

The Quest for Naturalness in Supersymmetry

LPC Topic of the Week

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Many thanks for the discussions with:

R. Cavanaugh, M. Mangano, F. Wuerthwein, T. Han and S. Su

Outline

- Introduction
- Naturalness in Supersymmetry
- SUSY searches at the LHC
 - Inclusive searches
 - Third generation studies
 - SUSY electroweak(ino) production
- Charginos/Neutralinos in the light of the Higgs boson
- Summary and conclusion

Observation of the Higgs boson at the LHC

Observation of the Higgs-like boson at the LHC ushers in a new era in particle physics

ATLAS (hep-ex: 1207.7214) and CMS (hep-ex: 1207.7235)

CMS Combined local significance

Expected: 5.8σ

Observed: 5.0σ

ATLAS Combined local significance

(WW, ZZ, $\gamma\gamma$)

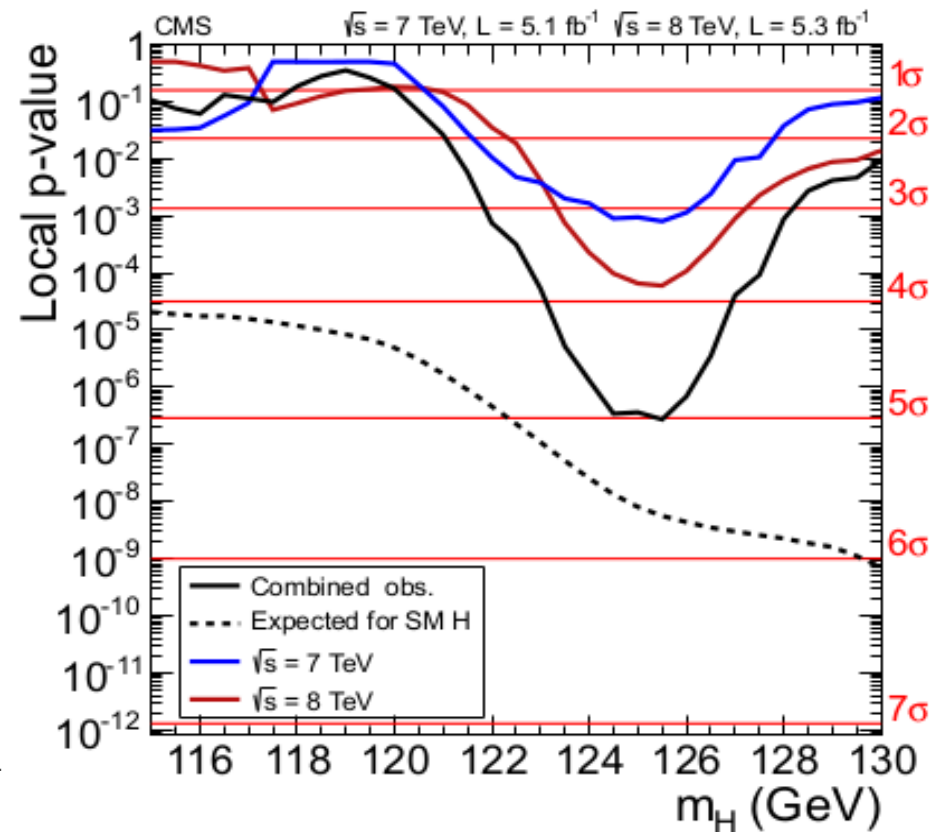
Expected: 4.9σ

Observed: 5.9σ

Combined Mass fit:

M (CMS) = 125.3 ± 0.4 (stat.) ± 0.5 (sys.) GeV

M (ATLAS) = 126.0 ± 0.4 (stat.) ± 0.4 (sys.) GeV

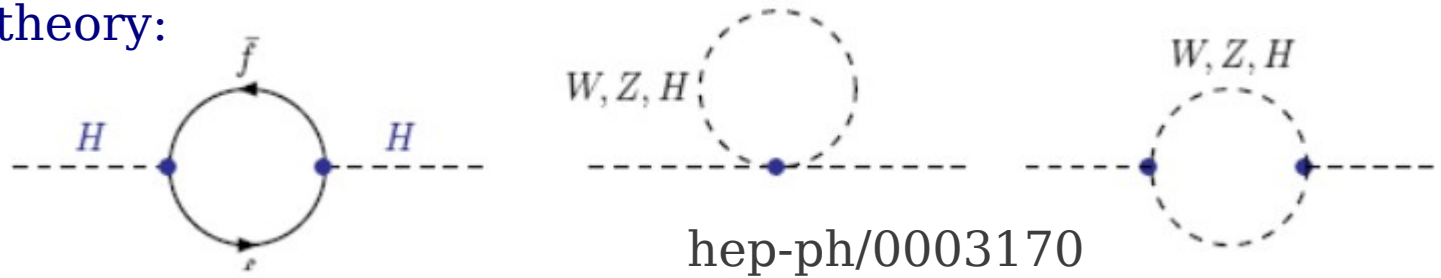


What does ~ 125 GeV Higgs indicate?

p-value: probability that background fluctuates to give an excess as large as the (average) signal size expected for a SM Higgs.

