

# On Study of Quarkonium Melting using Lattice Correlators

Saumen Datta

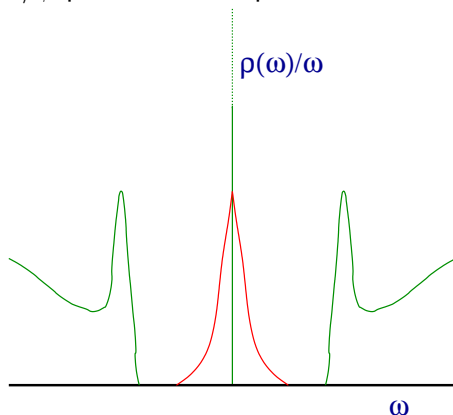
(Based on works with F. Karsch, P. Petreczky, I. Wetzorke, and  
with D. Banerjee)

Tata Institute of Fundamental Research, Mumbai

April 4, 2013

# Introduction

Problem at hand: find spectral function, e.g.,  $\rho_{\bar{c}\gamma c}(\omega)$  and the melting of the  $J/\psi$  peak in the dilepton channel.



From the lattice: look at Matsubara correlators

$$\Rightarrow \text{invert } G_{\bar{c}\gamma c}(\tau) = \int d\omega \rho(\omega) \frac{\cosh \omega(\tau - \frac{1}{2T})}{\sinh \frac{\omega}{2T}}.$$

# Charmonia using MEM

Most popular: maximum entropy method: use prior information to keep in check uncontrolled directions in search space.

default model dependence; sensitivity on other parameters.

To get stable  $\rho(\omega)$  from  $\mathcal{O}(20)$  data points: need to give ultraviolet information to get infrared.

First studies: 1S charmonia survive till quite deep in plasma, while the 1P states dissolve early.

Datta, Karsch, Petreczky, Wetzorke, 2004; Asakawa & Hatsuda, 2004

But correlators are completely consistent with other scenarios.

Umeda, 2007; Mocsy and Petreczky, 2007-2009.

Would be very useful to have other analysis techniques.

# Cuniberti et al approach

An interesting approach, using detailed analyticity properties of Green's functions.

Cuniberti, Micheli, Viano, 2001; Burnier, Laine, Mether, 2011

$$\rho(\omega) = \int_0^\infty dt \sin \omega t J^+(t), \quad J^+(t) = e^{-e^{-t}} \sum_{l=0}^{\infty} a_l L_l(2e^{-t})$$

Extraction of electric conductivity from UV-subtracted vector correlator data.

Burnier and Laine 2012

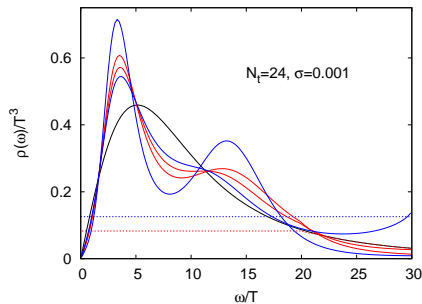
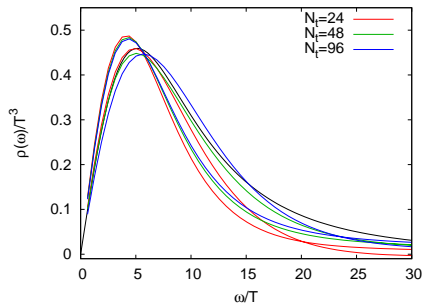
Can this work in future for quarkonia?

How do the two methods perform for per-mille accuracy correlator,  $N_t \sim 24 - 48$ ?

# Toy example: electric field correlator, simplified

Burnier, Laine, Mether, 2011

$$\frac{\rho(\omega)}{16\pi^2 T^3} = \frac{2\pi\omega}{(\omega^2 + 8\pi^2)^2}$$

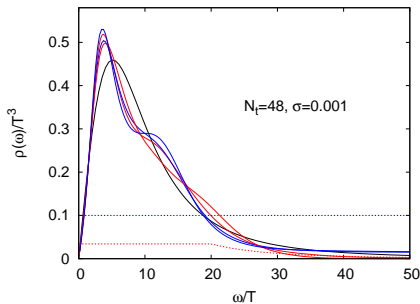
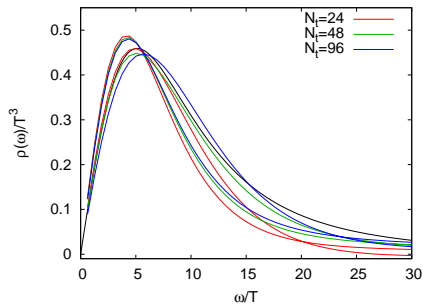


MEM: Brian's algorithm, with modified kernel to avoid problems at  $\omega \rightarrow 0$  (Aarts et al. '07).

