

Unsafe but Calculable

Ratio Observables in Perturbative QCD

Jesse Thaler



Based on 1307.1699 with Andrew Larkoski

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Outline

Ratio Observables: Ubiquitous & Unsafe

Ratios of Angularities: “Sudakov Safe”

Computational Results

Ratio Observables: Ubiquitous & Unsafe

For jet substructure, often dimensionless ratios

Planar Flow

$$\text{Pf} = \frac{4 \det S_{\perp}}{(\text{tr} S_{\perp})^2}$$

N-subjettiness

$$\tau_{N,N-1} = \frac{\tau_N}{\tau_{N-1}}$$

Energy Correlation Functions

$$C_N = \frac{\text{ECF}(N+1) \text{ECF}(N-1)}{\text{ECF}(N)^2}$$

Motivated by quasi-scale-invariant nature of QCD

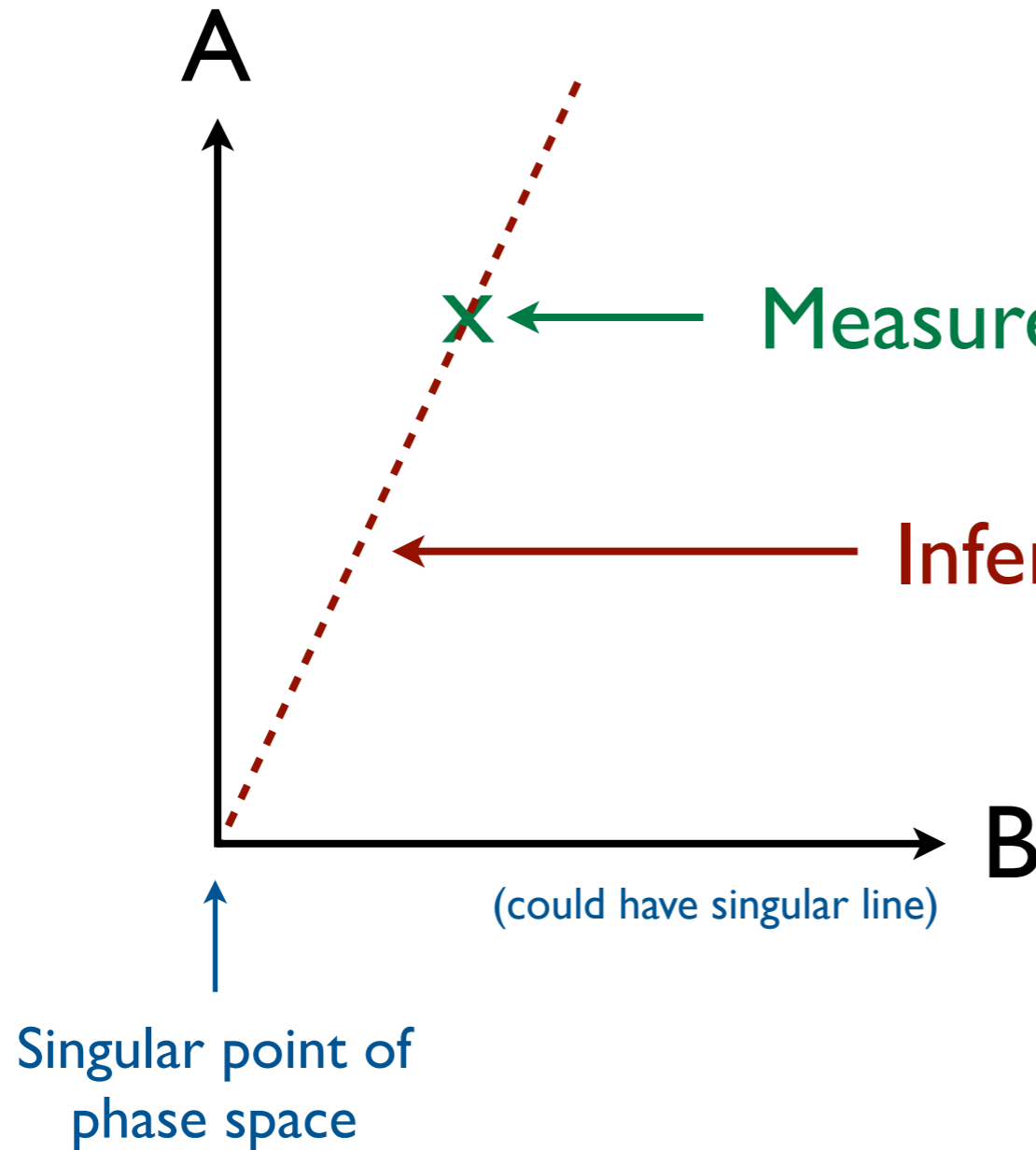
Surprising/Frustrating Fact:

$$\text{IRC Unsafe Ratio} = \frac{\text{IRC Safe Numerator}}{\text{IRC Safe Denominator}}$$

[observed by Soyez, Salam, Kim, Dutta, Cacciari]

WHAT?!

Safe/Safe = Unsafe?!



Measure: A (IRC Safe)
B (IRC Safe)

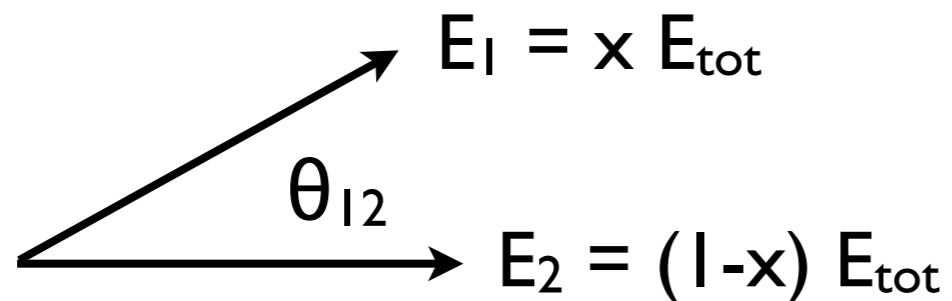
Infer: $R = A/B$ (IRC Unsafe)

Why unsafe?

Every R sensitive to singularity

Simple Example

Mass/Energy = Unsafe



$$\frac{m^2}{E^2} = x(1-x)(\theta_{12})^2$$

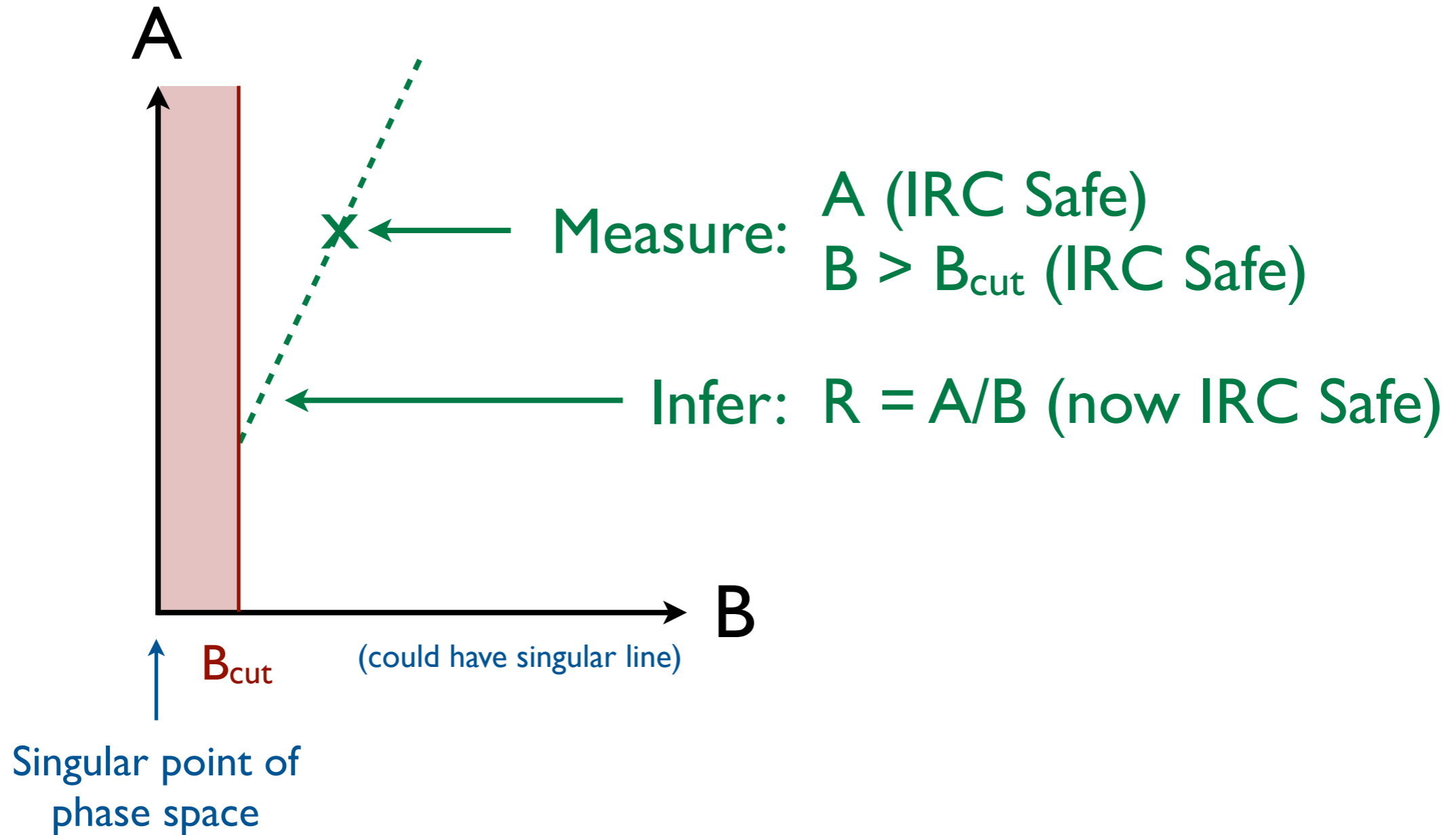
IRC Unsafe: Take soft limit $E_1, E_2 \rightarrow 0$ with x, θ_{12} fixed
Get arbitrary answer