Higgs mass corrections

Using SM as effective theory:

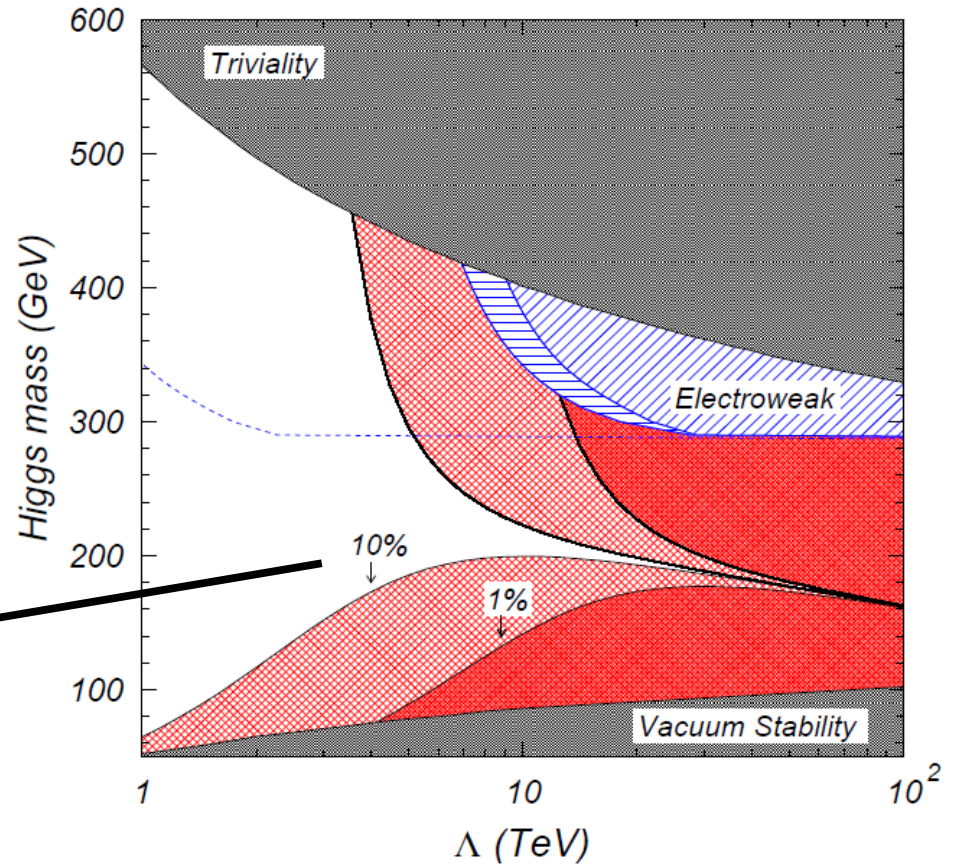


Corrections to the Higgs mass at one loop level.

$$\delta m_H^2 = \frac{3G_F}{4\sqrt{2}\pi^2} (4m_t^2 - 2m_W^2 - m_Z^2 - m_H^2) \Lambda^2$$

Either new physics appears at a scale Λ or there has to be a very delicate cancellation

Amount of fine-tuning



Fine tuning analogy



Giudice (arXiv:0801.2562)::

“The necessary accuracy needed to reproduce G_F/G_N is equal to the accuracy needed to balance a pencil as long as the solar system on a tip a millimeter wide.”

Is the Higgs mass protected by some symmetry?

First paper (to my knowledge) to explicitly search for SUSY: MARK-J at DESY

VOLUME 45, NUMBER 24

PHYSICAL REVIEW LETTERS

15 DECEMBER 1980

Experimental Study of Heavy Charged Leptons and Search for Scalar Partners of Muons at PETRA ($12 \text{ GeV} \leq E_{\text{c.m.}} \leq 36.7 \text{ GeV}$)

Data from the MARK-J detector on the reactions $e^+e^- \rightarrow \mu^+\mu^-$, $\tau^+\tau^-$ in the center-of-mass energy range from 12 to 36.7 GeV are presented. The μ , τ radii are shown to be $< 10^{-16}$ cm. A search has been made for the production of a new heavy lepton and for the production of spin-0 supersymmetric partners of the muon. 95%-confidence-level lower limits of 16 GeV for the mass of a new charged heavy lepton and 15 GeV for the mass of the scalar partners of the muon are obtained.

Naturalness in Supersymmetry

arXiv:1203.5539

$$\frac{1}{2}M_Z^2 = \frac{(m_{H_d}^2 + \Sigma_d) - (m_{H_u}^2 + \Sigma_u) \tan^2 \beta}{(\tan^2 \beta - 1)} - \mu^2$$

“Tuned” due to the Higgs mass - Colored sector

SUSY weak sector

- Individual terms on right side should be comparable in magnitude

- “Large” cancellations are “unnatural”

- $|\mu|$ can be a measure of naturalness

Σ - arises from radiative correction $\rightarrow \Sigma_u \sim \frac{3f_t^2}{16\pi^2} \times m_{\tilde{t}_i}^2 \left(\ln(m_{\tilde{t}_i}^2 / Q^2) - 1 \right)$

Stop mass

For, $\Sigma \approx 1/2M_Z^2 \rightarrow m_{\tilde{t}_i} \approx 500$ GeV

Assuming $\mu \sim 150$ (200) GeV \rightarrow Mass(stop) ~ 1 (1.5) TeV

Other heavier Higgs can easily be in the TeV mass range and is perfectly natural:

$$m_A^2 \simeq 2\mu^2 + m_{H_u}^2 + m_{H_d}^2 + \Sigma_u + \Sigma_d$$

Naturalness in Supersymmetry

R. Barbieri

The key equations:

$$\frac{m_h^2}{2} \approx -|\mu|^2 + m_u^2 + \dots$$

$$\delta m_u^2 \approx -\frac{3y_t^2}{8\pi^2} (m_{\tilde{t}_L}^2 + m_{\tilde{t}_R}^2 + A_t^2) \log M/m_{\tilde{t}}$$