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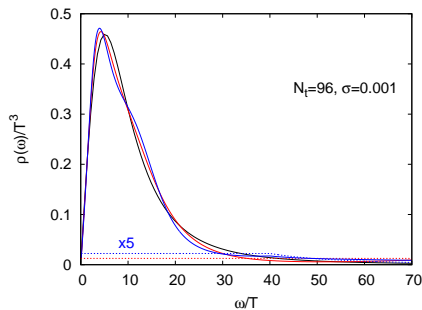
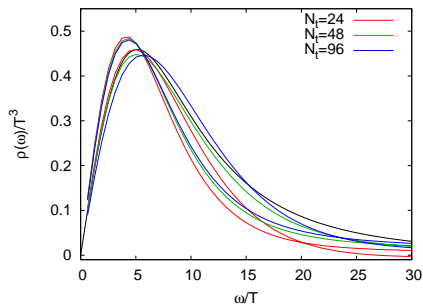


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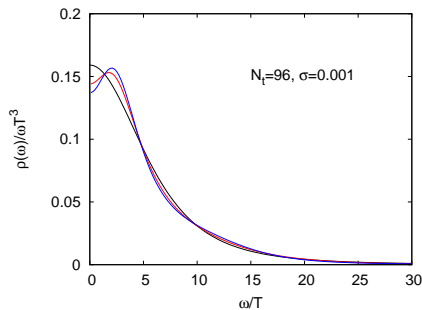
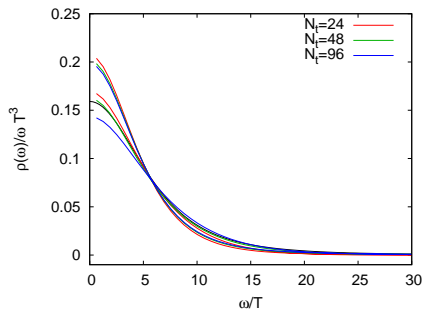


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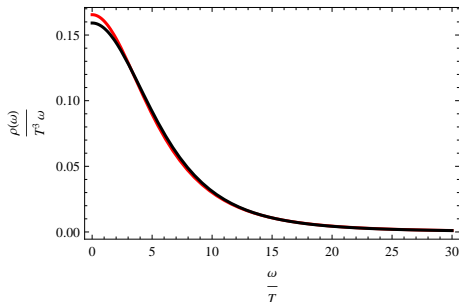


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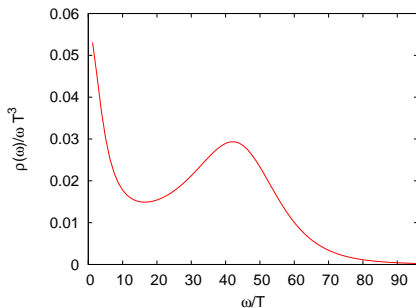
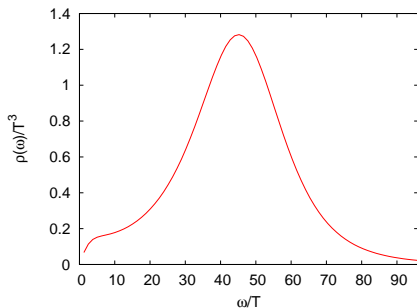
# Toy example: electric field correlator, simplified

For this simple example, Pade method also works well  $\Rightarrow$  fit  $G(\omega_n)$  with a rational fit form, and get  $\rho(\omega) = \frac{1}{\pi} \text{Im}G(\omega \rightarrow -i\omega)$ .



