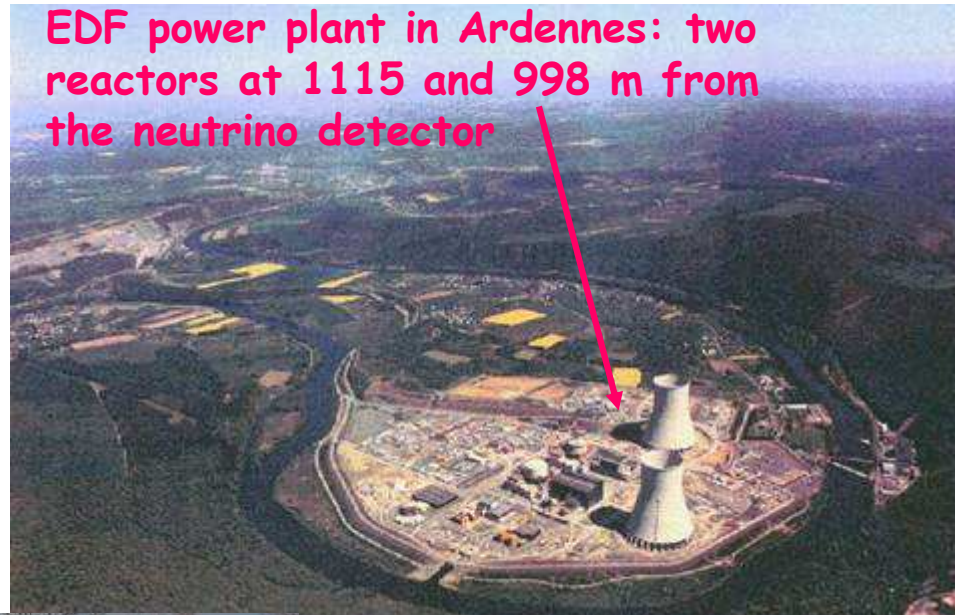
Icarus SPSLC 96/58 P304 19/12/1996

6. Conclusion

There is a substantial body of data leading to a theoretical prejudice which suggests that most probably the Gran Sasso and possibly the Jura locations, coupled with the SPS neutrino beam could be the real 'focal point' of the neutrino oscillation search. Spectacular $\nu_{\mu} \leftrightarrow \nu_{\tau}$ conversion is expected to be visible behind the Jura and a monumental $\nu_{\mu} \leftrightarrow \nu_e$ conversion is expected to be observed at the Gran Sasso position.



CHOOZ (the first long baseline experiment) 1997-1998



EDF power plant in Ardennes: two reactors at 1115 and 998 m from the neutrino detector

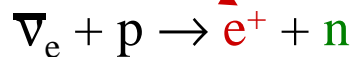
$\bar{\nu}_e \rightarrow \bar{\nu}_e$ (disappearance experiment at nuclear reactor)

$P_{th} = 8.5 \text{ GW}_{th}$, 1 detector at $L \sim 1 \text{ km}$, overburden equivalent to 300 m H_2O , Reactor neutrino flux known at 2.7 %, $L/E \sim 330 \text{ Km/GeV}$



Target: 5 ton liquid scintillator target with 0.09% Gadolinium

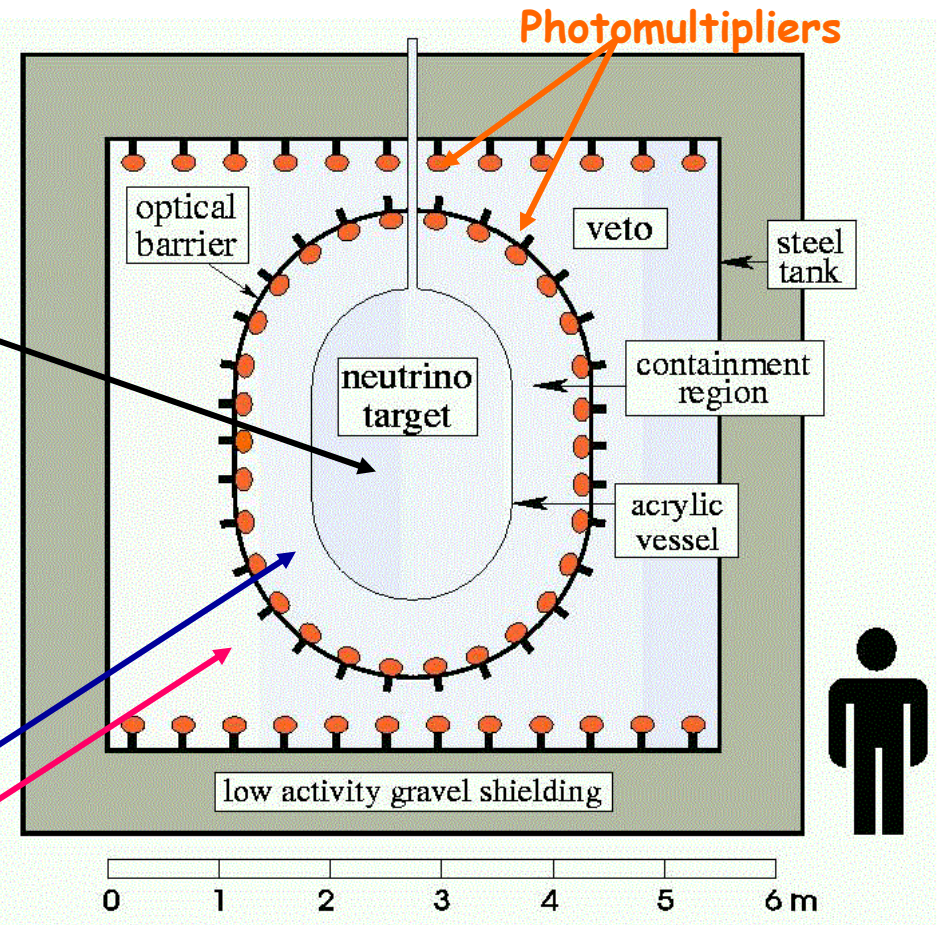
Prompt annihilation signal (γ rays)

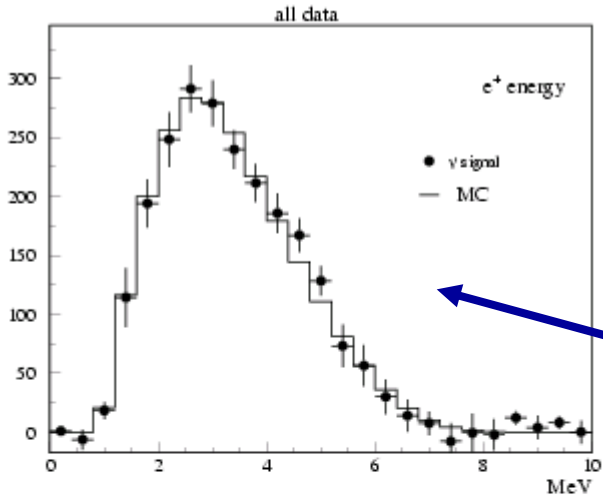


n capture on Gd after thermalization $\sim 30\mu\text{s}$

17 ton liquid scint. without Gd (containment of γ rays)

