

# Electroweak penguins with di-leptons

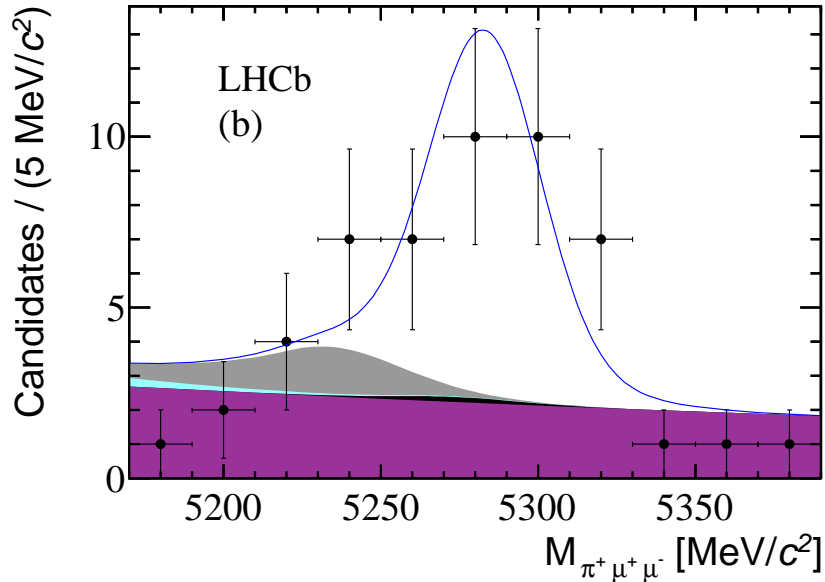
M. Kreps on behalf of the LHCb Collaboration

