Front-End electronics for particle detectors

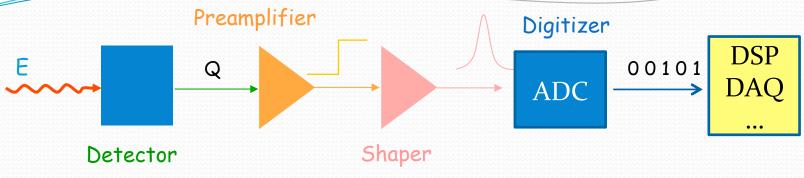
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Outline

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- Detectors and signals
- Noise basic principles
- Radiation damage
- Front-end schemes
- Charge-sensitive amplifier
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Introduction

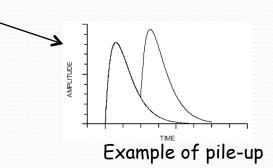


General readout architecture

- The particle deposits energy in a detecting medium Solice
 Liqui
- Energy is converted into an electrical signal: Q = KE
- The charge Q is typically small and must be amplified, in order to be measured and processed
- The preamplifier converts Q into a voltage
- The shaper provides gain and shape, according to the application and trying to optimize S/N
- The Digitizer converts the "analog" information into sequence of bits, for storage and processing

Purpose of Front-End Electronics

- 1. Acquire an electrical signal from the detector
- 2. Choose the gain and shaping time in order to optimize:
 - minimum detectable signal over the noise (maximize S/N)
 - energy measurements (linearity ...);
 - event rate (pile-up, ballistic deficit, ...);
 - time of arrival (time-walk, jitter ...);
 - radiation hardness/tolerance;
 - power consumption;
 - cost



Often the requirements are in conflict each other



The final design comes out as a compromise, according to the specific application:

- Triggering (focus on timing)
- Tracking (focus on minimum detect. signal)
- Energy measurement (focus on linearity, dynamic range ...)

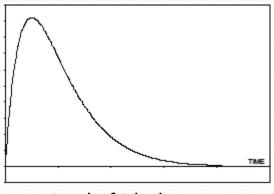
NOISE BASIC PRINCIPLES

Noise

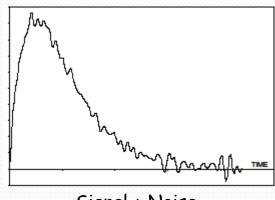
The precision of amplitude and timing measurements is limited by the NOISE

Definition

Noise is every undesirable signal superimposed to our signal of interest \rightarrow fluctuations on amplitude and time measurement



Signal of ideal system



Signal + Noise

1. External noise (interference)

It is generated by external sources (RF, ripple of power lines, ground loops ...)

<u>Can be eliminated</u> by proper shielding, cabling ...

2. Intrinsic noise

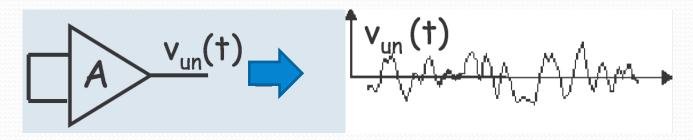
It is a property of detector and/or electronics

<u>Can be reduced</u> by proper design of front-end electronics

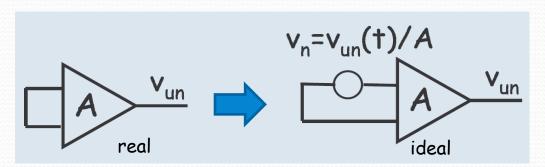
Intrinsic noise

The output voltage of a <u>real amplifier</u> is never constant, even if $V_{in} = 0$

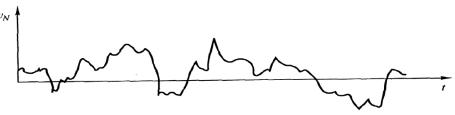
The fluctuations of $V_{un}(t)$ when $V_{in} = 0$ correspond to the <u>noise</u> of amplifier



The noise of a <u>real amplifier</u> can be attributed to a noise voltage source in input to an <u>ideal amplifier</u> (noiseless)



Intrinsic noise







 V_n has mean value = 0, but power \neq 0



We can define:

• Source of voltage noise:
$$v_n = \sqrt{v_n^2}(f)$$

• Source of current noise: $i_n = \sqrt{i_n^2(f)}$

$$i_n = \sqrt{i_n^2}(f)$$

A noise source is usually defined by its POWER SPECTRAL DENSITY: noise power per unit of bandwidth

If Power Spectral Density is constant → White Noise

Basic noise mechanisms

The fluctuation of the current is given by:
$$< di>^2 = (\frac{ne}{l} < dv>)^2 + (\frac{ev}{l} < dn>)^2$$

There are two basic mechanism contributing to noise:

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Velocity fluctuations → Thermal noise

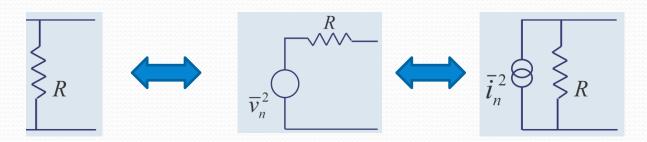
Number fluctuations  

Excess (or flicker, or "1/f") noise
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1. Thermal noise (Johnson noise)

It is typical of resistors

- Caused by the random thermal motion of charge carriers (electrons)
- Does not depends on a DC current



A real (noisy) resistor is equivalent to an ideal (noiseless) resistor + noise source (voltage or current)

Power spectral density:

$$S_{v}(f) = \frac{dv_{n}^{2}}{df} = 4kTR$$

$$S_i(f) = \frac{di_n^2}{df} = \frac{4kT}{R}$$

k = Boltzmann constant = $1.3806503 \times 10^{-23} \text{ J/K}$

T = absolute temperature

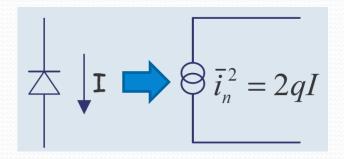
R = resistance



Does not depend on $f \rightarrow$ Thermal noise is a white noise

2. Shot noise

It is caused by fluctuations in the number of charge carriers, for example in the current flowing in a semiconductor diode of transistor, where e/h cross a potential barrier



$$S_i(f) = \frac{d\overline{i}_n^2}{df} = 2qI$$

Power spectral density: $S_i(f) = \frac{d\tilde{i}_n^2}{df} = 2qI$ does not depend on $f \to also$ shot noise is white (but a current I must be present)

Example: consider a reversed-biased diode, with leakage I = 1 nA

$$S_i(f) = \frac{d^{-\frac{2}{n}}}{df} = 2*1.6*10^{-19}*10^{-9} = 3.2*10^{-28}A^2/Hz$$



3. Flicker noise (1/f noise)

It is associated to random trapping and recombination of charge carriers in the semiconductors, typically caused by imperfections in the interface regions. It is also present in carbon resistors

Power spectral density:

$$S_{v}(f) = \frac{d\overline{v_n}^{-2}}{df} = K_f \frac{I^a}{f^b}$$

I is dc current $S_{v}(f) = \frac{d\overline{v_n}^2}{df} = K_f \frac{I^a}{f^b}$ $\begin{cases} K_f \text{ is a constant (vary from device to device)} \\ a \sim 0.5 \div 2 \\ b \sim 1 \end{cases}$ $S_{v}(f)$

It depends on f and clearly it is important at low frequencies

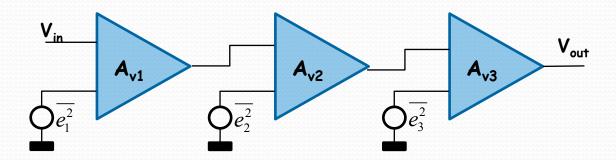
4. Burst noise (POPCORN noise)

Another low-frequency noise. It can be found in some integrated circuits and discrete transistor and is associated to contamination by ions of heavy metals (i.e. Au).

$$\frac{d\overline{i_b^2}}{df} = K_b \frac{I_b^c}{1 + (f/f_c)^2}$$

Log(f)

Intrinsic noise: importance of first stage



$$\begin{cases}
V_{out} = A_{v1} * A_{v2} * A_{v3} * V_{in} \\
\overline{e_{out}^2} = A_{v1}^2 * A_{v2}^2 * A_{v3}^2 * \overline{e_1^2} + A_{v2}^2 * A_{v3}^2 * \overline{e_2^2} + A_{v3}^2 * \overline{e_3^2}
\end{cases}$$



$$\left(\frac{Noise}{Signal}\right)^{2} = \left(\frac{\overline{e_{out}^{2}}}{V_{out}^{2}}\right) = \frac{\overline{e_{1}^{2}} + \frac{\overline{e_{2}^{2}}}{A_{V1}^{2}} + \frac{\overline{e_{3}^{2}}}{A_{V1}^{2} * A_{V2}^{2}}}{V_{in}^{2}}$$



