#### Marco Bonvini

**DESY Hamburg** 

LHCP, Barcelona, May 14, 2013



Work in collaboration with: Richard Ball, Stefano Forte, Simone Marzani, Giovanni Ridolfi arXiv:1303.3590









Convergence is slow! NNLO not definitive.

1

$$\sigma(\tau) = \tau \, \sigma_0 \sum_{ij} \int_{\tau}^1 \frac{dz}{z} \, \mathscr{L}_{ij}\left(\frac{\tau}{z}\right) C_{ij}(z, \alpha_s), \qquad \tau = \frac{m_H^2}{s}, \qquad z = \frac{m_H^2}{\hat{s}}$$

 $C_{ij}(z,\alpha_s) = \delta_{ig}\delta_{jg}\delta(1-z) + \alpha_s C_{ij}^{(1)}(z) + \alpha_s^2 C_{ij}^{(2)}(z) + \alpha_s^3 C_{ij}^{(3)}(z) + \dots$ 

$$\sigma(\tau) = \tau \,\sigma_0 \sum_{ij} \int_{\tau}^{1} \frac{dz}{z} \,\mathscr{L}_{ij}\left(\frac{\tau}{z}\right) C_{ij}(z,\alpha_s), \qquad \tau = \frac{m_H^2}{s}, \qquad z = \frac{m_H^2}{\hat{s}}$$
$$C_{ij}(z,\alpha_s) = \delta_{ig} \delta_{jg} \delta(1-z) + \alpha_s \, C_{ij}^{(1)}(z) + \alpha_s^2 \, C_{ij}^{(2)}(z) + \alpha_s^3 \, C_{ij}^{(3)}(z) + \dots$$

• NLO  $C_{ij}^{(1)}(z)$ :

- large  $m_t \ (\gg m_H)$  approximation
- full  $m_t$  dependence

[Dawson 1991; Djouadi, Spira, Zerwas 1991] [Spira, Djouadi, Graudenz, Zerwas 1995]

$$\sigma(\tau) = \tau \,\sigma_0 \sum_{ij} \int_{\tau}^{1} \frac{dz}{z} \,\mathscr{L}_{ij}\left(\frac{\tau}{z}\right) C_{ij}(z,\alpha_s), \qquad \tau = \frac{m_H^2}{s}, \qquad z = \frac{m_H^2}{\hat{s}}$$
$$C_{ij}(z,\alpha_s) = \delta_{ig} \delta_{jg} \delta(1-z) + \alpha_s \, C_{ij}^{(1)}(z) + \alpha_s^2 \, C_{ij}^{(2)}(z) + \alpha_s^3 \, C_{ij}^{(3)}(z) + \dots$$

- NLO  $C_{ij}^{(1)}(z)$ :
  - large  $m_t \ (\gg m_H)$  approximation
  - full  $m_t$  dependence

[Dawson 1991; Djouadi, Spira, Zerwas 1991] [Spira, Djouadi, Graudenz, Zerwas 1995]

### • NNLO $C_{ij}^{(2)}(z)$ :

• large  $m_t$  approximation [Harlander, Kilgore 2002; Anastasiou, Melnikov 2002] • expansion in  $m_H/m_t$  and in (1-z) [Harlander, Ozeren 2009] • expansion in  $m_H/m_t$  [Pak, Rogal, Steinhauser 2010]

• finite  $m_t$  small-z behavior [Marzani, Ball, Del Duca, Forte, Vicini 2008]

$$\sigma(\tau) = \tau \,\sigma_0 \sum_{ij} \int_{\tau}^{1} \frac{dz}{z} \,\mathscr{L}_{ij}\left(\frac{\tau}{z}\right) C_{ij}(z,\alpha_s), \qquad \tau = \frac{m_H^2}{s}, \qquad z = \frac{m_H^2}{\hat{s}}$$
$$C_{ij}(z,\alpha_s) = \delta_{ig} \delta_{jg} \delta(1-z) + \alpha_s \, C_{ij}^{(1)}(z) + \alpha_s^2 \, C_{ij}^{(2)}(z) + \alpha_s^3 \, C_{ij}^{(3)}(z) + \dots$$