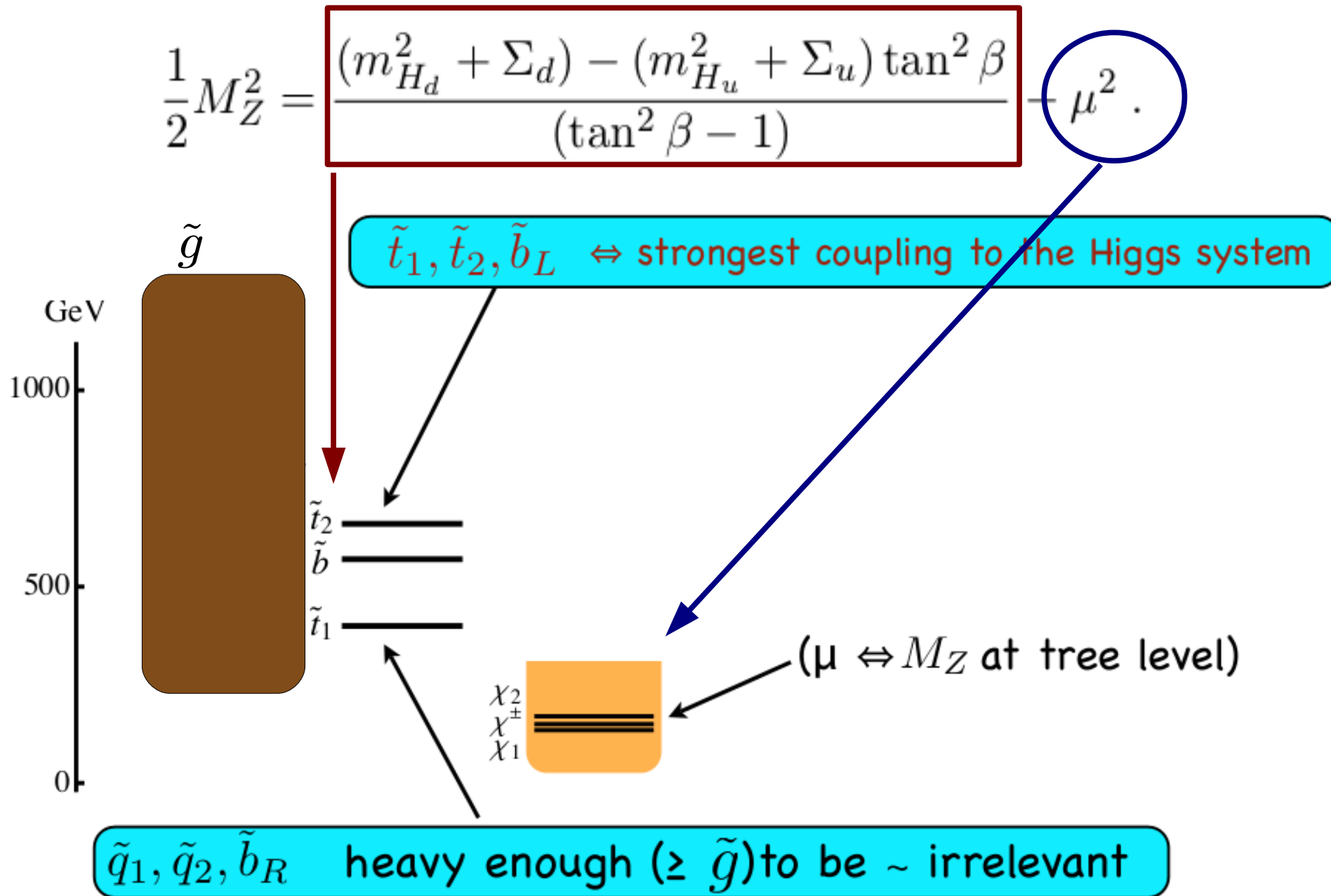
$$m_{\tilde{b}_L}$$

$$\delta m_{\tilde{t}}^2 \approx \frac{8\alpha_s}{3\pi} m_{\tilde{g}}^2 \log M/m_{\tilde{t}}$$

to be made more precise in any given SB-mediation scheme

see Dimopoulos, Giudice for SUGRA-mediation

Naturalness in Supersymmetry



Naturalness in Supersymmetry

Re-examine naturalness based on recent experimental results:

- Inclusive SUSY searches
- Search for SUSY third generation particle production
- Search for SUSY weak production