Physics Department



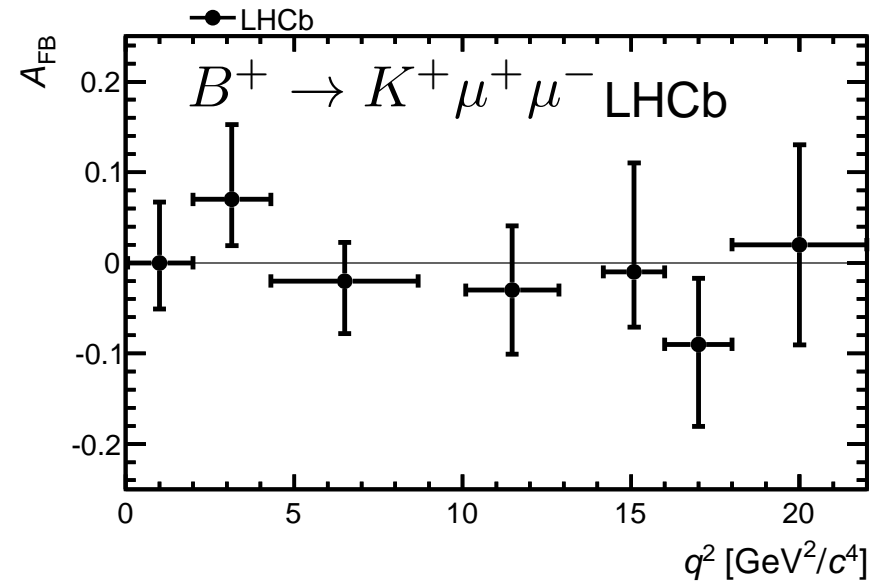
# Reminder of recent results

LHCb-PAPER-2012-020

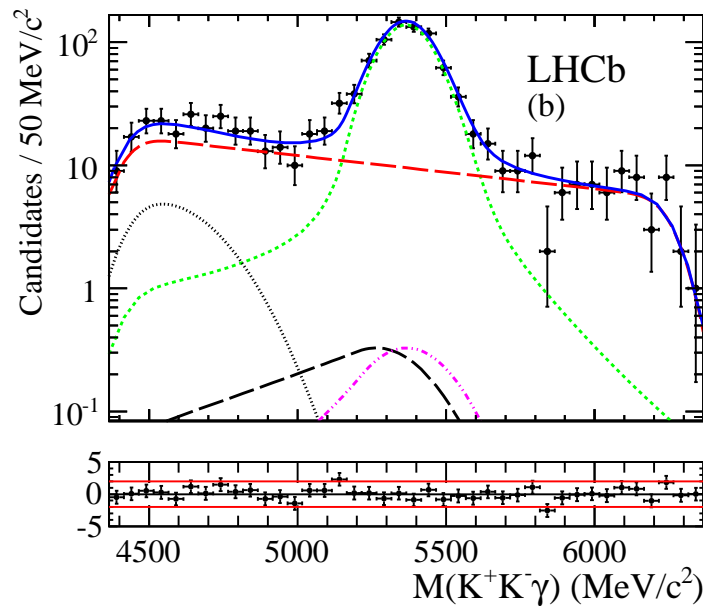


$B^+ \rightarrow \pi^+ \mu^+ \mu^-$

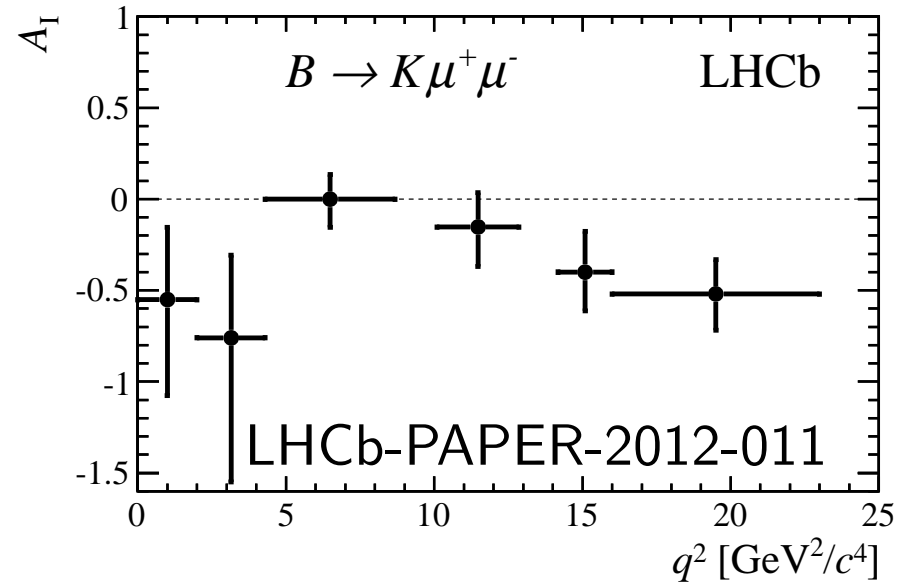
LHCb-PAPER-2012-024



$B_s^0 \rightarrow \phi \gamma$



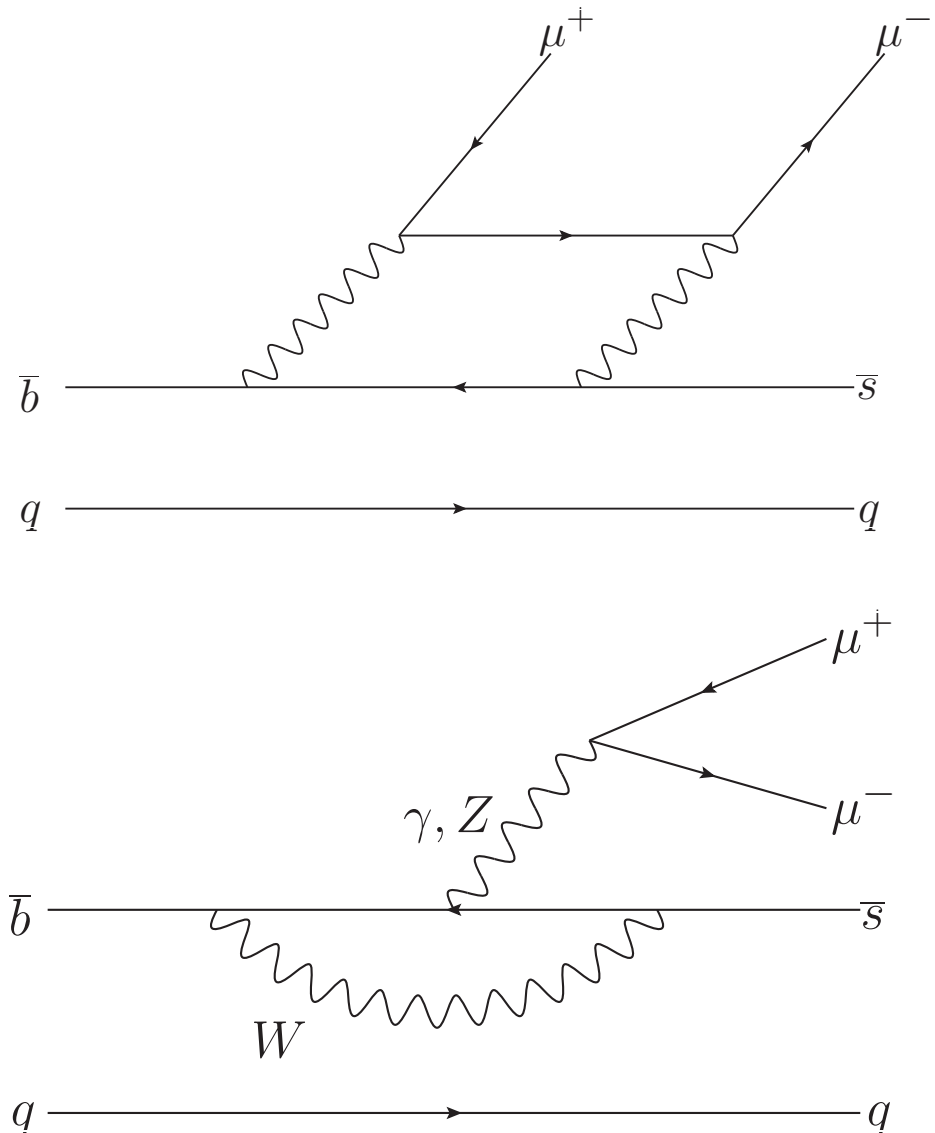
LHCb-PAPER-2012-019



LHCb-PAPER-2012-011

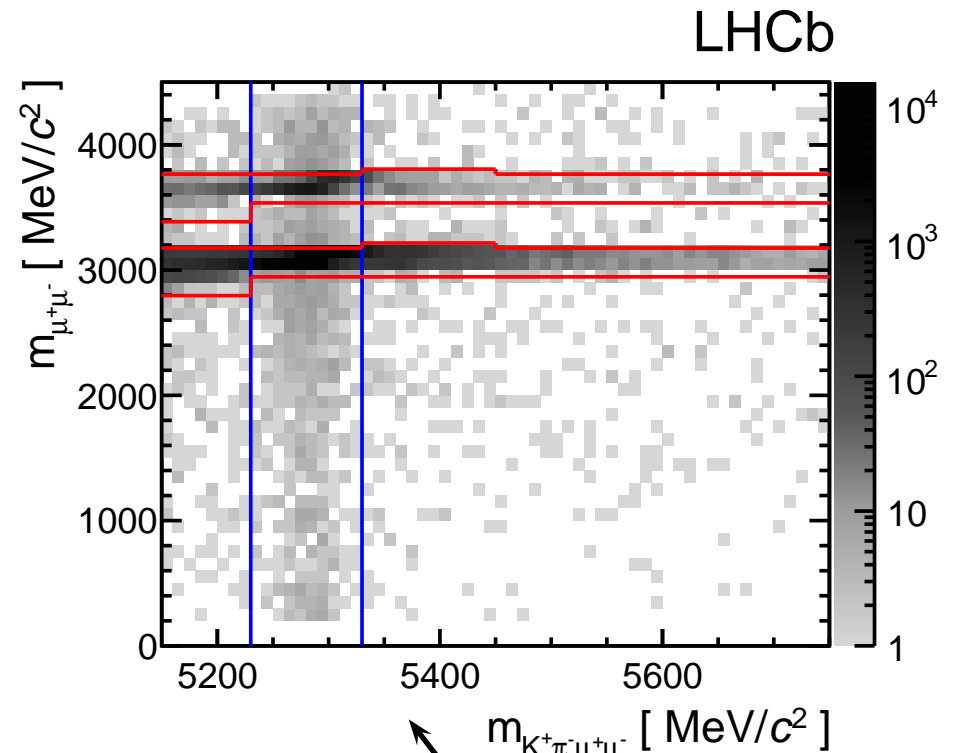
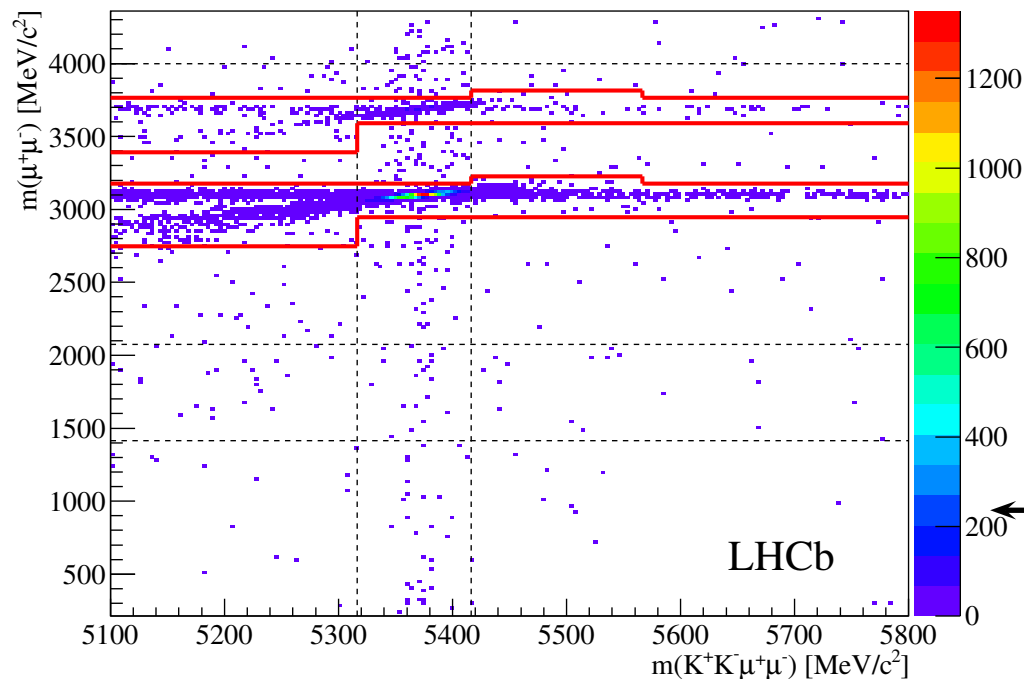
# Motivation for EWK penguins

- In the SM,  $b \rightarrow s$  FCNC decays
- SM BF of the order of  $10^{-6}$
- With angular analysis offers variety of observables
- Allows to test some underlying details of the NP
- Form factors make prediction of some observables less precise
- But many observables are free of form factor uncertainties
- Today, new results on
  - $B^0 \rightarrow K^{*0} \mu^+ \mu^-$
  - $B^0 \rightarrow K^{*0} e^+ e^-$
  - $B_s^0 \rightarrow \phi \mu^+ \mu^-$
- All use  $1 \text{ fb}^{-1}$  of data from 2011
- All results preliminary



# Selection with dimuons

- Uses BDT to combine kinematic, topological and PID inputs
- Trained on resonant ( $J/\psi$ ) signal in data
- Relatively complicated removal of  $c\bar{c}$  regions
- Decays via  $J/\psi$  used for normalization



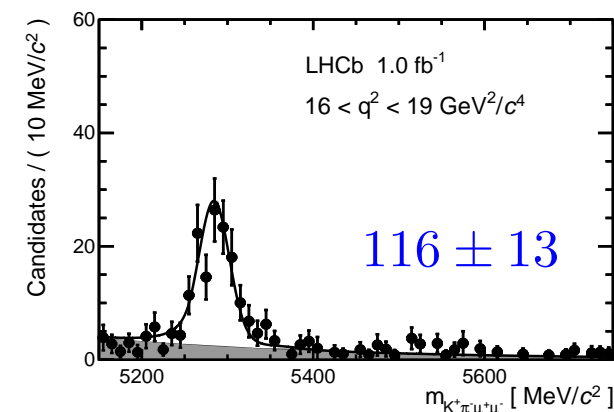
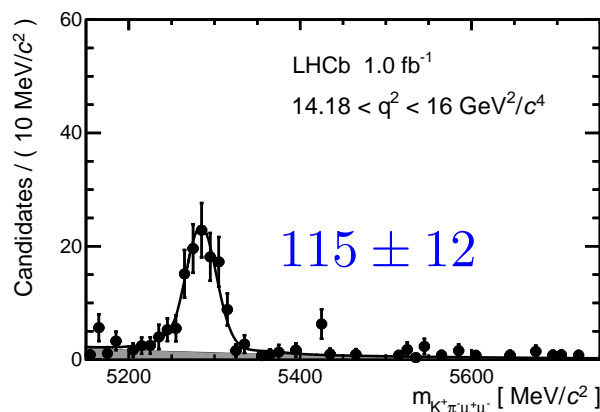
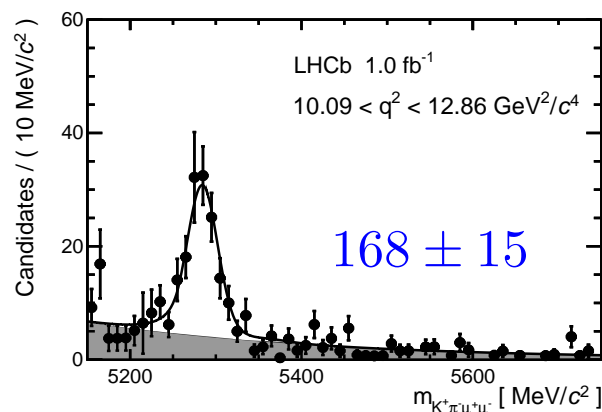
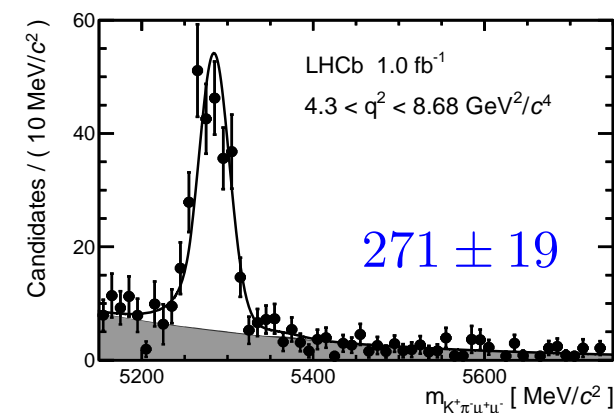
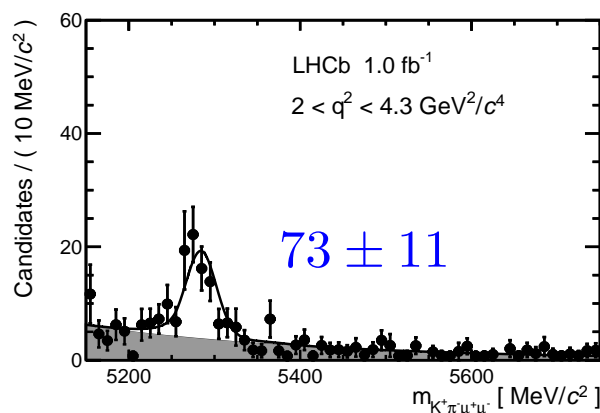
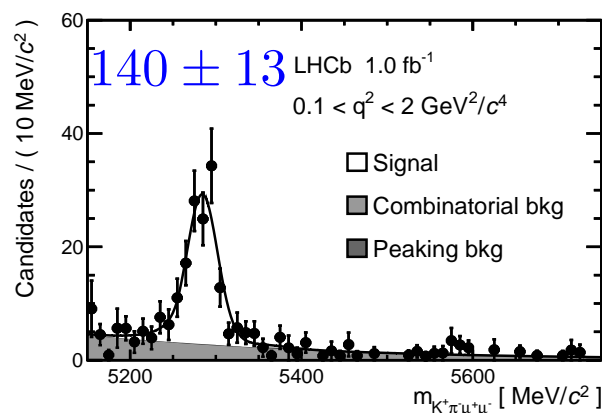
$B^0 \rightarrow K^{*0} \mu^+ \mu^-$   
LHCb-PAPER-2013-019

