

Time-dependent Ginzburg-Landau equations for rotating two-flavor color superconductors

K. M. Shahabasyan and M. K. Shahabasyan

Physics Department, Yerevan State University, Yerevan, Armenia

A. D. Sedrakian

Institute for Theoretical Physics, J. W. Goethe University,

D-60438 Frankfurt am Main, Germany

TIME-INDEPENDENT GL EQUATIONS FOR THE 2SC PHASE

$$\mathcal{F} = \mathcal{F}_n + \int \left[2\alpha dd^* + \beta (dd^*)^2 + 2\gamma \left| \left(\hbar \nabla + i \frac{q}{c} \mathbf{A}_x \right) d \right|^2 + \frac{1}{8\pi} (\text{rot } \mathbf{A}_x)^2 + \frac{1}{8\pi} (\text{rot } \mathbf{A}_y)^2 \right] dV, \quad (1)$$

$$\mathbf{A}_x = -\sin \theta_M \mathbf{A} + \cos \theta_M \mathbf{A}_8,$$

$$\mathbf{A}_y = \cos \theta_M \mathbf{A} + \sin \theta_M \mathbf{A}_8,$$

$$\cos \theta_M = \frac{\sqrt{3}g}{\sqrt{e^2 + 3g^2}}, \quad q = \sqrt{e^2 + 3g^2}/3$$

$$\frac{\delta \mathcal{F}}{\delta d^*} = 0, \quad \frac{\delta \mathcal{F}}{\delta \mathbf{A}_x} = \frac{\delta \mathcal{F}}{\delta \mathbf{A}_y} = 0.$$

$$\alpha d + \beta |d|^2 d - \gamma \left(\hbar \nabla + i \frac{q}{c} \mathbf{A}_x \right)^2 d = 0.$$

$$\mathbf{j}_x = 2i\hbar\gamma q (d^* \nabla d - d \nabla d^*) - 4\gamma \frac{q^2}{c} |d|^2 \mathbf{A}_x,$$

$$\mathbf{j}_y = 0,$$

$$\alpha = \frac{(T - T_c)}{T_c} \nu, \quad \beta = \frac{7\zeta(3)}{8(\pi k_B T_c)^2} \nu, \quad \gamma = \frac{p_F^2}{6\mu^2} \beta = \frac{7\zeta(3)}{16(\pi k_B T_c)^2} \frac{n_b}{\mu},$$

$$d_0 = \left(-\frac{\alpha}{\beta} \right)^{1/2} = \left[\frac{8}{7\zeta(3)} \left(\frac{T_c - T}{T_c} \right) \right]^{1/2} \pi k_B T_c.$$

$$\mathbf{A}_x \rightarrow \mathbf{A}'_x \equiv \mathbf{A}_x + \frac{2\mu}{3q} [\boldsymbol{\Omega} \times \mathbf{r}] \quad \mathbf{B}_L = -\frac{4\mu}{3q} \boldsymbol{\Omega}.$$

$$B_L \ll H_{c1}$$

TIME-DEPENDENT GL EQUATIONS FOR THE 2SC PHASE

$$-\Gamma \left(\frac{\partial}{\partial t} - \frac{iq}{\hbar} \phi \right) d = \alpha d + \beta |d|^2 d - \gamma_d \left(\hbar \nabla + \frac{iq}{c} \mathbf{A}_x \right)^2 d.$$

$$\gamma_d = \frac{\pi \nu D}{8 \hbar k_B T_c}, \quad \nu = \mu p_F / \pi^2 \hbar^3$$

is the density of states at the Fermi surface

$D = v_F^2 \tau_q / 3$, is the diffusion coefficient,

$\tau k_B T_c \gg \hbar$ weak pair breaking $\tau \gg 4.3 \cdot 10^{-19} s$.

$$\Gamma = \frac{\pi \hbar \nu}{8k_B T_c}$$

$$\tau_d = \frac{\Gamma}{|\alpha|} = \frac{\pi \hbar}{8k_B T_c} \left(1 - \frac{T}{T_c}\right)^{-1}$$

$$\mathbf{j}_d = 2i\hbar\gamma_d q [d^* \nabla d - d \nabla d^*] - \frac{4\gamma_d q^2}{c} |d|^2 \mathbf{A}_x - \sigma_q \left(\frac{1}{c} \frac{\partial \mathbf{A}}{\partial t} + \nabla \phi \right),$$

$$\tau_j = \frac{\sigma_q}{4\gamma_d q^2 |d|^2}$$

$$\phi_8 \rightarrow \phi'_8 = \phi_8 - \frac{\mu}{3q} \Omega^2 r^2$$

$$\mathbf{A}_x \rightarrow \mathbf{A}'_x = \mathbf{A}_x + \frac{2}{3} \frac{\mu c}{q} [\boldsymbol{\Omega} \times \mathbf{r}]$$

$$j_d = 2i\gamma_d \hbar q [d^* \nabla d - d \nabla d^*] - \frac{4q^2 \gamma_d}{c} |d|^2 \mathbf{A}'_x - \sigma_q \mathbf{E}'$$

$$\mathbf{E}' \equiv \frac{1}{c} \frac{\partial \mathbf{A}_x}{\partial t} + \nabla \phi_8 - \frac{2\mu}{3q} \Omega^2 \mathbf{r} \quad \phi_r = -\frac{\mu}{3q} \Omega^2 r^2$$

$$W = 2\Gamma \left| \frac{\partial d}{\partial t} - \frac{iq}{\hbar} \phi'_8 d \right|^2 + \sigma_q \mathbf{E}'^2$$

$$d = |d| e^{-i\chi}$$

$$W = 2\Gamma \left[\left(\frac{\partial |d|}{\partial t} \right)^2 + |d|^2 \left(\frac{\partial \chi}{\partial t} + \frac{q}{\hbar} \phi'_8 \right)^2 \right] + \sigma_q \mathbf{E}'^2,$$

$$\frac{2iq}{\hbar} \left(d \frac{\delta F}{\delta d} - d^* \frac{\delta F}{\delta d^*} \right) = \operatorname{div} \mathbf{j}_s,$$

$$-\operatorname{div} \mathbf{j}_s = \operatorname{div} \mathbf{j}_n = \sigma_q \operatorname{div} \mathbf{E}',$$

$$\lambda_E^2 \nabla^2 \phi'_8 - \frac{|d|^2}{d_0^2} \phi'_8 = 0,$$

$$\lambda_E = \left(\frac{\sigma_q \hbar^2}{4\Gamma' q^2 d_0^2} \right)^{1/2}, \quad \partial \chi / \partial t = 0 = \partial \mathbf{A} / \partial t.$$