Complaint: But if $E_1, E_2 \rightarrow 0$, no jet to measure

Solution: Require $E > E_{\text{cut}}$

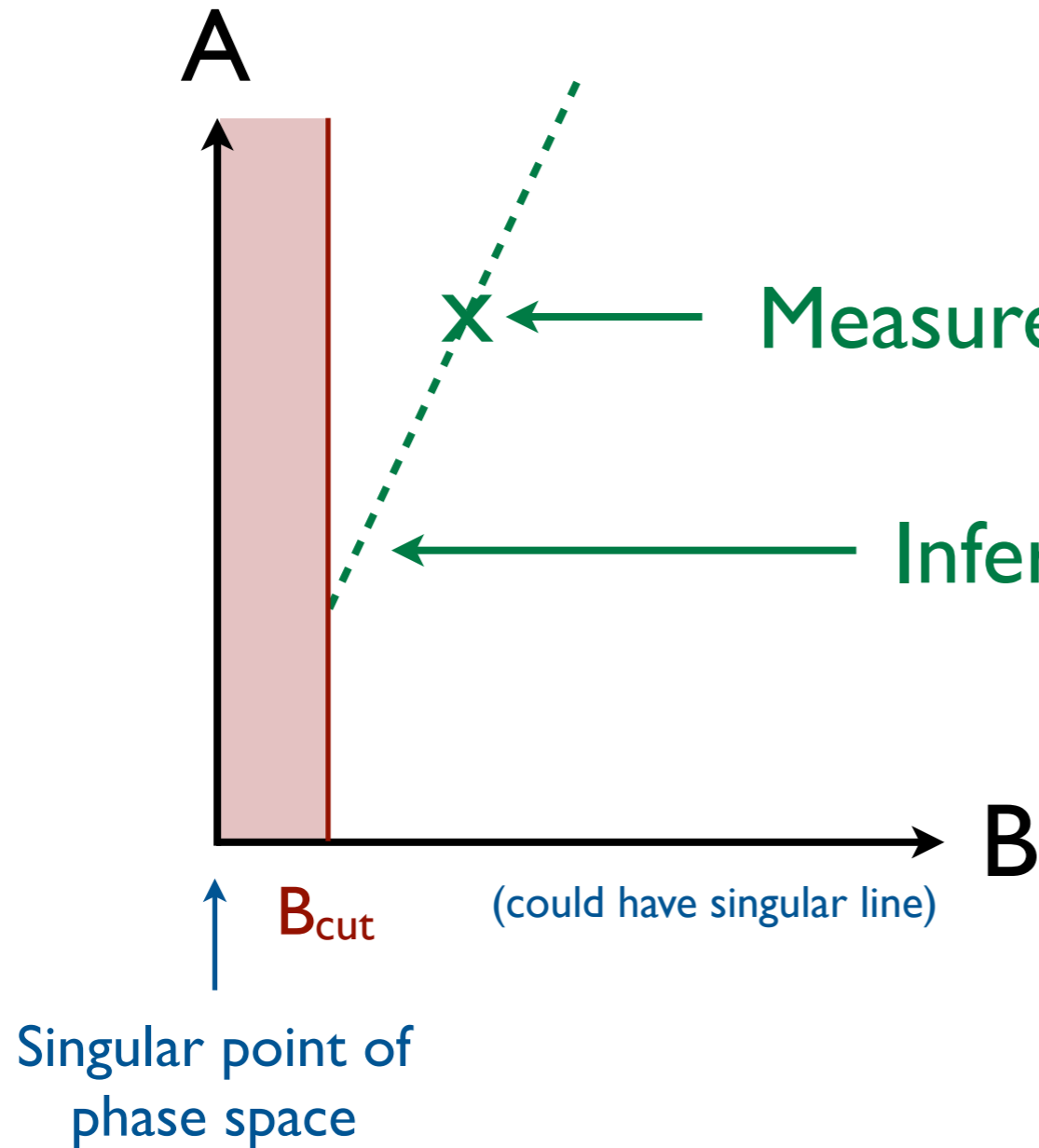
Ah, ok.

Have to put cut on the denominator

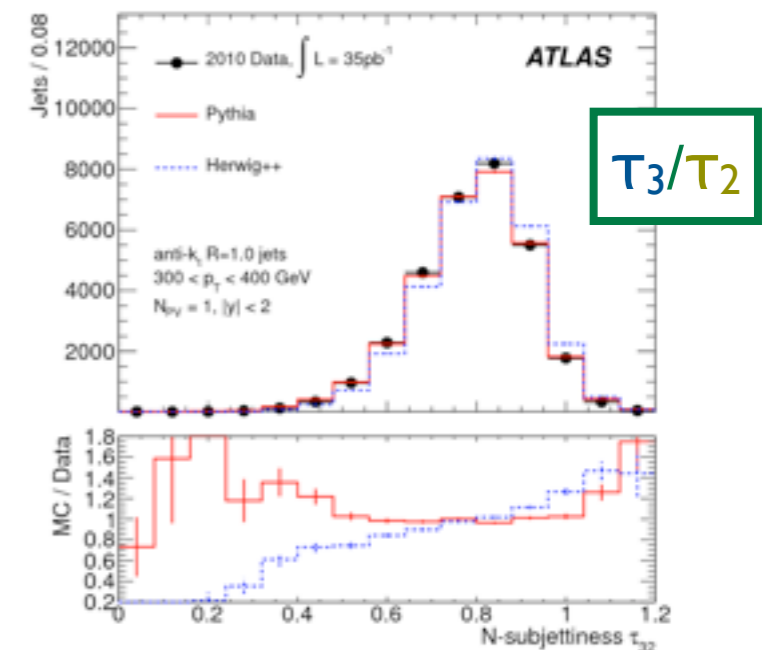


Ah, ok.

Have to put cut on the denominator



No cut on T_2 :
Strictly speaking
IRC unsafe



Surprising/Frustrating Fact:

$$\text{IRC Unsafe Ratio} = \frac{\text{IRC Safe Numerator}}{\text{IRC Safe Denominator}}$$

Ratios observables not calculable at any fixed order in α_s
(unless you cut on the denominator)

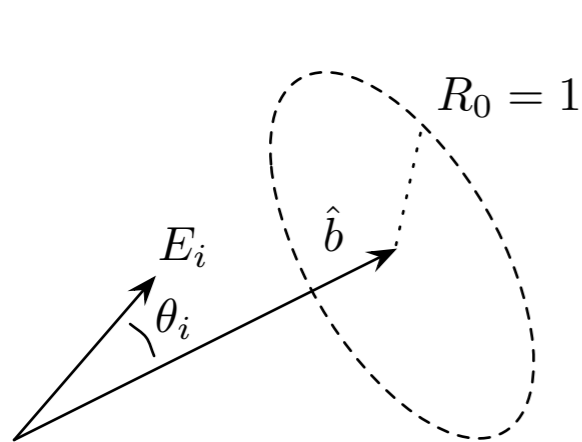
Surprising Resolution:

Ratios observables *are* calculable in QCD
if one accounts for *all orders* in α_s (even without a cut)

Corollary: Parton showers capture the dominant behavior of ratio observables

Ratios of Angularities

Simplest Test Case

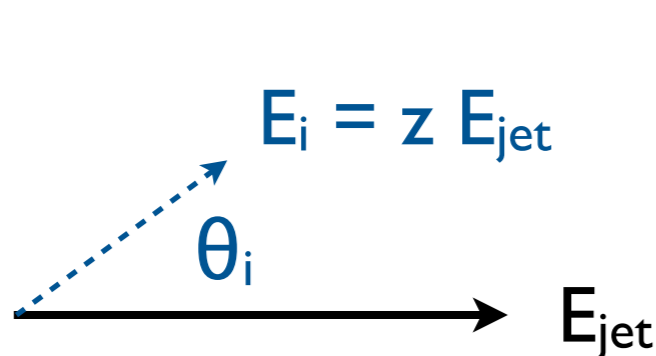


$$e_\beta = \frac{1}{E_{\text{jet}}} \sum_i E_i (\theta_i)^\beta$$

$\beta = 1$: Broadening
 $\beta = 2$: Thrust $\approx m^2/E^2$

Recoil subtlety: Measure angles with respect to “broadening axis”
 Alternative: Energy-energy correlation function with same β , i.e. $C_1^{(\beta)}$