Inclusive SUSY Searches

SUSY Search strategy

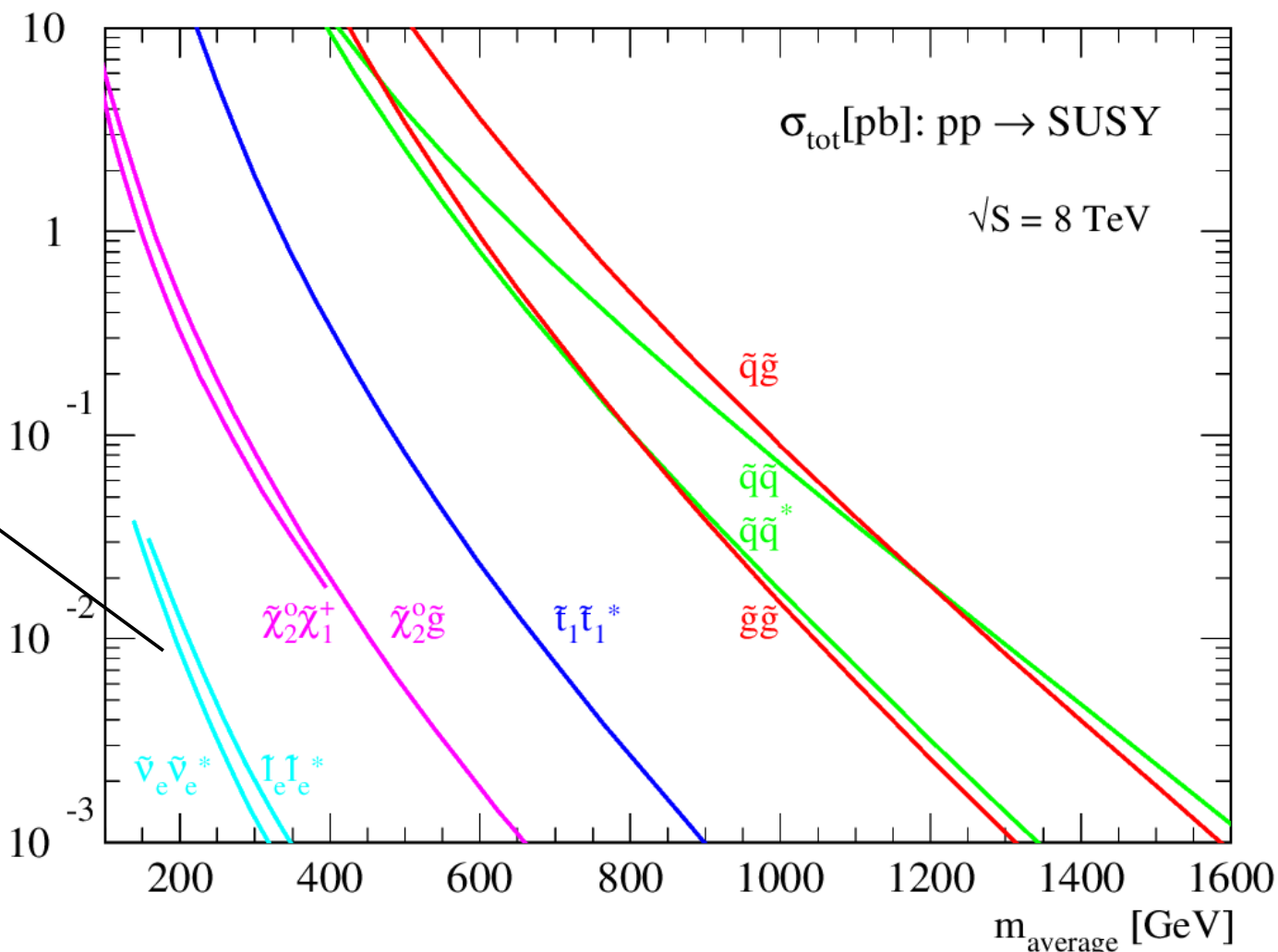
| 0-leptons | 1-lepton | OSDL | SSDL | ≥ 3 leptons | 2-photons | γ +lepton |
|------------|----------------------------|--------------------------------------|----------------------------------|------------------|-----------------------|-----------------------|
| Jets + MET | Single lepton + Jets + MET | Opposite-sign di-lepton + jets + MET | Same-sign di-lepton + jets + MET | Multi-lepton | Di-photon + jet + MET | Photon + lepton + MET |

Searches are defined (explore MET +X signatures):

- Categorized by the number of leptons in final state
- Generic missing energy signatures
- Many include jet requirements to be sensitive to strong production
- Direct stop/sbottom production using btag jets
- Direct electroweak production - Sensitive to leptonic final state

