

Yerevan, 18. 9. 2013

## The Modern Physics of Compact Stars and Relativistic Gravity

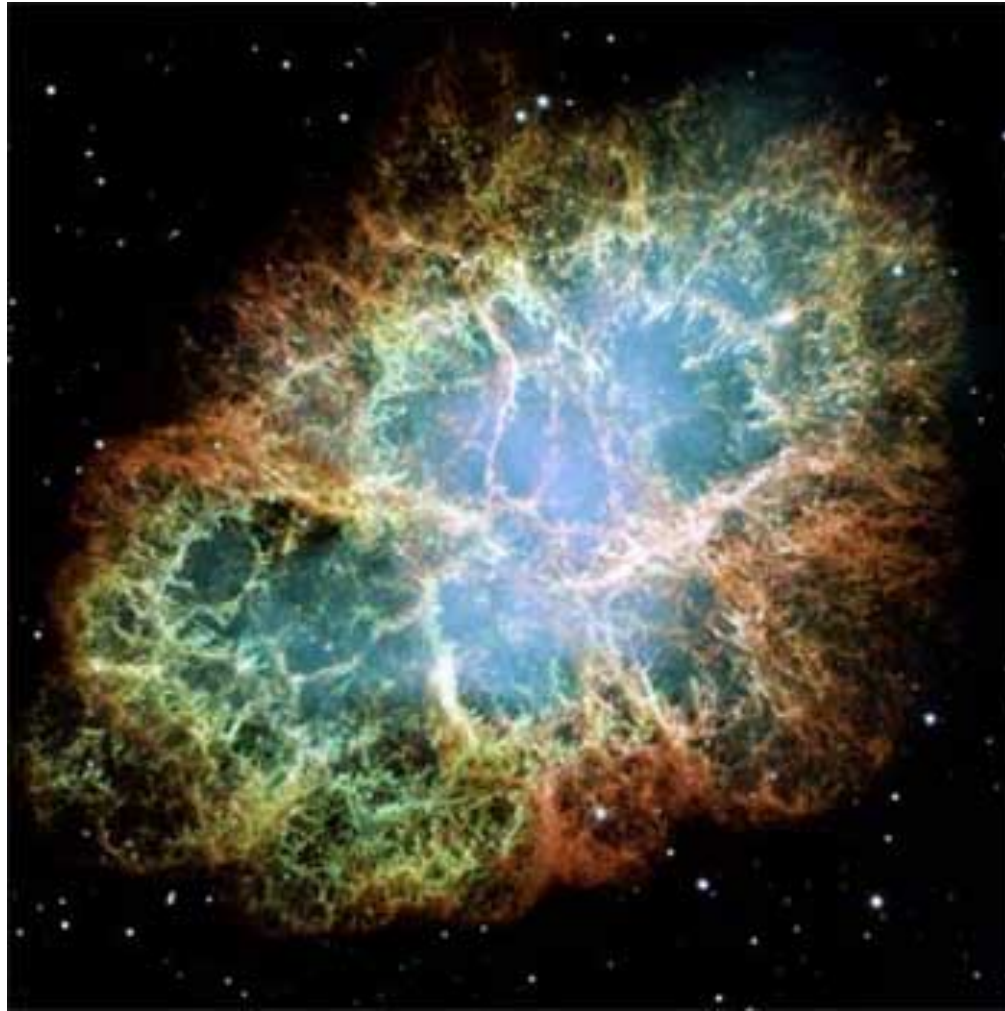
# Correlations in Nuclear Matter and the Symmetry Energy

Gerd Röpke, Rostock



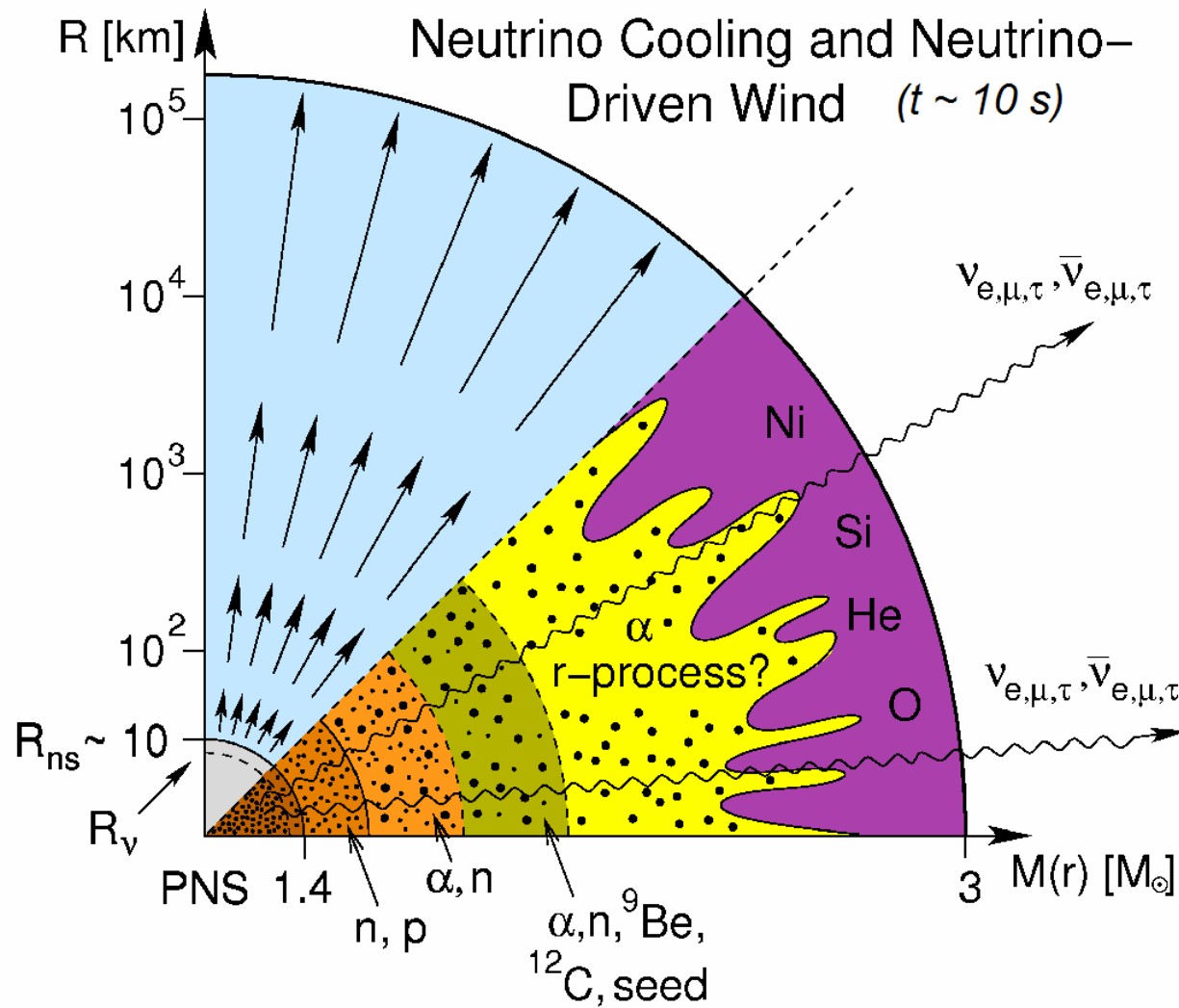
# Supernova

Crab nebula, 1054 China, PSR 0531+21



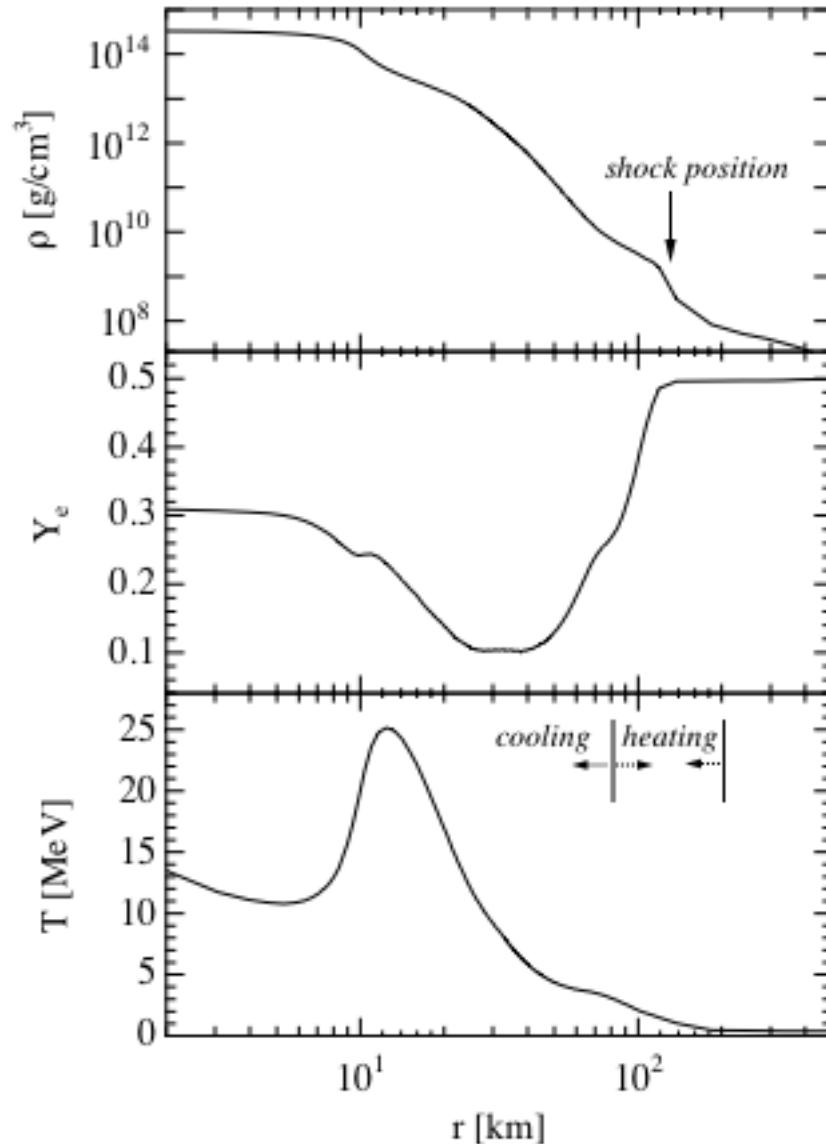
M1, the Crab Nebula. Courtesy of NASA/ESA

# Supernova explosion



T.Janka

# Core-collapse supernovae



Density.

electron fraction, and

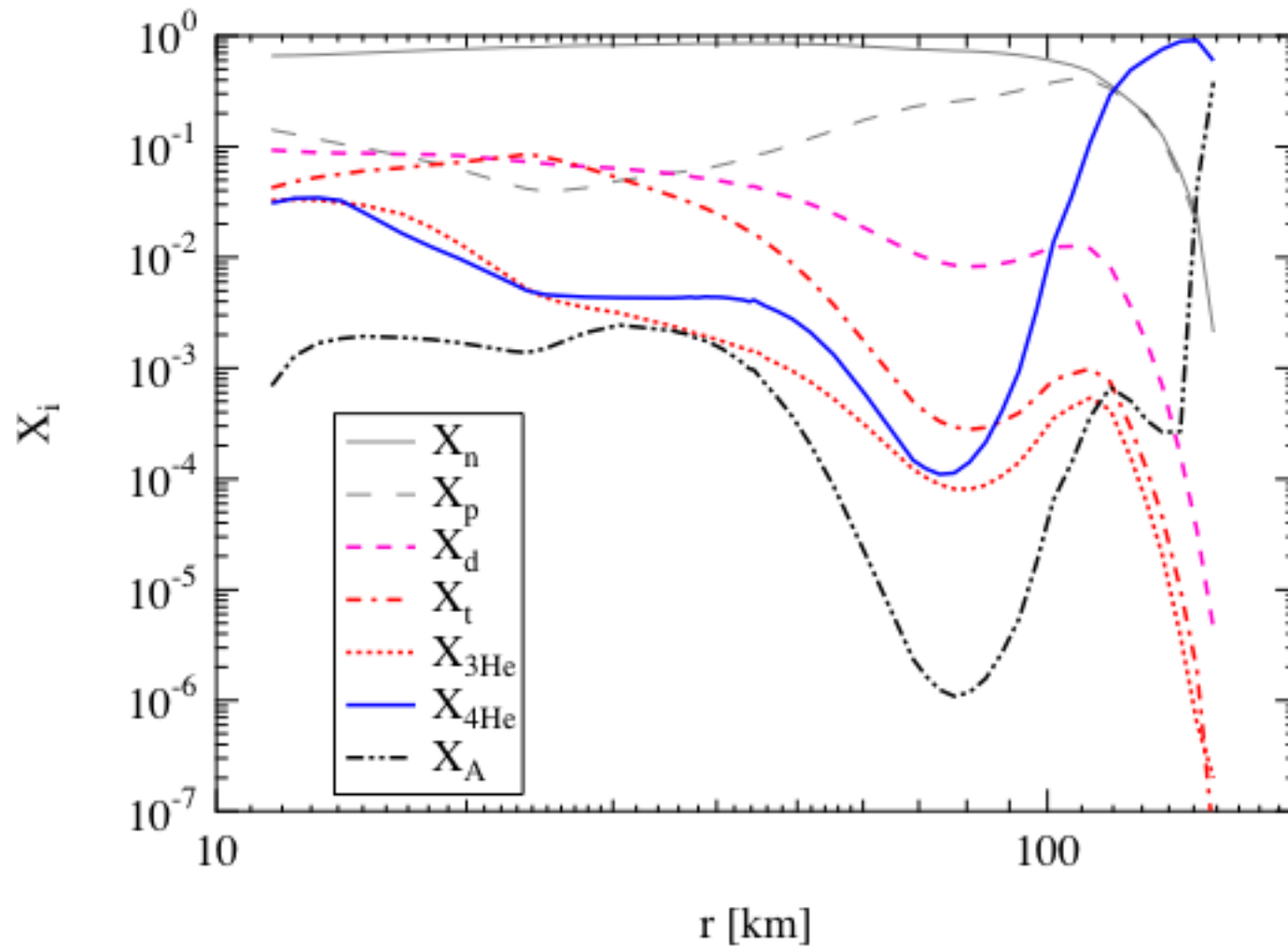
temperature profile

of a 15 solar mass supernova  
at 150 ms after core bounce  
as function of the radius.

Influence of cluster formation  
on neutrino emission  
in the cooling region and  
on neutrino absorption  
in the heating region ?

K.Sumiyoshi et al.,  
*Astrophys.J.* **629**, 922 (2005)

# Composition of supernova core



K. Sumiyoshi,  
G. R.,  
PRC 77,  
055804 (2008)

Mass fraction  $X$  of light clusters for a post-bounce supernova core

# Nuclear matter phase diagram

## Core collapse supernovae

### Relevant Parameters:

- **density:**

$$10^{-9} \lesssim \varrho/\varrho_{\text{sat}} \lesssim 10$$

with nuclear saturation density

$$\varrho_{\text{sat}} \approx 2.5 \cdot 10^{14} \text{ g/cm}^3$$

$$(n_{\text{sat}} = \varrho_{\text{sat}}/m_n \approx 0.15 \text{ fm}^{-3})$$

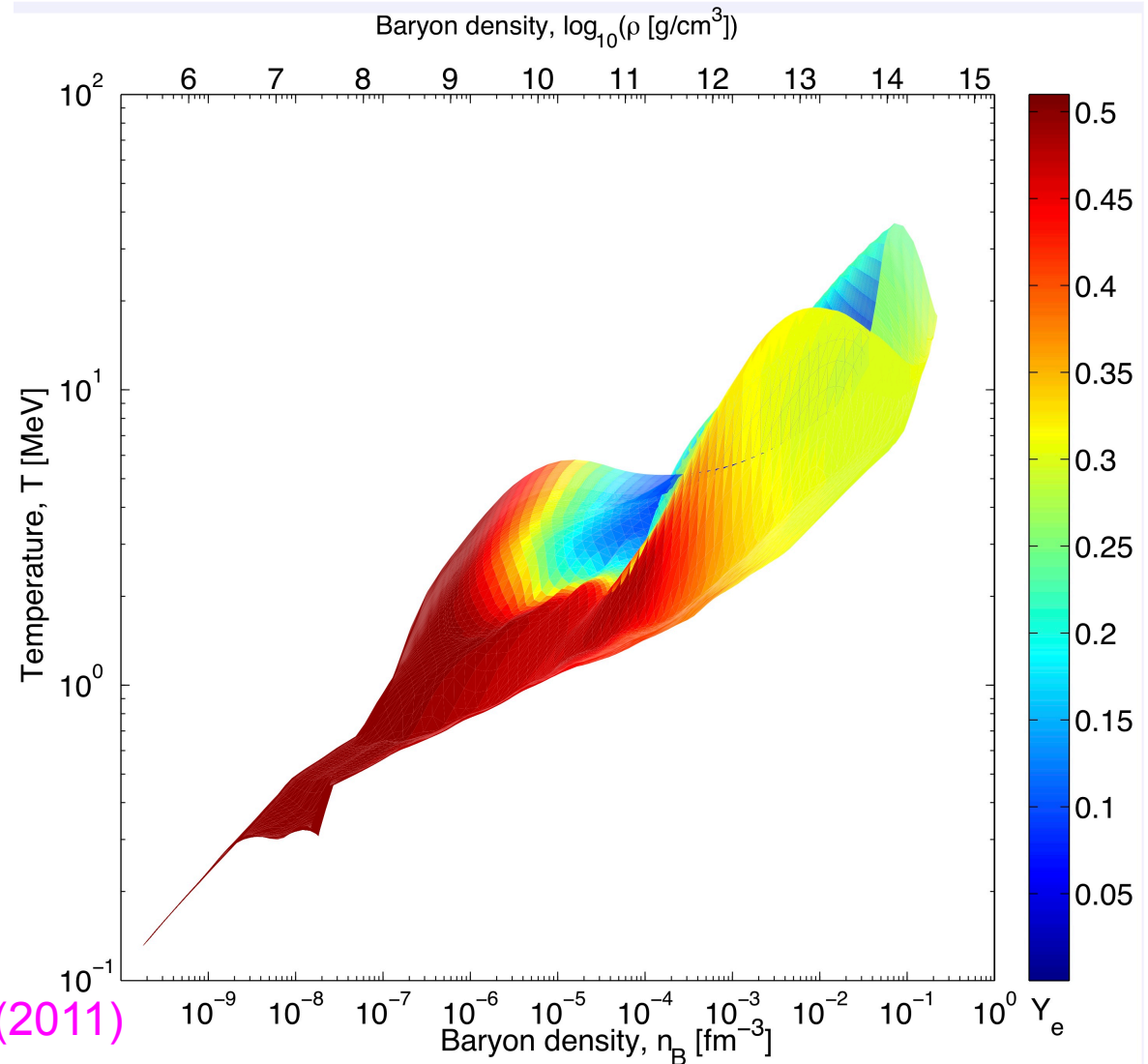
- **temperature:**

$$0 \text{ MeV} \leq k_B T \lesssim 50 \text{ MeV}$$

$$(\hat{=} 5.8 \cdot 10^{11} \text{ K})$$

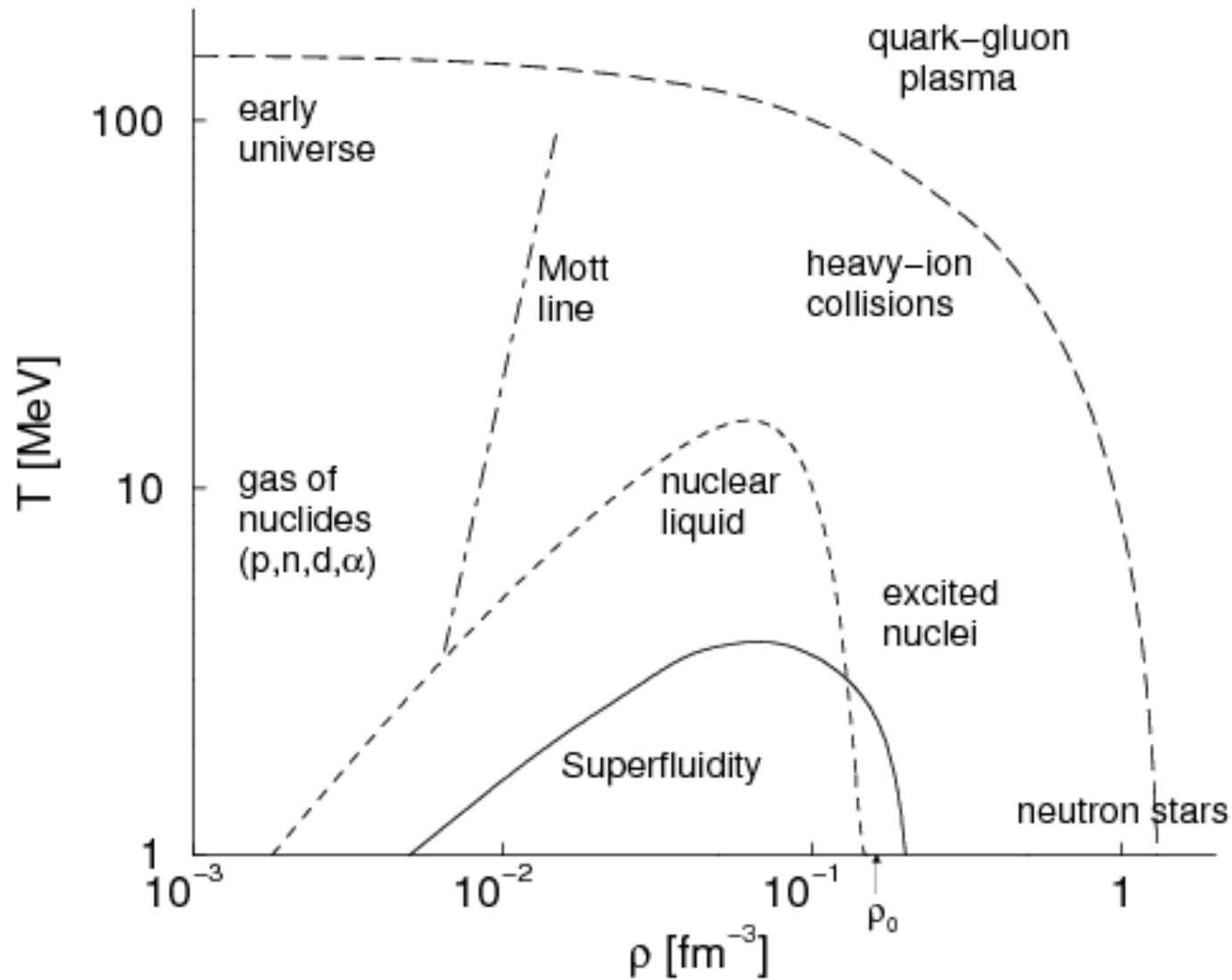
- **electron fraction:**

$$0 \leq Y_e \lesssim 0.6$$



T. Fischer et al., ApJS 194, 39 (2011)

## Symmetric nuclear matter: Phase diagram

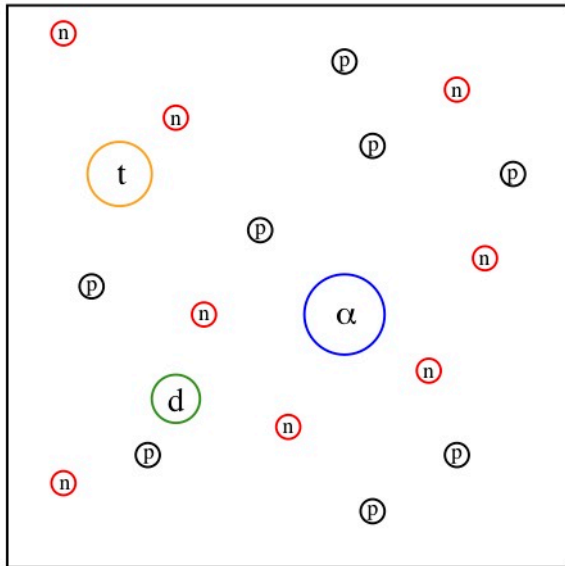


# Nuclear statistical equilibrium (NSE)

Chemical picture:

Ideal mixture of reacting components

Mass action law





## Ideal mixture of reacting nuclides

$$n_p(T, \mu_p, \mu_n) = \frac{1}{V} \sum_{A,\nu,K} Z_A f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

$$n_n(T, \mu_p, \mu_n) = \frac{1}{V} \sum_{A,\nu,K} (A - Z_A) f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

mass number  $A$ ,

charge  $Z_A$ ,

energy  $E_{A,\nu,K}$ ,

$\nu$  internal quantum number,

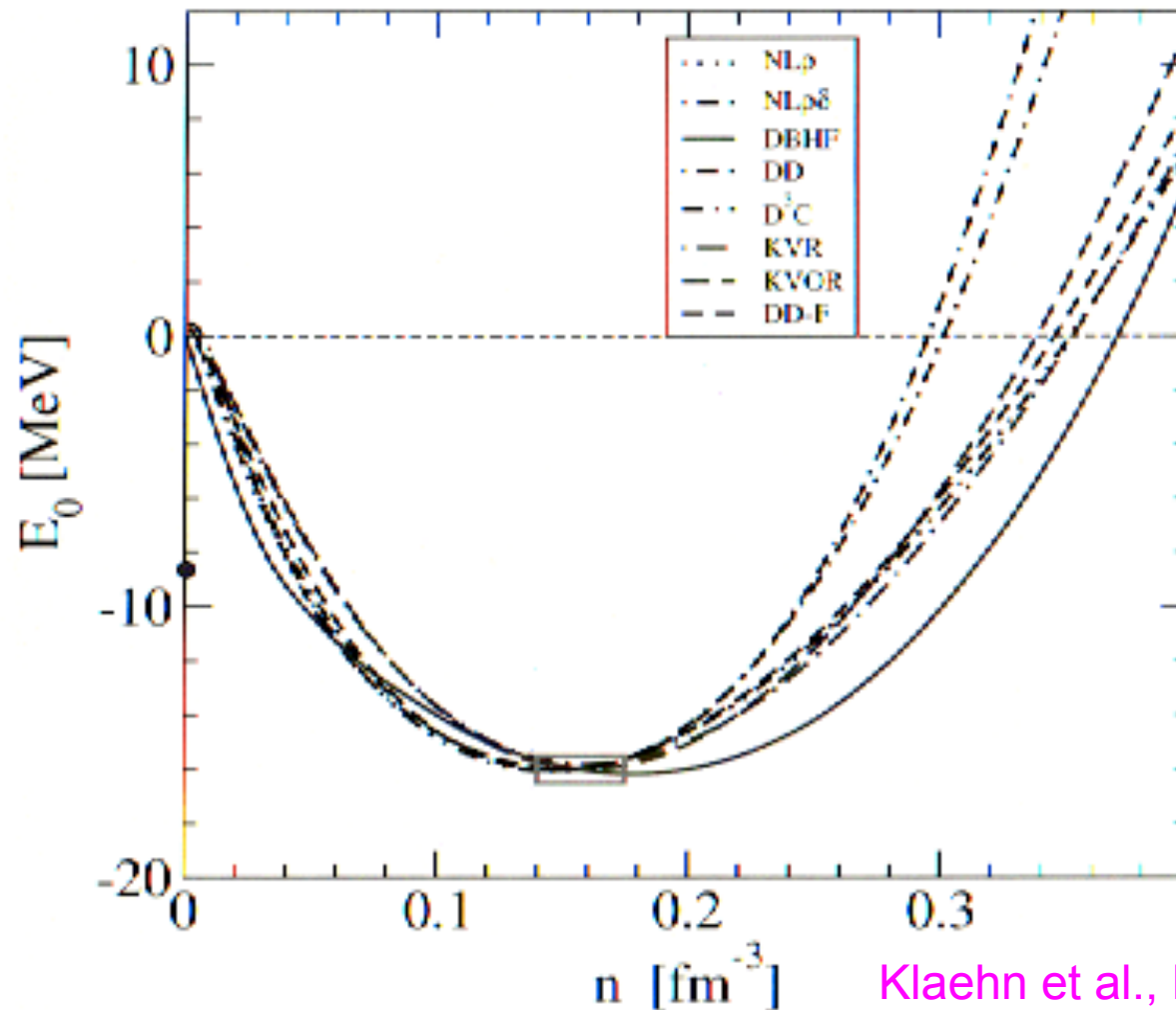
$K$ : center of mass momentum

$$f_A(z) = \frac{1}{\exp(z/T) - (-1)^A}$$

Nuclear Statistical Equilibrium  
(NSE)

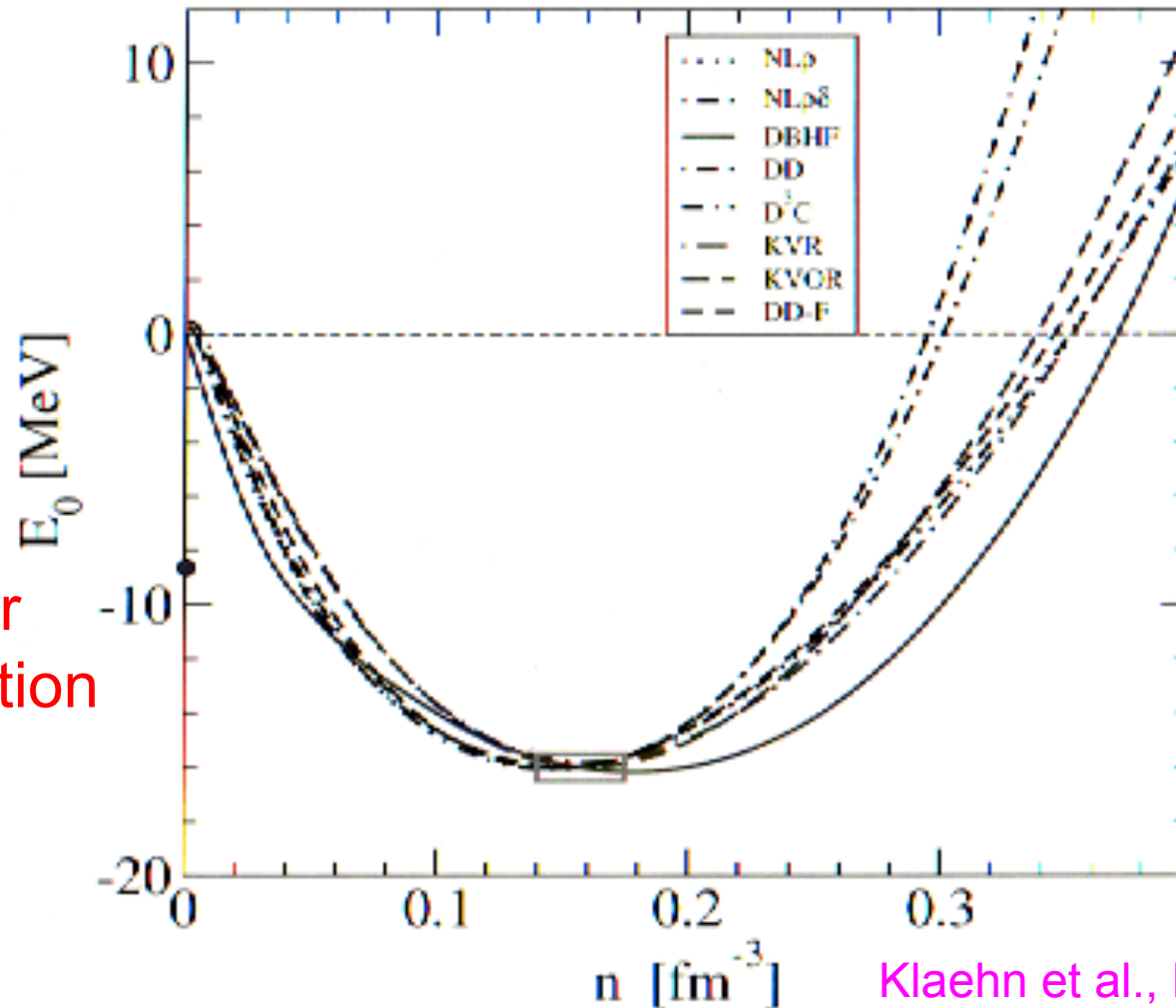
# Quasiparticle approximation for nuclear matter

## Equation of state for symmetric matter



# Quasiparticle approximation for nuclear matter

## Equation of state for symmetric matter



But:  
cluster  
formation

Incorrect  
low-density  
limit

Klaehn et al., PRC 2006

# Outline

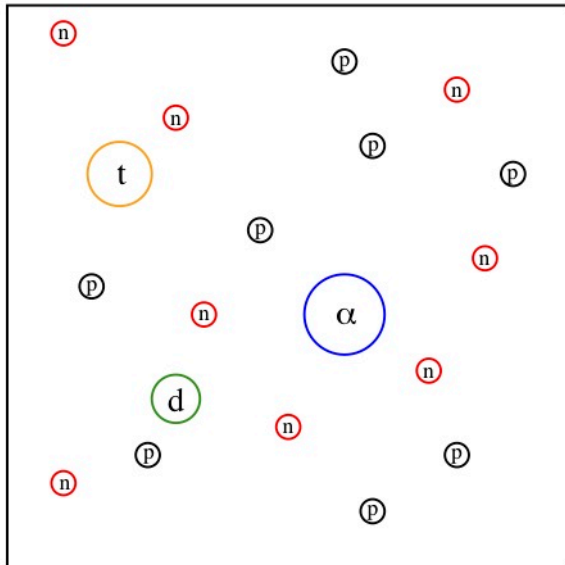
- Nuclear systems: Quasiparticle approach  
Brueckner, HFB; Skyrme, Relativistic Mean Field (RMF)
- Account of correlations in warm dense matter:  
two-particle (deuteron, pairing),  
four-particle (alpha-like) correlations, light elements
- Low-density regions: Nuclear Statistical Equilibrium (NSE)  
Hoyle-like states in light expanded nuclei,  
surface of nuclei, neck emission, alpha matter...
- Quantum statistical approach ( $n < 0.15 \text{ fm}^{-3}$ ,  $T < 20 \text{ MeV}$ )  
Equation of state, Beth-Uhlenbeck formula  
disappearance of clusters at high densities, Pauli blocking
- Experimental signatures  
Heavy Ion Collisions (HIC), Symmetry energy, SN explosions, ...

# Nuclear statistical equilibrium (NSE)

Chemical picture:

Ideal mixture of reacting components

Mass action law

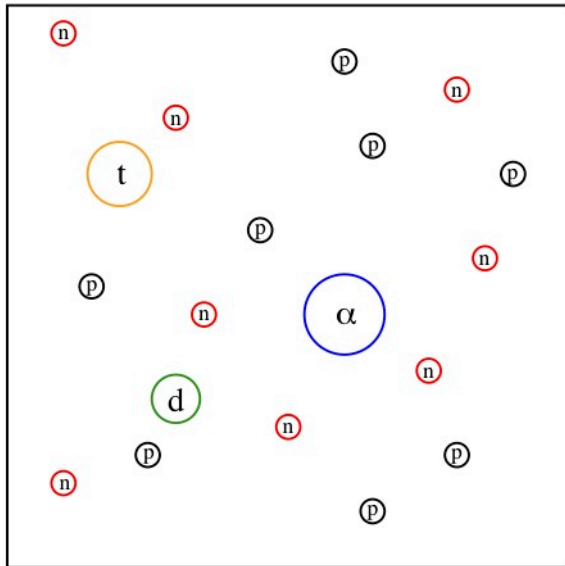


Interaction between the components  
internal structure: Pauli principle

# Nuclear statistical equilibrium (NSE)

Chemical picture:

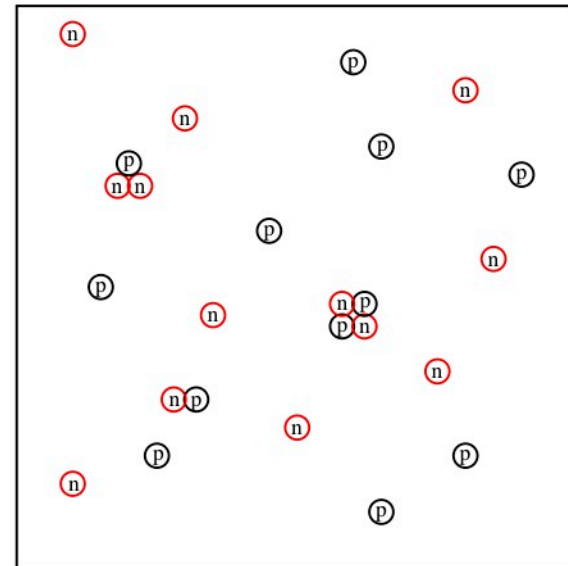
Ideal mixture of reacting components  
Mass action law



Interaction between the components  
internal structure: Pauli principle

Physical picture:

"elementary" constituents  
and their interaction



Quantum statistical (QS) approach,  
quasiparticle concept, virial expansion

# Many-particle theory

- equilibrium correlation functions

e.g. equation of state  $n(\beta, \mu) = \frac{1}{\Omega_0} \sum_1 \langle a_1^\dagger a_1 \rangle$

density matrix  $\langle a_1^\dagger a_1^\dagger \rangle = \int \frac{d\omega}{2\pi} e^{-i\omega t} f_1(\omega) A(1, 1', \omega)$

- Spectral function

$$A(1, 1', \omega) = \text{Im} [G(1, 1', \omega + i\eta) - G(1, 1', \omega - i\eta)]$$

- Matsubara Green function

$$G(1, 1', iz_\nu), \quad z_\nu = \frac{\pi\nu}{\beta} + \mu, \quad \nu = \pm 1, \pm 3, \dots$$

$$1 \equiv \{\mathbf{p}_1, \sigma_1, c_1\}, \quad f_1(\omega) = \frac{1}{e^{\beta(\omega - \mu)} + 1}, \quad \Omega_0 = \text{volume}$$

# Many-particle theory

- Dyson equation and self energy (homogeneous system)

$$G(1, iz_\nu) = \frac{1}{iz_\nu - E(1) - \Sigma(1, iz_\nu)}$$

- Evaluation of  $\Sigma(1, iz_\nu)$ :  
perturbation expansion, diagram representation

$$A(1, \omega) = \frac{2\text{Im } \Sigma(1, \omega + i0)}{[\omega - E(1) - \text{Re } \Sigma(1, \omega)]^2 + [\text{Im } \Sigma(1, \omega + i0)]^2}$$

approximation for  
self energy



approximation for  
equilibrium correlation functions

alternatively: simulations, path integral methods



# Different approximations

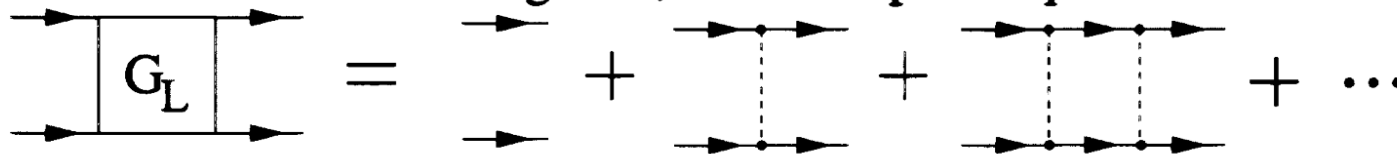
- Expansion for small  $\text{Im } \Sigma(1, \omega + i\eta)$

$$A(1, \omega) \approx \frac{2\pi\delta(\omega - E^{\text{quasi}}(1))}{1 - \frac{d}{dz} \text{Re } \Sigma(1, z)|_{z=E^{\text{quasi}} - \mu_1}} - 2\text{Im } \Sigma(1, \omega + i\eta) \frac{d}{d\omega} \frac{P}{\omega + \mu_1 - E^{\text{quasi}}(1)}$$

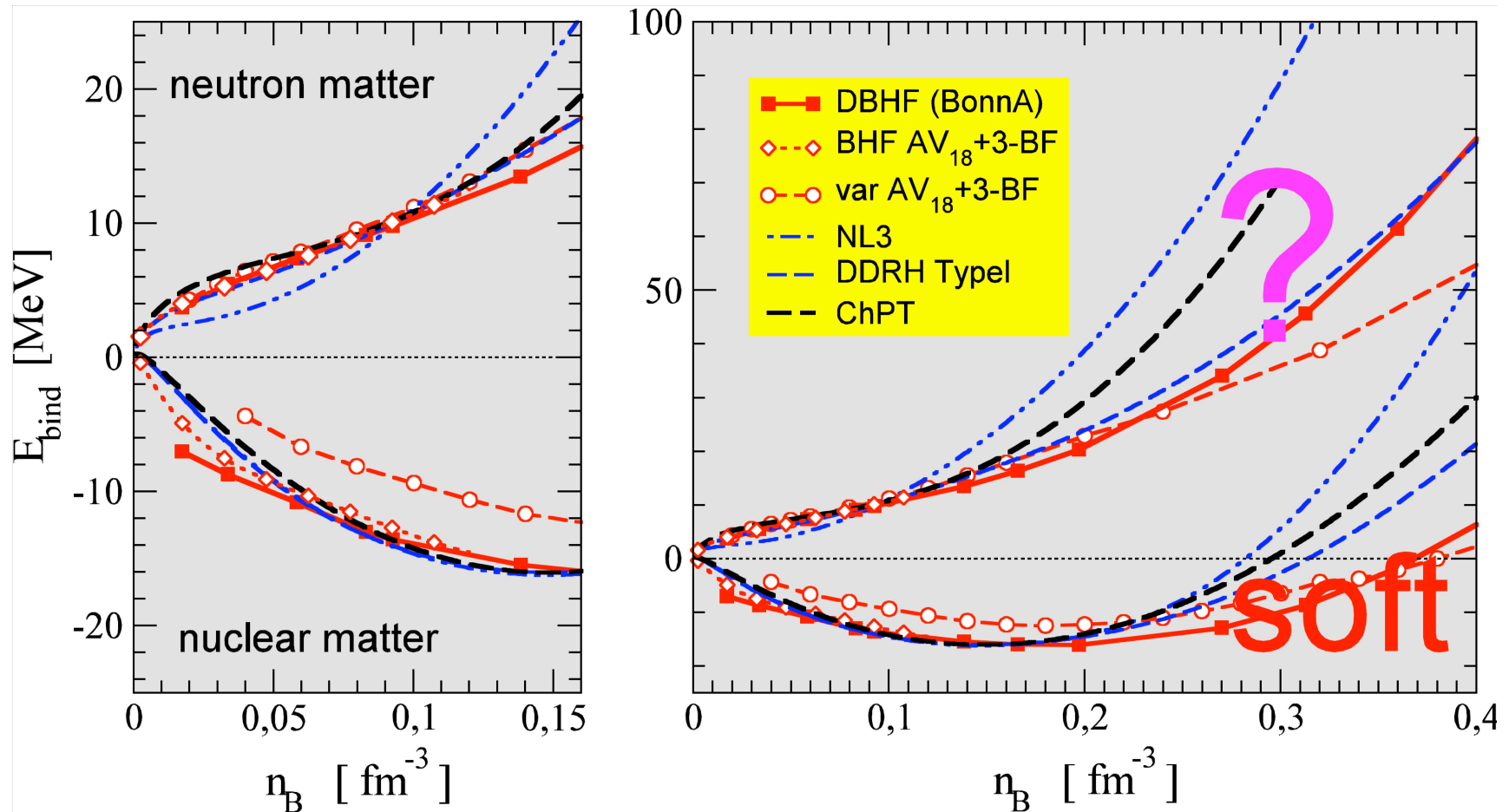
quasiparticle energy  $E^{\text{quasi}}(1) = E(1) + \text{Re } \Sigma(1, \omega)|_{\omega=E^{\text{quasi}}}$

- chemical picture: bound states  $\hat{=}$  new species

summation of ladder diagrams, Bethe-Salpeter equation

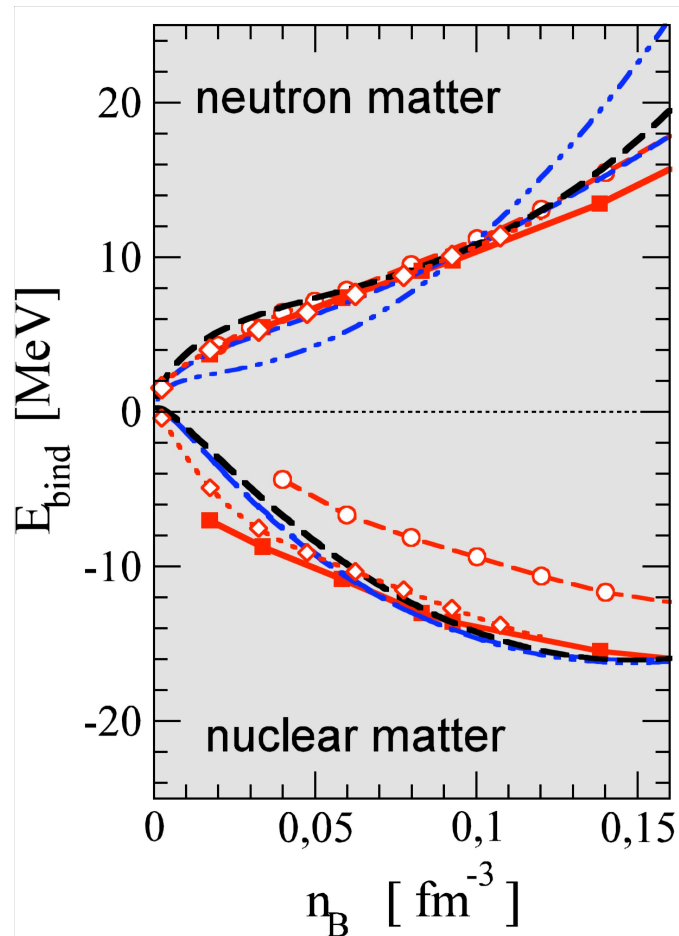


# Quasiparticle picture: RMF and DBHF



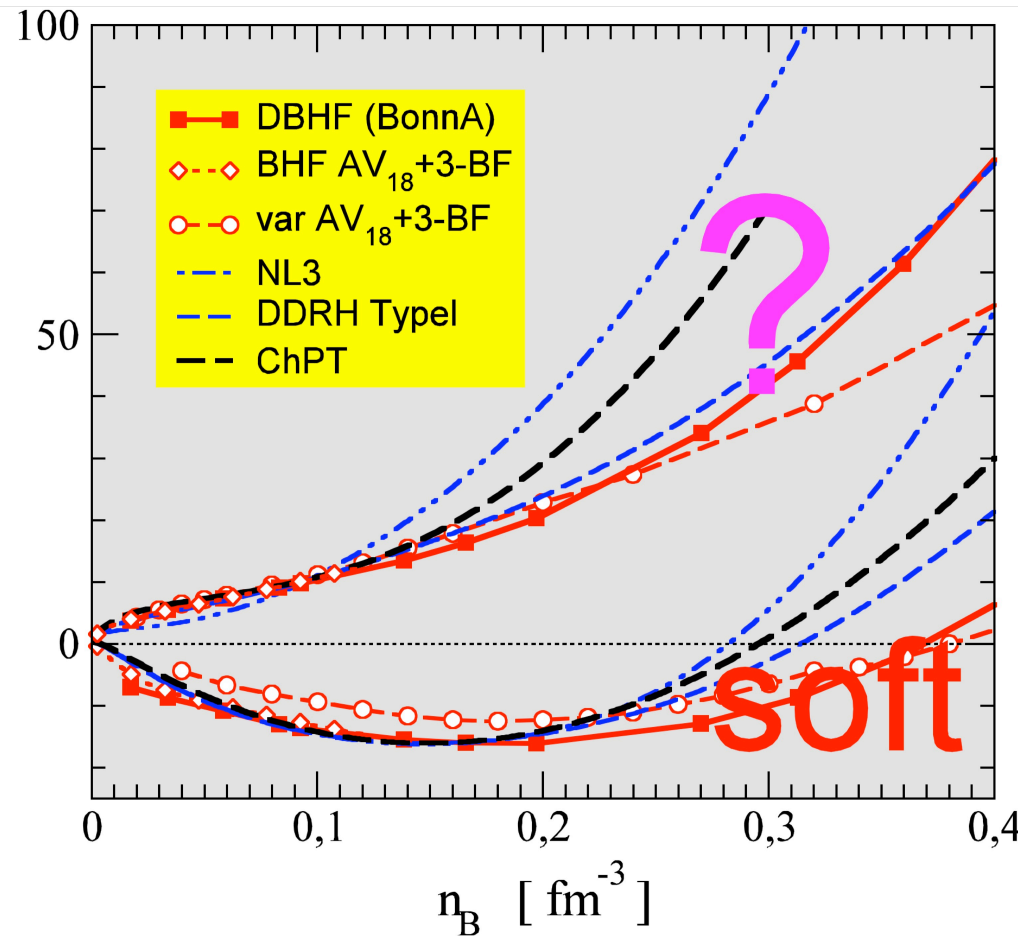
C. Fuchs et al.;  
 J. Margueron et al., Phys.Rev.C 76,034309 (2007)