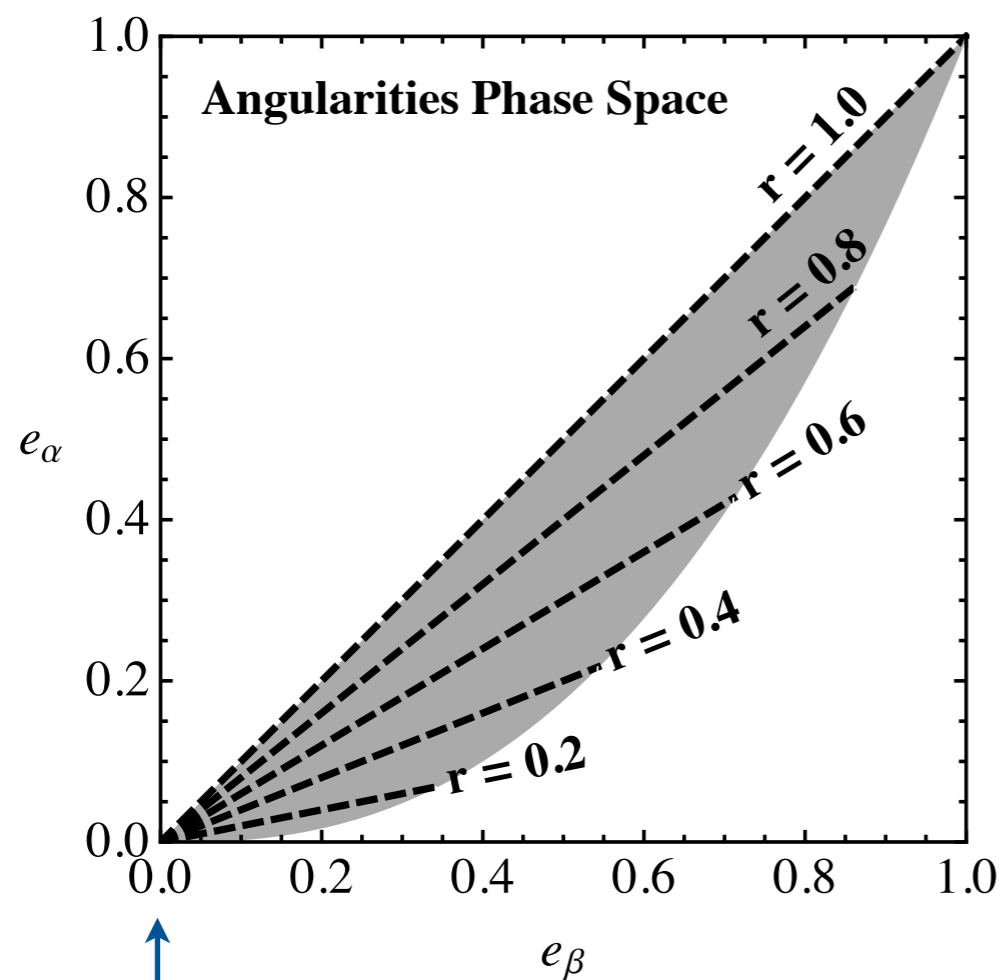
In soft & collinear limit (single emission):



$$e_\beta = z \theta^\beta \quad \text{IRC Safe}$$

$$r = \frac{e_\alpha}{e_\beta} = \theta^{\alpha-\beta} \quad \text{IRC Unsafe: Arbitrary answer in } z \rightarrow 0 \text{ limit}$$

Angularities Phase Space



$$e_\beta = \sum_i z_i (\theta_i)^\beta \quad r = \frac{e_\alpha}{e_\beta}$$

For $\alpha > \beta$ (set $R = 1$):

$$0 < e_\alpha < e_\beta \quad (0 < r < 1)$$

Single emission:

$$(e_\alpha)^\beta > (e_\beta)^\alpha$$

Singular point of
phase space

IRC Unsafe:

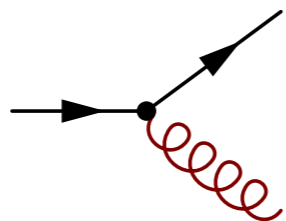
Each r has real emission singularity,
but virtual diagrams only have “ $r = 0$ ” singularity

Fixed-Order vs. Strongly-Ordered

$$\frac{d\sigma}{dr} = \int de_\alpha de_\beta \frac{d^2\sigma}{de_\alpha de_\beta} \delta\left(r - \frac{e_\alpha}{e_\beta}\right)$$



IRC Unsafe
Infinity at $O(\alpha_s)$



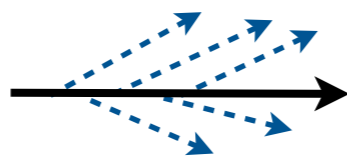
IRC Safe

“I can simultaneously measure e_α and e_β ”

vs.



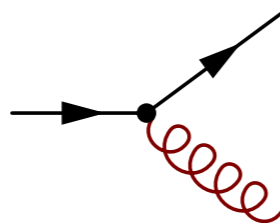
“Sudakov Safe”
Finite answer with
all orders in α_s



Fixed-Order vs. Strongly-Ordered

$$\frac{d\sigma}{dr} = \int de_\alpha de_\beta \frac{d^2\sigma}{de_\alpha de_\beta} \delta\left(r - \frac{e_\alpha}{e_\beta}\right)$$

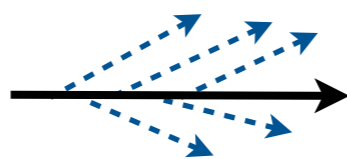
↑
IRC Unsafe
 Infinity at $O(\alpha_s)$



↑
IRC Safe
 “I can simultaneously measure e_α and e_β ”

vs.

↑
“Sudakov Safe”
 Finite answer with
 all orders in α_s

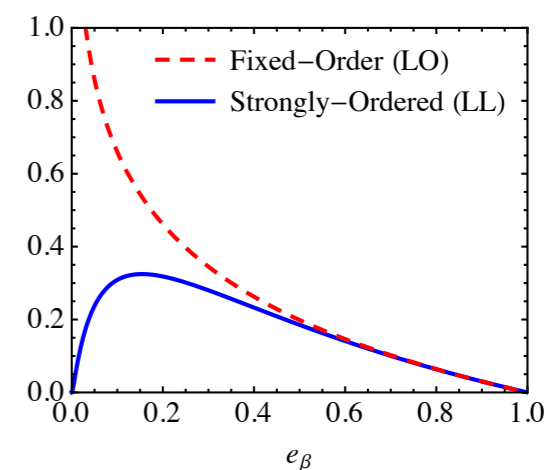


Two Options for a Finite Result:

Put explicit cut on e_β

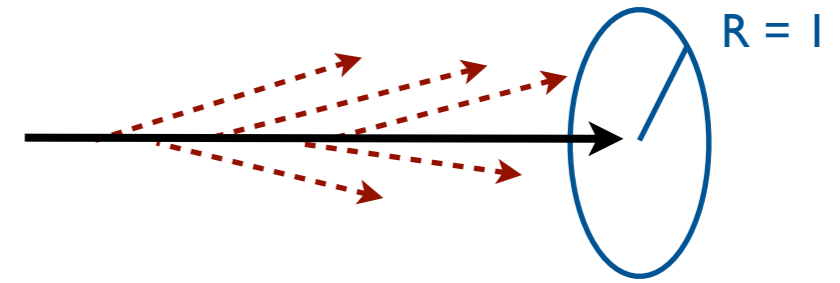
vs.

Sudakov factor
 effectively cuts on e_β



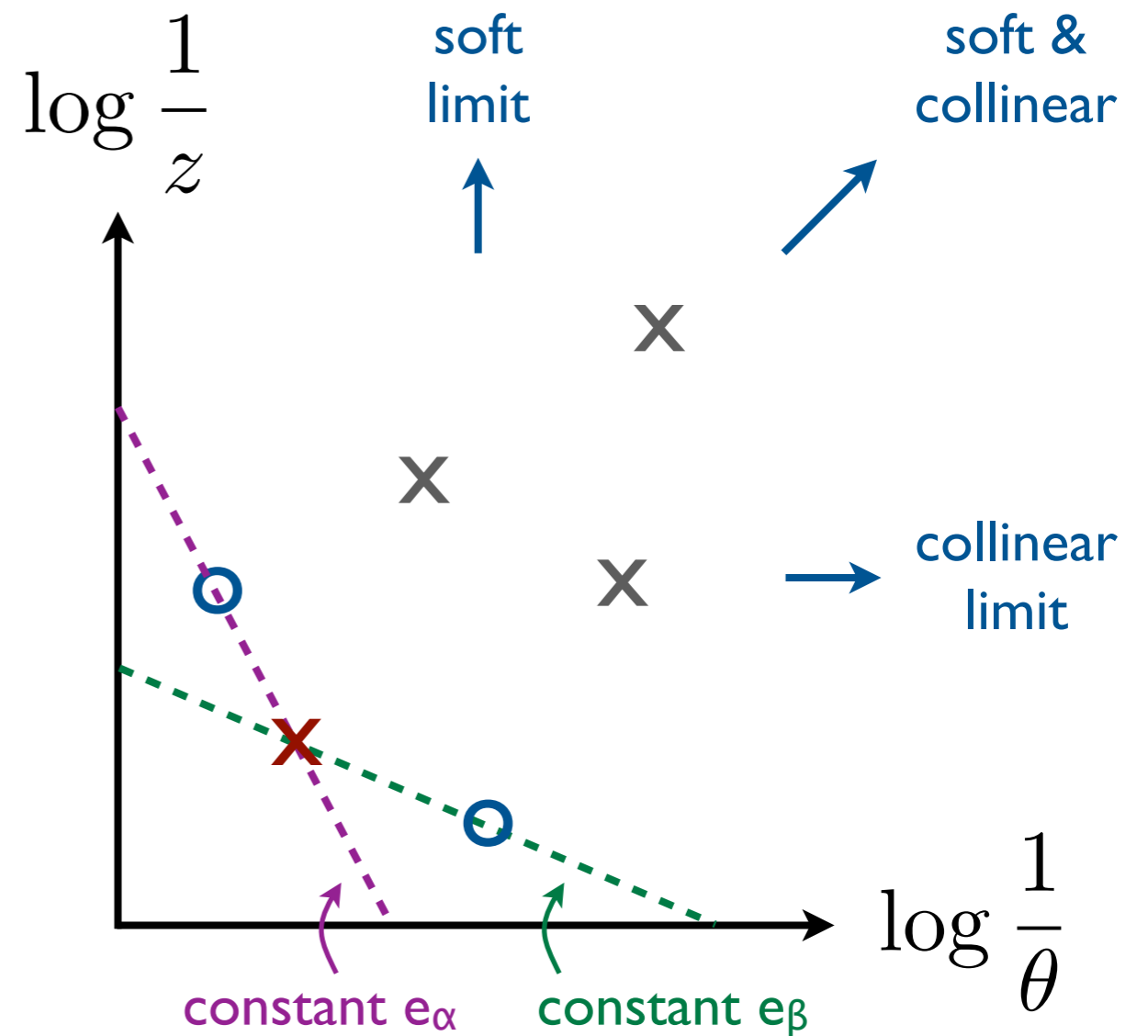
In Strongly-Ordered Limit

Recall introductory talk (set $R = 1$)



