

A Cocktail Model:
From Radiative ν Masses
to Dark Matter and Back

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In Collaboration with Michael Gustafsson & Maximiliano Rivera.

November 14th 2012



Neutrino Masses

⇒ *Neutrino Oscillations*  Direct Experimental Evidence
of Massive Neutrinos

$$\Delta m_{21}^2 \equiv m_2^2 - m_1^2 = 7.62_{-0.19}^{+0.19} \times 10^{-5} \text{eV}^2 \quad |\Delta m_{31}^2| \equiv |m_3^2 - m_1^2| = 2.53_{-0.1}^{+0.08} \times 10^{-3} \text{eV}^2$$



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⇒ $m_\nu = 0$ in the SM

→ NO Right-Handed Neutrinos



No Dirac Mass Term

→ NO Lepton Number Violation



No Majorana Mass Term

⇒ Can Generate Majorana M_ν in SM via Effective Operators

Lowest Order: Dim-5 (Unique)

$$\overline{L}_L^c \tilde{\phi}^* \tilde{\phi}^\dagger L_L + \text{h.c.}$$

Neutrino Masses

If we add Right-Handed Neutrinos to the SM...



Yukawa coupling for Neutrinos

$$\mathcal{L} \supset -\bar{L}_L \phi Y_l l_R - \bar{L}_L \tilde{\phi} Y_\nu \nu_R$$

(Dirac Mass Term)



Neutrino Masses

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(Dirac Mass Term)

However...

Neutrino Oscillation Data

+

$$m_{\nu_e} \equiv \sqrt{\sum_i |U_{ei}|^2 m_{\nu_i}^2} < 2 \text{ eV}$$



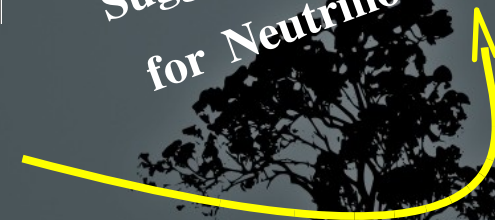
$$m_\nu < \text{eV}$$

Tritium β - Decay

$$Y_\nu \sim 10^{-11} \ll Y_l$$

Suggests Alternative Mechanism
for Neutrino Mass Generation

Why So Small?



“Natural” Neutrino Mass Generation

1 See-Saw Mechanism:

Type I See-Saw
(Fermion Singlet)

$$\mathcal{L} \supset -\bar{L}_L \phi Y_l l_R - \bar{L}_L \tilde{\phi} Y_\nu \nu_R - \frac{1}{2} \bar{\nu}_R^c M_R \nu_R$$

Majorana Mass Term

Lepton Number Breaking

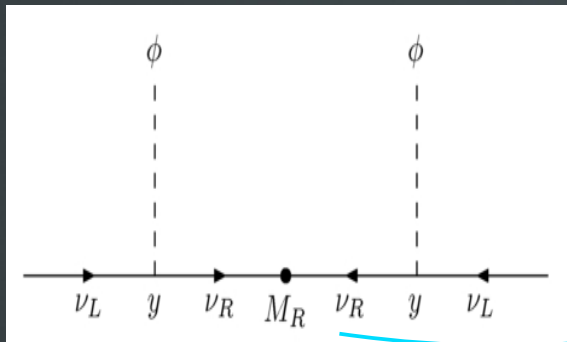
Majorana Neutrino
Mass Matrix

$$\begin{pmatrix} 0 & Y_\nu v \\ Y_\nu v & M_R \end{pmatrix} \rightarrow \begin{pmatrix} \frac{(Y_\nu v)^2}{M_R} & 0 \\ 0 & M_R \end{pmatrix}$$

$$m_\nu \sim (Y_\nu v)^2 / M_R$$

$$M_R \sim 10^{15} \text{ GeV} \rightarrow Y_\nu \sim 1$$

$$M_R \sim \text{TeV} \rightarrow Y_\nu \sim 10^{-6}$$



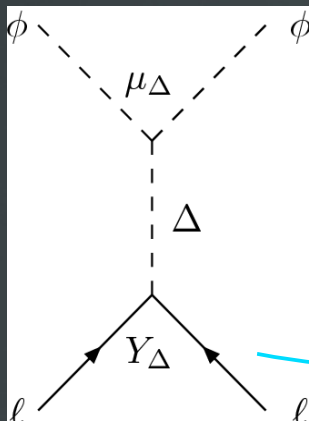
*D=5 Weinberg Operator
Generated at Tree Level*



“Natural” Neutrino Mass Generation

1 See-Saw Mechanism:

Type II See-Saw
(Scalar Triplet)



*D=5 Weinberg Operator
Generated at Tree Level*

$$\mathcal{L} \supset -\vec{\Delta}^\dagger M_\Delta^2 \vec{\Delta} + \overline{\tilde{L}}_L Y_\Delta (\vec{\sigma} \cdot \vec{\Delta}) L_L + \mu_\Delta \tilde{\phi}^\dagger (\vec{\sigma} \cdot \vec{\Delta})^\dagger \phi + \text{h.c.}$$

Lepton Number Breaking

$$m_\nu \sim Y_\Delta \mu_\Delta (v/M_\Delta)^2$$

...Type III See-Saw
(Fermion Triplet)

Typically $M_R, M_\Delta \gg 100 \text{ GeV}$
in See-Saw models

“Natural” Neutrino Mass Generation

② Radiative Neutrino Masses:

(Typically Involving Extended Scalar Sectors)

