

# Jet Tomography at Lower Energies and the temperature dependence of Jet transport coefficients

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Wayne State University

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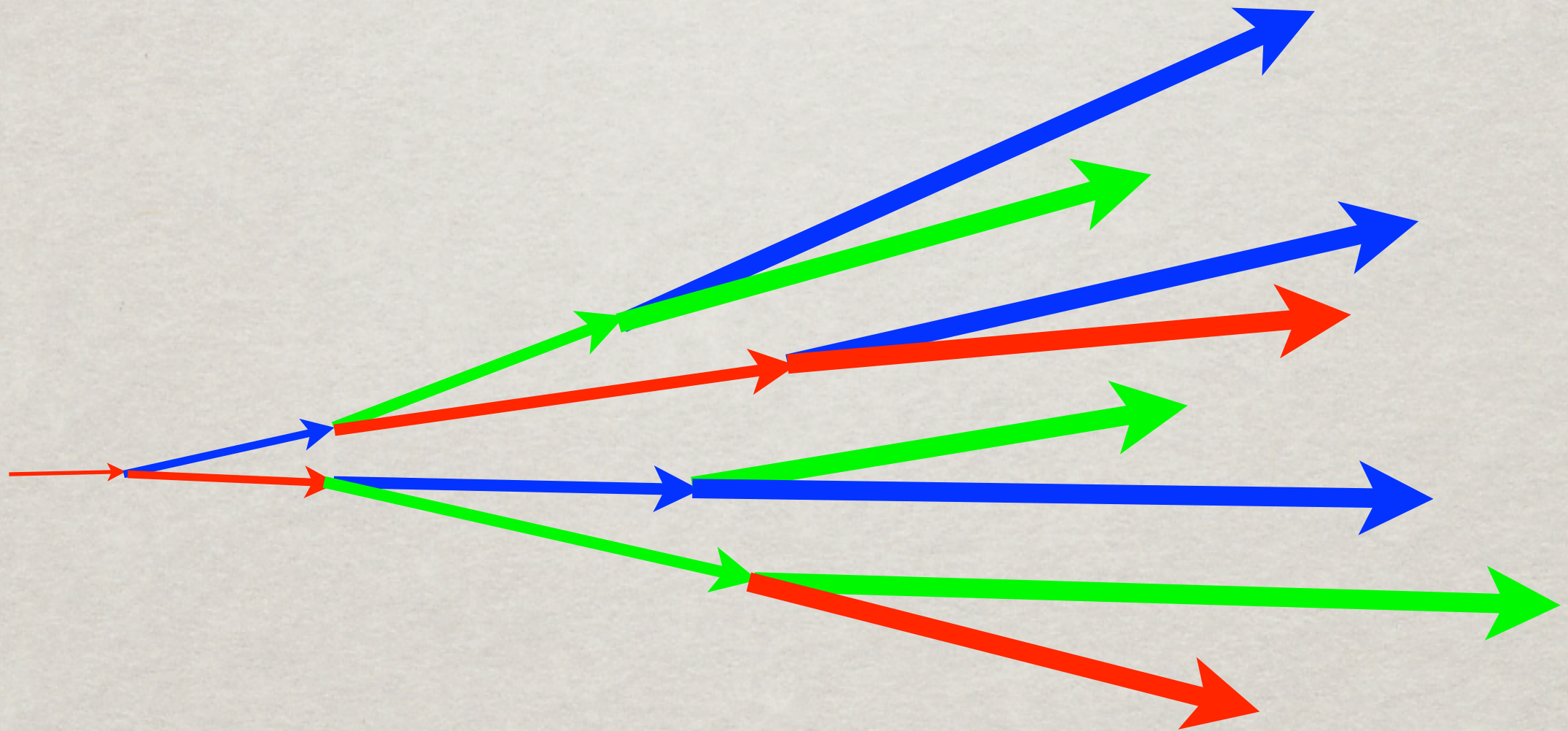
Annual meeting of the Division of Nuclear Physics of the APS, New Port Beach, CA,



