

# PRUNING AND MASS-DROP WITH ANALYTICAL METHODS

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with

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# Pruning

1. From an initial jet define pruning radius  $R_{\text{prune}} \sim m / p_t$
2. Re-cluster the jet, vetoing recombination for which

$$z = \frac{\min(p_{ti}, p_{tj})}{|\vec{p}_{ti} + \vec{p}_{tj}|} < z_{\text{cut}}$$

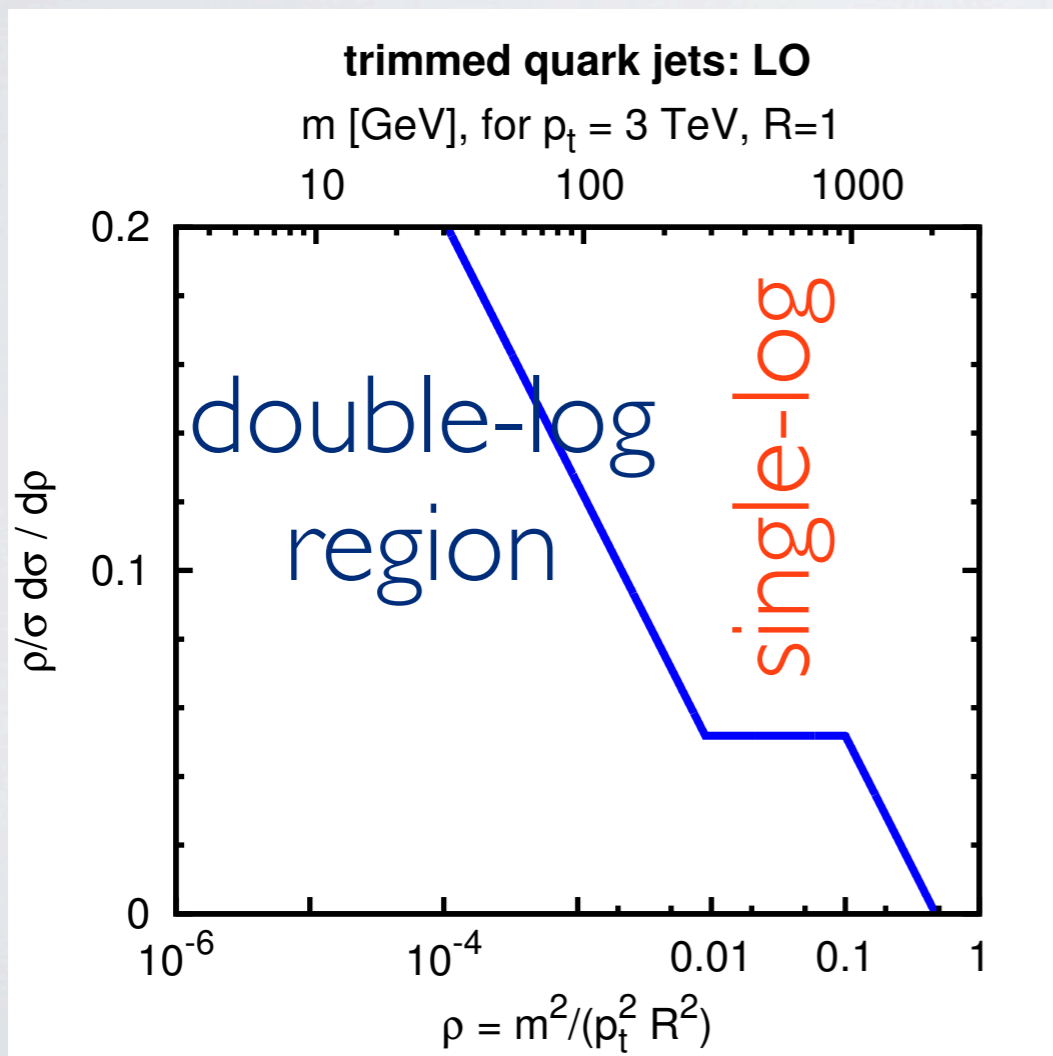
$$d_{ij} > R_{\text{prune}}$$

i.e. soft and wide angle

Ellis, Vermillion and Walsh (2009)

# LO calculation

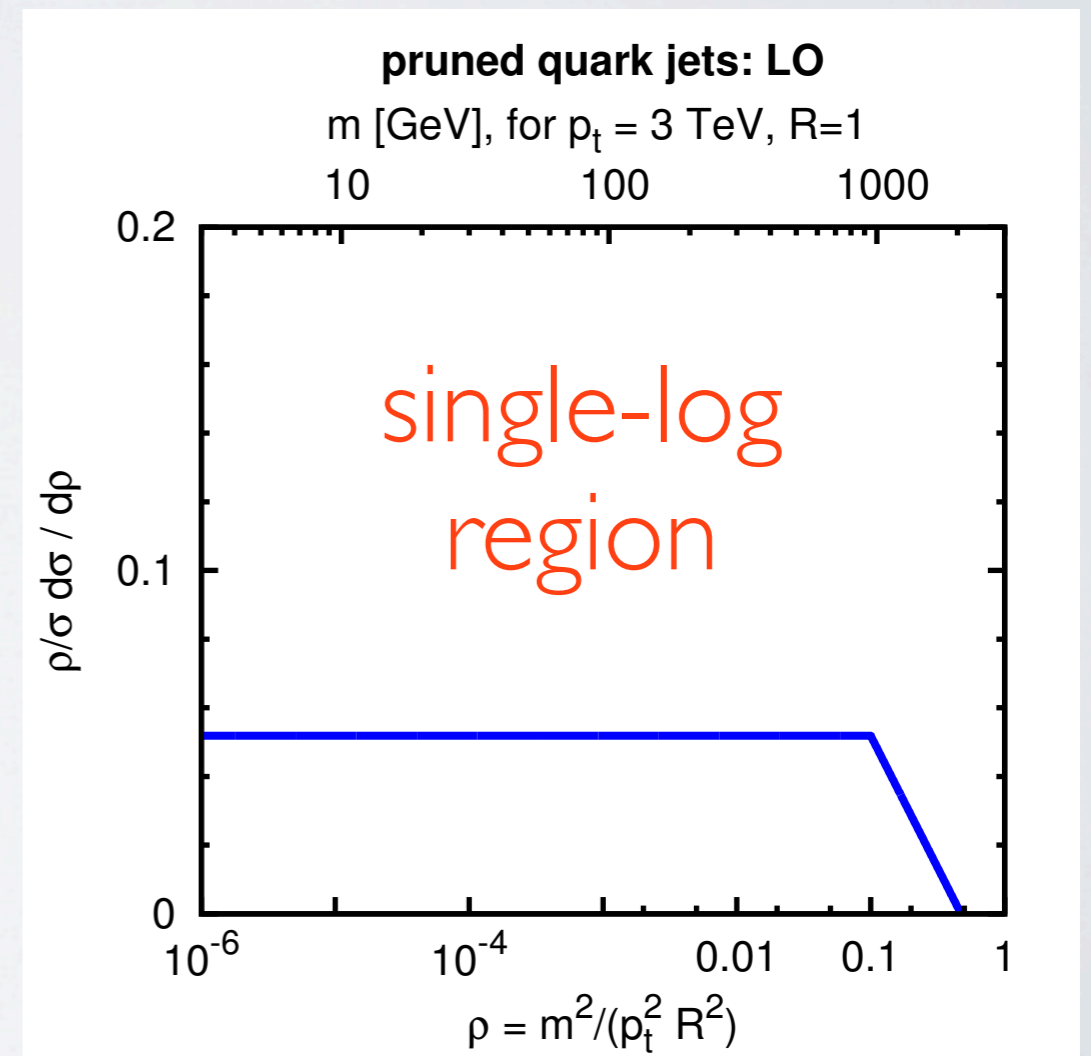
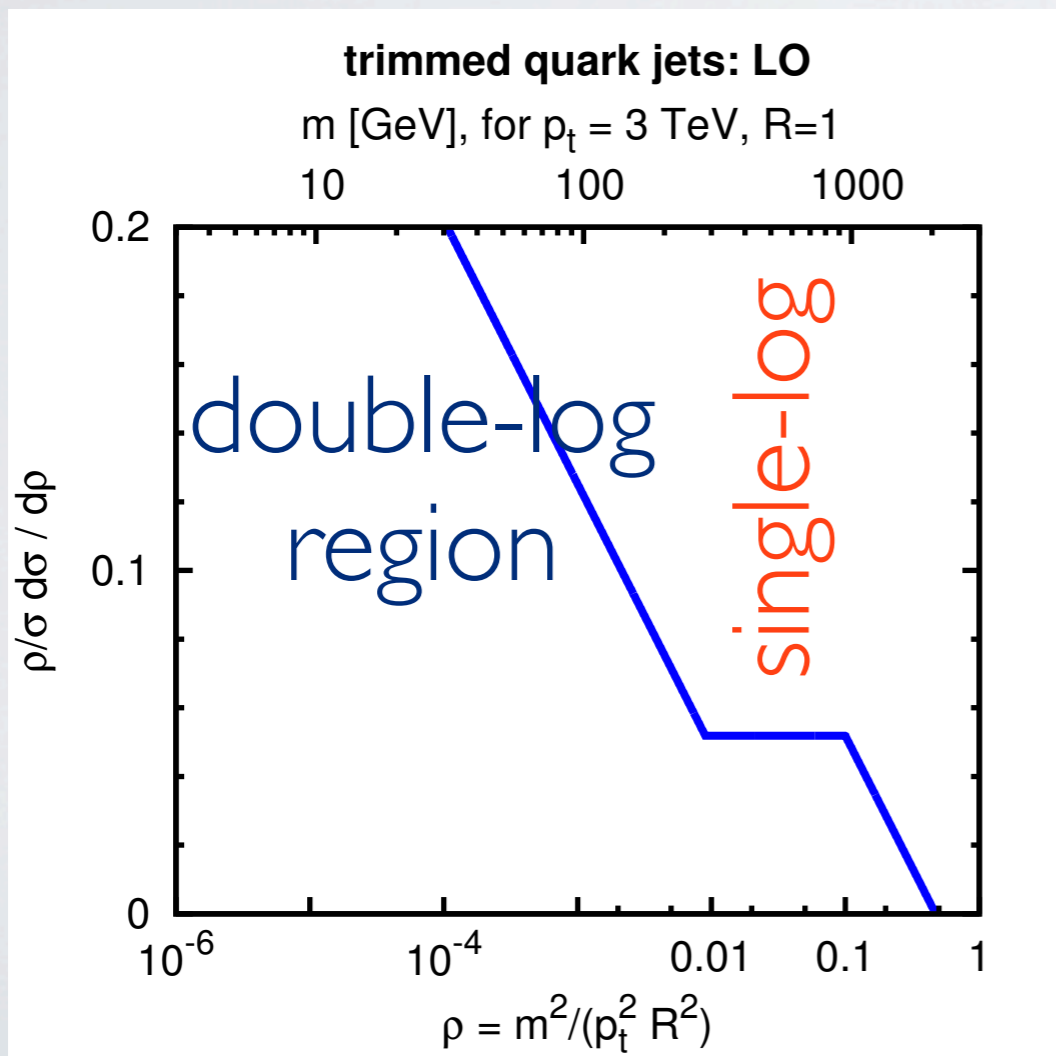
- LO calculation similar to trimming



log-counting for cumulant  
distribution  $\Sigma$

# LO calculation

- LO calculation similar to trimming
- The pruning radius is set dynamically:  $R_{\text{prune}} < d_{ij}$
- The 2 prongs are always tested for  $z_{\text{cut}}$ : single logs