- 1. It is necessary to decrease as much as possible the noise contribution e_1^2 of the first stage
- 2. It is necessary to increase the gain A_{v1} of the first stage because the noise contribution of next stages are divided by the gain of previous stages

Intrinsic noise: some practical rules

1. Uncorrelated noise sources must be added in quadrature

$$\overline{e_{tot}^2} = \overline{e_1^2} + \overline{e_2^2} + \overline{e_3^2} + \dots$$

2. In an amplifying chain, the noise generated in the first stage dominates

In first (and good) approximation, it is enough to evaluate (and decrease) the noise of the first stage

- 3. It is useful to represent a real (noisy) amplifier as an ideal (noiseless) amplifier with an equivalent noise source at its input: in this way the noise can be directly compared with input signal
- 4. In the case of particle detection systems, where the input is a charge Q, we use ENC: Equivalent Noise Charge: it is the signal magnitude which produces an output amplitude equal to rms noise

Representing the noise with ENC, we can directly compare the input charge with the noise introduced by our amplifier

The problem of radiation damage

The problem of radiation damage

When an electronic device is exposed to radiation, like in HEP experiments, there is a permanent or transient modification of the electrical properties of the active devices

- Fake signal
- Modification of memory content
- Degradation of performance
- Catastrophic failure
- Displacement damage: radiation (neutrons, protons, heavy ions...) change the arrangement of Si atoms in the crystal lattice \rightarrow the electronic characteristic are altered
- Ionization damage: charged particles produces transient currents (drift, diffusion) and entrapment of charge in SiO₂

Total dose (TID) \rightarrow Threshold shift, parasitic leakage currents, mobility degradation

Single Event Effects (SEE) \rightarrow temporary or permanent errors

Radiation-hard electronics: which technology?

- JFET (no gate oxide → no entrapment of charged carriers)
- Rad-hard CMOS (mainly used in the '90, very expensive)
- Deep submicron CMOS (scaling the oxide thickness the radimmunity is increased)
 - LHC (CMOS 0.25 μm)
 - LHC- Upgrade Phase1 (CMOS 0.13 μm)
 - HL-LHC (CMOS 0.065 μm)
- SiGe (RF bipolar transistors: high doping levels guarantees adequate current gain and reduce the sensitivity of the surface carrier concentration to radiation-induced charge in the oxide)

Monolithic technologies for custom applications

Bipolar

- > used mainly in the past
- > today few foundries
- > more speed and less power in analog applications
- > low integration

· BiCMOS

- > used mainly in the past
- > combines advantages of bipolar and CMOS
- > complex fabrication process

· Silicon on Insulator (SOI)

- > used (mainly in the past) for rad-hard applications
- > expensive

· GaAs

> not suitable for analog applications with high bandwidth

· SiGe

- > RF bipolar transistors: high speed
- > expensive
- > rad-hard

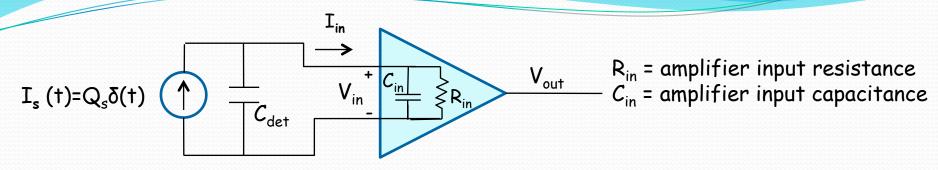
Most used technology is CMOS

CMOS technologies for Front-End electronics

- · Large part of Front-End electronics are developed in CMOS technology
- · Relatively "cheap" if recent/"old" techn. are used
- Using the "multiproject foundry runs", prototyping and small productions are very affordable
- \bullet Suitable to combine on the same chip analog section, digital part and $\mu\text{processors}$
- Very low power consumption
- The new deep submicron CMOS tech. (< 130 nm) are rad-hard and suitable for HL-LHC, ILC, Space applications

FRONT-END SCHEMES

Signal integration



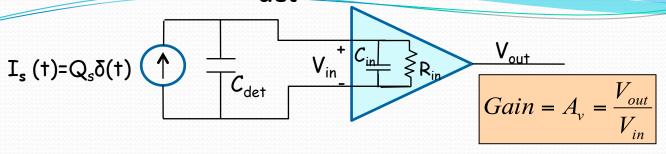
- The sensor signal is usually a short current pulse $\mathbf{I}_s(t) = \mathbf{Q} \cdot \delta(t)$ with duration ranging from few hundreds of ps, as in Si sensors and Resistive Plate Chambers to tens of μs , as in inorganic scintillators
- The physic quantity of interest is the deposited energy E, that is proportional to Q
- We must integrate I to have a measurement of E: $E \propto Q_S = \int I_S(t) dt$

WHERE to integrate?

<u>OPTIONS</u> (depending on charge collection time t_c and input time constant R_iC_t :

- 1. Detector capacitance $\rightarrow V_{in} \propto Q_s \rightarrow$ followed by voltage amplifier
- 2. Current sensitive amplifier $\rightarrow V_{out} \propto I_s \rightarrow$ followed by integrating Analog-to-Digital Converter
- 3. Charge sensitive amplifier $\rightarrow V_{out} \propto Q_s$

1. Integration on C_{det} (+ voltage amplifier)



If R_{in} is very big $\rightarrow \tau_{in} = R_{in}(C_{det} + C_{in})$ for discharging the sensor \rightarrow pulse duration (collection time)



the detector capacitance discharge slowly

$$I_{s}(t) \text{ is integrated on the total capacitance} \quad C_{t} = C_{\text{det}} + C_{\text{in}}$$

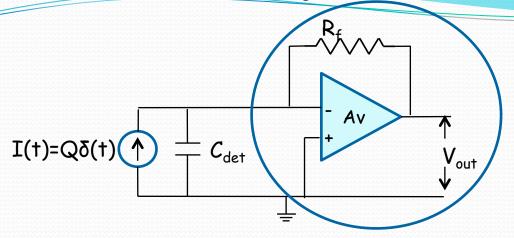
$$V_{in} = \frac{1}{C_{t}} \int I_{s} dt = \frac{Q_{s}}{C_{\text{det}} + C_{in}}$$

$$V_{out} = A_{v} \cdot V_{in} = A_{v} \cdot \frac{Q_{s}}{C_{\text{det}} + C_{in}}$$

In this method, V_{out} is proportional to Q_s , but it also depends on C_{det}

This is not desirable in the systems where C_{det} can vary: $\begin{cases} & \text{different strip length/width} \\ & \text{bias voltage} \end{cases}$

2. Current-sensitive amplifier

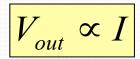


If R_{in} is small $\rightarrow \tau_{in} = R_{in}(C_{det} + C_{in})$ \leftarrow pulse duration (collection time)

The detector capacitance discharges rapidly \rightarrow the amplifier senses the current



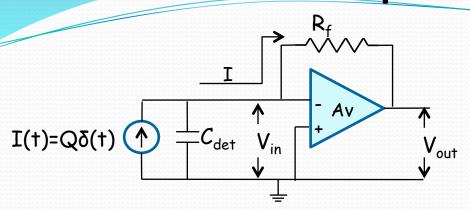
Using a transresistance amplifier (high gain operational amplifier with resistive feedback) we have:



In this method, V_{out} is proportional to I and does not depend on C_{det}

An integrating ADC can follow the amplifier and provide a digital word proportional to Q

Current-sensitive amplifier: the ideal case



Hypothesis:

- 1. input impedance of op-amp is ∞ (i.e. MOS gate)

 → all current flows in the feedback

 Vout

 2. A_v is very big

GAIN

$$\begin{cases} V_{\text{out}} = -A_{\text{v}}V_{\text{in}} \\ V_{\text{out}} - V_{\text{in}} = -R_{\text{f}}I \end{cases} \longrightarrow V_{out} + \frac{V_{out}}{A_{\text{v}}} = -R_{\text{f}}I \longrightarrow V_{out} = -R_{\text{f}}I \frac{A_{\text{v}}}{1 + A_{\text{v}}} \xrightarrow{A_{\text{v}}} 1 \longrightarrow V_{out} \approx -R_{\text{f}}I$$

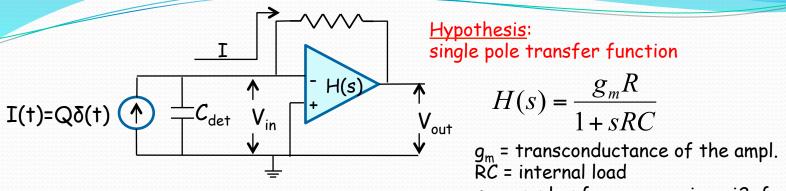
Main advantage: the shape of input signal is preserved → shaping stage not strictly necessary