# Jet modification in simple pictures

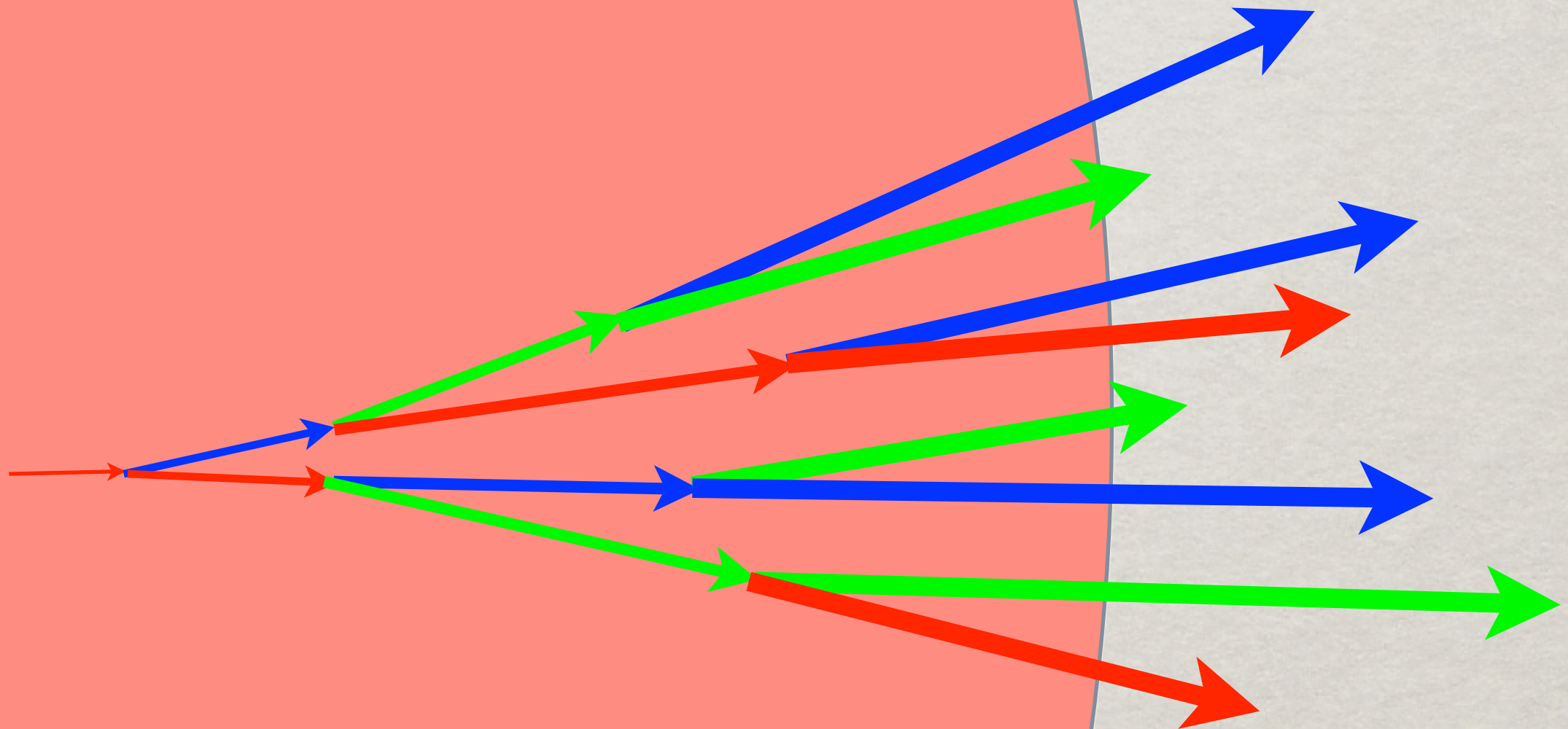


# Jet modification in simple pictures

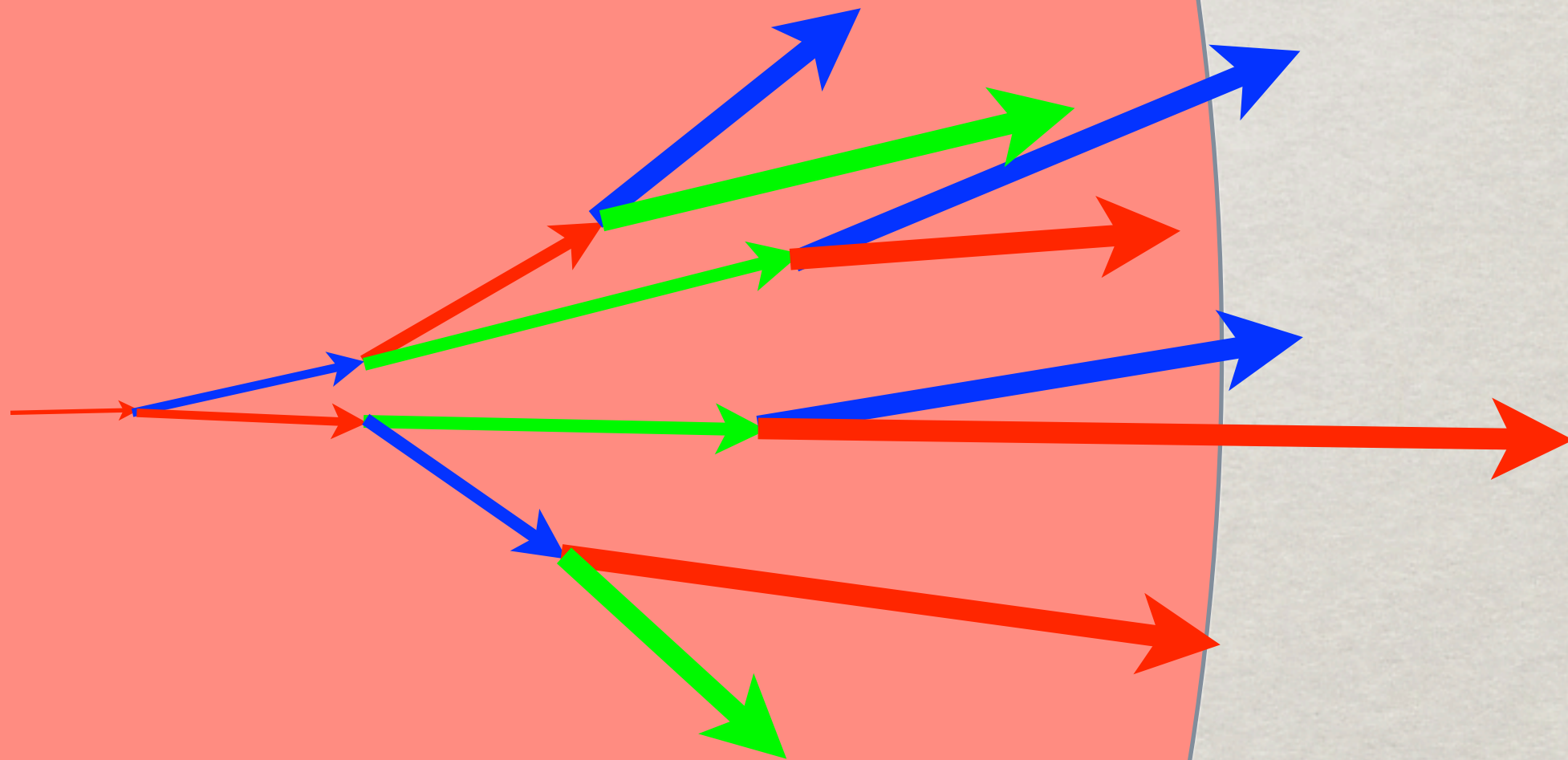




# Jet modification in simple pictures

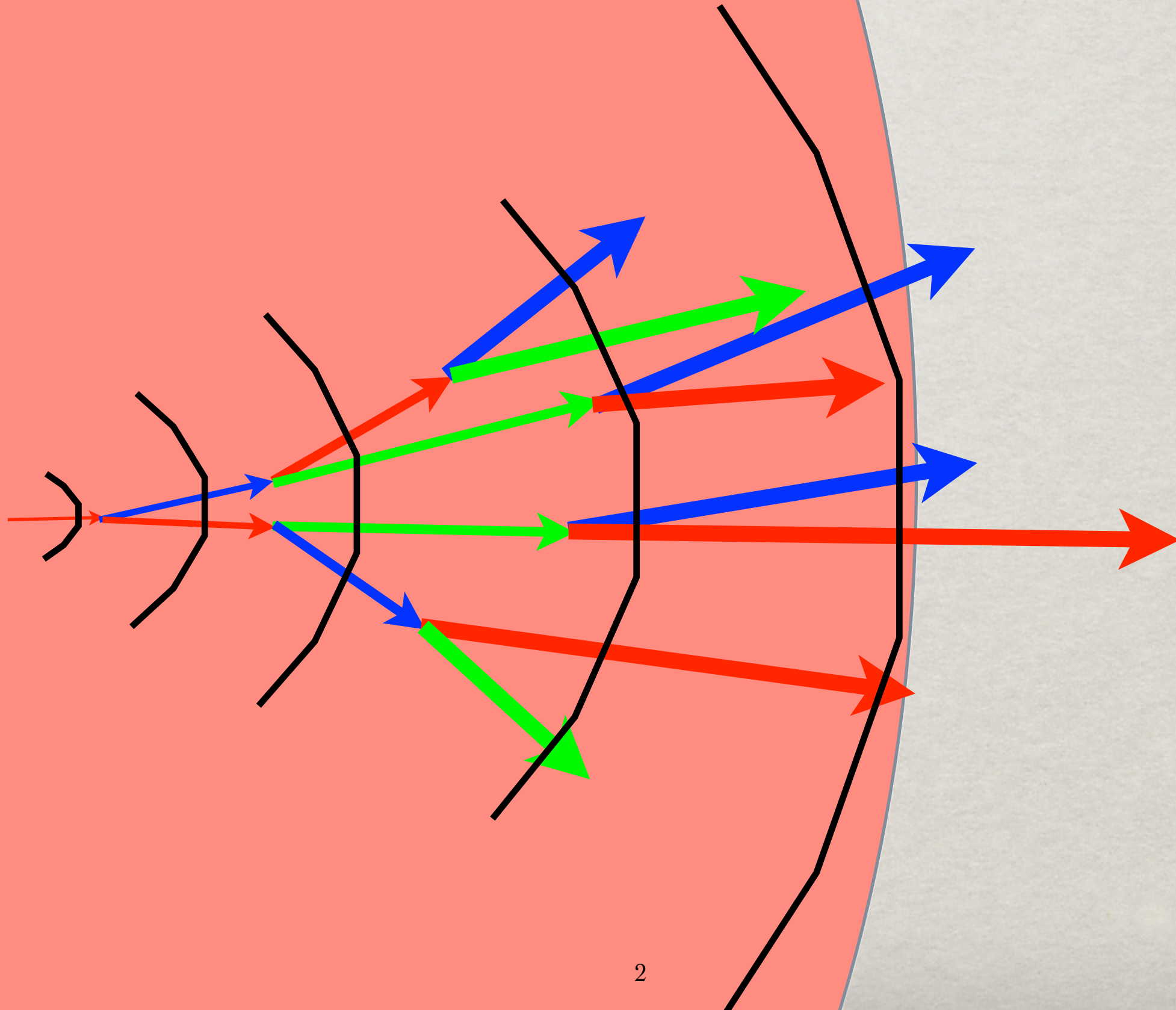


# Jet modification in simple pictures

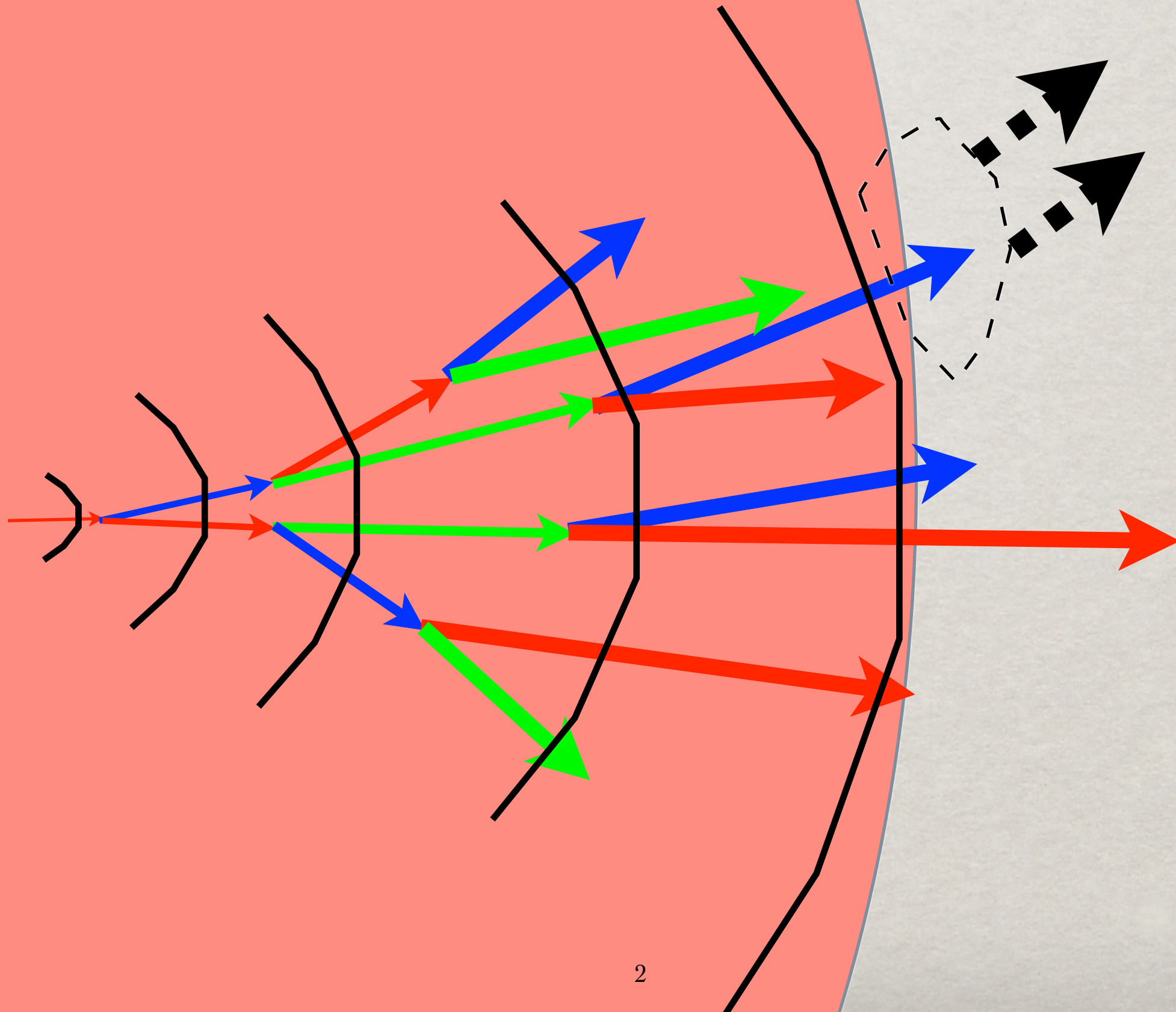




# Jet modification in simple pictures

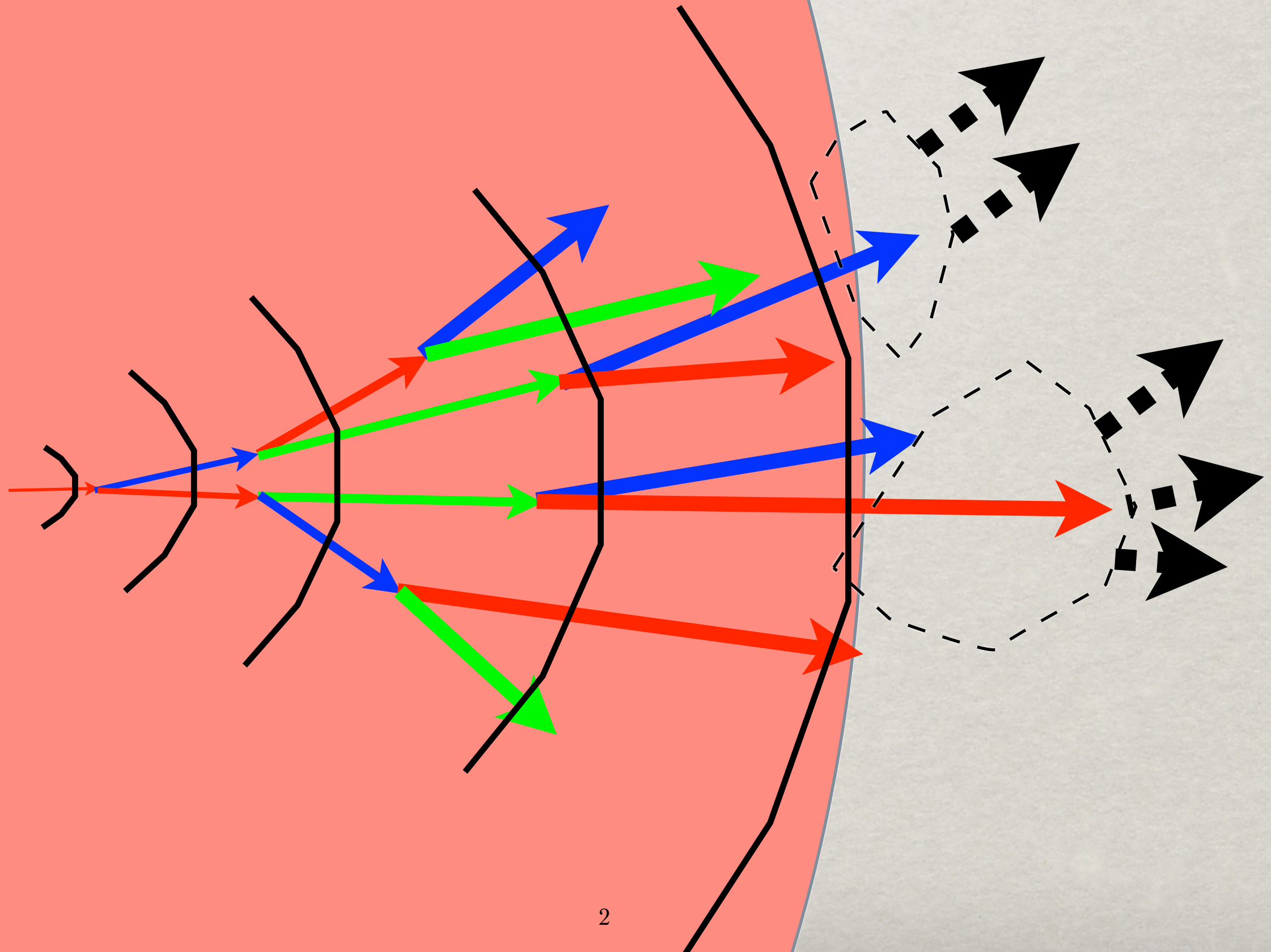


# Jet modification in simple pictures





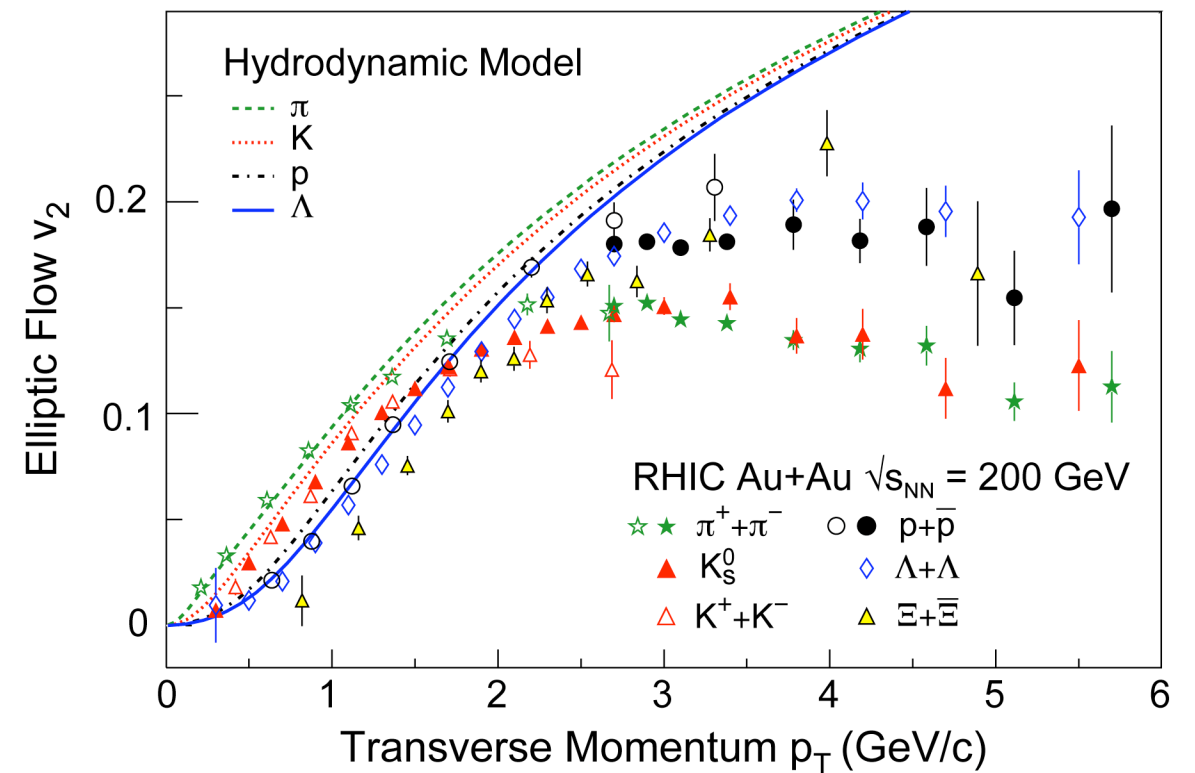
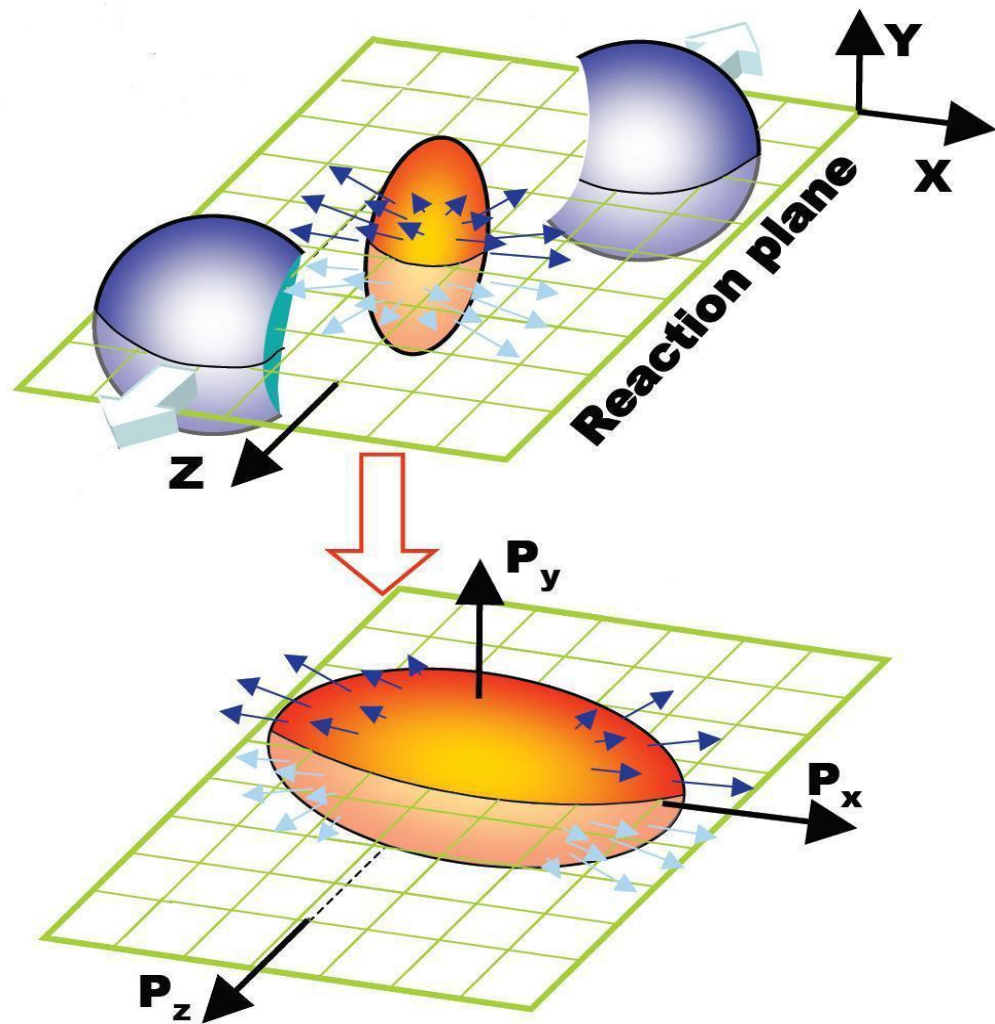
# Jet modification in simple pictures





# The current paradigm for high energy HIC

Soft medium described by viscous fluid dynamics,

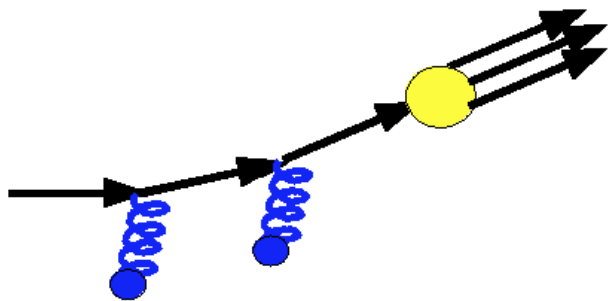


Ideal Hydro calculations  
by P. Houvinen, also  
Kolb and Heinz, Teaney, Lauret,  
Shuryak, and Nonaka and Bass



# Hard jets described by pQCD

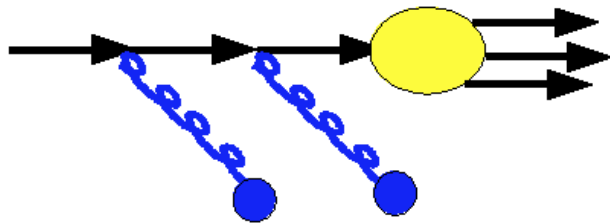
## with factorized transport coefficients



$$D\left(\frac{\vec{p}_h}{|\vec{p} + \vec{k}_\perp|}, m_J^2\right)$$

$$\hat{q} = \frac{\langle p_\perp^2 \rangle L}{L}$$

Transverse momentum  
diffusion rate



$$D\left(\frac{p_h}{p - k}, m_J^2\right)$$

$$\hat{e} = \frac{\langle p_z \rangle L}{L}$$

Elastic energy loss rate  
also diffusion rate  $e_2$



$$D_g\left(\frac{p_h}{p + k}, m_J^2\right)$$

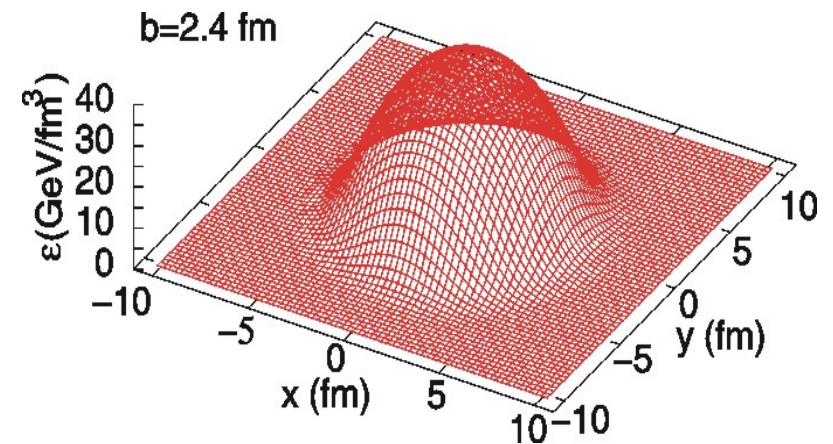
$$\hat{f} = \frac{\langle \Delta N \rangle L}{L}$$

Flavor ( $q \leftrightarrow g$ )  
diffusion rate



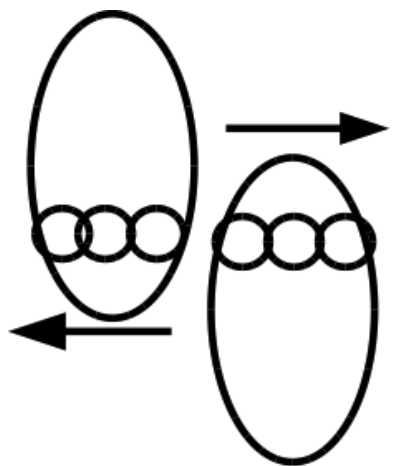
# Non-perturbative transport coefficients represent medium's influence in jet quenching calculations

Fit the  $\hat{q}$  for the initial T  
in the hydro in central coll.