# Quasiparticle picture: RMF and DBHF



But: cluster formation

Incorrect low-density limit



C. Fuchs et al.;

J.Margueron et al., Phys.Rev.C 76,034309 (2007)

# Different approximations

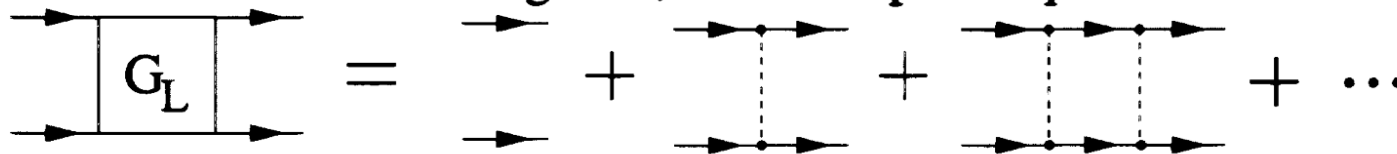
- Expansion for small  $\text{Im } \Sigma(1, \omega + i\eta)$

$$A(1, \omega) \approx \frac{2\pi\delta(\omega - E^{\text{quasi}}(1))}{1 - \frac{d}{dz} \text{Re } \Sigma(1, z)|_{z=E^{\text{quasi}} - \mu_1}} - 2\text{Im } \Sigma(1, \omega + i\eta) \frac{d}{d\omega} \frac{P}{\omega + \mu_1 - E^{\text{quasi}}(1)}$$

quasiparticle energy  $E^{\text{quasi}}(1) = E(1) + \text{Re } \Sigma(1, \omega)|_{\omega=E^{\text{quasi}}}$

- chemical picture: bound states  $\hat{=}$  new species

summation of ladder diagrams, Bethe-Salpeter equation



# Different approximations

low density limit:

$$G_2^L(12, 1'2', i\lambda) = \sum_{n\mathbf{P}} \Psi_{n\mathbf{P}}(12) \frac{1}{i\omega_\lambda - E_{n\mathbf{P}}} \Psi_{n\mathbf{P}}^*(12)$$

$$\Sigma = \text{[Diagram: a box labeled } T_2^L \text{ with a loop above it containing an arrow pointing left]}$$

$$n(\beta, \mu) = \sum_1 f_1(E^{\text{quasi}}(1)) + \sum_{2, n\mathbf{P}}^{\text{bound}} g_{12}(E_{n\mathbf{P}}) + \sum_{2, n\mathbf{P}} \int_0^\infty dk \delta_{\mathbf{k}, \mathbf{p}_1 - \mathbf{p}_2} g_{12}(E^{\text{quasi}}(1) + E^{\text{quasi}}(2)) 2 \sin^2 \delta_n(k) \frac{1}{\pi} \frac{d}{dk} \delta_n(k)$$

- generalized Beth-Uhlenbeck formula  
correct low density/low temperature limit:  
mixture of free particles and bound clusters

# Effective wave equation for the deuteron in matter

In-medium two-particle wave equation in mean-field approximation

$$\left( \frac{p_1^2}{2m_1} + \Delta_1 + \frac{p_2^2}{2m_2} + \Delta_2 \right) \Psi_{d,P}(p_1, p_2) + \sum_{p_1', p_2'} (1 - f_{p_1} - f_{p_2}) V(p_1, p_2; p_1', p_2') \Psi_{d,P}(p_1', p_2')$$

Add self-energy

Pauli-blocking

$$= E_{d,P} \Psi_{d,P}(p_1, p_2)$$

Thouless criterion

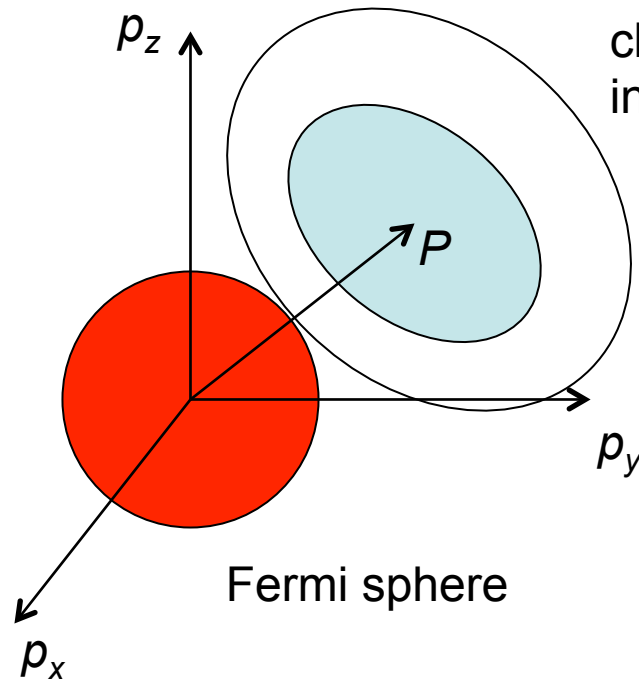
$$E_d(T, \mu) = 2\mu$$

Fermi distribution function

$$f_p = \left[ e^{(p^2/2m - \mu)/k_B T} + 1 \right]^{-1}$$

BEC-BCS crossover:  
Alm et al., 1993

# Pauli blocking – phase space occupation



cluster wave function (deuteron, alpha,...)  
in momentum space

$P$  - center of mass momentum

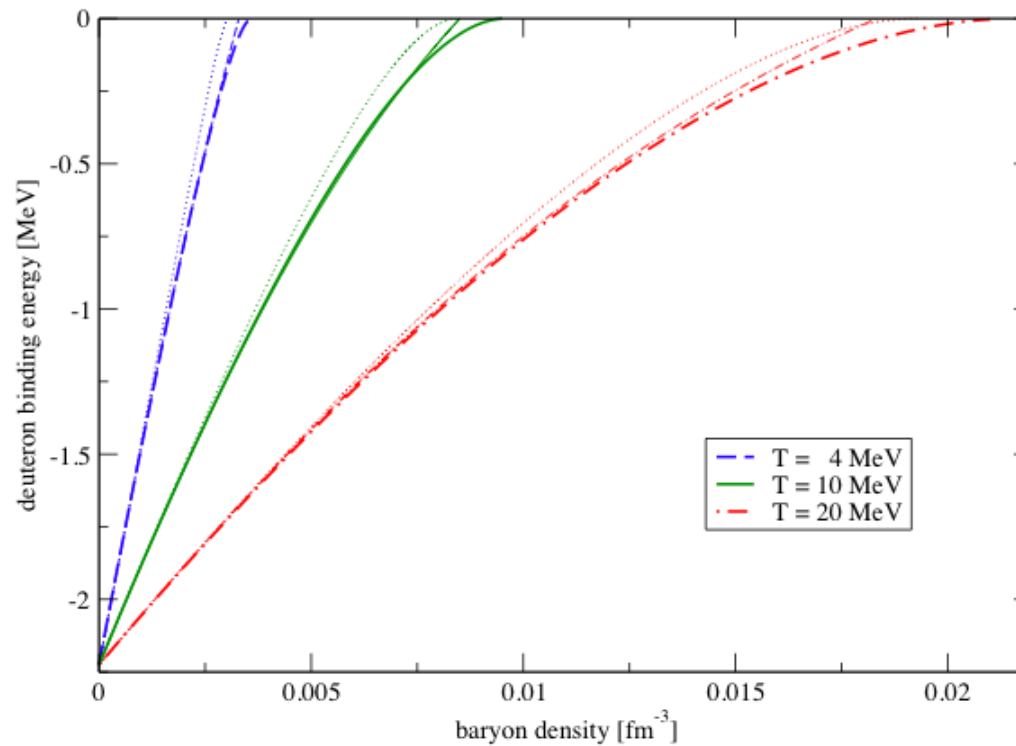
The Fermi sphere is forbidden,  
deformation of the cluster wave function  
in dependence on the c.o.m. momentum  $P$

momentum space

The deformation is maximal at  $P = 0$ .  
It leads to the weakening of the interaction  
(disintegration of the bound state).

# Shift of the deuteron binding energy

Dependence on nucleon density, various temperatures,  
zero center of mass momentum



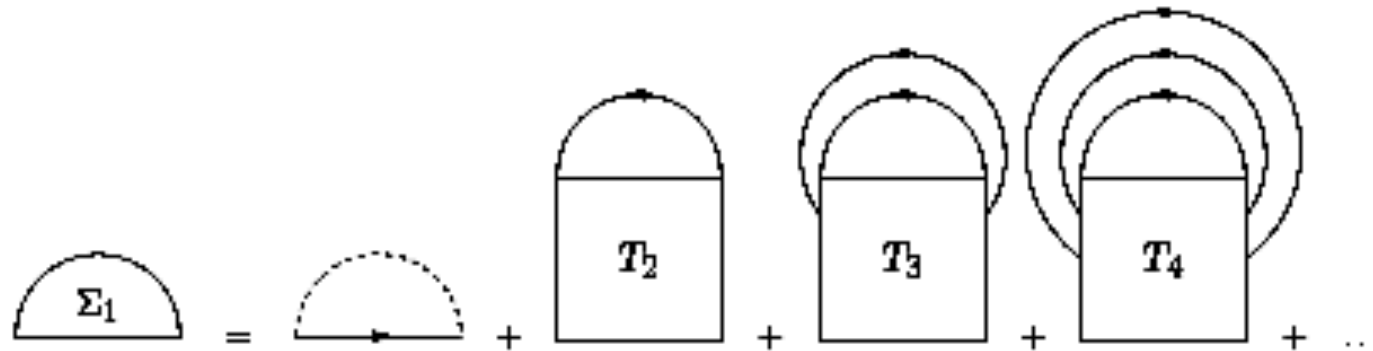
thin lines:

fit formula

G.R., NP A 867, 66 (2011)



# Cluster decomposition of the self-energy



T-matrices: bound states, scattering states  
Including clusters like new components  
chemical picture,  
mass action law, nuclear statistical equilibrium (NSE)

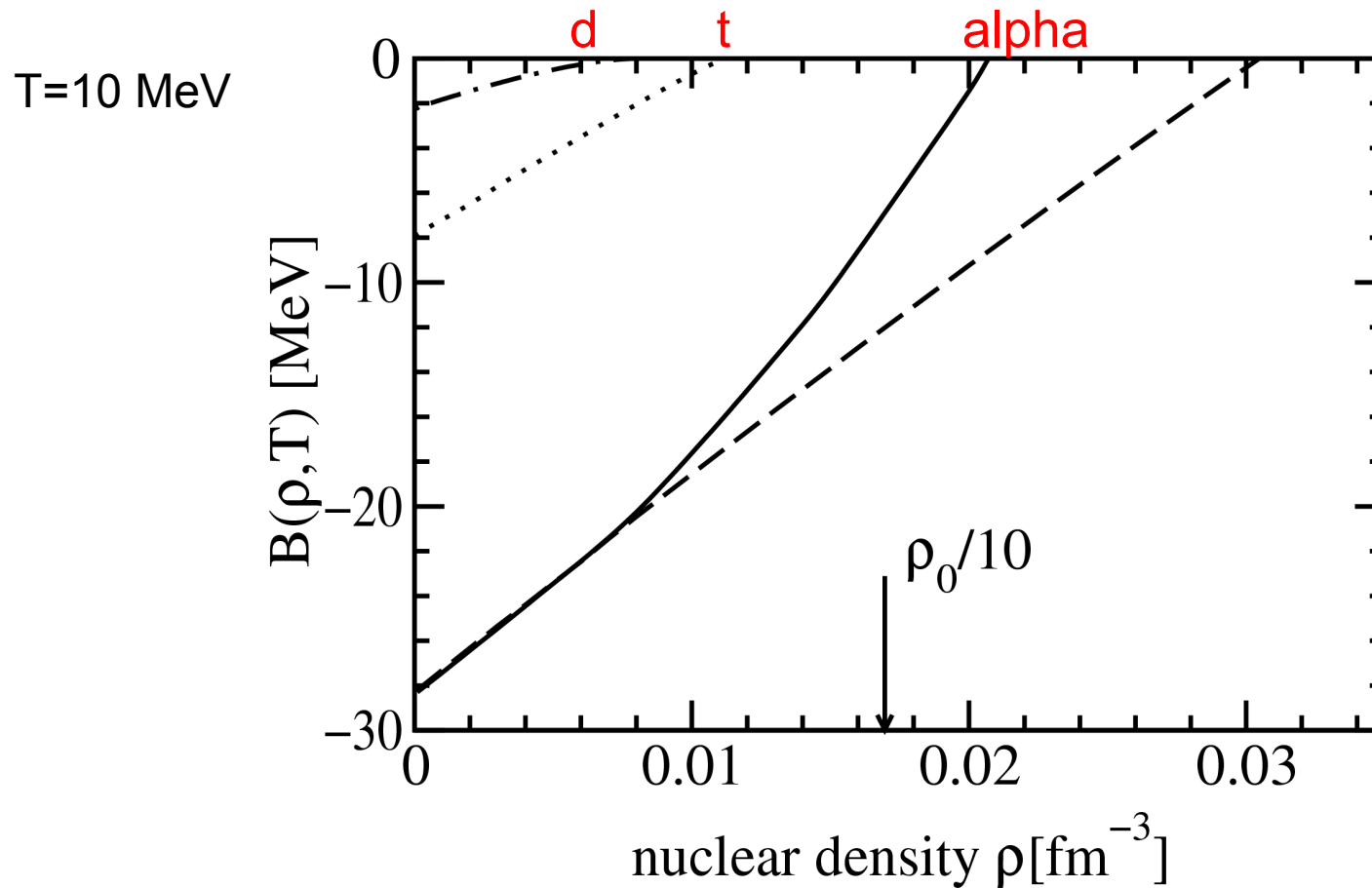
# Few-particle Schrödinger equation in a dense medium

4-particle Schrödinger equation with medium effects

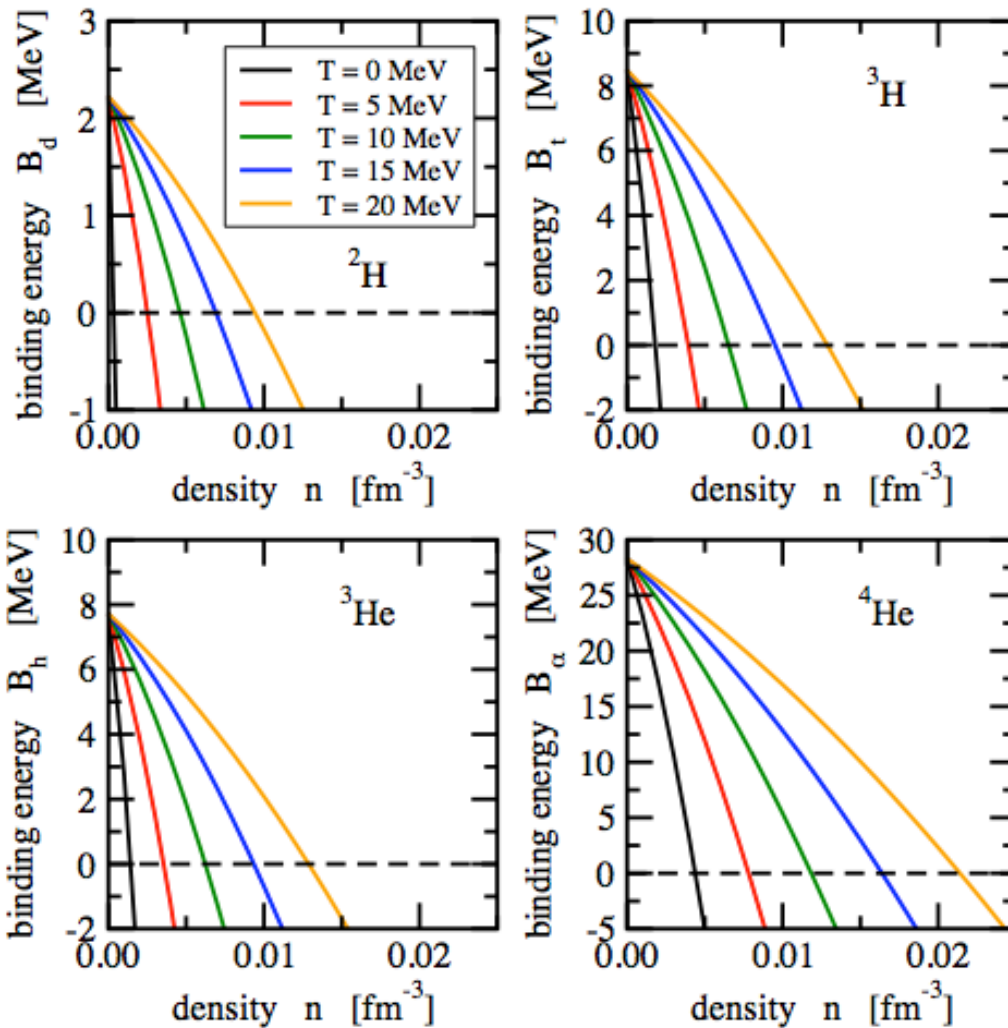
$$\begin{aligned} & \left( \left[ E^{HF}(p_1) + E^{HF}(p_2) + E^{HF}(p_3) + E^{HF}(p_4) \right] \right) \Psi_{n,P}(p_1, p_2, p_3, p_4) \\ & + \sum_{p_1', p_2'} (1 - f_{p_1} - f_{p_2}) V(p_1, p_2; p_1', p_2') \Psi_{n,P}(p_1', p_2', p_3, p_4) \\ & + \{ \text{permutations} \} \\ & = E_{n,P} \Psi_{n,P}(p_1, p_2, p_3, p_4) \end{aligned}$$

# In-medium shift of binding energies of clusters

Solution of the Faddeev-Yakubovskii equation with Pauli blocking



# Shift of Binding Energies of Light Clusters



Symmetric matter

G.R., PRC 79, 014002 (2009)  
S. Typel et al.,  
PRC 81, 015803 (2010)

# Composition of dense nuclear matter

$$n_p(T, \mu_p, \mu_n) = \frac{1}{V} \sum_{A,\nu,K} Z_A f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

$$n_n(T, \mu_p, \mu_n) = \frac{1}{V} \sum_{A,\nu,K} (A - Z_A) f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

mass number  $A$ ,

charge  $Z_A$ ,

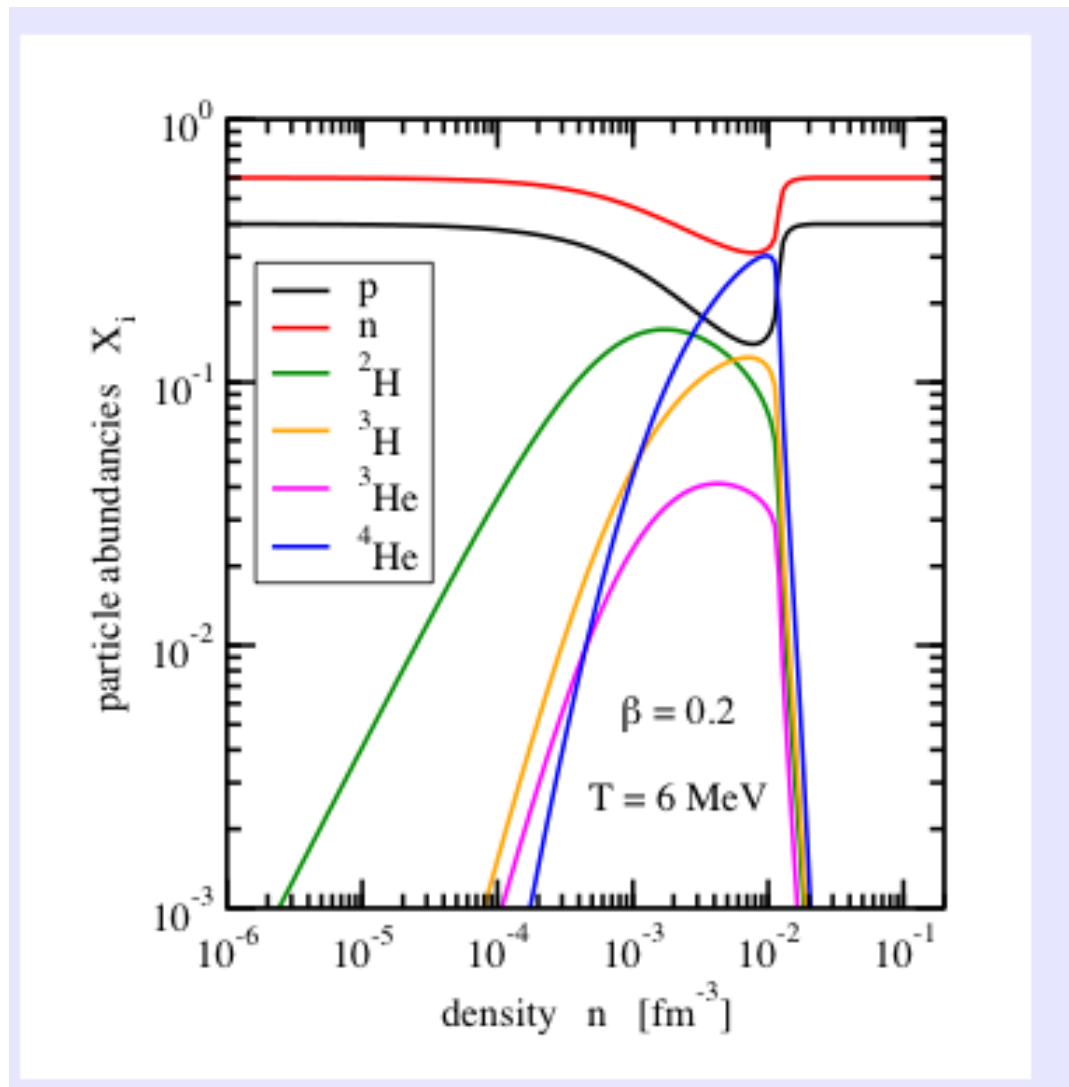
energy  $E_{A,\nu,K}$ ,

$\nu$ : internal quantum number,

$$f_A(z) = \frac{1}{\exp(z/T) - (-1)^A}$$

- Inclusion of excited states and continuum correlations
- Medium effects:
  - self-energy and **Pauli blocking shifts** of binding energies,
  - Coulomb corrections due to screening (Wigner-Seitz, Debye)

