In-medium QCD forces at high temperature

Yukinao Akamatsu (Nagoya/KMI)

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1. "INTRODUCTION"

Key members

• J/Ψ suppression







• J/ Ψ survival (up to 1.6Tc?)

Debye









Shannon 4

13/04/03

HQs and quarkonia in thermal QCD

Bayes?

New paradigm since '07

• Imaginary part of potential



et al

followed by other groups with path-integral approach lattice QCD+MEM approach pNRQCD EFT approach

$$V_{Q\bar{Q}}(R) = -\frac{g^2 C_F}{4\pi} \left(\omega_D + \frac{e^{-\omega_D R}}{R} + iT\varphi(\omega_D R) \right)$$

New members

• New does not mean young...



Feynman



Caldeira

Leggett

Quartet Magico-Fisica (di sistemi quantistici aperti)?

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2. STOCHASTIC POTENTIAL

Complex potential



Laine et al (07), Beraudo et al (08), Brambilla et al (10), Rothkopf et al (12)

• Stochastic interpretation

$$\begin{split} \Psi(t + \Delta t, R) &= \exp\left[-i\Delta t \left\{ V(R) + \Theta(t, R) \right\} \right] \Psi(t, R), \\ \left\langle \Theta(t, R) \right\rangle &= 0, \quad \left\langle \Theta(t, R) \Theta(t', R') \right\rangle = \Gamma(R, R') \delta_{tt'} / \Delta t, \\ &\Rightarrow i \frac{\partial}{\partial t} \left\langle \Psi(t, R) \right\rangle_T = \left\{ V(R) - \frac{i}{2} \Gamma(R, R) \right\} \left\langle \Psi(t, R) \right\rangle_T. \end{split}$$

Unitary evolution

Imaginary part from averaging the random phase

Akamatsu & Rothkopf (12)

Numerical simulation

Stochastic evolution



3. INFLUENCE FUNCTIONAL

Classical picture

• $M = \infty$ or $M < \infty$ matters



Key words: Forces, open quantum system, influence functional

Open quantum system

• Basics

{ sys = heavy quarks
 env = gluon, light quarks

Hilbert space $H_{tot} = H_{sys} \otimes H_{env}$

 $i\frac{d}{dt}\hat{\rho}_{tot}(t) = \left[\hat{H}_{tot}, \hat{\rho}_{tot}(t)\right]$

von Neumann equation

Trace out the environment

Reduced density matrix

Master equation

$$\hat{\rho}_{red}(t) = Tr_{env} [\hat{\rho}_{tot}(t)]$$
$$i \frac{d}{dt} \hat{\rho}_{red}(t) = ?$$

(Markovian limit)

$$\begin{array}{c} \begin{array}{c} \mathcal{O}[\varphi_{1}^{\text{ini}},\varphi_{2}^{\text{ini}}] & 1 & \varphi_{1},\eta_{1} \\ & \rho[\varphi_{1}^{\text{ini}},\varphi_{2}^{\text{ini}}] & 2 & \varphi_{2},\eta_{2} \end{array}$$

• Partition function
$$\begin{array}{c} \mathcal{O}[\psi_{1,2}q_{1,2}A_{1,2}]\rho[\psi_{1}^{*}q_{1}^{*}A_{1}|^{\text{ini}},\psi_{2}q_{2}A_{2}|^{\text{ini}}] \\ & \times \exp[iS[\psi_{1}]-iS[\psi_{2}]+i\int\psi_{1}\eta_{1}-i\int\psi_{2}\eta_{2}] \\ & \times \exp[iS[q_{1}A_{1}]-iS[q_{2}A_{2}]+ig\int j_{1}A_{1}-ig\int j_{2}A_{2}] \end{array}$$

$$\begin{array}{c} \rho_{\text{tot}} = \rho_{\text{env}}^{\text{eq}} \otimes \rho_{\text{sys}} \text{ Factorized initial density matrix} \\ & \to \rho_{\text{tot}}[\psi_{1}^{*}q_{1}^{*}A_{1}|^{\text{ini}},\psi_{2}q_{2}A_{2}|^{\text{ini}}] = \rho_{\text{env}}^{\text{eq}}[q_{1}^{*}A_{1}|^{\text{ini}},q_{2}A_{2}|^{\text{ini}}] \cdot \rho_{\text{sys}}[\psi_{1}^{*\text{ini}},\psi_{2}^{\text{ini}}] \end{array}$$
nfluence functional
$$\begin{array}{c} = Z_{qA}[j_{1},j_{2}] = \exp[iS_{\text{FV}}[j_{1},j_{2}]] \\ = \int \rho_{A}(f_{1},f_{2}) = \exp[iS_{\text{FV}}[j_{1},j_{2}]] \end{array}$$

$$= \exp\left[-g^{2}/2\int j_{1}G_{A}^{F}j_{1} + j_{2}G_{A}^{\tilde{F}}j_{2} - j_{1}G_{A}^{>}j_{2} - j_{2}G_{A}^{<}j_{1} + \int g^{3}G_{A}^{(3)}jjj + g^{4}G_{A}^{(4)}jjjj + \cdots\right]$$

