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Zurich^{UZH}



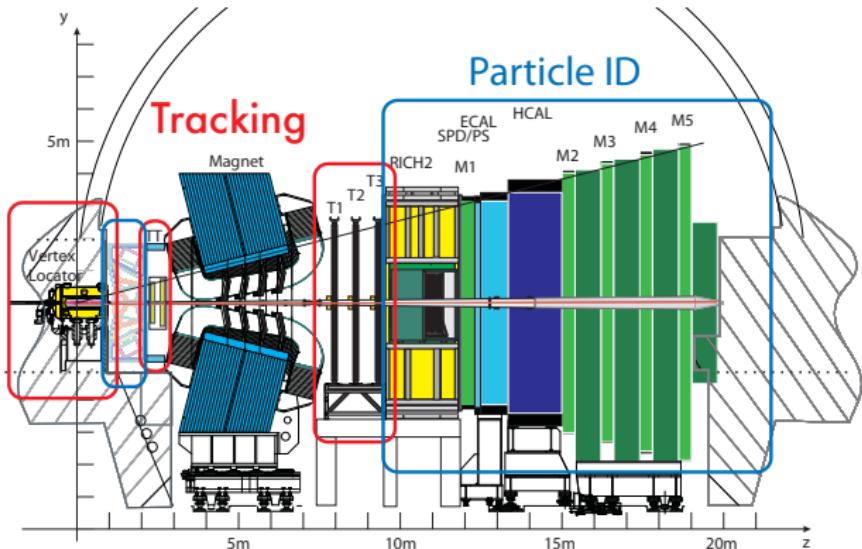
PHYSICS WITH EW PENGUINS AT LHCb

7TH INTERNATIONAL WORKSHOP ON THE CKM UNITARITY
TRIANGLE, CINCINNATI, OHIO, USA
28. SEPTEMBER - 2. OCTOBER 2012



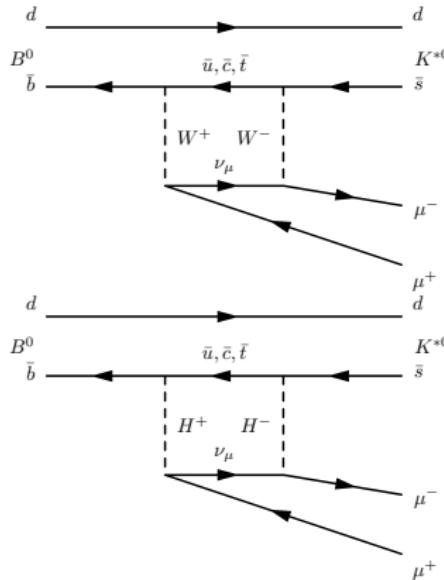
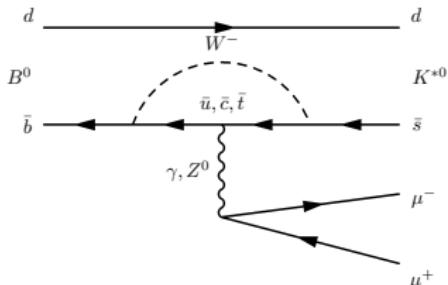
MICHEL DE CIAN
ON BEHALF OF THE LHCb COLLABORATION

THE LHCb DETECTOR



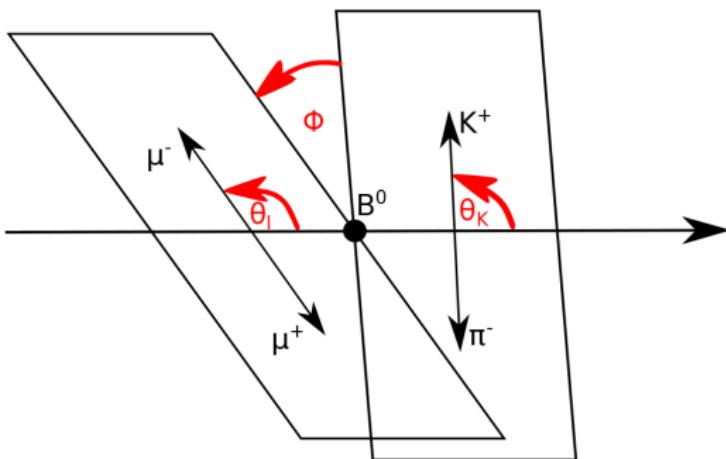
- LHCb covers a pseudorapidity $\eta = 2 - 5$.
- Excellent momentum resolution: $\Delta p/p = 0.4\% - 0.6\%$ in $5 - 140 \text{ GeV}/c$.
- $K - \pi$ separation up to $100 \text{ GeV}/c$.
- All presented analyses with $\approx 1 \text{ fb}^{-1}$ collected in 2011.