$$e_\beta = z \theta^\beta \quad r = \frac{e_\alpha}{e_\beta}$$

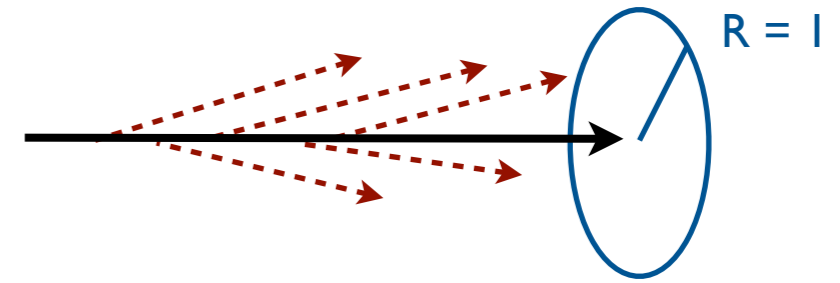
One emission (x) can set both e_α & e_β ,
or can be set by two different (o) emissions.



$$\log \frac{1}{e_\beta} = \log \frac{1}{z} + \beta \log \frac{1}{\theta}$$

In Strongly-Ordered Limit

Recall introductory talk (set $R = 1$)



$$e_\beta = z \theta^\beta \quad r = \frac{e_\alpha}{e_\beta}$$

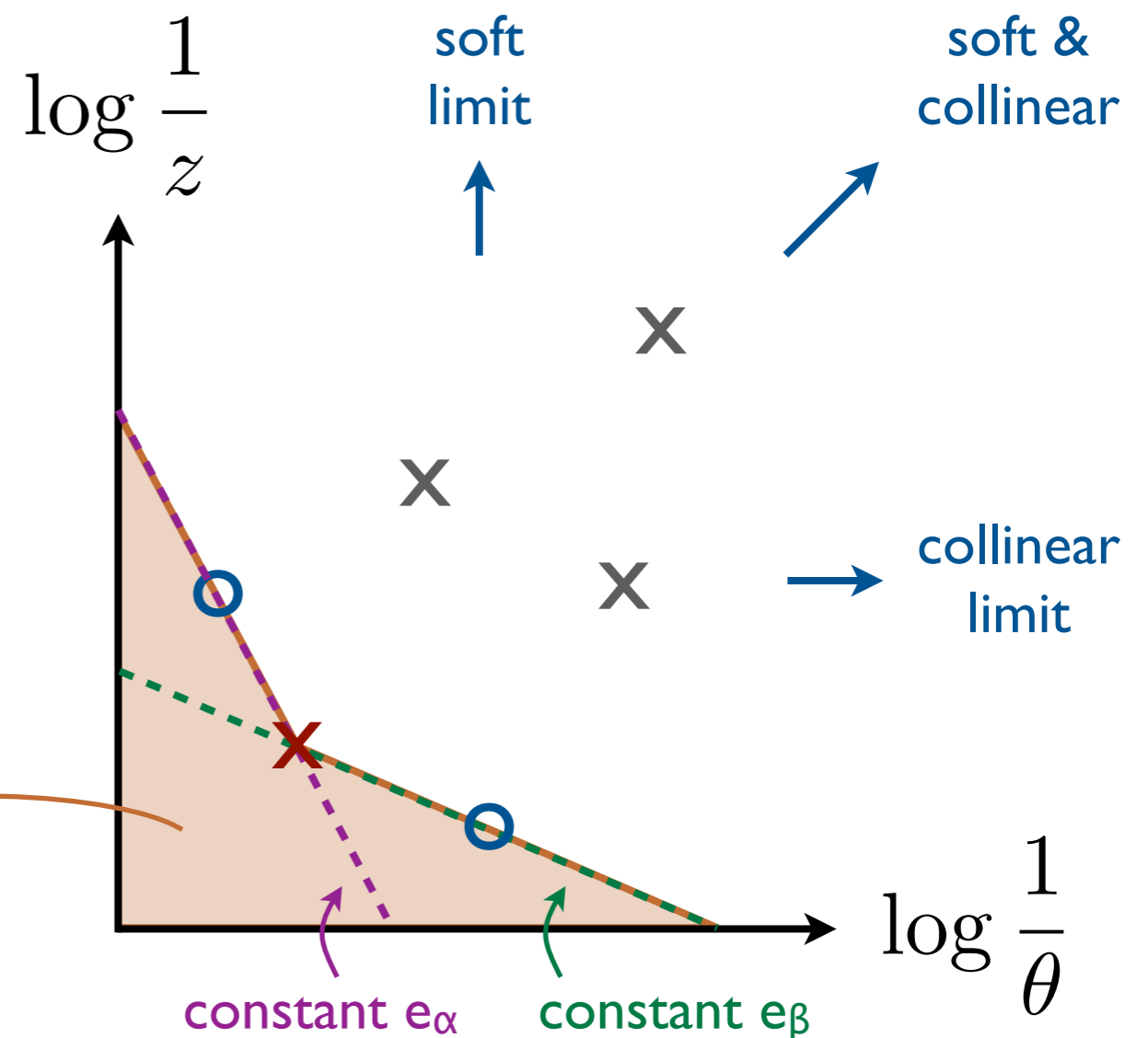
One emission (x) can set both e_α & e_β ,
or can be set by two different (o) emissions.

Sudakov Factor

(probability to get measurement below
a certain value of both e_α & e_β)

$$\Delta(e_\alpha, e_\beta) = e^{-\frac{2\alpha_s C_F}{\pi}}$$

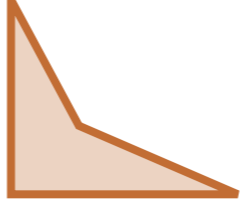
Captures leading logs of
 e_α & e_β distribution



$$\log \frac{1}{e_\beta} = \log \frac{1}{z} + \beta \log \frac{1}{\theta}$$

Turning the Crank

Sudakov Factor:

$$\Delta(e_\alpha, e_\beta) = \exp \left[-\frac{2\alpha_s C_F}{\pi} \right]$$


Double Differential Cross Section:

$$\frac{d^2\sigma^{\text{LL}}}{de_\alpha de_\beta} = \frac{\partial}{\partial e_\alpha} \frac{\partial}{\partial e_\beta} \Delta(e_\alpha, e_\beta)$$

Ratio Cross Section:

$$\frac{d\sigma^{\text{LL}}}{dr} = \int de_\alpha de_\beta \frac{d^2\sigma^{\text{LL}}}{de_\alpha de_\beta} \delta \left(r - \frac{e_\alpha}{e_\beta} \right)$$

Turning the Crank

Sudakov Factor:

$$\Delta(e_\alpha, e_\beta) = \exp \left[-\frac{\alpha_s}{\pi} C_F \left(\frac{1}{\beta} \log^2 e_\beta + \frac{1}{\alpha - \beta} \log^2 \frac{e_\alpha}{e_\beta} \right) \right]$$

Double Differential Cross Section:

$$\frac{d^2 \sigma^{\text{LL}}}{de_\alpha de_\beta} =$$

Ratio Cross Section:

$$\frac{d\sigma^{\text{LL}}}{dr} =$$

Turning the Crank

Sudakov Factor:

$$\Delta(e_\alpha, e_\beta) = \exp \left[-\frac{\alpha_s}{\pi} C_F \left(\frac{1}{\beta} \log^2 e_\beta + \frac{1}{\alpha - \beta} \log^2 \frac{e_\alpha}{e_\beta} \right) \right]$$

Double Differential Cross Section:

$$\frac{d^2 \sigma^{\text{LL}}}{de_\alpha de_\beta} = \left(\frac{2\alpha_s}{\pi} \frac{C_F}{\alpha - \beta} \frac{1}{e_\alpha e_\beta} + \frac{4\alpha_s^2}{\pi^2} \frac{C_F^2}{\beta(\alpha - \beta)^2} \frac{1}{e_\alpha e_\beta} \log \frac{e_\beta}{e_\alpha} \log \frac{e_\alpha^\beta}{e_\beta^\alpha} \right) \Delta(e_\alpha, e_\beta)$$

(Cross check: Reduces to known single differential)