90 ton liquid scint. cosmic rays



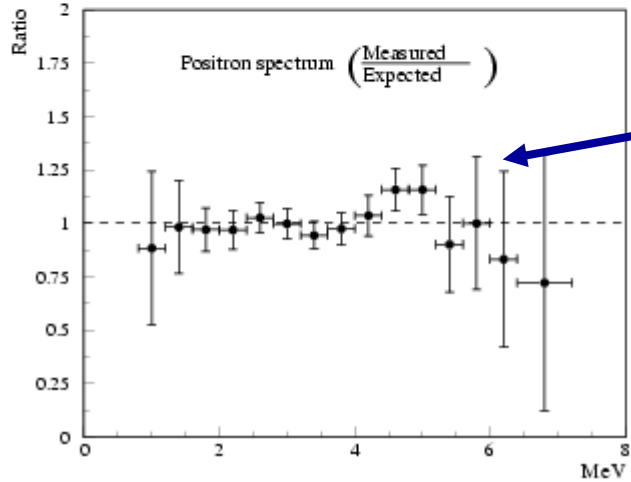
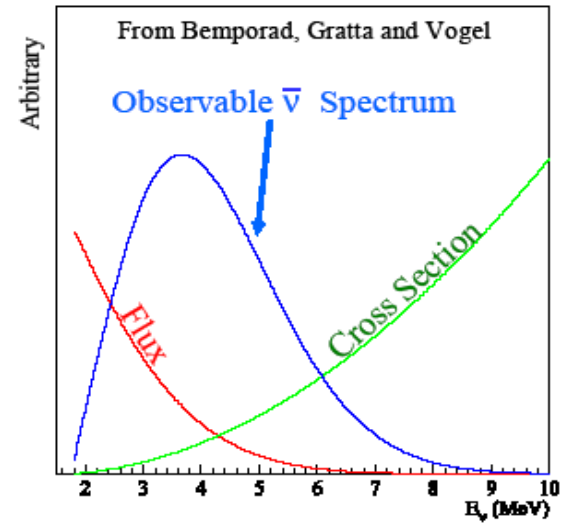


$$N(\bar{\nu}_e) \sim 2 \cdot 10^{20} \text{ s}^{-1} / \text{GW}_{\text{th}}$$

Signal ~ 25 events/day,
background (reactors off)
 ~ 1.2 events/day

Energy spectrum of the
positrons compared with the
predicted one (no oscillations)

$$E(\bar{\nu}_e) = E(e^+) + 1.8 \text{ MeV}$$

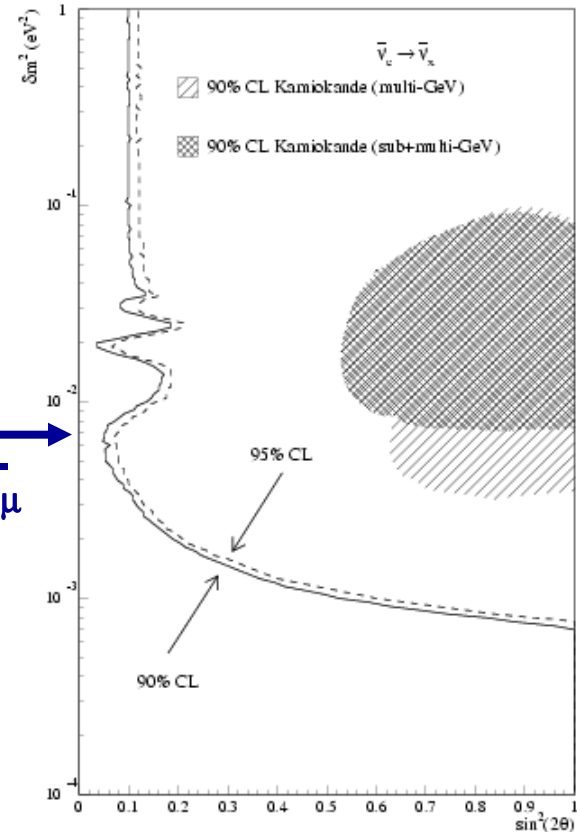


Ratio measured/expected

Integrated ratio =
 $1.01 \pm 0.028 \pm 0.027$

CHOOZ did not observe a
significant deficit of $\bar{\nu}_e$

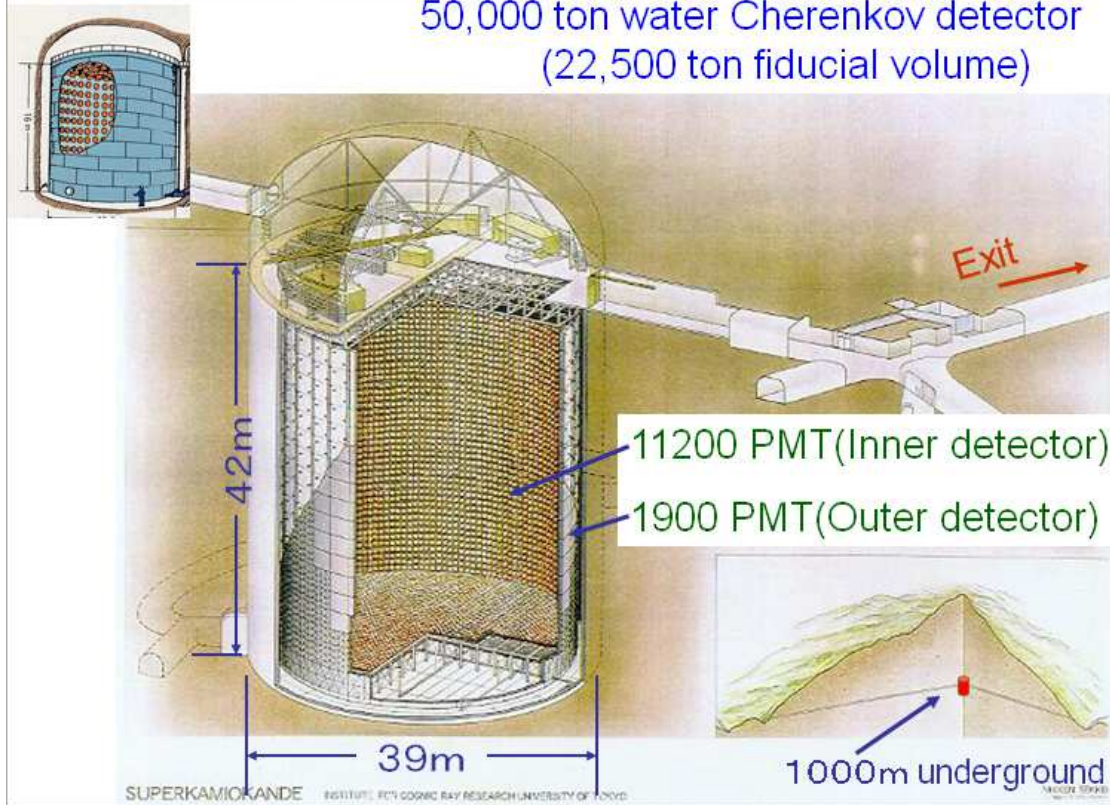
NO « monumental » $\bar{\nu}_e \rightarrow \bar{\nu}_\mu$
conversion



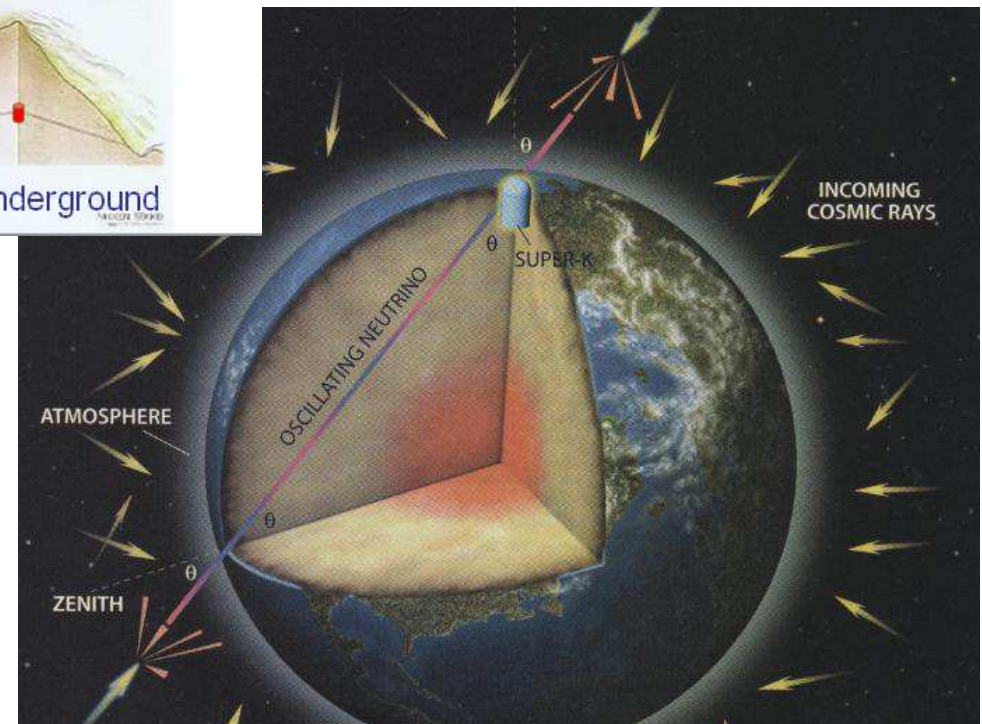
This result was published in 1998 before the Super-Kamiokande results and excluded the atmospheric neutrino anomaly interpretation in terms of $\nu_\mu \rightarrow \nu_e$ oscillations