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$



- Rare decay with $\mathcal{B} = (1.05^{+0.16}_{-0.13}) \times 10^{-6}$ [PDG]
- Decay only possible via penguin- or box diagrams, "new physics" can enter at the same level as SM physics.
- Four-particle final state: Plenty of observables in angular distributions.

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$: ANGULAR DISTRIBUTION (I)



- Decay can be fully described by three angles ($\cos \theta_\ell$, $\cos \theta_K$, ϕ) and the dimuon invariant mass (square) q^2 .

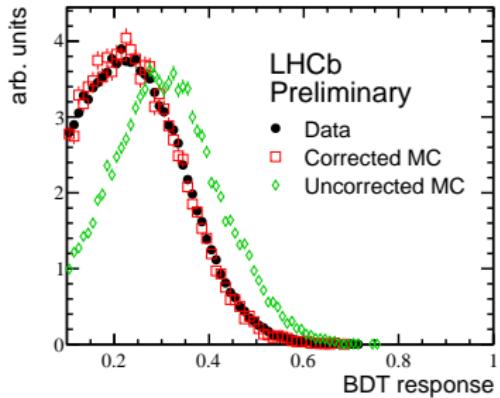
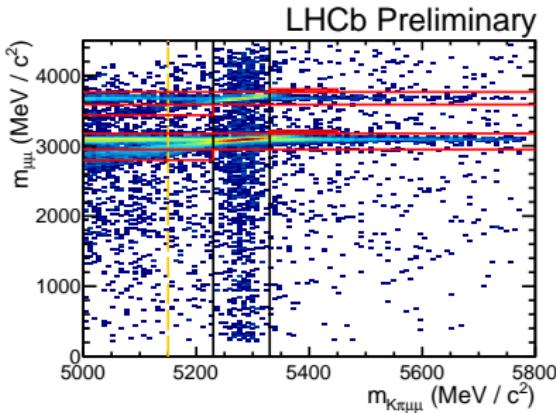
$B^0 \rightarrow K^{*0} \mu^+ \mu^-$: ANGULAR DISTRIBUTION (II)

- Apply "folding": $\phi \rightarrow \phi + \pi$ for $\phi < 0$: Cancels four terms.
- And leaves...

$$\begin{aligned} \frac{d^4(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell \, d \cos \theta_K \, d\phi \, dq^2} \propto & \quad \textcolor{red}{F_L} \cos^2 \theta_K + \frac{3}{4}(1 - \textcolor{red}{F_L})(1 - \cos^2 \theta_K) + \\ & \textcolor{red}{F_L} \cos^2 \theta_K (2 \cos^2 \theta_\ell) + \\ & \frac{1}{4}(1 - \textcolor{red}{F_L})(1 - \cos^2 \theta_K)(2 \cos^2 \theta_\ell - 1) + \\ & \textcolor{red}{S_3}(1 - \cos^2 \theta_K)(1 - \cos^2 \theta_\ell) \cos 2\phi + \\ & \frac{4}{3}\textcolor{red}{A_{FB}}(1 - \cos^2 \theta_K) \cos \theta_\ell + \\ & \textcolor{red}{S_9}(1 - \cos^2 \theta_K)(1 - \cos^2 \theta_\ell) \sin 2\phi \end{aligned}$$