- NLO  $C_{ij}^{(1)}(z)$ :
  - large  $m_t \ (\gg m_H)$  approximation
  - full  $m_t$  dependence

[Dawson 1991; Djouadi, Spira, Zerwas 1991] [Spira, Djouadi, Graudenz, Zerwas 1995]

### • NNLO $C_{ij}^{(2)}(z)$ :

- large  $m_t$  approximation [Harlander, Kilgore 2002; Anastasiou, Melnikov 2002]
- expansion in  $m_H/m_t$  and in (1-z) [Harlander, Ozeren 2009]
- expansion in  $m_H/m_t$ 
  - finite  $m_t$  small-z behavior [Marzani, Ball, De

# [Pak, Rogal, Steinhauser 2010]

[Marzani, Ball, Del Duca, Forte, Vicini 2008]

• NNNLO  $C_{ij}^{(3)}(z)$ :

- soft approximation in the large  $m_t$  limit [Moch, Vogt, 2005]
- large  $m_t$  as an expansion in (1-z) [Anastasiou, Duhr, Dulat, Mistlberger 2013+]

$$\sigma(\tau) = \tau \, \sigma_0 \sum_{ij} \int_{\tau}^{1} \frac{dz}{z} \, \mathscr{L}_{ij}\left(\frac{\tau}{z}\right) C_{ij}(z, \alpha_s), \qquad \tau = \frac{m_H^2}{s}, \qquad z = \frac{m_H^2}{\hat{s}}$$
$$C_{ij}(z, \alpha_s) = \delta_{ig} \delta_{jg} \delta(1-z) + \alpha_s \, C_{ij}^{(1)}(z) + \alpha_s^2 \, C_{ij}^{(2)}(z) + \alpha_s^3 \, C_{ij}^{(3)}(z) + \dots$$

- NLO  $C_{ij}^{(1)}(z)$ :
  - large  $m_t \ (\gg m_H)$  approximation
  - full  $m_t$  dependence

[Dawson 1991; Djouadi, Spira, Zerwas 1991] [Spira, Djouadi, Graudenz, Zerwas 1995]

# • NNLO $C_{ij}^{(2)}(z)$ :

- large  $m_t$  approximation [Harlander, Kilgore 2002; Anastasiou, Melnikov 2002]
- expansion in  $m_H/m_t$  and in (1-z) [Harlander, Ozeren 2009]
- ullet expansion in  $m_H/m_t$ 
  - finite  $m_t$  small-z behavior [Marz
- [Pak, Rogal, Steinhauser 2010] [Marzani, Ball, Del Duca, Forte, Vicini 2008]

- NNNLO  $C_{ij}^{(3)}(z)$ :
  - soft approximation in the large  $m_t$  limit [Moch, Vogt, 2005]
  - large  $m_t$  as an expansion in (1-z) [Anastasiou, Duhr, Dulat, Mistlberger 2013+]
- NNLO + NNLL resummation

[de Florian, Grazzini 2012]

$$\sigma(\tau) = \tau \,\sigma_0 \sum_{ij} \int_{\tau}^{1} \frac{dz}{z} \,\mathscr{L}_{ij}\left(\frac{\tau}{z}\right) C_{ij}(z,\alpha_s), \qquad \tau = \frac{m_H^2}{s}, \qquad z = \frac{m_H^2}{\hat{s}}$$
$$C_{ij}(z,\alpha_s) = \delta_{ig} \delta_{jg} \delta(1-z) + \alpha_s \, C_{ij}^{(1)}(z) + \alpha_s^2 \, C_{ij}^{(2)}(z) + \alpha_s^3 \, C_{ij}^{(3)}(z) + \dots$$

- NLO  $C_{ii}^{(1)}(z)$ :
  - large  $m_t \ (\gg m_H)$  approximation
  - full  $m_t$  dependence

[Dawson 1991; Djouadi, Spira, Zerwas 1991] [Spira, Djouadi, Graudenz, Zerwas 1995]