Super-Kamiokande detector

50,000 ton water Cherenkov detector
(22,500 ton fiducial volume)

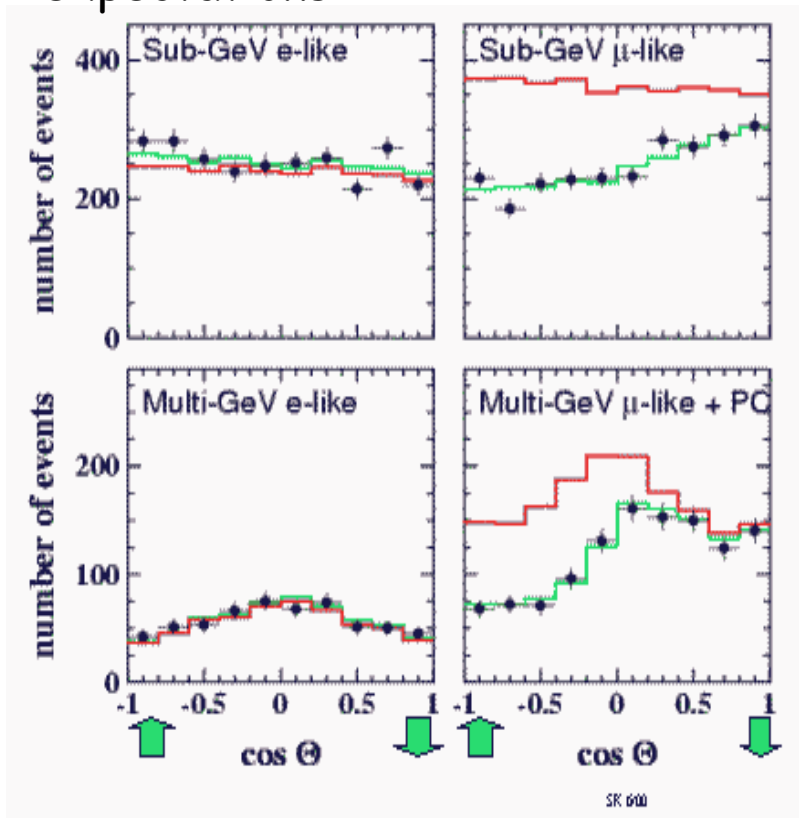


SuperKamiokande
1996- ... now



Neutrino 98 Conference in Takayama (June 1998)

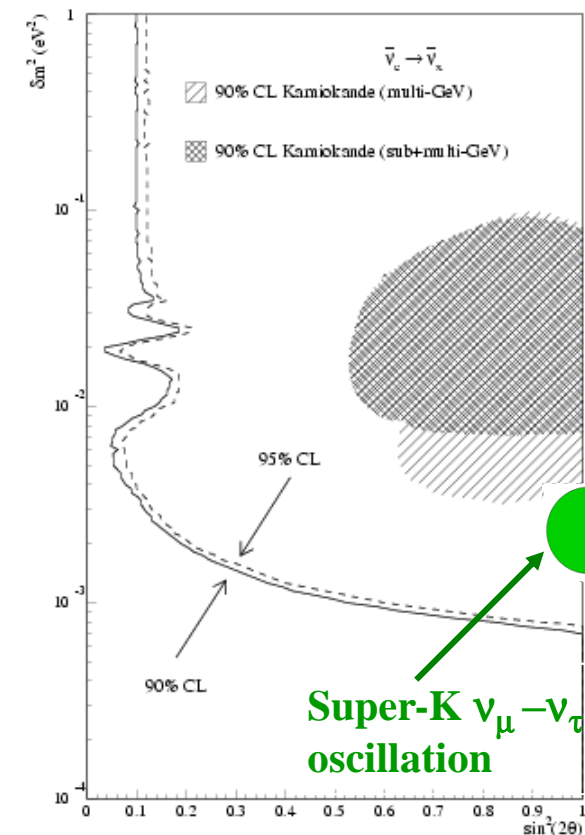
First results from Super-Kamiokande on atmospheric neutrinos, evidence of a zenith angle dependence of ν_μ disappearance, ν_e in agreement with expectations

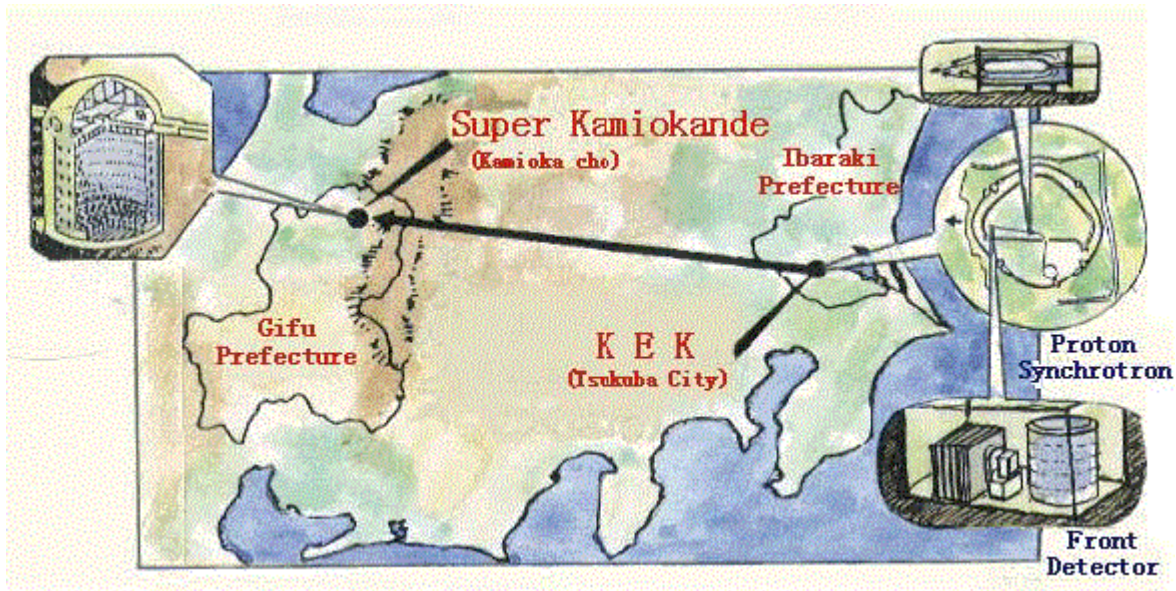


SK: Atmospheric neutrinos anomaly
 interpretable in terms of $\nu_\mu \rightarrow \nu_\tau$
 oscillations with a $\Delta m^2 \sim$ a few 10^{-3} eV^2

CHOOZ: no $\nu_\mu \rightarrow \nu_e$ oscillations, $\Theta_{13} < 11^\circ$

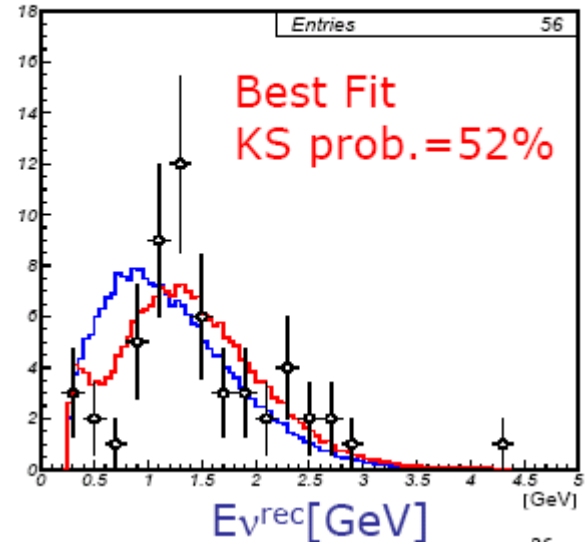
Neutrino oscillations start to be taken seriously as explanation of the atmospheric neutrinos anomaly
 Opens a campaign for a new generation of long baseline experiments to provide a final proof