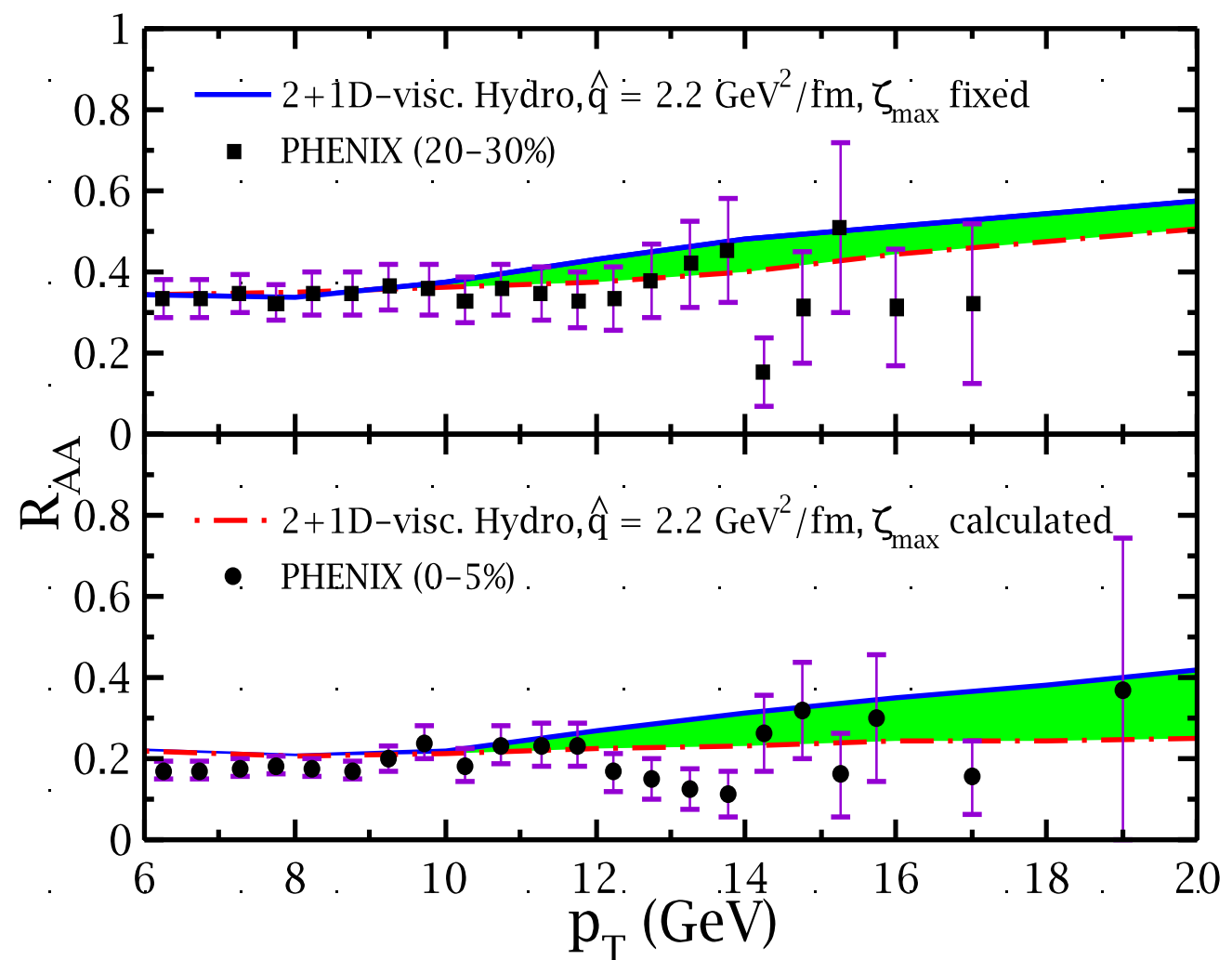


$$\hat{q}(\vec{r}, t) = \hat{q}_0 \frac{s(\vec{r}, t)}{s_0}$$

$$s_0 = s(T_0)$$



$$R_{AA} \sim \frac{\frac{dN_{AA}}{dp_T dy}}{N_{bin} \frac{dN_{pp}}{dp_T dy}}$$

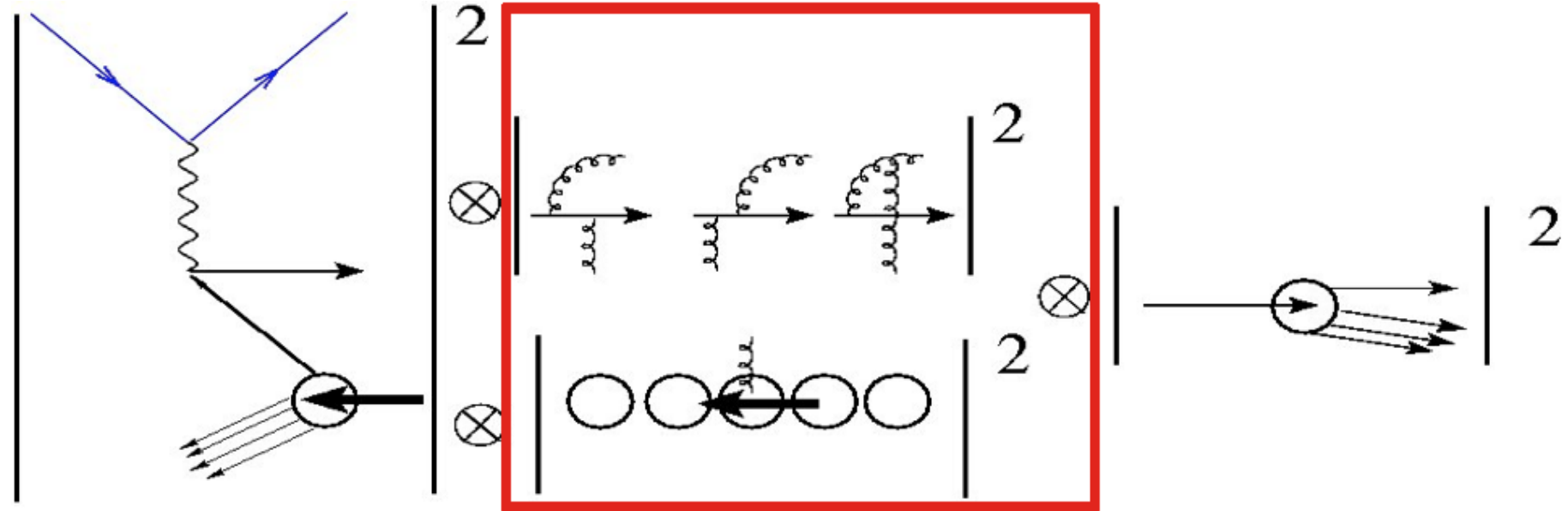




# An aside on 2 underlying assumptions

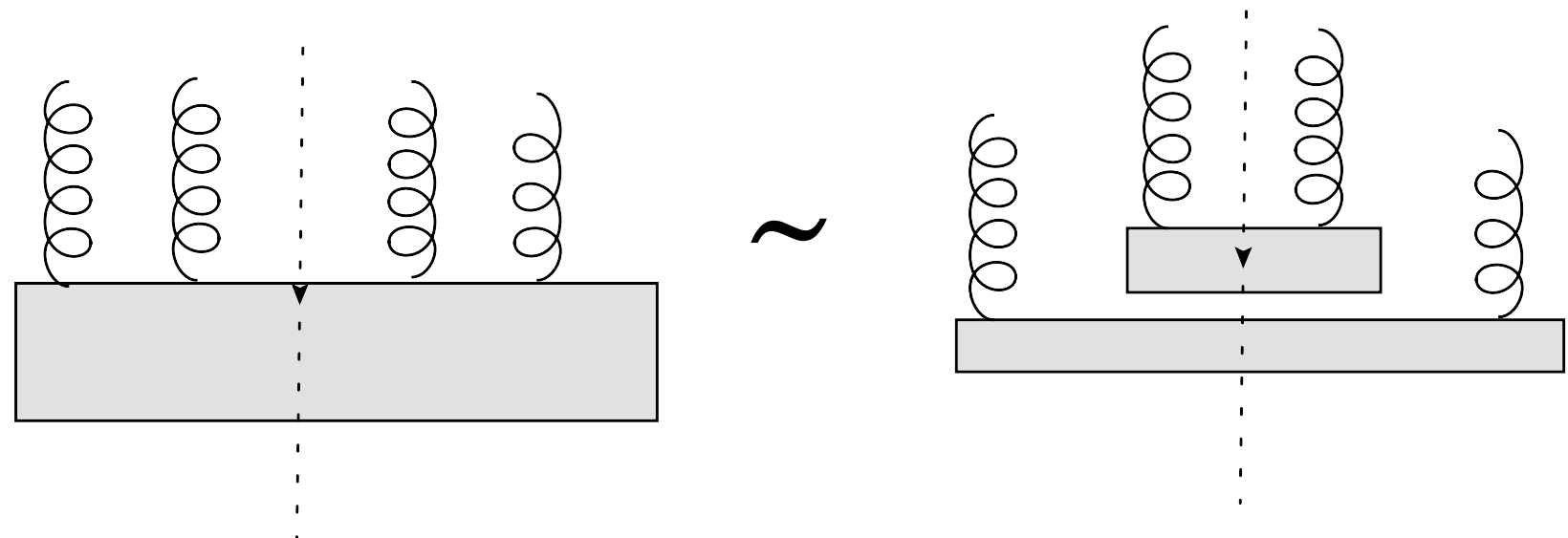
- 1) We will only consider very high  $Q^2$  jets

In the maximally factorized region



- 2) We have assumed short correlation lengths in the medium

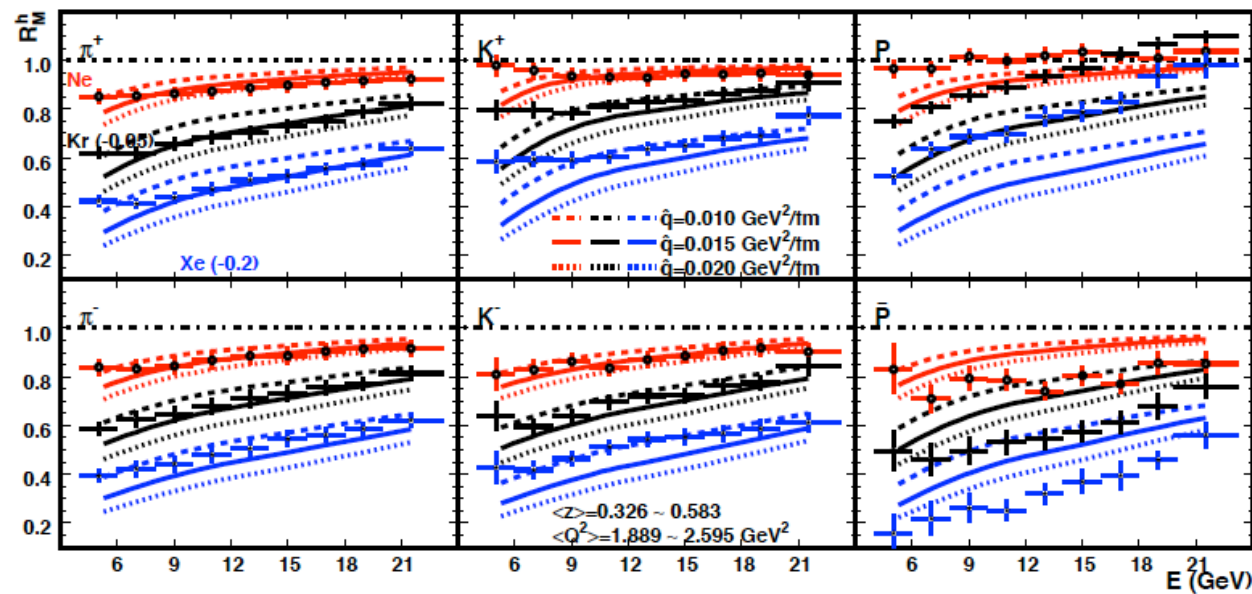
Multiple scattering  
 $2N$  point correlators  
 simplify to  
 $N$  2-point correlators



