

Direct CP-violation in nonleptonic charm decays: New Physics

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Outline

- Significance of CPV in Charm within and beyond SM
 - Quantify (parametrize) theory expectations of direct CPV in charm decays
- Δa_{CP} implications for weak scale NP
 - EFT & models

Why CP Violation in Charm?

- CPV in charm provides a unique probe of New Physics (NP)
 - sensitive to NP in the up sector
 - SM charm physics is CP conserving to first approximation (2 generation dominance, no hard GIM breaking)
- Common lore "any signal for CPV would be NP":
 - In the SM, CPV in mixing enters at $\mathcal{O}(V_{cb}V_{ub}/V_{cs}V_{us}) \sim 10^{-3}$
 - In the SM, direct CPV enters at $\mathcal{O}([V_{cb}V_{ub}/V_{cs}V_{us}] \alpha_s/\pi) \sim 10^{-4}$ (in singly Cabibbo suppressed decays)

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- CPV in decays (direct CPV)
 - Time-integrated CPV decay asymmetries to CP eigenstates

$$a_f \equiv \frac{\Gamma(D^0 \to f) - \Gamma(\bar{D}^0 \to f)}{\Gamma(D^0 \to f) + \Gamma(\bar{D}^0 \to f)}$$

• Focus on K^+K^- and $\pi^+\pi^-$ final states: $\Delta a_{CP} \equiv a_{K^+K^-} - a_{\pi^+\pi^-}$



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 $\Delta a_{CP}^{\text{World}} = -(0.68 \pm 0.15)\% \quad (\sim 4.3\sigma \text{ from 0})$

see talks by Ko, Parkes & Mattson

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...beyond natural expectation within the SM

Grossman et al., hep-ph/0609178 Cheng & Chiang, 1201.0785

Franco, Mishima & Silvestrini, 1203.3131 Li, Lu & Yu, 1203.3120

 but not possible to prove from first principles, and SM-like explanation Cannot be excluded
 Golden & Grinstein Phys. Lett. B 222 (1989)

Brod, Kagan & Zupan 1111.5000 Brod, Grossman, Kagan & Zupan 1203.6659 Feldmann, Nandi & Soni, 1202.3795 Bhattacharya, Gronau & Rosner, 1201.2351

see talks by Bhattacharya & Brod

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in terms of weak decay amplitudes

$$\lambda_q \equiv V_{cq}^* V_{uq} \,, \quad \lambda_d + \lambda_s + \lambda_b = 0$$

$$A_K = \lambda_s A_K^{(1)} + \lambda_b A_K^{(2)}$$
$$A_\pi = \lambda_d A_\pi^{(1)} + \lambda_b A_\pi^{(2)}$$

• naive expectation $A_f^{(2)}/A_f^{(1)} \sim \frac{\alpha_s(m_c)}{\pi}$ (assuming $m_c >> \Lambda_{QCD}$)

• SM expectations for $a_f^{\text{dir.}} \simeq 2r_f \sin \phi_f \sin \delta_f$ $f = K, \pi$ absolute ratio of interf. CPV & strong phase differences decay amplitudes • in terms of weak decay amplitudes $\lambda_q \equiv V_{cq}^* V_{uq}$, $\lambda_d + \lambda_s + \lambda_b = 0$ $A_K = \lambda_s A_K^{(1)} + \lambda_b A_K^{(2)}$ $r_f \propto \xi = |\lambda_b / \lambda_s| \simeq |\lambda_b / \lambda_d| \approx 0.0007$ $A_{\pi} = \lambda_d A_{\pi}^{(1)} + \lambda_b A_{\pi}^{(2)}$ • naive expectation $A_f^{(2)}/A_f^{(1)} \sim \frac{\alpha_s(m_c)}{\pi}$ $\phi_K^{SM} = \arg(\lambda_b/\lambda_s)$ $\approx -\arg(\lambda_b/\lambda_d) = -\phi_{\pi}^{SM}$ (assuming $m_c >> \Lambda_{OCD}$) $\approx 70^{\circ}$

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$$\Delta a_{CP} \approx (0.13\%) \text{Im}(\Delta R^{\text{SM}}), \qquad \Delta R^{\text{SM}} \equiv \frac{A_K^{(2)}}{A_K^{(1)}} + \frac{A_\pi^{(2)}}{A_\pi^{(1)}}$$





O(4-6) values of $|\Delta R^{SM}|$ needed

- Assume SM does not saturate the experimental value
- Parametrize NP contributions in EFT normalized to the effective SM scale

$$\mathcal{H}_{|\Delta c|=1}^{\text{eff}-\text{NP}} = \frac{G_F}{\sqrt{2}} \sum_i C_i^{\text{NP}} Q_i$$