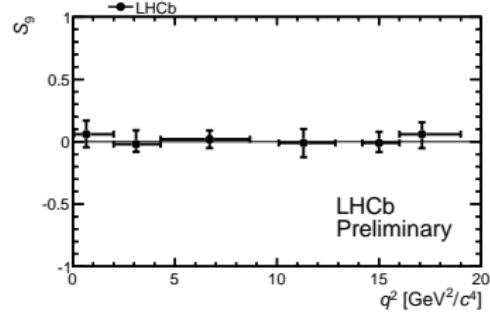
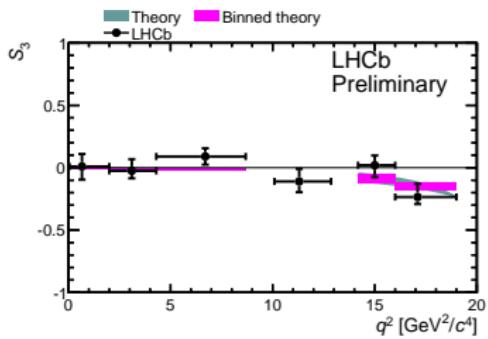
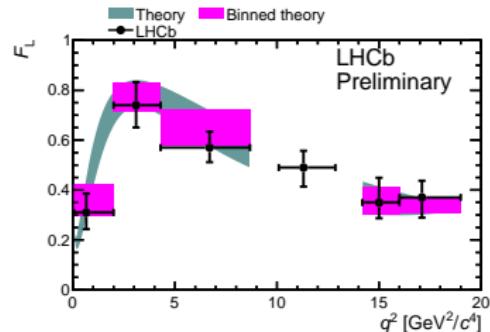
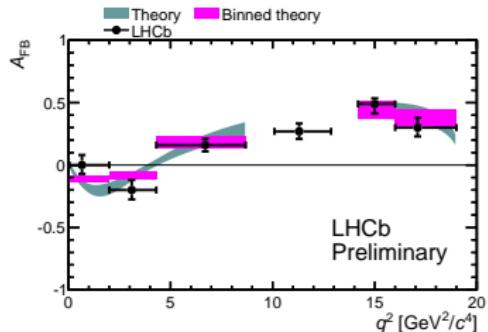
- This expression was simultaneously fitted to the angles and the invariant mass in 2011 dataset.
- The S_i expressions are the \mathcal{CP} averaged I_i expressions.
- Neglect lepton masses and S-wave contribution (\rightarrow systematics).

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$: EXPERIMENTAL ASPECTS



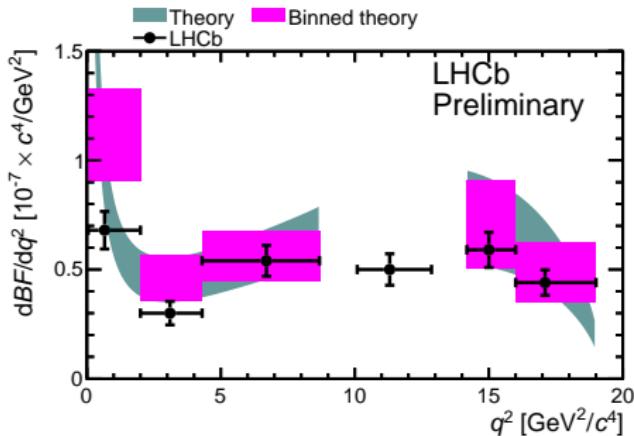
- Some experimental details:
 - Cut out $B^0 \rightarrow J/\psi K^{*0}$ and $B^0 \rightarrow \psi(2S)K^{*0}$, veto peaking background, e.g. $B_s^0 \rightarrow \phi \mu^+ \mu^-$.
 - Select signal events with a BDT.
 - Correct for acceptance effects with event-by-event correction using simulation.
 - Correct for simulation \leftrightarrow data differences with control channels (e.g. $J/\psi \rightarrow \mu^+ \mu^-$ for particle identification)

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$: ANGULAR OBSERVABLES



Theory predictions from [C. Bobeth et al., JHEP 07 (2011) 067] and references therein