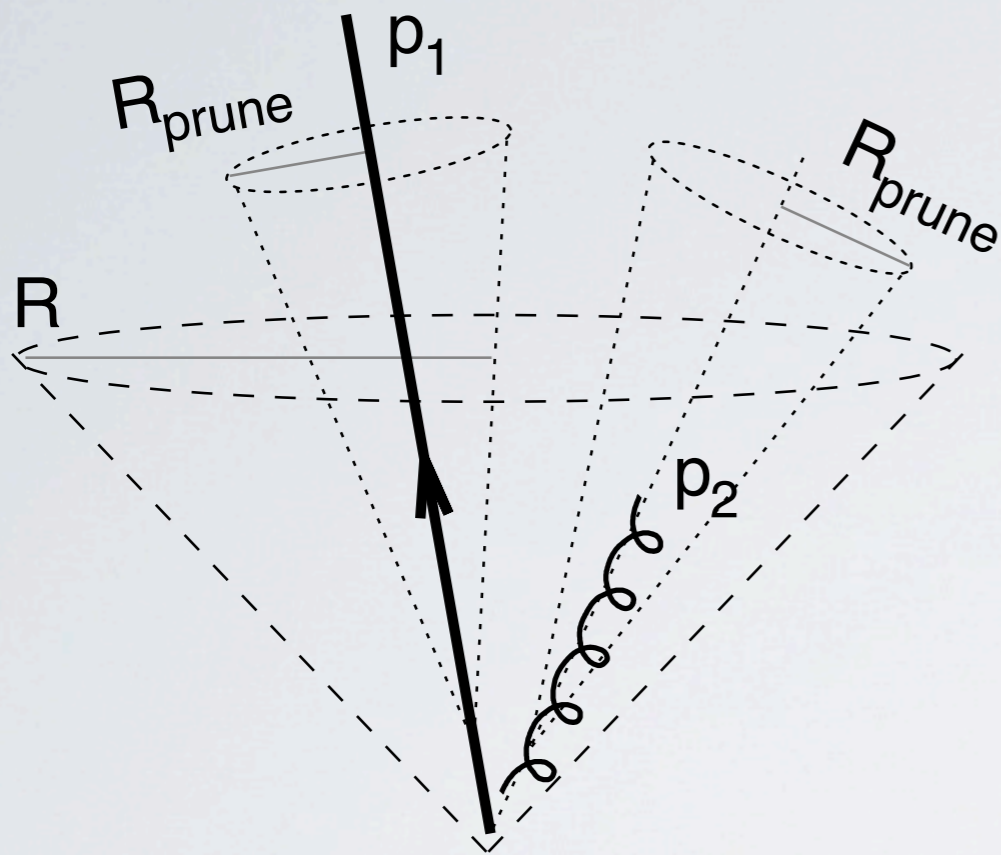


# Beyond LO

## What pruning is meant to do

Choose an  $R_{\text{prune}}$  such that different hard prongs ( $p_1, p_2$ ) end up in different hard subjects.

Discard any softer radiation.

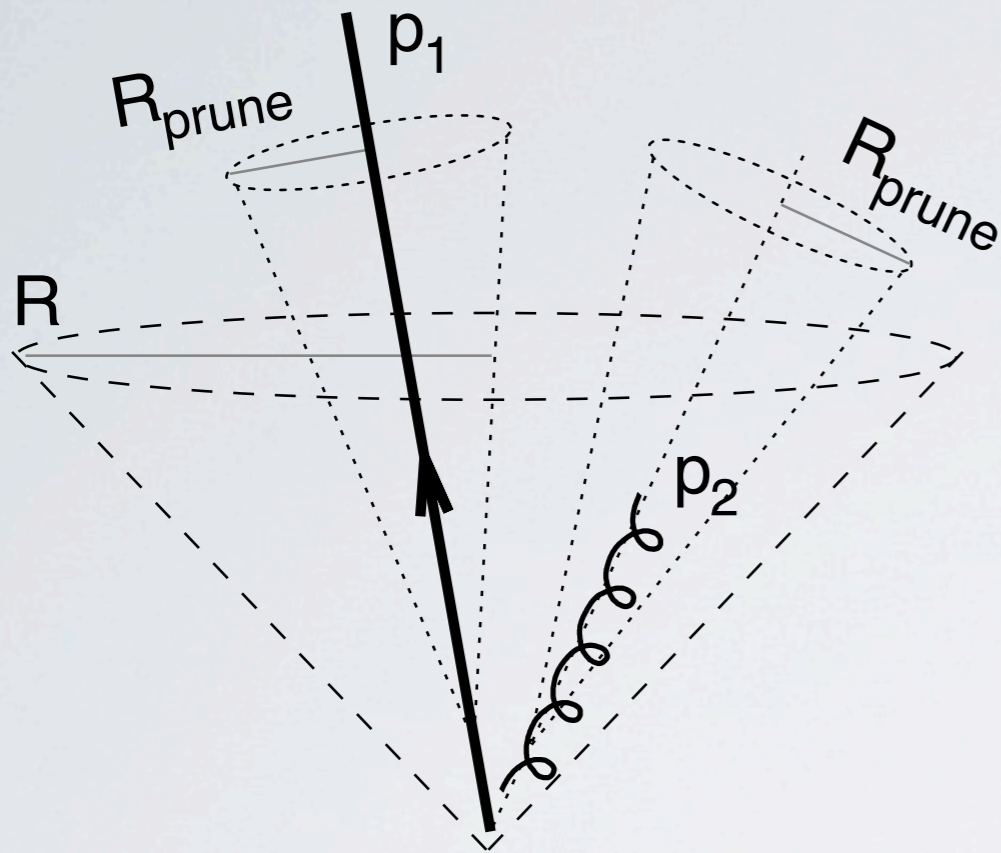


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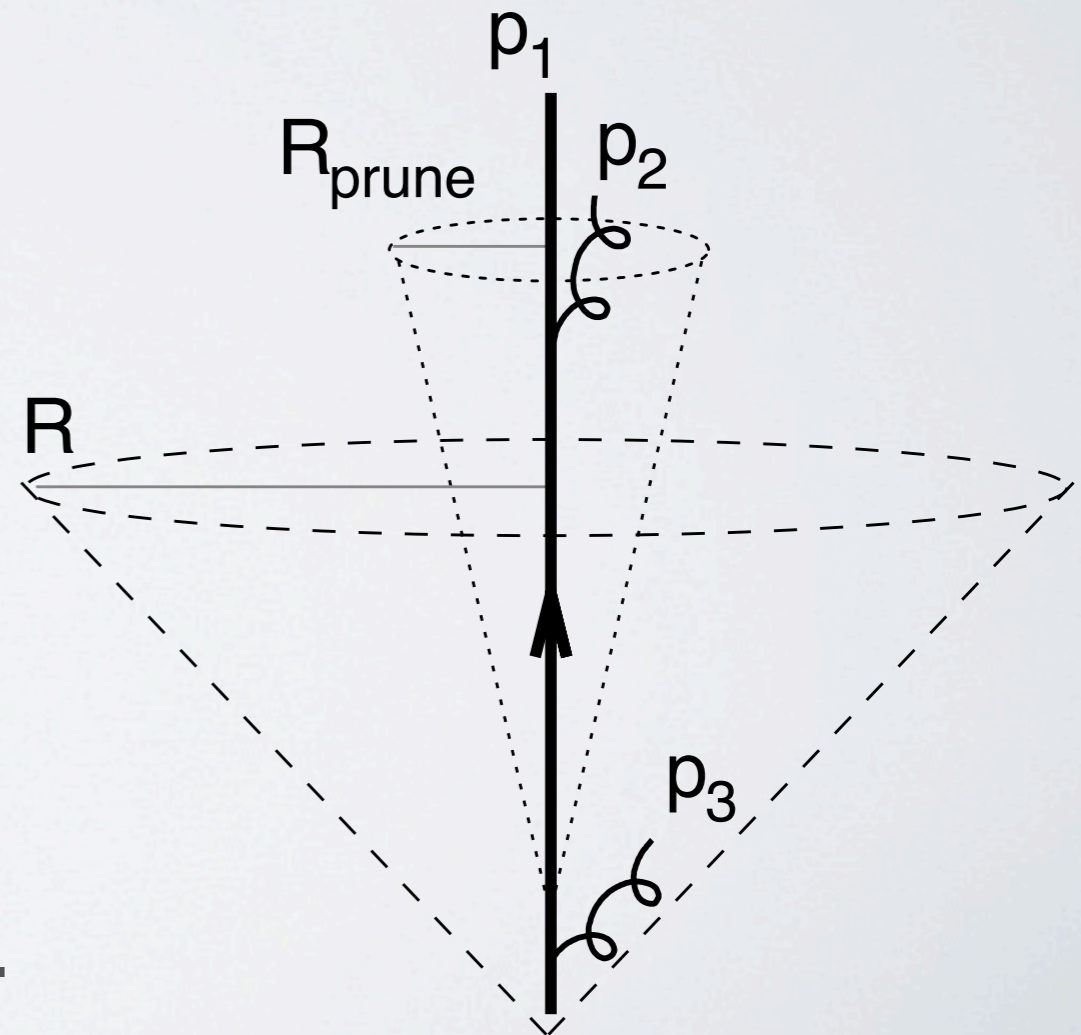
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## What pruning sometimes does

Chooses  $R_{\text{prune}}$  based on a soft  $p_3$  (dominates total jet mass), and leads to a single narrow subjet whose mass is also dominated by a soft emission ( $p_2$ , within  $R_{\text{prune}}$  of  $p_1$ , so not pruned away).