# Toy Example: Gunnarsson et al.

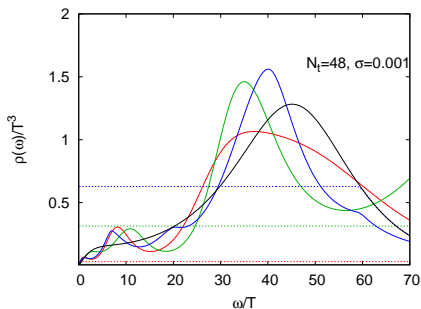
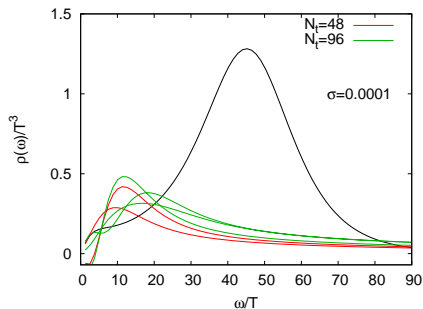
Unfortunately, Cuniberti's method does not work so well in more complicated examples.



Model for current-current correlator in 2D Hubbard model.

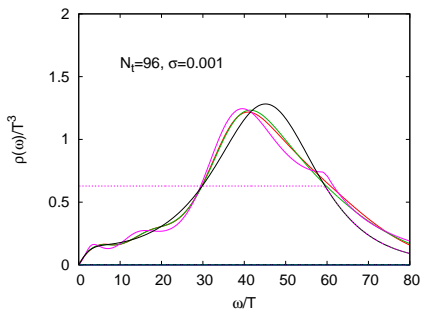
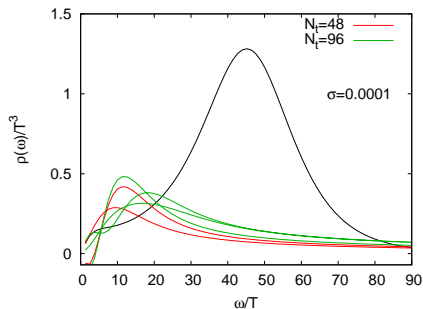
Gunnarsson, Haverkort, Sangiovanni, PR B 82, 165125 ('10)

# Toy Example: Gunnarsson et al, Contd.



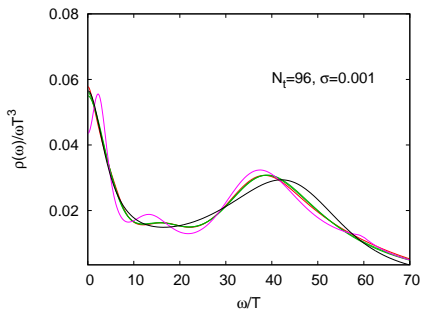
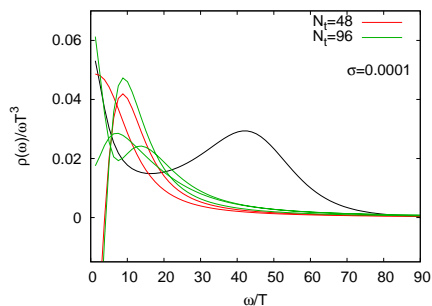
Similar results found for a two-peak spectral function.

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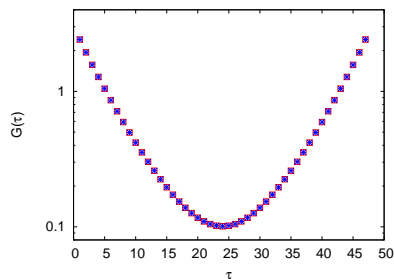
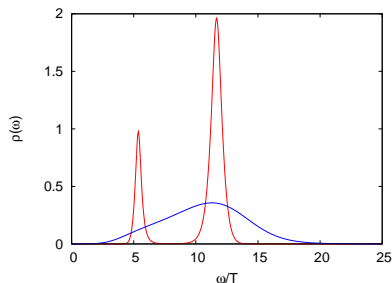


Similar results found for a two-peak spectral function.

# Current status

Cuniberti's method in current form probably suitable when the subtracted correlator has little structure, e.g., a diffusive peak.

Maximum entropy method: qualitative features reproduced, but prone to create structures : for  $N_t = 48$  or less, need to put in sufficient information, specially about the ultraviolet, and use any other reliable theoretical input.