Ratio Cross Section:

$$\frac{d\sigma^{\text{LL}}}{dr} =$$

Turning the Crank

Sudakov Factor:

$$\Delta(e_\alpha, e_\beta) = \exp \left[-\frac{\alpha_s}{\pi} C_F \left(\frac{1}{\beta} \log^2 e_\beta + \frac{1}{\alpha - \beta} \log^2 \frac{e_\alpha}{e_\beta} \right) \right]$$

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Ratio Cross Section:

$$\frac{d\sigma^{\text{LL}}}{dr} = \frac{\sqrt{\alpha_s C_F \beta}}{\alpha - \beta} \frac{1}{r} \left(1 - 2 \frac{\alpha_s}{\pi} \frac{C_F}{\alpha - \beta} \log^2 r \right) \left(\operatorname{erf} \left[\frac{\sqrt{\alpha_s C_F \beta}}{\sqrt{\pi}(\alpha - \beta)} \log r \right] + 1 \right) e^{-\frac{\alpha_s}{\pi} \frac{C_F}{\alpha - \beta} \log^2 r} - 2 \frac{\alpha_s}{\pi} \frac{C_F}{\alpha - \beta} \frac{\log r}{r} e^{-\frac{\alpha_s}{\pi} C_F \frac{\alpha}{(\alpha - \beta)^2} \log^2 r}$$

Turning the Crank

Ratio Cross Section:

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Expanded in small α_s :

$$\frac{d\sigma^{\text{LL}}}{dr} = \sqrt{\alpha_s} \frac{\sqrt{C_F \beta}}{\alpha - \beta} \frac{1}{r} + \mathcal{O}(\alpha_s)$$

Turning the Crank

Ratio Cross Section:

$$\frac{d\sigma^{\text{LL}}}{dr} = \frac{\sqrt{\alpha_s C_F \beta}}{\alpha - \beta} \frac{1}{r} \left(1 - 2 \frac{\alpha_s}{\pi} \frac{C_F}{\alpha - \beta} \log^2 r \right) \left(\operatorname{erf} \left[\frac{\sqrt{\alpha_s C_F \beta}}{\sqrt{\pi}(\alpha - \beta)} \log r \right] + 1 \right) e^{-\frac{\alpha_s}{\pi} \frac{C_F}{\alpha - \beta} \log^2 r} - 2 \frac{\alpha_s}{\pi} \frac{C_F}{\alpha - \beta} \frac{\log r}{r} e^{-\frac{\alpha_s}{\pi} C_F \frac{\alpha}{(\alpha - \beta)^2} \log^2 r}$$

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Not a valid
Taylor expansion

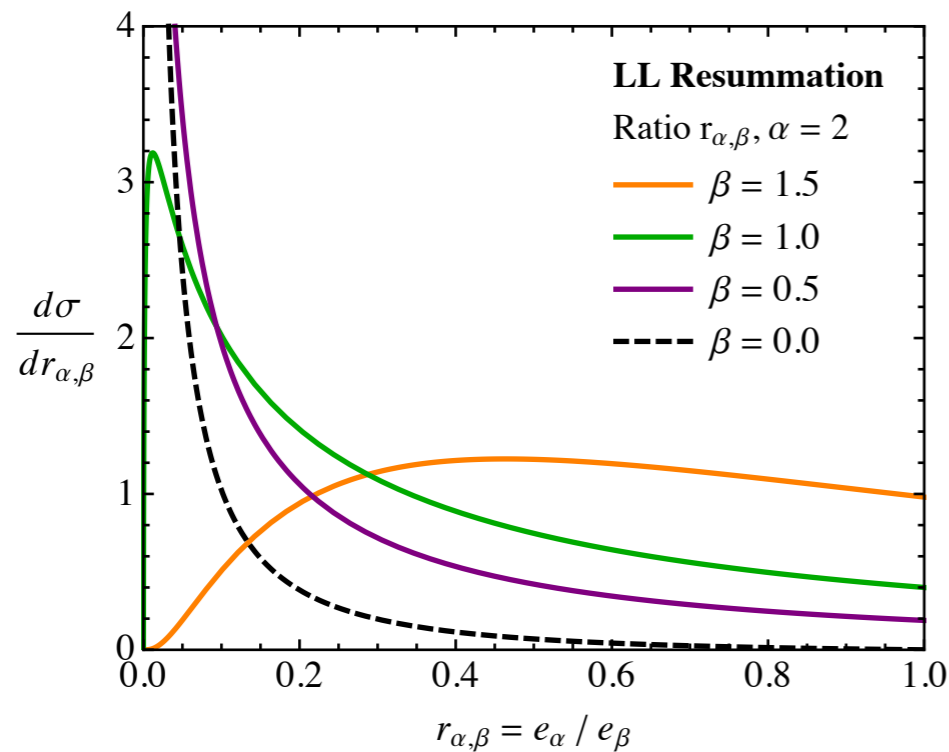
Finite cross section
for all r (even $r = 0$)

Unsafe... ..but Calculable

Computational Results

Leading-Log Result is Systematically Improvable

LL



Three representative ratios

$$\frac{e_2}{e_{1.5}} \quad \frac{e_2}{e_1} \quad \frac{e_2}{e_{0.5}}$$

Usual Angularity

$$\frac{e_2}{e_0} = e_2$$

strongly-ordered

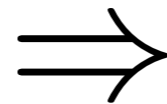
$$\text{LL} \Rightarrow (\alpha_s \log^2 r)^{n/2}$$