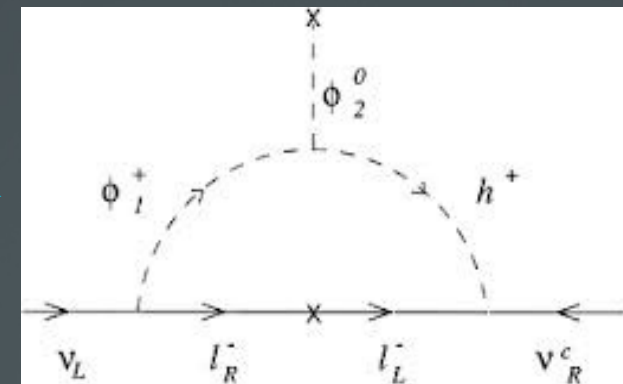
Zee Model

$$-\Delta\mathcal{L}_{Zee} = \kappa_1 \phi_1^T i\sigma_2 \phi_2 h^- + Y_{ab} \phi_1 \bar{l}_{L_a} l_{R_b} + f_{ab} \bar{l}_{L_a}^c i\sigma_2 l_{L_b} h^+ + \text{h.c.}$$

Lepton Number Breaking

Majorana Mass Matrix
for Left-Handed Neutrinos
Generated at 1-loop

A. Zee, *Phys. Lett. B* **93** (1980) 389



“Natural” Neutrino Mass Generation

② Radiative Neutrino Masses:

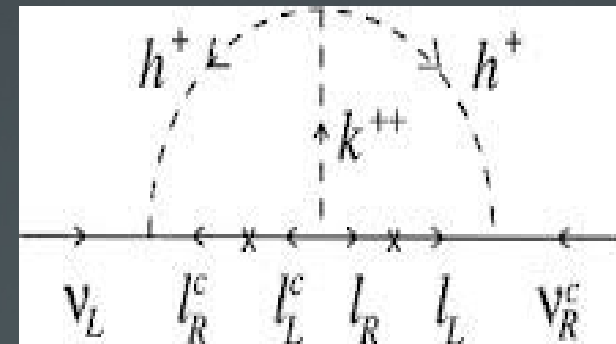
Babu Model

$$-\Delta\mathcal{L}_{\text{Babu}} = \kappa_2 h^- h^- k^{++} + C_{ab} \overline{l_{R_a}^c} l_{R_b} k^{++} + f_{ab} \overline{l_{L_a}^c}^T i\sigma_2 l_{L_b} h^+ + \text{h.c.}$$

Lepton Number Breaking

Majorana Mass Matrix
for Left-Handed Neutrinos
Generated at 2-loops

K. S. Babu, Phys. Lett. B **203** (1988) 132



Both Zee and Babu Models Generate the D = 5 Weinberg Operator at Loop Level

“Natural” Neutrino Mass Generation

② Radiative Neutrino Masses:

Other Models... (inspired by Zee/Babu)

C. S. Chen, C. Q. Geng & J. N. Ng, Phys. Rev. D **75** (2007) 053004

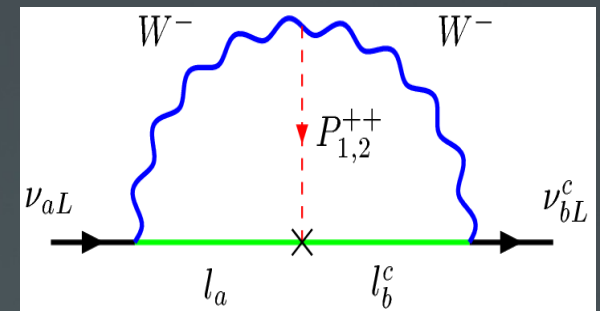
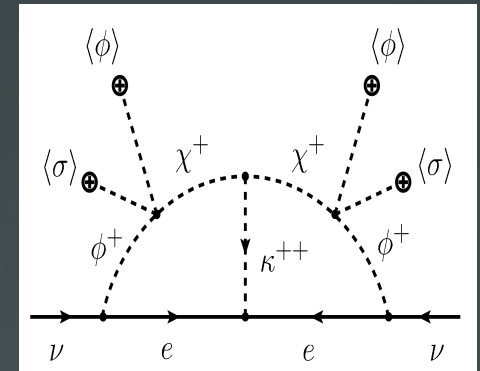
F. del Aguila, A. Aparici, A. Santamaria & J. Wudka, JHEP **1205** (2012) 133

...

Very Rich Phenomenology:

➔ *Possible Collider Signatures at LHC & ILC*

- ➔ New Charged States
- ➔ Same Sign Di-Leptons
- ➔ LFV Processes
- ...



“Natural” Neutrino Mass Generation

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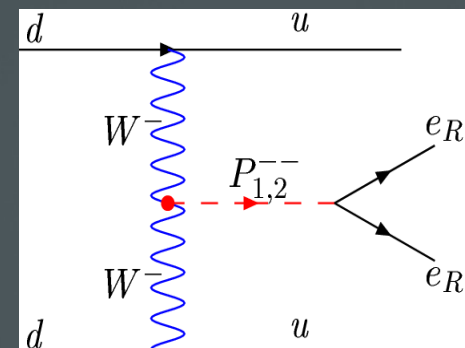
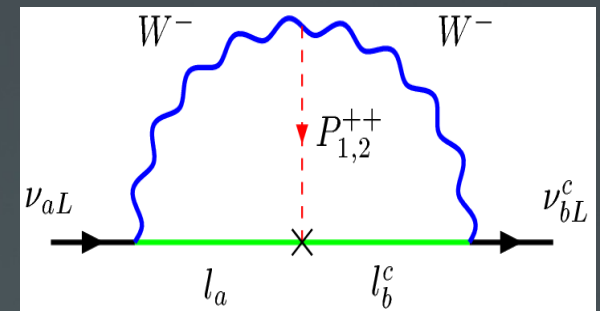
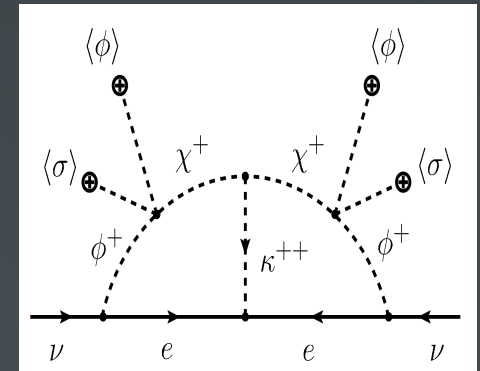
F. del Aguila, A. Aparici, A. Santamaria & J. Wudka, *JHEP* **1205** (2012) 133