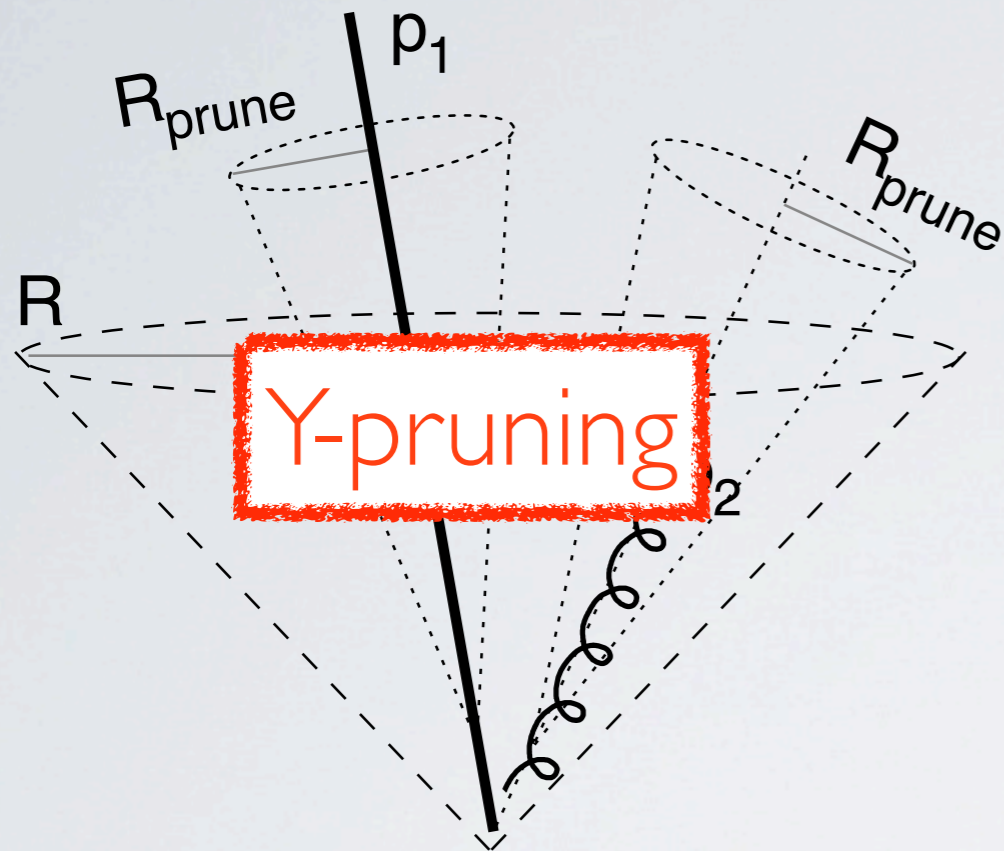


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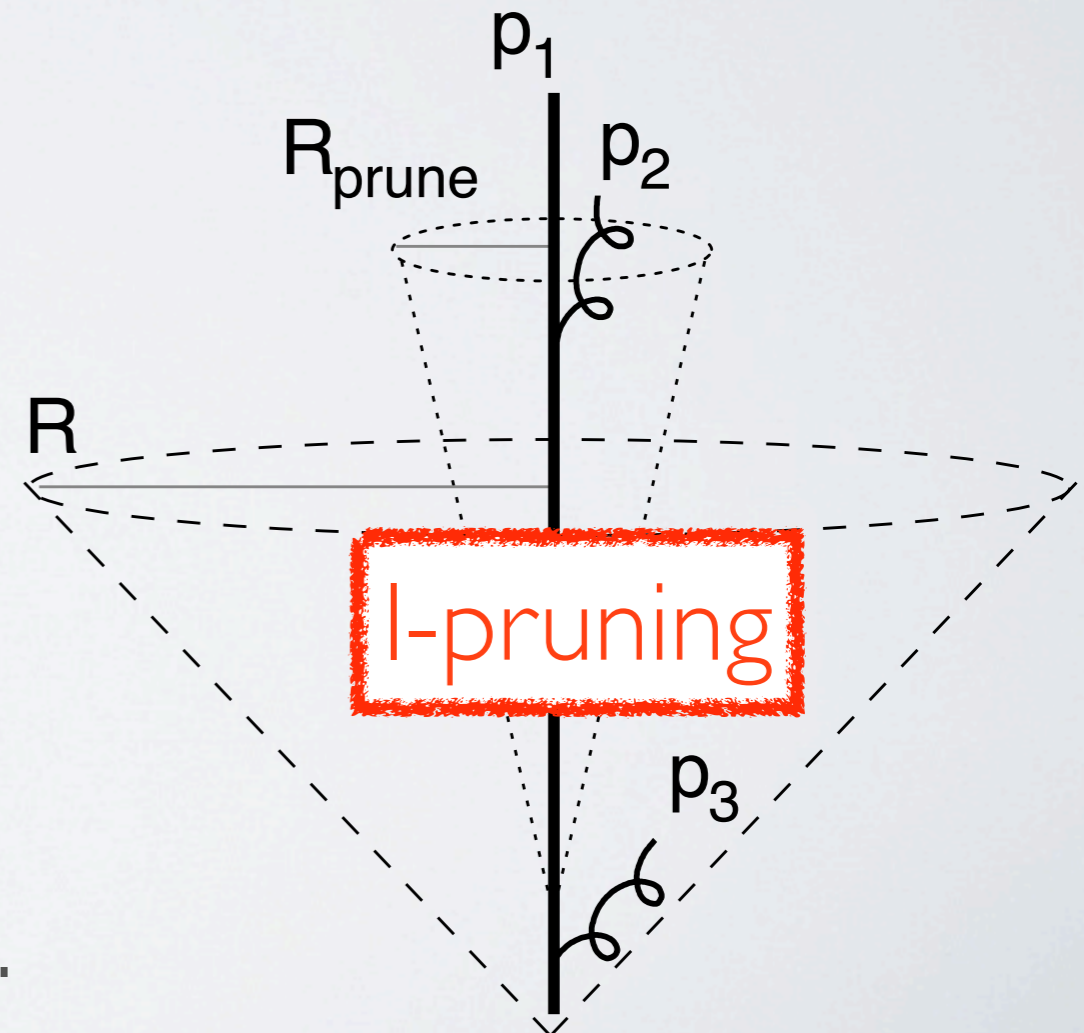
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# Structure beyond LO

- Pruning @ NLO  $\sim \alpha_s^2 L^4$  (like plain jet mass)
- Explicit calculation shows that I-pruning is active for  $\rho < z_{\text{cut}}^2$
- A simple modification: require at least one successful merging with  $\Delta R > R_{\text{prune}}$  and  $z > z_{\text{cut}}$  (Y-pruning)



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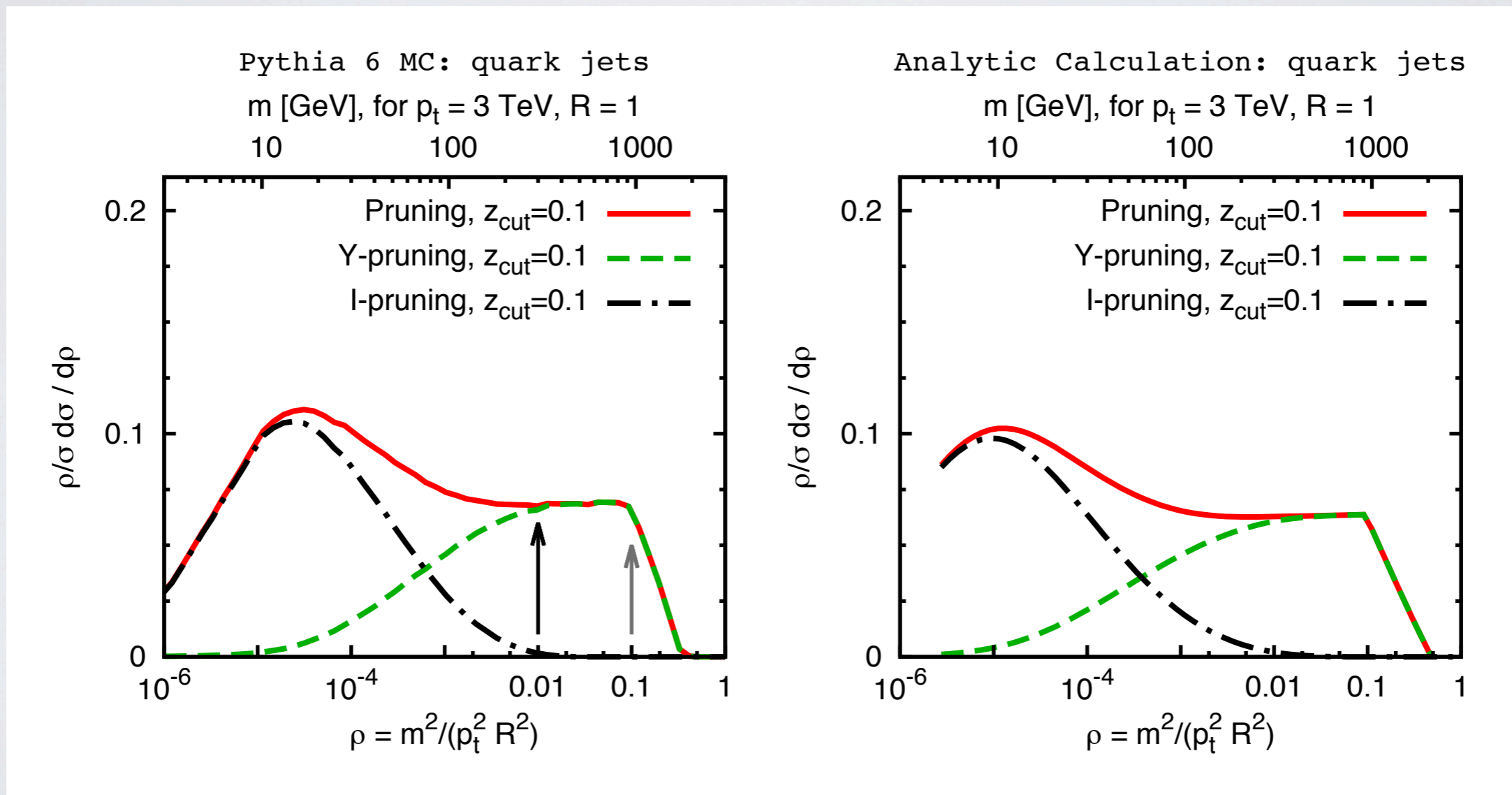
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- A simple modification: require at least one successful merging with  $\Delta R > R_{\text{prune}}$  and  $z > z_{\text{cut}}$  (**Y-pruning**)
  
- It is convenient to resum the two components separately
- **Y-pruning**: essentially Sudakov suppression of LO  $\sim \alpha_s^n L^{2n-1}$
- **I-pruning**: convolution between the pruned and the original mass  $\sim \alpha_s^n L^{2n}$

# All-order results

- Full Pruning: single-log region for  $z_{\text{cut}}^2 < \rho < z_{\text{cut}}$
- We control  $\alpha_s^n L^{2n}$  and  $\alpha_s^n L^{2n-1}$  in the expansion
- NG logs present but deferred to NNLO

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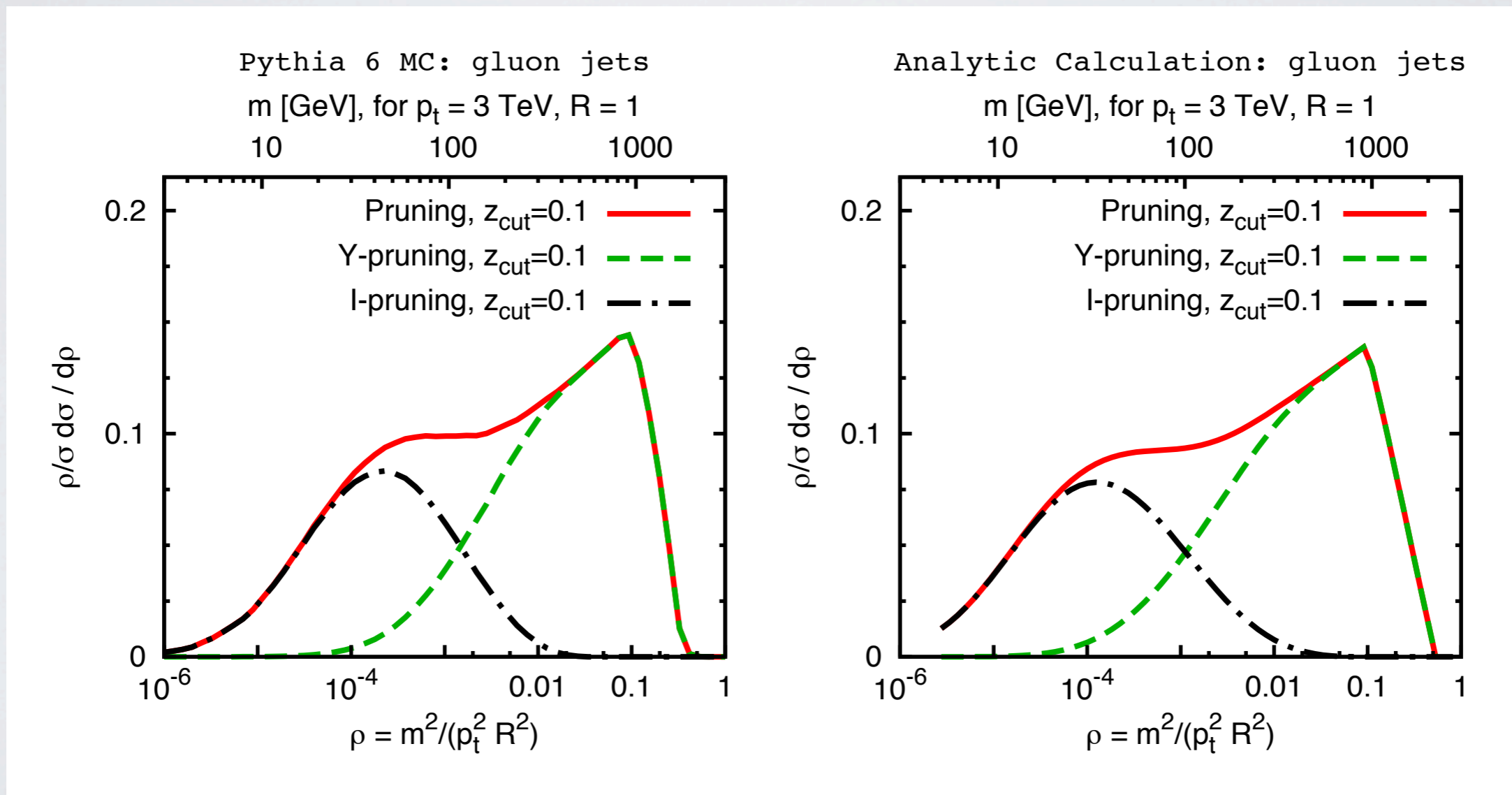
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# Mass Drop Tagger at LO

1. Undo the last stage of the C/A clustering. Label the two subjets  $j_1$  and  $j_2$  ( $m_1 > m_2$ )
2. If  $m_1 < \mu m$  (mass drop) and the splitting was not too asymmetric ( $y_{ij} > y_{\text{cut}}$ ), tag the jet.
3. Otherwise redefine  $j = j_1$  and iterate.