# The temperature dependence of $\hat{q}$

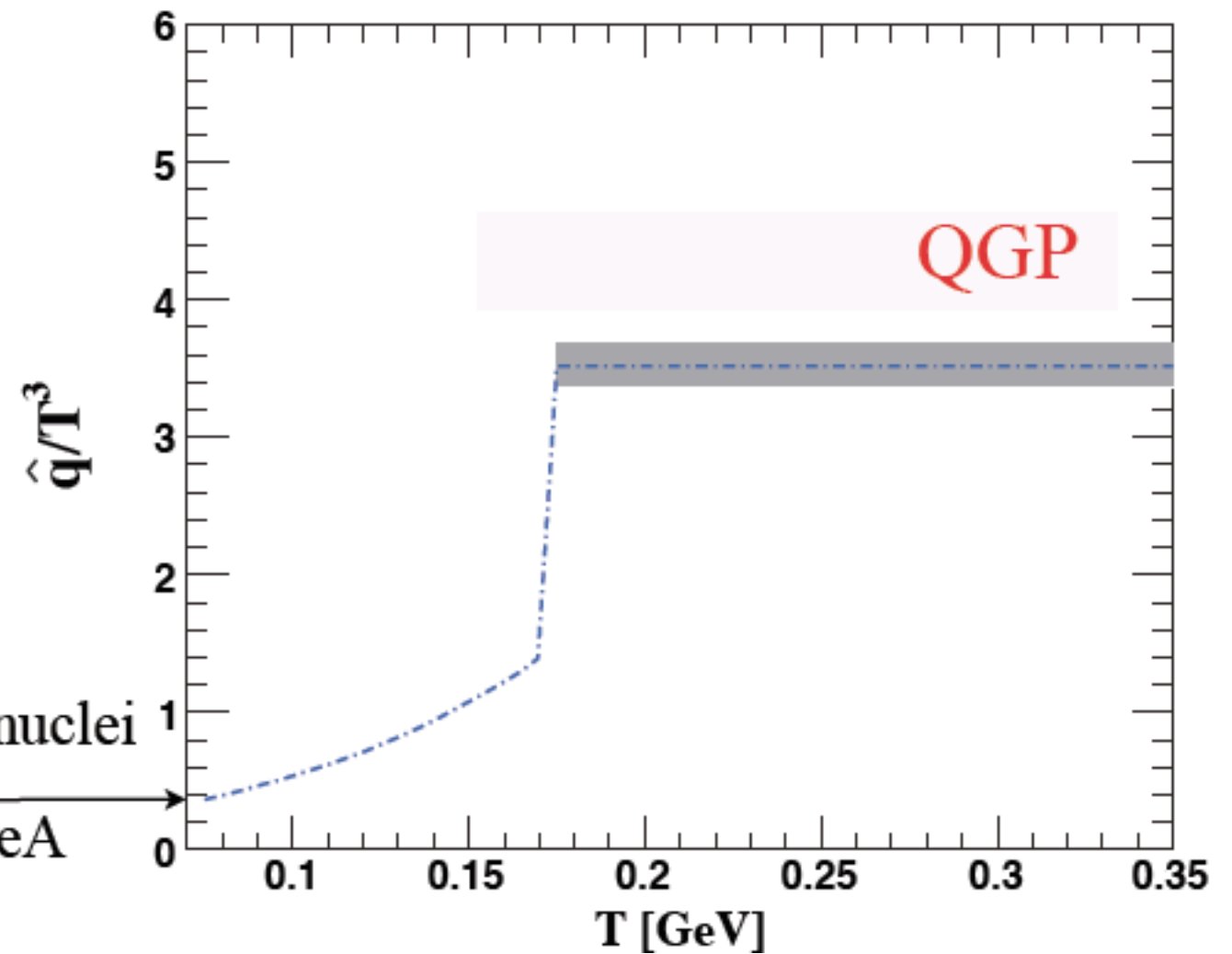
Partly known and partly guessed!

$$\hat{q} = \frac{p_{\perp}^2}{t} = \frac{2\pi^2 \alpha_s C_R}{N_c^2 - 1} \int d\tilde{t} \left\langle F^{\mu\alpha}(\tilde{t}) v_{\alpha} F_{\mu}^{\beta}(0) v_{\beta} \right\rangle$$



From Deng, Chang and Wang

cold nuclei  
DIS eA

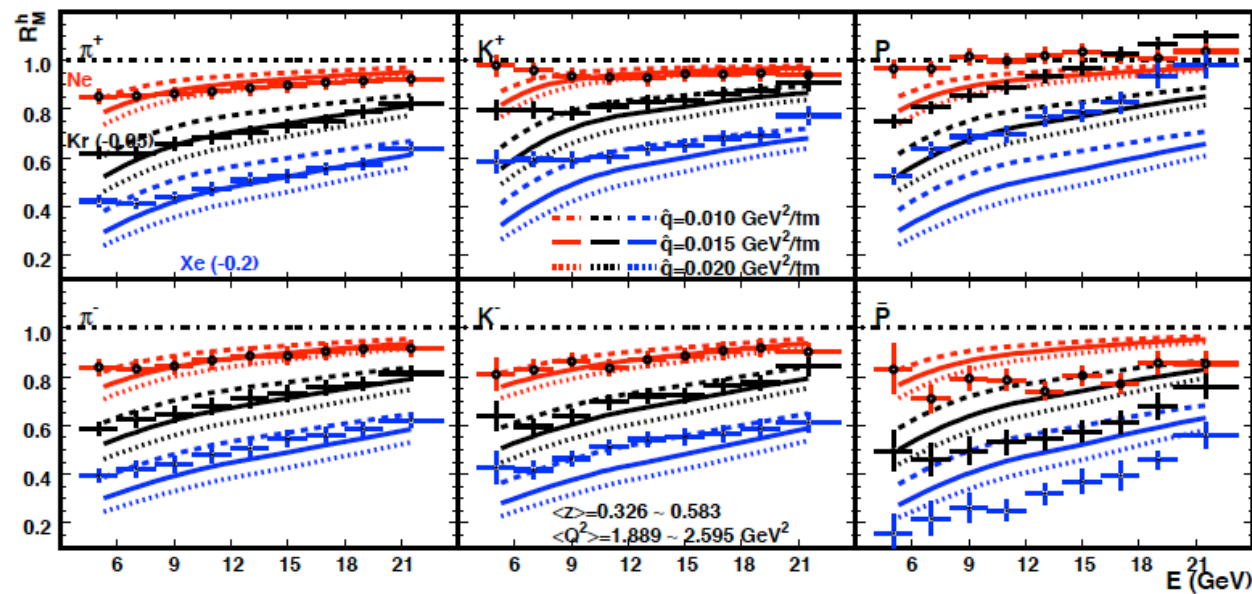


From Chen et. al.



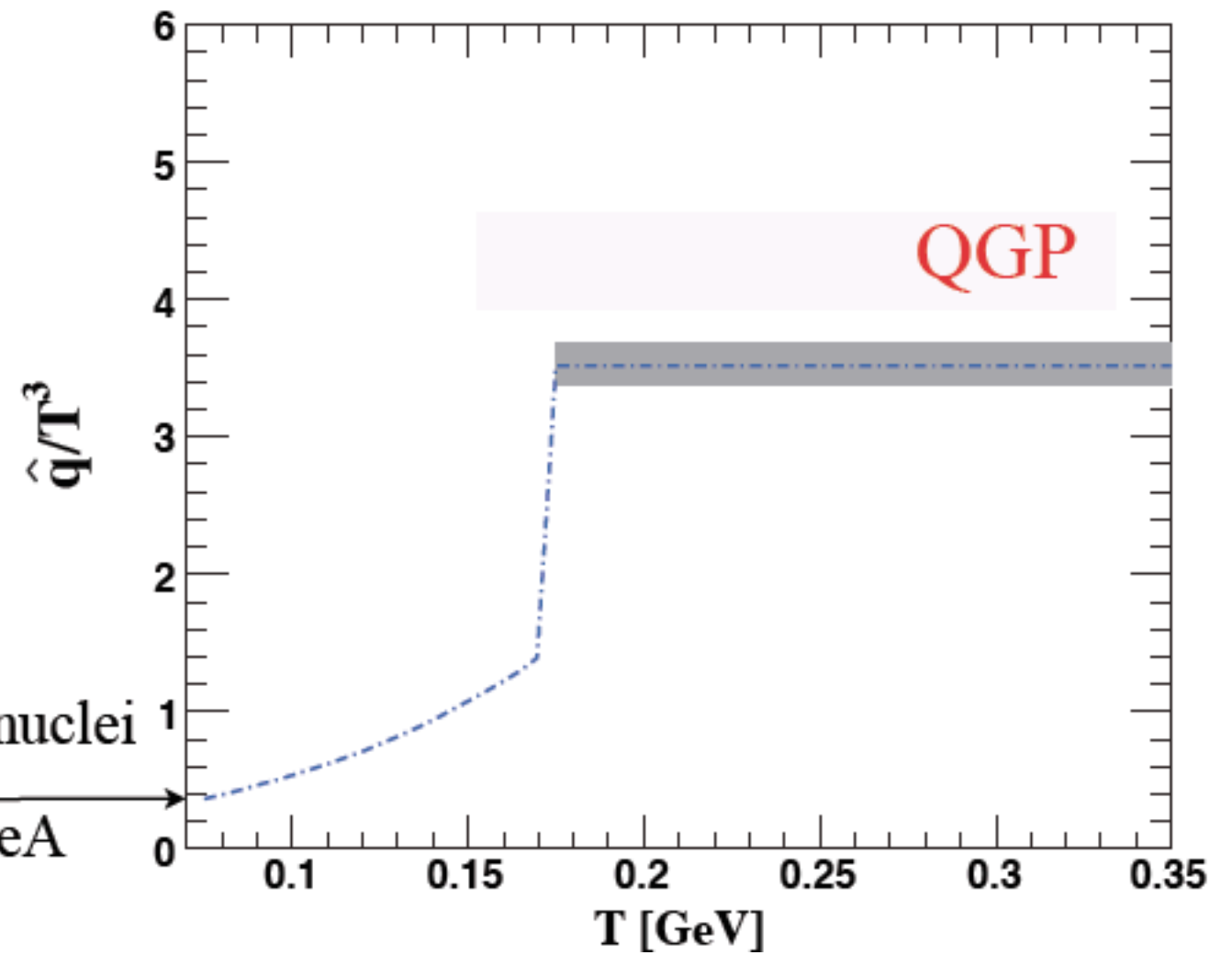
# The temperature dependence of $\hat{q}$

Partly known and partly guessed!



From Deng, Chang and Wang

cold nuclei  
DIS eA



From Chen et. al.

# The emerging picture in the temperature dependence of viscosity

PRL 97, 152303 (2006)

PHYSICAL REVIEW LETTERS

week ending  
13 OCTOBER 2006

## Strongly Interacting Low-Viscosity Matter Created in Relativistic Nuclear Collisions

Laszlo P. Csernai,<sup>1,2</sup> Joseph I. Kapusta,<sup>3</sup> and Larry D. McLerran<sup>4</sup>

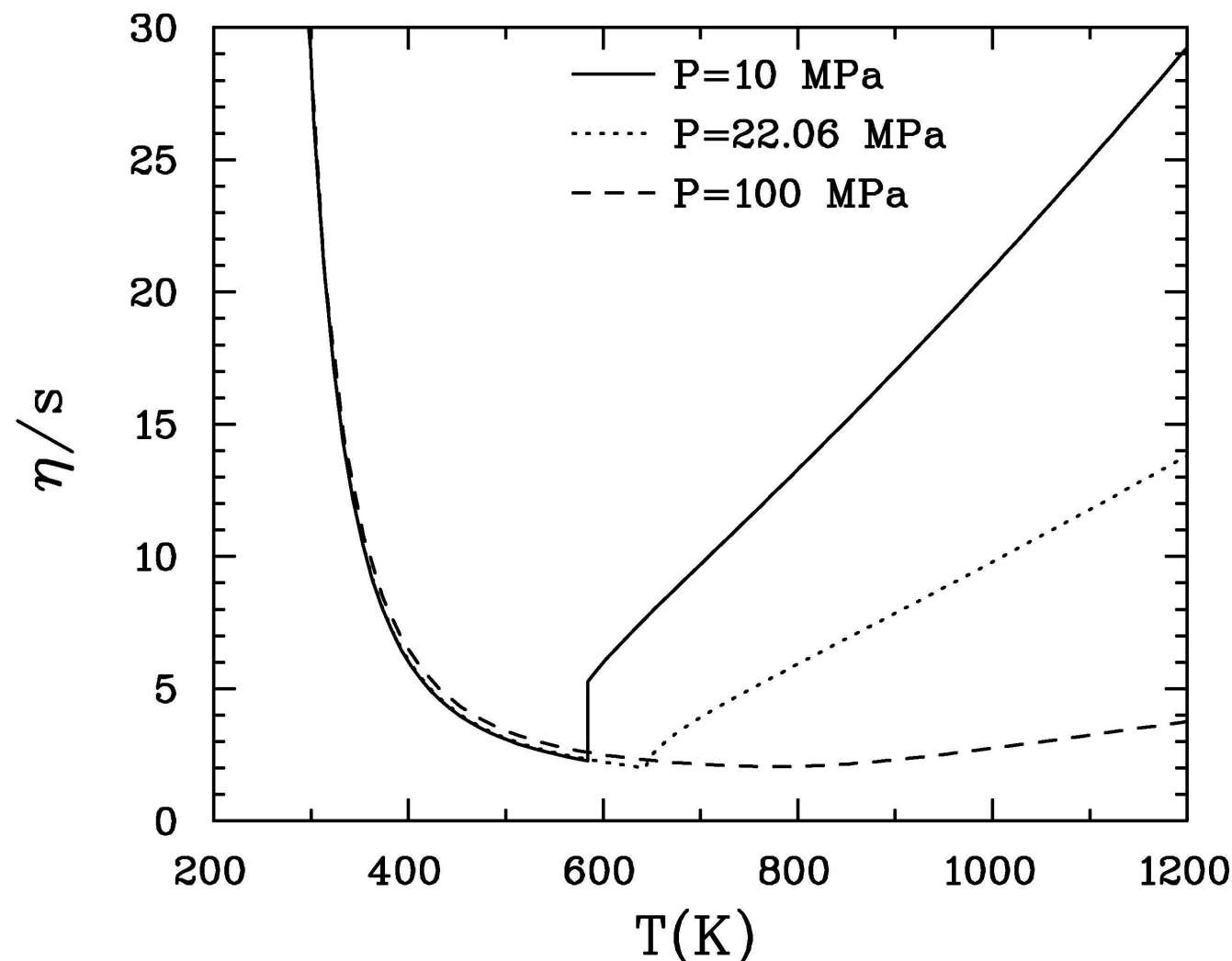
<sup>1</sup>*Section for Theoretical Physics, Department of Physics, University of Bergen, Allegaten 55, 5007 Bergen, Norway*

<sup>2</sup>*MTA-KFKI, Research Institute of Particle and Nuclear Physics, 1525 Budapest 114, P. O. Box 49, Hungary*

<sup>3</sup>*School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455, USA*

<sup>4</sup>*Nuclear Theory Group and Riken Brookhaven Center, Brookhaven National Laboratory, Bldg. 510A, Upton, New York 11973, USA*

(Received 12 April 2006; published 12 October 2006)





# Is there a relationship between $\eta$ and $\hat{q}$ ?

PRL **99**, 192301 (2007)

PHYSICAL REVIEW LETTERS

week ending  
9 NOVEMBER 2007

## Small Shear Viscosity of a Quark-Gluon Plasma Implies Strong Jet Quenching

Abhijit Majumder,<sup>1</sup> Berndt Müller,<sup>1</sup> and Xin-Nian Wang<sup>2</sup>

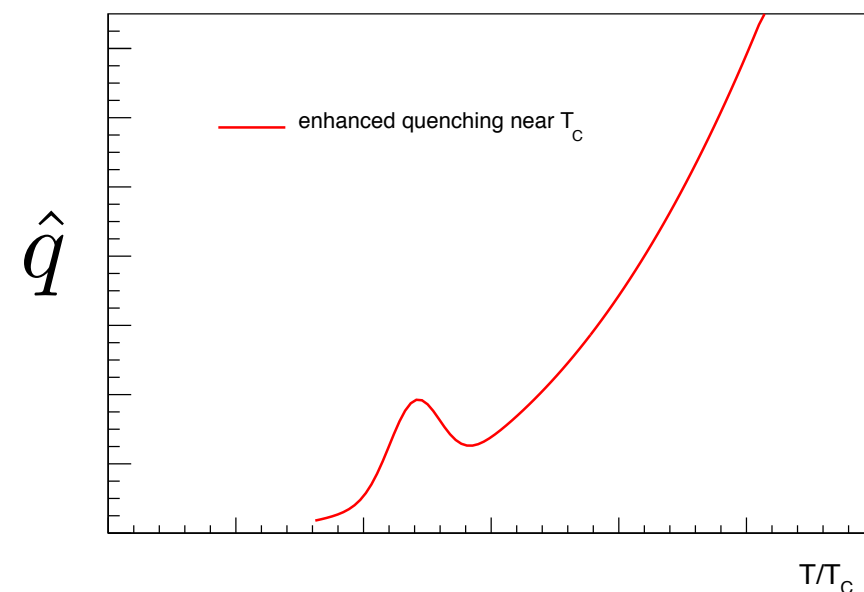
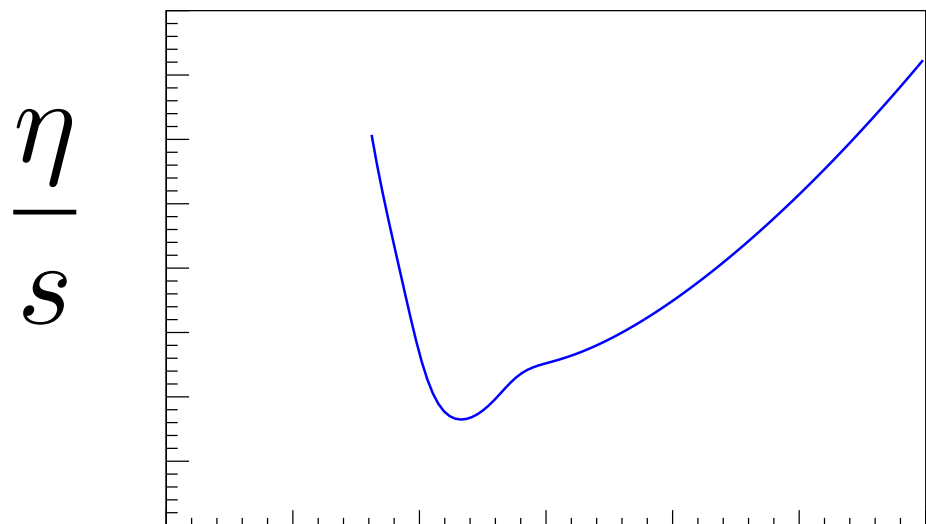
<sup>1</sup>*Department of Physics, Duke University, Durham, North Carolina 27708, USA*