...

Very Rich Phenomenology:

- ➔ Large Lepton Flavour Violation
- ➔ Large Neutrinoless Double β -Decay

LFV and $0\nu\beta\beta$ Much More Suppressed
in See-Saw Models (typically)



$0\nu\beta\beta$ Decay independent of m_ν

“Natural” Neutrino Mass Generation

Radiative Neutrino Mass Generation:

→ Attractive Alternative to See-Saw Mechanism

→ Very Rich (and Testable) Phenomenology



“Natural” Neutrino Mass Generation

Radiative Neutrino Mass Generation:

→ *Attractive Alternative to See-Saw Mechanism*

→ *Very Rich (and Testable) Phenomenology*

→ *Possible Connection between Neutrino Masses & Dark Matter*



Linking Neutrino Masses & WIMP Dark Matter

⇒ *WIMP Dark Matter from BSM Extended Scalar Sectors*

Singlet Higgs-Portal DM

J. McDonald, Phys. Rev. D **50** (1994) 3637

C. P. Burgess, M. Pospelov & T. ter Veldhuis, Nucl. Phys. B **619** (2001) 709

...

$$-\mathcal{L}_S \supset V(\phi) + \xi |\phi|^2 S^2$$

Inert Doublet Model

E. Ma, Phys. Rev. D **73** (2006) 077301

R. Barbieri, L. J. Hall and V. S. Rychkov, Phys. Rev. D **74** (2006) 015007

L. Lopez-Honorez, E. Nezri, J. F. Oliver and M. Tytgat, JCAP **1207** (2007) 028

...

...

$$-\mathcal{L}_D \supset V(|\phi_1|^2, |\phi_2|^2) + \lambda_4 \left| \phi_1^\dagger \phi_2 \right|^2 + \lambda_5 \left(\phi_1^\dagger \phi_2 \right)^2 + \text{h.c.}$$



Linking Neutrino Masses & WIMP Dark Matter

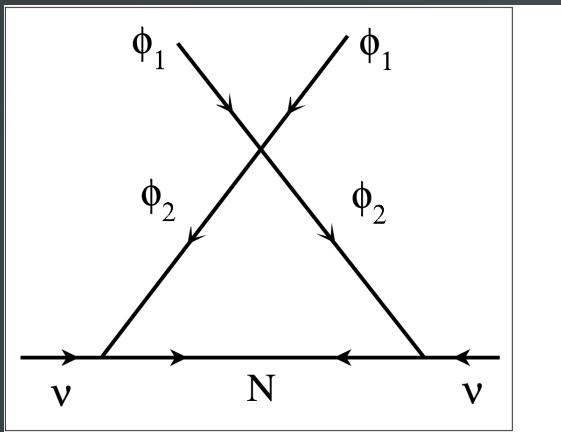
⇒ *WIMP Dark Matter from BSM Extended Scalar Sectors*

⇒ *Small ν Masses May be Connected to Dark Matter Stability*

*E. Ma, Phys. Rev. D **73** (2006) 077301*

*L. Krauss, S. Nasri and M. Trodden, Phys. Rev. D **67** (2003) 085002*

$$-\mathcal{L} \supset V(|\phi_1|^2, |\phi_2|^2) + \lambda_4 |\phi_1^\dagger \phi_2|^2 + \lambda_5 (\phi_1^\dagger \phi_2)^2 + \frac{1}{2} \overline{\nu_R^c} M_R \nu_R + \overline{L}_L \phi_1 Y_l l_R + \overline{L}_L \tilde{\phi}_1 Y_\nu \nu_R + \overline{L}_L \tilde{\phi}_2 Y_\nu \nu_R + \text{h.c.}$$



Radiative (1-loop) See-Saw

$N (\nu_R)$ or *Neutral component of ϕ_2*

Fermionic



Bosonic

DM candidate

Linking Neutrino Masses & WIMP Dark Matter

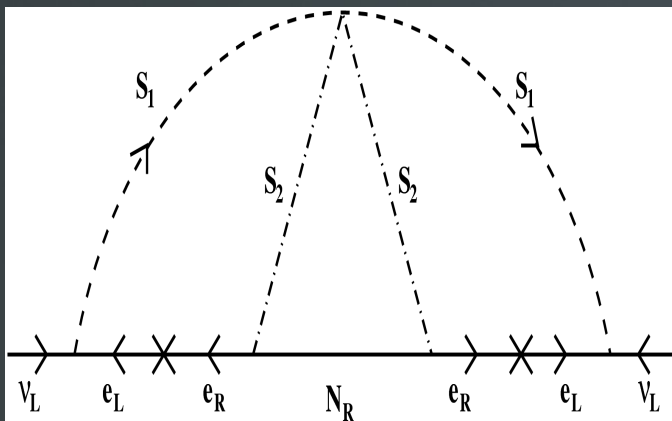
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*E. Ma, Phys. Rev. D **73** (2006) 077301*