INPUT IMPEDANCE

$$Z_{in} = \frac{V_{in}}{I} = \frac{R_f}{1 + A_v}$$

 $Z_{in} = \frac{V_{in}}{I} = \frac{R_f}{1+A}$ Since A_v is very big, $Z_{in} \to 0$ \longrightarrow Our hypothesis that τ_{in} « pulse duration is confirmed

Current-sensitive amplifier: the realistic case



Hypothesis:

single pole transfer function

$$H(s) = \frac{g_m R}{1 + sRC}$$

 $s = complex frequency = jw = j2\pi f$

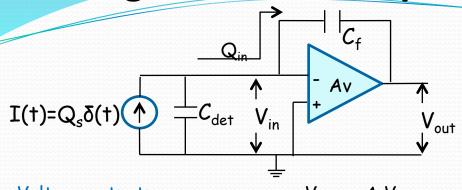
$$V_{out}(s) = -R_f I \cdot \frac{H(s)}{1 + H(s)} = -R_f I \cdot \frac{g_m R}{(1 + g_m R) + sRC}$$
At low frequency
$$V_{out}(s) = -R_f I \cdot \frac{g_m R}{1 + g_m R} \approx -R_f I$$

$$Z_{in}(s) = \frac{V_{in}}{I} = \frac{R_f}{1 + H(s)} = R_f \cdot \frac{1 + sRC}{(1 + g_m R) + sRC} \approx R_f \cdot \frac{s + \frac{1}{RC}}{s + \frac{g_m}{C}}$$

$$\Rightarrow \begin{cases} \text{Zero} = 1/RC \\ \text{Pole} = g_m/C \end{cases}$$

The presence of a zero in $Z_{in} \rightarrow$ inductive behavior \rightarrow possible instabilities (oscillations) for some values of C_{det}

3. Charge-sensitive amplifier (CSA): the ideal case



Hypothesis:

- 1. input impedance of op-amp is ∞ (i.e. MOS gate) → all current flows in the feedback
- 2. A_v is very big

Voltage output:

$$V_{out} = -A_v V_{in}$$

Voltage difference across C_f : $V_f = V_{in} - V_{out} = (A_v + 1)V_{in}$

$$V_f = V_{in} - V_{out} = (A_v + 1)V_{in}$$

Charge deposited on
$$C_f$$
: $Q_f = C_f V_f = C_f (A_v + 1) V_{in} = Q_{in}$ (for Hypothesis 1)

Effective input capacitance (seen by the sensor): $C_{in} = Q_{in}/V_{in} = C_f(A_v+1)$

GAIN (Charge Sensitivity):

$$CS = \frac{V_{out}}{Q_{in}} = -\frac{A_{v}V_{in}}{C_{f}(A_{v}+1)V_{in}} = -\frac{A_{v}}{C_{f}(A_{v}+1)} \approx -\frac{1}{C_{f}}$$
(A_v >> 1)

BUT ... not all the charge goes in the amplifier and is measured: a small fraction remains on C_{det} !!!

Charge transfer efficiency:

$$\frac{Q_{in}}{Q_{S}} = -\frac{Q_{in}}{Q_{det} + Q_{in}} = \frac{1}{1 + \frac{Q_{det}}{Q_{in}}} = \frac{1}{1 + \frac{C_{det}}{C_{in}}} \approx 1$$

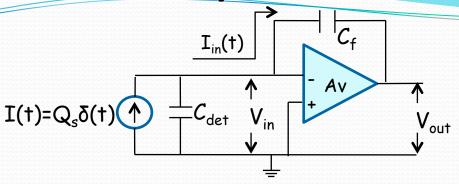
(if $C_{in} = C_f(A_v+1) \gg C_{det}$)

Example:
$$C_{det} = 10 \text{ pF}$$
 $A_v = 10^3 C_f = 1 \text{ pF} \rightarrow C_{in} = 1 \text{ nF}$ $Q_{in}/Q_s = 0.99$



$$Q_{in}/Q_{s} = 0.99$$

Charge-sensitive amplifier: the time response



In the frequency domain:

$$V_{out}(\omega) = -A_{v}V_{in}(\omega) \quad \text{(assuming Av constant and } \Rightarrow \infty)$$

$$V_{out}(\omega) - V_{in}(\omega) = -Z_{f}(\omega) \cdot I_{in}(\omega) = -\frac{I_{in}(\omega)}{j\omega C_{f}}$$

$$V_{out}(\omega) - V_{in}(\omega) = -\frac{I_{in}(\omega)}{j\omega C_{f}} \left[\frac{1}{1 + \frac{1}{A_{v}}} \right] \approx -\frac{I_{in}(\omega)}{j\omega C_{f}}$$

$$V_{out}(\omega) = -\frac{I_{in}(\omega)}{j\omega C_f} \left(\frac{1}{1 + \frac{1}{A_v}}\right) \approx -\frac{I_{in}(\omega)}{j\omega C_f}$$

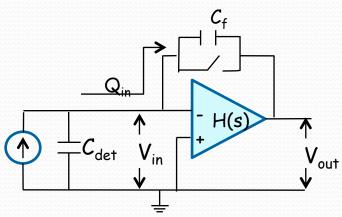
$$I_{\mathrm{in}}(t) = Q_{\mathrm{in}}\delta(t) \Rightarrow I_{\mathrm{in}}(\omega) = Q_{\mathrm{in}} \Rightarrow V_{out}(\omega) \approx -\frac{I_{in}}{j\omega C_f} \text{ in time domain}$$

$$V_{out}(t) \approx -\frac{1}{C_f} \int I_{in}\delta(t)dt = -\frac{Q_{in}}{C_f}u(t)$$

$$\downarrow V_{out}(t) \approx -\frac{1}{C_f} \int I_{in}\delta(t)dt = -\frac{Q_{in}}{C_f}u(t)$$
Step function

Charge-sensitive amplifier: the reset

Pulsed RESET

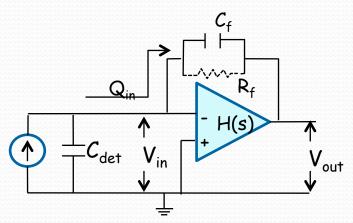


- \bullet The reset switch allows the removal of charge stored in $C_{\rm f}$
- The switch can be closed periodically or driven by some control signal

Drawbacks:

- · Dead time
- · Switch noise
- Leakage current

Continuous RESET

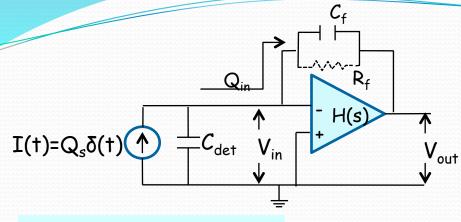


- The resistor $R_{\rm f}$ continuously discharges $\mathcal{C}_{\rm f}$ after the pulse
- Discharge time constant R_fC_f

Drawbacks:

- Additional parallel noise
- Long tail → Risk of pile-up

Charge-sensitive amplifier: the realistic case



- 1. Resistor R_f used to discharge C_f. Since it is a source of parallel noise (inject noise current into input noise), it must made very large to contribution to noise.
 - 2. Real amplifier (finite bandwidth and gain, C_1)

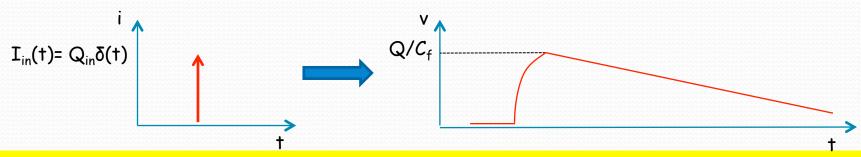
$$\frac{V_{out}(s)}{I(s)} = \frac{\frac{-g_m}{C_L C_T}}{\left(s + \frac{1}{R_f C_f}\right) \left(s + \frac{1}{R_i C_T}\right)}$$

$$2 \text{ poles}$$

$$\begin{cases} p_1 = \frac{1}{R_f C_f} \text{ (Low freq)} \\ p_2 = \frac{1}{R_i C_T} = \frac{\omega_0 C_f}{C_T} \text{ (High freq)} \end{cases}$$

$$\tau_1 = R_f C_f$$

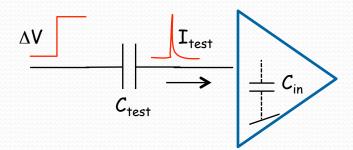
$$\tau_2 = R_i C_T = \frac{C_T}{\omega_0 C_f} \text{ Rise time constant}$$



- The fall time depends on the feedback: can be very large, since R_f must be very high for low noise (>> 1 $M\Omega$)
- The rise time depends on the input time constant, thus
 - Ri must be small to have short rise time
 - ω_0 : the amplifier GBW must be very large
 - $C_T \rightarrow C_d$: the rise time increase with detector capacitance