$B_s^0 \rightarrow \phi \mu^+ \mu^-$   
LHCb-PAPER-2013-017

# $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ — $q^2$ bins

- Extended maximum likelihood fit
- Signal: double Crystal ball function fixed from  $J/\psi$  decays
- Bg: combinatorial (exponential),  $B_s^0 \rightarrow K^{*0} \mu\mu$ ,  $B_s^0 \rightarrow \phi\mu\mu$  and  $B^0 \rightarrow J/\psi K^{*0}$  from MC

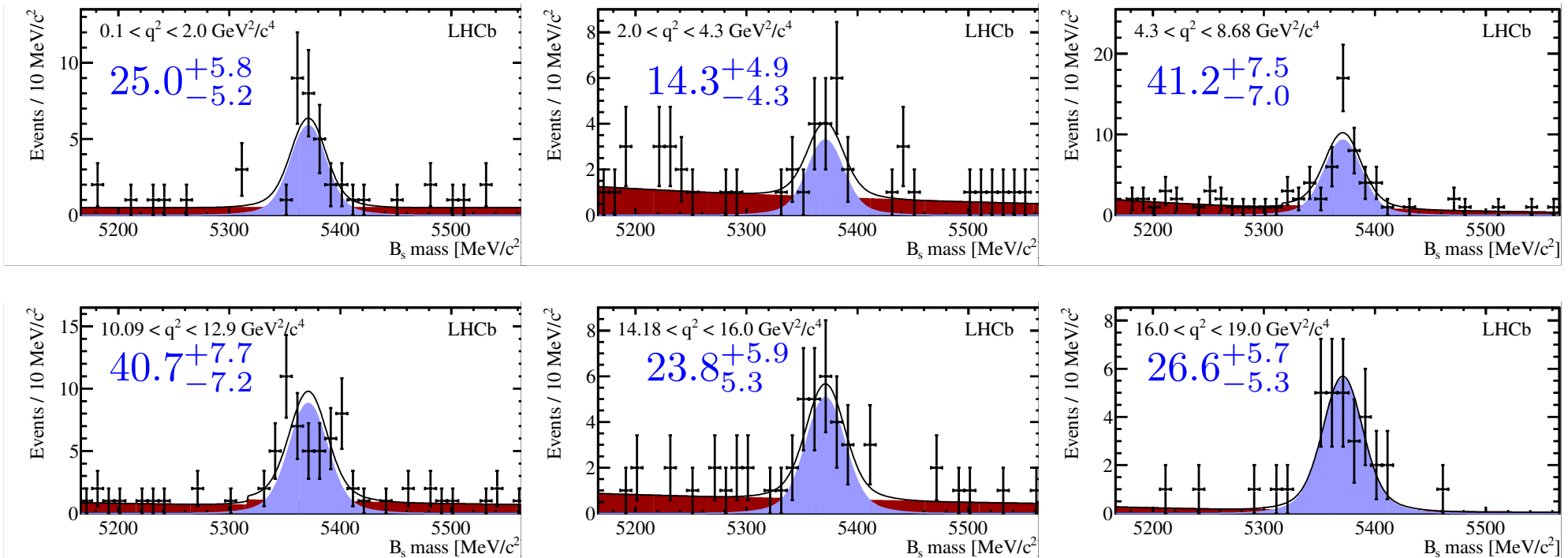
LHCb-PAPER-2013-019



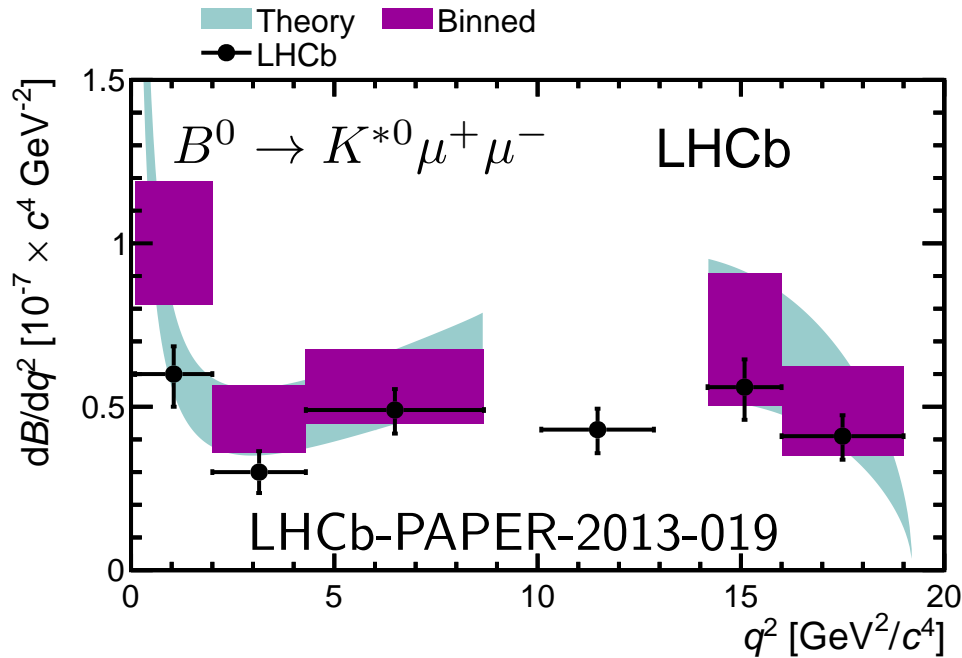
# $B_s^0 \rightarrow \phi \mu^+ \mu^-$ — $q^2$ bins

- Extended maximum likelihood fit
- Signal: double Gaussian function fixed from  $J/\psi$  decays
- Bg: combinatorial (exponential)
- Clear signal in all  $q^2$  bins
- First time we see signal in all  $q^2$  bins