*L. Krauss, S. Nasri and M. Trodden, Phys. Rev. D **67** (2003) 085002*

$$-\mathcal{L} \supset V(|S_1|^2, |S_2|^2) + \lambda_4 |S_1^+ S_2^-|^2 + \lambda_5 (S_1^+ S_2^-)^2 + f_{ab} \overline{l_{L_a}^c}^T i\sigma_2 l_{L_b} S_1^+ + g_a \nu_R S_2^+ l_{R_a} + \frac{1}{2} \overline{\nu_R^c} M_R \nu_R + \text{h.c.}$$



$N(\nu_R)$ is a DM candidate

3-loop Majorana Mass!



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→ *Z_2 Symmetry (DM) Responsible for Small M_ν (Loop Suppression)*

→ *In Both Models, $D = 5$ Weinberg Operator → Leading Contribution to M_ν*



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→ *ν_R as DM Candidate:*

→ No Direct Detection signal

→ Indirect Detection signal

Linking ν Masses to WIMP Dark Matter
May Influence DM Phenomenology

Neutrino Masses & WIMP Dark Matter

⇒ *Naturally Small ν Masses from BSM Extended Scalar Sectors*
+ *(inspired by Zee/Babu)*
WIMP Dark Matter

→ *Z_2 Symmetry (DM) Responsible for Small M_ν*

→ *Rich Phenomenology (LFV, $0\nu\beta\beta$, Colliders...)*

→ *Properties of Neutrino Sector Affected in Dramatic Way*



Neutrino Masses & WIMP Dark Matter

- ⇒ *Naturally Small ν Masses from BSM Extended Scalar Sectors*
+ *(inspired by Zee/Babu)*
WIMP Dark Matter

$$\Phi_1 = \begin{pmatrix} G^+ \\ \frac{h+iG_0}{\sqrt{2}} \end{pmatrix} + \begin{pmatrix} 0 \\ v \end{pmatrix} \quad \Phi_2 = \begin{pmatrix} \Lambda^+ \\ \frac{H_0+iA_0}{\sqrt{2}} \end{pmatrix} \quad S^+ \quad \rho^{++}$$



Neutrino Masses & WIMP Dark Matter

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(inspired by Zee/Babu)

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Inert Doublet Model



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Z_2

+

-

-

+



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$$-\Delta\mathcal{L} = C_{ab} \overline{l_{R_a}^c} l_{R_b} \rho^{++} + V(|\Phi_1|^2, |\Phi_2|^2, |S|^2, |\rho|^2) + \lambda_4 \left| \Phi_1^\dagger \Phi_2 \right|^2 + \frac{\lambda_5}{2} \left(\Phi_1^\dagger \Phi_2 \right)^2$$

$$+ \kappa_1 \Phi_1^T i\sigma_2 \Phi_2 S^- + \lambda_{\rho S} \Phi_1^T i\sigma_2 \Phi_2 S^+ \rho^{--} + \kappa_2 S^- S^- \rho^{++} + \text{h.c.}$$

For $\kappa_1 \neq 0$, S- Λ Mixing

$$\begin{aligned} \Lambda^+ &= c_\theta H_1^+ - s_\theta H_2^+ \\ S^+ &= s_\theta H_1^+ + c_\theta H_2^+ \end{aligned}$$

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(inspired by Zee/Babu)

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$$-\Delta\mathcal{L} = C_{ab} \overline{l_{R_a}^c} l_{R_b} \rho^{++} + V(|\Phi_1|^2, |\Phi_2|^2, |S|^2, |\rho|^2) + \lambda_4 \left| \Phi_1^\dagger \Phi_2 \right|^2 + \frac{\lambda_5}{2} \left(\Phi_1^\dagger \Phi_2 \right)^2 \\ + \kappa_1 \Phi_1^T i\sigma_2 \Phi_2 S^- + \lambda_{\rho S} \Phi_1^T i\sigma_2 \Phi_2 S^+ \rho^{--} + \kappa_2 S^- S^- \rho^{++} + \text{h.c.}$$

Lepton Number Breaking

Need $C_{ab} \neq 0, \lambda_5 \neq 0, (\kappa_1 \neq 0)$

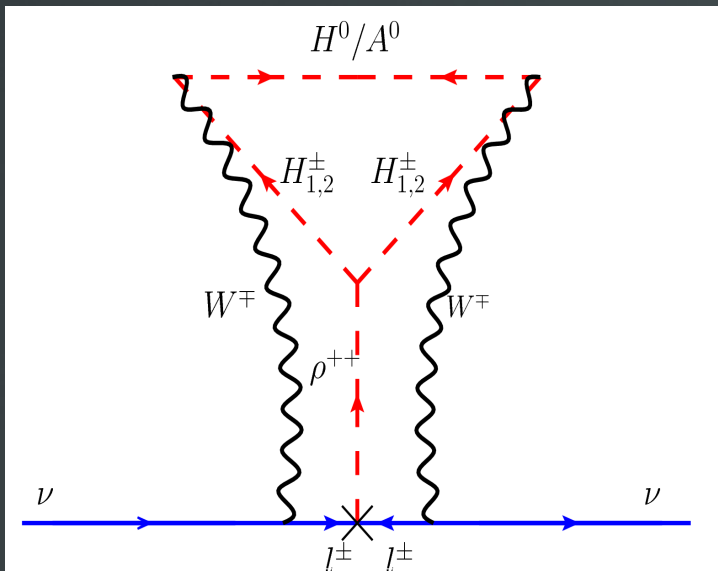
Need $\lambda_{\rho S}$ or $\kappa_2 \neq 0$