Butterworth, Davison, Rubin and Salam (2008)

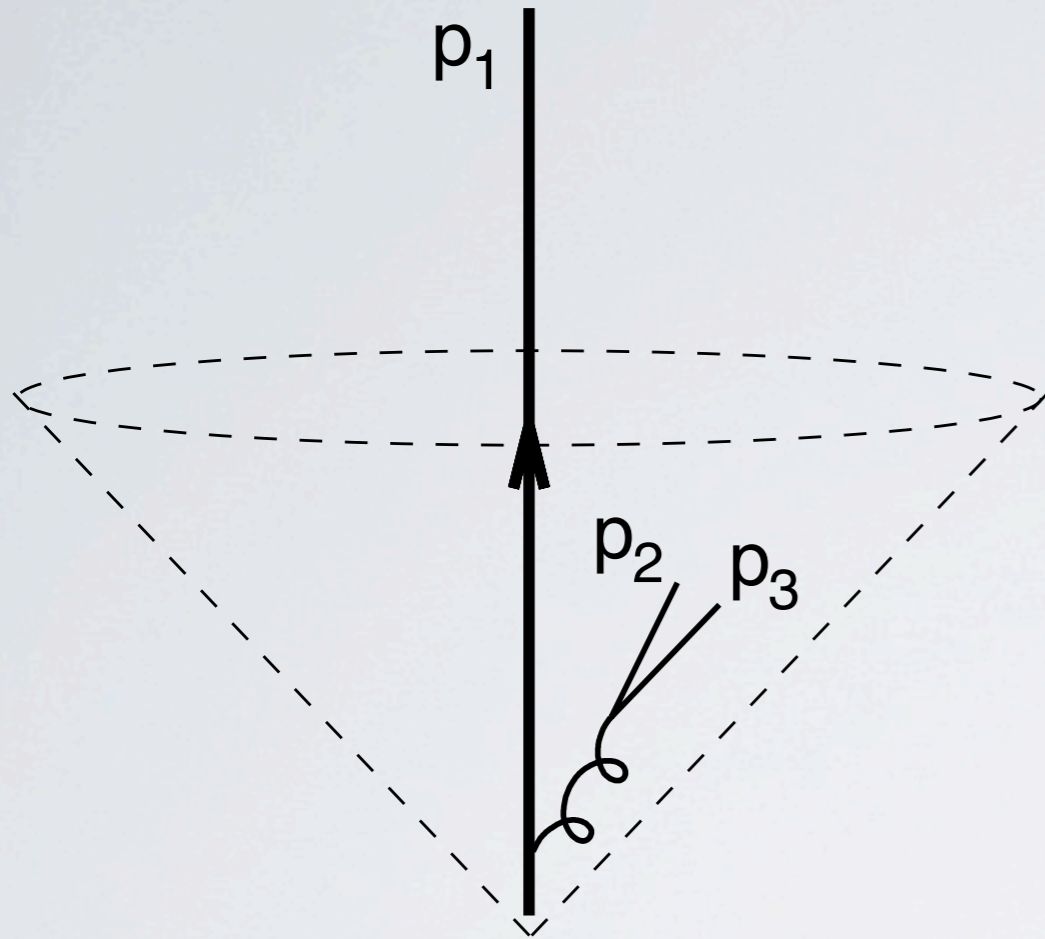
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In the small- $y_{\text{cut}}$  limit the result is identical to LO pruning:  
single-log distribution

# Problems beyond LO



## What MDT does wrong:

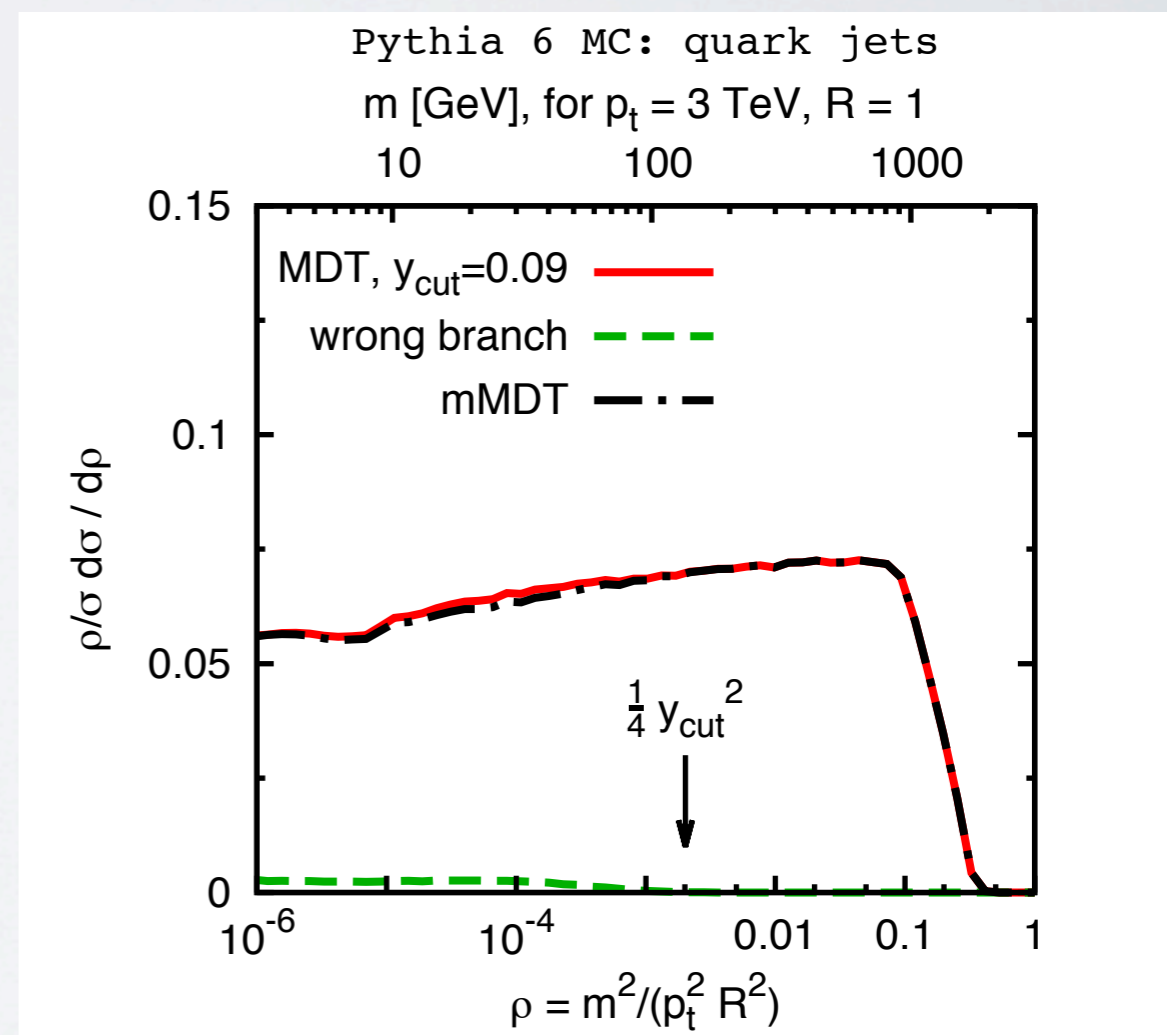
If the  $y_{ij}$  condition fails, MDT iterates on the more massive subjet. It can follow a soft branch ( $p_2 + p_3 < y_{\text{cut}} p_{\text{tjet}}$ ), when the “right” answer was that the (massless) hard branch had no substructure

- This can be considered a flaw of the tagger
- It worsens the logarithmic structure:  $\alpha_s^2 L^3$
- It makes all-order treatment difficult
- It calls for a modification

# Modified Mass Drop Tagger

1. Undo the last stage of the C/A clustering.  
Label the two subjets  $j_1$  and  $j_2$  ( $m_1 > m_2$ )
2. If  $m_1 < \mu m$  (mass drop) and the splitting was not too asymmetric ( $y_{ij} > y_{\text{cut}}$ ), tag the jet.
3. Otherwise redefine  $j$  to be the subjet with highest transverse mass and iterate.

- In practice the soft-branch contribution is very small
- However, this modification makes the all-order structure particularly interesting





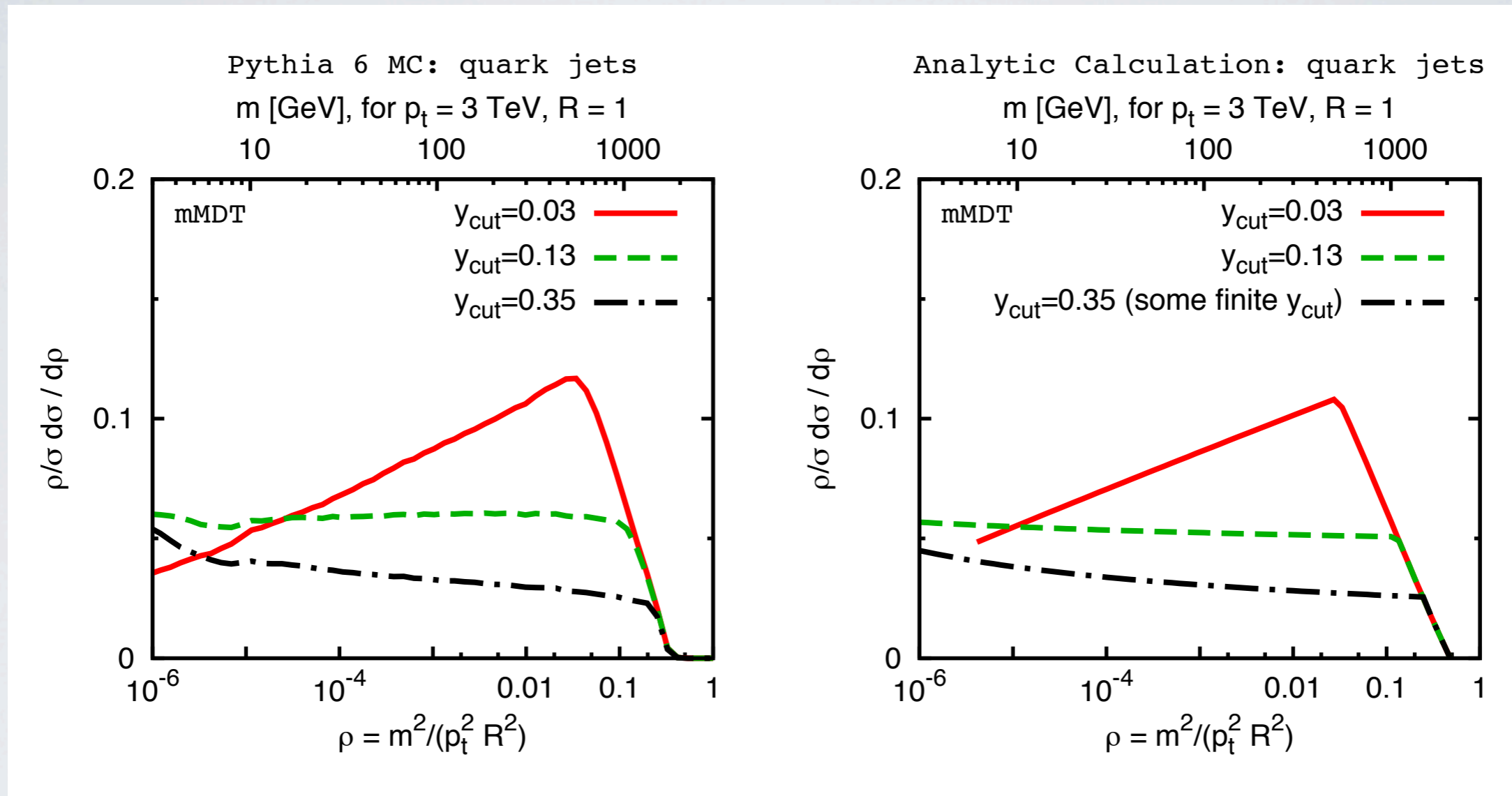
# All-order structure of mMDT

- In the small  $y_{\text{cut}}$  limit, it is just the exponentiation of LO
- The mMDT has **single logs to all orders** (i.e.  $\alpha_s^n L^n$ )
- This potentially extends the validity of FO calculations