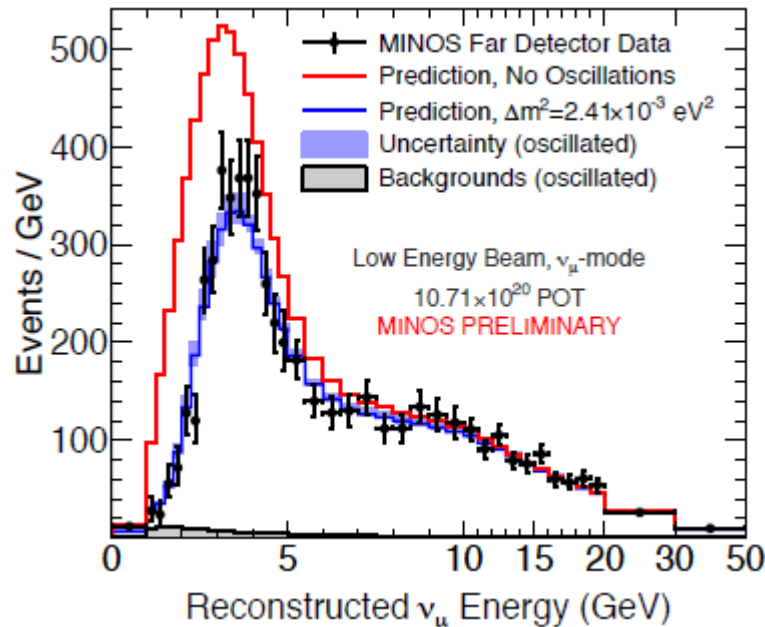


K2K results in 2004

- $N_{SK}^{obs} = 108$
- $N_{SK}^{exp} (best\ fit) = 104.8$



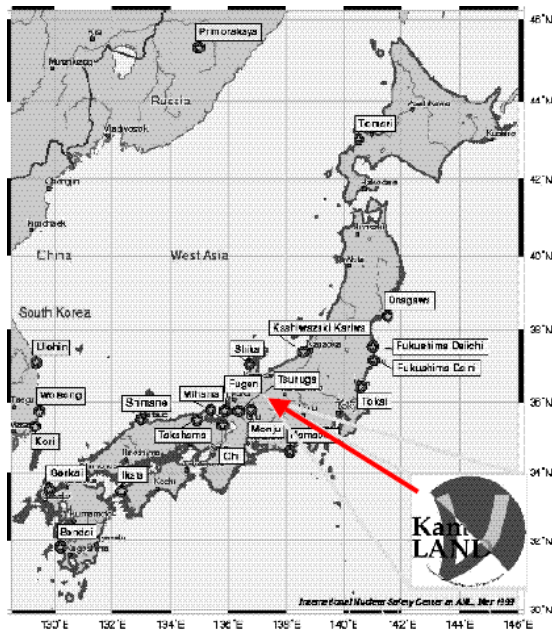
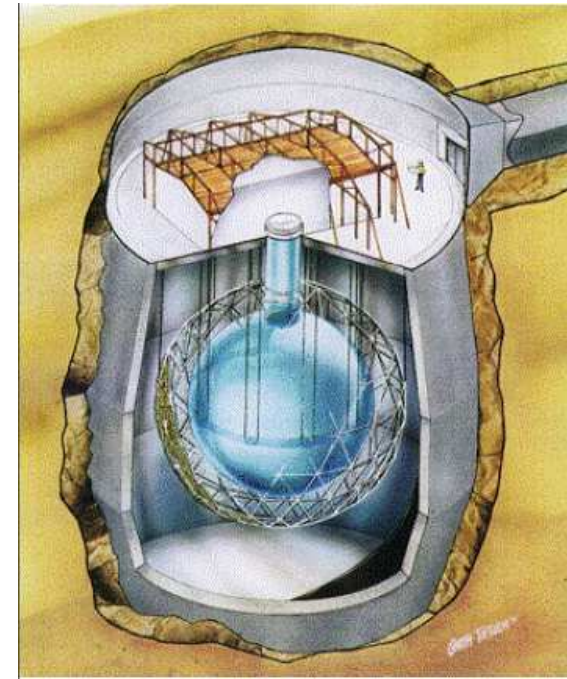
MINOS (U.S.) results in 2006



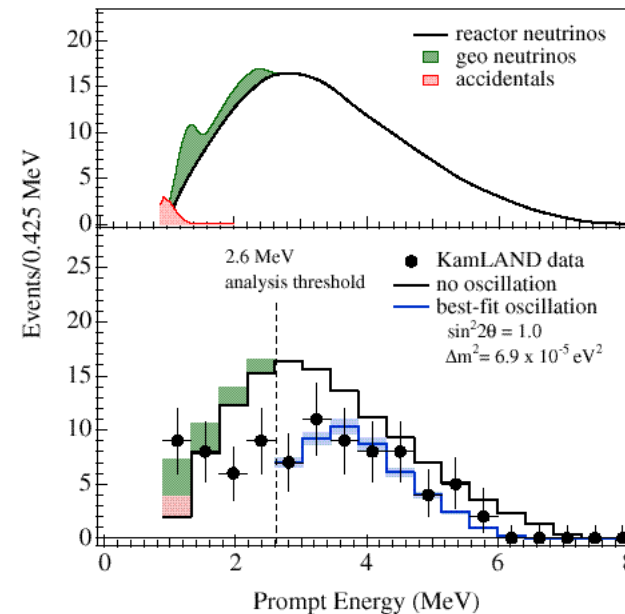
The final proof for solar neutrinos:

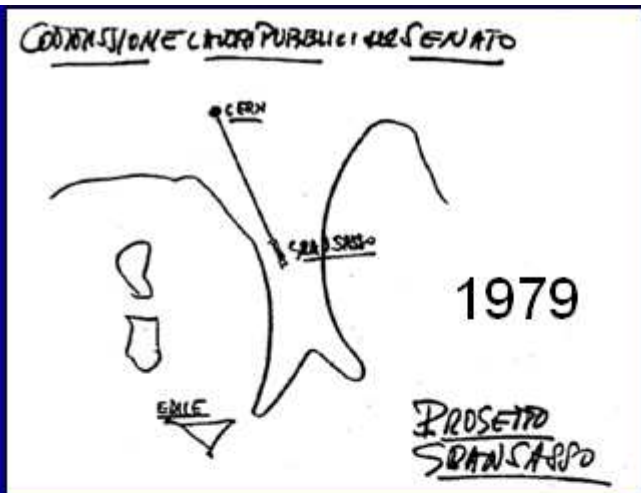
2001: SNO 1000 tons of heavy water, sensitive to neutral current reactions \rightarrow measure the total neutrino flux independently from their flavor
(NC) $\nu + d \rightarrow \nu + p + n$

The total neutrino flux agrees with the SSM!
 Electron neutrinos change into other neutrinos



2002: Kamland reactor experiment
 1000 ton liquid scintillator reproduces the solar neutrino oscillations on earth using antineutrinos of far reactors (on average 180 km)



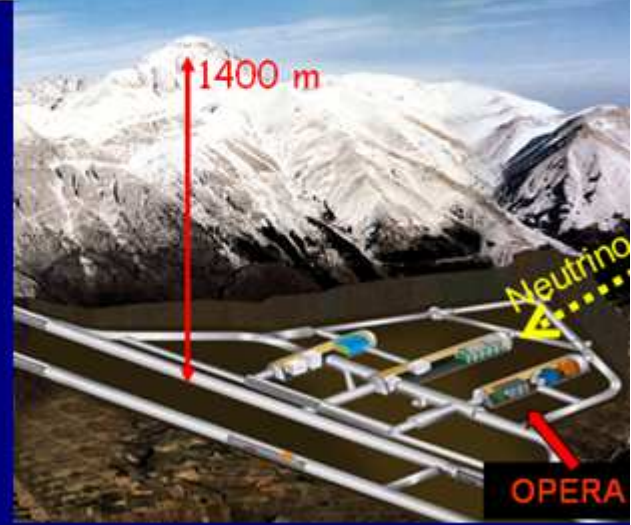


Cern Neutrinos to Gran Sasso

- Unambiguous evidence for $\nu_\mu \rightarrow \nu_\tau$ oscillations in the region of atmospheric neutrinos by looking for ν_τ appearance in a pure ν_μ beam
- Search for the subleading $\nu_\mu \rightarrow \nu_e$ oscillations