Neutrino Masses & WIMP Dark Matter

⇒ Naturally Small ν Masses from BSM Extended Scalar Sectors

+
WIMP Dark Matter

(inspired by Zee/Babu)



ν Masses generated at 3-Loops
via the "Cocktail Diagram"

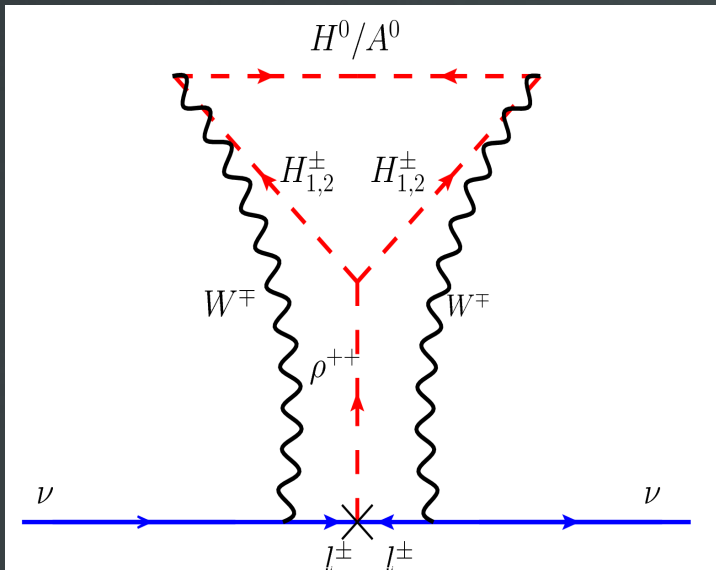
$$m_{\nu_{ab}} \simeq \frac{\text{Sin}^2(2\theta)}{(16\pi^2)^3} \frac{m_{A_0}^2 - m_{H_0}^2}{m_{A_0} m_{H_0}} \left(\frac{m_{H_1}^2 - m_{H_2}^2}{m_{H_1} m_{H_2}} \right)^2 \frac{(\kappa_2 + \lambda_\rho S v)}{m_\rho^2} m_{l_a} m_{l_b} C_{ab}$$

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→ Leading Contribution to M_ν from D = 9 Operator (D = 5 Operator Sub-Leading)

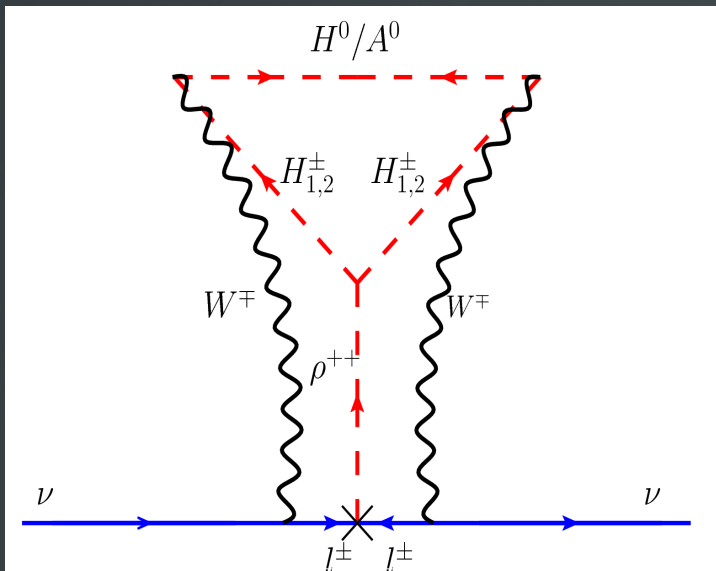
WEINBERG OPERATOR
at 5-LOOPS

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"GIM" Mechanism: $m_{A_0}^2 - m_{H_0}^2 \sim \lambda_5$, $m_{H_1}^2 - m_{H_2}^2 \sim \kappa_1$

Experimental Constrains & Model Predictions

① *Neutrino Masses & Mixings*

$$m_\nu = U^T m_\nu^D U$$



Experimental Constrains & Model Predictions

① Neutrino Masses & Mixings

$$m_\nu = U^T m_\nu^D U$$

$$U = \begin{pmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{pmatrix} \times \begin{pmatrix} c_{13}c_{12} & -c_{23}s_{12} - s_{23}c_{12}s_{13}e^{i\delta} & s_{23}s_{12} - c_{23}c_{12}s_{13}e^{i\delta} \\ c_{13}s_{12} & c_{23}c_{12} - s_{23}s_{12}s_{13}e^{i\delta} & -s_{23}c_{12} - c_{23}s_{12}s_{13}e^{i\delta} \\ s_{13}e^{-i\delta} & s_{23}c_{13} & c_{23}c_{13} \end{pmatrix}$$

2 Majorana Phases

3 Mixing Angles, 1 CP Phase

$$m_\nu^D = \begin{pmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{pmatrix}$$

3 Neutrino Masses

$$s_{12} = \sin(\theta_{12}), s_{13} = \sin(\theta_{13}), s_{23} = \sin(\theta_{23})$$
$$c_{12} = \cos(\theta_{12}), c_{13} = \cos(\theta_{13}), c_{23} = \cos(\theta_{23})$$