LHCb-PAPER-2013-017

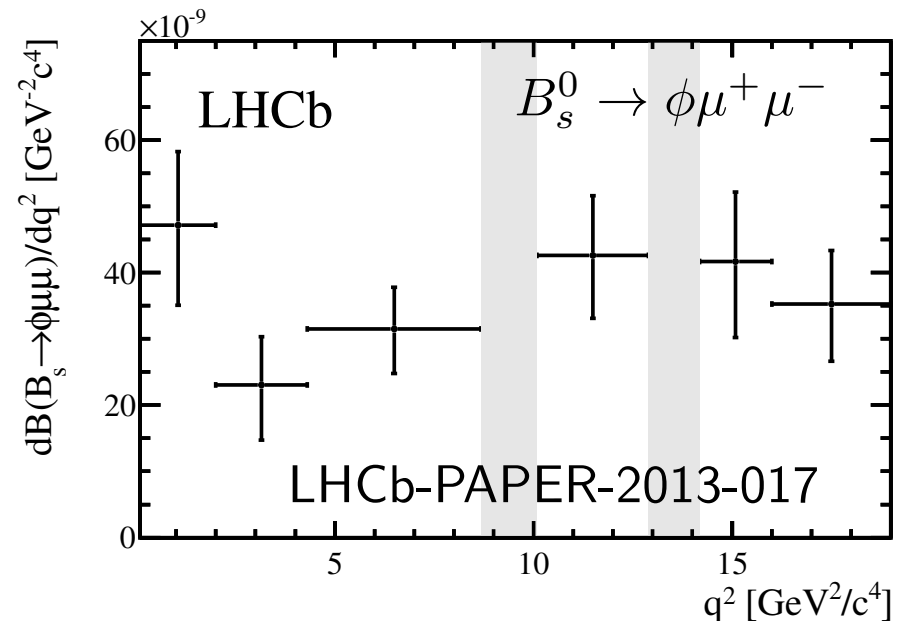


# Differential BF

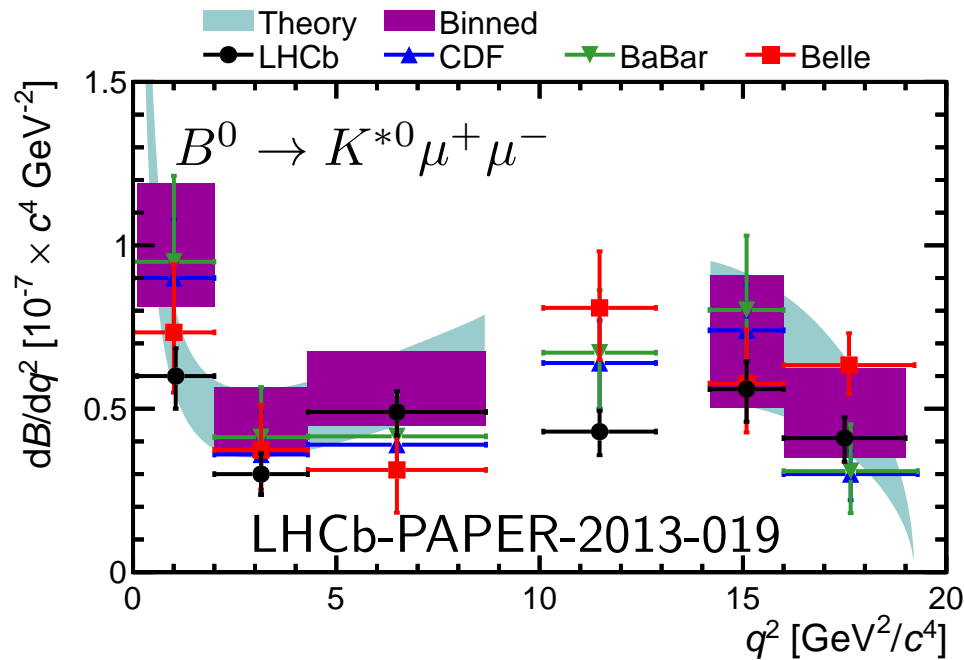


- $$dB/dq^2 = \frac{1}{\Delta q^2} \frac{N_{sig}}{N_{norm}} \epsilon_{rel} B_{norm}$$
- $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  agrees with SM
- Two decays have similar shape
- For  $B_s^0 \rightarrow \phi \mu \mu$  we measure integrated rate  
 $(7.07_{-0.59}^{+0.64} \pm 0.17 \pm 0.71) \times 10^{-7}$

- Syst. uncertainties dominated by normalization BF
- Other contributions from Bg PDFs and data-MC differences
- Large fraction of systematic uncertainty correlated across  $q^2$  bins

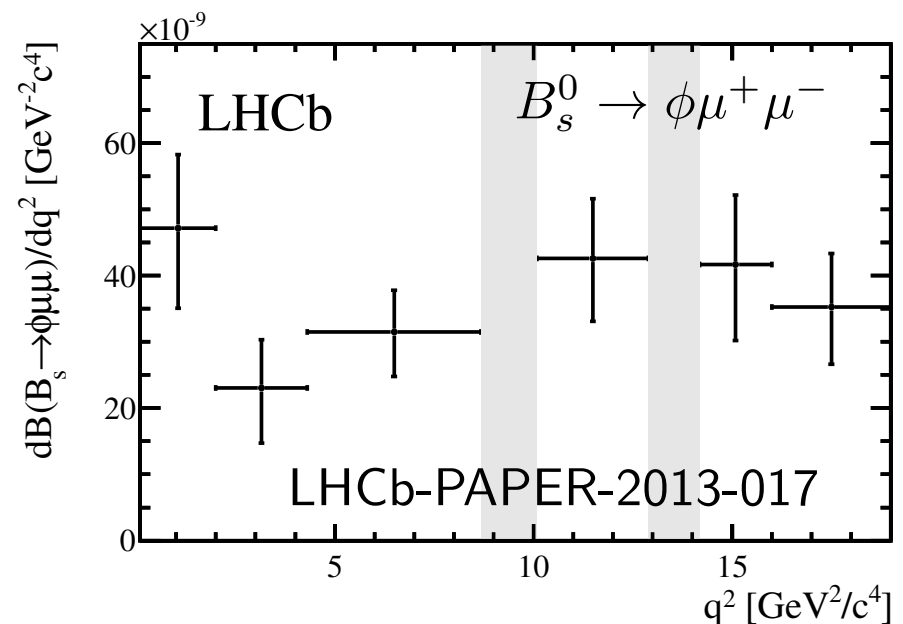


# Differential BF



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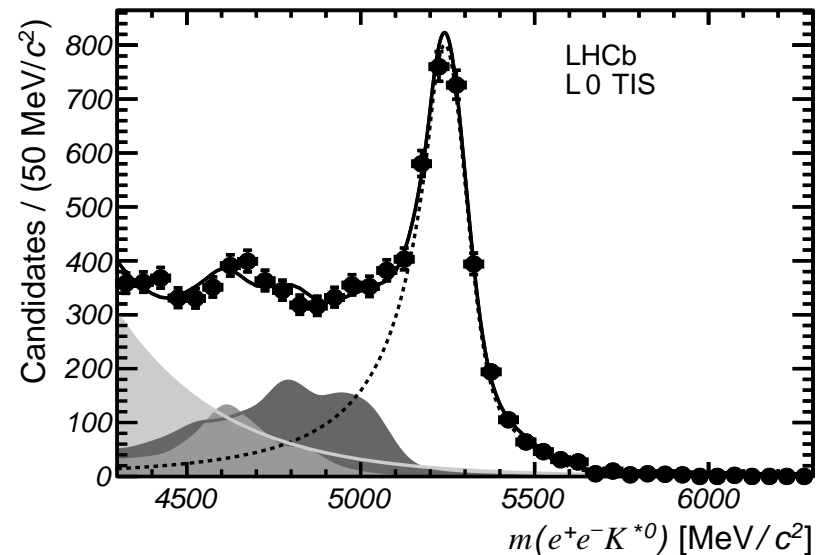
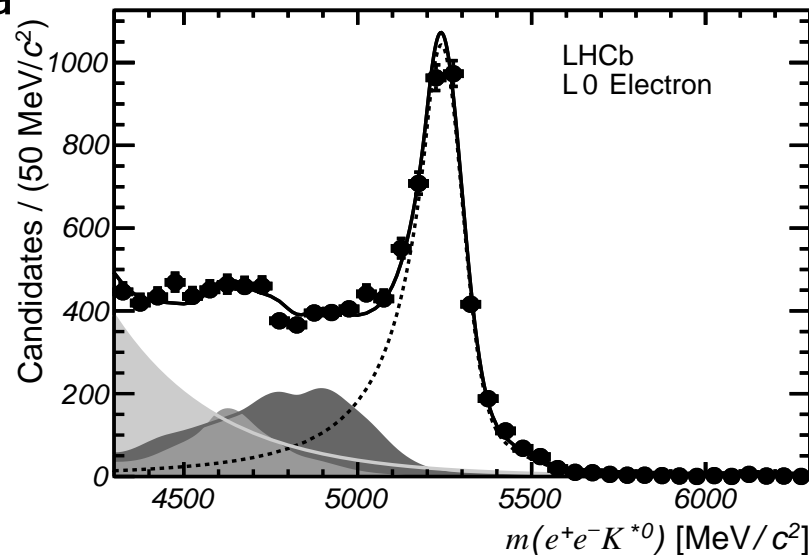
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# Low $q^2$ region