# • NNLO $C_{ii}^{(2)}(z)$ :

- large  $m_t$  approximation [Harlander, Kilgore 2002; Anastasiou, Melnikov 2002]
- expansion in  $m_H/m_t$  and in (1-z)[Harlander, Ozeren 2009] [Pak, Rogal, Steinhauser 2010]
- expansion in  $m_H/m_t$ 
  - finite  $m_t$  small-z behavior [Marzani, Ball, Del Duca, Forte, Vicini 2008]
- NNNLO  $C_{ii}^{(3)}(z)$ : ightarrow 
  ightarro
  - soft approximation in the large  $m_t$  limit [Moch, Vogt, 2005]
  - large  $m_t$  as an expansion in (1-z) [Anastasiou, Duhr, Dulat, Mistlberger 2013+]

#### NNLO + NNLL resummation

#### [de Florian, Grazzini 2012]

# Ingredients of our N<sup>3</sup>LO prediction

gg channel only:

$$egin{aligned} C_{gg}(z,lpha_s) \ &\simeq \ C_{
m soft}(z,lpha_s) \ + \ C_{
m high-energy}(z,lpha_s) \ &z 
ightarrow 1 \qquad z 
ightarrow 0 \end{aligned}$$

### Ingredients of our N<sup>3</sup>LO prediction

gg channel only:

$$C_{gg}(z, \alpha_s) \simeq C_{\text{soft}}(z, \alpha_s) + C_{\text{high-energy}}(z, \alpha_s)$$
  
 $z \rightarrow 1 \qquad z \rightarrow 0$ 

Outline:

- construction of  $C_{\text{soft}}(z, \alpha_s)$  and  $C_{\text{high-energy}}(z, \alpha_s)$
- comparison against known NLO and NNLO
- results at NNNLO

### Ingredients of our N<sup>3</sup>LO prediction

gg channel only:

 $\begin{array}{ll} C_{gg}(z,\alpha_s) \ \simeq \ C_{\rm soft}(z,\alpha_s) \ + \ C_{\rm high-energy}(z,\alpha_s) \\ \\ z \rightarrow 1 & z \rightarrow 0 \\ \\ N \rightarrow \infty & N \rightarrow 1 \end{array}$ 

Outline:

- construction of  $C_{\text{soft}}(z, \alpha_s)$  and  $C_{\text{high-energy}}(z, \alpha_s)$
- comparison against known NLO and NNLO
- results at NNNLO

Mellin space: 
$$C_{gg}(N, \alpha_s) = \int_0^1 dz \ z^{N-1} C_{gg}(z, \alpha_s)$$

Leading Log poles in  $C_{gg}(N, \alpha_s)$ :

$$\frac{\alpha_s^k}{(N-1)^k}$$



Leading Log poles in  $C_{gg}(N, \alpha_s)$ :  $\frac{\alpha_s^k}{(N-1)^k}$   $\left(\alpha_s^k \frac{\log^{k-1} z}{z}\right)$ 

$$C_{\mathsf{high-energy}}(N,\alpha_s) = \sum_{k_1,k_2 \ge 0} c_{k_1,k_2} \left(\frac{m_H}{m_t}\right) \left[\gamma_+^{k_1}\right] \left[\gamma_+^{k_2}\right]$$

 $\gamma_+(N)$ : DGLAP anomalous dimension (largest eigenvalue) [Marzani, Ball, Del Duca, Forte, Vicini 2008]

Leading Log poles in  $C_{gg}(N, \alpha_s)$ :  $\frac{\alpha_s^k}{(N-1)^k} \qquad \left(\alpha_s^k \frac{\log^{k-1} z}{z}\right)$ 

$$C_{\mathsf{high-energy}}(N,\alpha_s) = \sum_{k_1,k_2 \ge 0} c_{k_1,k_2}(\frac{m_H}{m_t}) \left[\gamma_+^{k_1}\right] \left[\gamma_+^{k_2}\right]$$

 $\gamma_+(N)$ : DGLAP anomalous dimension (largest eigenvalue) [Marzani, Ball, Del Duca, Forte, Vicini 2008]

Properties:

• valid at LL (with running coupling effects)