In following plots,
dashed always corresponds to
previous level of accuracy

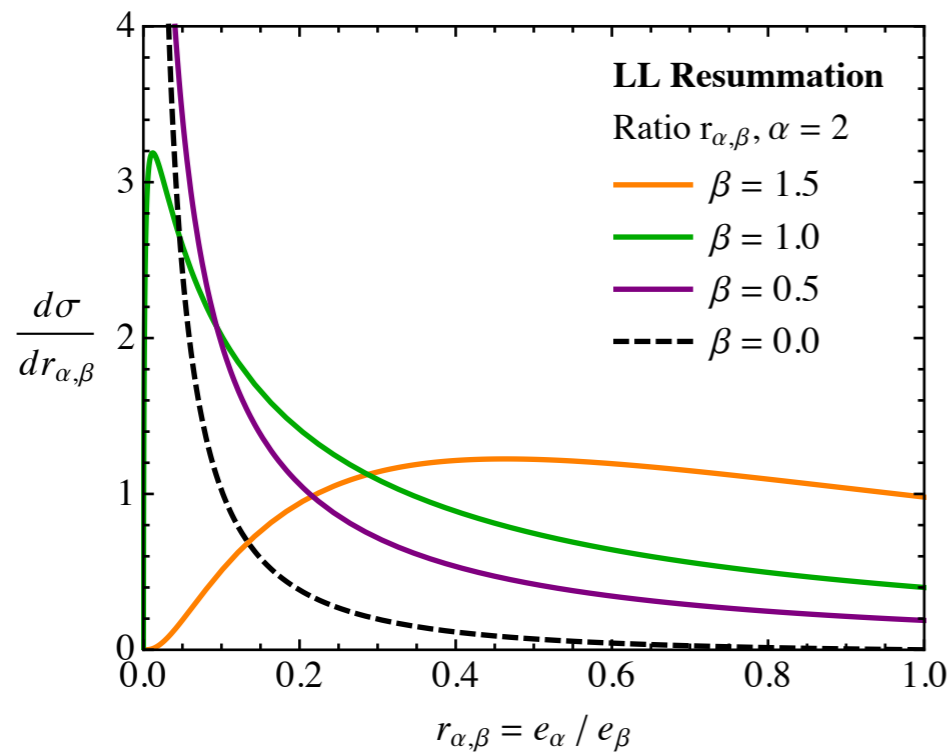
Fixed-Order Matching

Mild impact from exact matrix element

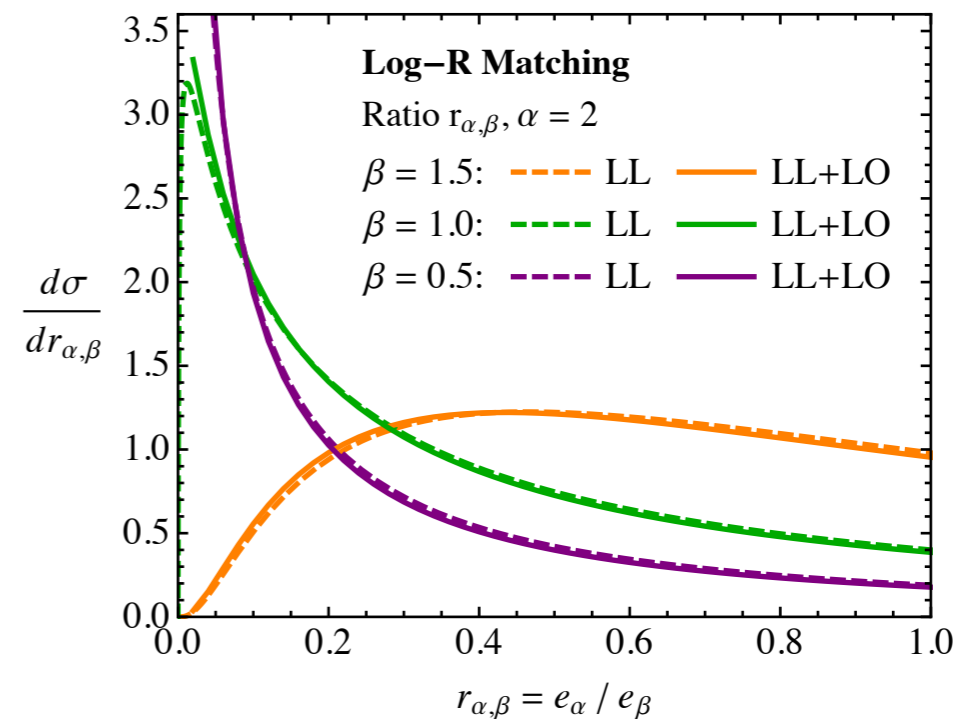
LL



LL+LO



strongly-ordered

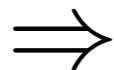


tree-level matching

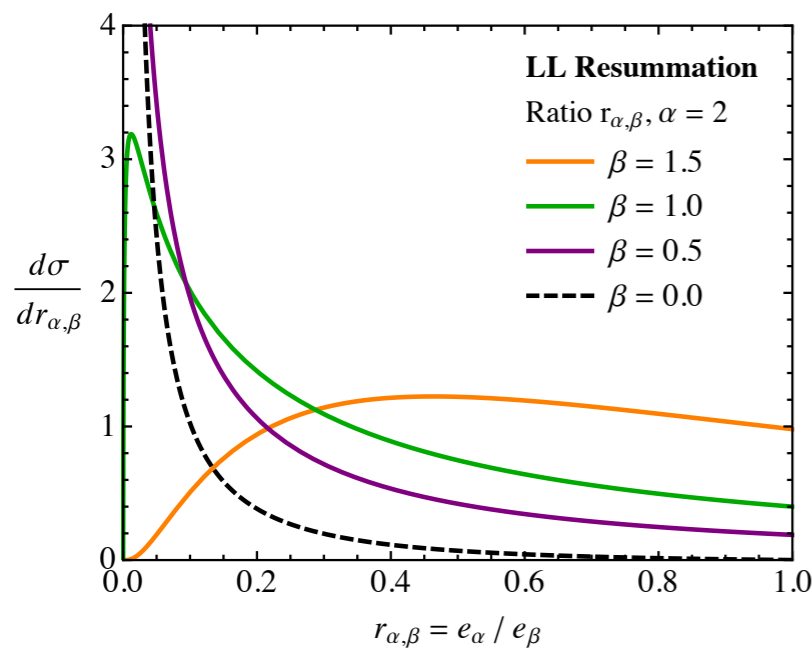
Higher-Order Resummation

Large impact from running coupling: “Modified Leading Log”

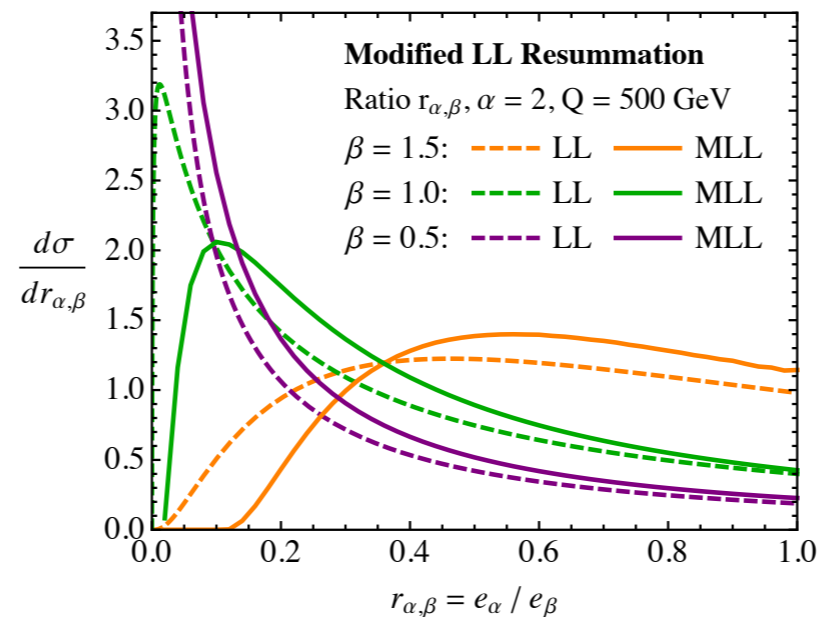
LL



MLL



strongly-ordered

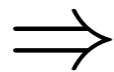


running α_s ,
subleading splitting functions

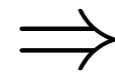
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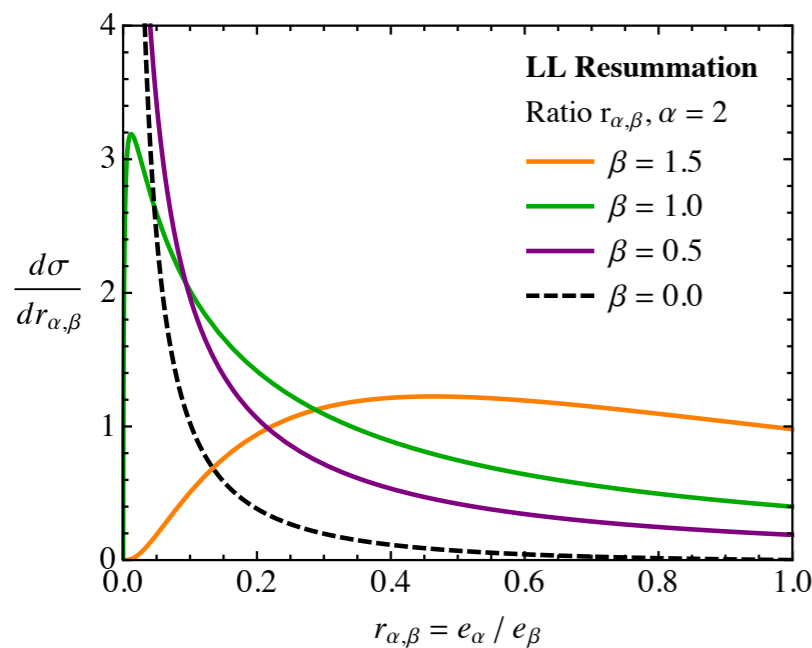
LL



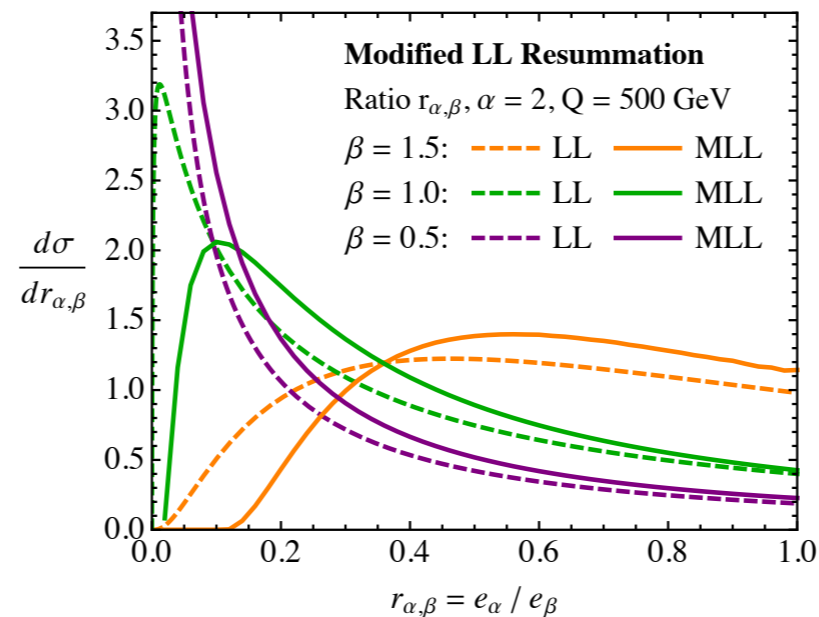
MLL



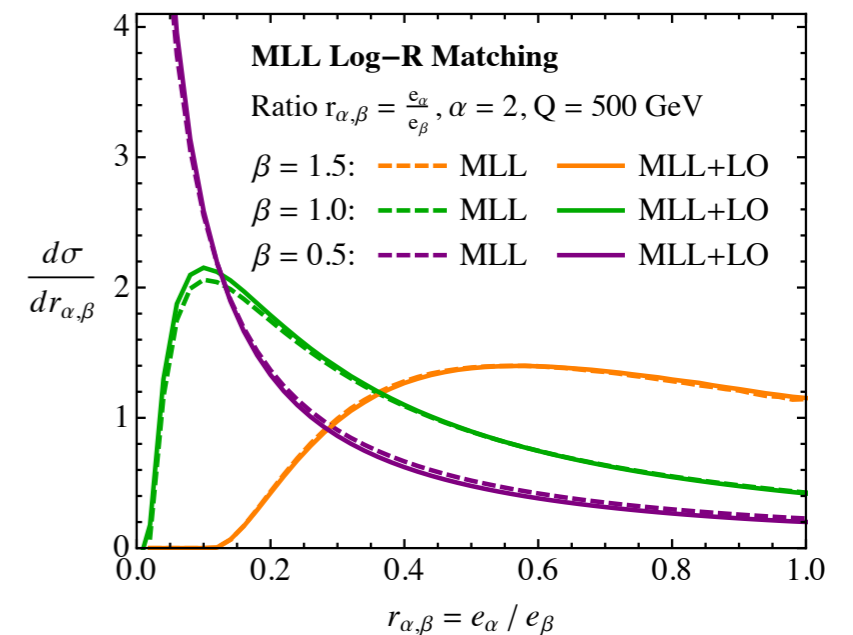
MLL+LO



strongly-ordered



running α_s ,
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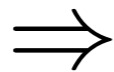