<sup>2</sup>*Nuclear Science Division, MS 70R0319, Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

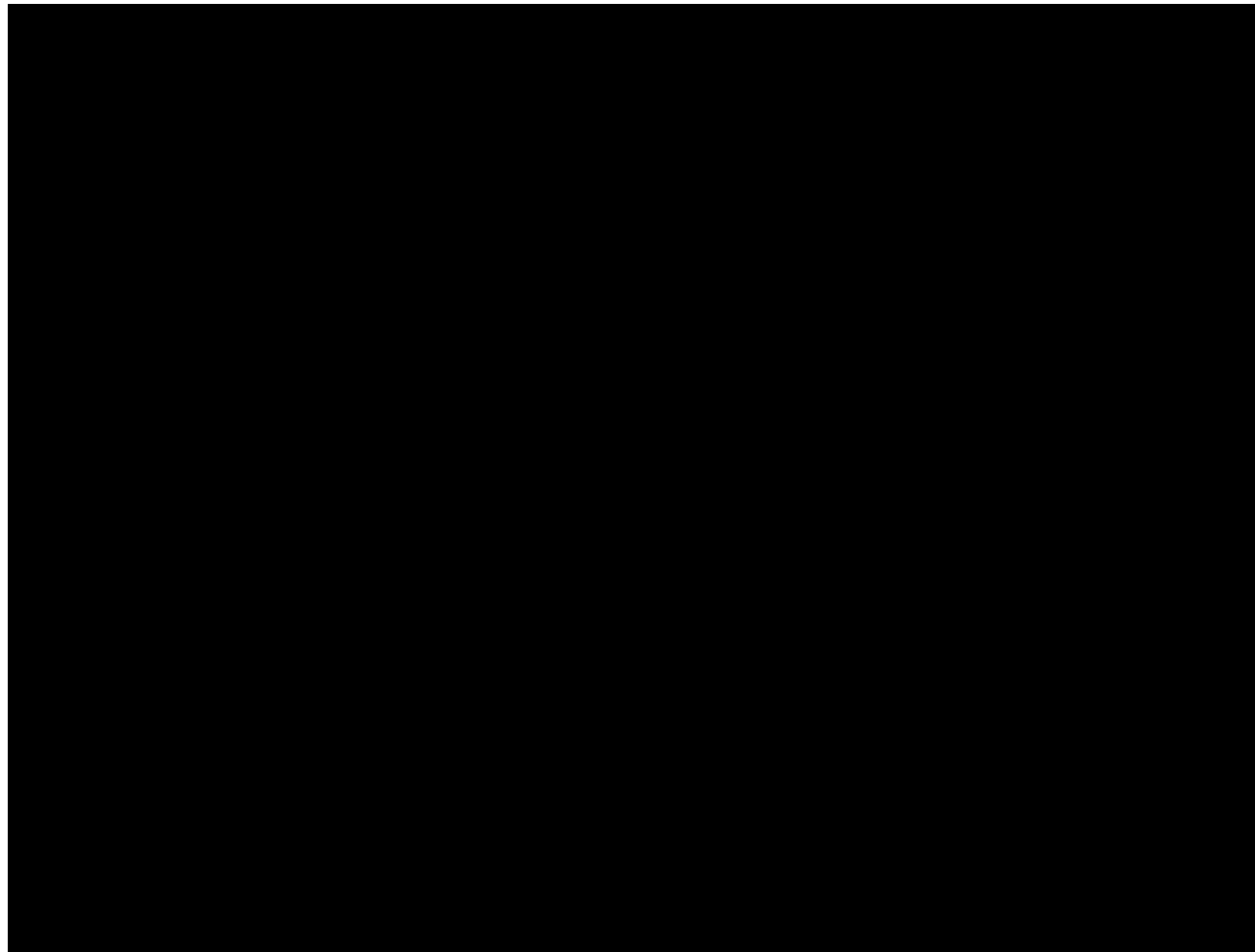
(Received 10 March 2007; revised manuscript received 13 June 2007; published 7 November 2007)

$$\frac{\eta}{s} \sim \frac{T^3}{\hat{q}}$$

For a weakly coupled medium, proportionality constant  $\sim 1$

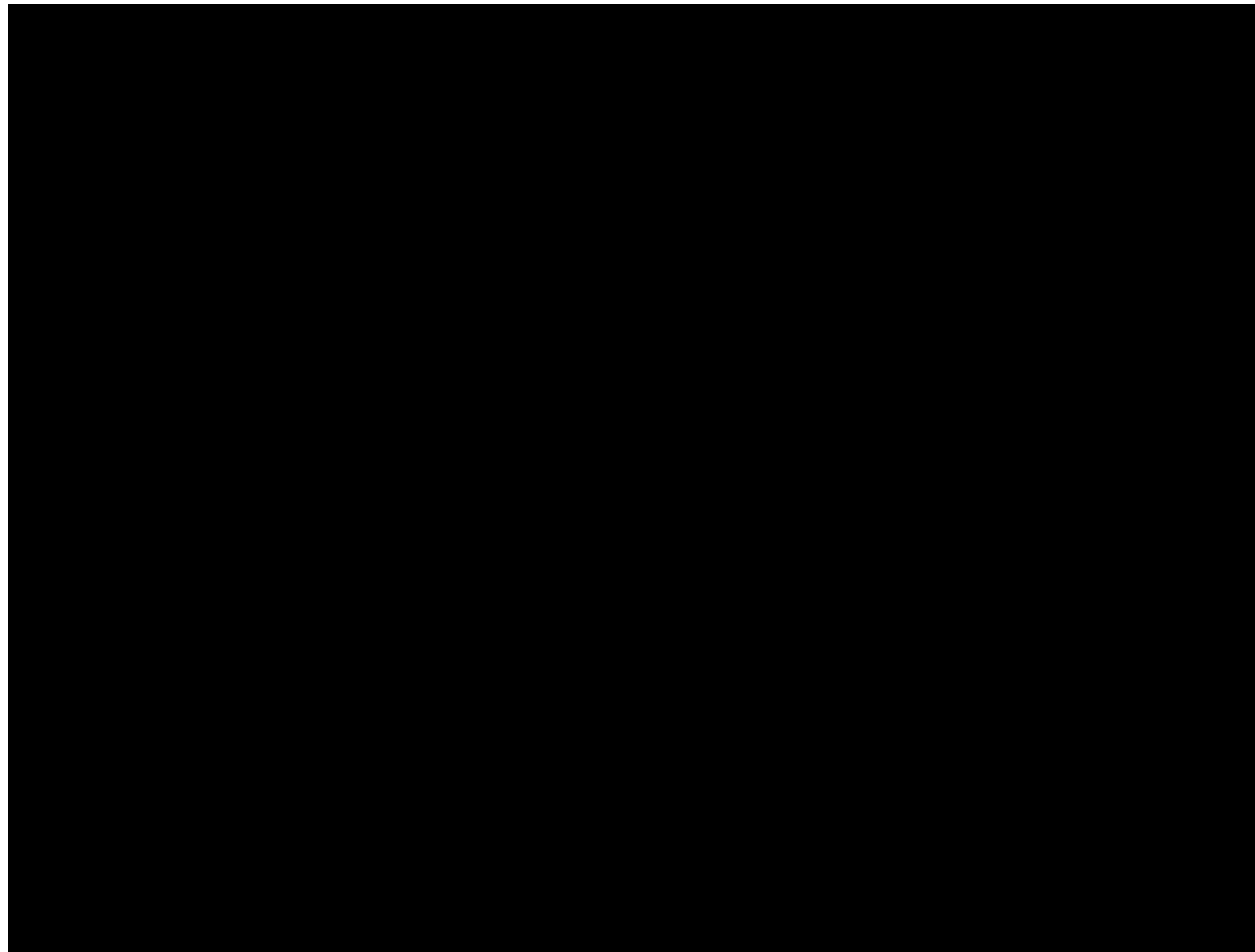


Where else in nature does something similar happen



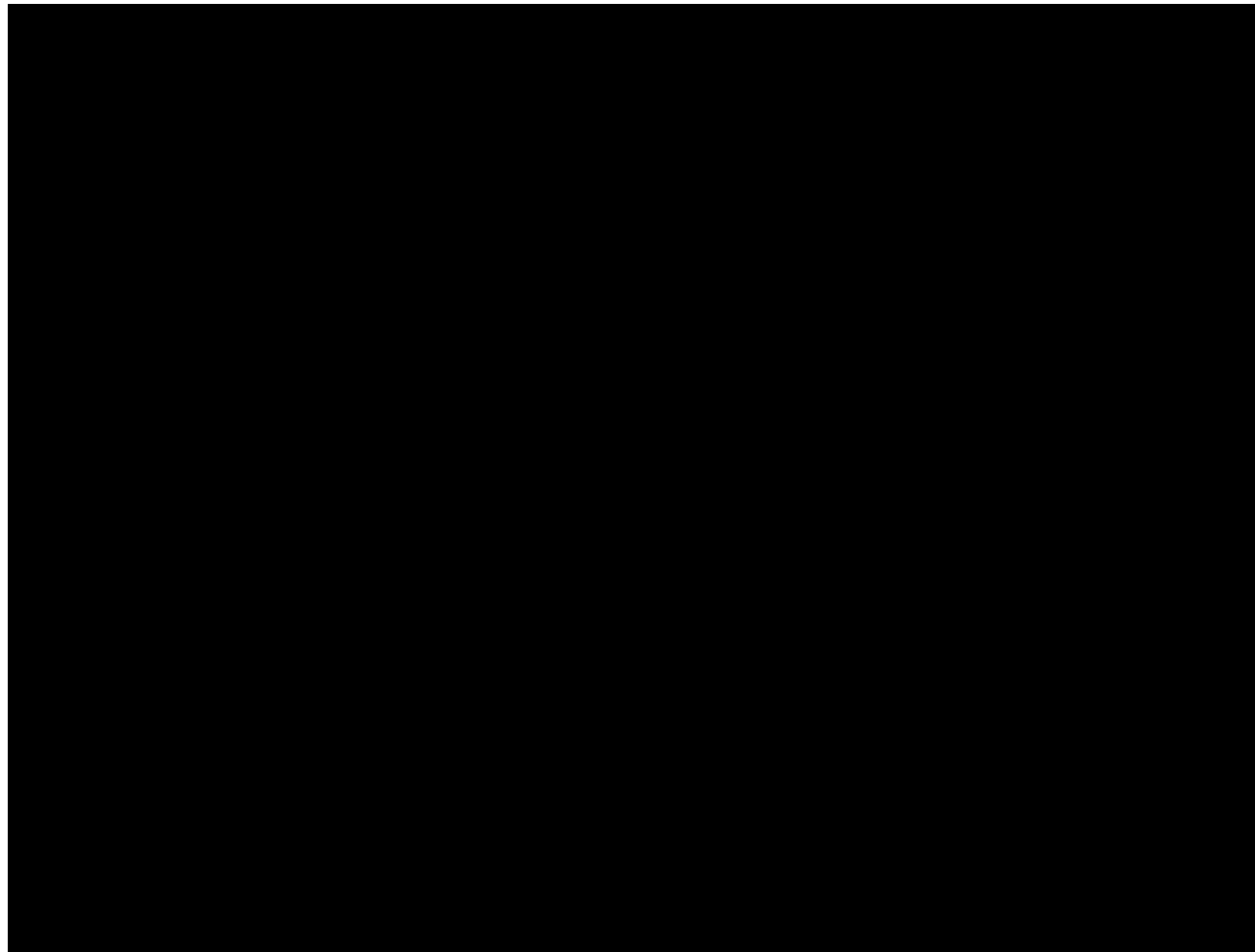


Where else in nature does something similar happen



Pretty much everywhere !

Where else in nature does something similar happen

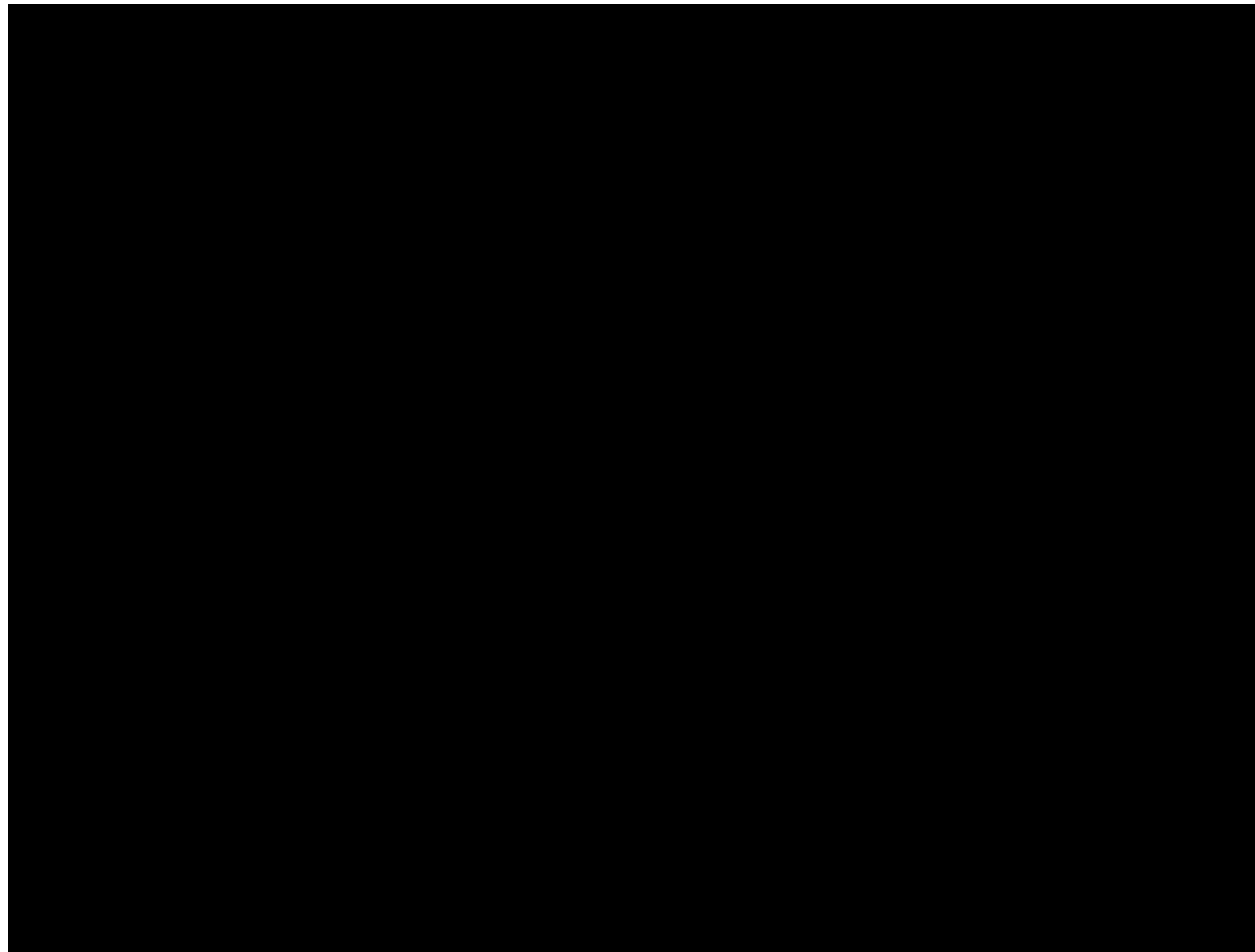


Pretty much everywhere !