Leading Log poles in  $C_{gg}(N, \alpha_s)$ :  $\frac{\alpha_s^k}{(N-1)^k}$   $\left(\alpha_s^k \frac{\log^{k-1} z}{z}\right)$ 

$$C_{\mathsf{high-energy}}(N,\alpha_s) = \sum_{k_1,k_2 \ge 0} c_{k_1,k_2}(\frac{m_H}{m_t}) \left[\gamma_+^{k_1}\right] \left[\gamma_+^{k_2}\right]$$

 $\gamma_+(N)$ : DGLAP anomalous dimension (largest eigenvalue) [Marzani, Ball, Del Duca, Forte, Vicini 2008]

Properties:

- valid at LL (with running coupling effects)
- we use an expansion of  $\gamma_+(N)$  to NLL

Leading Log poles in  $C_{gg}(N, \alpha_s)$ :  $\frac{\alpha_s^k}{(N-1)^k}$   $\left(\alpha_s^k \frac{\log^{k-1} z}{z}\right)$ 

$$C_{\mathsf{high-energy}}(N,\alpha_s) = \sum_{k_1,k_2 \ge 0} c_{k_1,k_2}(\frac{m_H}{m_t}) \big[\gamma_+^{k_1}\big] \big[\gamma_+^{k_2}\big]$$

 $\gamma_+(N)$ : DGLAP anomalous dimension (largest eigenvalue) [Marzani, Ball, Del Duca, Forte, Vicini 2008]

Properties:

- valid at LL (with running coupling effects)
- we use an expansion of  $\gamma_+(N)$  to NLL
- momentum conservation  $C_{\text{high-energy}}(N=2,\alpha_s)=0$

## Soft part: $C_{soft}$

#### From soft-gluon (threshold) resummation

$$C_{gg}(N,\alpha_s) \stackrel{N \to \infty}{=} g_0\left(\alpha_s, \frac{m_H}{m_t}\right) \exp\left[\frac{1}{\alpha_s}g_1(\alpha_s \ln N) + g_2(\alpha_s \ln N) + \alpha_s g_3(\alpha_s \ln N) + \dots\right]$$

### Soft part: $C_{soft}$

#### From soft-gluon (threshold) resummation

 $C_{gg}(N,\alpha_s) \stackrel{N \to \infty}{=} g_0\left(\alpha_s, \frac{m_H}{m_t}\right) \exp\left[\frac{1}{\alpha_s}g_1(\alpha_s \ln N) + g_2(\alpha_s \ln N) + \alpha_s g_3(\alpha_s \ln N) + \dots\right]$ 

in z-space linear combination of

of 
$$\left(\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right)_+$$

### Soft part: $C_{soft}$

#### From soft-gluon (threshold) resummation

 $C_{gg}(N,\alpha_s) \stackrel{N \to \infty}{=} g_0\left(\alpha_s, \frac{m_H}{m_t}\right) \exp\left[\frac{1}{\alpha_s}g_1(\alpha_s \ln N) + g_2(\alpha_s \ln N) + \alpha_s g_3(\alpha_s \ln N) + \dots\right]$ 

+

in 
$$z$$
-space linear combination of  $\left(rac{\log^k \log rac{1}{z}}{\log rac{1}{z}}
ight)$ 

Our improvements:

• we use the correct logs (from kinematics):

$$\left(\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right)_+$$

### Soft part: C<sub>soft</sub>

#### From soft-gluon (threshold) resummation

 $C_{gg}(N,\alpha_s) \stackrel{N \to \infty}{=} g_0\left(\alpha_s, \frac{m_H}{m_t}\right) \exp\left[\frac{1}{\alpha_s}g_1(\alpha_s \ln N) + g_2(\alpha_s \ln N) + \alpha_s g_3(\alpha_s \ln N) + \dots\right]$ 

in z-space linear combination of 
$$\left(\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right)$$

#### Our improvements:

• we use the correct logs (from kinematics):

$$\left(\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right)_+$$

• we supply each emission with the Altarelli-Parisi splitting

$$P_{gg} = \frac{C_A}{\pi} \frac{1 - 2z + 3z^2 - 2z^3 + z^4}{z (1 - z)} \qquad (z < 1)$$

(avoiding double counting)

[Krämer, Laenen, Spira 1997]