- Beam: CNGS (1999)
 - ν_τ appearance experiments at LNGS
 - No near detectors needed in appearance mode
- CNGS1 (2000)

CNGS2 (2002)



$$N_{\tau} \approx 1.61 \sin^2(2\theta) (\Delta m^2)^2 L^2 \int_{3.5 \text{ GeV}}^{E_{\text{max}}} \Phi_{\mu}(E) \frac{\sigma_{\tau}(E)}{E^2} dE$$

Dependence of the events rate on $(\Delta m^2)^2$

Signal constant as a function of L for $L \ll \lambda$
 Φ_{μ} contains a factor $1/L^2$

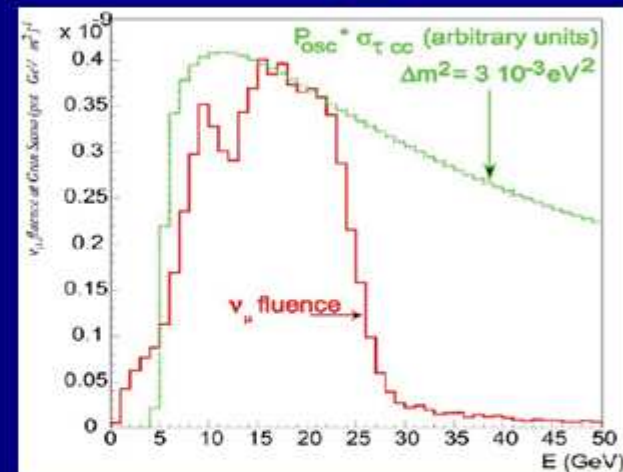
CNGS fluxes

$\langle E(\nu_{\mu}) \rangle$	17 GeV
L	730 km
L/E	43 Km/GeV
$(\nu_e + \bar{\nu}_e) / \nu_{\mu \text{ CC}}$	0.87%
$\nu_{\mu} / \nu_{\mu \text{ CC}}$	2.1%
ν_{τ} prompt	negligible

Quantity to be optimized playing with the beam spectrum:

ν_{μ} flux vs $\sigma(\tau)/E^2$

→ produce the max. number of ν_{τ} CC



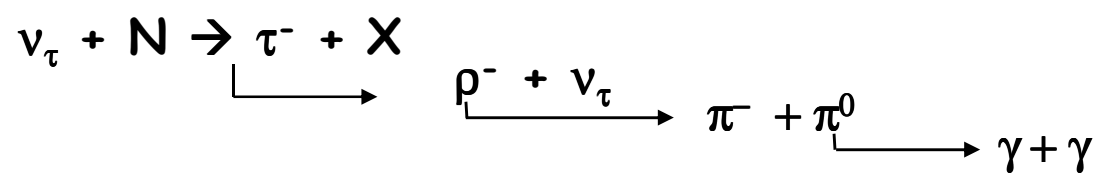
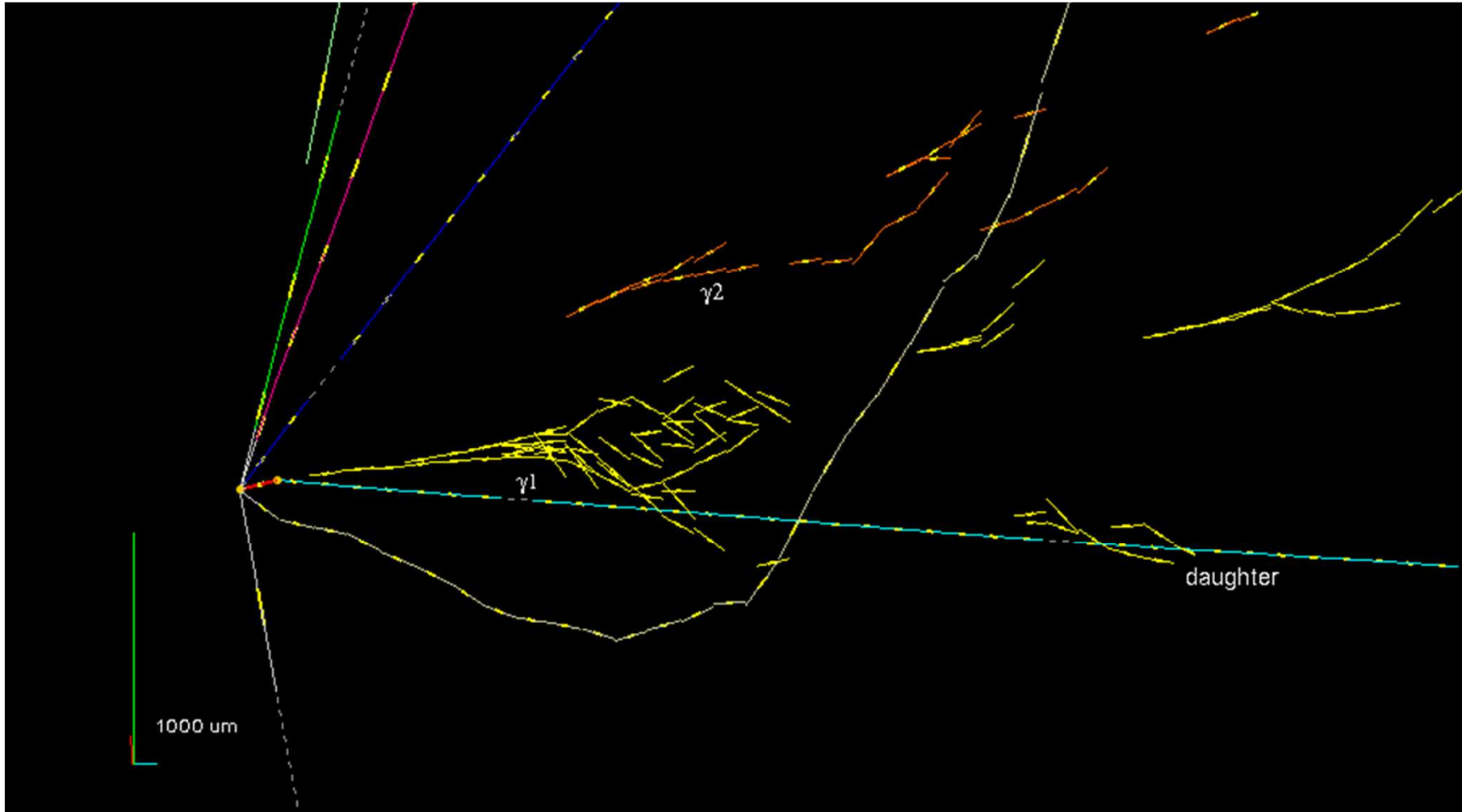
- Nominal beam performance ($4.5 \cdot 10^{19}$ pot/y)
- OPERA Target mass of 1.25 kton

→ Expected number of interactions in 5 years :
 ~ 23600 ν_{μ} CC+NC
 ~ 170 $\nu_e + \bar{\nu}_e$ CC
 ~ 115 ν_{τ} CC ($\Delta m^2 = 2.5 \times 10^3 \text{ eV}^2$)

After efficiencies, 8 tau decays are expected, with <1 background events

First OPERA ν_τ candidate
(single hadronic prong τ decay)

<http://arxiv.org/abs/1006.1623>
Physics Letters B (PLB-D-10-00744)



Visible tau decay topology
with kink and two gammas

Neutrino velocity measurement

Time

Neutrinos production time (CERN)
Neutrinos interaction time (OPERA detector)

Space

Accurate determination of the distance (Geodesy)

→ Massive application of metrology techniques in HEP

2009-2011 measurement with standard CNGS beam → Blind analysis: “box” opened after assessment of delays, previously fixed to arbitrary values:

- High neutrino energy - high statistics ~15000 events
- Precise measurement of neutrino time distribution at CERN through proton waveforms
- Sophisticated timing system: ~1 ns CNGS-OPERA synchronization
- Calibrations techniques of CNGS and OPERA timing chains: ~ 1 ns level
- Measurement of baseline by global geodesy: 20 cm accuracy over 730 km (longest neutrino baseline actually available)

→ Result: ~10 ns overall accuracy on TOF with similar stat. and sys. errors


GPS common-view mode



R. Garwin
one of the fathers of GPS

Standard GPS operation:

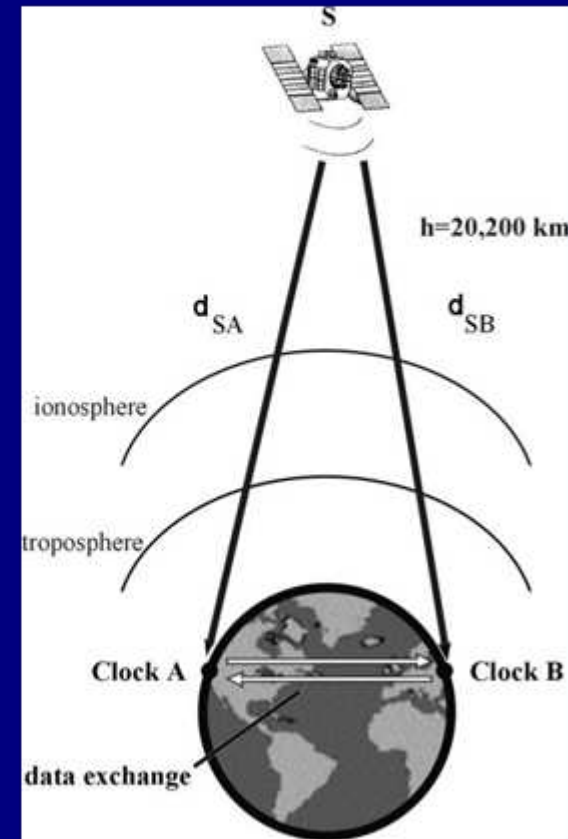
resolves x, y, z, t with ≥ 4 satellite observations

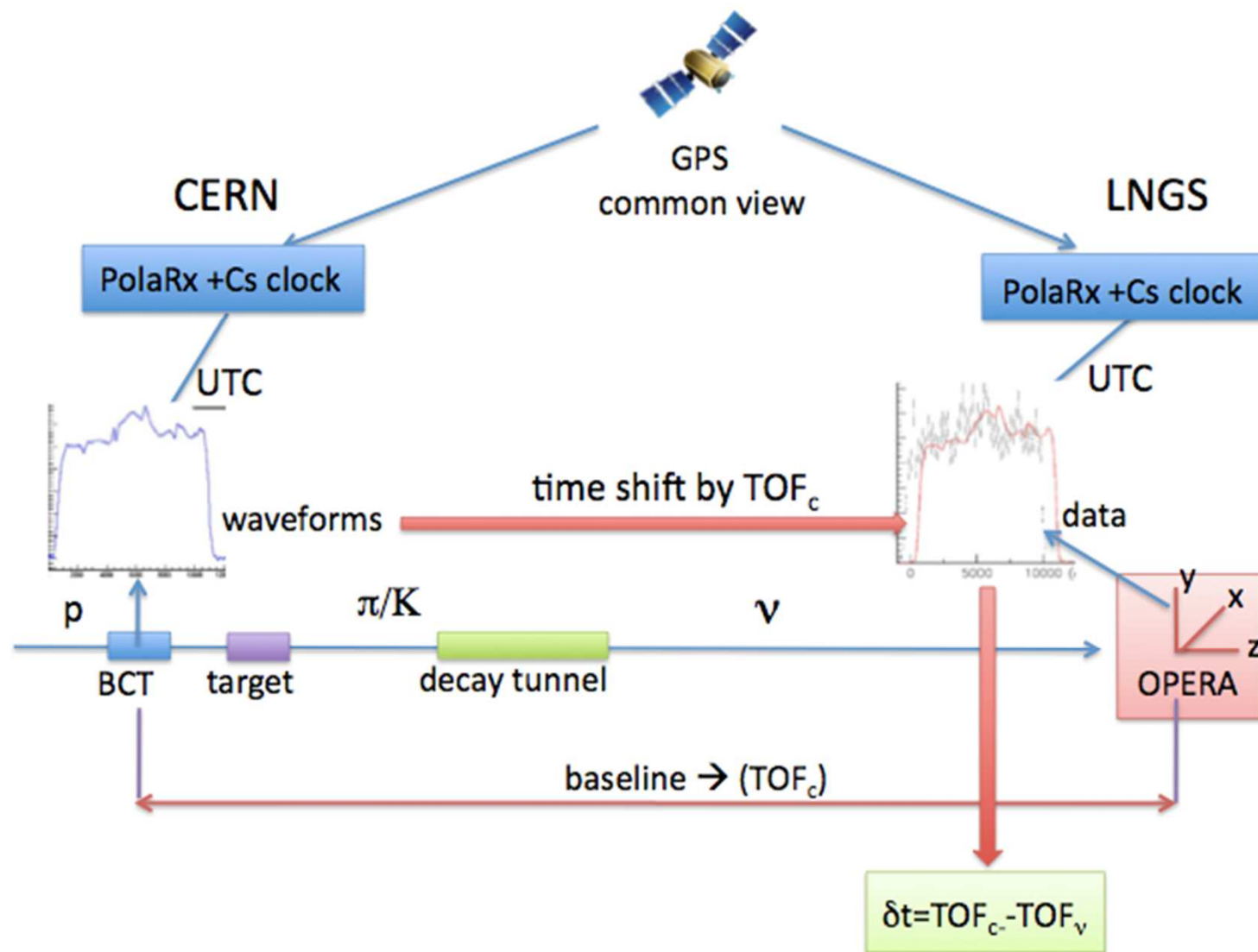
Common-view mode (the same satellite for the two sites, for each comparison): 

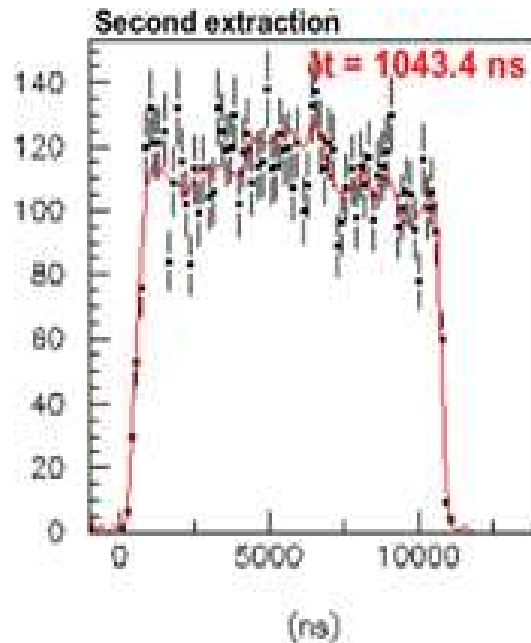
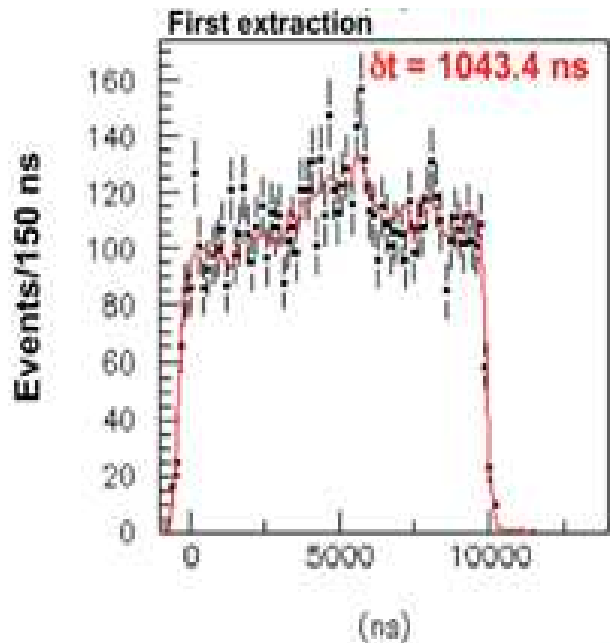
- x, y, z known from former dedicated measurements
- determine time differences of local clocks (both sites) w.r.t. the satellite, by offline data exchange
- $730 \text{ km} \ll 20000 \text{ km}$ (satellite height) \rightarrow similar paths in ionosphere