Influence functional

Stochastic potential • LO pQCD, NR limit, slow dynamics (finite in $M \rightarrow \infty$) $S_{1+2} = S_{kin}^{NR}[Q_{1(c)}] - S_{kin}^{NR}[Q_{2(c)}] + S_{FV}^{LONR}[j_1, j_2] + \cdots$ $S_{\rm FV}^{\rm LONR}[j_1, j_2] = \begin{bmatrix} -\frac{1}{2} \int_{t, \vec{x}, \vec{y}} (\rho_{1a}, \rho_{2a})_{(t, \vec{x})} \begin{bmatrix} V(\vec{x} - \vec{y}) & -iD(\vec{x} - \vec{y}) \\ -iD(\vec{x} - \vec{y}) & -V^*(\vec{x} - \vec{y}) \end{bmatrix} \begin{pmatrix} \rho_{1a} \\ \rho_{2a} \end{pmatrix}_{(t, \vec{y})}$ $\left. -\int_{t,\vec{x},\vec{y}} \left\{ \frac{\vec{\nabla} D(\vec{x}-\vec{y})}{4T} \cdot \left(\vec{j}_{1a}(t,\vec{x})\rho_{2a}(t,\vec{y}) + \rho_{1a}(t,\vec{x})\vec{j}_{2a}(t,\vec{y}) \right) \right\}$ Drag force $-g^{2}\left\{\overline{G}_{00\ ab}^{R}(\vec{x}-\vec{y})+i\overline{G}_{00\ ab}^{>}(\vec{x}-\vec{y})\right\} \equiv V(\vec{x}-\vec{y})\delta_{ab}$ (vanishes in $M \rightarrow \infty$) $-g^{2}\overline{G}_{00\ ab}^{>}(\vec{x}-\vec{y}) \equiv D(\vec{x}-\vec{y})\delta_{ab} = \mathrm{Im}V(\vec{x}-\vec{y})\delta_{ab}$

4. REAL-TIME DYNAMICS

Functional differential equation

• Path integral \rightarrow Schroedinger equation

 $\rho_{\rm red}[t, Q_{1(c)}^*, \tilde{Q}_{2(c)}^*] \sim \int_{-\infty}^{t, Q_{1(c)}^*, \tilde{Q}_{2(c)}^*} D[Q_{1(c)}^{(*)}, \tilde{Q}_{2(c)}^{(*)}] \rho_{\rm sys}[Q_{1(c)}^{*\rm ini}, \tilde{Q}_{2(c)}^{*\rm ini}] \\ \times \exp\left[iS_{\rm NR}[Q_{1(c)}^{(*)}] - iS_{\rm NR}[\tilde{Q}_{2(c)}^{(*)}] + iS^{\rm LOFV}[j_1, j_2] + \cdots\right]$

 $i\frac{\partial}{\partial t}\rho_{\rm red}[t,Q_{1(c)}^*,\tilde{Q}_{2(c)}^*] = H_{1+2}^{\rm func}[Q_{1(c)}^*,\tilde{Q}_{2(c)}^*]\rho_{\rm red}[t,Q_{1(c)}^*,\tilde{Q}_{2(c)}^*]$

Density matrix for a few HQs

• Remember coherent states

$$\begin{split} \rho_{\text{red}} \left[t, Q_{1(c)}^*, \widetilde{Q}_{2(c)}^* \right] &= \left\langle Q_{1(c)}^* \left| \hat{\rho}_{\text{red}}(t) \right| \widetilde{Q}_{2(c)}^* \right\rangle \\ \left\langle Q_{1(c)}^* \right| &= \left\langle \Omega \left| \exp\left[-\int_{\vec{x}} \left\{ \hat{Q} Q_1^* + \hat{Q}_c Q_{1(c)}^* \right\} \right] \right\rangle & \longrightarrow & \frac{\delta}{\delta Q_1^*(\vec{x})} \left\langle Q_{1(c)}^* \right| \right|_{Q_{1(c)}^* = 0} = \left\langle \Omega \left| \hat{Q}(\vec{x}) \right\rangle \\ \left| \widetilde{Q}_{2(c)}^* \right\rangle &= \exp\left[-\int_{\vec{x}} \left\{ \widetilde{Q}_2^* \widehat{Q}^\dagger + \widetilde{Q}_{2c}^* \widehat{Q}_c^\dagger \right\} \right] \left| \Omega \right\rangle & \frac{\delta}{\delta \widetilde{Q}_2^*(\vec{x})} \left| \widetilde{Q}_{2(c)}^* \right\rangle \right|_{\widetilde{Q}_{2(c)}^* = 0} = -\hat{Q}^\dagger(\vec{x}) \left| \Omega \right\rangle \end{split}$$