tree-level matching

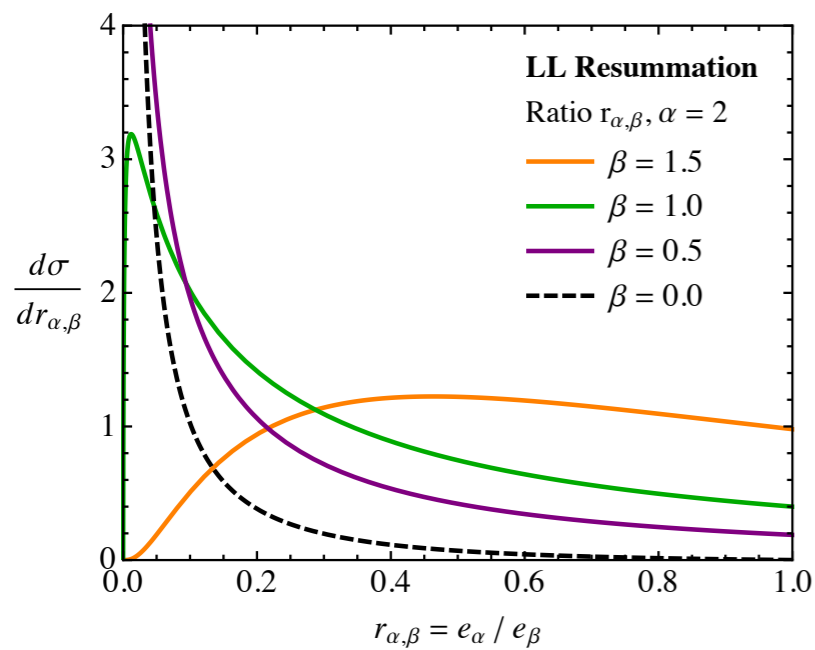
Multiple Emissions

Change in tail region from Monte-Carlo-style resummation

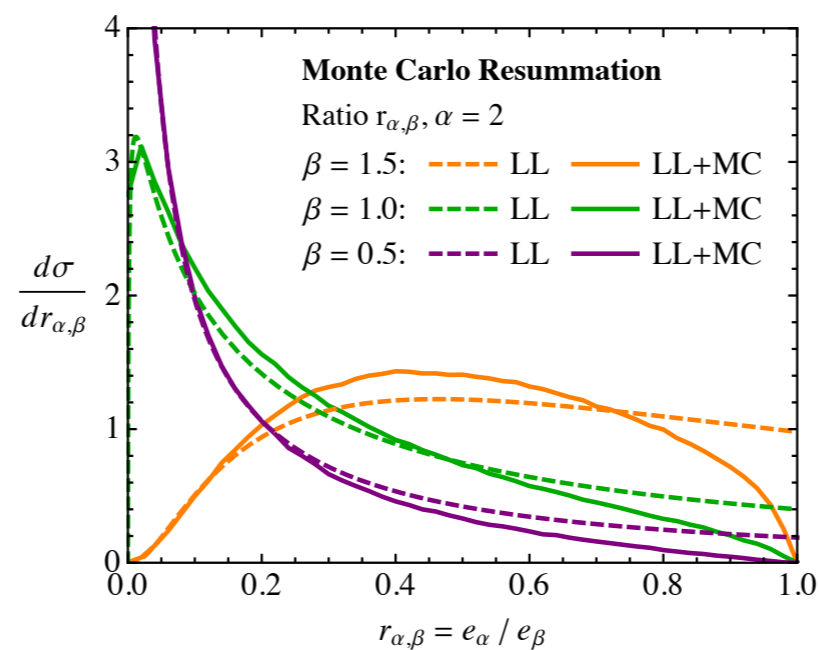
LL



LL+MC



strongly-ordered



multiple emissions

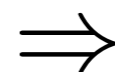
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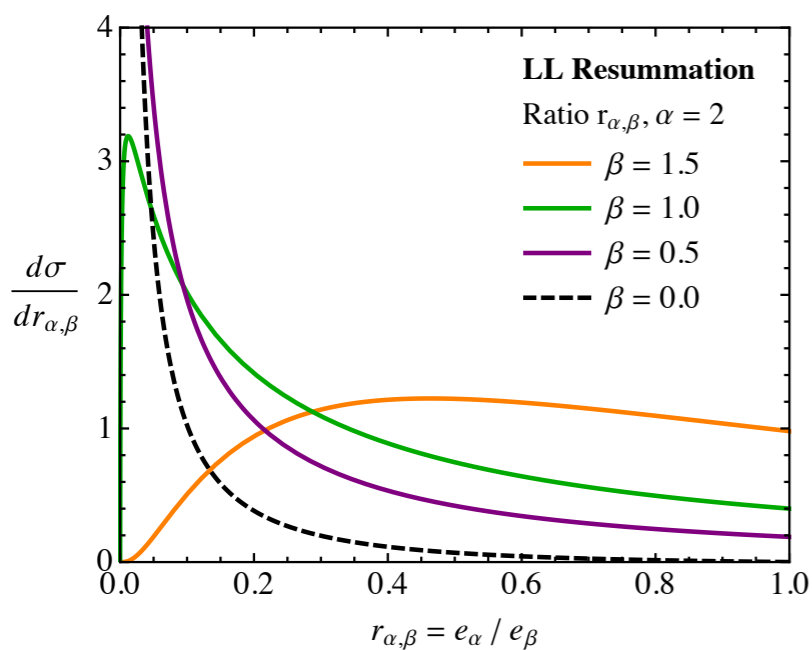
LL