SUSY sparticle production at the LHC

EWK-inos have low cross section



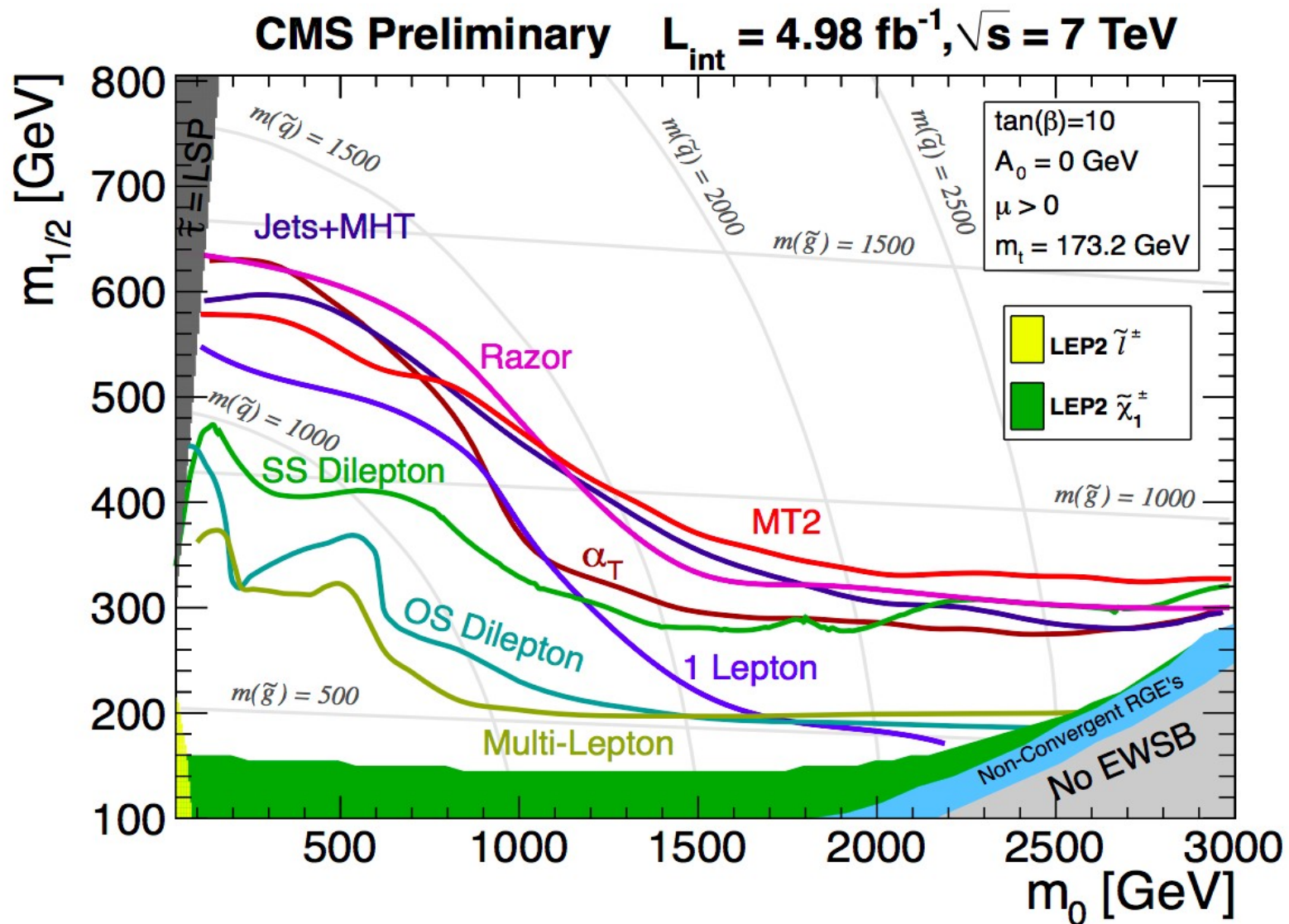
Largest cross sections are due to colored sources (squarks and gluinos)