• most general dim 6 Hamiltonian at $\mu < m_{W,t}$

$$Q_{1}^{q} = (\bar{u}q)_{V-A} (\bar{q}c)_{V-A}$$

$$Q_{2}^{q} = (\bar{u}_{\alpha}q_{\beta})_{V-A} (\bar{q}_{\beta}c_{\alpha})_{V-A},$$

$$Q_{5}^{q} = (\bar{u}c)_{V-A} (\bar{q}q)_{V+A},$$

$$Q_{6}^{q} = (\bar{u}_{\alpha}c_{\beta})_{V-A} (\bar{q}_{\beta}q_{\alpha})_{V+A},$$

$$Q_{7} = -\frac{e}{8\pi^{2}} m_{c} \bar{u}\sigma_{\mu\nu} (1+\gamma_{5}) F^{\mu\nu} c,$$

$$Q_{8} = -\frac{g_{s}}{8\pi^{2}} m_{c} \bar{u}\sigma_{\mu\nu} (1+\gamma_{5}) T^{a} G_{a}^{\mu\nu} c,$$

$$+ \text{Ops. with V} \leftrightarrow \text{A}$$

x 5 q \overline{q} flavor structures

- Assume SM does not saturate the experimental value
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$$\Delta a_{CP} \approx (0.13\%) \operatorname{Im}(\Delta R^{\mathrm{SM}}) + 9 \sum_{i} \operatorname{Im}(C_{i}^{\mathrm{NP}}) \operatorname{Im}(\Delta R_{i}^{\mathrm{NP}}) \qquad \Delta R_{i}^{\mathrm{NP}} \equiv \frac{G_{F}}{\sqrt{2}} \sum_{f=\pi,K} \frac{\langle Q_{i} \rangle_{f}}{A_{f}^{(1)}}$$

10

 \sim

• for
$$\operatorname{Im}(C_i^{\operatorname{NP}}) = \frac{v^2}{\Lambda^2}$$
 : $\frac{(10 \text{ TeV})^2}{\Lambda^2} = \frac{(0.61 \pm 0.17) - 0.12 \operatorname{Im}(\Delta R^{\operatorname{SM}})}{\operatorname{Im}(\Delta R^{\operatorname{NP}})}$

Are such contributions allowed by other flavor constraints?

• In EFT can be estimated via "weak mixing" of operators

- Important constraints expected from D-D mixing and direct CPV in $K^0 \rightarrow \pi^+ \pi^-$ (ϵ'/ϵ)
- Quadratic NP contributions
 - either chirally suppressed...
 - ... or highly UV sensitive



Isidori, J.F.K, Ligeti & Perez

1111.4987



- SM quark flavor symmetry $\ \mathcal{G}_F = SU(3)_Q imes SU(3)_U imes SU(3)_D$
 - two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^{\dagger})_{t\!\!/r}$, $\mathcal{A}_d \equiv (Y_d Y_d^{\dagger})_{t\!\!/r}$

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 - in the 2-gen limit single source of CPV: $J\equiv i[{\cal A}_u,\,{\cal A}_d]$ Gedalia, Mannelli & Perez 1002.0778, 1003.3869
 - invariant under SO(2) rotations between up-down mass bases



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- SU(2)_Q breaking NP $\mathcal{O}_L = \left[(X_L)^{ij} \overline{Q}_i \gamma^{\mu} Q_j \right] L_{\mu}$

$$\operatorname{Im}(X_L^u)_{12} = \operatorname{Im}(X_L^d)_{12} \propto \operatorname{Tr}(X_L \cdot J) .$$



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 - SM 3-gen case characterized by SU(3)/SU(2) breaking pattern by Y_{b,t} Kagan et al., 0903.1794
 - 3-gen X_L can be decomposed under SU(2)_Q, constrained separately (barring cancelations)
 - SM breaking of residual SU(2)_Q suppressed by m_c/m_t , m_s/m_b , θ_{13} , θ_{23} (charm and kaon sectors dominated by 2-gen physics)

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 - Implication (1): direct correspondence between Δa_{CP} and ε'/ε (no weak loop suppression)
 - constraint on SU(3)_Q breaking NP contributions

Gedalia, J.F.K, Ligeti & Perez 1202.5038

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- Implication (2): bounds on degeneracy in SUSY alignment models
 - CPV constraints less stringent than previously thought see talk by G. Perez