Experimental Constrains & Model Predictions

① Neutrino Masses & Mixings

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$$U = \begin{pmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{pmatrix} \times \begin{pmatrix} c_{13}c_{12} & -c_{23}s_{12} - s_{23}c_{12}s_{13}e^{i\delta} & s_{23}s_{12} - c_{23}c_{12}s_{13}e^{i\delta} \\ c_{13}s_{12} & c_{23}c_{12} - s_{23}s_{12}s_{13}e^{i\delta} & -s_{23}c_{12} - c_{23}s_{12}s_{13}e^{i\delta} \\ s_{13}e^{-i\delta} & s_{23}c_{13} & c_{23}c_{13} \end{pmatrix}$$

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2 Majorana Phases

3 Mixing Angles, 1 CP Phase

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Global Fit to Neutrino Oscillation Data:

D. V. Forero, M. Tortola and J. W. F. Valle, Phys. Rev. D **86** (2012) 073012

$$\Delta m_{21}^2 \equiv m_2^2 - m_1^2 = 7.62_{-0.19}^{+0.19} \times 10^{-5} \text{eV}^2 \quad |\Delta m_{31}^2| \equiv |m_3^2 - m_1^2| = 2.53_{-0.1}^{+0.08} \times 10^{-3} \text{eV}^2$$

$$s_{12}^2 = 0.320_{-0.017}^{+0.015} \quad s_{23}^2 = 0.49_{-0.05}^{+0.08} \quad s_{13}^2 = 0.026_{-0.004}^{+0.003}$$

m_1, α_1, α_2 Not Constrained

δ Beyond Exp. Sensitivity

“Large” θ_{13}

Experimental Constrains & Model Predictions

① Neutrino Masses & Mixings

Cocktail Model
Prediction



$$m_{\nu_{ab}} \propto m_{l_a} m_{l_b} C_{ab} \simeq \begin{pmatrix} 0 & 0 & m_{e\tau} \\ 0 & m_{\mu\mu} & m_{\mu\tau} \\ m_{e\tau} & m_{\mu\tau} & m_{\tau\tau} \end{pmatrix}$$

Approximate ν Mass Texture

Global Fit to Neutrino Oscillation Data:

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"Large" θ_{13}



Experimental Constrains & Model Predictions

① Neutrino Masses & Mixings

Cocktail Model
Prediction



$$m_{\nu_{ab}} \propto m_{l_a} m_{l_b} C_{ab} \simeq \begin{pmatrix} 0 & 0 & m_{e\tau} \\ 0 & m_{\mu\mu} & m_{\mu\tau} \\ m_{e\tau} & m_{\mu\tau} & m_{\tau\tau} \end{pmatrix}$$

Approximate ν Mass Texture

→ *Texture Predicts Normal Hierarchy*

Global Fit to Neutrino Oscillation Data:

*D. V. Forero, M. Tortola and J. W. F. Valle, Phys. Rev. D **86** (2012) 073012*

$$\Delta m_{21}^2 \equiv m_2^2 - m_1^2 = 7.62_{-0.19}^{+0.19} \times 10^{-5} \text{eV}^2 \quad |\Delta m_{31}^2| \equiv |m_3^2 - m_1^2| = 2.53_{-0.1}^{+0.08} \times 10^{-3} \text{eV}^2$$

$$s_{12}^2 = 0.320_{-0.017}^{+0.015} \quad s_{23}^2 = 0.49_{-0.05}^{+0.08} \quad s_{13}^2 = 0.026_{-0.004}^{+0.003}$$

m_1, α_1, α_2 Not Constrained

δ Beyond Exp. Sensitivity

"Large" θ_{13}



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Approximate \mathcal{V} Mass Texture

→ Texture Imposes $Re(m_{ee}) \simeq 0, Im(m_{ee}) \simeq 0, Re(m_{e\mu}) \simeq 0, Im(m_{e\mu}) \simeq 0$

Prediction for $m_1, \alpha_1, \alpha_2, \delta$

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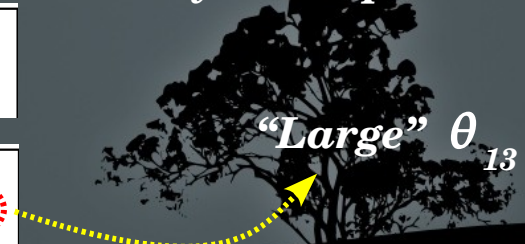
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"Large" θ_{13}



Experimental Constrains & Model Predictions

① Neutrino Masses & Mixings

Cocktail Model
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Prediction for $\theta_{13}!!$