$$- \text{ Single HQ} \qquad \rho_{Q}(t, \vec{x}, \vec{y}) \propto \langle \Omega | \hat{Q}(\vec{x}) \hat{\rho}_{\text{red}}(t) \hat{Q}^{\dagger}(\vec{y}) | \Omega \rangle$$
$$= -\frac{\delta}{\delta Q_{1}^{*}(\vec{x})} \frac{\delta}{\delta \widetilde{Q}_{2}^{*}(\vec{y})} \rho_{\text{red}} \left[t, Q_{1(c)}^{*}, \widetilde{Q}_{2(c)}^{*} \right]_{Q_{1(c)}^{*} = \widetilde{Q}_{2(c)}^{*} = 0}$$

$$\rho_{QQ_c}(t, \vec{x}_1, \vec{x}_2, \vec{y}_1, \vec{y}_2), \cdots$$

 $\langle Q^*_{l(c)}$

Physical process

• Scatterings in t-channel



Master equation

 Single HQ $i\partial_t \rho_Q(t, \vec{x}, \vec{y}) = \left\{ \left(a - a^*\right)M + \left(-\frac{\nabla_x^2 - \nabla_y^2}{2M}\right) \right\} \rho_Q(t, \vec{x}, \vec{y})$ (color traced) $+C_F \left\{-iD(\vec{x}-\vec{y}) + \frac{\vec{\nabla}_x D(\vec{x}-\vec{y})}{4T} \cdot \frac{\nabla_x - \nabla_y}{iM}\right\} \rho_Q(t,\vec{x},\vec{y})$ $a = 1 + \frac{C_{\rm F}}{2M} \lim_{r \to 0} V^{(T>0)}(r), \quad V^{(T>0)}(r) = V(r) - V^{(T=0)}(r)$ $\begin{cases} \frac{d}{dt} \langle \vec{x} \rangle = \frac{\langle \vec{p} \rangle}{M}, & \frac{d}{dt} \langle \vec{p} \rangle = -\frac{\gamma}{2MT} \langle \vec{p} \rangle, \\ \frac{d}{dt} \langle E \rangle = -\frac{\gamma}{MT} \left(\langle E \rangle - \frac{3T}{2} \right). & \text{Langev} \\ \gamma = \frac{C_F}{3} \nabla^2 D(x) \Big|_{x=0} = \frac{g^2 C_F}{9} \int \frac{d^3 k}{(2\pi)^3} k^2 \tilde{G}_{00,aa}^{>}(\omega = 0, k) \end{cases}$ Langevin dynamics Moore & Teaney (05)

HQs and quarkonia in thermal QCD

Complex potential

• QQbar

$$\frac{\partial_t \left\langle \Psi_{ij}(t;\vec{x},\vec{y}) \right\rangle_T}{\left\langle \Psi_{ij}(t;\vec{x},\vec{y}) \right\rangle_T} = \frac{\partial_t \left\langle J_{ij}(t;\vec{x},\vec{y})J_{kl}^{\dagger}(0;\vec{x}_0,\vec{y}_0) \right\rangle_T}{\left\langle J_{ij}(t;\vec{x},\vec{y})J_{kl}^{\dagger}(0;\vec{x}_0,\vec{y}_0) \right\rangle_T} = V_{ij}(\vec{x}-\vec{y})$$

Projection onto singlet state

Debye screening w/ imaginary part (complex potential)

$$V_{\text{singlet}}(R) = 2(a-1)M - C_F V(R) = -\frac{g^2 C_F}{4\pi} \left(\omega_D + \frac{e^{-\omega_D R}}{R} + iT\varphi(\omega_D R)\right)$$

Laine et al (07), Beraudo et al (08), Brambilla et al (10)

Stochastic process

• *M*=∞

$$\begin{split} &\exp\left[iS_{\rm FV}^{\rm LONR}[j_1,j_2]\right] \\ &= \exp\left[-i/2\int_{t,\vec{x},\vec{y}} \operatorname{Re}V(\vec{x}-\vec{y})\left\{\rho_{1a}(t,\vec{x})\rho_{1a}(t,\vec{y})-\rho_{2a}(t,\vec{x})\rho_{2a}(t,\vec{y})\right\}\right] \\ &\times \left\langle \exp\left[-i\int_{t,\vec{x},\vec{y}}\xi_a(t,\vec{x})\left\{\rho_{1a}(t,\vec{x})-\rho_{2a}(t,\vec{x})\right\}\right]\right\rangle_{\xi} \\ &\left\langle \xi_a(t,\vec{x})\xi_b(s,\vec{y})\right\rangle = -\delta_{ab}\delta(t-s)D(\vec{x}-\vec{y}) \end{split}$$

Debye screening + fluctuation (stochastic potential)

Akamatsu & Rothkopf (12)

5. SUMMARY

So far and beyond

- LO perturbation
 - Heavy quarks in QGP as open quantum system
 - Unified description of forces
 - Scatterings in t-channel
- NLO perturbation
 Gluo-dissociation
- Phenomenology at RHIC/LHC