• In EFT can be estimated via "weak mixing" of operators



- Important constraints expected from D- \overline{D} mixing and direct CPV in $K^0 \rightarrow \pi^+ \pi^-$ (ϵ'/ϵ)
 - LL 4q operators: excluded
 - LR 4q operators: ajar potentially visible effects in D- \overline{D} and/or ϵ'/ϵ

Model example: Hochberg, Nir, 1112.5268

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• **RR 4q operators:** unconstrained in EFT - UV sensitive contributions? Model example: Da Rold et al., 1208.1499

Dipole operators only weakly constrained (edm's)

Δacp in (enter favorite NP model name)

Before LHCb result,

DCPV in charm not on top of NP theorists expectations

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In last year, situation has improved considerably

Δa_{CP} in <u>SUSY Models</u>

• Left-right up-type squark mixing contributions

 $\left|\Delta a_{CP}^{\rm SUSY}\right|\approx 0.6\% \left(\frac{\left|{\rm Im}\,(\delta_{12}^u)_{LR}\right|}{10^{-3}}\right) \left(\frac{{\rm TeV}}{\tilde{m}}\right)$

- contributions to $\Delta F=2$ helicity suppressed
- requires large trilinear (A) terms, non-trivial flavor in UV

$$\operatorname{Im} \left(\delta_{12}^{u}\right)_{LR} \approx \frac{\operatorname{Im}(A) \ \theta_{12} \ m_{c}}{\tilde{m}} \approx \left(\frac{\operatorname{Im}(A)}{3}\right) \left(\frac{\theta_{12}}{0.3}\right) \left(\frac{\operatorname{TeV}}{\tilde{m}}\right) 0.5 \times 10^{-3}$$

Grossman, Kagan & Nir, hep-ph/0609178 Giudice, Isidori & Paradisi, 1201.6204 Hiller, Hochberg, Nir, 1204.1046

 \widetilde{g} \widetilde{u} $\widetilde{\delta}_{12}^{u})_{LR}$

Giudice, Isidori & Paradisi, 1201.6204 see also Keren-Zur et al., 1205.5803

Δa_{CP} in <u>Warped Extra-Dim. Models</u>

• Anarchic flavor with bulk Higgs

$$\begin{split} \left| \Delta a_{CP}^{\text{chromo}} \right|_{\text{RS}} &\simeq 0.6\% \times \left(\frac{\mathcal{O}_{\beta}}{0.1} \right) \left(\frac{Y_5}{6} \right) \left(\frac{3 \text{ TeV}}{m_{\text{KK}}} \right) \\ & \text{Higgs wavefunction overlap} \end{split}$$
• requires very large 5D Yukawas
• helps to avoid D-D mixing constraints
Gedalia et al., 0906.1879
$$\int \frac{\mathcal{O}_{\text{KK}}}{\mathcal{O}_{\text{KK}}} \propto \frac{1}{Y_5^2} \\ \text{• implies low UV cut-off} \quad \frac{1}{2} \lesssim Y_5 \lesssim \frac{4\pi}{\sqrt{N_{KK}}} \end{split}$$

Delaunay, J.F.K., Perez & Randall 1207.0474



Agashe, Azatov & Zhu, 0810.1016 Csaki et al., 0907.0474

• Can be mapped to 4D partial compositeness models (for localized Higgs $\mathcal{O}_{\beta} \sim 1$)

 (α) $(\mathbf{x})^2$ $(\mathbf{a} - \mathbf{x})^2$

Keren-Zur et al., 1205.5803

Δa_{CP} and <u>4th Generation</u>

• 3-gen CKM non-unitarity and b' penguins

Feldmann, Nandi & Soni 1202.3795

$$\Delta a_{CP} \propto 4 \operatorname{Im} \left[\frac{\lambda_{b'}}{\lambda_d - \lambda_s} \right] \simeq \frac{2 \sin \theta_{14} \sin \theta_{24} \sin(\delta_{14} - \delta_{24})}{\sin \theta_{12}}$$

No parametric enhancement allowed due to existing ΔF=2 CPV bounds

Nandi & Soni, 1011.6091 Buras et al., 1002.2126

- Effects comparable to SM still allowed
- Similar conclusions for generic mixing with vector-like quarks

Grossman, Kagan & Nir hep-ph/0609178 Altmannshofer et al. 1202.2866

Generic Implications for Experiment

• correlations with EDM's, rare top & down-type quark processes Giudice, Iside Very model dependent

Giudice, Isidori & Paradisi, 1201.6204 Hochberg & Nir, 1112.5268 Altmannshofer et al., 1202.2866