m_1, α_1, α_2 Not Constrained

δ Beyond Exp. Sensitivity

"Large" θ_{13}



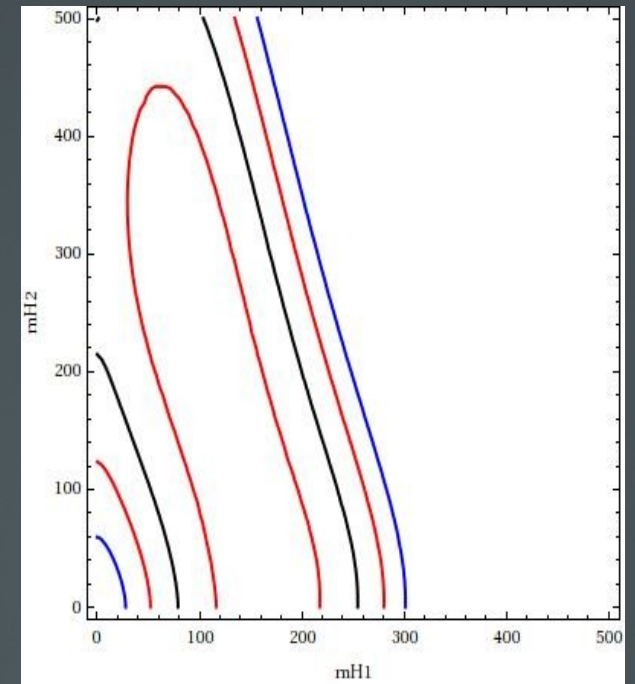
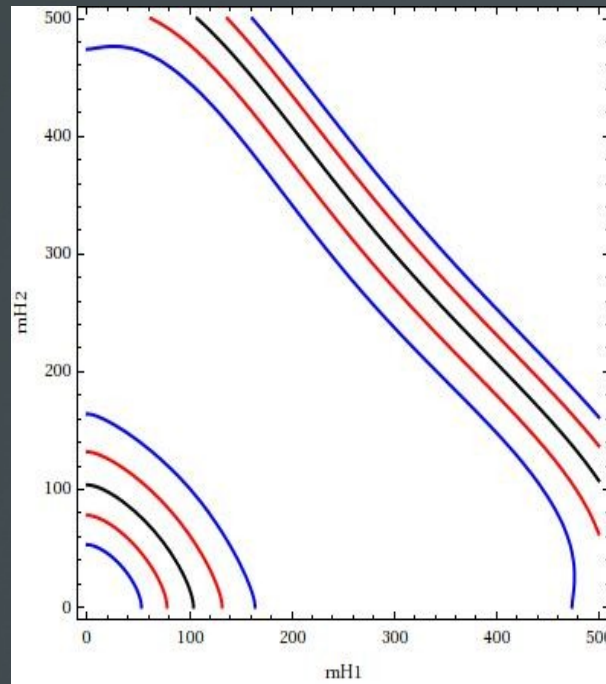
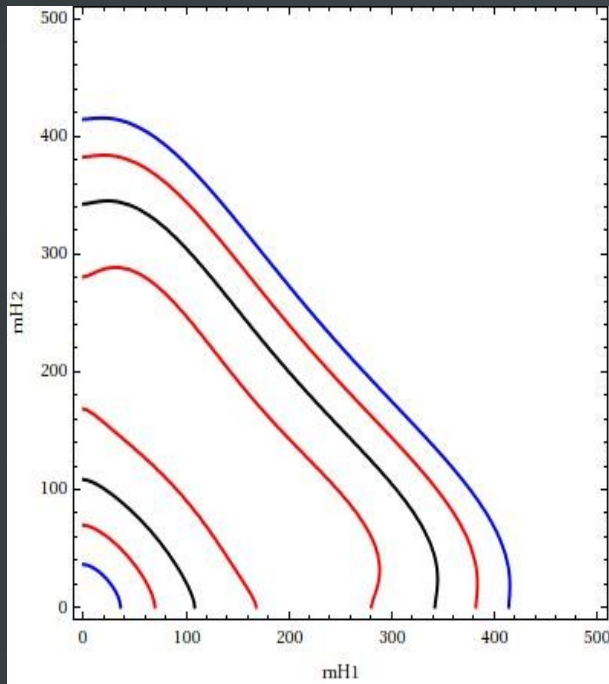
Experimental Constrains & Model Predictions

② Electroweak Precision Constrains

$H_0, A_0, H_1^+, H_2^+, \rho^{++}$ Contribute to EWPO (S, T Parameters)

$$\Delta T \sim \frac{1}{24 \pi^2 \alpha_{EM} v^2} [(m_{H_1}^2 - m_{A_0}^2)(m_{H_1}^2 - m_{H_0}^2) + (m_{H_2}^2 - m_{A_0}^2)(m_{H_2}^2 - m_{H_0}^2)] + \Delta T_\rho$$

$$m_{H_0} = m_{DM} = 65 \text{ GeV}, m_\rho = 5 \text{ TeV}$$



$$m_{A_0} = 200 \text{ GeV}, \theta = \pi/4$$

$$m_{A_0} = 300 \text{ GeV}, \theta = \pi/4$$

$$m_{A_0} = 200 \text{ GeV}, \theta = \pi/6$$

Experimental Constrains & Model Predictions

③ Lepton Flavour Violation Constrains

$\mu \rightarrow 3e$	$C_{e\mu} C_{ee} < 1.2 \times 10^{-11} \text{GeV}^{-2} m_\rho^2$
$\tau \rightarrow 3e$	$C_{e\tau} C_{ee} < 1.3 \times 10^{-8} \text{GeV}^{-2} m_\rho^2$
$\tau \rightarrow 3\mu$	$C_{\mu\tau} C_{\mu\mu} < 1.2 \times 10^{-8} \text{GeV}^{-2} m_\rho^2$
$\tau^- \rightarrow \mu^+ e^- e^-$	$C_{\mu\tau} C_{ee} < 9.3 \times 10^{-9} \text{GeV}^{-2} m_\rho^2$
$\tau^- \rightarrow e^+ e^- \mu^-$	$C_{e\tau} C_{e\mu} < 1.7 \times 10^{-8} \text{GeV}^{-2} m_\rho^2$
$\tau^- \rightarrow \mu^+ \mu^- e^-$	$C_{\mu\tau} C_{e\mu} < 1.8 \times 10^{-8} \text{GeV}^{-2} m_\rho^2$
$\tau^- \rightarrow e^+ \mu^- \mu^-$	$C_{e\tau} C_{\mu\mu} < 1.0 \times 10^{-8} \text{GeV}^{-2} m_\rho^2$
$\mu \rightarrow e\gamma$	$\sum C_{l\mu} C_{le} < 4.7 \times 10^{-9} \text{GeV}^{-2} m_\rho^2$
$\tau \rightarrow e\gamma$	$\sum C_{l\tau} C_{le} < 1.05 \times 10^{-6} \text{GeV}^{-2} m_\rho^2$