# Light Cluster Abundances



S. Typel et al.,  
PRC 81, 015803 (2010)

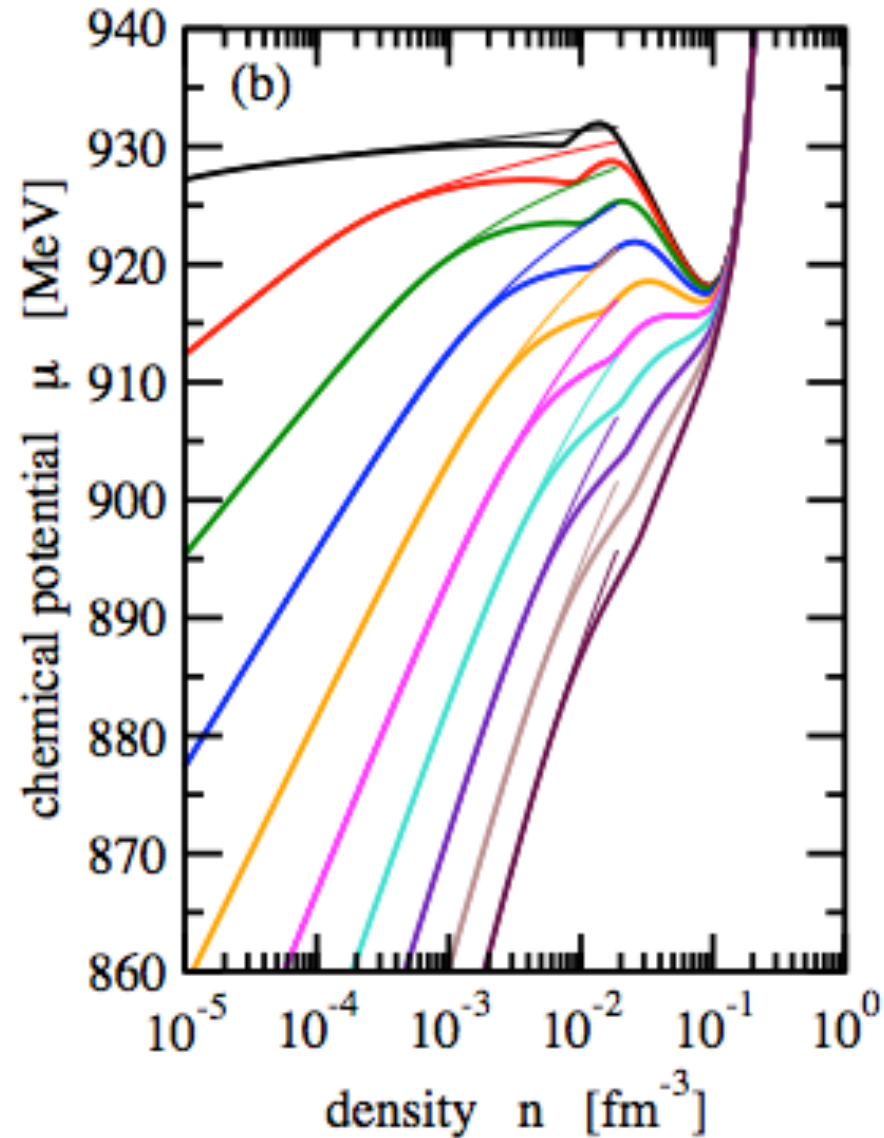
# Chemical potential of symmetric matter

Isotherms

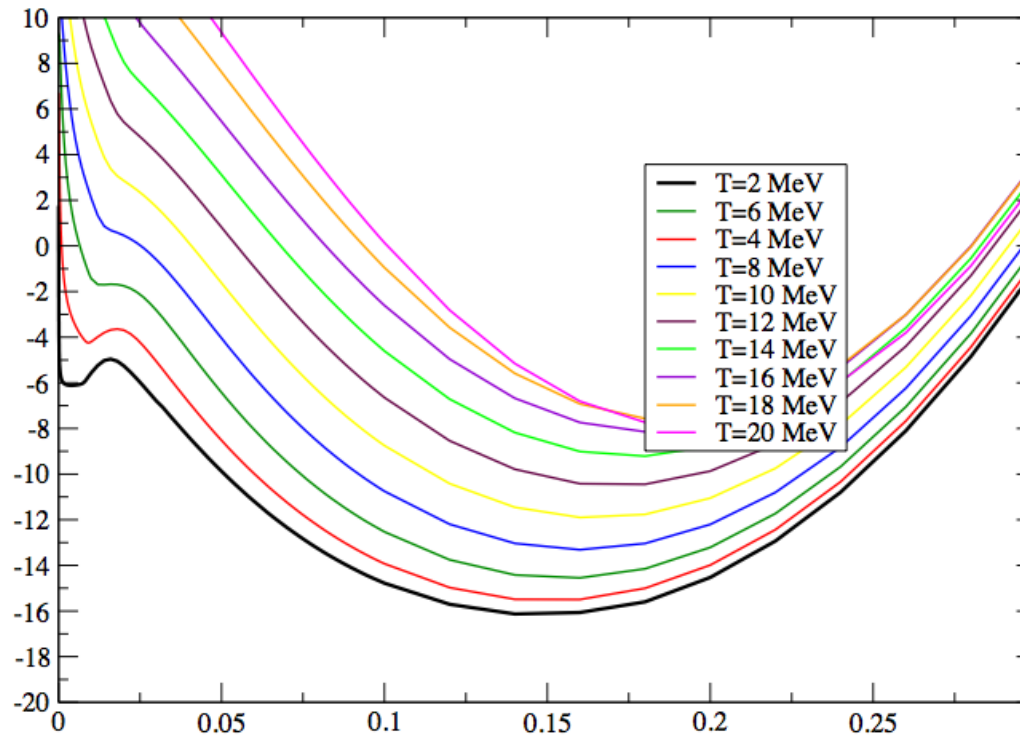
T[MeV]

2  
4  
6  
8  
10  
12  
14  
16  
18  
20

thin lines: NSE



# Internal energy per nucleon



Quantum  
statistical  
approach:

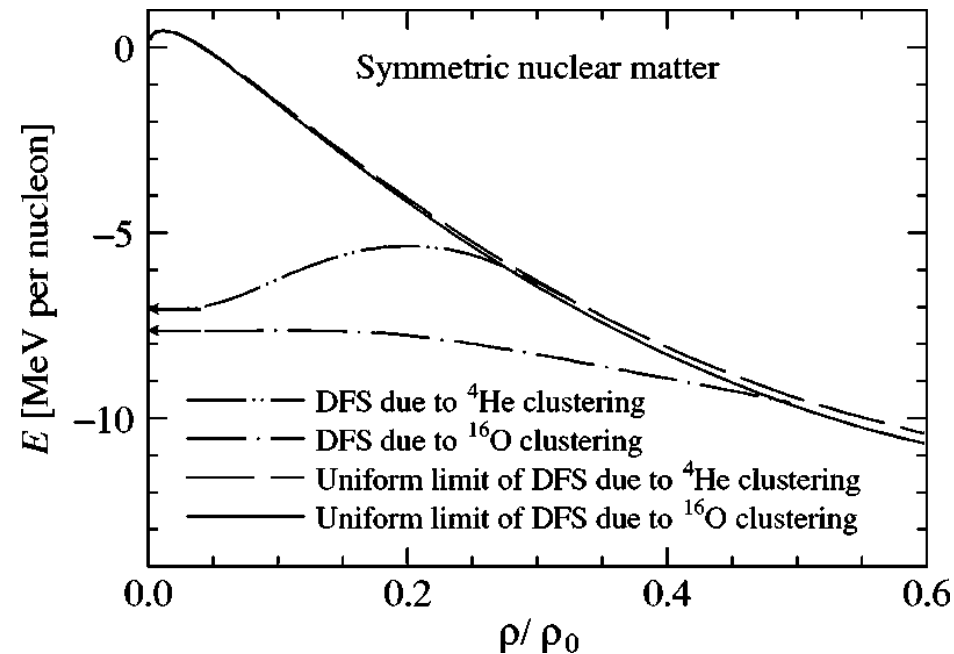
Cluster ?

Condensate?

EOS for symmetric matter - low density region?



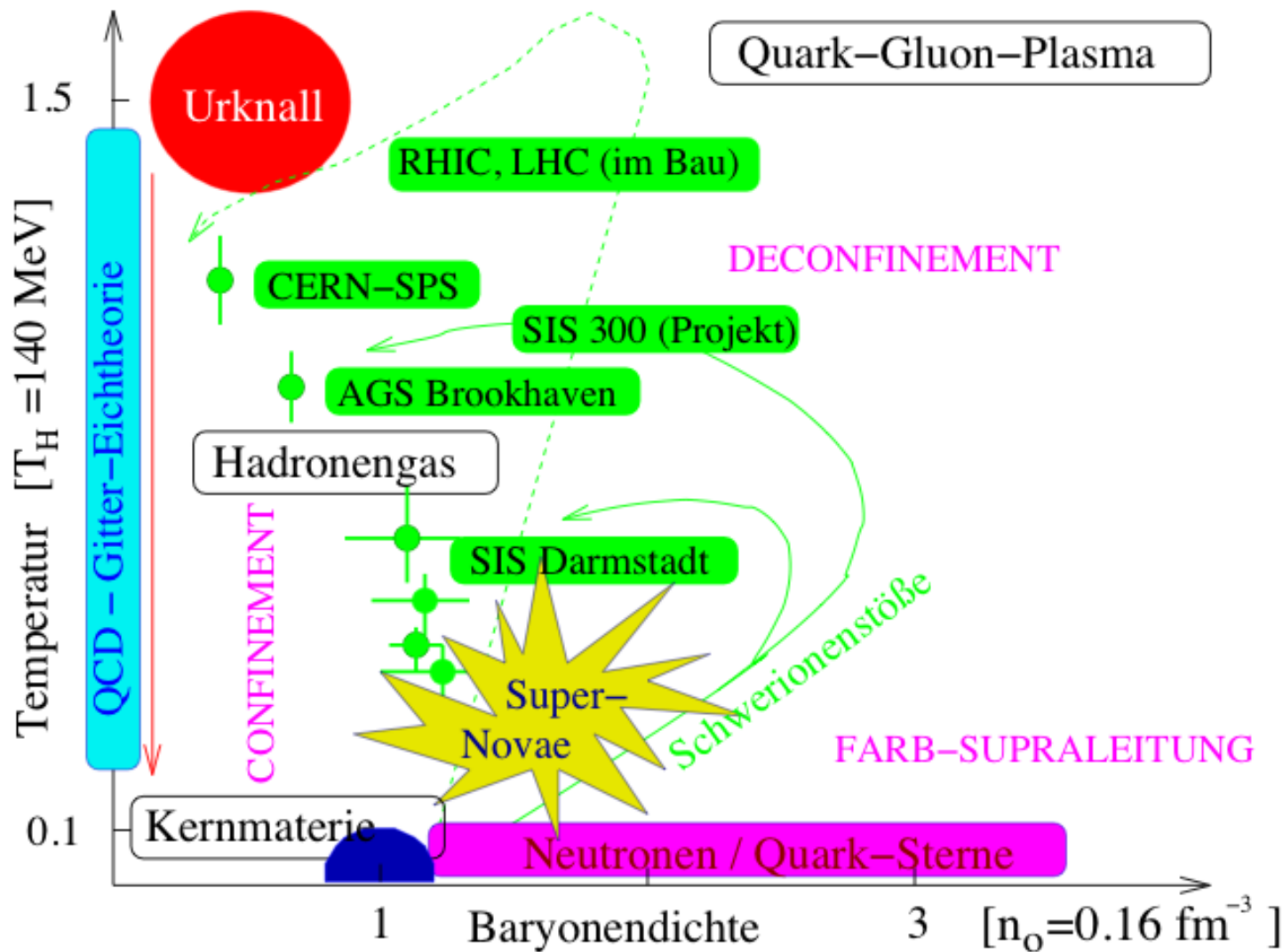
# Clustering phenomena in nuclear matter below the saturation density



•FIG. 8. Energy curves of DFSs due to  $\alpha$  and  $^{16}\text{O}$  clustering in  
•the symmetric nuclear matter by the use of the  $\text{BB}\text{-sB4d}$  force. The  
•density of matter is normalized by the saturation density of the  
•uniform matter with the Fermi sphere,  $\rho_0=0.206\text{ fm}^{-3}$ . The presentation  
•of the curves is similar to that in Fig. 4.

Hiroki Takemoto et al.,  
PR C 69, 035802 (2004)

# Phase diagram nuclear matter



# Application to Heavy Ion Reactions

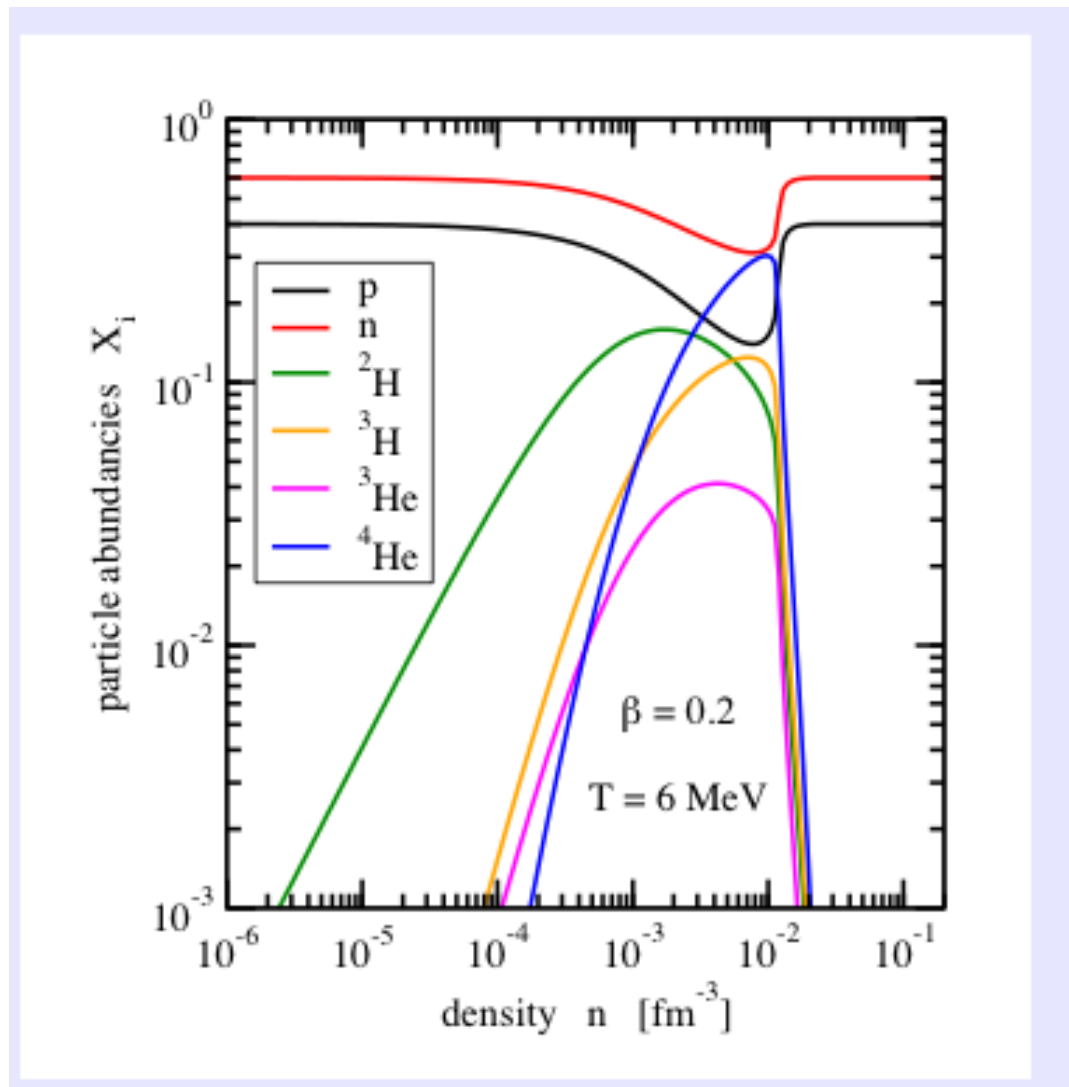
- Test the EOS  
(NSE, virial,... at low densities,  
Skyrme, DBHF, RMF,... near saturation)
- Unifying quantum statistical approach, [medium effects](#), Mott effect
- [Symmetry energy](#)
- [Bose enhancement?](#)

Nimrod @ TAMU,  
40Ar + 112,124Sn,  
64Zn + 112,124Sn; 47 A MeV

[Yields of p, \(n\), d, t, <sup>3</sup>He, <sup>4</sup>He,...](#)

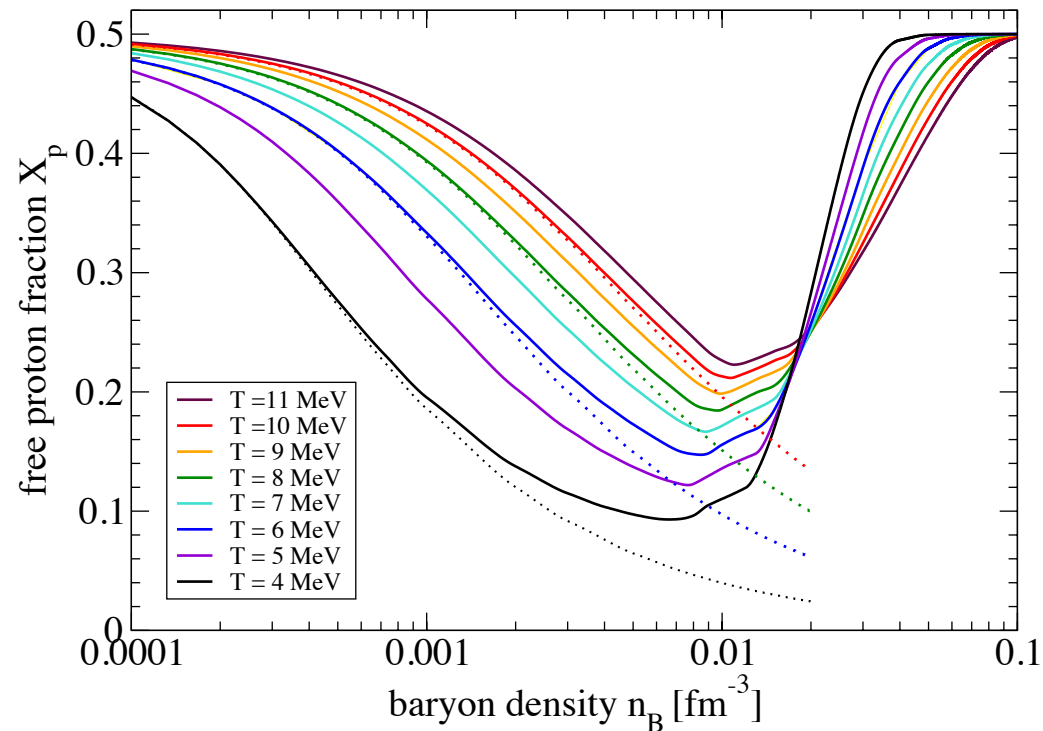
Open questions: freeze-out model or dynamical transport models?  
Identification of the source?

# Light Cluster Abundances



S. Typel et al.,  
PRC 81, 015803 (2010)

# Pauli blocking in symmetric matter



Free proton fraction as function of density and temperature in symmetric matter. QS calculations (solid lines) are compared with the NSE results (dotted lines). **Mott effect** in the region  $n_{\text{saturation}}/5$ .

# EOS at low densities from HIC

PRL 108, 172701 (2012)

PHYSICAL REVIEW LETTERS

week ending  
27 APRIL 2012

## Laboratory Tests of Low Density Astrophysical Nuclear Equations of State

L. Qin,<sup>1</sup> K. Hagel,<sup>1</sup> R. Wada,<sup>2,1</sup> J. B. Natowitz,<sup>1</sup> S. Shlomo,<sup>1</sup> A. Bonasera,<sup>1,3</sup> G. Röpke,<sup>4</sup> S. Typel,<sup>5</sup> Z. Chen,<sup>6</sup> M. Huang,<sup>6</sup> J. Wang,<sup>6</sup> H. Zheng,<sup>1</sup> S. Kowalski,<sup>7</sup> M. Barbui,<sup>1</sup> M. R. D. Rodrigues,<sup>1</sup> K. Schmidt,<sup>1</sup> D. Fabris,<sup>8</sup> M. Lunardon,<sup>8</sup> S. Moretto,<sup>8</sup> G. Nebbia,<sup>8</sup> S. Pesente,<sup>8</sup> V. Rizzi,<sup>8</sup> G. Viesti,<sup>8</sup> M. Cinausero,<sup>9</sup> G. Prete,<sup>9</sup> T. Keutgen,<sup>10</sup> Y. El Masri,<sup>10</sup> Z. Majka,<sup>11</sup> and Y. G. Ma<sup>12</sup>

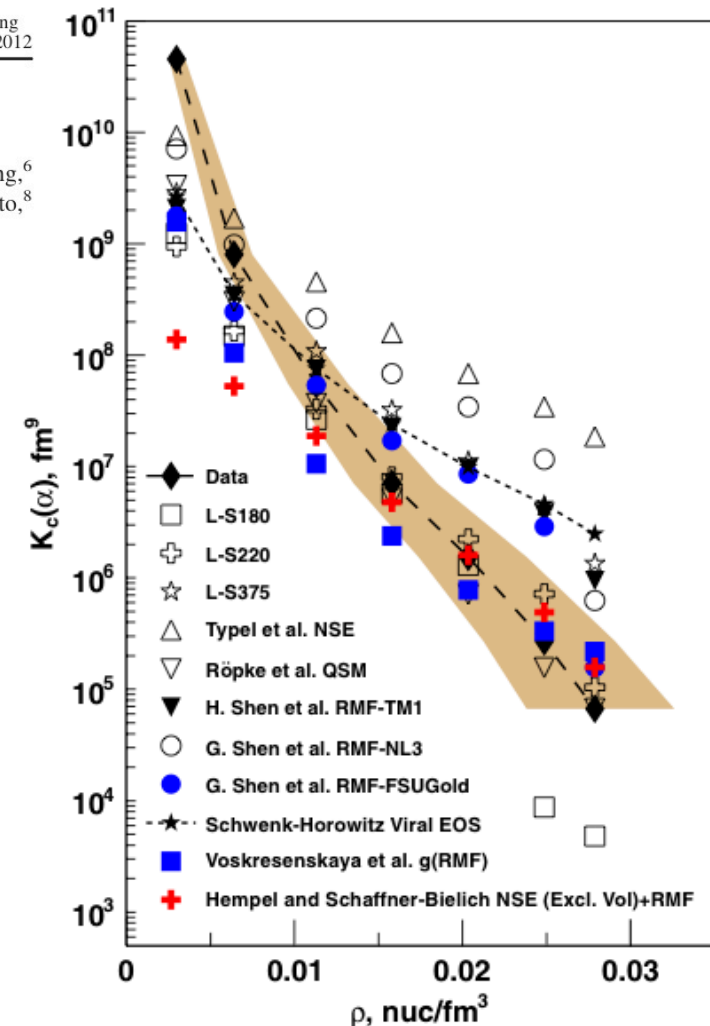
Yields of clusters from HIC: p, n, d, t, h,  $\alpha$

chemical constants

$$K_c(A, Z) = \rho_{(A,Z)} / [(\rho_p)^Z (\rho_n)^N]$$

Bose enhancement?