- decouple them to study them individually

Stops and sbottoms have similar cross section, except at low sbottom masses

Inclusive SUSY searches

Landscape today - CMS Study

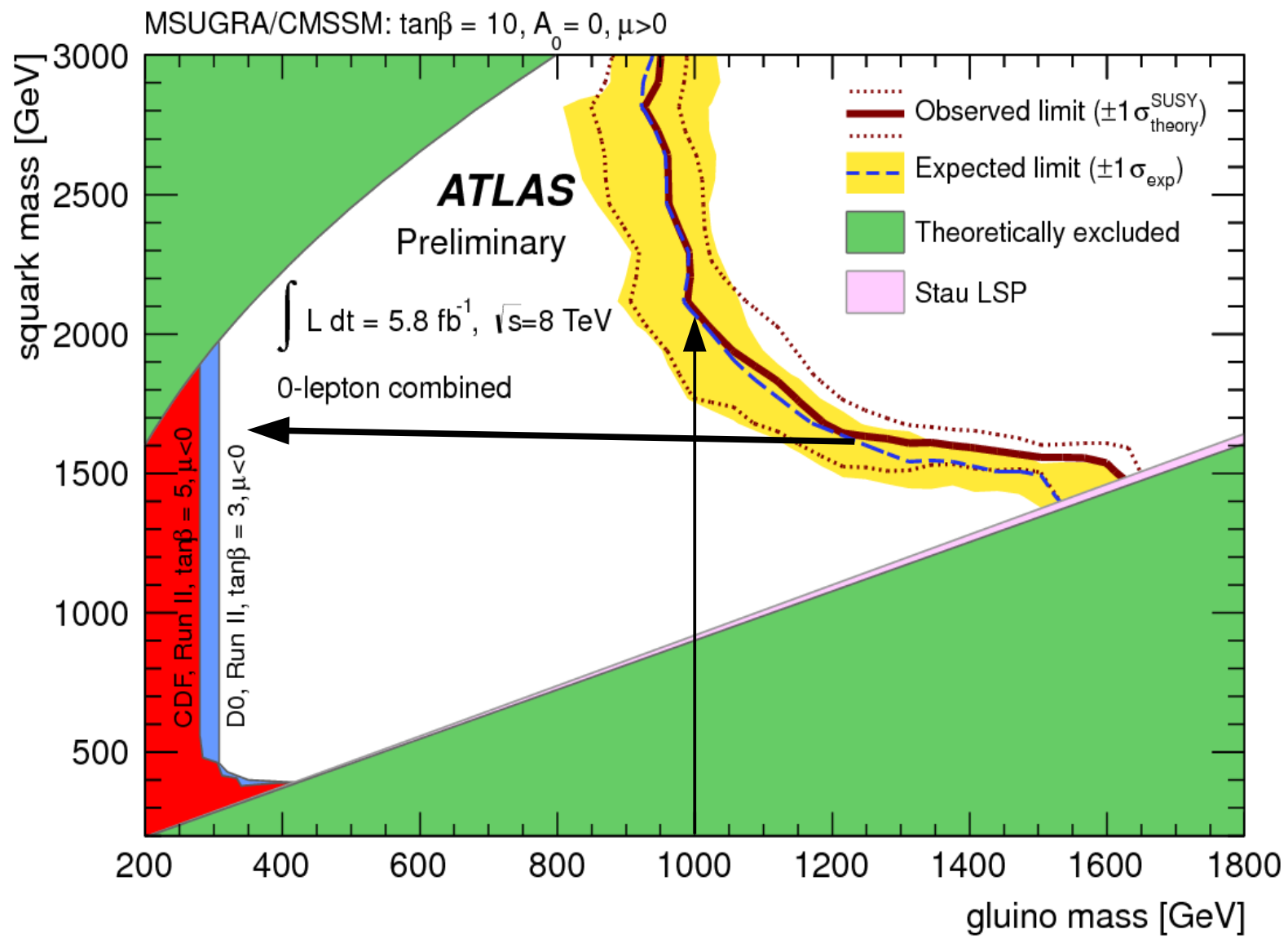


Impressive variety of inclusive SUSY searches - model dependent

Inclusive SUSY searches

Landscape today - ATLAS study

ATLAS-CONF-2012-109

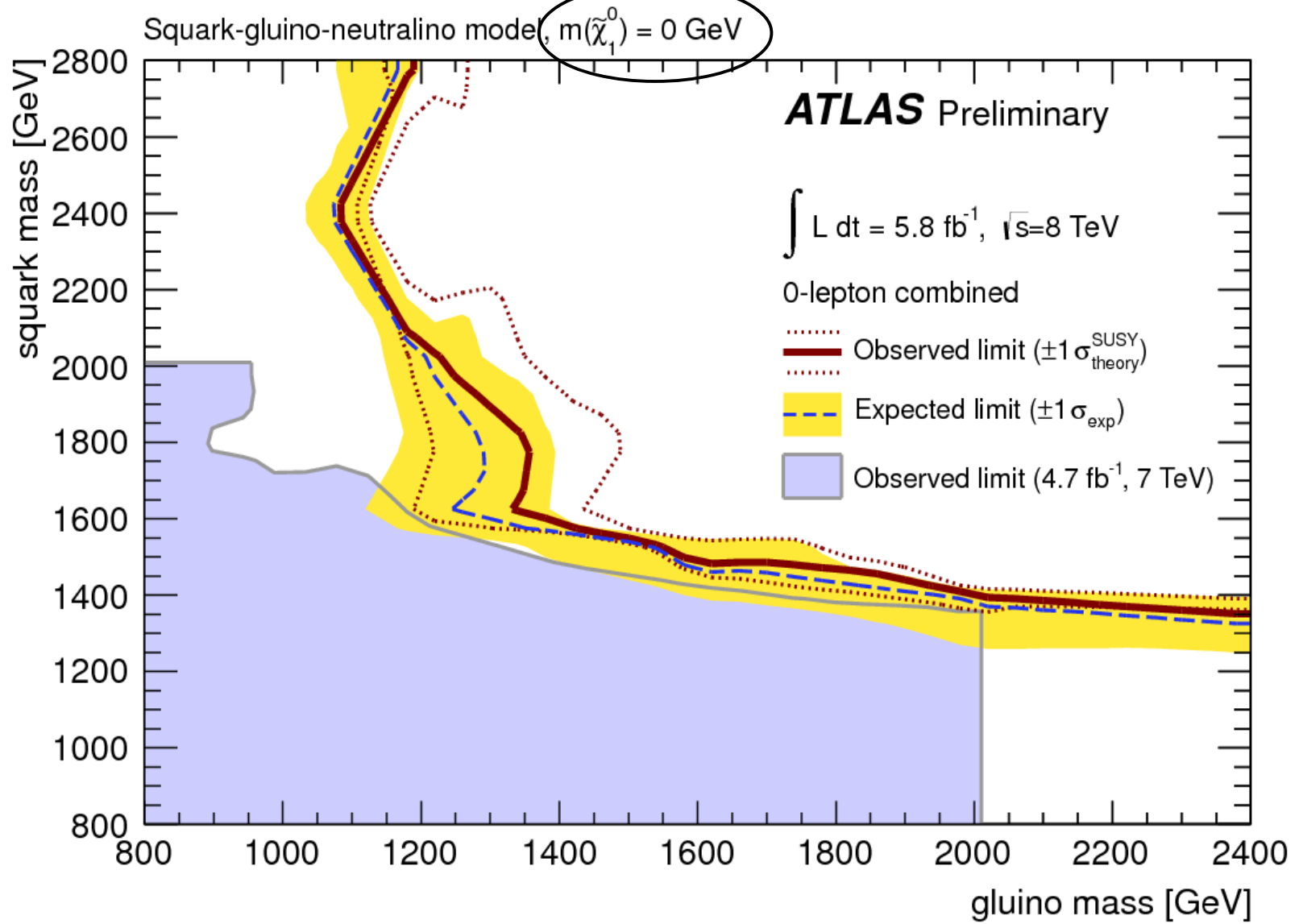


Impressive variety of inclusive SUSY searches - model dependent

Inclusive SUSY searches

Landscape today - ATLAS study

ATLAS-CONF-2012-109



Impressive variety of inclusive SUSY searches - model dependent

Inclusive search for squark and gluinos

ATLAS results (CMS similar)