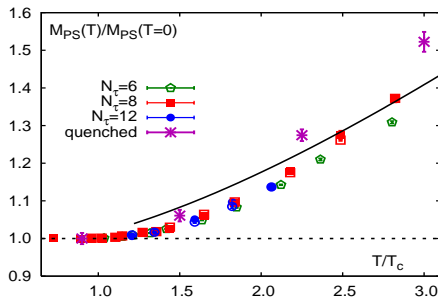


# Spatial “screening” correlator

Can one look at “screening correlators” instead?

$$G(z, T) = \int \frac{dk_z}{\pi} e^{ik_z z} \int d\omega \frac{\rho(\omega, k_{\perp} = 0, k_z)}{\omega}$$

Study of correlator easier, but interpretation in terms of spectral function more difficult.

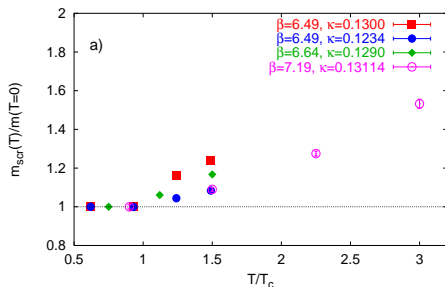
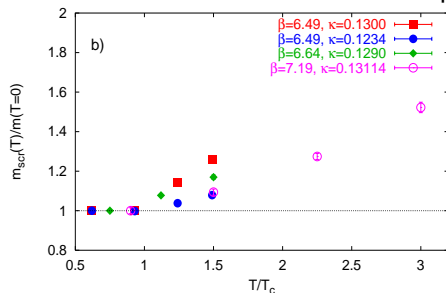


Indicative of dissolution of  $J/\psi$  around  $T \sim 1.5 T_c$ .

Karsch, Laermann, Mukherjee, Petreczky, PRD 85('12) 114501

# Quenched study

Similar behavior observed earlier in quenched QCD, where spatial correlator was seen to show temperature effect already at  $1.1 T_c$ .

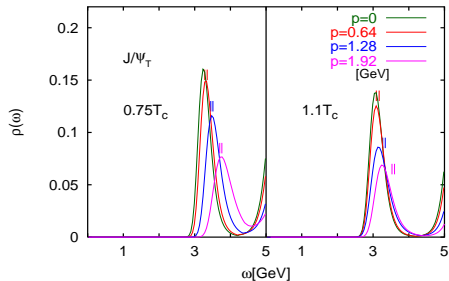
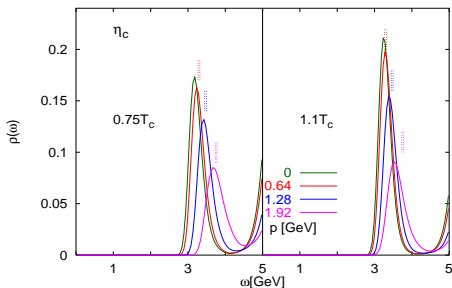


Datta, Karsch, Petreczky, Wetzorke, PRD ('04)

Change in screening mass does not directly imply ground state modification, e.g., one can have a spectral function peak  $\delta(\omega^2 - A^2(T)\vec{p}^2) \Rightarrow m_{\text{scr}}(T) = m/A(T)$



# Momentum modification?

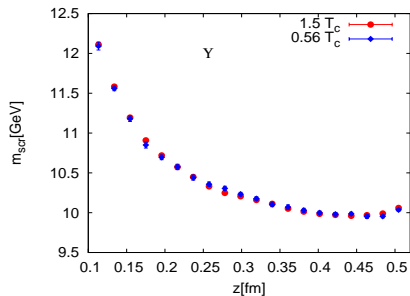
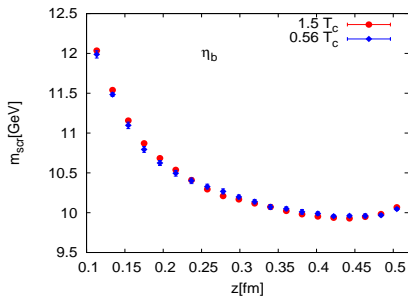


Datta, Karsch, Petreczky, Wetzorke, J.Phys. G 31 ('05) S351.

One can study more detailed properties of the screening correlator to isolate the cause.