Symmetry energy



# Cluster yields in HIC

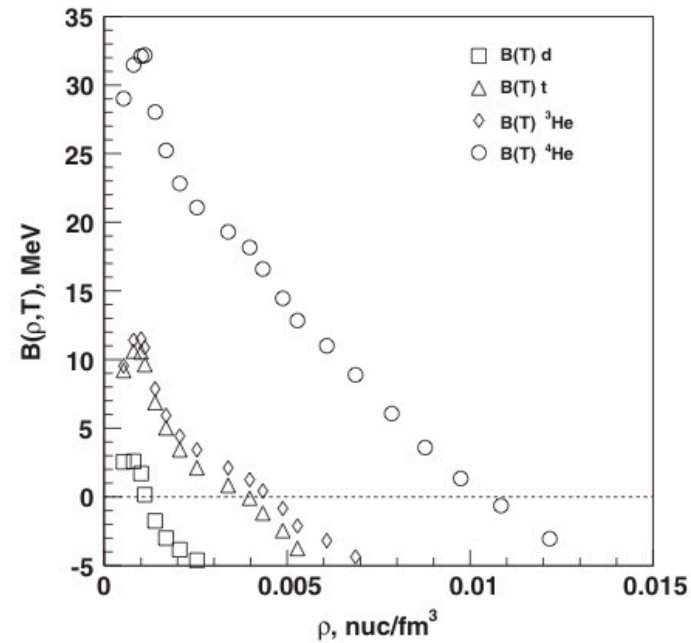
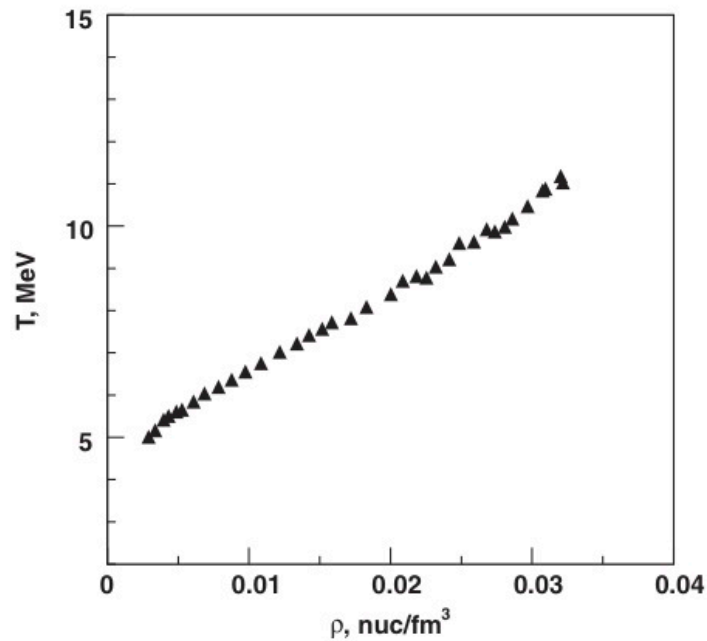
PRL 108, 062702 (2012)

PHYSICAL REVIEW LETTERS

week ending  
10 FEBRUARY 2012

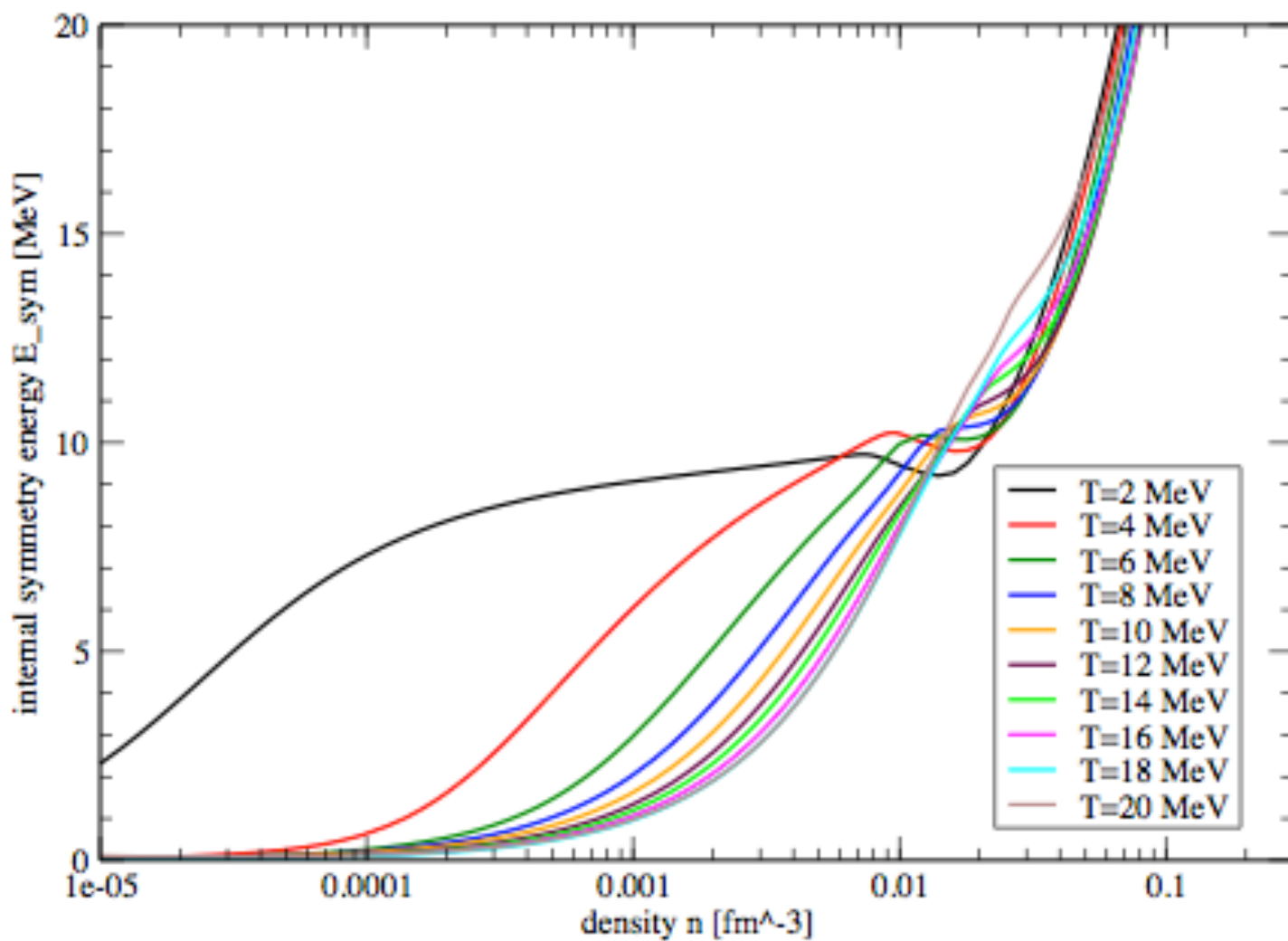
## Experimental Determination of In-Medium Cluster Binding Energies and Mott Points in Nuclear Matter

K. Hagel,<sup>1</sup> R. Wada,<sup>2,1</sup> L. Qin,<sup>1</sup> J. B. Natowitz,<sup>1</sup> S. Shlomo,<sup>1</sup> A. Bonasera,<sup>1,3</sup> G. Röpke,<sup>4</sup> S. Typel,<sup>5</sup> Z. Chen,<sup>2</sup> M. Huang,<sup>2</sup> J. Wang,<sup>2</sup> H. Zheng,<sup>1</sup> S. Kowalski,<sup>6</sup> C. Bottosso,<sup>1</sup> M. Barbui,<sup>1</sup> M. R. D. Rodrigues,<sup>1</sup> K. Schmidt,<sup>1</sup> D. Fabris,<sup>7</sup> M. Lunardon,<sup>7</sup> S. Moretto,<sup>7</sup> G. Nebbia,<sup>7</sup> S. Pesente,<sup>7</sup> V. Rizzi,<sup>7</sup> G. Viesti,<sup>7</sup> M. Cinausero,<sup>8</sup> G. Prete,<sup>8</sup> T. Keutgen,<sup>9</sup> Y. El Masri,<sup>9</sup> and Z. Majka<sup>10</sup>



in-medium binding energies

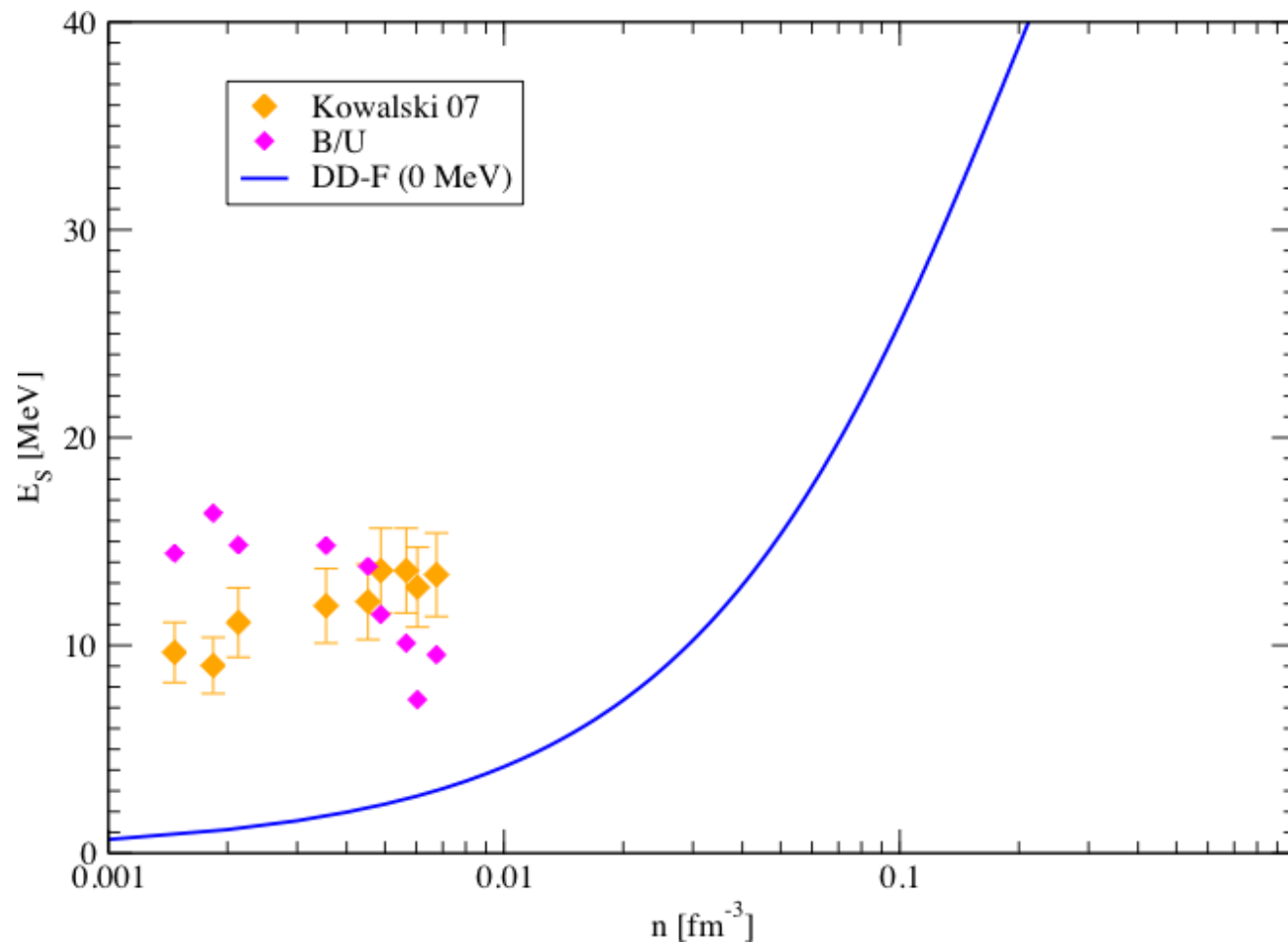
# Internal symmetry energy





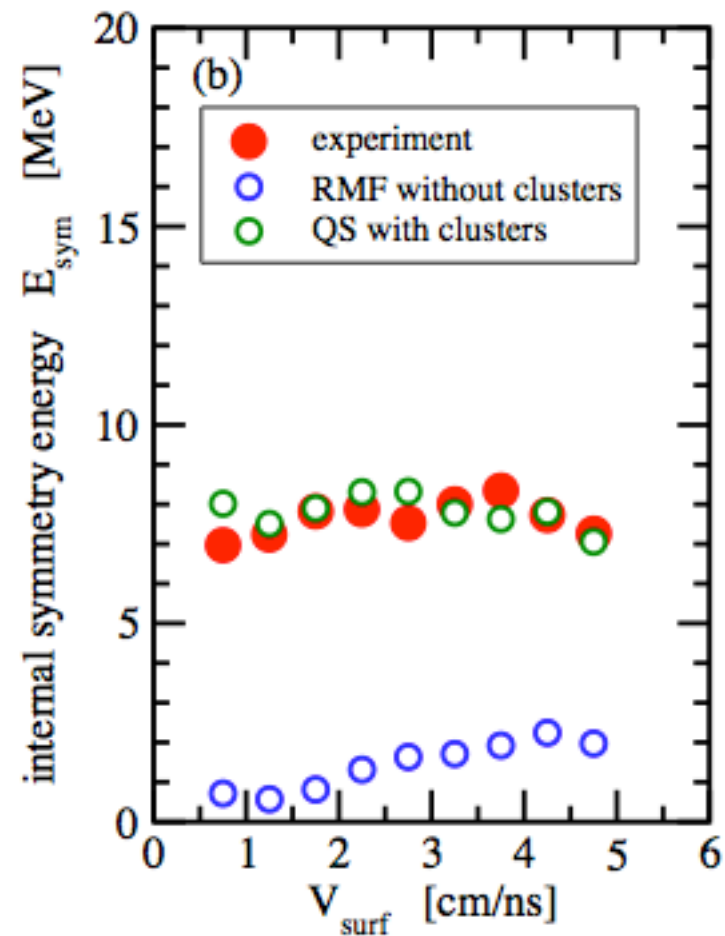
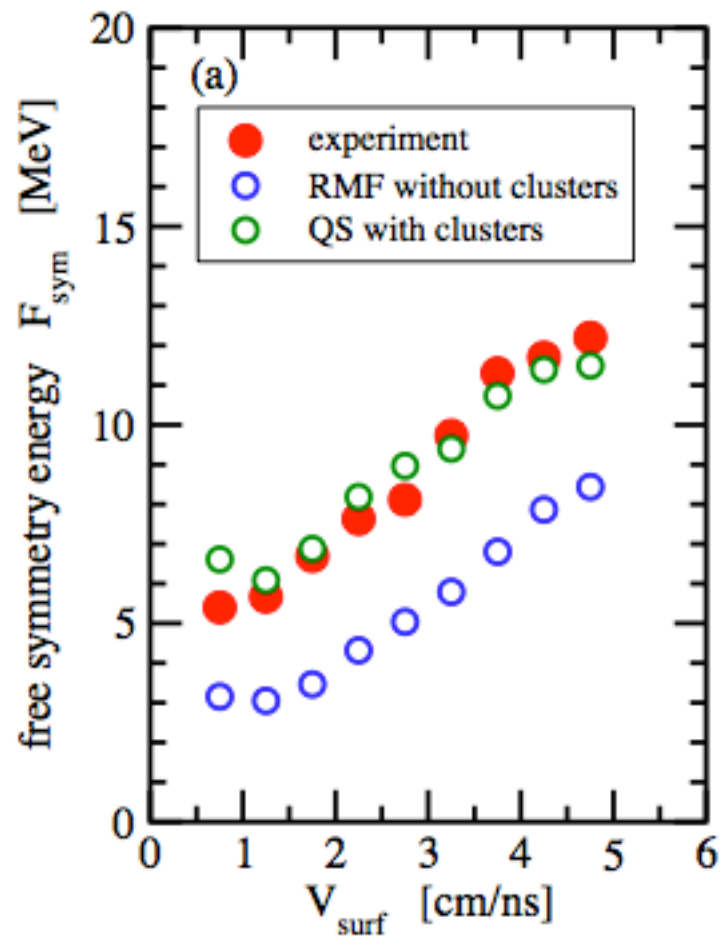
# Symmetry energy

Heavy-ion collisions, spectra of emitted clusters,  
temperature (3 - 10 MeV), free energy

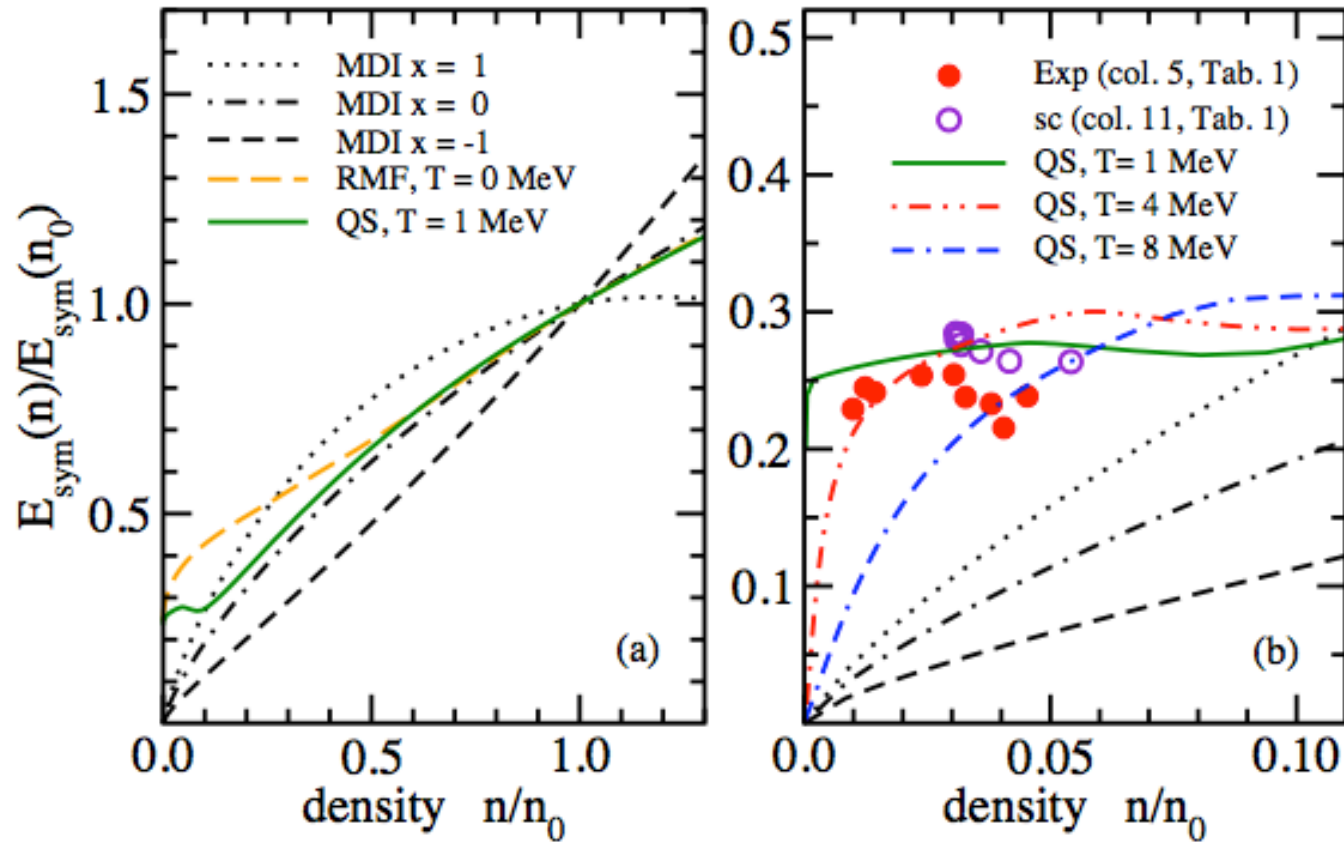


S. Kowalski et al.,  
PRC **75**, 014601  
(2007)

# Symmetry energy, comparison experiment with theories



# Symmetry Energy



Scaled internal symmetry energy as a function of the scaled total density.

MDI: Chen et al., QS: quantum statistical, Exp: experiment at TAMU

J.Natowitz et al. PRL, May 2010

# Problems:

1. Inclusion of larger Clusters in the equation of state
2. HFB and cluster formation/quartetting
3. EoS at low temperatures/low densities
4. Inclusion of larger nuclei
5. Influence of the symmetry energy on beta equilibrium
  - Transition from  $\alpha$ -matter to nucleon matter  
Is there a region of metastability?

# Problems:

1. Inclusion of larger Clusters in the equation of state
2. HFB and cluster formation/quartetting
3. EoS at low temperatures/low densities
4. Inclusion of larger nuclei
5. Influence of the symmetry energy  
on beta equilibrium:  
Diploma thesis work, Armen Sedrakian, 1987?