Critical opalescence



Where else in nature does something similar happen

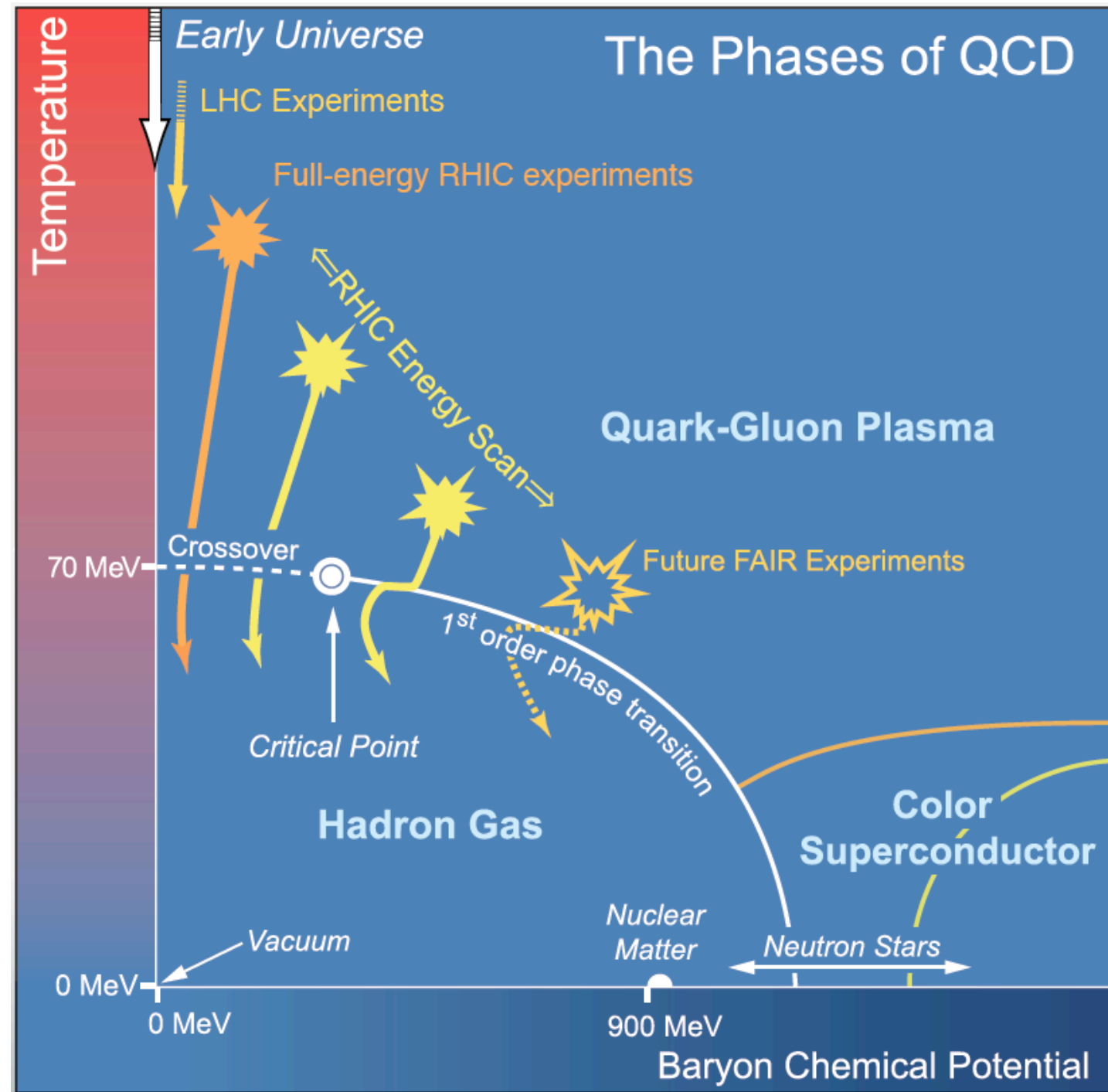


Pretty much everywhere !

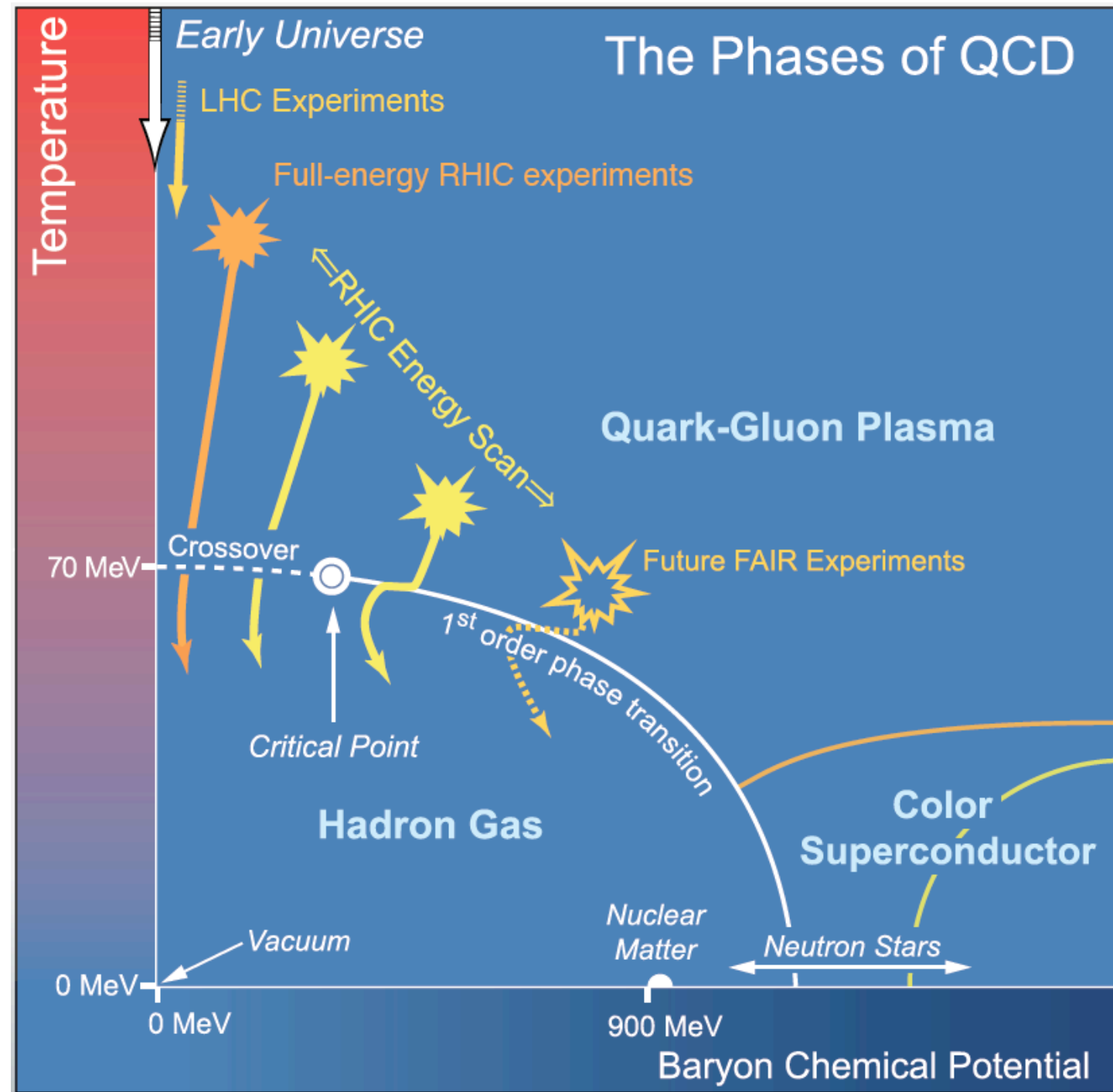
Critical opalescence

Does QCD show Critical opalescence ?

# But this happens at the critical point



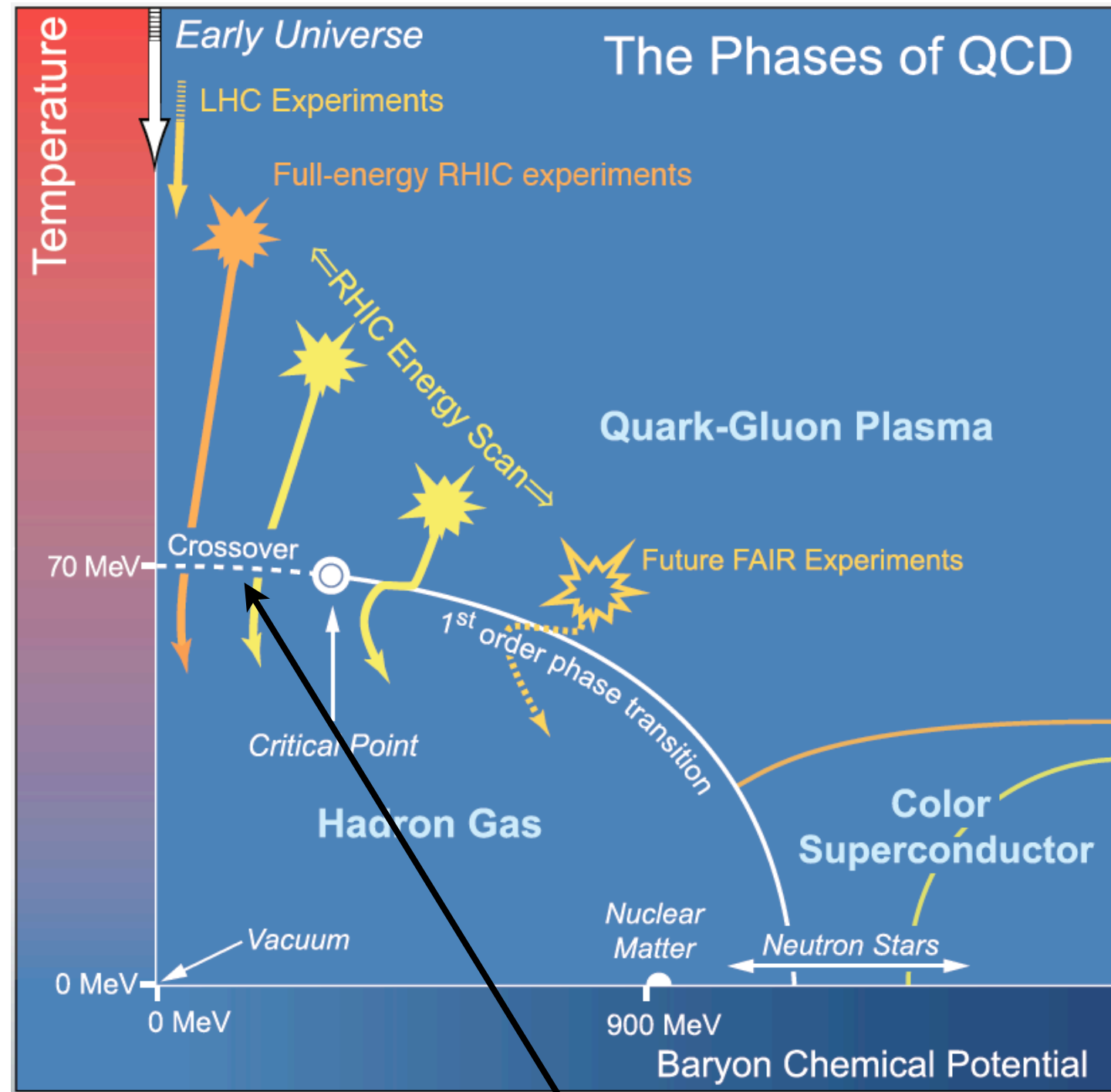
# But this happens at the critical point



There may be remnant extremal behavior at smaller  $\mu_B$

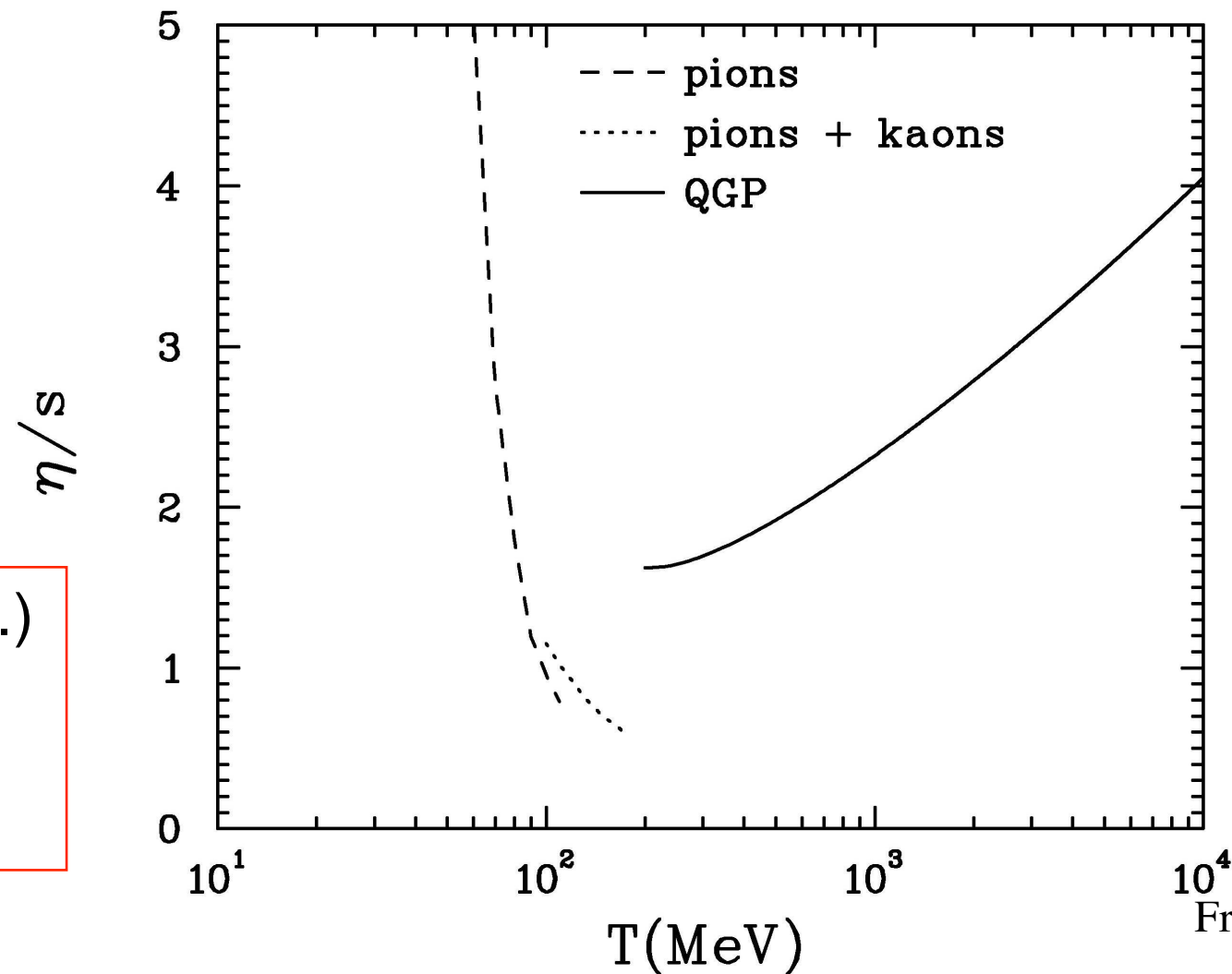


# But this happens at the critical point



There may be remnant extremal behavior at smaller  $\mu_B$

# Remnants of extremal behavior at $\mu \rightarrow 0$



Low T (*Prakash et al.*)  
using experimental  
data for 2-body  
interactions.