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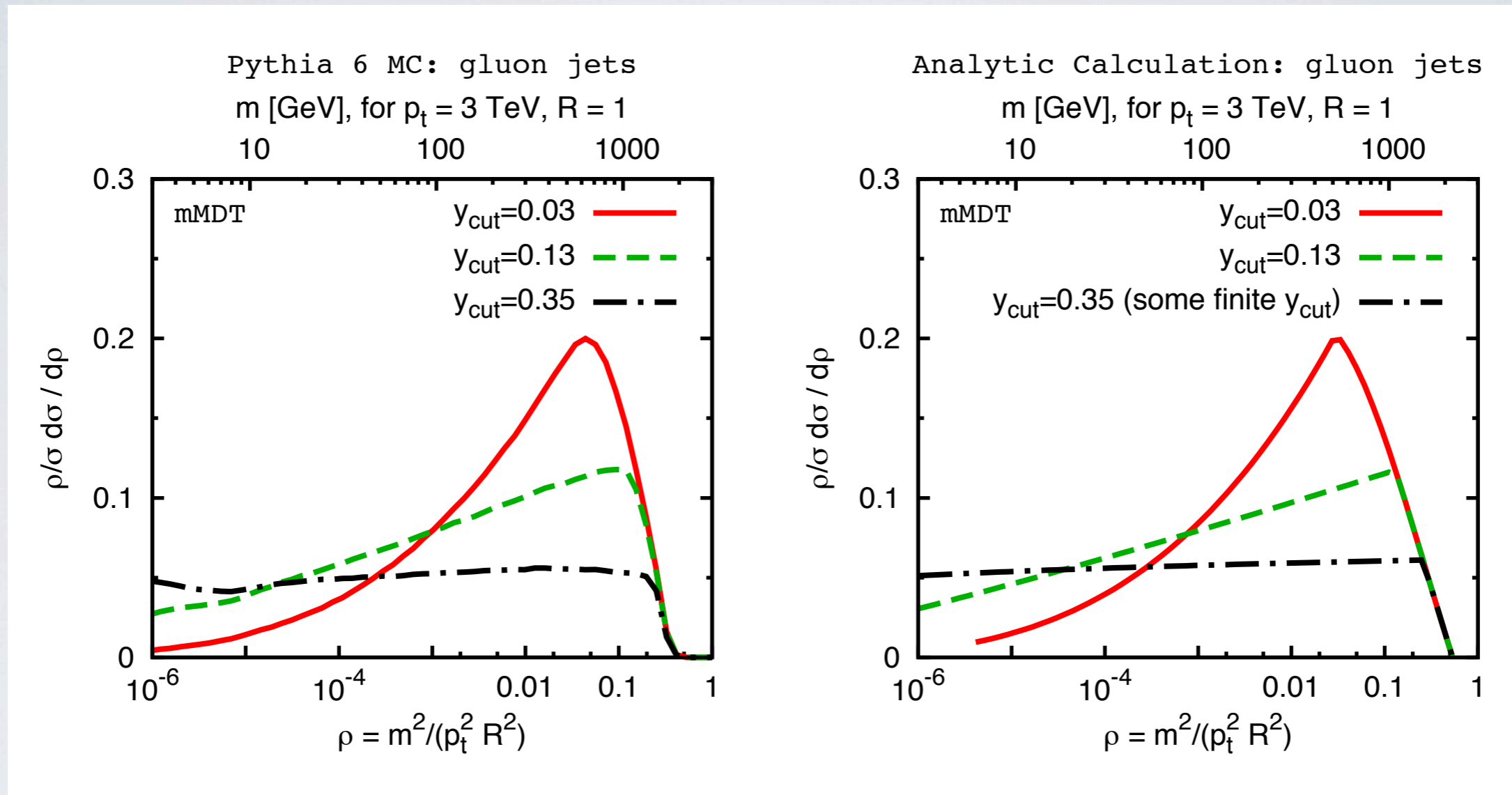
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- The mMDT has **single logs to all orders** (i.e.  $\alpha_s^n L^n$ )
- This potentially extends the validity of FO calculations
  - These single logs are of collinear origin
  - Remarkable consequence: mMDT is free of non-global logs!

# Comparison to MC



Remarkable agreement !

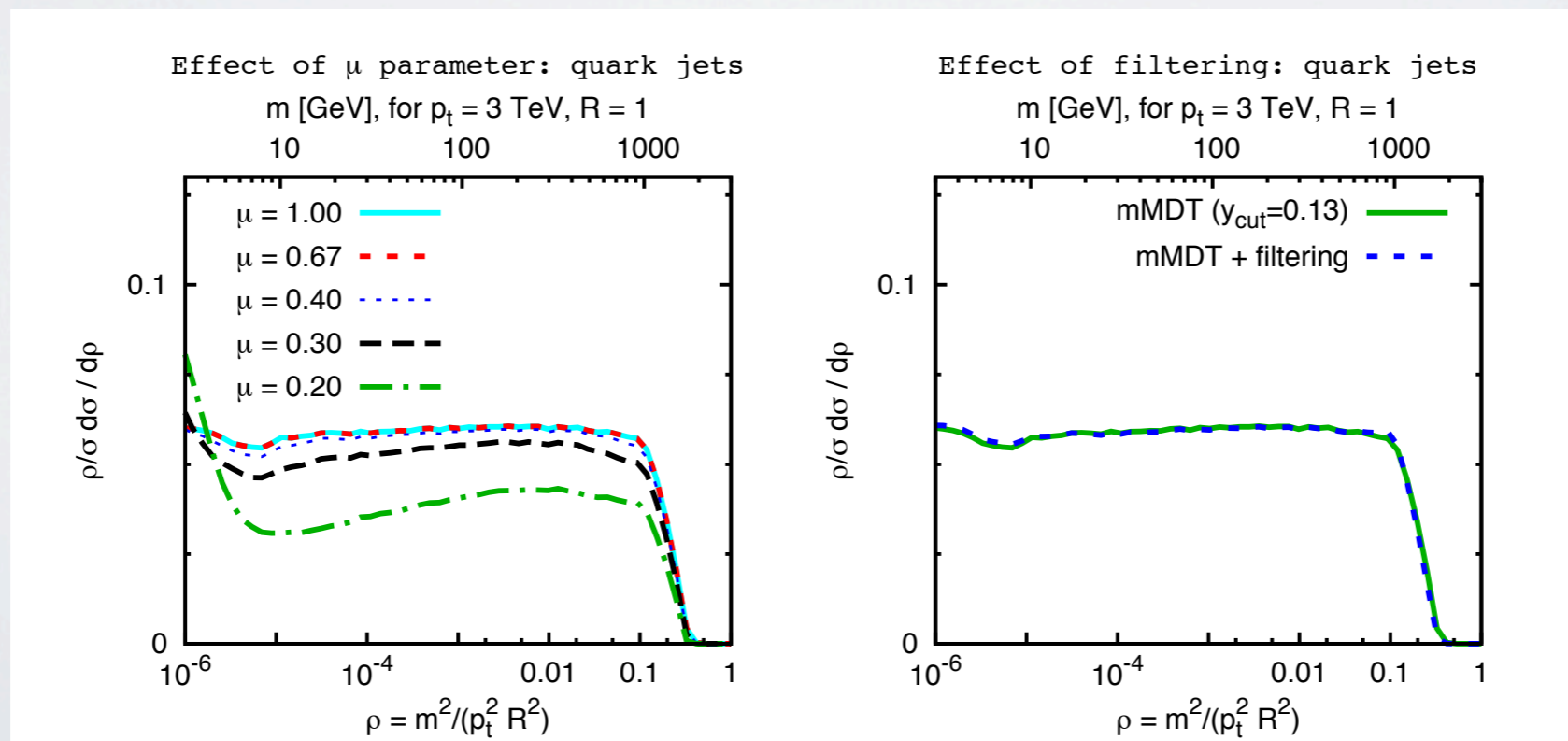
# Comparison to MC



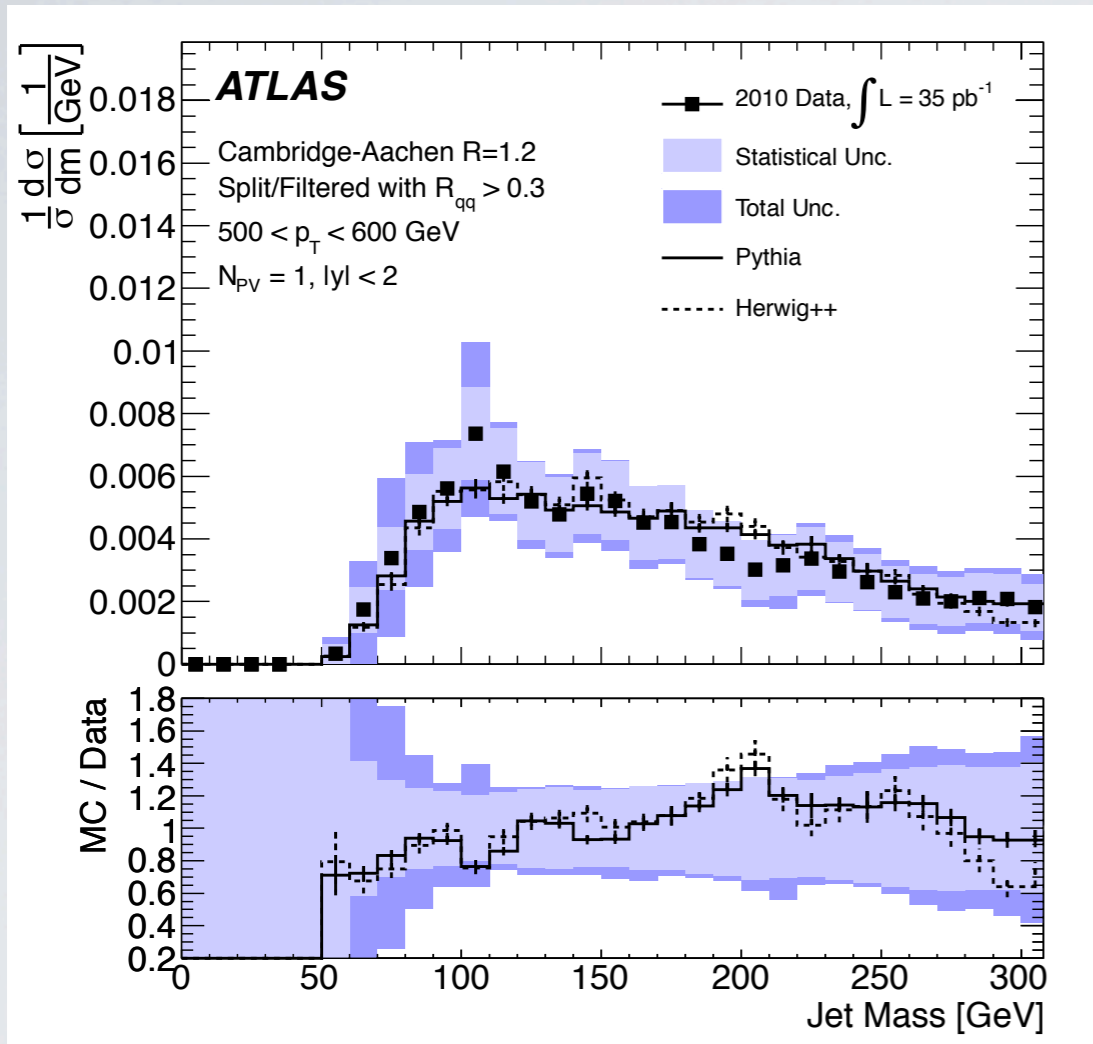
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# Other properties of mMDT