# Astrophysical Applications

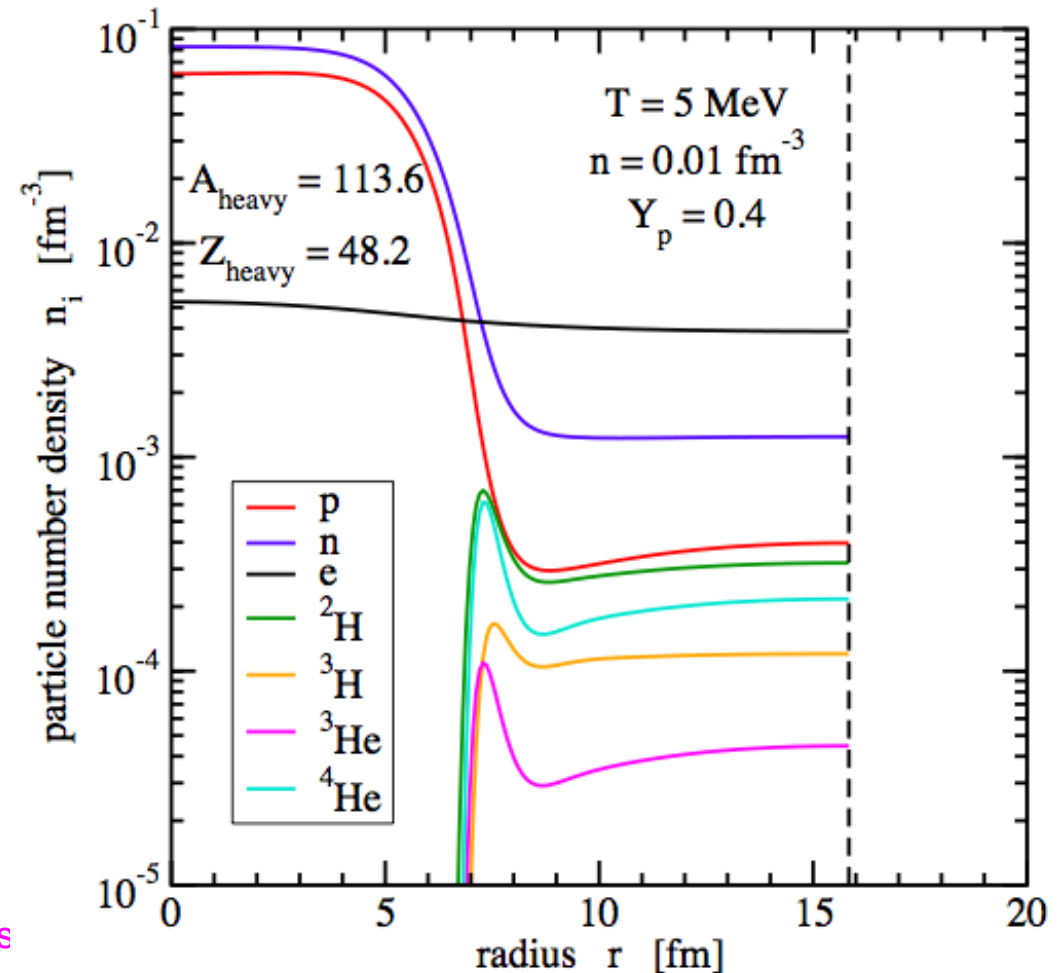
- Supernova explosions
- Neutrino transport
- Neutron star structure
- Equation of state (EOS)
- Composition
- Transport properties (cross sections)

# $\alpha$ cluster in astrophysics

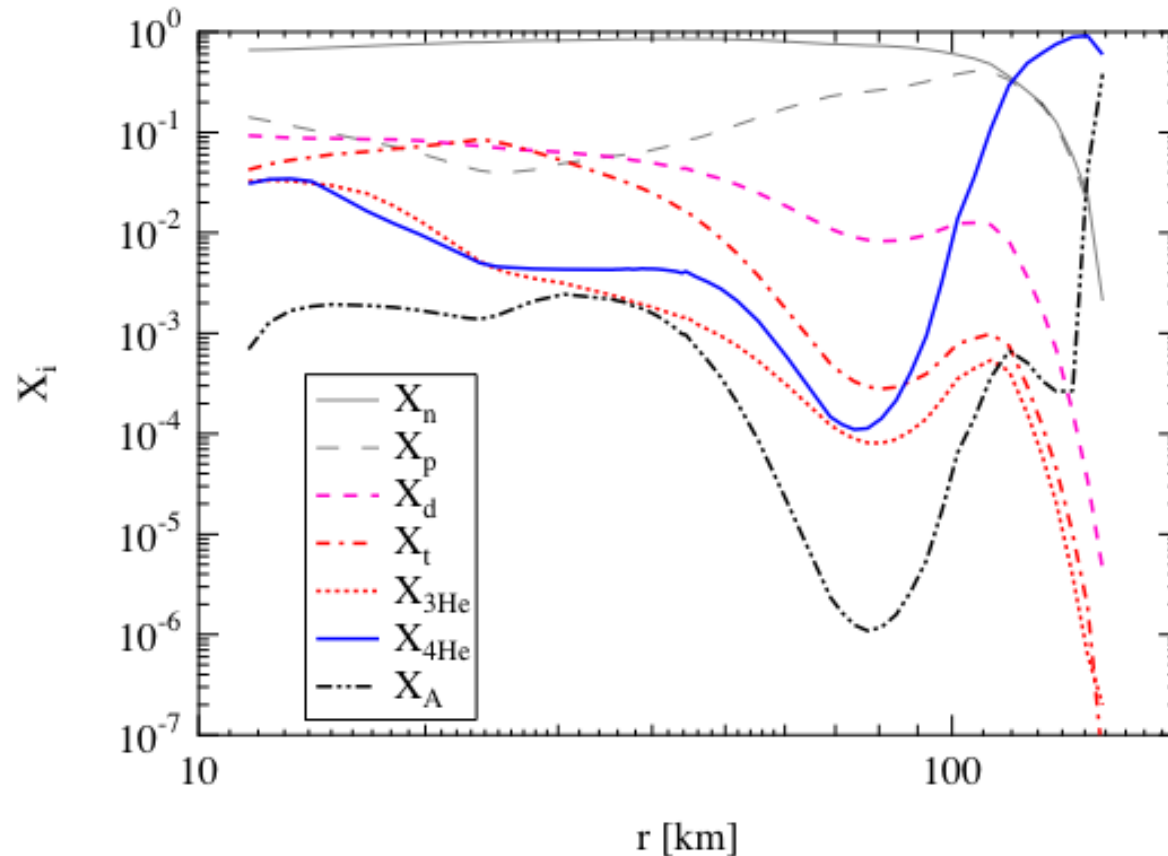
Crust of neutron stars

Protons in droplets  
(heavy nuclei)

$\alpha$ -cluster outside,  
at the surface,  
condensate?



# Composition of supernova core



Mass fraction  $X$   
of light clusters  
for a post-bounce  
supernova core

K. Sumiyoshi,  
G. R.,  
PRC 77,  
055804 (2008)

S. Heckel, P. P. Schneider and A. Sedrakian,  
Light nuclei in supernova envelopes: a quasiparticle gas model  
Phys. Rev. C **80**, 015805 (2009).



# In-medium modification of transport properties of dense matter

- D. Blaschke, G. Röpke, H. Schulz, A.D. Sedrakian,  
Nuclear in-medium effect on the thermal conductivity and viscosity of  
neutron star matter  
PL B 338, 111 (1994)
- D. Blaschke, G. Röpke, H. Schulz, A.D. Sedrakian, D. Voskresensky,  
Nuclear in-medium effects and neutrino emissivity of neutron stars.  
M. N. R. A. S. 273, 596 (1995)

# Summary

- Correlations (cluster formation, quantum condensates) are essential in low-density matter. They are suppressed with increasing density (Pauli blocking).
- The low-density limit of the nuclear matter EoS can be rigorously treated. The [Beth-Uhlenbeck virial expansion](#) is a benchmark. Larger nuclei and pasta structures must be treated in future works.
- An [extended quasiparticle approach](#) can be given for single nucleon states and nuclei. In a first approximation, [self-energy](#) and [Pauli blocking](#) is included. An interpolation between low and high densities is possible.
- Compared with the standard quasiparticle approach, significant changes arise in the low-density limit due to clustering. Examples are [Bose-Einstein condensation \(quartetting\)](#), and the behavior of the [symmetry energy](#).
- Finite Systems (nuclei): [Clusters](#) in low-density regions, quantum condensates. Possibly preformed clusters at surface.

# Thanks

to D. Blaschke, C. Fuchs, Y. Funaki, H. Horiuchi,  
J. Natowitz, T. Klaehn, Z. Ren, S. Shlomo, P. Schuck,  
A. Sedrakian, K. Sumiyoshi, A. Tohsaki, S. Typel,  
H. Wolter, Z. Xu, T. Yamada, B. Zhou  
for collaboration

to you

for attention

D.G.

# Supernova

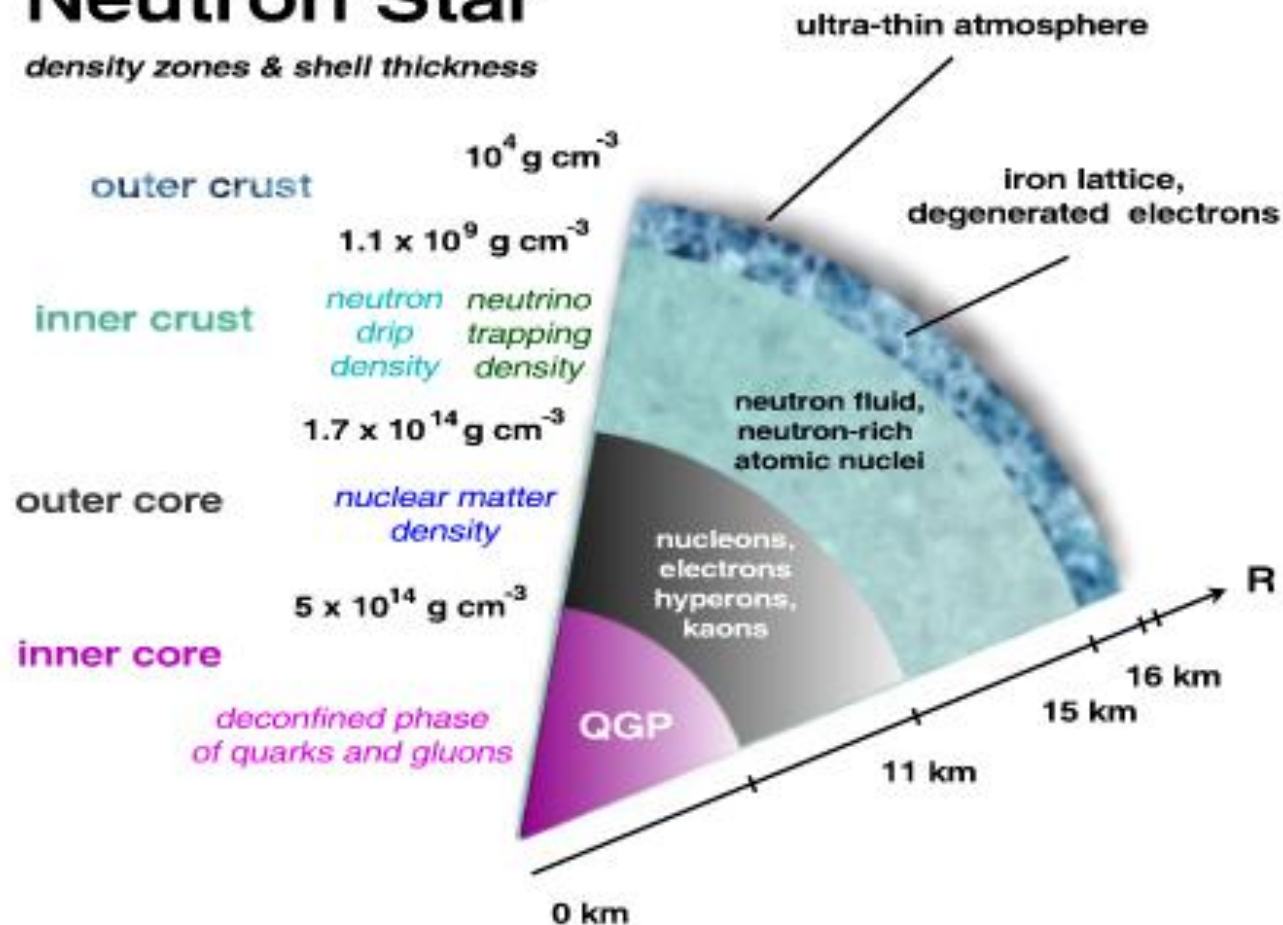
Crab nebula, 1054 China, PSR 0531+21



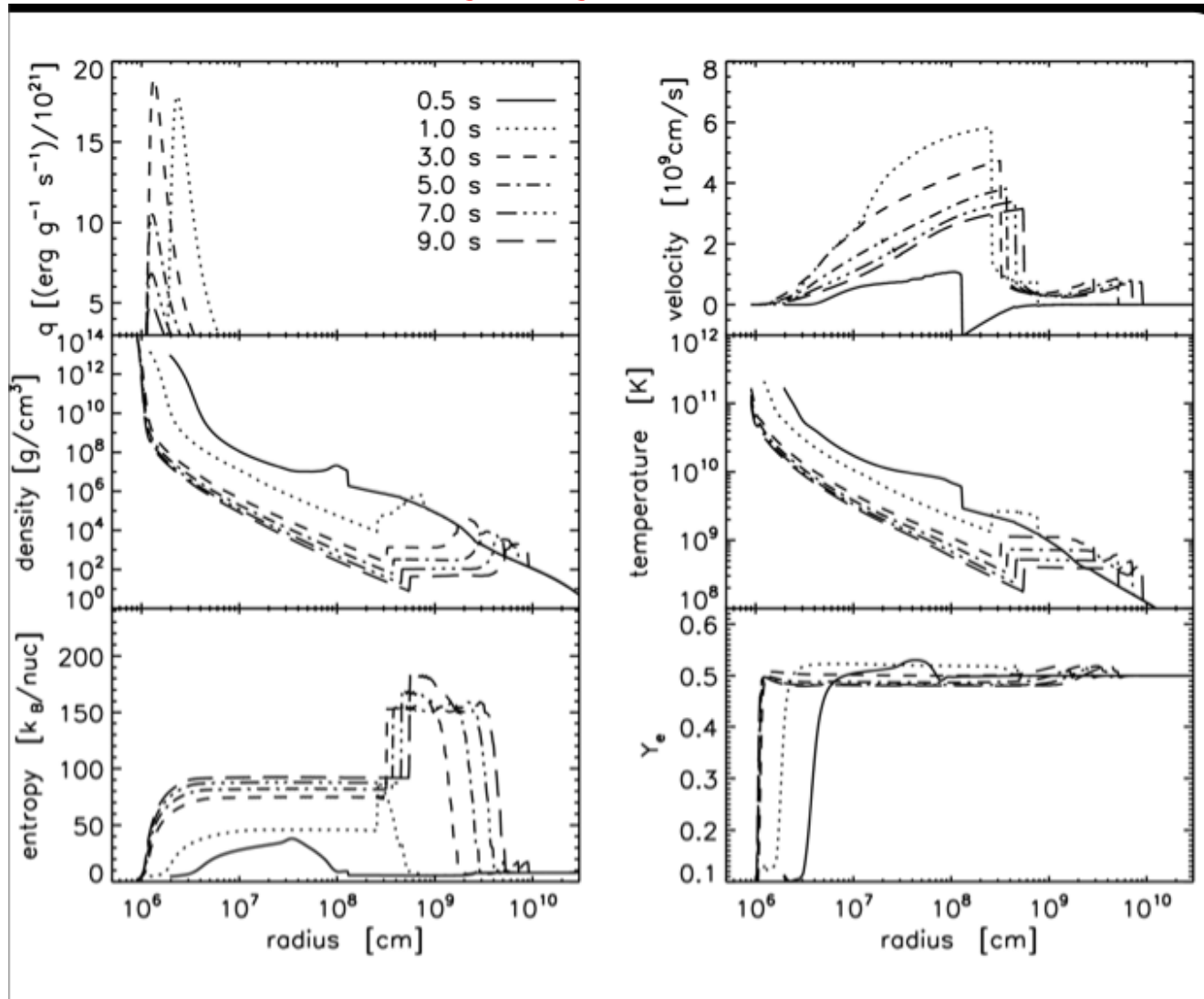
# Structure of a Neutron star

## Neutron Star

*density zones & shell thickness*



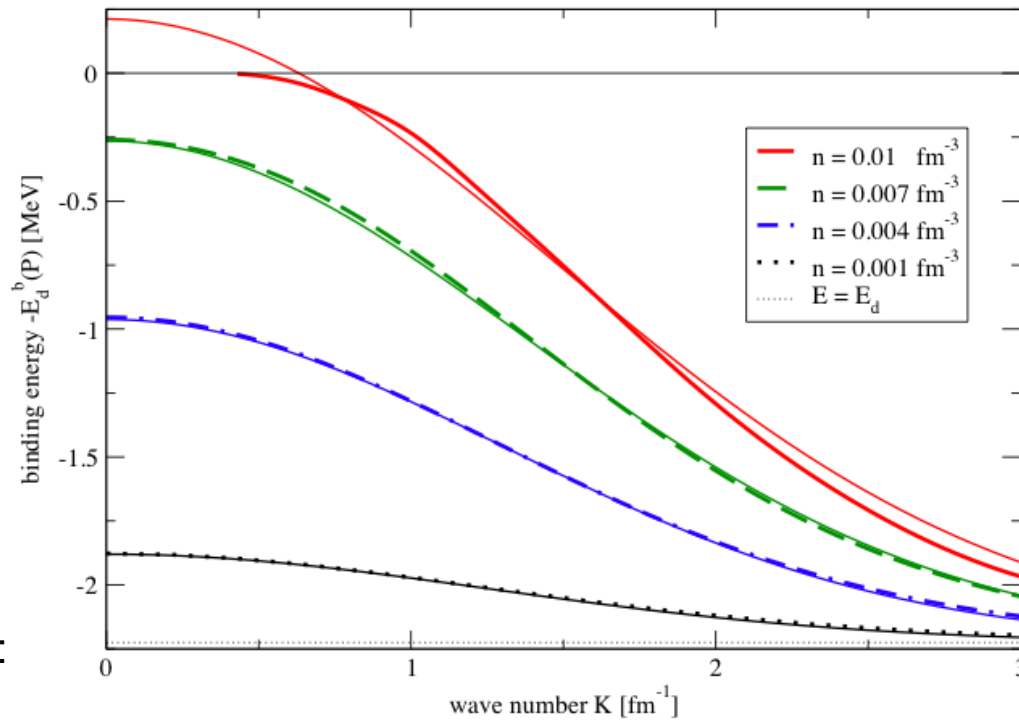
# Supernova collapse: spherically symmetric simulations



A. Arcones et al.  
Neutrino driven winds,  
PRC 78, 015806 (08)

# Shift of the deuteron binding energy

Dependence on center of mass momentum, various densities,  $T=10$  MeV



thin lines:

fit formula

# Deuteron quasiparticle properties

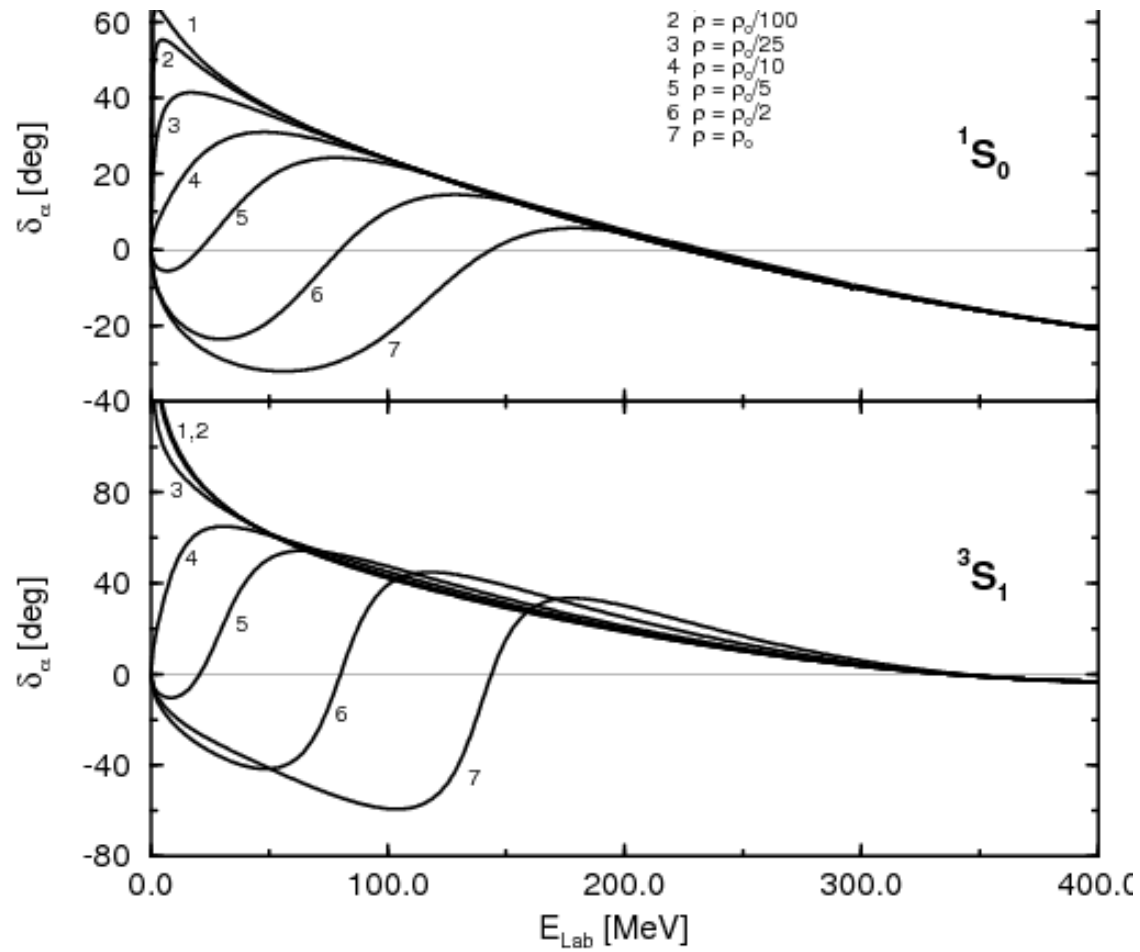
$$E_d^{\text{qu}}(P) = E_d^{\text{free}} + \Delta E_d + \frac{\hbar}{2m_d^*} P^2 + O(P^4)$$
$$E_d^{\text{free}} = -2.225\text{MeV}$$

$$\Delta E_d^{\text{Pauli}}(T, n_B, \alpha) = \delta E_d^{(0)}(T, \alpha) n_B + O(n_B^2)$$
$$\frac{m_d^*}{m_d}(T, n_B, \alpha) = 1 + \delta m_d^{(0)}(T, \alpha) n_B + O(n_B^2)$$

T [MeV]	delta E [MeV fm <sup>3</sup> ]	delta m <sup>*</sup> [fm <sup>3</sup> ]
10	364.3	21.3
4	712.9	87.1



# Scattering phase shifts in matter



# Application to Heavy Ion Reactions

- Test the EOS  
(NSE, virial,... at low densities,  
Skyrme, DBHF, RMF,... near saturation)
- Unifying quantum statistical approach, medium effects, Mott effect
- Symmetry energy
- Bose enhancement?

Nimrod @ TAMU,  
40Ar + 112,124Sn,  
64Zn + 112,124Sn; 47 A MeV

Open questions: freeze-out model or dynamical transport models?  
Identification of the source? - yields of p, (n), d, t, 3He, 4He,...