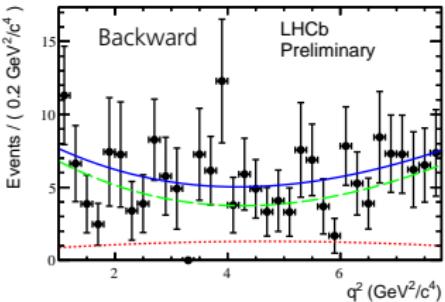
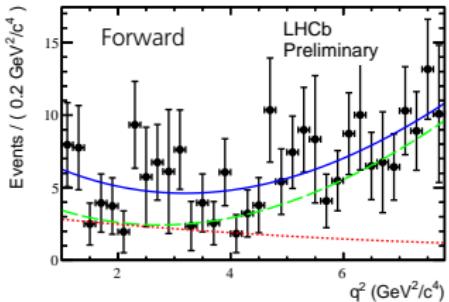
$B^0 \rightarrow K^{*0} \mu^+ \mu^-$: DIFFERENTIAL BRANCHING FRACTION



- Differential branching fraction determined with normalisation to $B^0 \rightarrow J/\psi K^{*0}$.

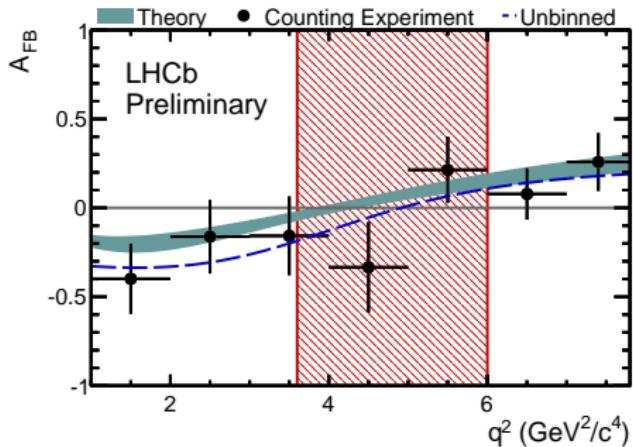
Theory prediction from [C. Bobeth et al., JHEP 07 (2011) 067] and references therein

MEASURING THE ZERO-CROSSING POINT IN A_{FB} OF $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (I)



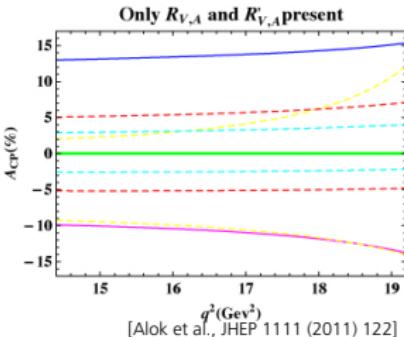
- Zero-crossing point of A_{FB} has a very clean prediction, as the form factors cancel (to first order).
- Zero-crossing point was extracted using "unbinned counting" technique:
 - Split dataset in "forward" and "backward" events (with respect to $\cos \theta_\ell$).
 - Perform a 2D unbinned extended maximum likelihood fit to (q^2, mass) for forward and backward.
- Extract $A_{FB} = \frac{N_F \cdot PDF_F(q^2) - N_B \cdot PDF_B(q^2)}{N_F \cdot PDF_F(q^2) + N_B \cdot PDF_B(q^2)}$

MEASURING THE ZERO-CROSSING POINT IN A_{FB} OF $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (II)