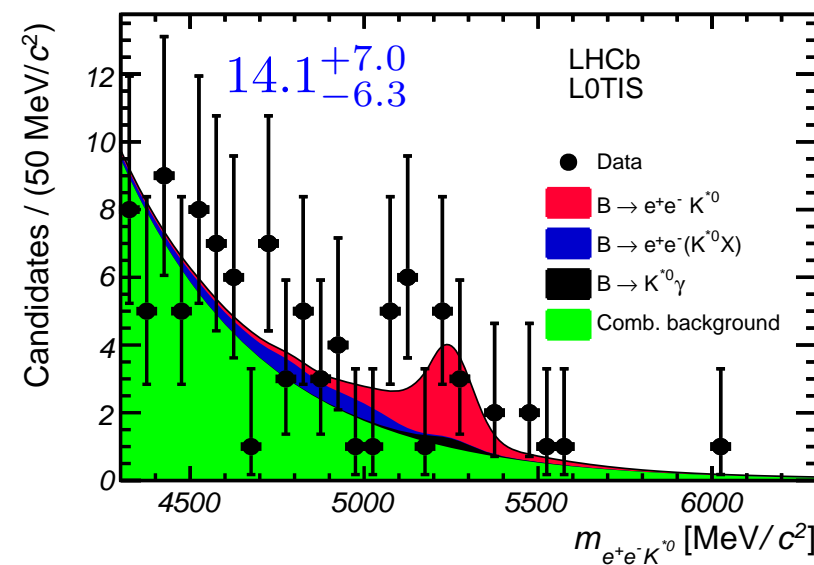
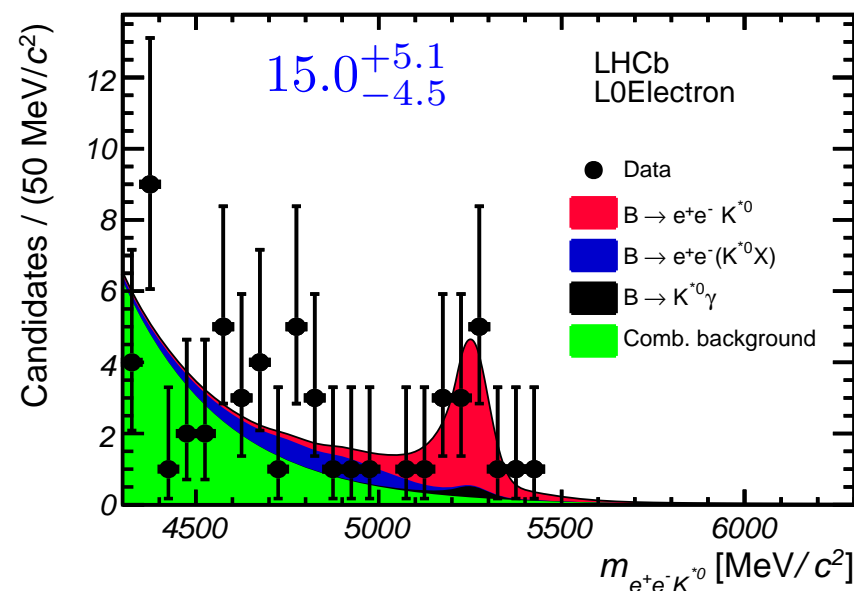
- Low  $q^2$  region particularly sensitive to  $\gamma$  polarization
- Decay  $B^0 \rightarrow K^{*0} e^+ e^-$  can reach lower  $q^2$  as decay with dimuons
- Never observed on its own at low  $q^2$
- Before moving to angular analysis, measure differential BF in  $30 < m(\ell^+ \ell^-) < 1000$  MeV
- Below 30 MeV angles hard to measure
- Analysis more challenging by significantly larger brehmstrahlung
- ⇒ Try to recover radiated photons
- Use of BDT for selection, trained on simulated signal and bg from data



# $B^0 \rightarrow K^{*0} e^+ e^-$ at low $q^2$

- Main physics background from  $B^0 \rightarrow K^{*0} \gamma$  with  $\gamma$  conversion
- Extended maximum likelihood fit
  - Signal: double Crystal ball
  - Combinatorial bg: exponential
  - Partially reconstructed bg from simulation
- Significance  $4.6\sigma$
- Systematic uncertainties well below statistical
- Dominant contributions from normalization BF and fit procedure
- Resulting  $B^0 \rightarrow K^{*0} e^+ e^-$  BF  
 $(3.1^{+0.9}_{-0.8} \text{ } ^{+0.2}_{-0.3} \pm 0.3) \times 10^{-7}$   
 in  $m(\ell^+ \ell^-)$  within 30–1000 MeV

LHCb-PAPER-2013-005



# Angular analysis

- Angular distribution

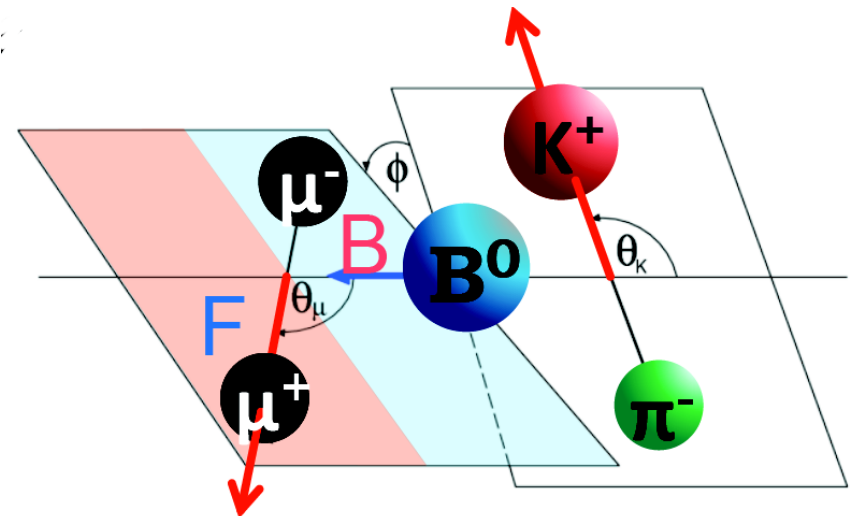
$$\begin{aligned}
 & S_1^s \sin^2 \theta_K + S_1^c \cos^2 \theta_K + \\
 & S_2^s \sin^2 \theta_K \cos 2\theta_\ell + S_2^c \cos^2 \theta_K \cos 2\theta_\ell + \\
 & S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + \\
 & S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi + S_6 \cos^2 \theta_K \cos \theta_\ell + \\
 & A_7 \sin 2\theta_K \sin \theta_\ell \sin \phi + A_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + \\
 & A_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi
 \end{aligned}$$

- Helicity angles  $\theta_\ell$  and  $\theta_K$  defined by  $\mu^+$  ( $\mu^-$ ) and  $K$  for  $B^0$  ( $\bar{B}^0$ )

- Angle  $\phi$  is given by

$$\sin \phi = (\hat{p}_{\mu^+} \times \hat{p}_{\mu^-}) \cdot \hat{p}_{K^{*0}}$$

$$\sin \phi = (\hat{p}_{\mu^-} \times \hat{p}_{\mu^+}) \cdot \hat{p}_{\bar{K}^{*0}}$$



**$B^0 \rightarrow K^{*0} \mu \mu$  decay plane**

# Angular analysis

- Assuming  $q^2 \gg 4m_\mu^2$  and transforming  $\phi$  as  $\phi + \pi$  if  $\phi < 0$  one can write

$$F_L \cos^2 \theta_K + \frac{3}{4}(1 - F_L)(1 - \cos^2 \theta_K) - F_L \cos^2 \theta_K (2 \cos^2 \theta_\ell - 1) + \frac{1}{4}(1 - F_L)(1 - \cos^2 \theta_K)(2 \cos^2 \theta_\ell - 1) + S_3(1 - \cos^2 \theta_K)(1 - \cos^2 \theta_\ell) \cos 2\hat{\phi} + \frac{4}{3}A_{FB}(1 - \cos^2 \theta_K) \cos \theta_\ell + A_9(1 - \cos^2 \theta_K)(1 - \cos^2 \theta_\ell) \sin 2\hat{\phi}$$

- The observables have to satisfy:  
 $|A_{FB}| \leq \frac{3}{4}(1 - F_L), \quad |A_9| \leq \frac{1}{2}(1 - F_L), \quad |S_3| \leq \frac{1}{2}(1 - F_L)$
- Those boundary conditions introduce non-trivial effects
- The statistical uncertainties estimated using Feldman-Cousins technique
- At lowest  $q^2$  bin one needs also correction from breakdown of  $q^2 \gg 4m_\mu^2$  assumption

# $B_s^0$ angular analysis

- After adding decay rate for  $B_s$  and  $\bar{B}_s$  it looks very similar to  $B^0$  angular distribution
- Main difference is that  $B_s \rightarrow \phi \mu^+ \mu^-$  is not self-tagging
- No sensitivity to  $A_{FB}$  unless production flavour is tagged
- Due to small statistics, flavour tagging not really possible
- Also full angular analysis difficult with current statistics
- Look to 1D distributions instead

$$\theta_K : \frac{3}{4}(1 - F_L)(1 - \cos^2 \theta_K) + \frac{3}{2}F_L \cos^2 \theta_K,$$

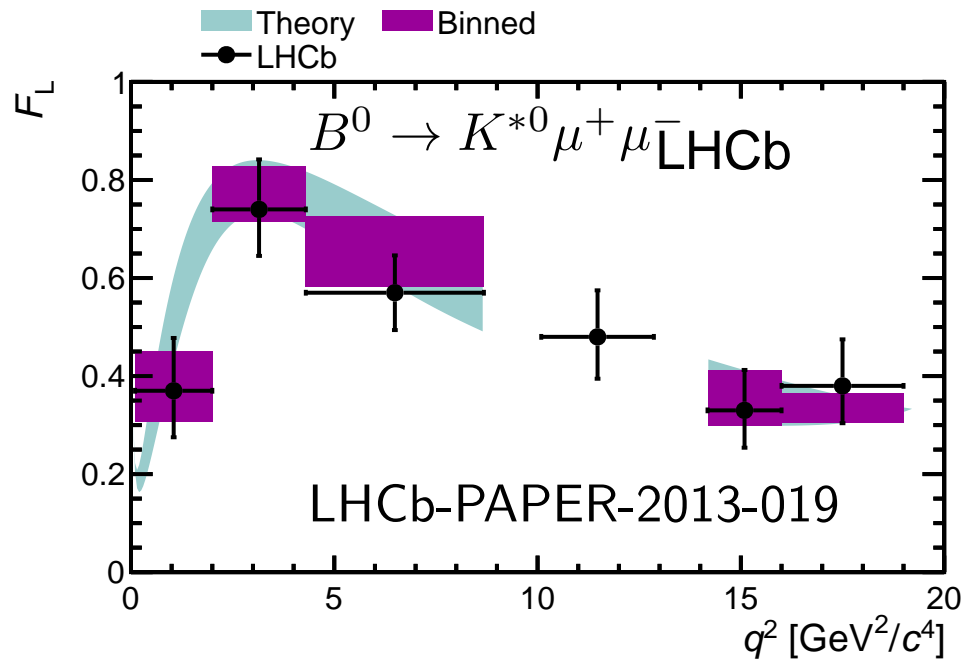
$$\theta_\ell : \frac{3}{8}(1 - F_L)(1 + \cos^2 \theta_\ell) + \frac{3}{4}F_L(1 - \cos^2 \theta_\ell) + \frac{3}{4}A_6^s \cos \theta_\ell,$$

$$d\phi : \frac{1}{2\pi} + \frac{1}{2\pi}S_3 \cos 2\phi + \frac{1}{2\pi}A_9 \sin 2\phi$$