# Mott points from cluster yields

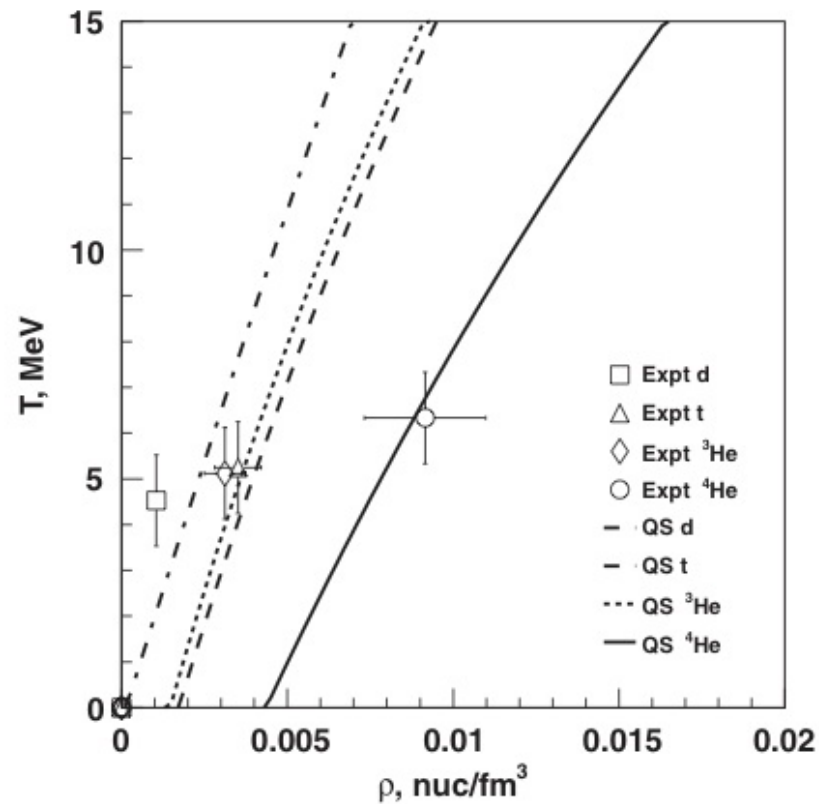
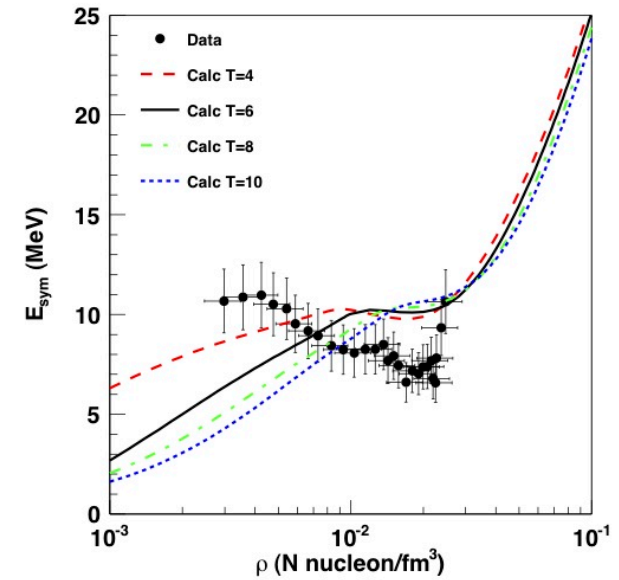
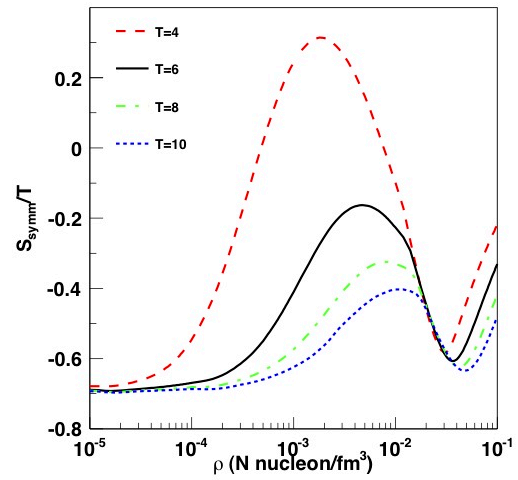
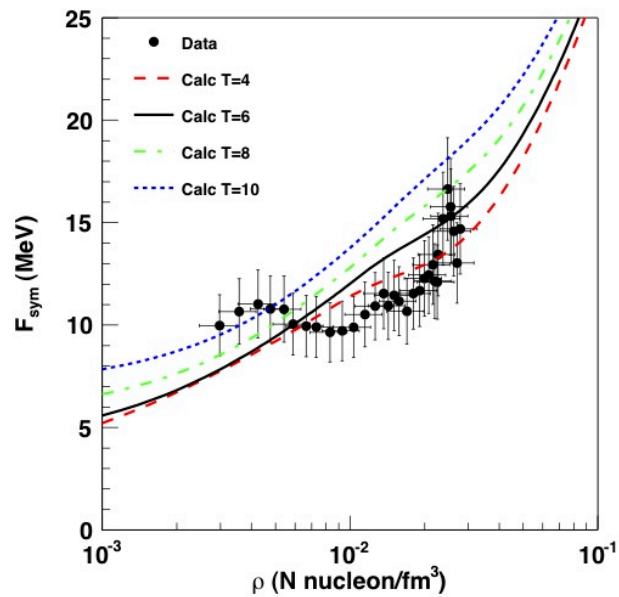


FIG. 3. Comparison of experimentally derived Mott point densities and temperatures with theoretical values. Symbols represent the experimental data. Estimated errors on the temperatures are 10% and on the densities 20%. Lines show polynomial fits to the Mott points presented in Ref. [1].

K. Hagel et al., PRL 108, 062702 (2012)

# Free symmetry energy



symmetry entropy

Internal symmetry energy

R. Wada et al., Phys. Rev. C 85, 064618 (2012).

# Internal energy per nucleon

Isotherms

$T$ [MeV]

20

18

16

14

12

10

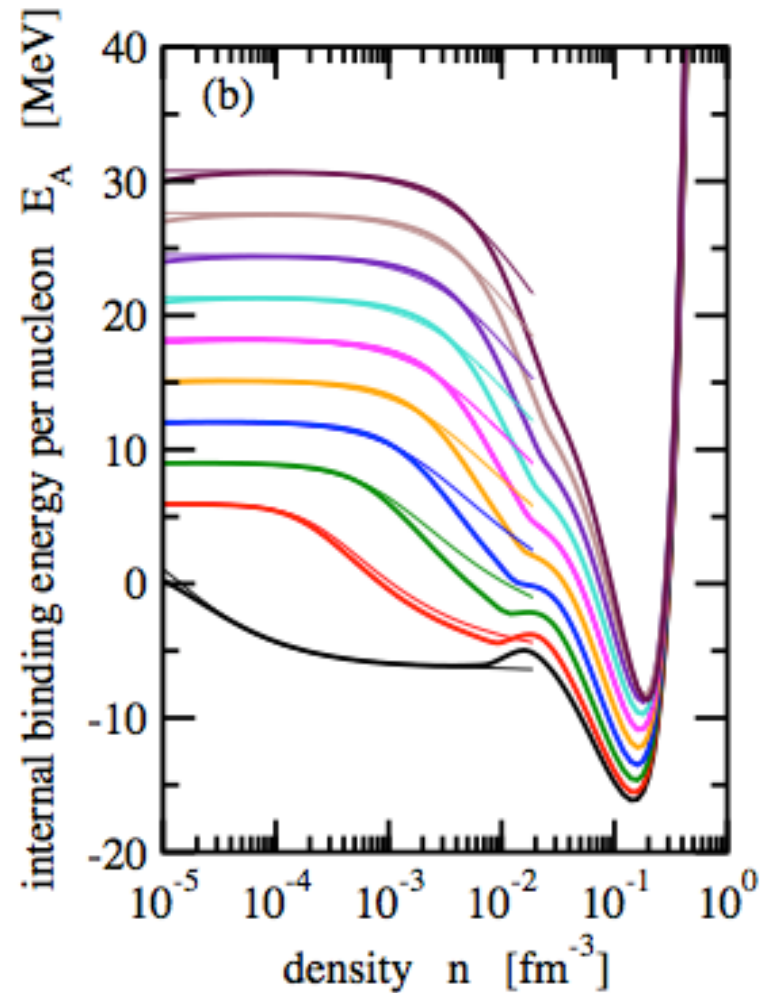
8

6

4

2

thin lines: NSE



# Composition of dense nuclear matter

$$n_p(T, \mu_p, \mu_n) = \frac{1}{V} \sum_{A,\nu,K} Z_A f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

$$n_n(T, \mu_p, \mu_n) = \frac{1}{V} \sum_{A,\nu,K} (A - Z_A) f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

mass number  $A$

charge  $Z_A$

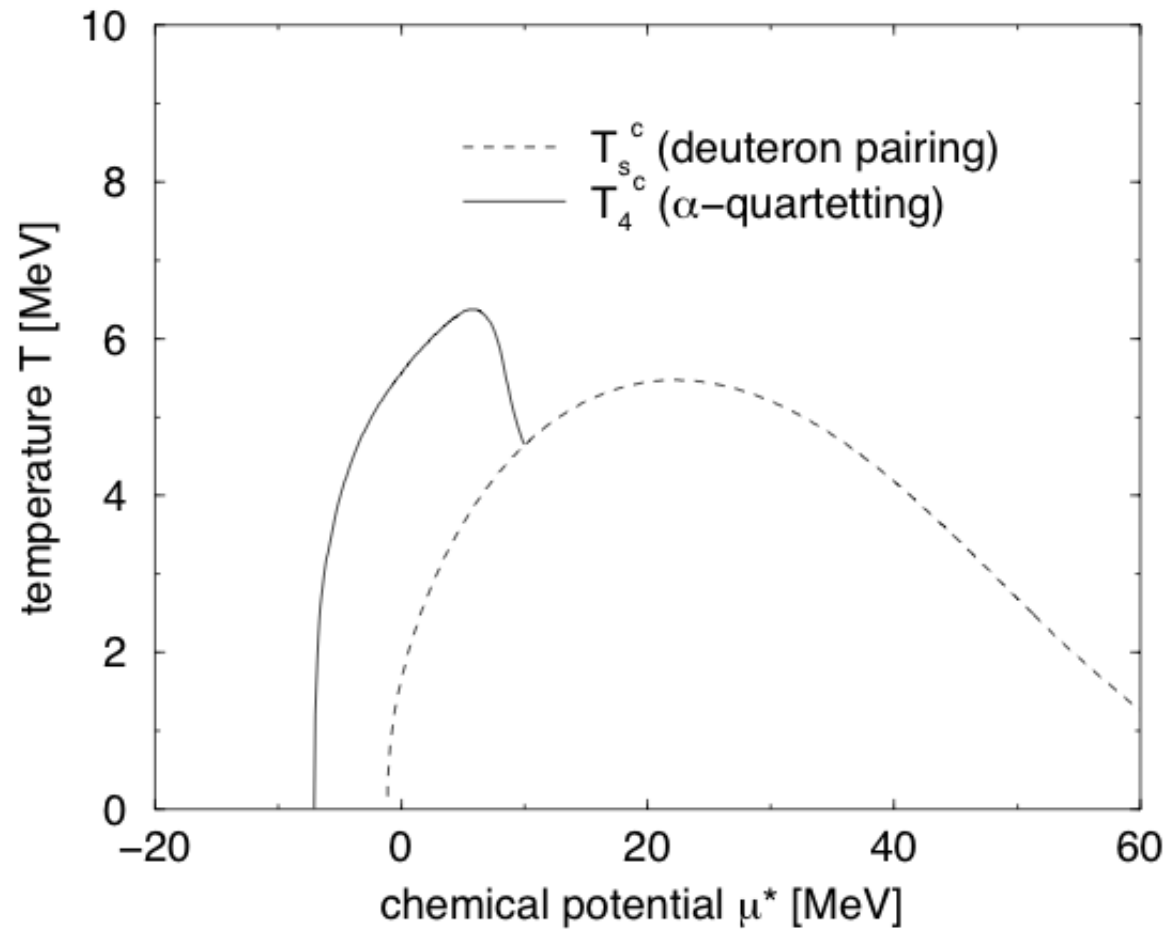
energy  $E_{A,\nu,K}$

$\nu$ : internal quantum number

$$f_A(z) = \frac{1}{\exp(z/T) - (-1)^A}$$

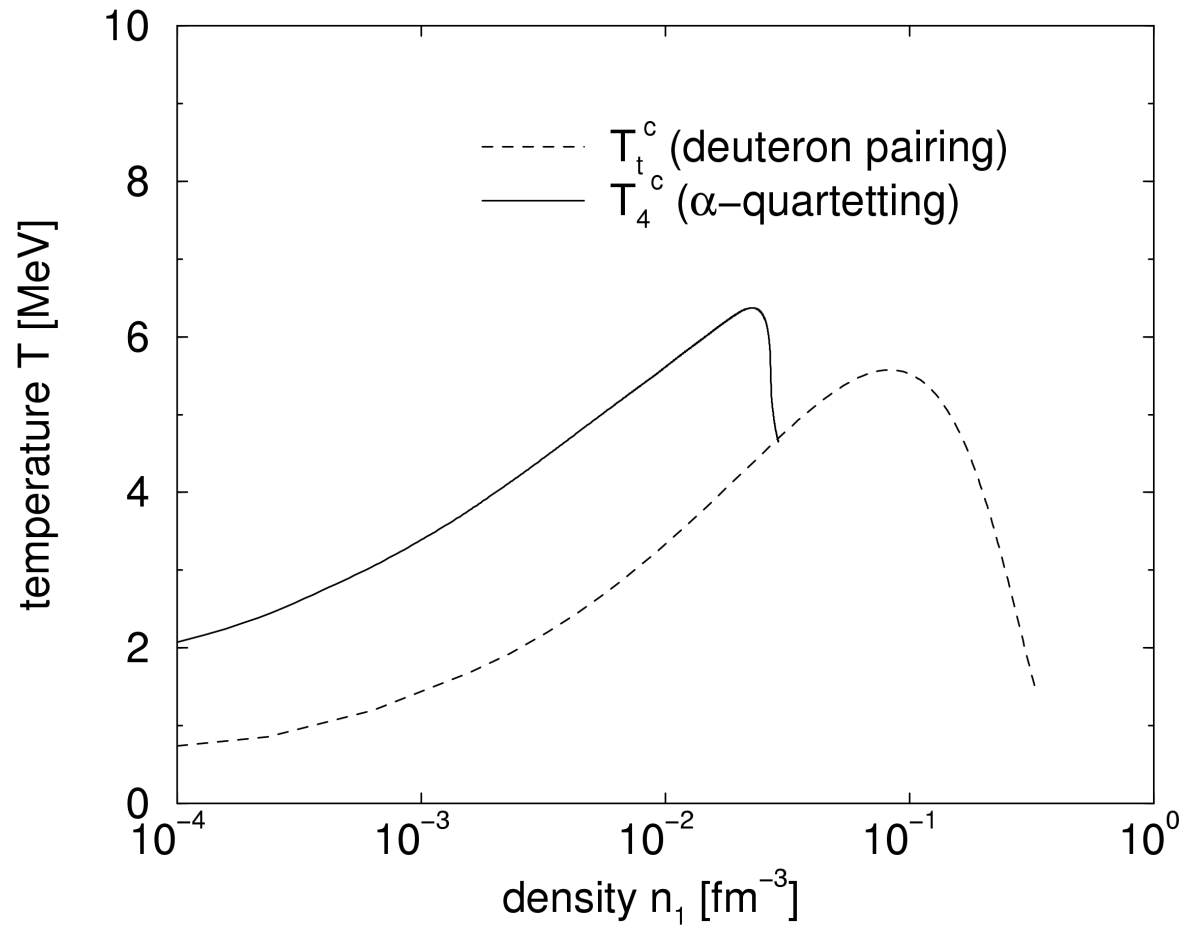
- Inclusion of excited states and continuum correlations
- Medium effects:  
**self-energy** and **Pauli blocking shifts** of binding energies,  
Coulomb corrections due to screening (Wigner-Seitz, Debye)
- Bose-Einstein condensation

# $\alpha$ -cluster-condensation (quartetting)



G.Röpke, A.Schnell, P.Schuck, and P.Nozieres, PRL 80, 3177 (1998)

# $\alpha$ -cluster-condensation (quartetting)



G.Röpke, A.Schnell, P.Schuck, and P.Nozieres, PRL 80, 3177 (1998)



# Effective wave equation for the deuteron in matter

In-medium two-particle wave equation in mean-field approximation

$$\left( \frac{p_1^2}{2m_1} + \Delta_1 + \frac{p_2^2}{2m_2} + \Delta_2 \right) \Psi_{d,P}(p_1, p_2) + \sum_{p_1', p_2'} (1 - f_{p_1} - f_{p_2}) V(p_1, p_2; p_1', p_2') \Psi_{d,P}(p_1', p_2')$$

Add self-energy

Pauli-blocking

$$= E_{d,P} \Psi_{d,P}(p_1, p_2)$$

Thouless criterion

$$E_d(T, \mu) = 2\mu$$

Fermi distribution function

$$f_p = \left[ e^{(p^2/2m - \mu)/k_B T} + 1 \right]^{-1}$$

BEC-BCS crossover:  
Alm et al., 1993



# Self-conjugate $4n$ nuclei

$^{12}\text{C}$ :

$0^+$  state at 0.39 MeV above the  $3\alpha$  threshold energy:  
 $\alpha$  cluster interact predominantly in relative  $S$  waves,  
gaslike structure

$\alpha$ -particle condensation in low-density nuclear matter  
( $\rho \leq \rho_0/5$ )

$n\alpha$  cluster condensed states

-- a general feature in  $N = Z$  nuclei?

# Self-conjugate 4n nuclei

$n\alpha$  nuclei:  ${}^8\text{Be}$ ,  ${}^{12}\text{C}$ ,  ${}^{16}\text{O}$ ,  ${}^{20}\text{Ne}$ ,  ${}^{24}\text{Mg}$ , ...

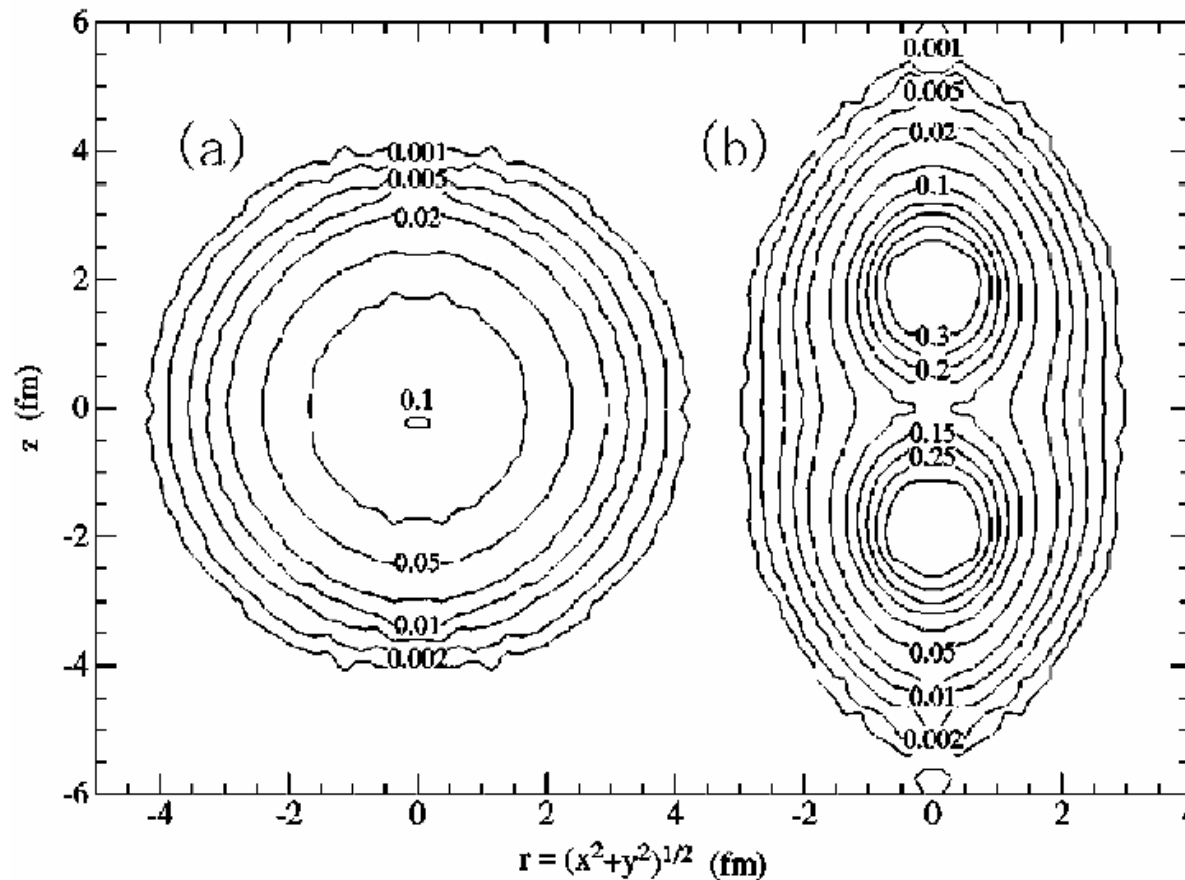
Single-particle shell model, or

Cluster type structures

ground state, excited states

$n\alpha$  break up at the threshold energy  $E_{n\alpha}^{\text{thr}} = nE_{\alpha}$

# Alpha cluster structure of Be 8



R.B. Wiringa et al.,  
PRC 63, 034605 (01)

Contours of constant density, plotted in cylindrical coordinates, for  $^8\text{Be}(0^+)$ .  
The left side is in the laboratory frame while the right side is in the intrinsic frame.

# Results

		$E_k$ (MeV)	$E_{exp}$ (MeV)	$E_k - E_{n\alpha}^{thr}$ (MeV)	$(E - E_{n\alpha}^{thr})_{exp}$ (MeV)	$\sqrt{\langle r^2 \rangle}$ (fm)	$\sqrt{\langle r^2 \rangle}_{exp}$ (fm)
$^{12}\text{C}$	$k = 1$	-85.9	-92.16 ( $0_1^+$ )	-3.4	-7.27	2.97	2.65
	$k = 2$	-82.0	-84.51 ( $0_2^+$ )	+0.5	0.38	4.29	
	$E_{3\alpha}^{thr}$	-82.5	-84.89				
$^{16}\text{O}$	$k = 1$	-124.8 (-128.0)*	-127.62 ( $0_1^+$ )	-14.8 (-18.0)*	-14.44	2.59	2.73
	$k = 2$	-116.0	-116.36 ( $0_3^+$ )	-6.0	-3.18	3.16	
	$k = 3$	-110.7	-113.62 ( $0_5^+$ )	-0.7	-0.44	3.97	
	$E_{4\alpha}^{thr}$	-110.0	-113.18				
					-0.17	+0.1	

Tabelle 1: Comparison of the generator coordinate method calculations with experimental values.  $E_{n\alpha}^{thr} = nE_\alpha$  denotes the threshold energy for the decay into  $\alpha$ -clusters, the values marked by \* correspond to a refined mesh.