Calibration

A common technique used for calibration and test is to inject a voltage step through a capacitor



Ideally: $Q_{test} = C_{test} * \Delta V$

But we must consider the dynamic input capacitance C_{in} of preamp $\rightarrow C_{test}$ and C_{in} are in series

$$Q_{test} = C_{eq} * \Delta V = \frac{1}{\frac{1}{C_{test}} + \frac{1}{C_{in}}} * \Delta V = \frac{C_{test}}{1 + \frac{C_{test}}{C_{in}}} * \Delta V$$

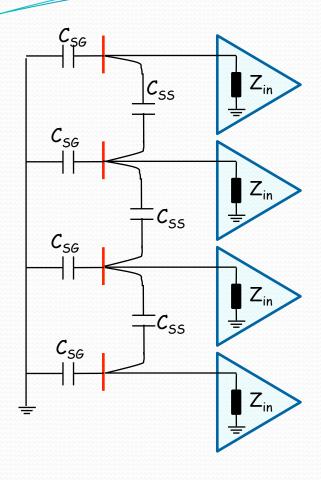
But usually $C_{\rm in} \gg C_{\rm test} \rightarrow C_{\rm test}/C_{\rm in} \sim 10^{3*}$

 $Q_{test} \sim C_{test} * \Delta V$

Eaxmple: $C_{test} = 1 \text{ pF}$, $\Delta V = 1 \text{ V} \rightarrow Q_{test} = 1 \text{ pC}$

*In fact, in case of CSA, $C_{in} = Q_{in}/V_{in} = C_f(A_v+1)$

Input impedance vs crosstalk



← Strip or pad

In strip or pixel detectors, where there are many adjacent channels, we must consider the following capacitive coupling:

- Strip or pad vs ground C_{SG}
- Inter-strip capacitance C_{SS}

If
$$Z_{in} \gg Z_{ss} = \frac{1}{\omega C_{SS}}$$

the charge induced on one strip is coupled into the adjacent channels through \mathcal{C}_{SS}

The number of affected signal depends on $\frac{C_{SS}}{C}$

If
$$Z_{in} \ll Z_{ss} = \frac{1}{\omega C_{SS}}$$

most part of the charge flows into the amplifier and only small part is coupled into the adjacent channels through C_{SS}

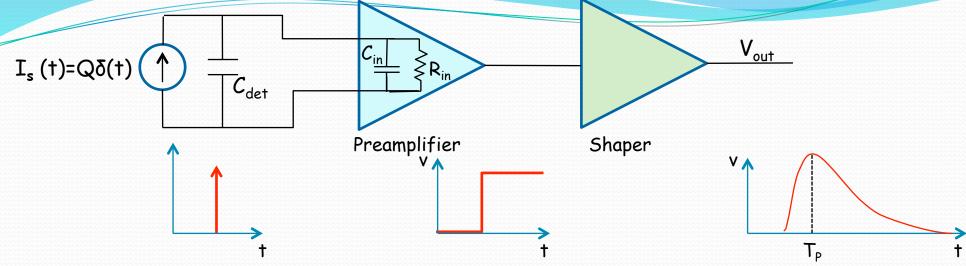
Summary:

low input impedance →

- Short rise time
- Small cross-talk

NOISE FILTERING: SHAPERS

Pulse shaping



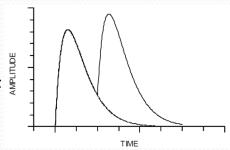
<u>Preamplifier = input amplifier</u> It is usually located close to detector and must have enough gain to make negligible the effects of induced noise. Typical example: Charge Sensitive Amplifier

Shaper Two conflicting objectives:

- 1. Improve the signal-to-noise ratio S/N, restricting the bandwidth (defining the peaking time T_P)
- 2. Tail the shape to improve the double-pulse resolution and avoid pile-up effect

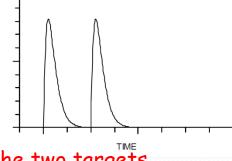
Slower pulse:

- · Less noise
- Pile-up (distortion of amplitude measurement)



Faster pulse:

- More noise
- Double-pulse resolution



The choice of the shaper (Tp, shape) derives from a compromise between the two targets

Noise through filters

$$\overline{v_n^2} - H(j\omega) - \overline{v_u^2} = \overline{v_n^2} * |H(j\omega)|^2$$

$$\omega = 2\pi f$$

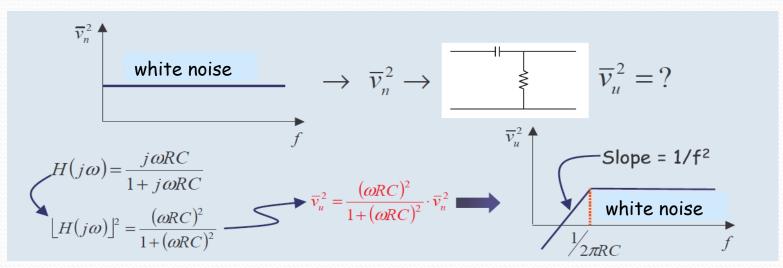
Noise power spectrum at output of a filter with transfer function $H(j\omega)$ is equal to input power spectrum multiplied by squared transfer function

The total noise depends on the bandwidth of the system. Since spectral noise components are non-correlated, we must integrate the noise power over the frequency range of the system

$$v_{on}^{2} = \int_{0}^{\infty} \overline{v_{un}^{2}} d\omega = \int_{0}^{\infty} \overline{v_{n}^{2}} * |H(\omega)|^{2} d\omega$$

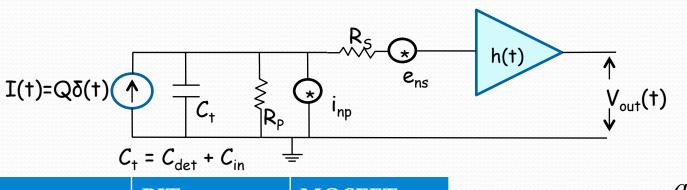
- The total noise increases with bandwidth
- Small bandwidth → large rise-times → less noise
- High bandwidth → fast pulse → more noise

Example: white noise source connected to high-pass filter



Optimum filter

In order to study the ENC and find the optimum filter (transfer function) of our amplifying system, it is convenient to represent our chain with a noiseless amplifier, with transfer function h(t) and all noise sources at its input, represented by R_s and R_p (we are considering only white noise source, not 1/f for the moment)



$\overline{e_n^2} = 4KTR_S$
$\frac{1}{i^2}$ 4KT
$\iota_n = \overline{R_P}$

	ВЈТ	MOSFET
$R_{\rm s}$	$1/(2g_m)$	$2/(3g_{\rm m})$
R_{p}	$2h_{FE}/g_{m}$	$2KT/(qI_G) \sim 0$

in general
$$R_S = \frac{a_n}{g_m}$$
 $a_n = \begin{cases} 0.5 \text{ in BJT} \\ 0.7 \text{ in Mosfet} \end{cases}$ $g_m = \text{conductance} = \frac{\partial I}{\partial V}$

It is possible to demonstrate that:

$$ENC^{2} = 2KTR_{S}C_{t}^{2}\int\left[\frac{d}{dt}h(t)\right]^{2}dt + \frac{2KT}{R_{p}}\int\left[h(t)\right]^{2}dt \longrightarrow ENC^{2} = 2KTR_{S}C_{t}^{2}\left[\int\left[h'(t)\right]^{2}dt + \frac{1}{\tau_{C}^{2}}\int\left[h(t)\right]^{2}dt\right]$$

$$\text{where } \tau_{C} = C_{t}\sqrt{R_{p}R_{S}}$$

$$\text{noise corner time constant}$$

Capacitive matching

- Parallel noise depends mainly by "external" factors (Feedback resistor, detector bias and leakage)
- Series noise depends on amplifier characteristics (Rs $\rightarrow g_m$, C_{in})

with proper design and dimensioning of preamp we can optimize $\mathsf{ENC}_{\mathsf{S}}$

$$ENC_{s}^{2} = 4KTR_{S}C_{t}^{2}\frac{1}{t_{m}} = 4KT\frac{a_{n}}{g_{m}}(C_{det} + C_{in})^{2}\frac{1}{t_{m}} = 4KTa_{n}C_{det}\frac{\tau_{A}}{t_{m}}\left[\sqrt{\frac{C_{det}}{C_{in}}} + \sqrt{\frac{C_{in}}{C_{det}}}\right]^{2}$$

The minimum value is when $C_{\text{det}} = C_{\text{in}}$

 $\begin{cases} a_n = 0.7 \text{ in MOS} \\ \tau_A = \frac{C_{in}}{g_m} \end{cases}$