- Subsequently fit all three distributions

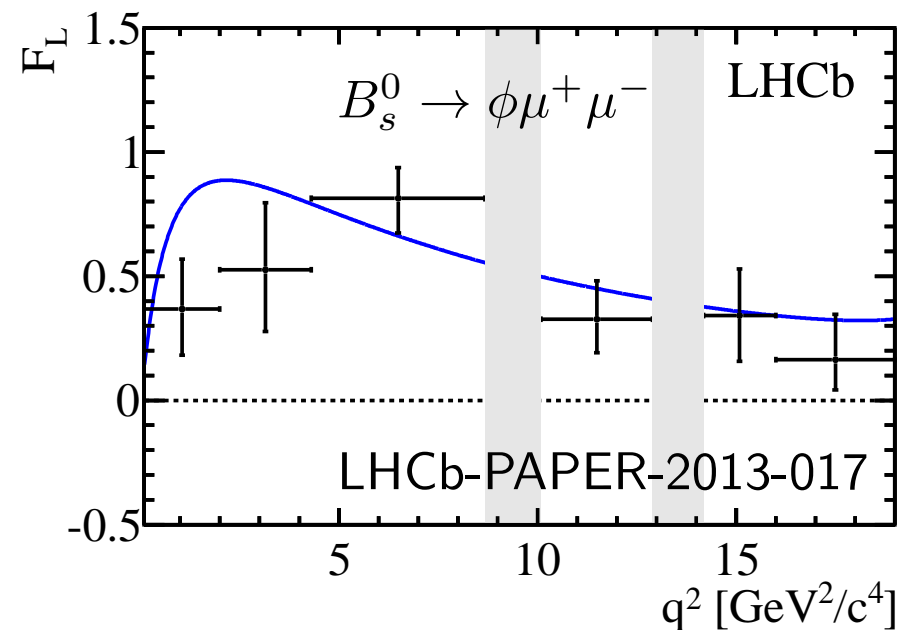
# Systematic uncertainties

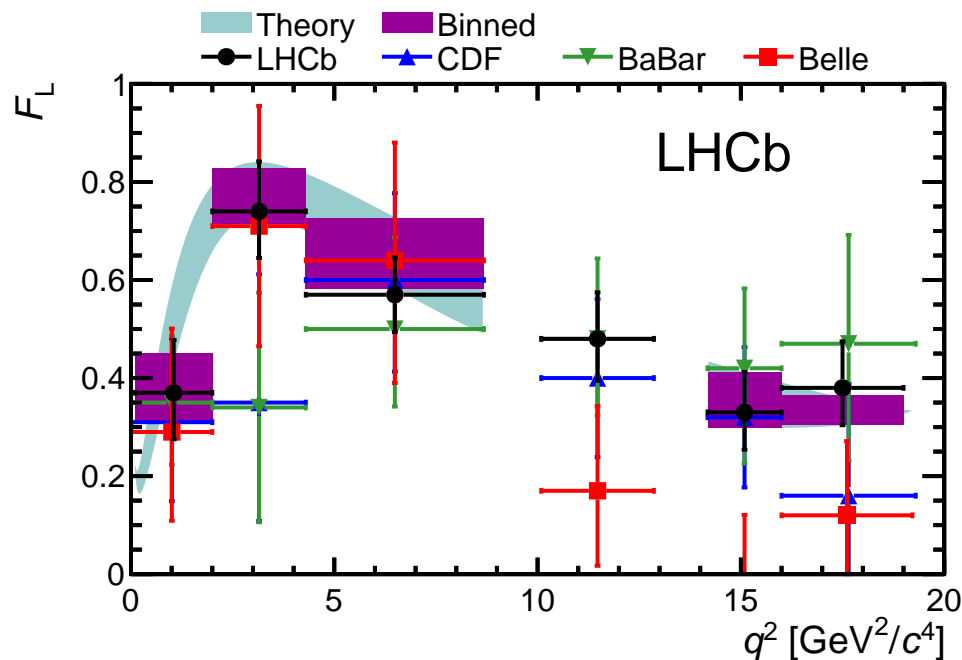
- Several sources investigated (not all for  $B_s$  due to small statistics)
  - Angular acceptance (dominant for  $B_s$ )
  - Data-MC differences
  - Mass model
  - Peaking backgrounds
  - S-wave (spin-0  $K\pi$ )
  - $B^0 \leftrightarrow \bar{B}^0$  mis-id from  $K \leftrightarrow \pi$  mis-id
  - $B^0, \bar{B}^0$  production and detection asymmetries
- Generally no single dominant contribution across all angular observables
- Many of the sources investigated are negligible
- Typical size for  $B_s$  of order 20% – 40% of statistical uncertainty
- For  $B^0$  typically of order 5% – 20% of statistical uncertainty



- $F_L$  is fraction of longitudinal  $K^{*0}/\phi$
- Theory affected by form factors
- Good agreement between experiment and SM

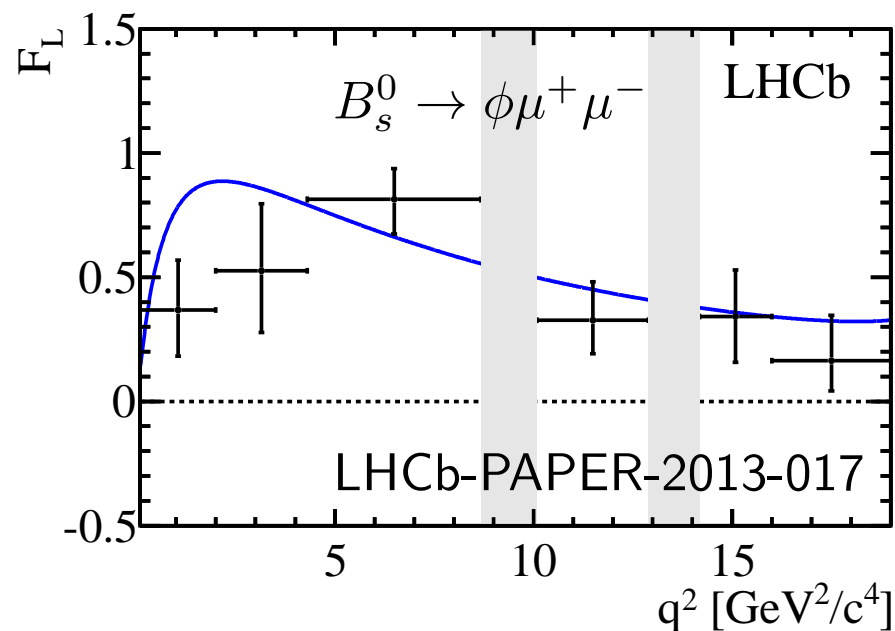
- First angular analysis of  $B_s^0 \rightarrow \phi \mu^+ \mu^-$
- Within limited statistics consistent with SM
- I would also say that  $B^0$  and  $B_s$  look consistent





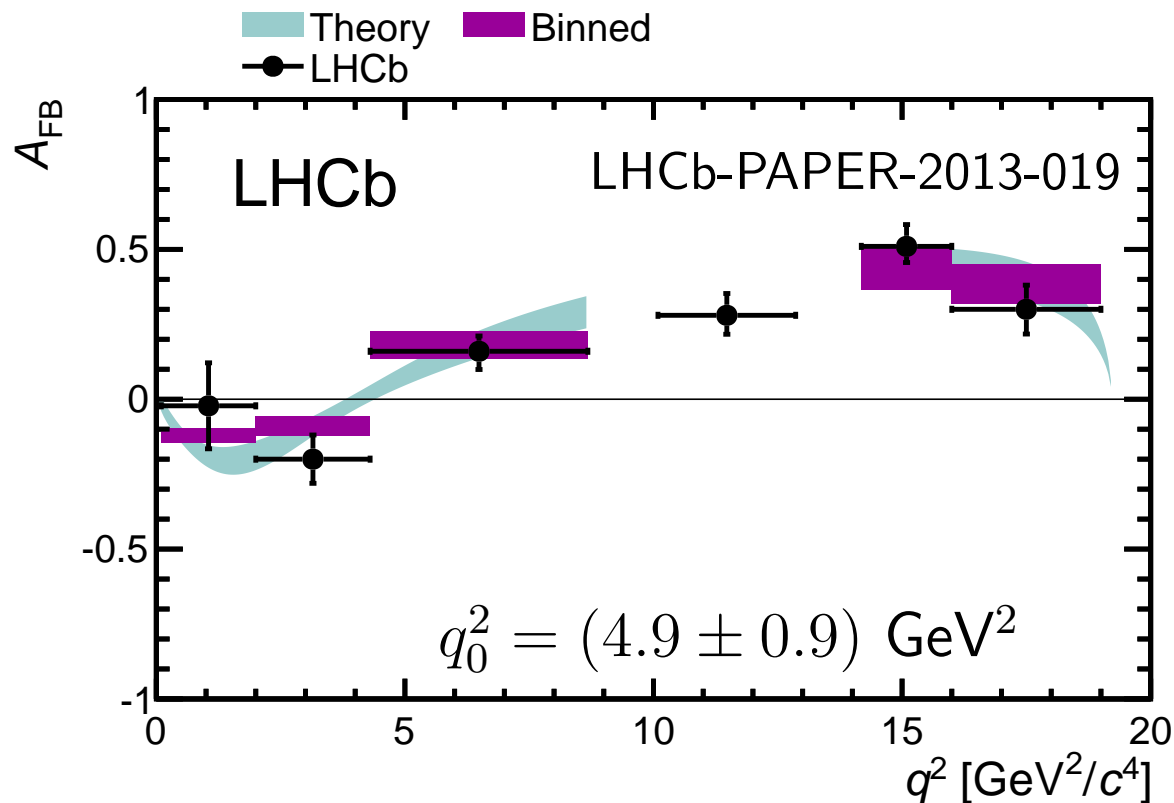
- First angular analysis of  $B_s^0 \rightarrow \phi \mu^+ \mu^-$
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- $F_L$  is fraction of longitudinal  $K^{*0}/\phi$
- Theory affected by form factors
- Good agreement between experiment and SM
- Consistent but more precise than previous measurements