- Flatness of the background is a desirable property (data-driven analysis, side bands)
- $y_{\text{cut}}$  can be adjusted to obtain it (analytic relation)
- Role of  $\mu$ , not mentioned so far
- It contributes to subleading logs and has small impact if not too small ( $\mu > 0.4$ )
- Filtering only affects subleading ( $N^{\text{filt}}\text{LL}$ ) terms

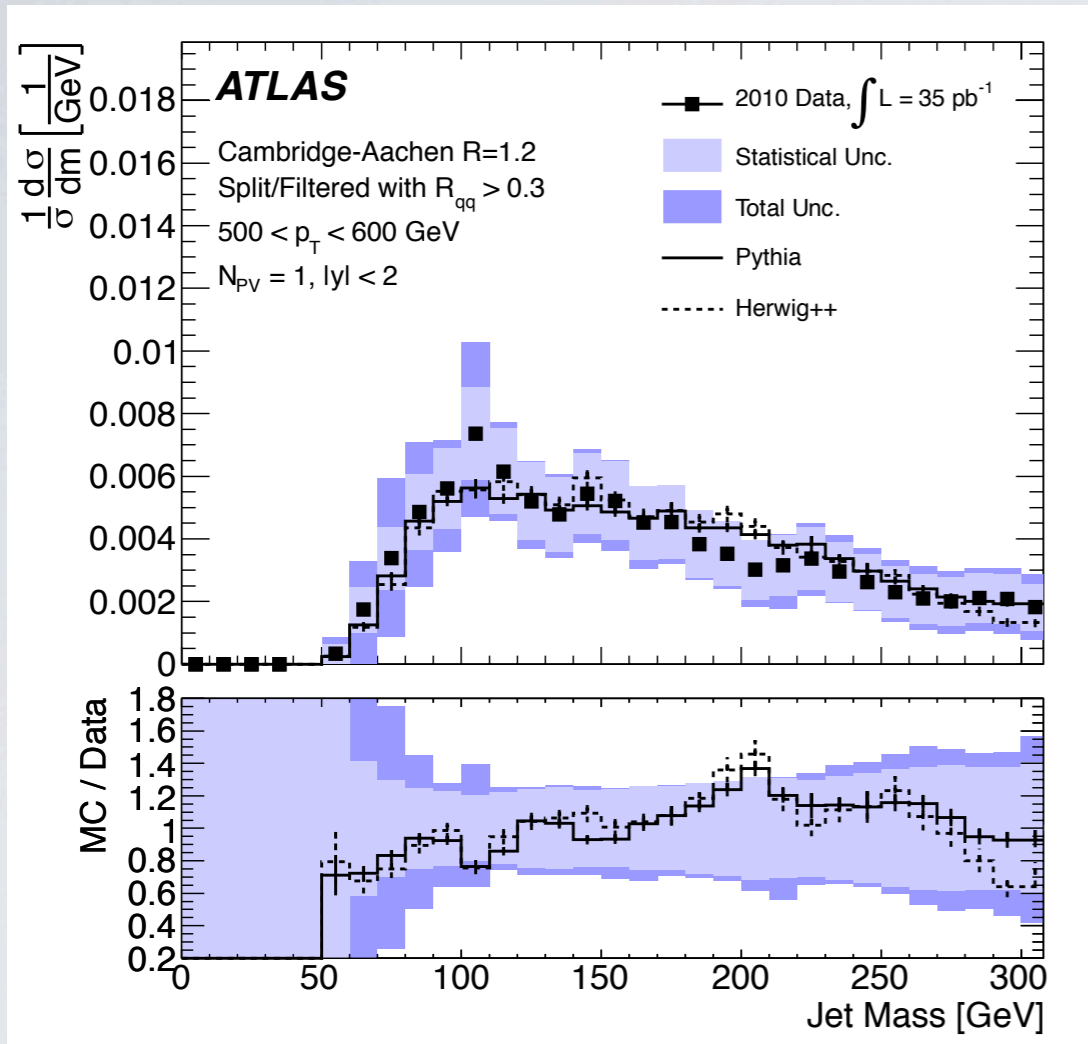


# ATLAS MDT



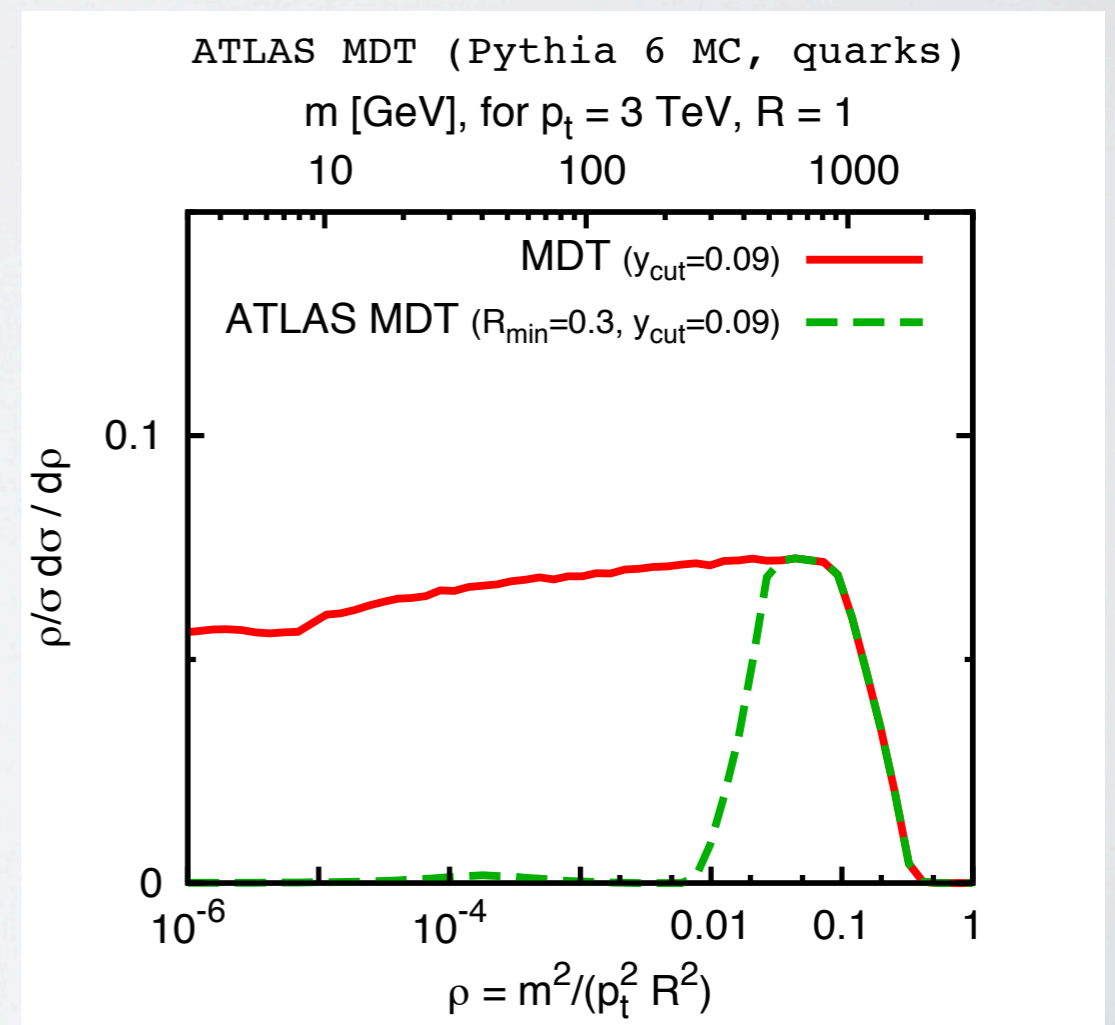
- ATLAS measured the jet mass with MDT
- Different version of the tagger with  $R_{\min}=0.3$  between the prongs

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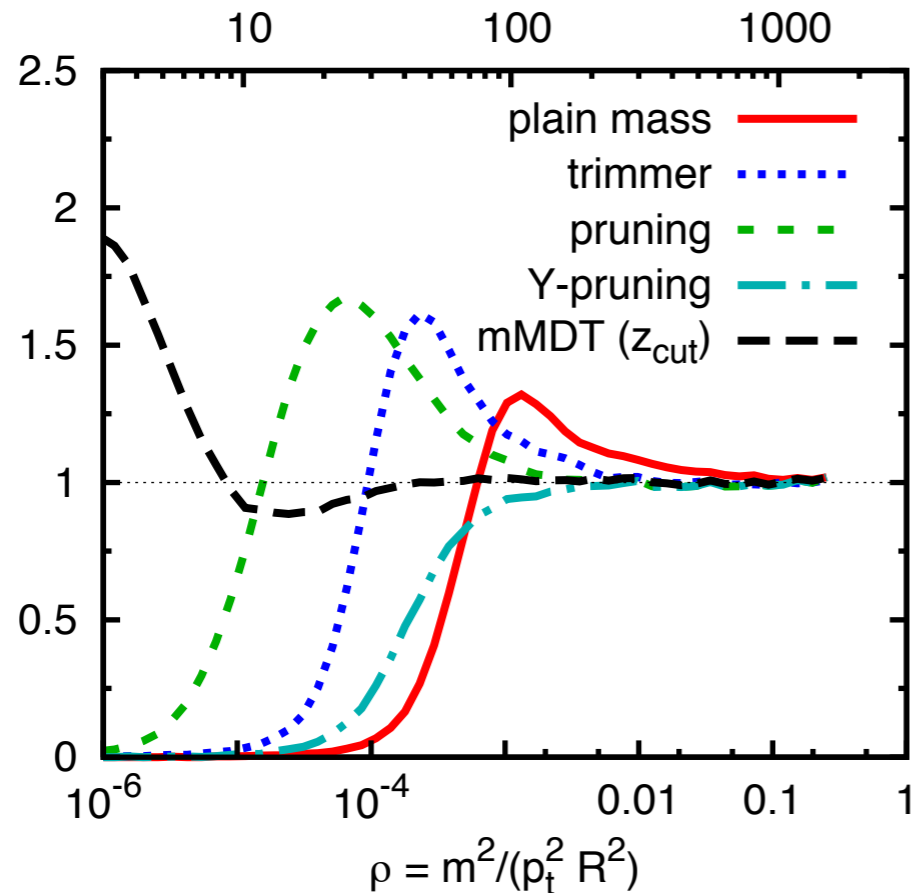
- ATLAS measured the jet mass with MDT
- Different version of the tagger with  $R_{\min}=0.3$  between the prongs

- This cut significantly changes the tagger's behaviour: mass minimum
- The single-log region is reduced (and can even disappear)
- We hope that future studies will be able to avoid this