 $C_{gg}(N,\alpha_s) = 1 + \alpha_s C^{(1)}(N) + \alpha_s^2 C^{(2)}(N) + \alpha_s^3 C^{(3)}(N) + \dots$ 



Marco Bonvini

Higgs production in gluon fusion beyond NNLO

$$C_{gg}(N,\alpha_s) = 1 + \alpha_s C^{(1)}(N) + \alpha_s^2 C^{(2)}(N) + \alpha_s^3 C^{(3)}(N) + \dots$$



Marco Bonvini

Higgs production in gluon fusion beyond NNLO

$$C_{gg}(N,\alpha_s) = 1 + \alpha_s C^{(1)}(N) + \alpha_s^2 C^{(2)}(N) + \alpha_s^3 C^{(3)}(N) + \dots$$



Marco Bonvini

Higgs production in gluon fusion beyond NNLO

 $C_{gg}(N,\alpha_s) = 1 + \alpha_s C^{(1)}(N) + \alpha_s^2 C^{(2)}(N) + \alpha_s^3 C^{(3)}(N) + \dots$ 









N-soft: [Moch, Vogt, 2005], expanded resummation [de Florian, Grazzini 2012]

Marco Bonvini

### Conclusions

• We are predicting the inclusive Higgs N<sup>3</sup>LO cross section using

$$C_{gg}^{(3)}(z) \simeq C_{\rm soft}^{(3)}(z) + C_{\rm high-energy}^{(3)}(z)$$

- exact  $m_t$  dependence
- improved soft approximation (kinematical logs and AP splittings)
- high-energy behavior at LL plus some NLL elements

### Conclusions

• We are predicting the inclusive Higgs N<sup>3</sup>LO cross section using

$$C_{gg}^{(3)}(z) \simeq C_{\rm soft}^{(3)}(z) + C_{\rm high-energy}^{(3)}(z)$$

- exact  $m_t$  dependence
- improved soft approximation (kinematical logs and AP splittings)
- high-energy behavior at LL plus some NLL elements
- We find  $(m_H = 125 \text{ GeV}, \text{ LHC at } 8 \text{ TeV})$ 
  - $\bullet$  an increase of  $\sim 17\%$  wrt the NNLO cross section
  - dramatic stabilization of scale dependence

### Conclusions

• We are predicting the inclusive Higgs N<sup>3</sup>LO cross section using

$$C_{gg}^{(3)}(z) \simeq C_{\rm soft}^{(3)}(z) + C_{\rm high-energy}^{(3)}(z)$$

- exact  $m_t$  dependence
- improved soft approximation (kinematical logs and AP splittings)
- high-energy behavior at LL plus some NLL elements
- We find  $(m_H = 125 \text{ GeV}, \text{ LHC at } 8 \text{ TeV})$ 
  - $\bullet$  an increase of  $\sim 17\%$  wrt the NNLO cross section
  - dramatic stabilization of scale dependence
- Public code: http://www.ge.infn.it/~bonvini/higgs/

• We are predicting the inclusive Higgs N<sup>3</sup>LO cross section using

$$C_{gg}^{(3)}(z) \simeq C_{\rm soft}^{(3)}(z) + C_{\rm high-energy}^{(3)}(z)$$

- exact  $m_t$  dependence
- improved soft approximation (kinematical logs and AP splittings)
- high-energy behavior at LL plus some NLL elements
- We find  $(m_H = 125 \text{ GeV}, \text{LHC at } 8 \text{ TeV})$ 
  - $\bullet$  an increase of  $\sim 17\%$  wrt the NNLO cross section
  - dramatic stabilization of scale dependence
- Public code: http://www.ge.infn.it/~bonvini/higgs/
- What's next?
  - the  $\delta(1-z)$  term at order  $\alpha_s^3$
  - subleading high-energy terms
  - other channels (at NNLO the qg channel gives a 10% contribution)
  - resummation

# Backup slides





### Scale dependence



### Scale dependence



# K-factors (NNLO pdfs)



# K-factors (NNLO pdfs)



# K-factors (NNLO pdfs)



 $13 {\rm ~TeV}$ 



Higgs hadron-level cross section

### Soft approximations



### Soft approximations



### Soft approximations



# Saddle point