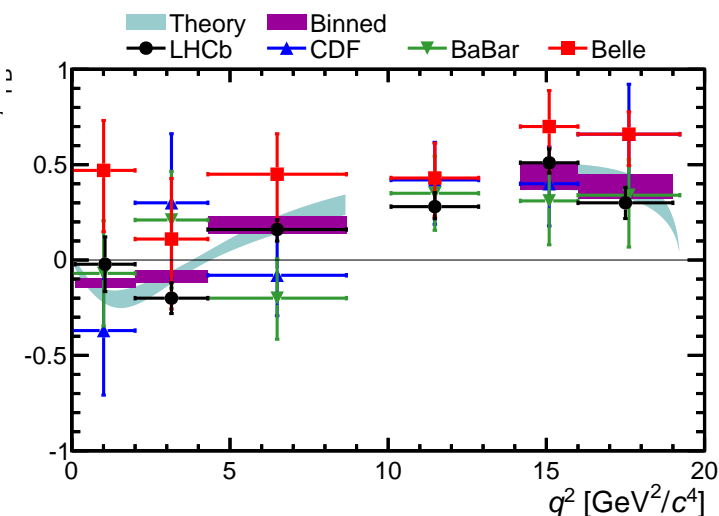


$$B^0 \rightarrow K^{*0} \mu^+ \mu^- \quad A_{FB}$$

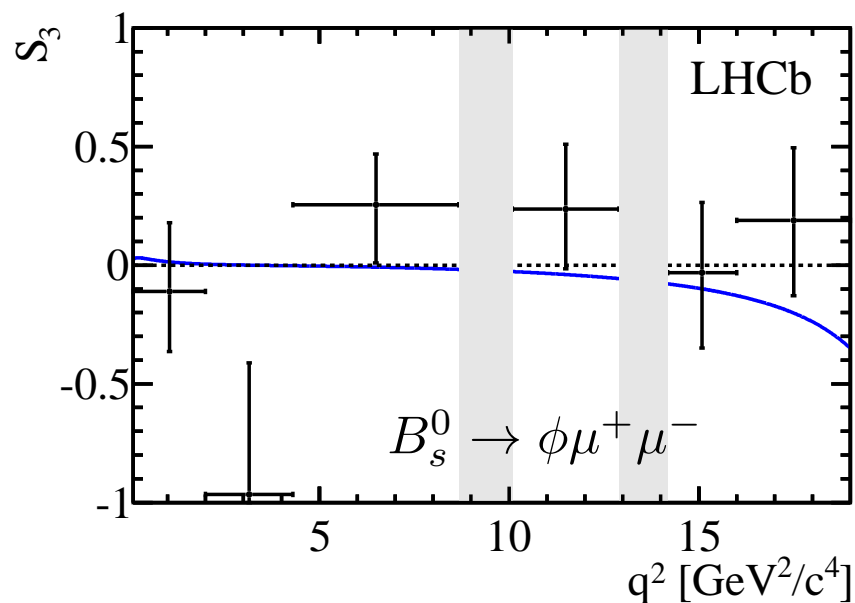
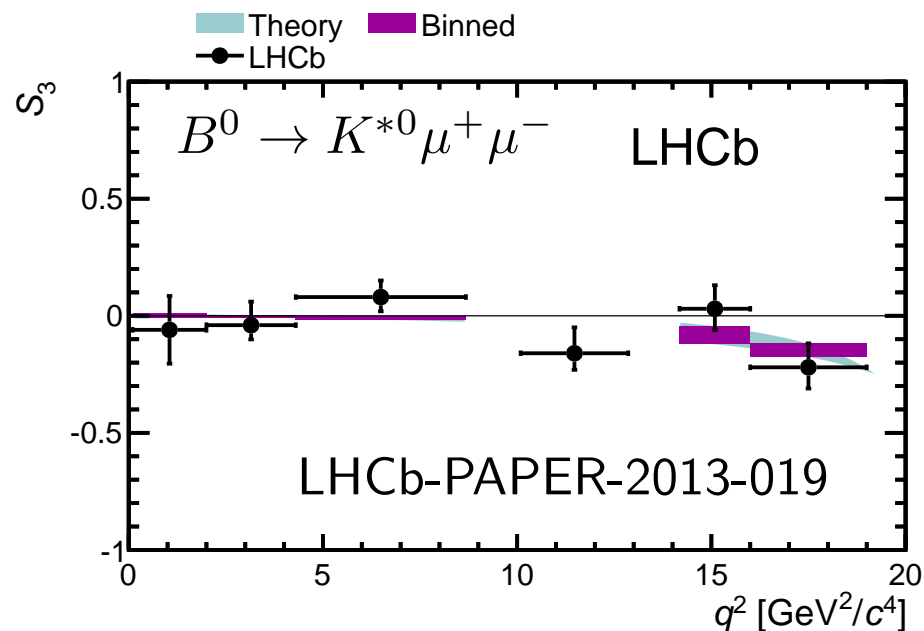


For  $B_s^0 \rightarrow \phi \mu \mu$   
no sensitivity without  
flavour tagging

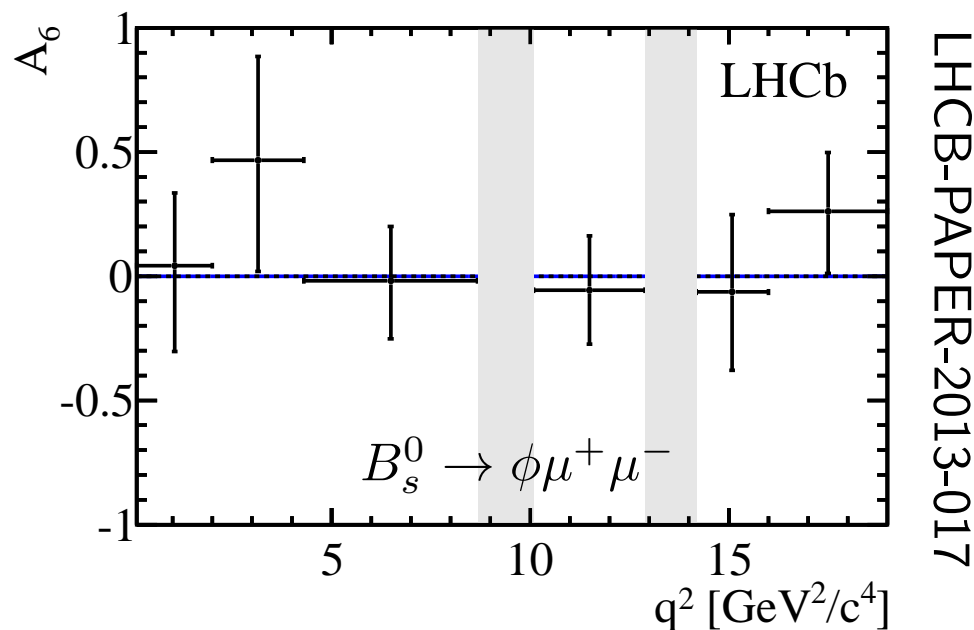
- Probably best known of angular observables
- Early measurements caused some excitements
- With increased statistical power of LHCb, fully consistent with SM