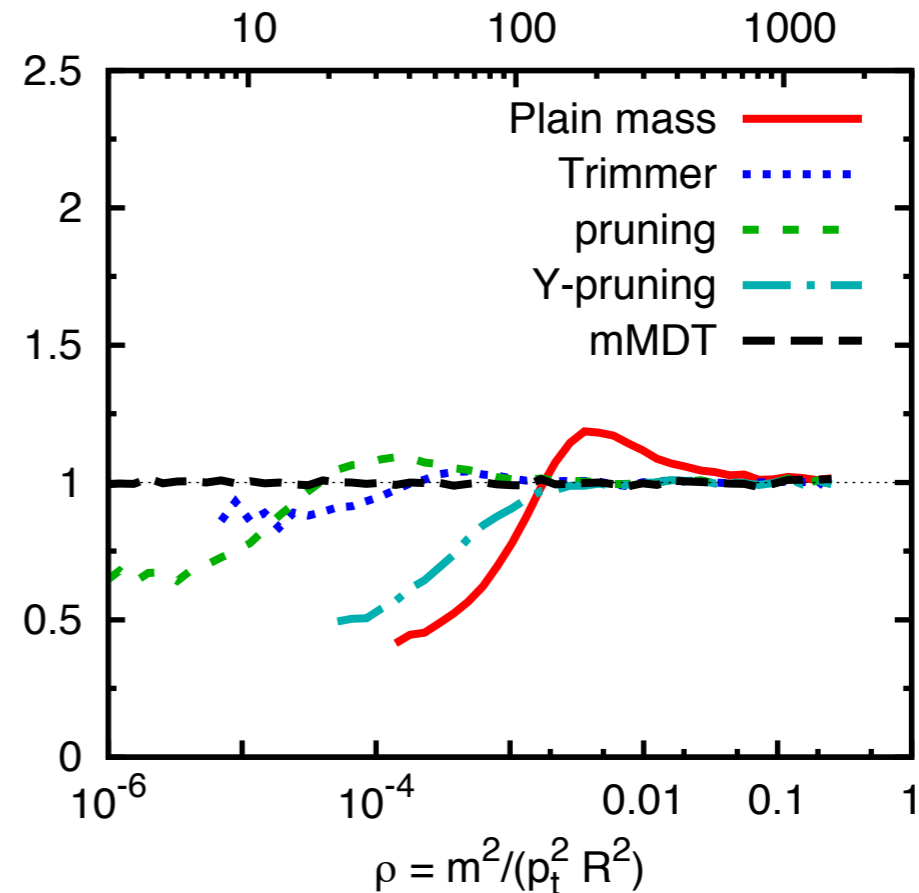


# Non-perturbative effects

## hadron / parton



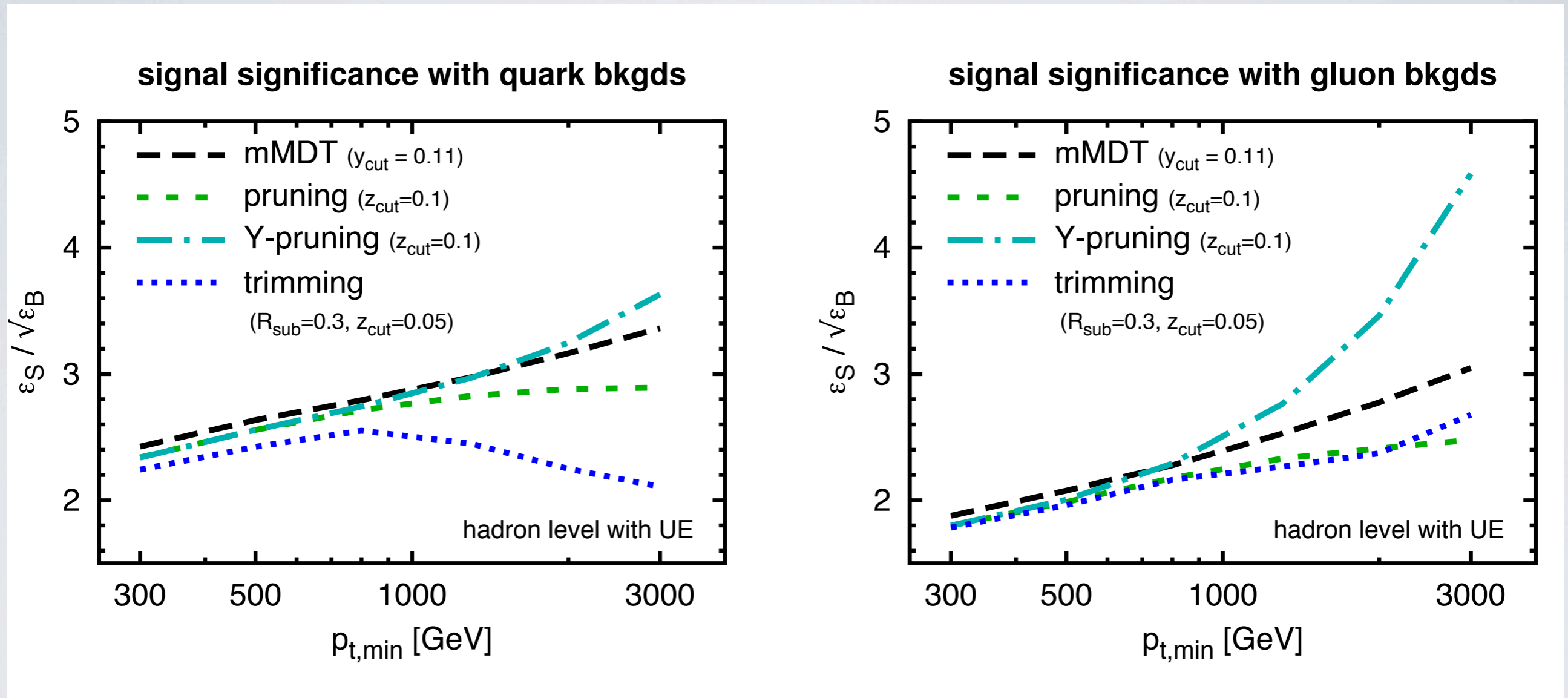
## wUE / noUE



- Most taggers have reduced sensitivity to NP physics
- mMDT particularly so (it's the most calculable)
- Y-pruning sensitive to UE because of the role played by the fat jet mass



# Performances for finding signals (Ws)



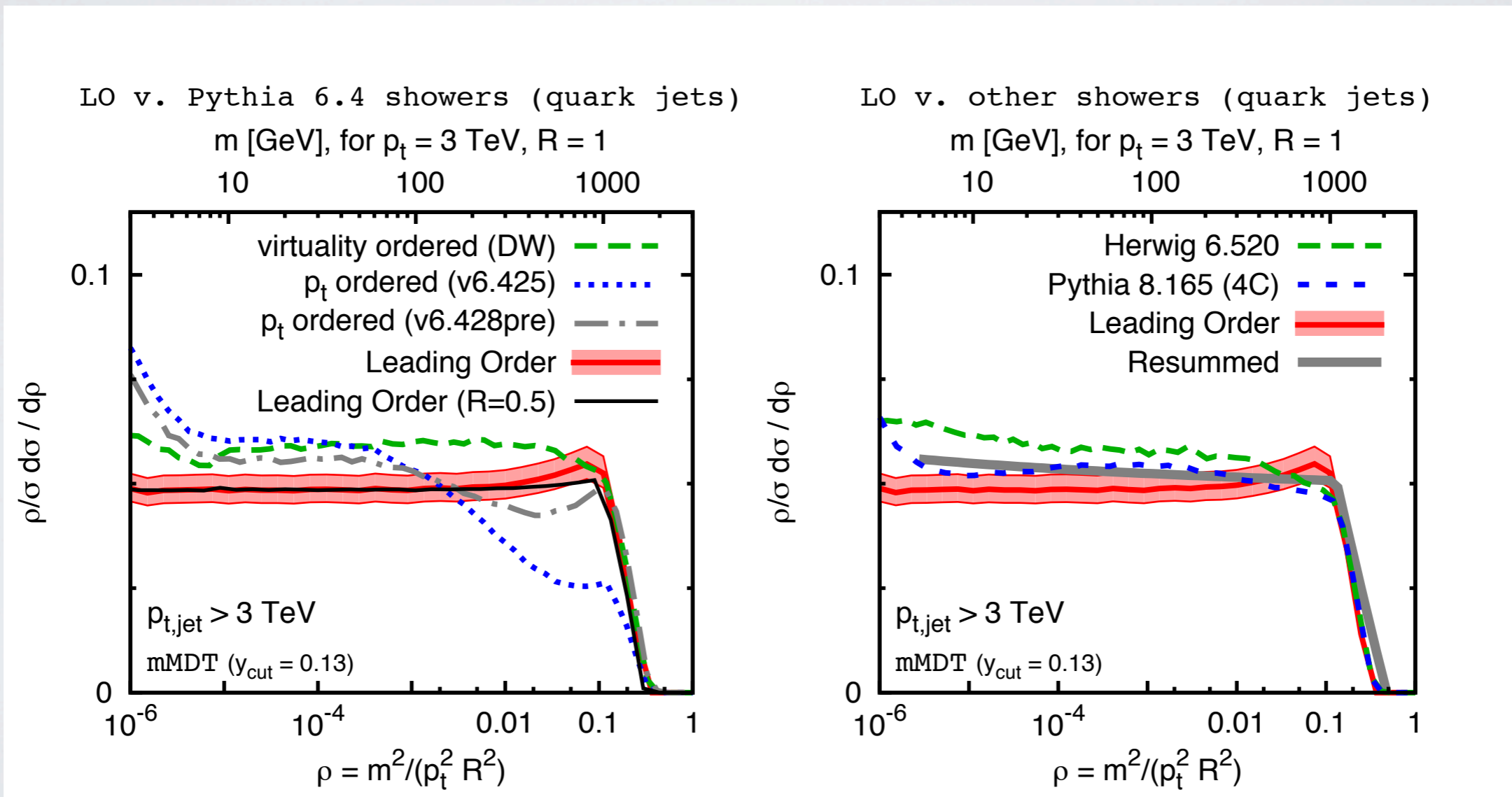
Y-pruning gives a visible improvement

# In summary ...

- Analytic studies of the taggers reveal their properties
- Particularly useful if MCs don't agree

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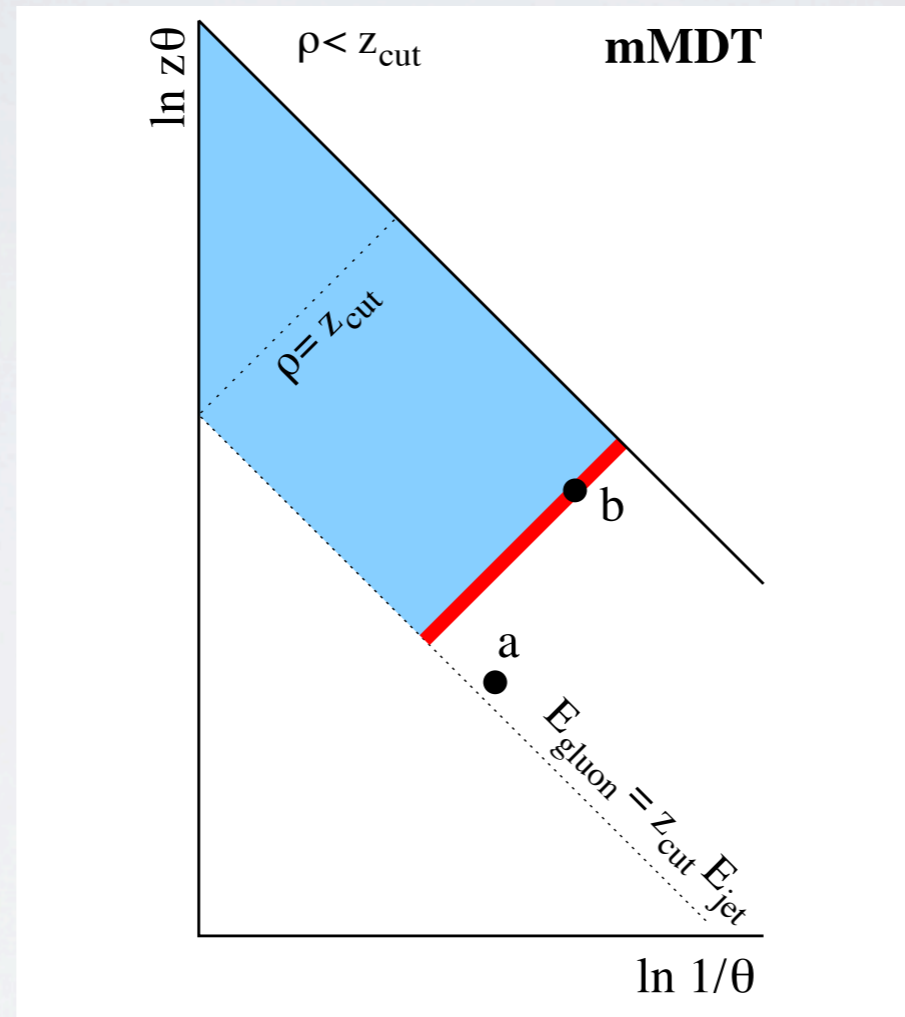
- Analytic studies of the taggers reveal their properties
- Particularly useful if MCs don't agree
- They also indicate how to develop better taggers
- Y-pruning:
  - improved log behaviour wrt pruning ( $\alpha_s^n L^{2n-1}$  vs  $\alpha_s^n L^{2n}$ )
  - better rejection of QCD background
- mMDT:
  - exceptionally simple structure (single logs, no non-global)
  - reduced sensitivity to non-perturbative physics