Most Stringent Constrains are on $m_{e\tau}$ and $m_{\mu\mu}$:

→ ① + ③ Set $m_\rho^2 > \text{TeV}$, $m_{A_0}^2 - m_{H_0}^2 \sim V^2$, $m_{H_1}^2 - m_{H_2}^2 \sim V^2$ and $C_{e\tau} \sim 1$

Need Heavy ρ^{++} & Large Mass Splittings

Experimental Constrains & Model Predictions

④ *Collider Phenomenology*

→ *Modification of $h \rightarrow \gamma\gamma$ Branching Ratio*

→ *H_0, A_0, H_1^+, H_2^+ States Accesible at LHC*



Experimental Constrains & Model Predictions

④ *Collider Phenomenology*

→ *Modification of $h \rightarrow \gamma\gamma$ Branching Ratio*

→ *H_0, A_0, H_1^+, H_2^+ States Accesible at LHC*

⑤ *Neutrinoless Double β -Decay*

→ *$m_{ee} \sim 0$, But Still Potentially Large $0\nu\beta\beta$!!*

→ *Leading Contribution to $0\nu\beta\beta$ Independent of M_ν*

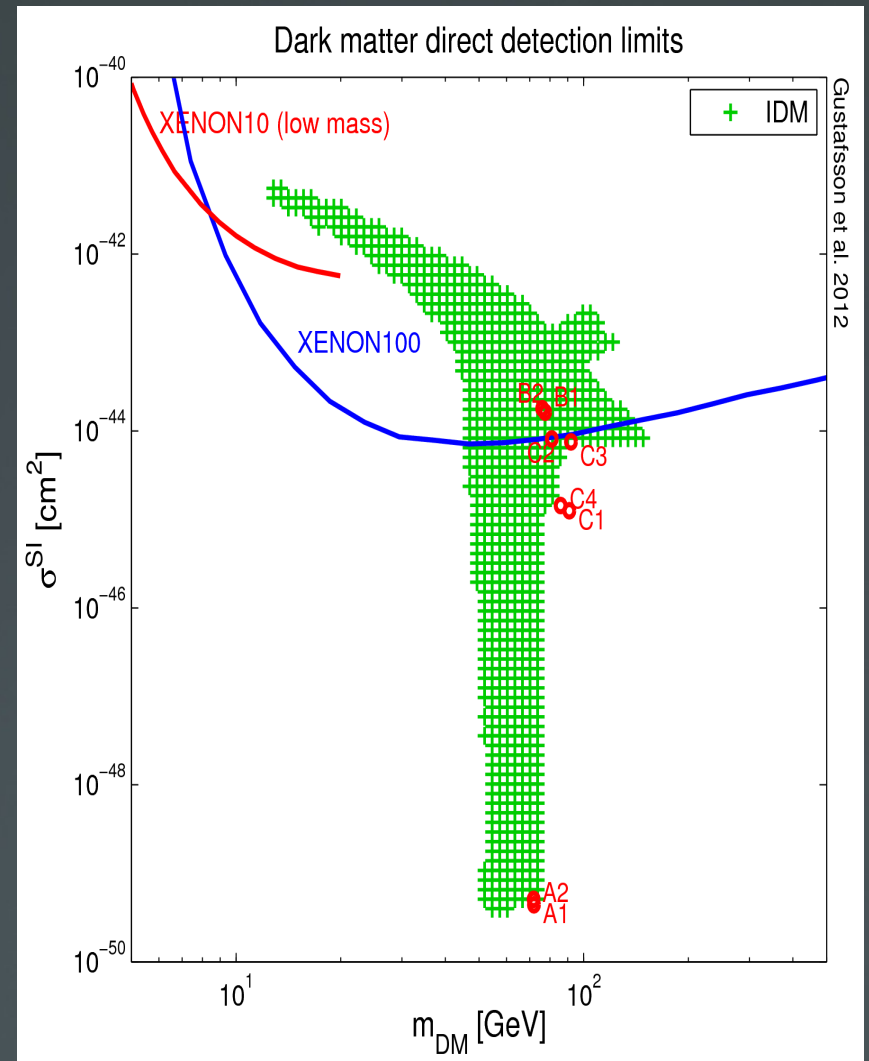


Experimental Constrains & Model Predictions

⑥ *Dark Matter (IDM)*

→ *Large Mass Splitting $m_{A_0}^2 - m_{H_0}^2$
Disfavours Co-Annihilation*

→ *Direct Detection Bounds Impose m_{DM}
Close to Higgs Resonance or W Threshold*



Conclusions

Small ν Masses May be Connected to Dark Matter Stability

Properties of ν Sector and DM Sector Affected by this Connection:

→ ν_R *as DM Candidate* (Ma; Krauss, Nasri, Trodden)

NO Dark Matter Direct Detection Signal

→ *Cocktail Model*

→ *Predicts Normal ν Mass Hierarchy & $0.011 < s_{13}^2 < 0.035$*

→ *LFV Close to Current Experimental Limits*

→ *Collider Signatures at LHC & ILC (New charged States)*

→ *(Possibly) Large $0\nu\beta\beta$*

→ *DM Co-Annihilation Disfavoured ($m_{DM} \sim 60 - 65 \text{ GeV?}$)*