Twin geodetic GPS receiver PolaRx2e + Cs clock \rightarrow Time-transfer ($\sim 1 \text{ ns}$ accuracy)

Standard technique used in TAI (International Atomic Time) by BIPM



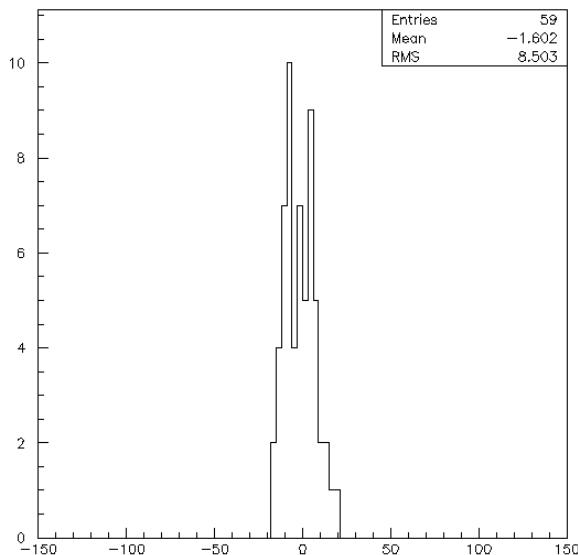
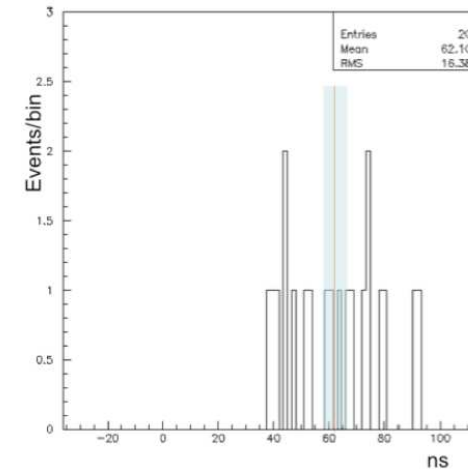




Neutrino data vs protons: after delay determination (2009-2011) ~15000 neutrinos (September 2011)

UNBLIND: $\delta t = \text{TOF}_c - \text{TOF}_{\nu}$
 $(57.8 \pm 7.8 \text{ (stat.)} - 5.9 + 8.3 \text{ (sys.)}) \text{ ns}$

November 2011 (20 neutrinos measured individually)
 $\delta t = (62.1 \pm 3.7 \text{ (stat.)}) \text{ ns}$



December-February 2012
 Two unknown sources of bias identified

May 2012 (59 neutrinos measured individually)

$\delta t = (-1.6 \pm 1.1 \text{ (stat.)}^{+6.1}_{-3.7} \text{ (sys.)}) \text{ ns}$

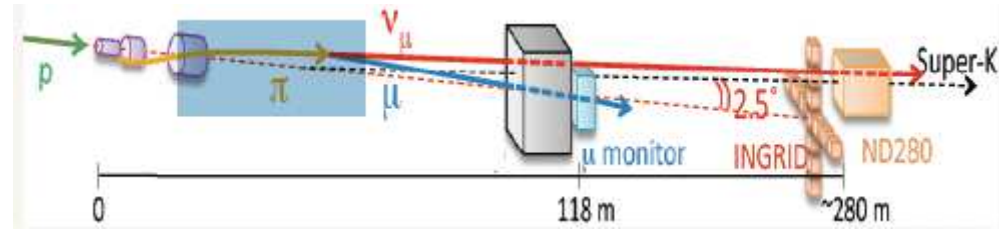
$\frac{v-c}{c} = \frac{\delta t}{\text{TOF}_c - \delta t} = (-0.7 \pm 0.5 \text{ (stat.)}^{+2.5}_{-1.5} \text{ (sys.)}) \times 10^{-6}$

The search for θ_{13} :

If the angle θ_{13} is different than zero at the same L/E of the atmospheric neutrino oscillation there should be a subleading oscillation between the muonic neutrino and the electron neutrino

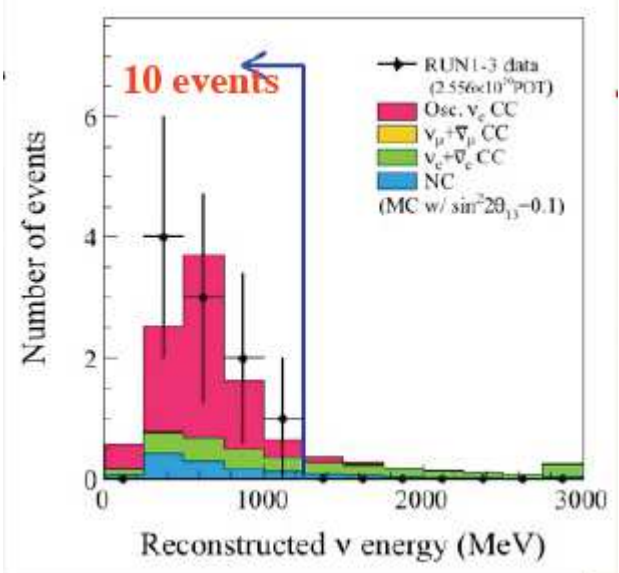


The determination θ_{13} of is fundamental for the next steps in neutrino physics (study of CP violation)



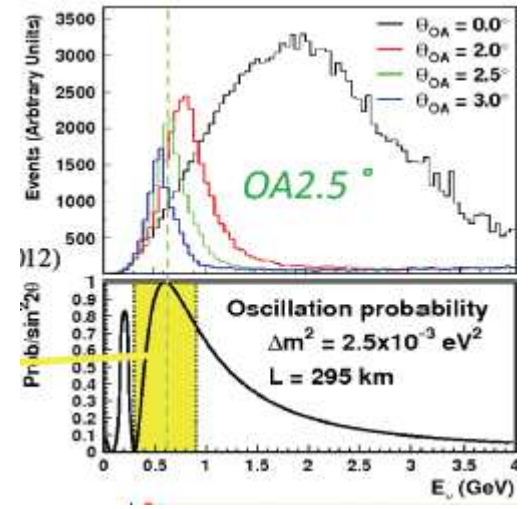
The off-axis neutrino beam

First hint of electron neutrino appearance signal in March 2011 before the earthquake



10 events of electron neutrino appearance observed with a background of 2.7

3.2 sigmas



2012 the Daya Bay experiment



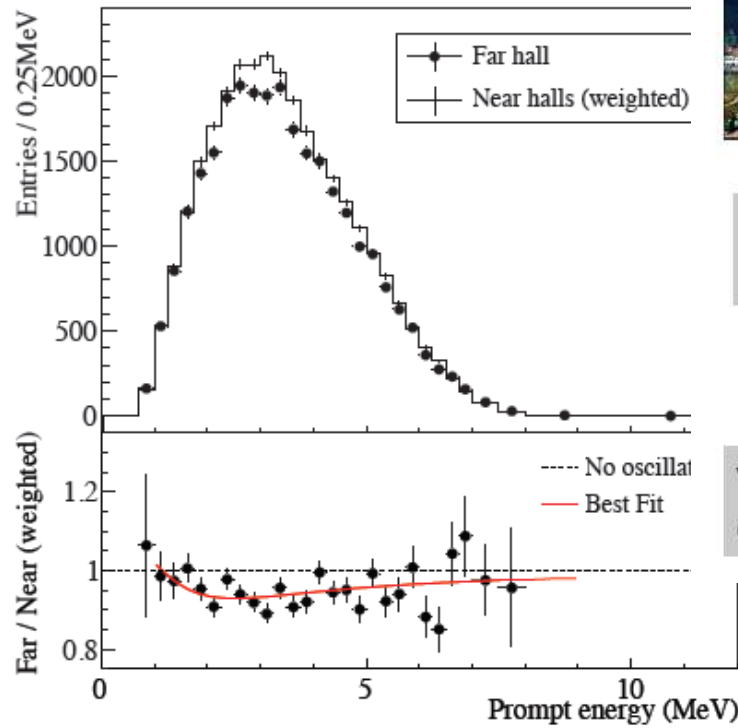
Adjacent mountains with horizontal access provide 860 (250) m.w.e cosmic shielding.