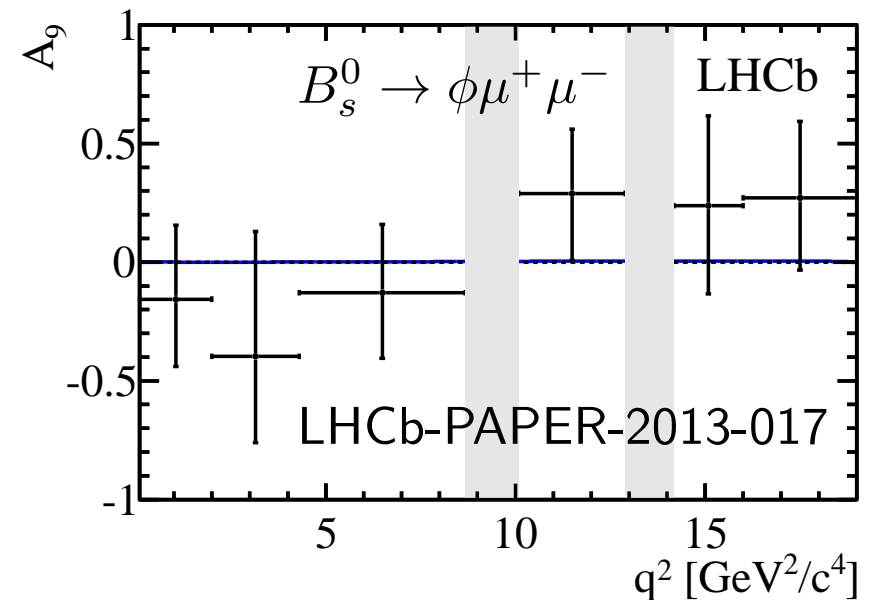
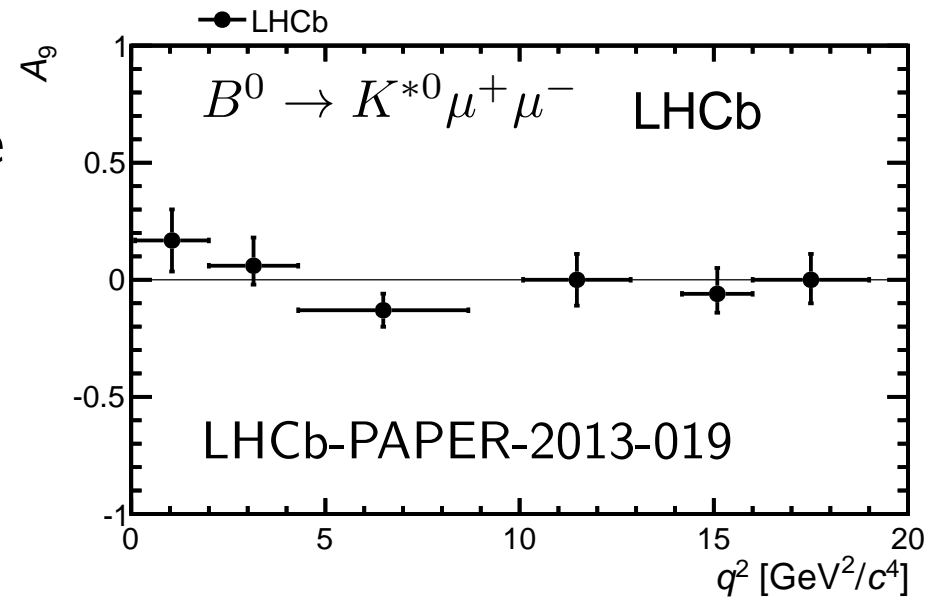
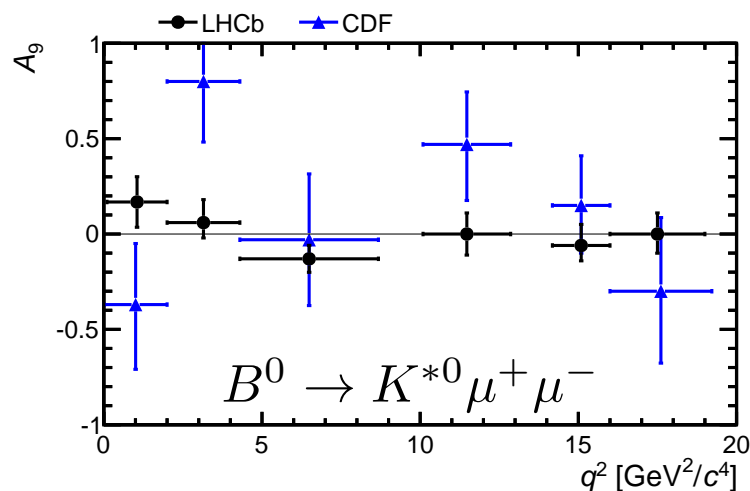
# $S_3$ and $A_6$



- $S_3$  is asymmetry between two transverse  $K^{*0}$  amplitudes (averaged  $B^0$  and  $\bar{B}^0$ )
- $A_6$  is asymmetry between  $B_s^0$  and  $\bar{B}_s^0$  of real part of transverse amplitudes interference
- Both agree with SM

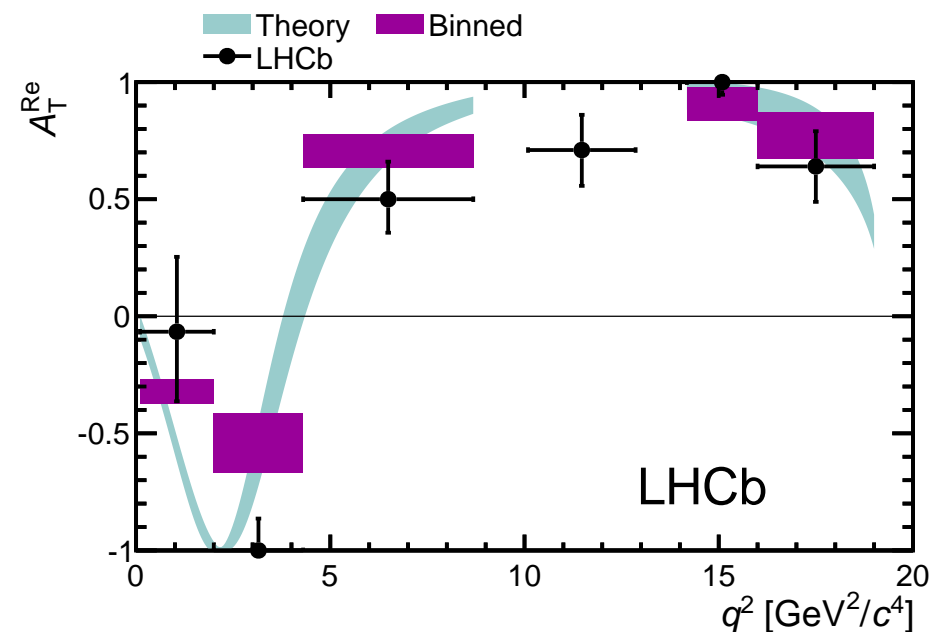
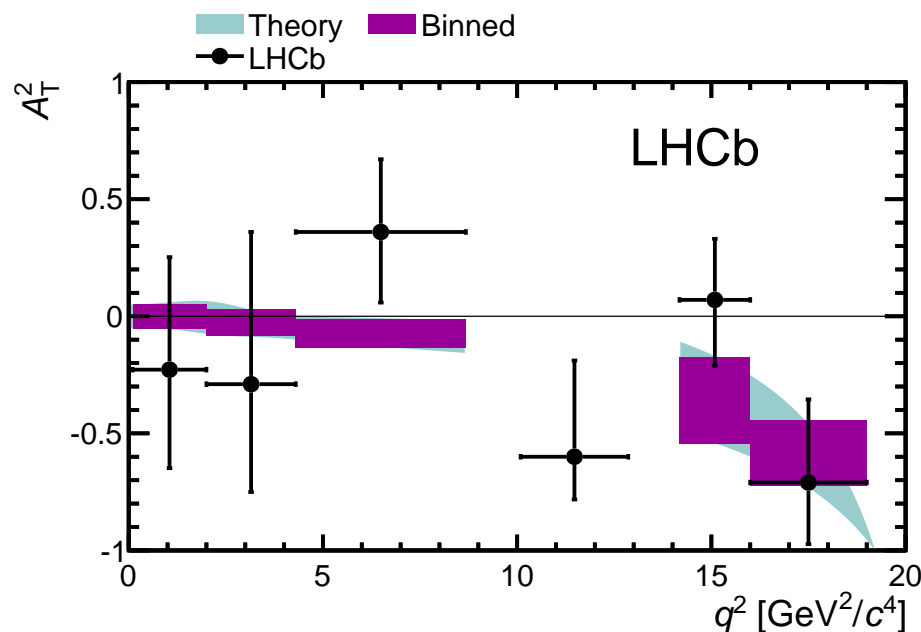


- T-odd asymmetry from interference of transverse  $K^{*0}$  amplitudes
- Sensitive when strong phase would make  $CP$  violation zero
- High sensitivity to right handed currents



# $B^0$ reparametrisation

- One can define alternative set of observables by  $S_3 = \frac{1}{2}(1 - F_L)A_T^2$  and  $A_{FB} = \frac{3}{4}(1 - F_L)A_T^{Re}$
  - In large recoil limit the observables  $A_T^2$  and  $A_T^{Re}$  have reduced form-factor uncertainties
  - It is not easy to handle correlations while maintaining solid meaning of uncertainties
- Extract also  $A_T^2$  and  $A_T^{Re}$



- Presented three new analyses of rare EWK decays
- First time we did multidimensional angular analysis of  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  (LHCb-PAPER-2013-019)
- First convincing signal for  $B^0 \rightarrow K^{*0} e^+ e^-$  at low  $q^2$  (LHCb-PAPER-2013-005)
- First angular analysis of  $B_s \rightarrow \phi \mu^+ \mu^-$  (LHCb-PAPER-2013-017)
- All results are worlds best
- All consistent with SM
- But placing strong constraints on new physics
- Papers to be submitted soon
  
- $B^0$  theory based on JHEP 07 (2011) 067, NP B612 (2001) 25, PR D70 (2004) 114005, EPJ C71 (2011) 1635, PR D71 (2005) 014029, JHEP 11 (2008) 032
- $B_s^0$  prediction from JHEP 01 (2009) 019, JHEP 0807 (2008) 106, PR D71 (2005) 014029