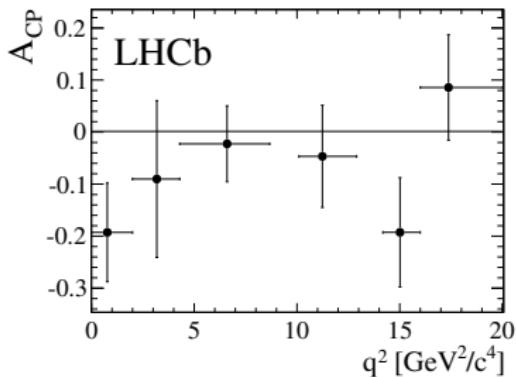
- Standard Model theory predicts zero-crossing in $4.0 - 4.3 \text{ GeV}^2/c^4$ (central values)
e.g. [C. Bobeth et al., JHEP 1201 (2012) 107][M. Beneke et al., Eur. Phys. J. C41 (2005), 173][A. Ali et al., Eur. Phys. J. C47 (2006) 625]
- LHCb preliminary result: $4.9^{+1.1}_{-1.3} \text{ GeV}^2/c^4$

\mathcal{CP} -ASYMMETRY IN $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (I)



- Form $\mathcal{A}_{CP} = \frac{\Gamma(\bar{B}^0 \rightarrow \bar{K}^{*0} \mu^+ \mu^-) - \Gamma(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{\Gamma(\bar{B}^0 \rightarrow \bar{K}^{*0} \mu^+ \mu^-) + \Gamma(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}$
- Predicted to be $\mathcal{O}(10^{-3})$ in SM, very clean prediction due to form factor suppression.
- Asymmetry up to 15% in certain models.
- Use the same corrections / selection / binning-scheme as for the $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ angular analysis.
- Additional difficulty: Measured is
$$\mathcal{A}_{raw} = \mathcal{A}_{CP} + \mathcal{A}_{\text{Detector}} + \kappa \mathcal{A}_{\text{Production}}$$
- $\kappa = \frac{\int_0^\infty \epsilon(t) e^{-\Gamma t} \cos \Delta m t \, dt}{\int_0^\infty \epsilon(t) e^{-\Gamma t} \cosh \frac{\Delta m t}{2} \, dt}$

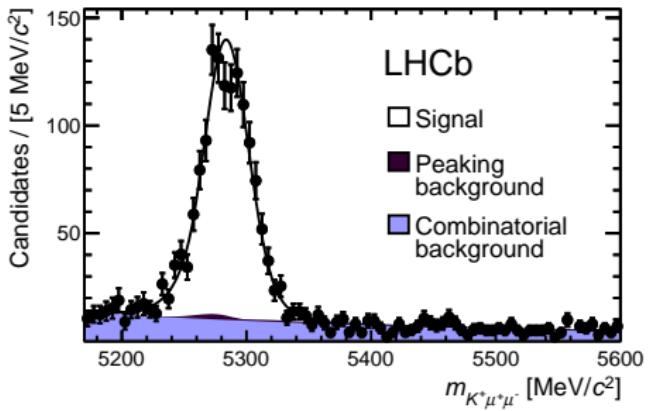
\mathcal{CP} -ASYMMETRY IN $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (II)



- $\mathcal{A}_{\text{Detector}}$: Detector asymmetries cancel when taking the average between the two magnet polarities.
- $\mathcal{A}_{\text{Detector}}/\mathcal{A}_{\text{Production}}$: Use $B^0 \rightarrow J/\psi K^{*0}$ as a control channel.
- $\mathcal{A}_{\mathcal{CP}} \approx \mathcal{A}_{\text{raw}}(B^0 \rightarrow K^{*0} \mu^+ \mu^-) - \mathcal{A}_{\text{raw}}(B^0 \rightarrow J/\psi K^{*0})$
- Residual differences due to kinematical differences are accounted for in the systematic uncertainty.
- $\mathcal{A}_{\mathcal{CP}} = -0.072 \pm 0.040 \text{ (stat)} \pm 0.005 \text{ (sys)}$