Daya Bay Ling Ao I + II

6 commercial reactor cores with 17.4 GW_{th} total power.

6 Antineutrino Detectors (ADs) give 120 tons total target mass.

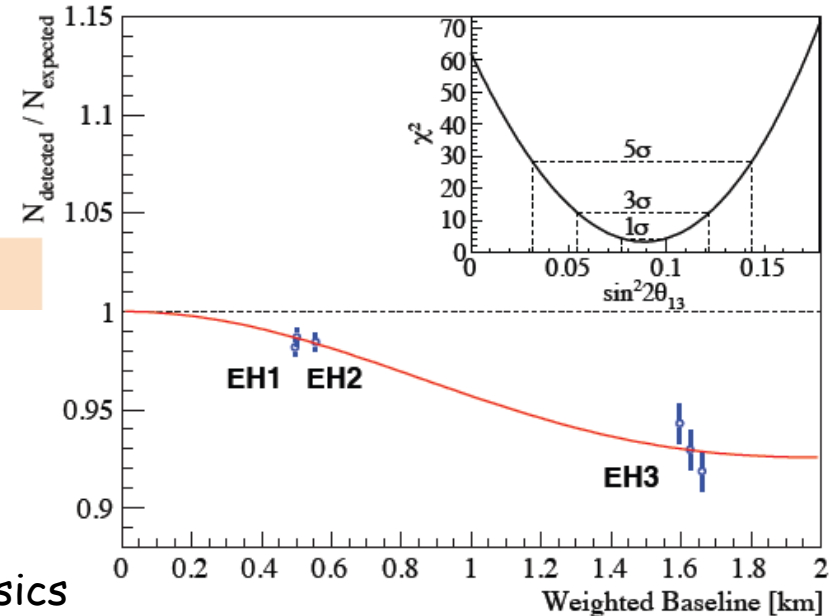
Via GPS and modern theodolites, relative detector-core positions known to 3 cm.



$R = 0.944 \pm 0.007 \text{ (stat)} \pm 0.003 \text{ (syst)}$

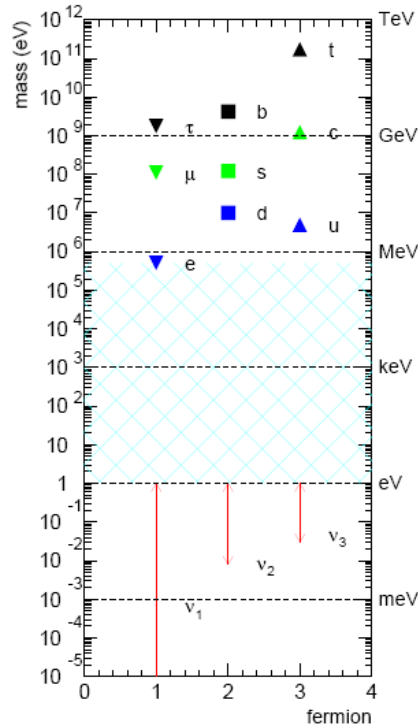
$\sin^2 2\theta_{13} = 0.089 \pm 0.010 \text{ (stat)} \pm 0.005 \text{ (syst)}$

The value of $\sin^2 2\theta_{13}$ is just below the CHOOZ limit
 Confirmed by other two reactor experiments:
 Double-CHOOZ and RENO
 A huge boost to future neutrino oscillations physics



Neutrinos: a window beyond the S.M. → G.U.T.

Fundamental questions related to a deeper description of physics and to the evolution of the universe



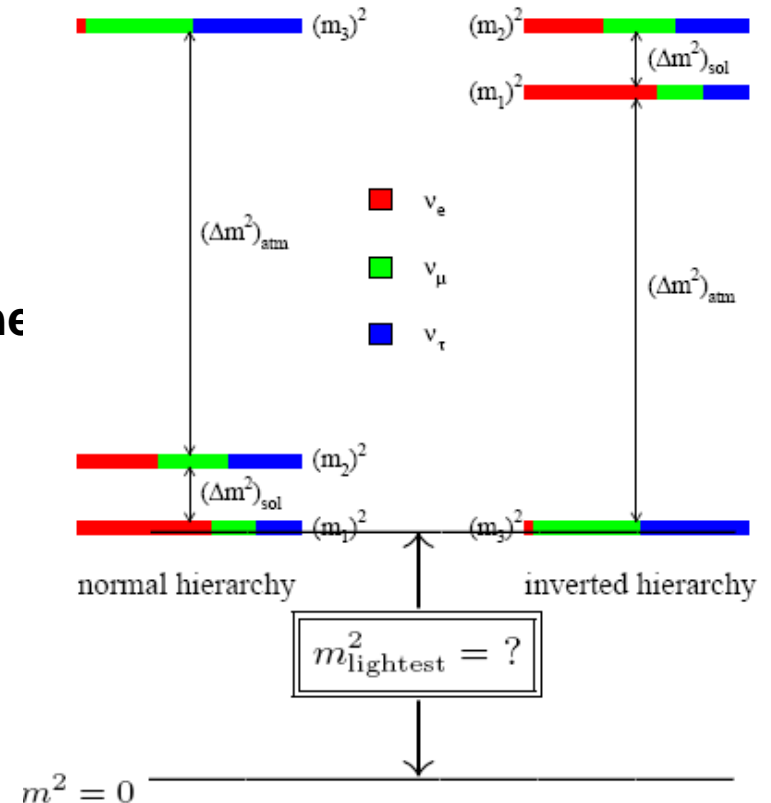
- Why are neutrino masses so small ?
- Why is the mixing matrix so different than the one of the quarks ?

What is this very strange puzzle suggesting us ?

$$V_{MNS} \sim \begin{pmatrix} 0.8 & 0.5 & 0.2 \\ 0.4 & 0.6 & 0.7 \\ 0.4 & 0.6 & 0.7 \end{pmatrix} \quad V_{CKM} \sim \begin{pmatrix} 1 & 0.2 & 0.001 \\ 0.2 & 1 & 0.01 \\ 0.001 & 0.01 & 1 \end{pmatrix}$$

- Which is the mass of the lightest state
- Are neutrinos Majorana particles ?
- Which is the hierarchy of the mass eigenstates ?
- Is there CP violation in the neutrino sector ?

CP violation in the neutrino sector can explain the matter/antimatter asymmetry in the universe



Several projects are being proposed and discussed for the next steps
 An example in Europe submitted to CERN in June 2012

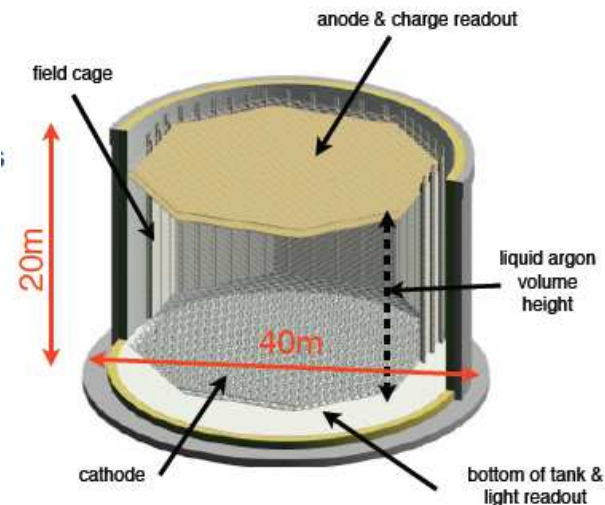


LBNO:

Beam from CERN to Finland over 2300 km

Profits of the study of underground sites performed by the european program LAGUNA

50 kton Liquid Argon detector for the determination of neutrino mass hierarchy, search for CP violation, search for proton decay and for supernova neutrinos



I hope you will find neutrino physics interesting

Thanks for your attention