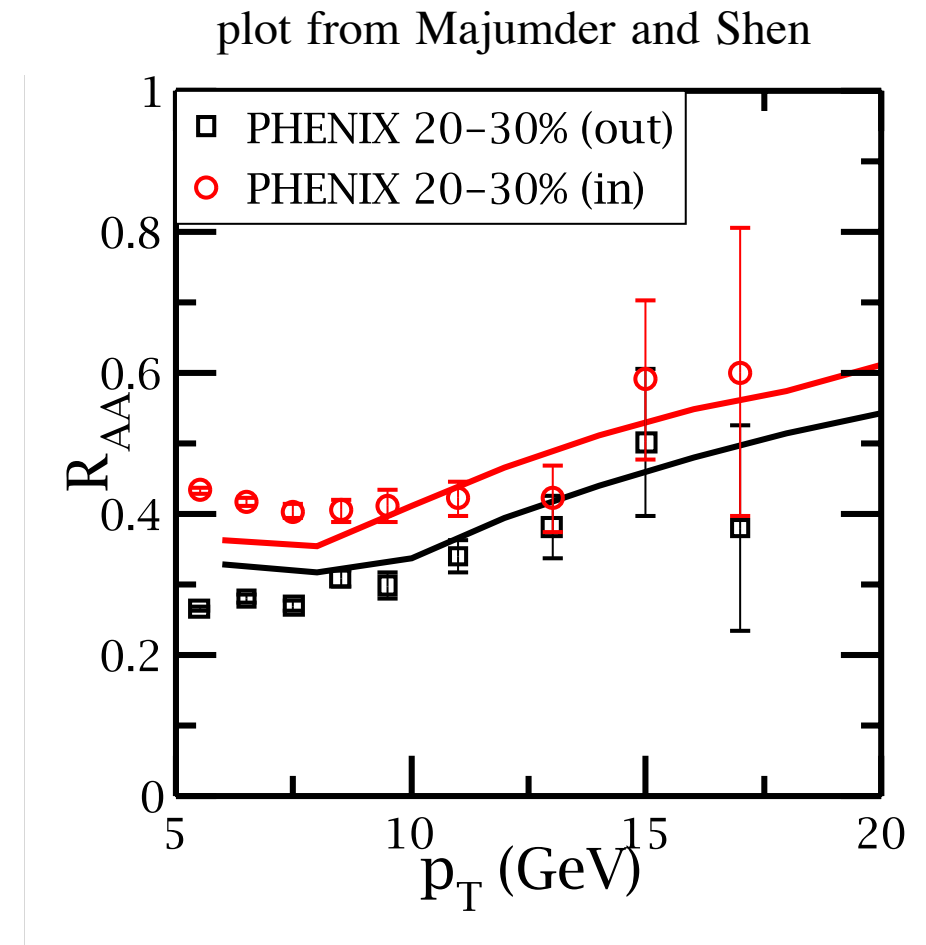
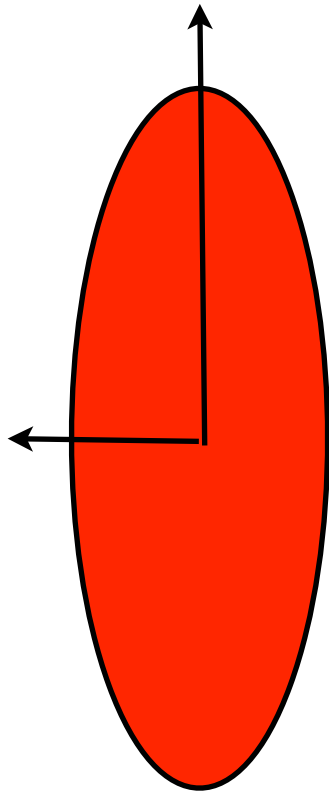
High T (*Yaffe et al.*)  
using perturbative  
QCD.

From talk by J. Kapusta at QM2006

QCD matter at RHIC and LHC is far from the critical point  
Yet one still expects a minimum, at least theoretically!  
Not inconsistent with any bulk measurement at RHIC/LHC

# Any reason to expect a bump in $\hat{q}/T^3$ ?

Look at the  $R_{AA}$  vs reaction plane



At low  $p_T$  hard to describe the reaction plane dependence

Note: at low  $p_T$  many higher power corrections become important  
these have so far been ignored



# A non-monotonic behavior in $\hat{q}/T^3$

PRL **102**, 202302 (2009)

PHYSICAL REVIEW LETTERS

week ending  
22 MAY 2009

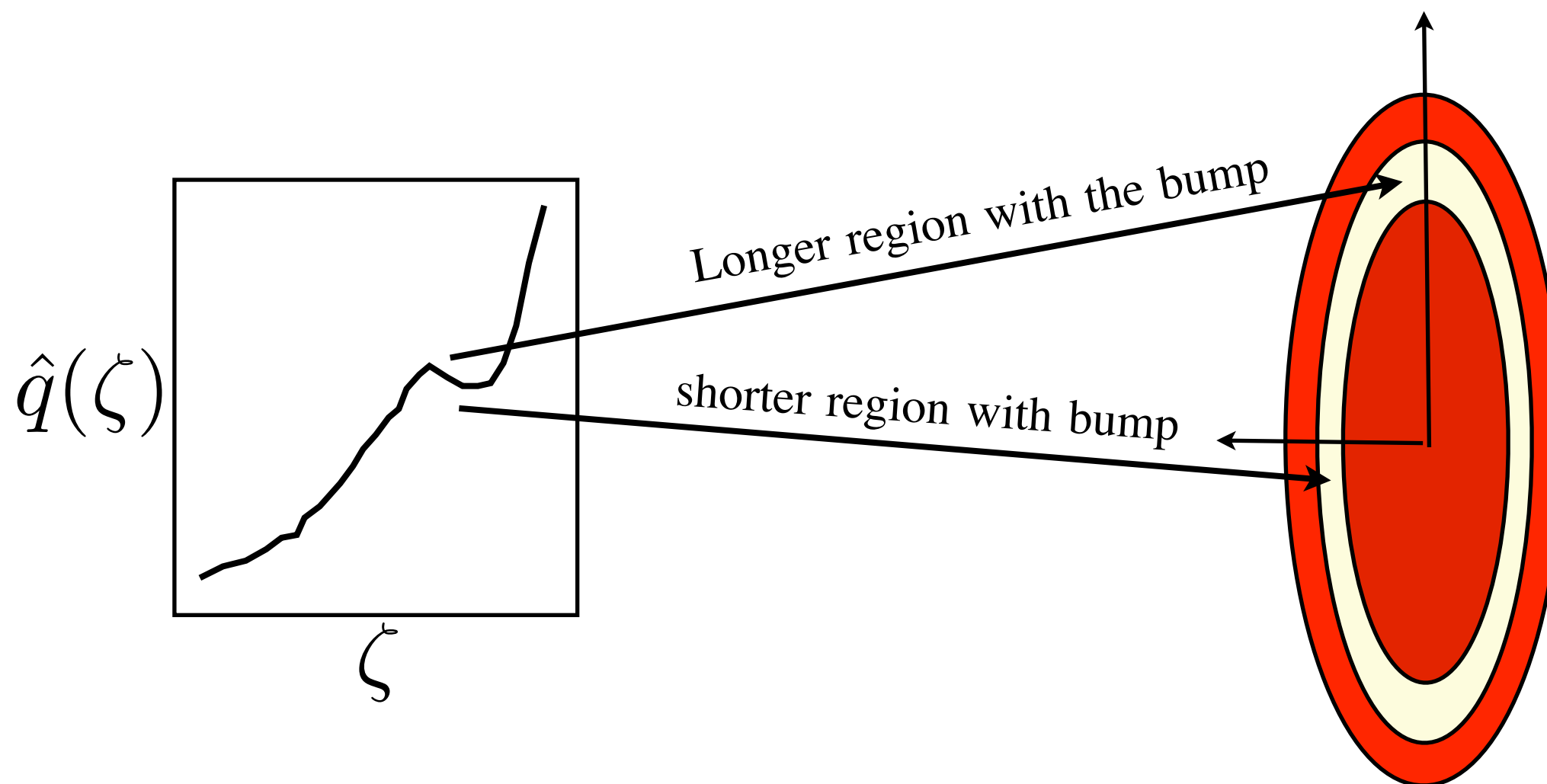
## Angular Dependence of Jet Quenching Indicates Its Strong Enhancement near the QCD Phase Transition

Jinfeng Liao<sup>1,2,\*</sup> and Edward Shuryak<sup>1,†</sup>

<sup>1</sup>*Department of Physics and Astronomy, State University of New York, Stony Brook, New York 11794, USA*

<sup>2</sup>*Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

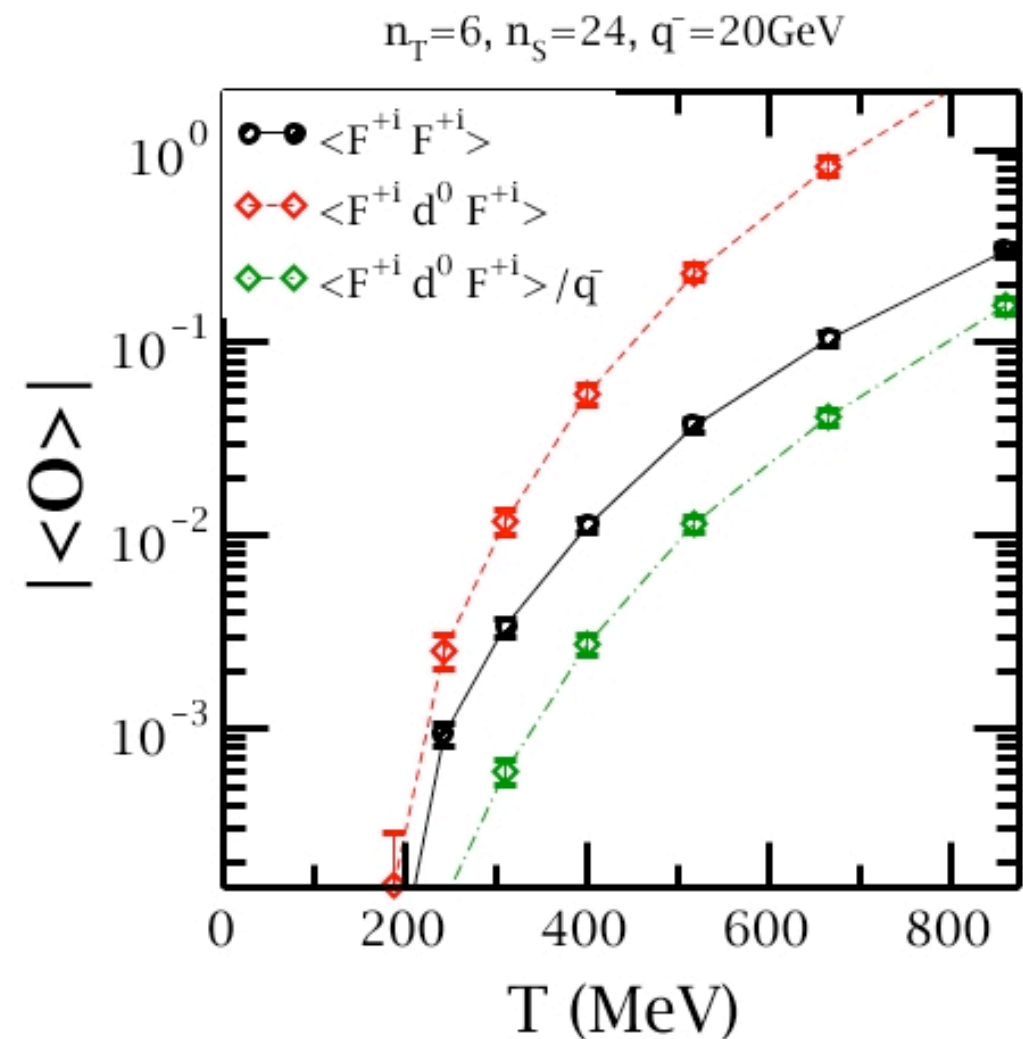
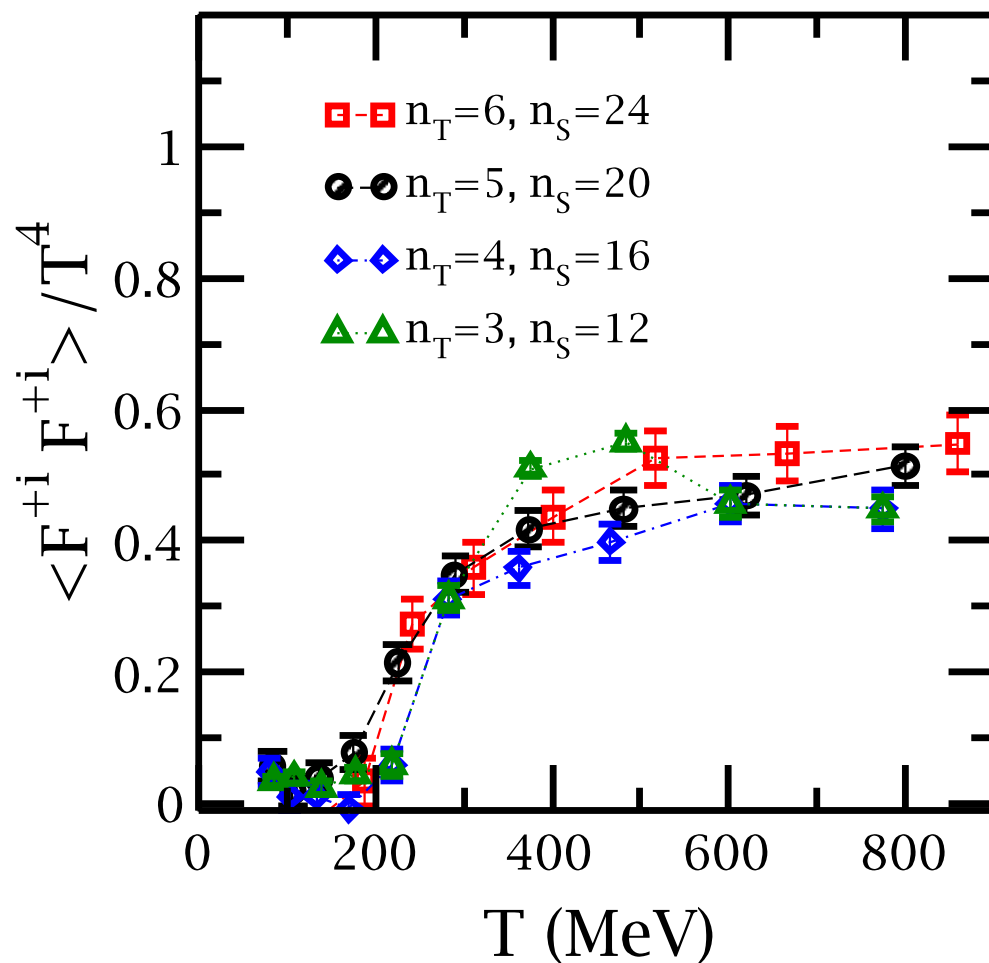
(Received 22 October 2008; revised manuscript received 19 February 2009; published 22 May 2009)