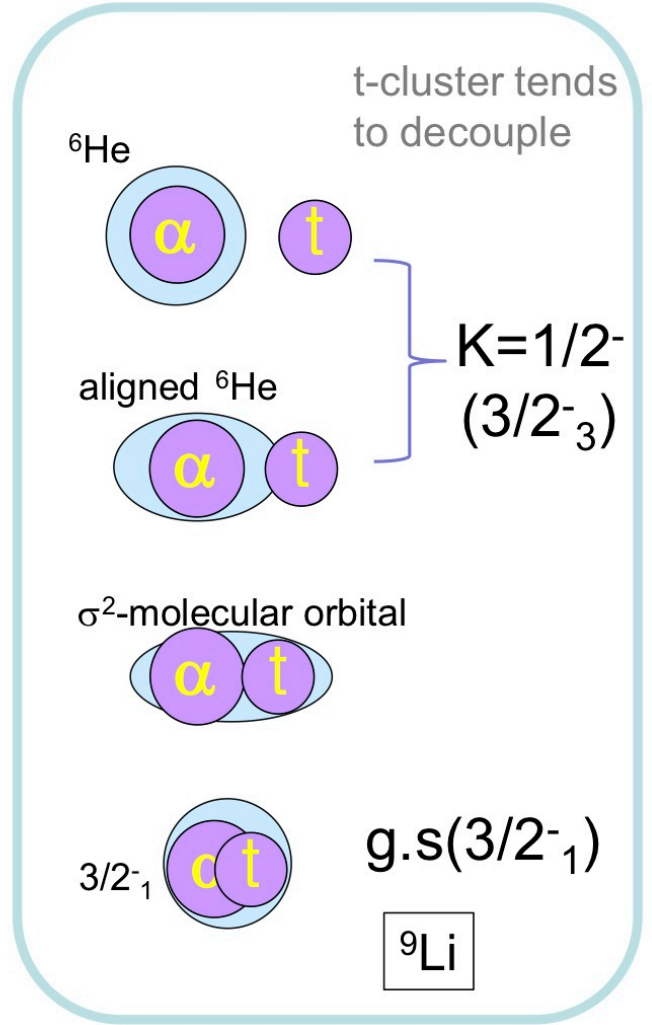
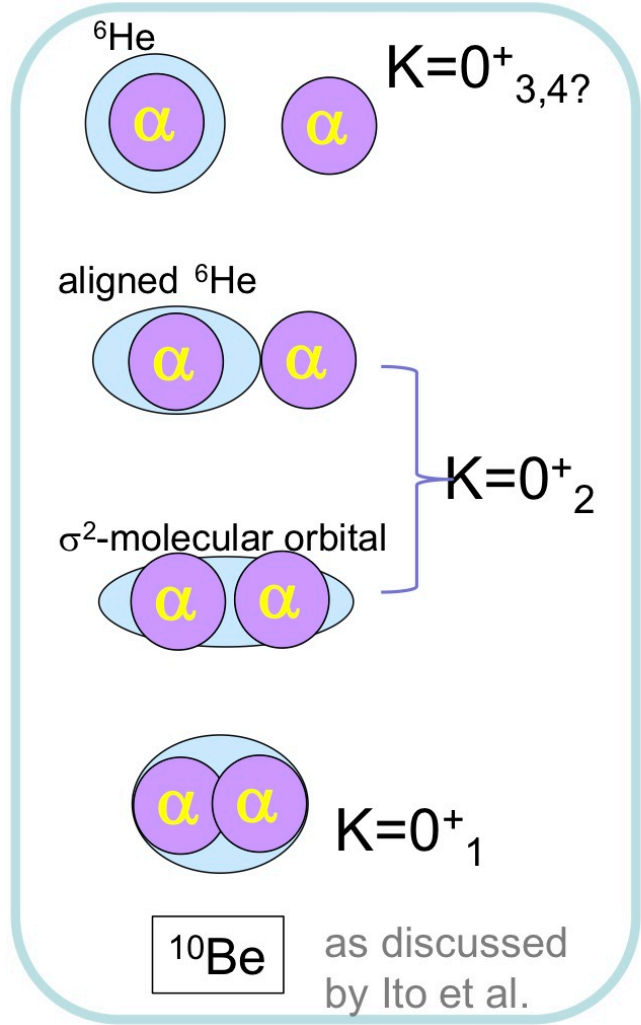
# Excited light nuclei

## Cluster structures in $^{10}\text{Be}$ and $^9\text{Li}$

Yoshiko Kanada-En'yo  
Cluster2012, Debrecen

decreasing density  
deuterons?  
systematics in weakly bound light elements  
light clusters in neutron matter

Inter-cluster distance (d) ↑

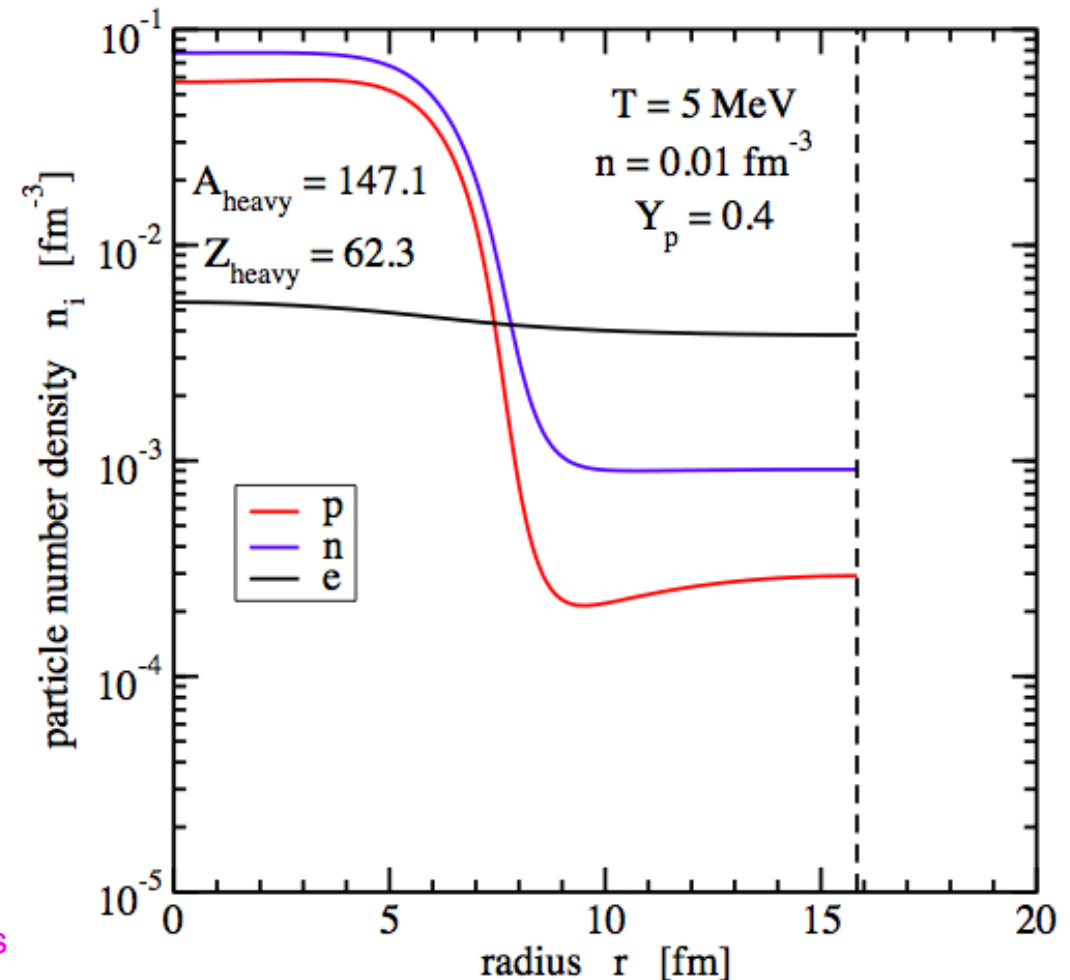


# $\alpha$ cluster in astrophysics

Crust of neutron stars

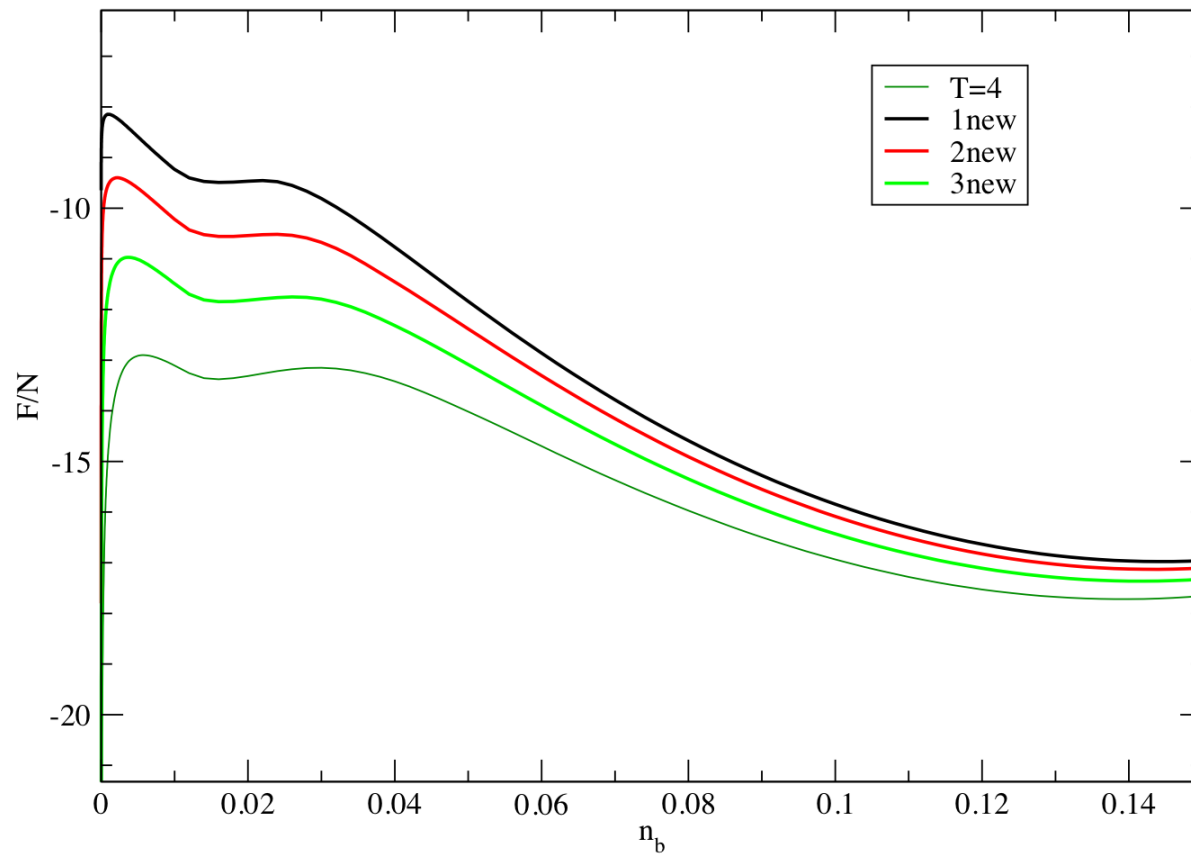
Protons in droplets  
(heavy nuclei)

$\alpha$ -cluster outside,  
at the surface,  
condensate?





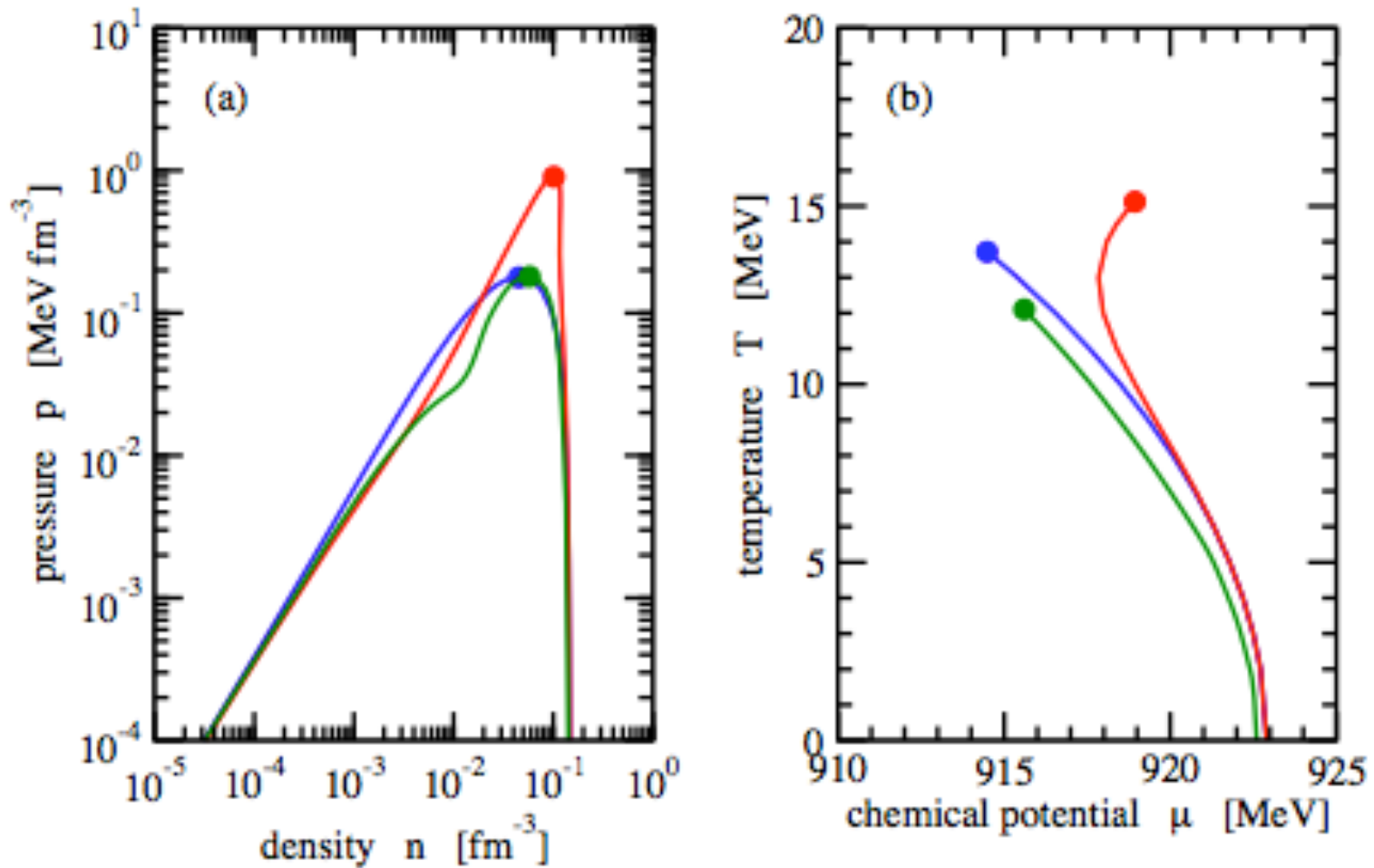
# Free energy per nucleon



preliminary

correlated  
medium

# Liquid-vapor phase transition

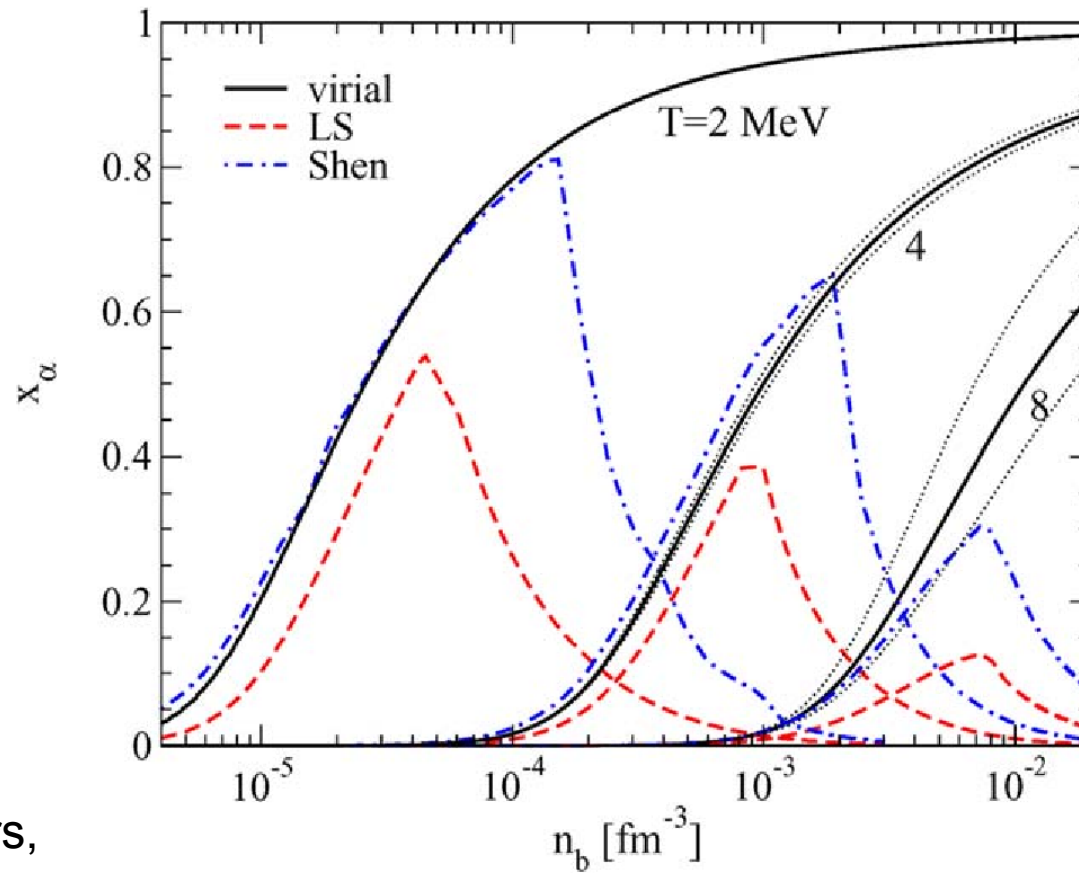


blue: no light cluster, green: with light clusters, QS, red: cluster-RMF

S. Typel et al., PRC 81, 015803 (2010)

# Alpha-particle fraction in the low-density limit

symmetric matter,  $T=2, 4, 8$  MeV



LS, Shen:  
higher clusters,  
excluded volume

C.J.Horowitz, A.Schwenk, Nucl. Phys. A 776, 55 (2006)



Available online at [www.sciencedirect.com](http://www.sciencedirect.com)

SciVerse ScienceDirect

Nuclear Physics A 897 (2013) 70–92

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PHYSICS A

[www.elsevier.com/locate/nuclphysa](http://www.elsevier.com/locate/nuclphysa)

## Cluster-virial expansion for nuclear matter within a quasiparticle statistical approach

G. Röpke<sup>a,\*</sup>, N.-U. Bastian<sup>a</sup>, D. Blaschke<sup>b,c</sup>, T. Klähn<sup>b</sup>, S. Typel<sup>d</sup>,  
H.H. Wolter<sup>e</sup>

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<sup>b</sup> *Instytut Fizyki Teoretycznej, Uniwersytet Wrocławski, pl. M. Borna 9, 50-204 Wrocław, Poland*

<sup>c</sup> *Bogoliubov Laboratory for Theoretical Physics, JINR Dubna, Joliot-Curie str. 6, 141980 Dubna, Russia*

<sup>d</sup> *GSI Helmholtzzentrum für Schwerionenforschung GmbH, Theorie, Planckstraße 1, D-64291 Darmstadt, Germany*

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### Abstract

Correlations in interacting many-particle systems can lead to the formation of clusters, in particular bound states and resonances. Systematic quantum statistical approaches allow to combine the nuclear statistical equilibrium description (law of mass action) with mean-field concepts. A chemical picture, which treats the clusters as distinct entities, serves as an intuitive concept to treat the low-density limit. Within a generalized Beth–Uhlenbeck approach, the quasiparticle-virial expansion is extended to include arbitrary clusters, where special attention must be paid to avoid inconsistencies such as double counting. Correlations are suppressed with increasing density due to Pauli blocking. The contribution of the continuum to the virial coefficients can be reduced by considering clusters explicitly and introducing quasiparticle energies. The cluster-virial expansion for nuclear matter joins known benchmarks at low densities with those near saturation density.

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# Cluster - mean field approximation

Cluster (A) interacting with a distribution of clusters (B) in the medium,  
fully antisymmetrized

$$\sum_{1' \dots A'} \{ H_A^0(1 \dots A, 1' \dots A') + \sum_i \Delta_i^{A,mf} \delta_{k,k'} + \frac{1}{2} \sum_{i,j} \Delta V_{ij}^{A,mf} \delta_{l,l'} - E_{AvP} \delta_{k,k'} \} \psi_{AvP}(1' \dots A') = 0$$

self-energy

$$\Delta_1^{A,mf}(1) = \sum_2 V(12,12)_{ex} f^*(2) + \sum_{BvP} \sum_{2 \dots B'} f_B(E_{BvP}) \sum_i V_{1i}(1i,1'i') \psi_{BvP}^*(1 \dots B) \psi_{BvP}(1' \dots B')$$

effective interaction

$$\Delta V_{12}^{A,mf} = -\frac{1}{2} [f^*(1) + f^*(2)] V(12,1'2') - \sum_{BvP} \sum_{2^* \dots B''} f_B(E_{BvP}) \sum_i V_{1i} \psi_{BvP}^*(22^* \dots B^*) \psi_{BvP}(2'2'' \dots B'')$$

phase space occupation  $f^*(1) = f_1(1) + \sum_{BvP} \sum_{2 \dots B} f_B(E_{BvP}) |\psi_{BvP}(1 \dots B)|^2$

# Self-consistent RPA

Two-time cluster Matsubara Green's functions

$$\begin{aligned} G_{\alpha\beta}^{\tau-\tau'} &= -\langle Tr A_{\alpha}(\tau) A_{\beta}^{\dagger}(\tau') \rangle \\ &= -Tr \left[ \rho_G T_{\tau} e^{\tau K} A_{\alpha} e^{-(\tau-\tau')K} A_{\beta}^{\dagger} e^{-\tau'K} \right] \end{aligned}$$

Equation of motion method

$$\begin{aligned} -\frac{\partial}{\partial \tau} G_{\alpha\beta}^{\tau-\tau'} &= \delta_{\tau-\tau'} \langle [A_{\alpha}, A_{\beta}^{\dagger}] \rangle - \langle Tr [A_{\alpha}, K]^{\tau} A_{\beta}^{\dagger}(\tau') \rangle \\ &= \delta_{\tau-\tau'} N_{\alpha\beta} + \sum_{\gamma} \int d\tau_1' \mathcal{H}_{\alpha\gamma}^{\tau-\tau_1'} G_{\gamma\beta}^{\tau_1'-\tau'} . \end{aligned}$$

Effective Hamiltonian is split into an instantaneous and a dynamic part

$$\begin{aligned} \mathcal{H}_{\alpha\beta}^{\tau-\tau'} &= \sum_{\beta'} \left\{ \delta_{\tau-\tau'} \langle [[A_{\alpha}, K], A_{\beta'}^{\dagger}] \rangle - \langle Tr [A_{\alpha}, K]^{\tau} [K, A_{\beta'}^{\dagger}]^{\tau'} \rangle_{irr} \right\} N_{\beta'\beta}^{-1} \\ &\equiv \mathcal{H}_{\alpha\beta}^{(0)} \delta_{\tau-\tau'} + \mathcal{H}_{\alpha\beta}^{(r)\tau-\tau'} . \end{aligned}$$

J.Dukelsky, G. Roepke, and P.Schuck, NPA 628, 17 (1998)  
P. Schuck, D.S. Delion, J.Dukelsky, and G. Roepke, in preparation

# Deuteron quasiparticle properties

$$E_d^{\text{qu}}(P) = E_d^{\text{free}} + \Delta E_d + \frac{\hbar^2}{2m_d^*} P^2 + O(P^4)$$

$$\Delta E_d^{\text{Pauli}}(T, n_B, \alpha) = \delta E_d^{(0)}(T, \alpha) n_B + O(n_B^2)$$

$$\frac{m_d^*}{m_d}(T, n_B, \alpha) = 1 + \delta m_d^{(0)}(T, \alpha) n_B + O(n_B^2)$$

T [MeV]	delta E [MeV fm <sup>3</sup> ]	delta m* [fm <sup>3</sup> ]
10	364.3	21.3
4	712.9	87.1

$$E_d^{\text{free}} = -2.225\text{MeV}$$

G.R., PRC 79, 014002 (2009)

# Nuclear matter properties

binding energy per nucleon near saturation:

$$\frac{E}{A}(n, \beta) = \frac{\varepsilon}{n} = a_V + \frac{K}{18}x^2 + \frac{K'}{162}x^3 + \beta^2 \left( J + \frac{L}{3}x + \dots \right) + \dots$$

with  $x = (n - n_{\text{sat}})/n_{\text{sat}}$ , asymmetry  $\beta = 1 - 2Y_p$  and

## nuclear matter parameters

- $n_{\text{sat}}$  saturation density
- $a_V$  bulk energy
- $K$  incompressibility
- $K'$  skewness
- $J$  bulk symmetry energy
- $L$  symmetry energy slope

S. Typel, 2012