Input transistor capacitance must be matched to detector capacitance

$$ENC^{2}_{s_Opt} = 16KTa_{n}C_{\det}\frac{\tau_{A}}{t_{m}}$$

ENC for Pseudo-Gaussian Shapers

Introducing also 1/f noise ...

$$ENC = \sqrt{\frac{F_v C_{in}^2 v_e^2}{T_{peak}} + F_i i_e^2 T_{peak}} + F_f K_f C_{in}^2$$

$$\sqrt{\frac{F_v C_{in}^2 v_e^2}{T_{peak}}} + F_i i_e^2 T_{peak} + F_f K_f C_{in}^2$$

$$\sqrt{\frac{F_v C_{in}^2 v_e^2}{T_{peak}}} + F_i i_e^2 T_{peak} + F_f K_f C_{in}^2$$

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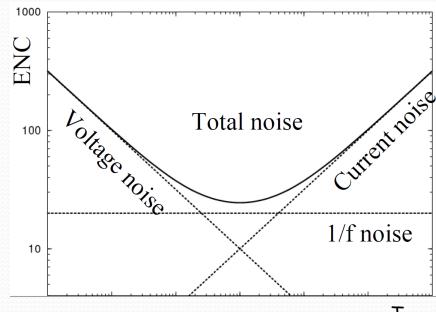
$$\sqrt{\frac{F_v C_{in}^2 v_e^2}{T_{peak}}} + F_i i_e^2 T_{peak} + F_f K_f C_{in}^2$$

$$\sqrt{\frac{F_v C_{in}^2 v_e^2}{T_{peak}}} + F_i i_e^2 T_{peak} + F_f K_f C_{in}^2$$

- F_v, F_i, F_f specific shape factors
- v_e^2 , i_e^2 voltage and current white noise densities
- K_f "1/f" noise constant
- C_{in} total input capacitance
- T_{peak} peaking time

- Voltage (serial) noise \downarrow when T_{peak} ↑
 Current (paralle) noise \downarrow when T_{peak} ↓
 1/f noise does not depend on T_{peak}

The best T_{peak} is that for ENC_S = ENC_P

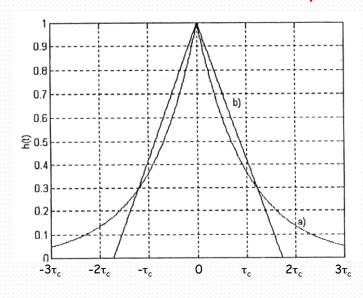


Optimum filter

What is the best h(t) that minimizes ENC?

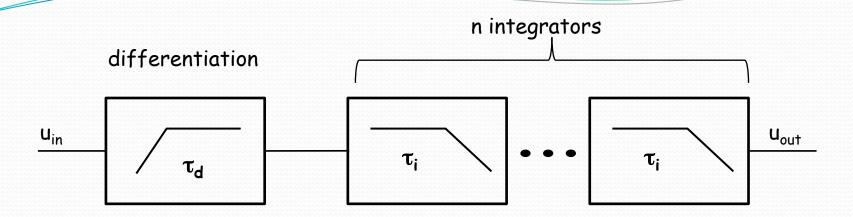
It is possible to demonstrate that
$$h_{opt}(t) = \exp\left(-\frac{|t|}{\tau_c}\right)$$
 \longrightarrow $ENC_{opt}^2 = 2KTR_S \frac{C_t^2}{\tau_c} = 2KTC_t \sqrt{\frac{R_s}{R_p}}$

This function is known as cusp or matched filter (curve a in the figure)



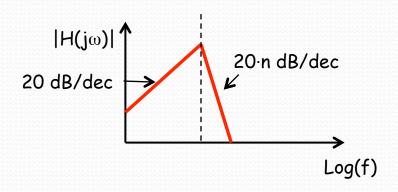
The cusp filter is not practically feasible, but can be approximated by triangular shapers (curve b) or Pseudo-Gaussian shaper

Pseudo-Gaussian (or Semi-Gaussian) shaper



- 1. A high-pass filter, that makes the derivative of the input pulse and introduces the decay time τ_d
- 2. n low-pass filters, that limits the bandwidth (and the noise) making the integral of the signal and limiting the rise time τ_i n is the order of the filter

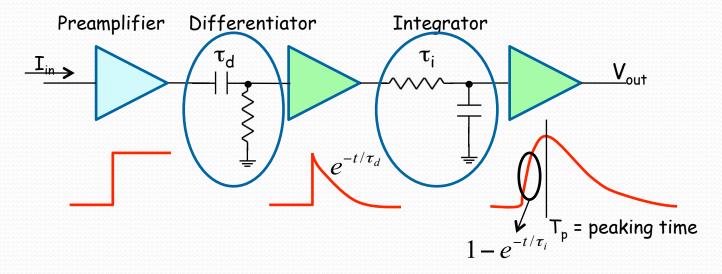
$$H(s) = \frac{u_{out}(s)}{u_{in}(s)} = \frac{s\tau_d}{(1+s\tau_d)} \frac{1}{(1+s\tau_i)^n}$$
 20 dB/dec



Simple shaper: CR-RC

The simplest Pseudo-Gaussian filter is the CR-RC shaper because:

- The high-pass filter is made with CR network
- 2. The low-pass filter is made with RC network



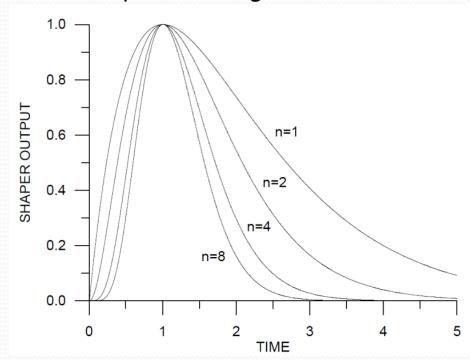
- This shaper is called CR-RC because the high-pass filter is made with CR network, while the low-pass filter with a RC network
- The noise is 36% worse than "optimum filter" with the same time constants

Shaper: CR-RCⁿ

The shapers are often more complicated, with multiple (n) integrators \rightarrow CR-RCⁿ

- Same peaking time if $\tau_n = \tau_{(n=1)}/n$
- With same peaking time
 - 1. More symmetrical
 - 2. Faster return to baseline
 - 3. Improved rate capability

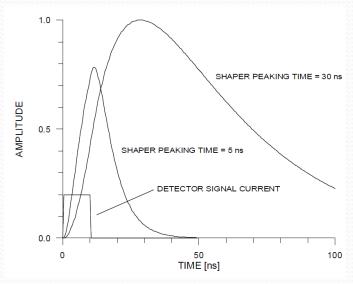
Shaper type	$\mathbf{F_v}$	$\mathbf{F_i}$
CR-RC	0.92	0.92
CR-RC ²	0.84	0.63
CR-RC ³	0.95	0.51
CR-RC ⁴	0.99	0.45
CR-RC ⁵	1.11	0.4
CR-RC ⁶	1.16	0.36
CR-RC ⁷	1.27	0.34



• Current noise decreases with shaping order

Ballistic deficit

Ballistic Deficit is a <u>Loss in Pulse Height</u> if the peaking time T_p of the shaper is shorter than the detector collection time or, more in general, the rise time of its input pulse



In fact, not all the charge is collected by the amplifier because it starts to discharge before the detector signal reaches its peak

Consequences:

- Loss of useful signal
- Increase of ENC (or decrease of S/N)

The shaping time must be carefully chosen, as a compromise among different factors:



- Short T_p: higher ENC, ballistic deficit but high sustainable event rate
- Long T_p: lower ENC but risk of pile-up

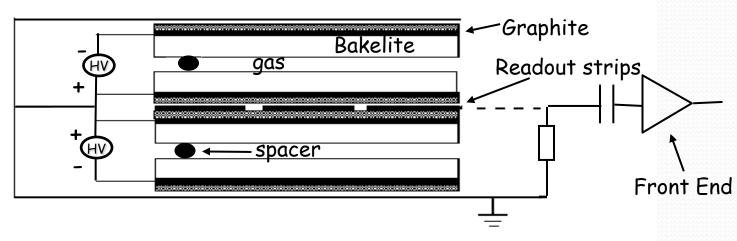
Some examples of Front-End amplifier

- 1. Current sensitive amplifier: Front-end chip of Resistive Plate Chambers (RPC) for the CMS Experiment, at CERN
- 2. Charge sensitive amplifier: Front-end chip of Cylindrical GEM (CGEM) for the KLOE Experiment, at Frascati INFN LAB
- 3. VFAT3/GdSP