ATLAS SUSY Searches* - 95% CL Lower Limits (Status: SUSY 2012)

| Search Category | Search Description | Lower Limit | Assumptions |
|---------------------------------------|--|-------------|---|
| Inclusive searches | MSUGRA/CMSSM : 0 lep + j's + $E_{T,miss}$ | 1.50 TeV | $\tilde{q} = \tilde{g}$ mass |
| | MSUGRA/CMSSM : 1 lep + j's + $E_{T,miss}$ | 1.24 TeV | $\tilde{q} = \tilde{g}$ mass |
| | Pheno model : 0 lep + j's + $E_{T,miss}$ | 1.18 TeV | \tilde{g} mass ($m(\tilde{q}) < 2$ TeV, light $\tilde{\chi}_1^0$) |
| | Pheno model : 0 lep + j's + $E_{T,miss}$ | 1.38 TeV | \tilde{q} mass ($m(\tilde{g}) < 2$ TeV, light $\tilde{\chi}_1^0$) |
| | Gluino med. $\tilde{\chi}_1^\pm (\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^\pm)$: 1 lep + j's + $E_{T,miss}$ | 900 GeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) < 200$ GeV, $m(\tilde{\chi}_1^\pm) = \frac{1}{2}(m(\tilde{\chi}_1^0) + m(\tilde{g}))$) |
| | GMSB : 2 lep (OS) + j's + $E_{T,miss}$ | 1.24 TeV | \tilde{g} mass ($\tan\beta < 15$) |
| | GMSB : 1-2 τ + 0-1 lep + j's + $E_{T,miss}$ | 1.20 TeV | \tilde{g} mass ($\tan\beta > 20$) |
| 3rd gen. squarks gluino mediated | GGM : $\gamma\gamma + E_{T,miss}$ | 1.07 TeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) > 50$ GeV) |
| | $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$ (virtual b) : 0 lep + 1/2 b-j's + $E_{T,miss}$ | 900 GeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) < 300$ GeV) |
| | $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$ (virtual b) : 0 lep + 3 b-j's + $E_{T,miss}$ | 1.02 TeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) < 400$ GeV) |
| | $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$ (real b) : 0 lep + 3 b-j's + $E_{T,miss}$ | 1.00 TeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) = 60$ GeV) |
| | $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ (virtual t) : 1 lep + 1/2 b-j's + $E_{T,miss}$ | 710 GeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) < 150$ GeV) |
| | $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ (virtual t) : 2 lep (SS) + j's + $E_{T,miss}$ | 850 GeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) < 300$ GeV) |
| | $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ (virtual t) : 3 lep + j's + $E_{T,miss}$ | 760 GeV | \tilde{g} mass (any $m(\tilde{\chi}_1^0) < m(\tilde{g})$) |
| | $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ (virtual t) : 0 lep + multi-j's + $E_{T,miss}$ | 1.00 TeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) < 300$ GeV) |
| | $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ (virtual t) : 0 lep + 3 b-j's + $E_{T,miss}$ | 940 GeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) < 50$ GeV) |
| | $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ (real t) : 0 lep + 3 b-j's + $E_{T,miss}$ | 820 GeV | \tilde{g} mass ($m(\tilde{\chi}_1^0) = 60$ GeV) |
| 3rd gen. squarks direct production | $b\bar{b}, b_1 \rightarrow b\tilde{\chi}_1^0$: 0 lep + 2-b-jets + $E_{T,miss}$ | 480 GeV | b mass ($m(\tilde{\chi}_1^0) < 150$ GeV) |
| | $b\bar{b}, b_1 \rightarrow t\tilde{\chi}_1^\pm$: 3 lep + j's + $E_{T,miss}$ | 380 GeV | \tilde{g} mass ($m(\tilde{\chi}_1^\pm) = 2m(\tilde{\chi}_1^0)$) |
| | $t\bar{t}$ (very light), $t \rightarrow b\tilde{\chi}_1^\pm$: 2 lep + $E_{T,miss}$ | 135 GeV | \tilde{t} mass ($m(\tilde{\chi}_1^0) = 45$ GeV) |
| | $t\bar{t}$ (light), $t \rightarrow b\tilde{\chi}_1^\pm$: 1/2 lep + b-jet + $E_{T,miss}$ | 120-173 GeV | \tilde{t} mass ($m(\tilde{\chi}_1^0) = 45$ GeV) |
| | $t\bar{t}$ (heavy), $t \rightarrow t\tilde{\chi}_1^0$: 0 lep + b-jet + $E_{T,miss}$ | 380-465 GeV | \tilde{t} mass ($m(\tilde{\chi}_1^0) = 0$) |
| | $t\bar{t}$ (heavy), $t \rightarrow t\tilde{\chi}_1^0$: 1 lep + b-jet + $E_{T,miss}$ | 230-440 GeV | \tilde{t} mass ($m(\tilde{\chi}_1^0) = 0$) |
| | $t\bar{t}$ (heavy), $t \rightarrow t\tilde{\chi}_1^0$: 2 lep + b-jet + $E_{T,miss}$ | 298-305 GeV | \tilde{t} mass ($m(\tilde{\chi}_1^0) = 0$) |
| EW direct | $t\bar{t}$ (GMSB) : $Z(\rightarrow ll) + b$ -jet + $E_{T,miss}$ | 310 GeV | \tilde{t} mass ($115 < m(\tilde{\chi}_1^0) < 230$ GeV) |
| | $l_1 l_1, l_1 \rightarrow l\tilde{\chi}_1^0$: 2 lep + $E_{T,miss}$ | 93-180 GeV | \tilde{l} mass ($m(\tilde{\chi}_1^0) = 0$) |
| | $\tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow l\nu(l\nu) \rightarrow l\nu\tilde{\chi}_1^0$: 2 lep + $E_{T,miss}$ | 120-330 GeV | $\tilde{\chi}_1^\pm$ mass ($m(\tilde{\chi}_1^0) = 0, m(\tilde{l}, \tilde{\nu}) = \frac{1}{2}(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$) |
| | $\tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow 3l(l\nu\nu) + \nu + 2\tilde{\chi}_1^0$: 3 lep + $E_{T,miss}$ | 60-500 GeV | $\tilde{\chi}_1^\pm$ mass ($m(\tilde{\chi}_1^\pm) = m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0) = 0, m(\tilde{l}, \tilde{\nu})$ as above) |