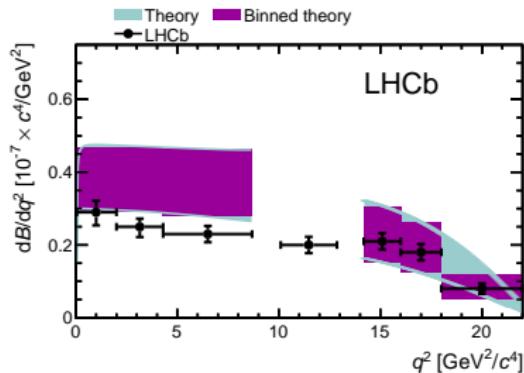


MEASUREMENT OF $B^+ \rightarrow K^+ \mu^+ \mu^-$



- Rare decay, $\mathcal{B} = (4.36 \pm 0.15 \pm 0.18) \times 10^{-7}$ [arXiv:1209.4284]
- Use a loose preselection and a BDT for the final selection. Training on $B^+ \rightarrow J/\psi K^+$ (signal) and $B^+ \rightarrow K^+ \mu^+ \mu^-$ sidebands (background).
- Cut out resonant regions of $B^+ \rightarrow J/\psi K^+$ and $B^+ \rightarrow \psi(2S)K^+$.
- Remaining peaking background accounted for in the fit.

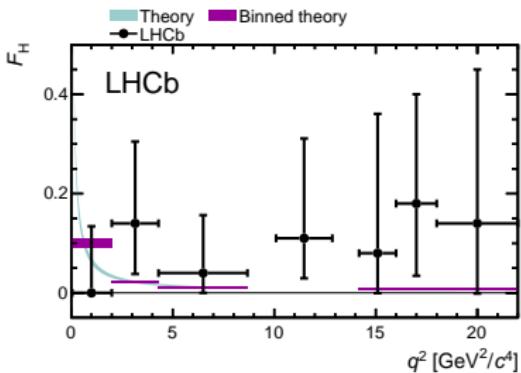
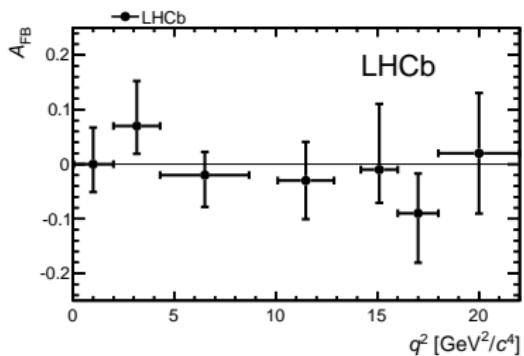
MEASUREMENT OF $B^+ \rightarrow K^+ \mu^+ \mu^-$: BRANCHING FRACTION



Theory predictions from [C. Bobeth et al., JHEP 07 (2011) 067] and [C. Bobeth et al., JHEP 01 (2012) 107]

- Determine branching fraction in 7 bins of q^2 , using $B^+ \rightarrow J/\psi K^+$ as a normalisation channel and accounting for differences in the efficiencies.
- $\mathcal{B} = (4.36 \pm 0.15 \pm 0.18) \times 10^{-7}$, taking excluded charmonium resonance regions into account.
- World's best measurement of $\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)$.

MEASUREMENT OF $B^+ \rightarrow K^+ \mu^+ \mu^-$: ANGULAR ANALYSIS

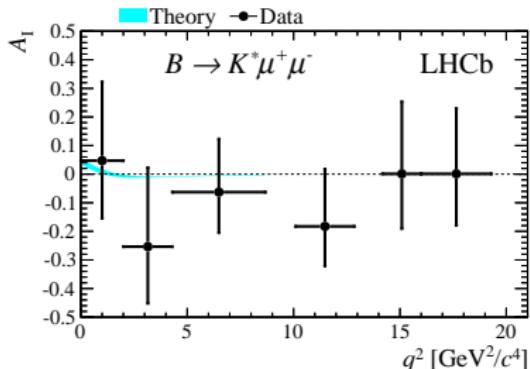
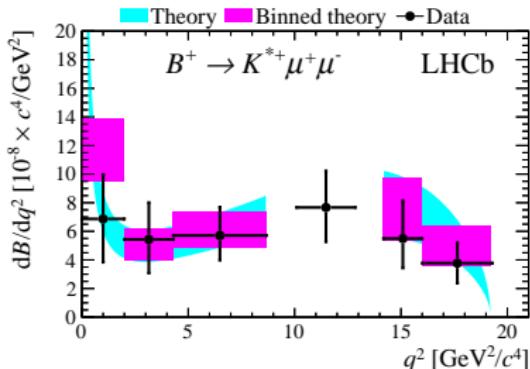