- NP explanations of Δa_{CP} via <u>chromo-magnetic dipole operators</u>
 - generically predict EM dipoles rare radiative charm decays

$$D^0 \rightarrow X\gamma$$
 $D^0 \rightarrow Xe^+e^-$

Delaunay, J.F.K., Perez & Randall 1207.0474

Expected NP rates few orders below SM LD contributions

possibility to access CPV observables

Isidori & J.F.K., 1205.3164 Fajfer & Kosnik, 1208.0759 Cappiello, Cata & D' Ambrosio, 1209.4235

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- NP explanations of Δa_{CP} via $\Delta I=3/2$ contributions
 - SM contributions to $A_f^{(2)}$ purely $\Delta I=1/2$

Grossman, Kagan & Zupan, 1204.3557

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No CPV expected in pure \Delta I = 3/2 decays
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see also later talk by Grossman

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 - larger than naive SM expectations...however, SM explanation cannot be excluded from first principles
- If NP, points towards new flavor structures in uR sector at the TeV scale
- More experimental observables could clarify the picture Grossman, Kagan & Zupan, 1204.3557 Isidori & J.F.K., 1205.3164
 - (CPV in) rare radiative charm decays sensitive to NP in dipole ops.
 - CPV in isospin related 2-, 3-body modes can test $\Delta I=3/2$ NP

Backup

• CPV in Mixing $|D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle$

$$m \equiv \frac{m_1 + m_2}{2}, \qquad \qquad \Gamma \equiv \frac{\Gamma_1 + \Gamma_2}{2},$$
$$x \equiv \frac{m_2 - m_1}{\Gamma}, \qquad \qquad y \equiv \frac{\Gamma_2 - \Gamma_1}{2\Gamma}.$$

• Experimentally accessible mixing quantities:

• x,y (CP conserving) Cannot be estimated accurately within SM NP contributions are predictable

flavor specific time-dependent CPV decay asymmetries [sensitive to q/p]

$$a_f(t) \equiv \frac{\Gamma(D^0(t) \to f) - \Gamma(\bar{D}^0(t) \to f)}{\Gamma(D^0(t) \to f) + \Gamma(\bar{D}^0(t) \to f)},$$

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• CPV in Mixing

Isidori, Nir & Perez 1002.0900

	Bounds on Λ (TeV)		Bounds on c_{ij} ($\Lambda = 1$ TeV)		
Operator	Re	Im	Re	Im	Observables
$(\bar{s}_L \gamma^\mu d_L)^2$	9.8×10^2	1.6×10^{4}	9.0×10^{-7}	3.4×10^{-9}	$\Delta m_K; \varepsilon_K$
$(\bar{s}_R d_L)(\bar{s}_L d_R)$	1.8×10^{4}	3.2×10^5	6.9×10^{-9}	2.6×10^{-11}	$\Delta m_K; \varepsilon_K$
$(\bar{c}_L \gamma^\mu u_L)^2$	1.2×10^{3}	2.9×10^{3}	5.6×10^{-7}	1.0×10^{-7}	$\Delta m_D; q/p , \phi_D$
$(\bar{c}_R u_L)(\bar{c}_L u_R)$	6.2×10^{3}	1.5×10^{4}	5.7×10^{-8}	1.1×10^{-8}	$\Delta m_D; q/p , \phi_D$
$(\bar{b}_L \gamma^\mu d_L)^2$	5.1×10^2	9.3×10^2	3.3×10^{-6}	1.0×10^{-6}	$\Delta m_{B_d}; S_{B_d \to \psi K}$
$(\bar{b}_R d_L)(\bar{b}_L d_R)$	1.9×10^{3}	3.6×10^{3}	5.6×10^{-7}	1.7×10^{-7}	$\Delta m_{B_d}; S_{B_d \to \psi K}$
$(\bar{b}_L \gamma^{\mu} s_L)^2$	1.1×10^{2}	1.1×10^{2}	7.6×10^{-5}	7.6×10^{-5}	Δm_{B_s}
$(\bar{b}_R s_L)(\bar{b}_L s_R)$	3.7×10^2	3.7×10^2	1.3×10^{-5}	1.3×10^{-5}	Δm_{B_s}

$$x, y \sim 1\%$$



Imply significant constraints on CPV NP contributions, second only to kaon sector

- SM quark flavor symmetry $\ \mathcal{G}_F = SU(3)_Q imes SU(3)_U imes SU(3)_D$
 - two sources of breaking: $\mathcal{A}_u \equiv (Y_u Y_u^{\dagger})_{t\!\!/r}, \qquad \mathcal{A}_d \equiv (Y_d Y_d^{\dagger})_{t\!\!/r}$
 - Implication (1): direct correspondence between Δa_{CP} and ε'/ε (no weak loop suppression)
 - constraint on SU(3)_Q breaking NP: $\Delta a_{CP}^{NP} \lesssim 4 \times 10^{-4}$ Gedalia, J.F.K, Ligeti & Perez 1202,5038
 - Similarly for rare semileptonic decays:

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 - Implication (2): bounds on degeneracy in SUSY alignment models
 - SUSY effects in flavor ~ masses, splittings (degeneracy); mixing angles (e.g. $\bar{q}_i \tilde{q}_j \tilde{g}$)

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 - D & K mixing said to imply that alignment not viable w/o degeneracy (Nir & Seiberg, hep-ph/9304307) Blum et al., 0903.2118
 - based on assumption of ~ maximal CPV in K & D mixing
 - not actually attainable in alignment due to CPV universality

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 - Implication (2): bounds on degeneracy in SUSY alignment models
 - viable SUSY spectra can be generated from complete anarchy at moderate mediation scales (*M_S*) (SUSY QCD RGE)



On Universality of CPV in SU(3) breaking [NP)

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- SM quark flavor symmetry $\ \mathcal{G}_F = SU(3)_Q imes SU(3)_U imes SU(3)_D$
 - two sources of breaking: $\mathcal{A}_{u}^{30} = (Y_{u}Y_{u}^{\dagger})_{tr}$, $\mathcal{A}_{d} \equiv (Y_{d}Y_{d}^{\dagger})_{tr}$ • Implication (2): bounds on degeneracy in SUSY alignment models = 0.10• viable SUSY spectra can be generated from complete anarchy at moderate mediation scales = 0.10 = 0
 - surprising mass hierarchies still viable, e.g. $L_{\alpha}^{\log_{\lambda} \alpha}$

$$m_{\tilde{g}}=1.3\,\mathrm{TeV},\,m_{\tilde{Q}_1}=550\,\mathrm{GeV},\,m_{\tilde{Q}_2}=950\,\mathrm{GeV}$$

Important implications for LHC searches



Generic Implications for Experiment

• NP explanations of Δa_{CP} via <u>chromo-magnetic dipole operators</u>

 $|\Delta a_{CP}^{\rm NP}| \approx -1.8 |{\rm Im}[C_8^{\rm NP}(m_c)]| ,$

Grossman, Kagan & Nir, hep-ph/0609178 Giudice, Isidori & Paradisi, 1201.6204 (estimate of matrix element in QCD fact.)

• generically predict EM dipoles

 $|\text{Im}[C_7^{\text{NP}}(m_c)]| \approx |\text{Im}[C_8^{\text{NP}}(m_c)]| \approx 0.4 \times 10^{-2}$. (QCD RGE evolution with TeV NP)

Isidori & J.F.K., 1205.3164

- possibility to access CPV observables in $D^0 \rightarrow \pi \pi \gamma$, $KK\gamma$
 - in SM CPV expected similar as in $D^0 \rightarrow \pi \pi$, KK
 - large strong phases natural for LD SM contributions

$$|a_{(\rho,\omega)\gamma}|^{\max} = 0.04(1) \left| \frac{\operatorname{Im}[C_7(m_c)]}{0.4 \times 10^{-2}} \right| \left[\frac{10^{-5}}{\mathcal{B}(D \to (\rho,\omega)\gamma)} \right]^{1/2} \lesssim 10\%$$

(smaller effects also in $D^0 \rightarrow KK\gamma$ with m_{KK} around Φ mass)

Generic Implications for Experiment

• NP explanations of Δa_{CP} via $\Delta I=3/2$ contributions

Grossman, Kagan & Zupan, 1204.3557

• SM contributions to $A_{K}^{(d)}$, $A_{\pi}^{(s)}$ purely $\Delta I=1/2$

No CPV expected in pure $\Delta I = 3/2$ decays

 $\Gamma(D^+ \to \pi^+ \pi^0) - \Gamma(D^- \to \pi^- \pi^0) = 0 \qquad \text{(up to small isospin breaking)}$

- nonzero difference would point towards CPV $\Delta I=3/2$ NP contributions
- decay amplitude sum-rules even in presence isospin breaking

$$\frac{1}{\sqrt{2}} |A_{\pi^+\pi^-} - \bar{A}_{\pi^-\pi^+}| \neq |A_{\pi^0\pi^0} - \bar{A}_{\pi^0\pi^0}|, \quad \Longrightarrow \quad \text{signal of } \Delta I=3/2 \text{ CPV NP}$$

• experimentally accessible with time-dependent measurements (also Dalitz plot analyses in $D \rightarrow 3\pi$, $D \rightarrow KK\pi$)