# A Lattice calculation of $\hat{q}$

Long story short: can analytically continue  $\hat{q}$  to euclidean space and evaluate as a series

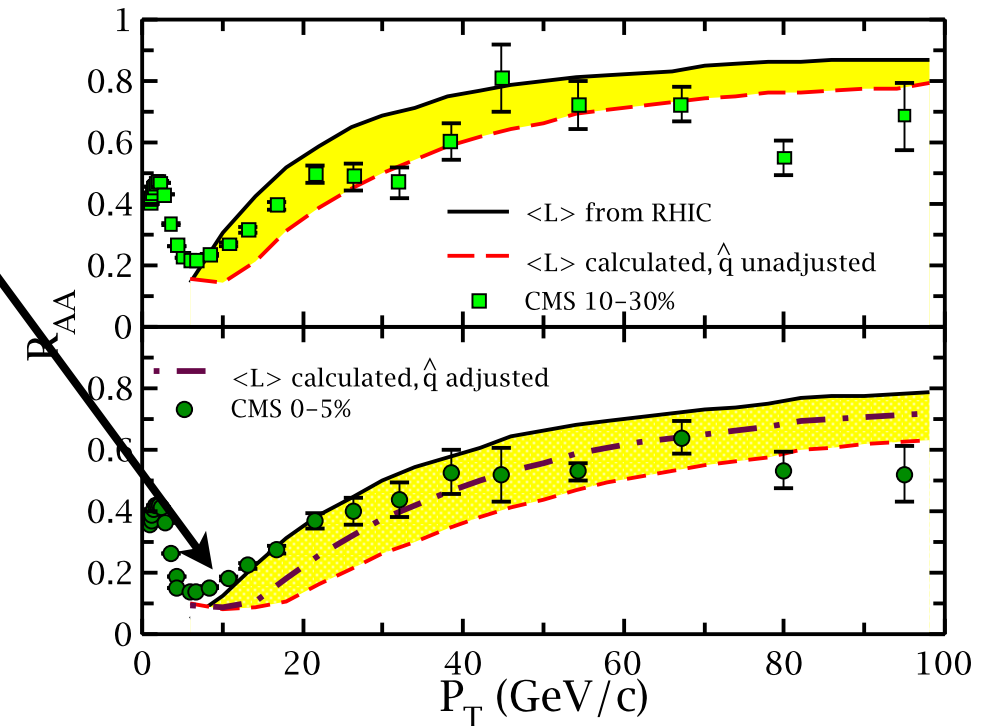
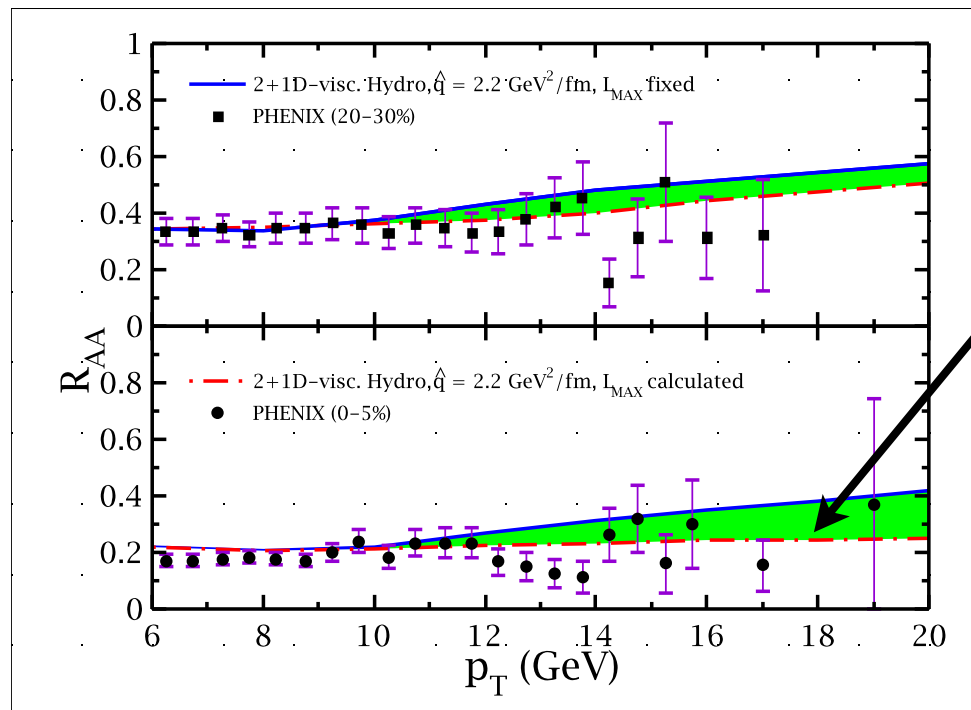
$$\langle M | F_{\perp}^{+\mu} \sum_{n=0}^{\infty} \left( \frac{-q \cdot i\mathcal{D} - \mathcal{D}_{\perp}^2}{2q^- Q_0} \right)^n F_{\perp, \mu}^{+} | M \rangle$$



Note: quenched  $SU(2)$ , results not inconsistent with a bump above  $T_C$ .

# What exactly are we looking for?

Note: The QGP at RHIC and LHC already very opaque to most jets



Jets tend to disintegrate as they propagate through the QGP

We need very specific range of parameters to see a  
maximal scaled opacity ( $\hat{q}/T^3$ ).



# A couple of things to keep in mind

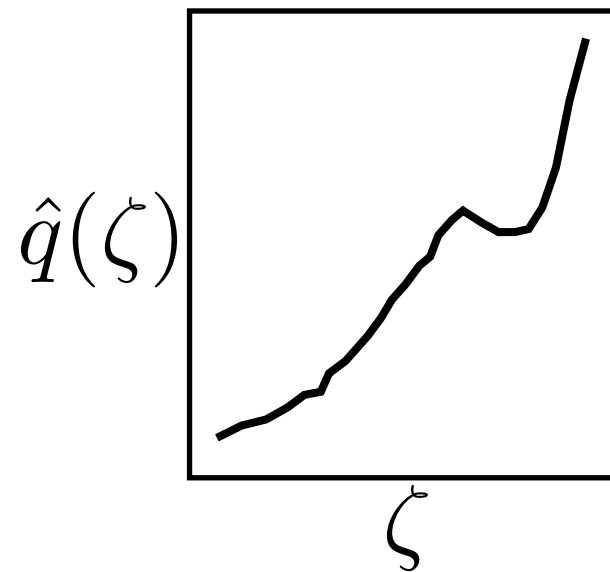
(I) Jet quenching is sensitive to

$$\int_{\text{Origin}}^{\text{Exit}} d\zeta q(\zeta)$$

And not to the scaled quantity

$$\int_{\text{Origin}}^{\text{Exit}} d\zeta \frac{q(\zeta)}{T^3(\zeta)}$$

Thus, we are looking for a wiggle in a steeply falling curve



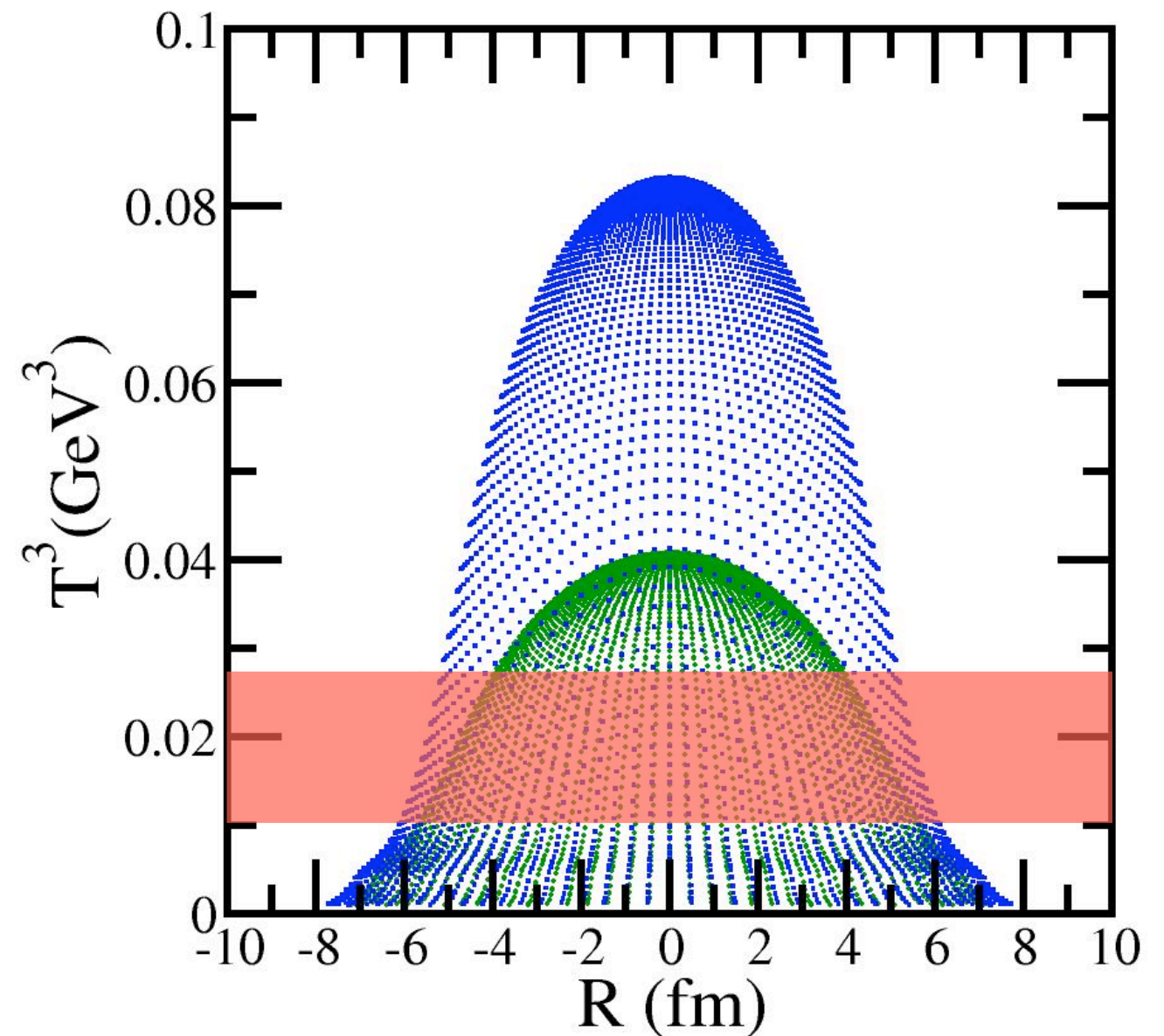
(II) At very high  $p_T$  a small  $R_{AA}$  does not mean the modification to the jet is as large as at lower  $p_T$

If there is a wiggle at  $T_C$ , RHIC is a better place to look

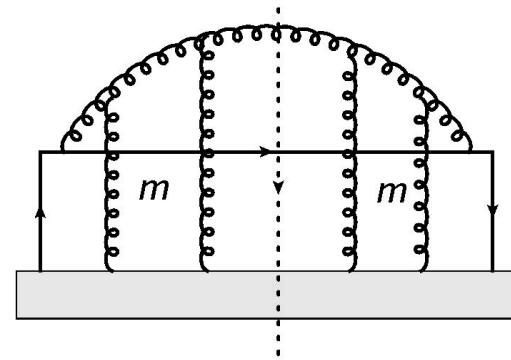
$T^3$  profiles from  
the OSU 2+1 D hydro, 0-5% evts

At LHC, region of non-monotonic  
behavior suppressed by much  
larger  $\hat{q}$  values at earlier times

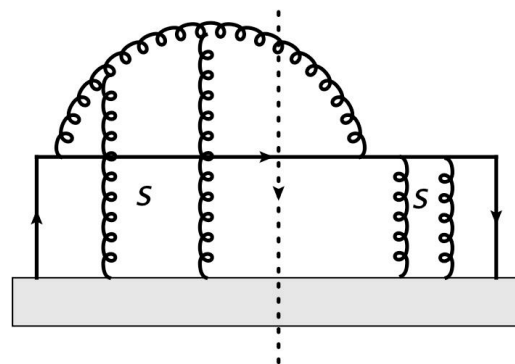
Region stretches for about  
2fm/c at RHIC and 1fm/c at LHC



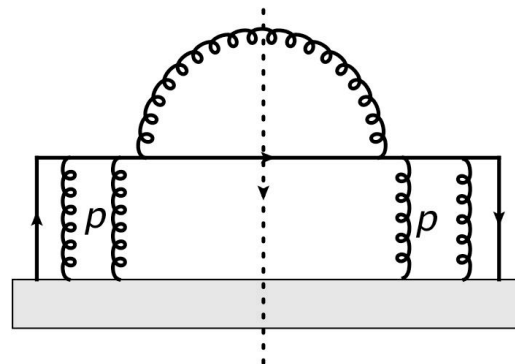
# What kind of jets are best suited for probing this region



$$\sim C_A^m \int dy \frac{l_{\perp}^2}{l_{\perp}^2} P(y) \int d\zeta \frac{2\hat{q}}{l_{\perp}^2} \left[ 2 - 2 \cos \left( \frac{l_{\perp}^2}{2E} \zeta \right) \right]$$



$$\sim - \left( \frac{C_A}{2} \right)^s \int dy \frac{l_{\perp}^2}{l_{\perp}^2} y P(y) \int d\zeta \frac{\hat{q}}{l_{\perp}^2} \left[ 2 - 2 \cos \left( \frac{l_{\perp}^2}{2E} \zeta \right) \right]$$



$$\sim (C_F)^p \int dy \frac{l_{\perp}^2}{l_{\perp}^2} y^2 P(y) \int d\zeta \frac{\hat{q}Q}{l_{\perp}^2} \left[ 2 - 2 \cos \left( \frac{l_{\perp}^2}{2E} \zeta \right) \right]$$

Results from  
a multiple  
scattering  
single emission  
calculation

The modification is controlled by the parameter

$$\frac{\hat{q}L}{Q^2}$$

If this is too small then jets not modified

If its too big then jets are completely quenched, ideal value  $\sim 0.1$



# Some estimates!

In the region of the bump,  $q_q \sim 0.5 \text{ GeV}^2/\text{fm}$

$$q_g \sim 1 \text{ GeV}^2/\text{fm}$$

Length is about 2fm at RHIC, thus  $qL \sim 1-2 \text{ GeV}^2$

Thus we need a  $Q^2 \sim 10 - 20 \text{ GeV}^2$

If we want the jet to emit once in this region then

$$\begin{aligned} \text{Formation time} &\sim \frac{E}{Q^2} \sim 1 \text{ fm} = 5 \text{ GeV}^{-1} \\ &\Rightarrow E \sim 50 - 100 \text{ GeV} \end{aligned}$$

At LHC, length of region is like 1fm, then  $qL \sim 0.5-1 \text{ GeV}^2$

thus for a  $Q^2 \sim 5 - 10 \text{ GeV}^2$ , need  $E \sim 25 - 50 \text{ GeV}$

# Virtuality driven MCs

These are very hand wavy estimates

However, we now have the technology to  
test these with virtuality driven MC on a medium  
with a bump in  $\hat{q}$

Results will appear soon!

Back up