Theory predictions from [C. Bobeth et al., JHEP 07 (2011) 067] and [C. Bobeth et al., JHEP 01 (2012) 107]

- $\frac{1}{\Gamma} \frac{d\Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}{d \cos \theta_\ell} = \frac{3}{4}(1 - F_H)(1 - \cos^2 \theta_\ell) + \frac{1}{2}F_H + A_{FB} \cos \theta_\ell$
- Acceptance correction using simulation in q^2 and $\cos \theta_\ell$.
- Simultaneous fit to mass and $\cos \theta_\ell$. Background modeled with second-order polynomial in angles.



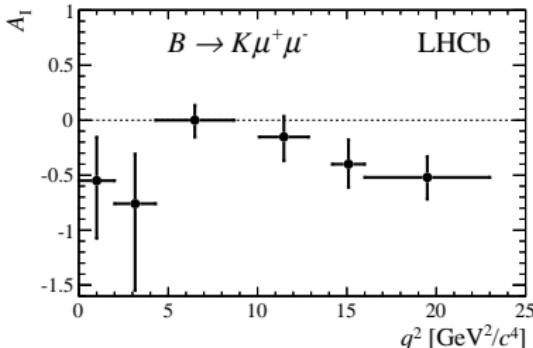
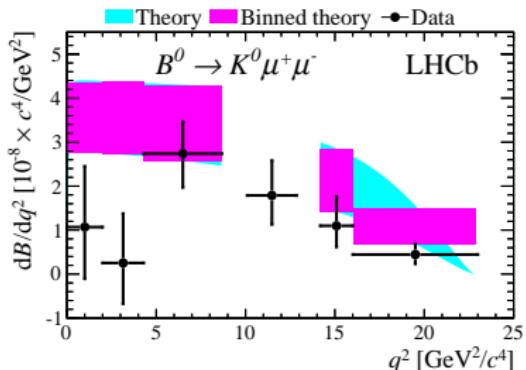
ISOSPIN ASYMMETRY (I)



Theory predictions from [C. Bobeth et al., JHEP 01 (2012) 107] and [M. Beneke et al., Nucl. Phys. B612 (2001) 25-58]

- Measure "Isospin asymmetry": $\frac{\Gamma(B^0 \rightarrow K^{*0} \mu^+ \mu^-) - \Gamma(B^+ \rightarrow K^{*+} \mu^+ \mu^-)}{\Gamma(B^0 \rightarrow K^{*0} \mu^+ \mu^-) + \Gamma(B^+ \rightarrow K^{*+} \mu^+ \mu^-)}$
- Predicted to be very small.
- Use $B \rightarrow J/\psi K^*$ as a normalisation channel.
- $B^+ \rightarrow K^{*+} \mu^+ \mu^-$ is challenging due to K_s^0 reconstruction from $K^{*+} \rightarrow K_s^0 \pi^+$.
- Results for $B \rightarrow K^* \mu^+ \mu^-$ agree well with prediction.