CMS RPC Front-End Preamp

Detector Characteristics

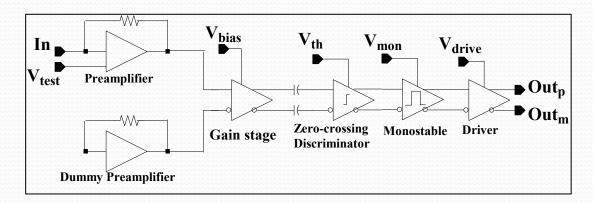
- ·Resistive Plate Chambers are gaseous parallel-plate detectors combining good spatial resolution with a time resolution comparable to that of scintillators (~ ns)
- They are suitable for fast space-time particle tracking as required for the muon trigger at the LHC experiments



Cross sectional view of a double-gap RPC

CMS RPC Front-End

- •RPC timing information is crucial for unambiguous assignment of the event to the related bunch crossing
- · We have developed an 8-channel front-end chip in BiCMOS technology



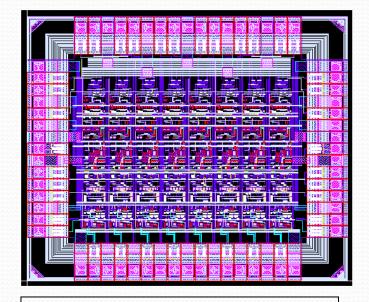
Single channel block diagram

Technology: 0.8 µm BiCMOS of AMS

8 channels

Power supplies: +5 V; GND

Power consumption: ~ 45 mW/channel



Dimensions: 2.9 mm X 2.6 mm

64 I/O pads

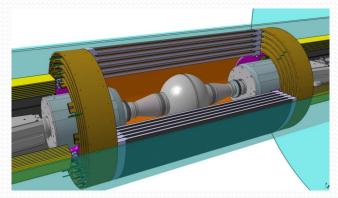
Package: PQFP 80

Ref.: F. Loddo et al., New developments on front-end electronics for the CMS Resistive Plate Chambers, Nucl. Instr. & Meth. A 456 (2000) 143-149

KLOE Inner Tracker Front-End

KLOE is an experiment at DAFNE accelerator, in Frascati INFN National Laboratories (Italy)

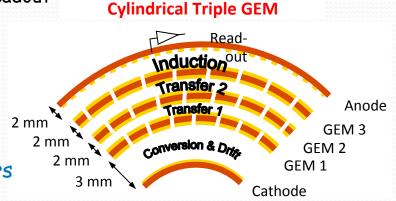
An Inner Tracker GEM-based will be inserted around the interaction point to improve the vertex resolution by a factor 3 (with respect to present detector, without IT)



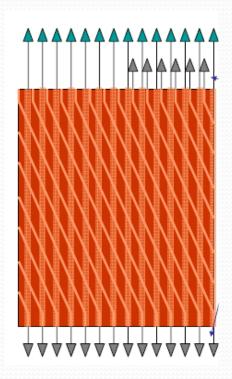
Gas Electron Multiplier (GEM) technology has been adopted, for its low material budget

- \circ 4 independent tracking layers for a fine vertex reconstruction of K_S and η
- \circ 200 μ m σ_{r_0} and 500 μ m σ_{z_0} spatial resolutions with XV readout
- o 700 mm active length
- from 130 to 220 mm radii
- \circ 1.8% X_0 total radiation length in the active region
- 5 kHz/cm² rate capability

Realized with <u>Cylindrical TRIPLE_GEM</u> detectors



KLOE IT Readout



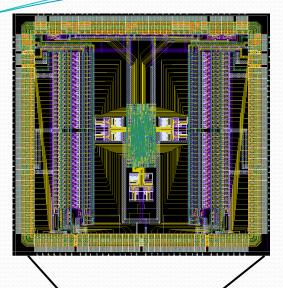
The Anode is shared out in a 2 dimensions layout having a XV geometry, X and V strips read out is shared out at both ends.

Rotated by roughly 40°, such strip's net provides 2D positioning for the particle passing through the layer.



KLOE-2 Inner Tracker

KLOE IT Front-End: GASTONE64



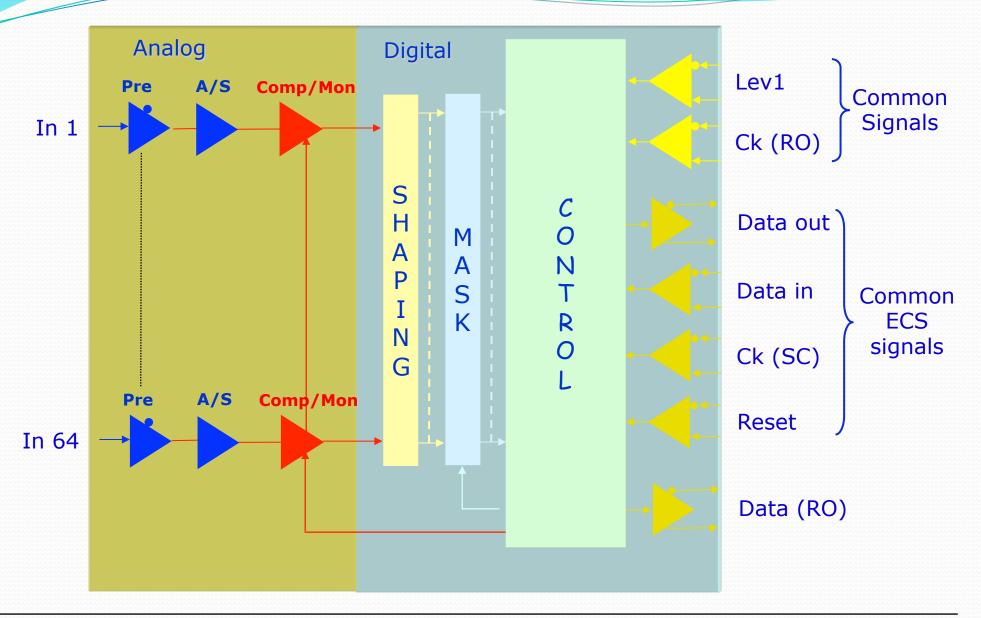


N. channels	64	
Technology	CMOS 0.35 µm	
Chip dimensions	4.5 X 4.5 mm ²	
Input impedance	120 Ω	
Charge sensitivity	16 mV/fC (Cdet = 100 pF)	
Peaking time	~90 ns (Cdet=100 pF)	
Crosstalk	< 3%	
ENC	800 e- + 40 e-/pF	
Power consumption	~ 6 mW/ch	
Readout	Serial LVDS (100 MBps)	

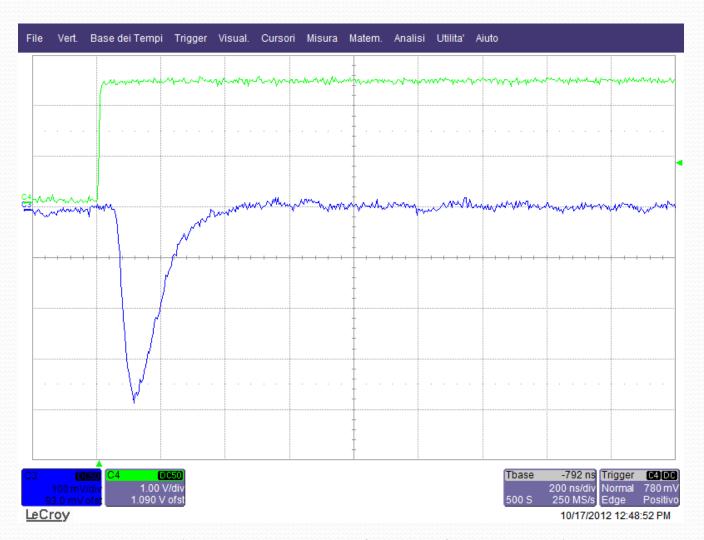
A. Balla et al., A new cylindrical GEM inner tracker for the upgrade of the KLOE experiment, Nucl. Phys. Proc. Suppl. 215:76-78,2011

A. Balla et al., GASTONE: A new ASIC for the cylindrical GEM inner tracker of KLOE experiment at DAFNE, Nucl. Instr. & Meth. A 604 (2009) 23-25