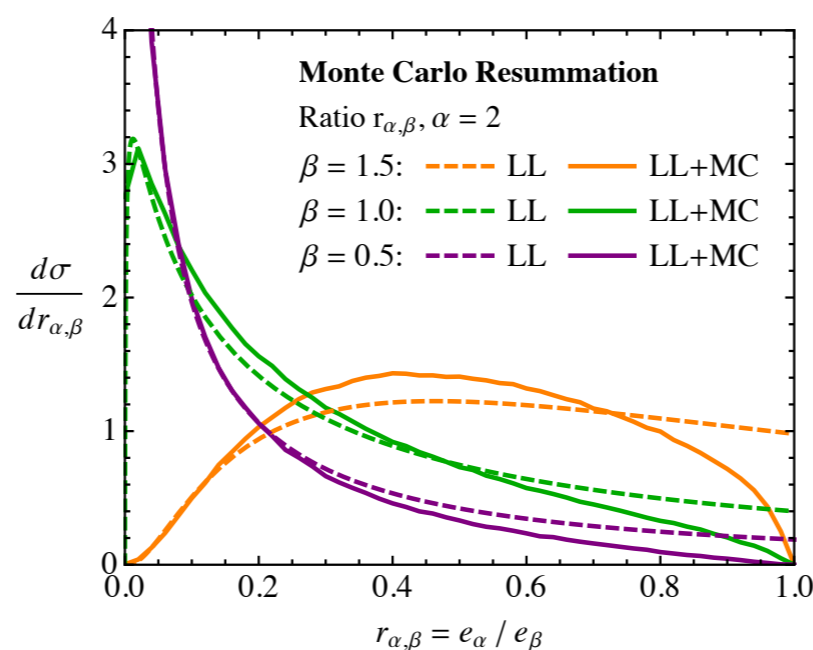
LL+MC



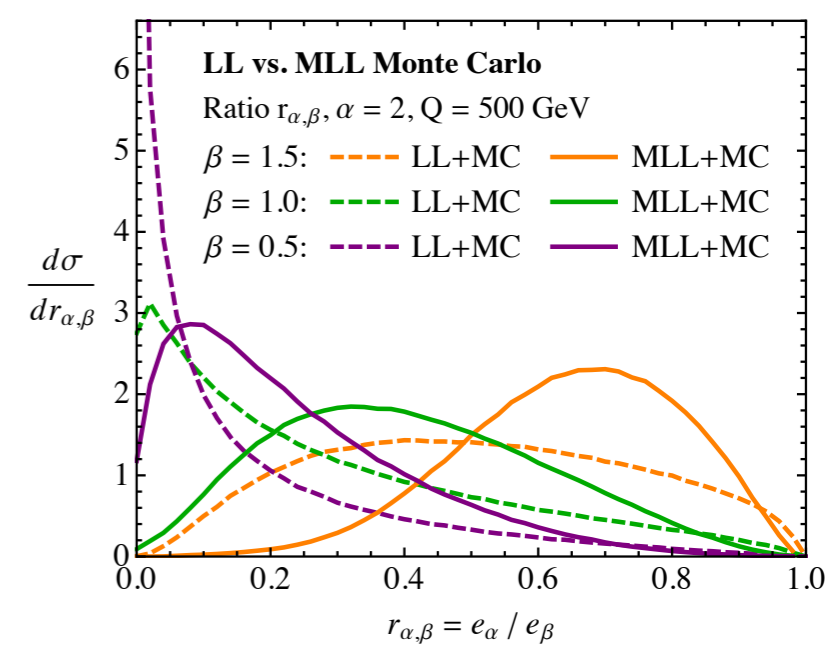
MLL+LO+MC



strongly-ordered



multiple emissions



running α_s ,
AP splitting functions

includes tree-level matching

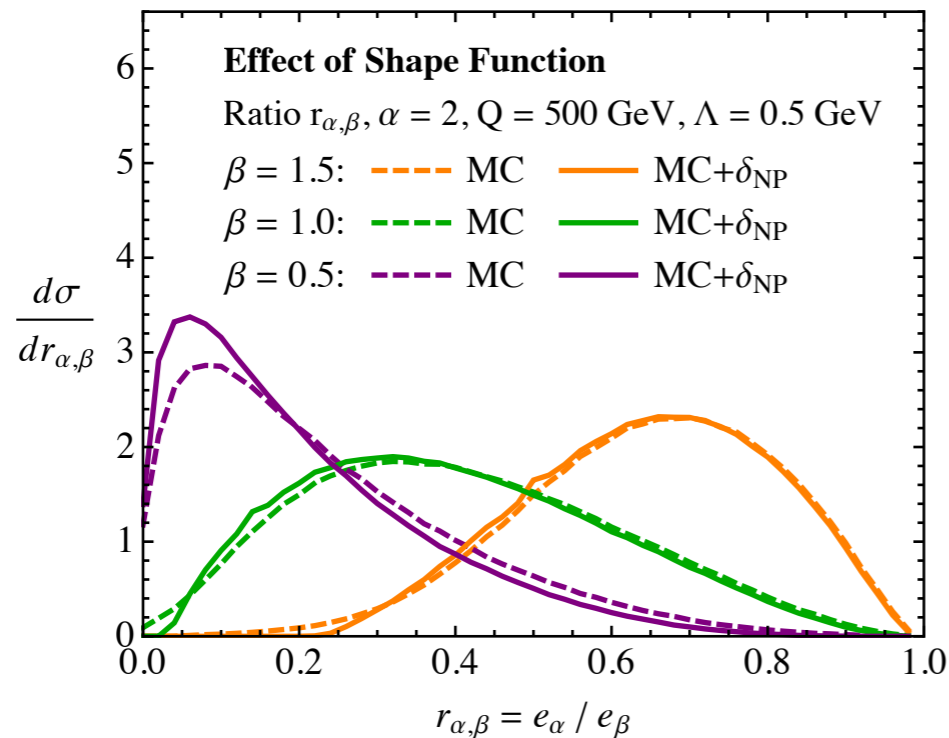
Hadronization Corrections

Power Corrections vs. Hadronization Model

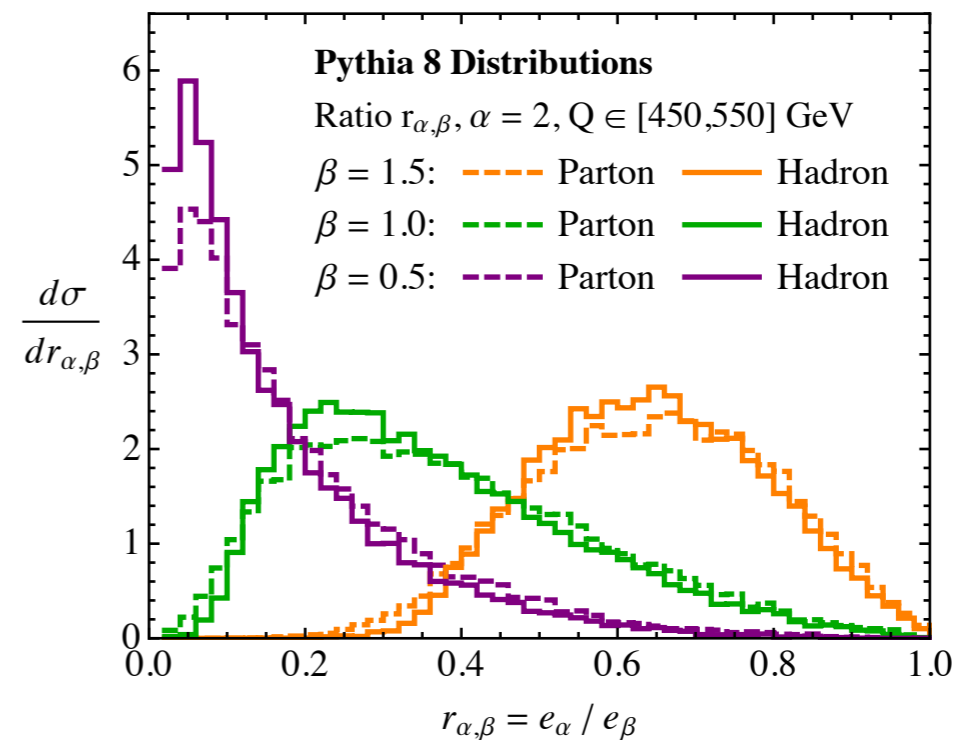
MLL+LO+MC+ δ_{NP}

vs.

Pythia 8



power corrections

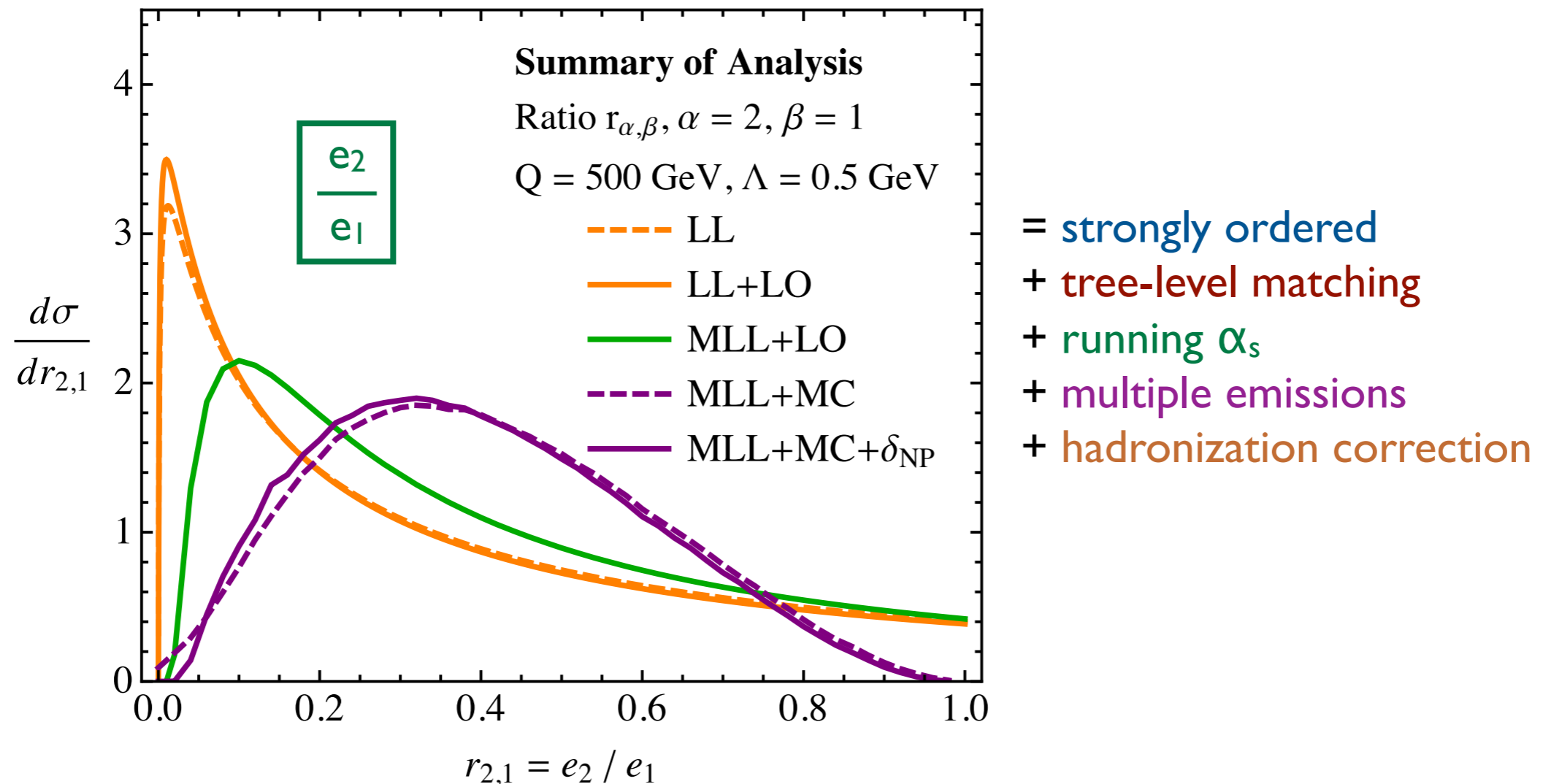


hadronization model

momentum conservation,
color coherence, ...

Building up a Realistic Distribution

All of these (and more) in existing parton shower programs



Expect ratio observables to be reasonably well-described in Monte Carlo

Summary

Ratio Observables: Ubiquitous & Unsafe

*Even if A, B are IRC safe, $R = A/B$ has singularities in the denominator B
Cannot calculate to any fixed order in α_s*

Ratios of Angularities: “Sudakov Safe”

*Sudakov factor regulates singularities
Ratio distribution well-defined with all orders in α_s
Amusing expansion in $\sqrt{\alpha_s}$*

Computational Results

*Can systematically improve strongly-ordered limit
Many improvements already in parton shower Monte Carlos
Ongoing work to calculate 2-prong ratio observables like τ_2/τ_1*

Broader Lesson

If an IRC unsafe observable is at all related to an IRC safe one, then probably a way to calculate

Hadronic Cross Sections = PDFs \otimes IRC Safe

Track-based Observables = Track Functions \otimes IRC Safe

Ratio Observables = IRC Safe / IRC Safe