BACKUP SLIDES

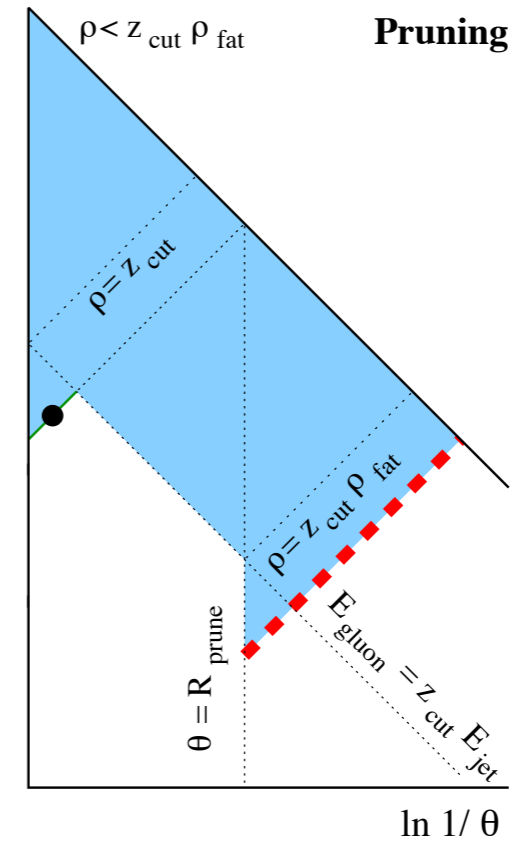
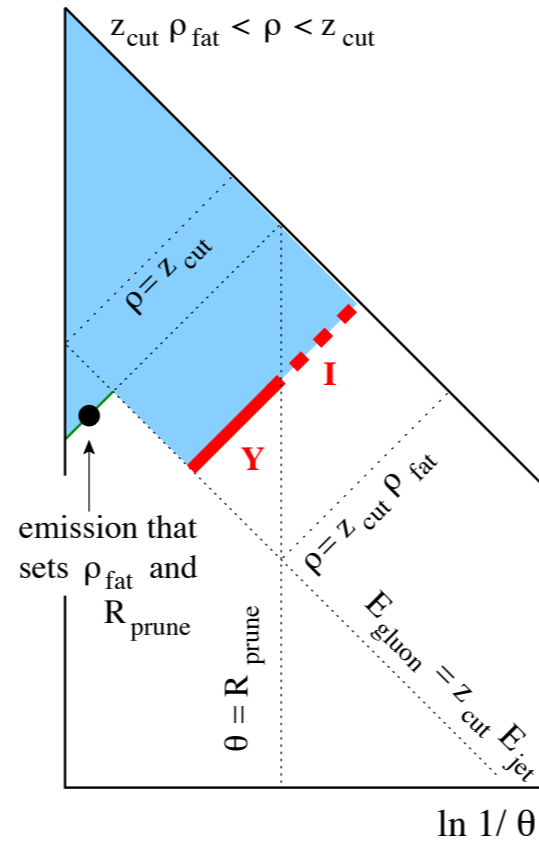
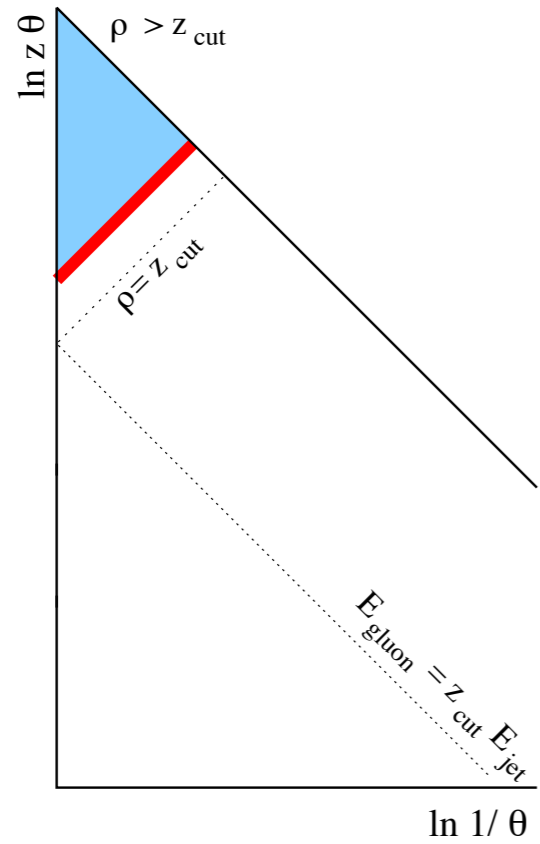
# Summary table

	highest logs	transition(s)	Sudakov peak	NGLs	NP: $m^2 \lesssim$
plain mass	$\alpha_s^n L^{2n}$	—	$L \simeq 1/\sqrt{\bar{\alpha}_s}$	yes	$\mu_{\text{NP}} p_t R$
trimming	$\alpha_s^n L^{2n}$	$z_{\text{cut}}, r^2 z_{\text{cut}}$	$L \simeq 1/\sqrt{\bar{\alpha}_s} - 2 \ln r$	yes	$\mu_{\text{NP}} p_t R_{\text{sub}}$
pruning	$\alpha_s^n L^{2n}$	$z_{\text{cut}}, z_{\text{cut}}^2$	$L \simeq 2.3/\sqrt{\bar{\alpha}_s}$	yes	$\mu_{\text{NP}} p_t R$
MDT	$\alpha_s^n L^{2n-1}$	$y_{\text{cut}}, \frac{1}{4}y_{\text{cut}}^2, y_{\text{cut}}^3$	—	yes	$\mu_{\text{NP}} p_t R$
Y-pruning	$\alpha_s^n L^{2n-1}$	$z_{\text{cut}}$	(Sudakov tail)	yes	$\mu_{\text{NP}} p_t R$
mMDT	$\alpha_s^n L^n$	$y_{\text{cut}}$	—	no	$\mu_{\text{NP}}^2 / y_{\text{cut}}$

# Lund diagrams for mMDT

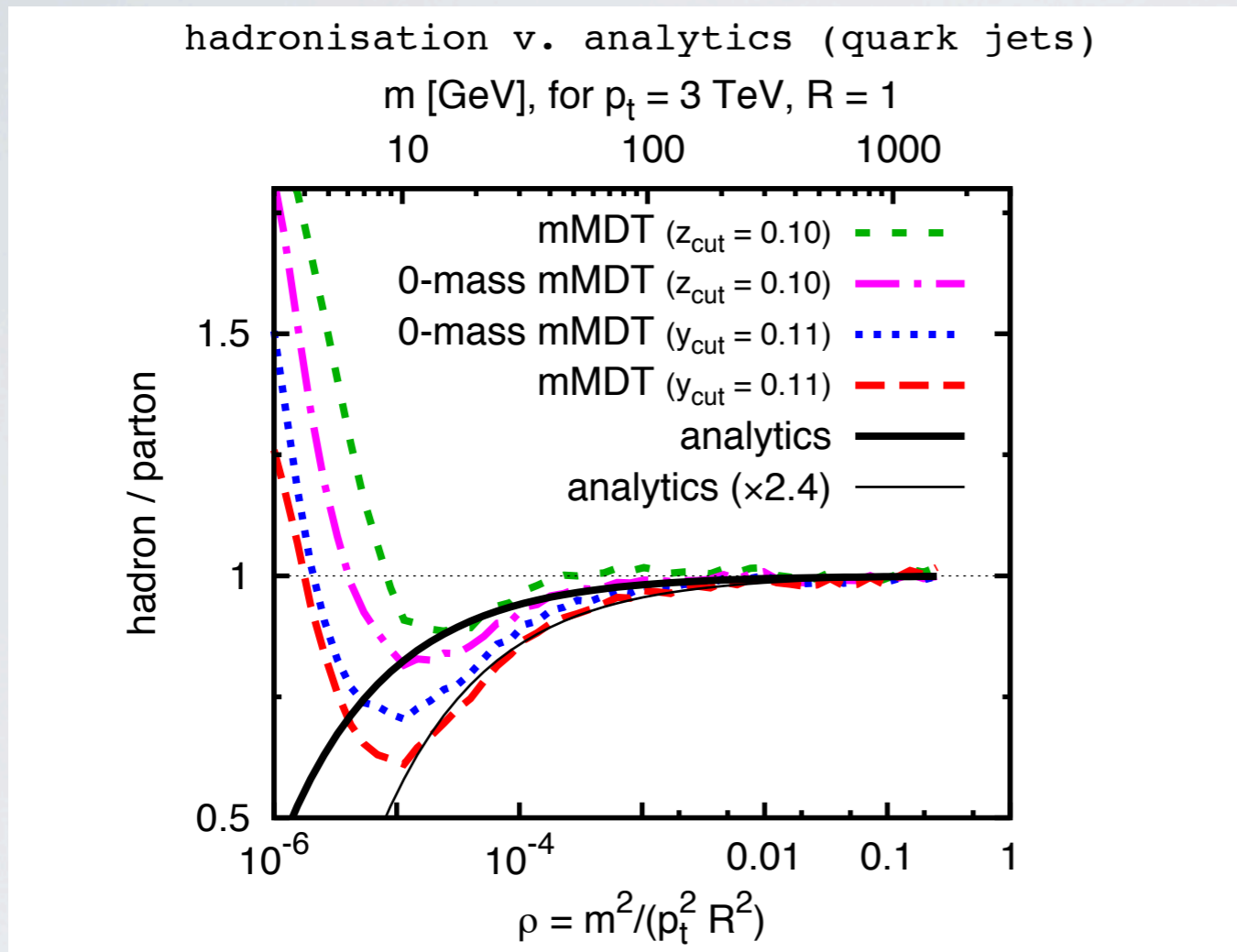


# Lund diagrams for pruning





# Hadronisation effects for mMDT



- Hadronisation produces:
1. a shift in the squared jet mass
  2. a shift in the jet's (or prong's) momentum

Same power behaviour but with competing signs:

$$\frac{d\sigma^{\text{NP}}}{dm} \simeq \frac{d\sigma^{\text{PT}}}{dm} \left[ 1 + a \frac{\Lambda_{\text{NP}}}{m} \right]$$

# Examples of NLO checks