KLOE IT – GASTONE64

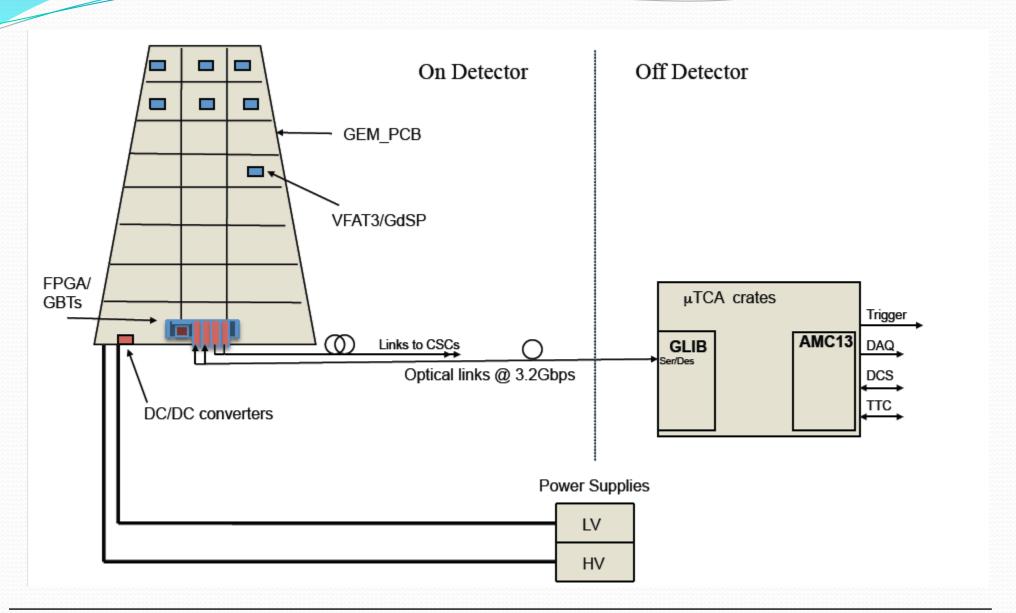


KLOE IT – GASTONE64



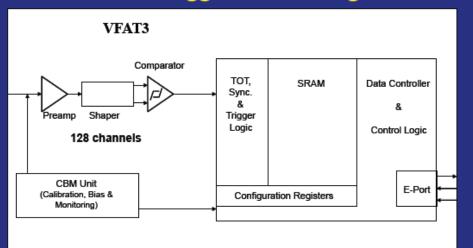
Shaper response for 20 fC input pulse

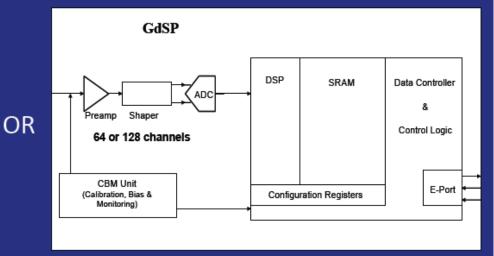
VFAT3: a proposed FE chip for CMS GEM



VFAT3

2 Trigger & Tracking Front-end architectures considered.





VFAT3:

Front-end with programmable shaping time.

Internal calibration.

Binary memory

Interface directly to GBT @ 320Mbps.

Designed for high rate (10kHz/cm^2 depending on segmentation)

Tech.: CMOS 130 nm

GdSP:

Similar to VFAT3 except has an ADC / channel instead of a comparator.

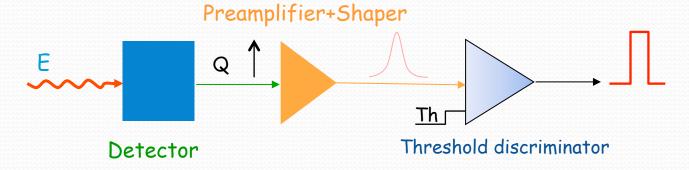
Internal DSP allows subtraction of background artifacts enabling a clean signal discrimination.

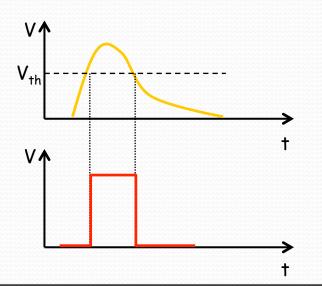
Time measurement

Time measurement

Often the purpose of the system Detector + Front End Electronics is time measurements

The simplest scheme is the following:





Leading edge or Threshold discriminator (comparator): when the signal crosses a threshold, the output goes from "low" to "high" level: we get a "time tag"

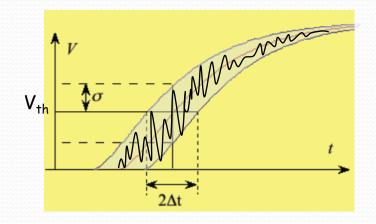
Accuracy of timing measurement is limited by:

- 1. Jitter
- 2. Time walk

Time measurement and noise: the jitter

Noise has an impact also in time measurements:

uncertainty in the time of crossing threshold -> Jitter



$$\Delta t = \frac{\sigma_{noise}}{dV/dt}$$
slope

How to decrease jitter? \rightarrow Conflicting conditions:

increase slope → increase bandwidth

As usual ... find compromise

Example:
$$V = V_{\text{max}} (1 - e^{-t/\tau})$$

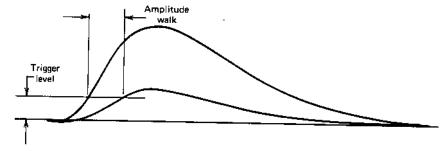
$$(\text{If } t \ll \tau)$$

$$\Delta t \approx \frac{\sigma_{noise} \tau}{V_{\text{max}}}$$

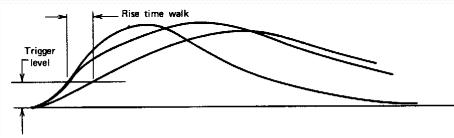
Rise time (10%-90%) = 2.2τ

The time walk

In the <u>leading edge discriminators</u>, two pulses with identical shape and time of occurrence, but different amplitude cross the same threshold in different times ($\Delta T = \text{time walk}$)



Even if the input amplitude is constant, time walk can still occur if the shape (rise time) of the pulse changes (for example, for changes in the charge collection time)



The sensitivity of leading edge triggering to time walk is minimized by <u>setting the threshold as low as possible</u> but it must be compatible with noise level and the discrimination point should be in a region where the slope is steep to minimize jitter

Time walk correction:

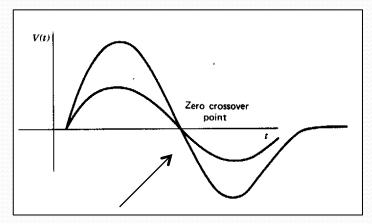
- <u>Software</u>: measure the pulse amplitude and apply correction to timing
- Hardware: instead of leading edge triggering, use
 - 1. Crossover timing
 - 2. Constant Fraction timing

Crossover timing

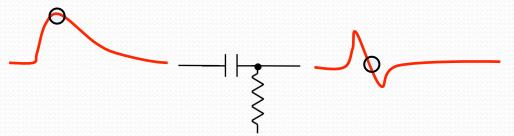
The crossover timing can greatly reduce the magnitude of the amplitude time walk

Hypothesis:

• the output of the shaper is a bipolar pulse and the time of crossing from the positive to the negative side of the axis (zero-crossing) is independent of the pulse amplitude



If the output of shaper is unipolar, but the peaking time is constant, adding a differentiator (C-R network) we get a bipolar pulse crossing the zero in correspondence of the signal peak

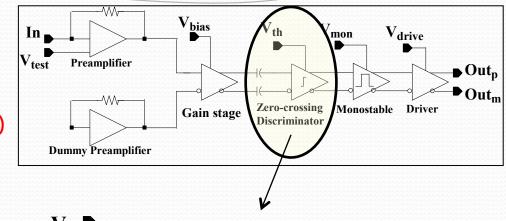


This method reduce amplitude walk, but usually jitter is larger than leading edge triggering

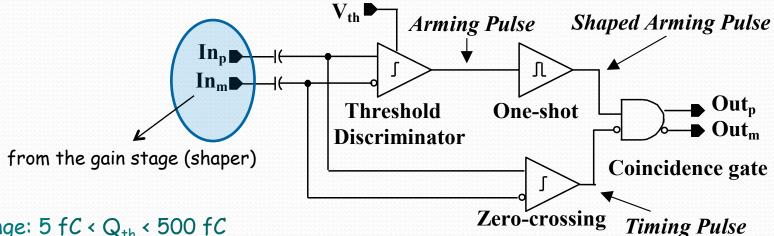
Example of crossover timing: CMS RPC Front-End

Zero-crossing discriminator:

- Threshold Discriminator (Charge selection)
- •Zero-Crossing Discriminator (Time reference)
- ·One-shot
- ·Coincidence gate