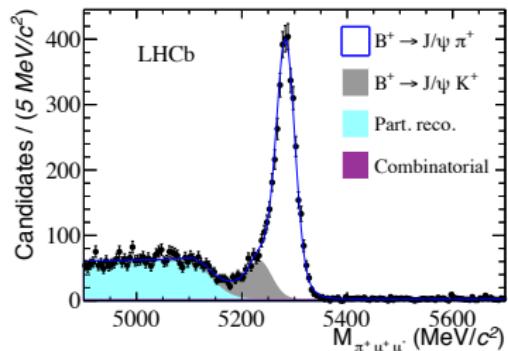
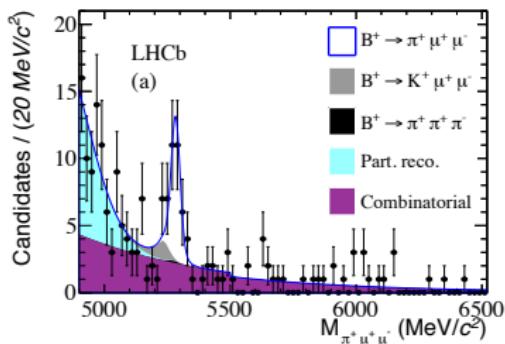
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- Measure "Isospin asymmetry": $\frac{\Gamma(B^0 \rightarrow K^0 \mu^+ \mu^-) - \Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\Gamma(B^0 \rightarrow K^0 \mu^+ \mu^-) + \Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}$
- Predicted to be very small.
- Use $B \rightarrow J/\psi K$ as a normalisation channel.
- $B^0 \rightarrow K^0 \mu^+ \mu^-$ is challenging due to K_s^0 reconstruction.
- A_I shows a 4.4σ deviation from 0, driven by low $\mathcal{B}(B^0 \rightarrow K^0 \mu^+ \mu^-)$.

BRANCHING FRACTION MEASUREMENT OF $B^+ \rightarrow \pi^+ \mu^+ \mu^-$



- Fit four distributions simultaneously: $B^+ \rightarrow J/\psi K^+$, misidentified $B^+ \rightarrow J/\psi K^+$, $B^+ \rightarrow K^+ \mu^+ \mu^-$, $B^+ \rightarrow \pi^+ \mu^+ \mu^-$
- Measured branching fraction: $\mathcal{B} = (2.3 \pm 0.6(\text{stat}) \pm 0.1(\text{sys})) \times 10^{-8}$
- In good agreement with SM expectation: $(1.96 \pm 0.21) \times 10^{-8}$.

[S. Hai-Zen et al., 2008 Commun. Theor. Phys. 50 696]

- Also determine: $R = \frac{\mathcal{B}(B^+ \rightarrow \pi^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)} = f^2 \frac{|V_{td}|^2}{|V_{ts}|^2}$
- Which leads to: $\frac{|V_{td}|}{|V_{ts}|} = 0.266 \pm 0.035 \text{ (stat)} \pm 0.007 \text{ (sys)}$

SUMMARY

- EW penguins are a very active area in LHCb.
- $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ is a "golden-channel" and allows measurement of many observables: angular analysis, zero-crossing point of A_{FB} , differential branching fraction, Isospin asymmetry, \mathcal{CP} asymmetry.
- $B^+ \rightarrow K^+ \mu^+ \mu^-$: Angular analysis, differential branching fraction and Isospin asymmetry.
- Discovery of $B^+ \rightarrow \pi^+ \mu^+ \mu^-$, rarest B decay ever observed.
- 2011+2012 ($\approx 3.2 \text{ fb}^{-1}$) data will allow more precision, more observables and (hopefully) conclusions on discrepancies.

Backupz

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$: FULL ANGULAR DISTRIBUTION

- If we neglect lepton masses and S-wave component, the angular distribution of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ is:

$$\frac{d^4\Gamma}{d\cos\theta_\ell\,d\cos\theta_K\,d\phi\,dq^2} \propto I_1^s \sin^2\theta_K + I_1^c \cos^2\theta_K \\ + (I_2^s \sin^2\theta_K + I_2^c \cos^2\theta_K) \cos 2\theta_\ell \\ + I_3 \sin^2\theta_K \sin^2\theta_\ell \cos 2\phi \\ + I_4 \sin 2\theta_K \sin 2\theta_\ell \cos\phi \\ + I_5 \sin 2\theta_K \sin\theta_\ell \cos\phi \\ + (I_6^s \sin^2\theta_K + I_6^c \cos^2\theta_K) \cos\theta_\ell \\ + I_7 \sin 2\theta_K \sin\theta_\ell \sin\phi \\ + I_8 \sin 2\theta_K \sin 2\theta_\ell \sin\phi \\ + I_9 \sin^2\theta_K \sin^2\theta_\ell \sin 2\phi$$

- The I_i depend on q^2 and contain the (transversity) amplitudes of the K^{*0} .
- All I_i are observables, 8 of them are independent.