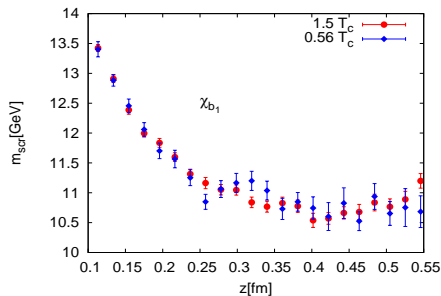
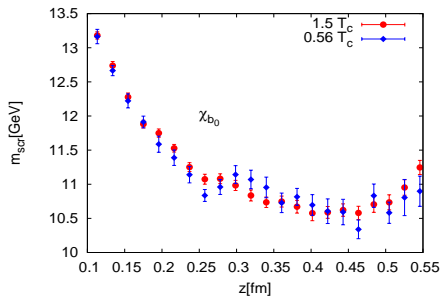
# Screening of $\bar{b}b$ charges

What does the screening correlator tell us about  $\bar{b}b$  mesons?



Expectedly, no modification of 1S screening length at  $1.5 T_c$ .

# Screening of $\bar{b}b$ charges Contd.

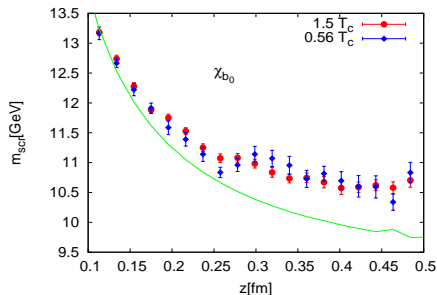
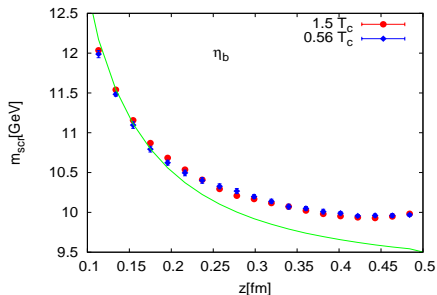


Interestingly, no modification of  $\chi_b$  screening length at  $1.5 T_c$  either.

Counterintuitive, if  $\chi_b$  is dissolved?

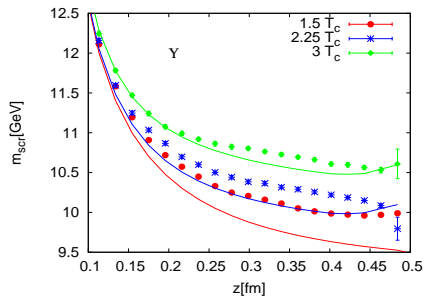
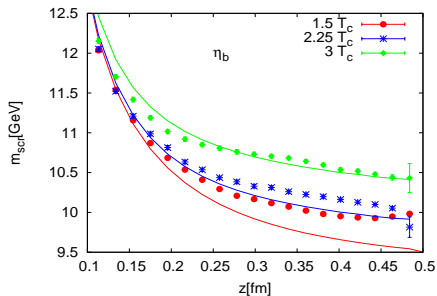
# Screening of $\bar{b}b$ charges

The behavior of  $m(z)$  very different from free theory.



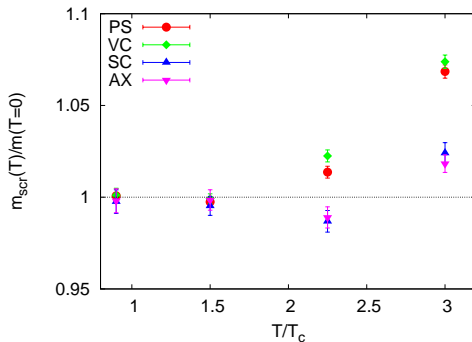
# Screening of $1S \bar{b}b$ above $T_c$

At  $2.25 T_c$ , the  $1S$  states also show a temperature dependence.



( $m_b^{\text{bare}}$  tuned to give agreement at  $3 T_c$ .)

# Combined plot of screening mass of $\bar{b}b$ sources



- ▶ Study of quarkonia melting on lattice is a nontrivial problem.
- ▶ Combination of fine lattices and maximum entropy based analysis is the leading strategy now.

Kaczmarek et al; Aarts, Skullerud, et al.

But important to explore, in particular, other analysis strategies, and utilize the other theoretical results.

- ▶ In particular, it is important to look at both the temporal and spatial correlation functions.

This may be particularly important for bottomonia.