Discriminator

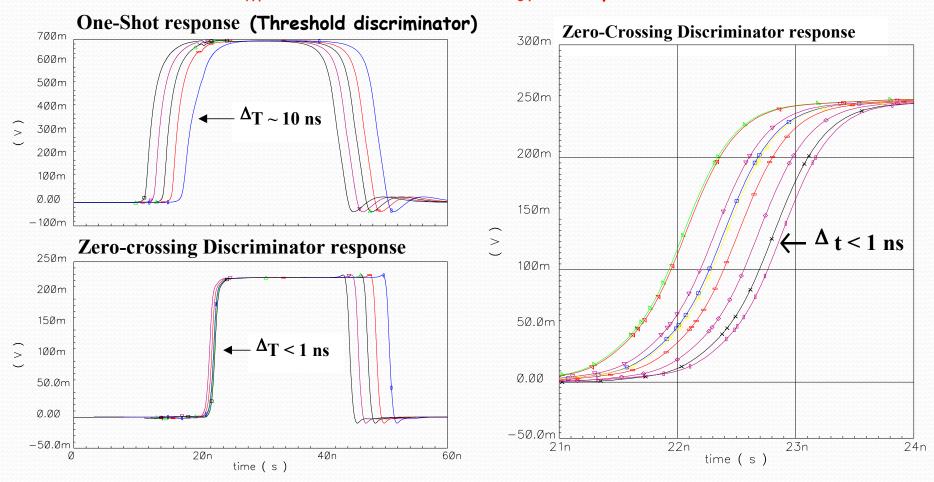


- Threshold range: 5 fC < Q_{th} < 500 fC
- Threshold uniformity: 1.5 fC rms
- Shaped Arming pulse width: 20 ns
- · Power: 8 mW

Example of crossover timing: CMS RPC Front-End

Zero-crossing discriminator response in the dynamic range

$$Q_{th} = 20 fC$$
 $1 fC < Q_{ov} < 20 pC$



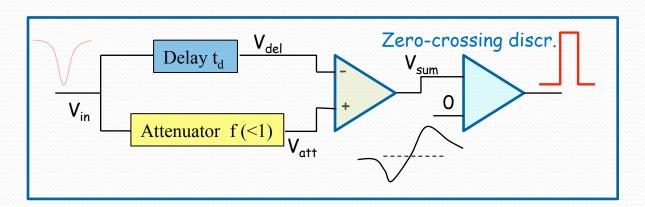
The Constant Fraction Timing

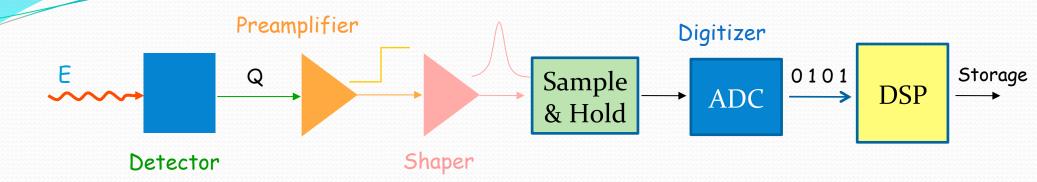
- •It is empirically found that the best leading edge timing characteristics are obtained when the threshold is set at about 10-20% of the pulse amplitude
- These observations have led to the development of the Constant Fraction Timing, that produce an output timing pulse a fixed time after the leading edge of the pulse has reached <u>a constant fraction</u> of the peak amplitude
- This point is independent of pulse amplitude for all pulses of constant shape, but with lower jitter than zero-crossing

Making the sum of

- inverted and delayed signal, with t_d > t_{rise}
- attenuated signal

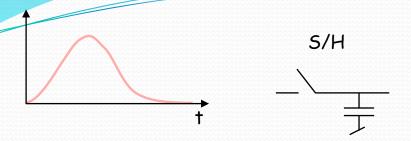
we get a bipolar signal, whose zero-crossing time is independent of pulse amplitude and corresponds to the time at which the pulse reaches the fraction \mathbf{f} of its final amplitude



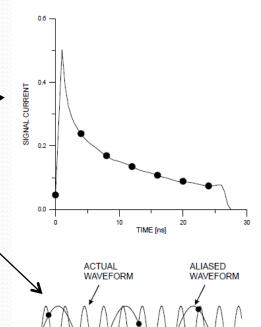


Another approach is to perform operations on the signal in the digital domain

- Baseline restoration
- Tail cancellation
- Filtering
- Zero suppression
- •
- 1. Sample the amplifier output (at a proper sampling rate)
- 2. Convert the analog value of the samples into digital word
- 3. Make operations on digital words using processors or Digital Signal Processors



If the sampling rate is too low, we loose information: high frequency components are *aliased* to lower frequencies and the reconstructed signal differs from the original one



Nyquist condition

The sampling rate must be $\geq 2 \times \text{highest signal frequency}$

Advantages:

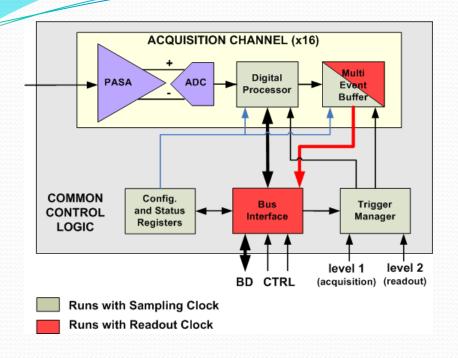
- great flexibility in implementing filter functions
- possibility to implement complex functions (impractical in hardware)
- easy to change filter parameters
- adaptive filtering can be used to compensate for pulse shape variations

Drawbacks:

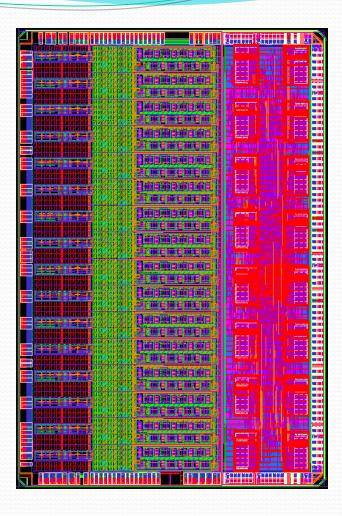
- increased complexity
- require fast and precise ADC
- · large silicon area required
- power consumption

Currently, there are many progress in the development of fast, low-power ADC, so the interest in DSP applications is growing and becoming more and more attractive

Example of Front-End + DSP: SAltro16



- Front-end for TPC, GEM, Micromega ...
- · Technology: CMOS 130 nm
- 16 channel demonstrator developed in 2009-2010
- Specification & architecture: Luciano Musa
- Design coordination: Paul Aspell
- Designers: Massimiliano De Gaspari, Hugo França-Santos, Eduardo Garcia



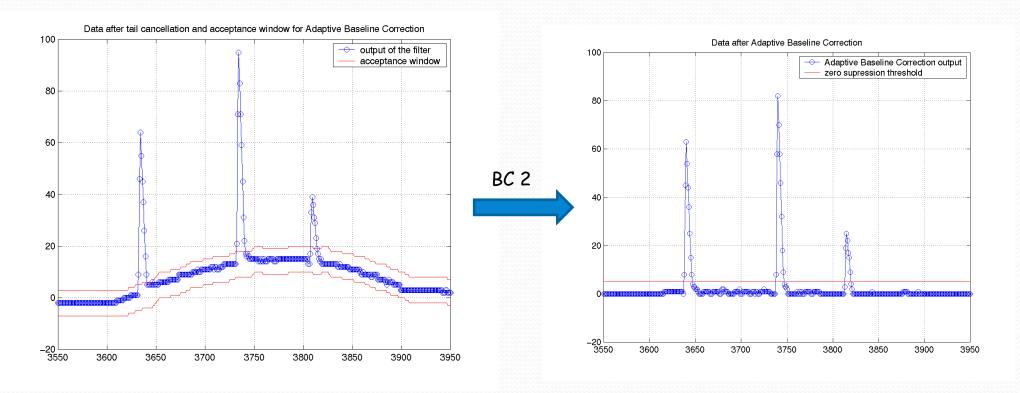
Example of Front-End + DSP: SAltro16

The Digital Signal Processing operations



- Baseline correction 1: removes systematic offsets that may have been introduced by noise pickup, clock interferences etc
- Common Mode Rejection: calculates the common mode pickup across all the channels and subtracts it
- Tail cancellation: removes the distortion of the signal shape caused by undershoot
- Baseline correction 2: reduce low frequency baseline shifts
- Zero suppression: removes from the data flow the samples smaller than a programmable threshold

Example of Front-End + DSP: SAltro16



After baseline correction, a fixed threshold can be applied safely

Summary

- The choice and design of Front-End electronics is crucial to obtain the desired energy and/or time resolution
- The technology strongly depends on the radiation environment
- \bullet ENC increases with detector capacitance and can be minimized matching \mathcal{C}_{det} with preamplifier input capacitance
- •The choice of pulse shape (and peaking time) comes out as a compromise between S/N optimization and double pulse resolution
- The shapers are built commonly with CR-RCⁿ filters
- · Depending on the event rate, baseline restoration may be needed
- Also digital shapers (using DSP) can be used, to enhance the flexibility and improve some optimizations (common mode rejections, baseline restoration ...) but require fast and precise ADC
- When the main goal is the time resolution, the Constant Fraction Timing provides the best results in terms of time walk, but requires higher circuital complexity respect to the simpler Leading Edge Timing and to the Zero-crossing Timing

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Thank you!