$\int L dt = (1.00 - 5.8) \text{ fb}^{-1}$
 $\sqrt{s} = 7, 8 \text{ TeV}$

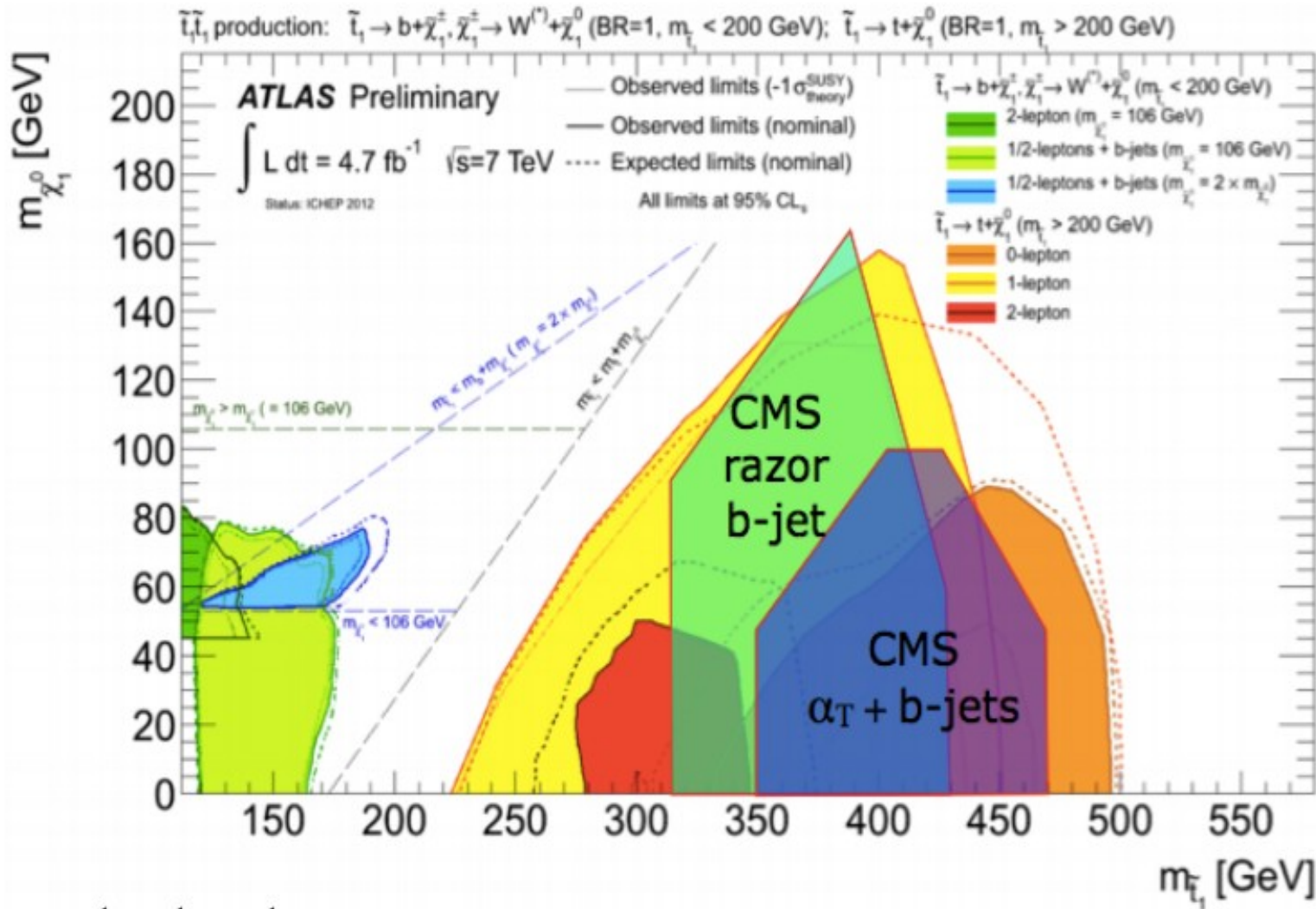
ATLAS Preliminary

Both ATLAS and CMS experiments have excluded SUSY colored production up to the TeV scale (with assumptions, not including direct stop/sbottom productions)

Search for SUSY third generation particle production

Third generation - stop searches

Direct stop production: $\tilde{t} \rightarrow t\tilde{\chi}_1^0; b\tilde{\chi}_1^\pm$



Daniele Alves
 "Implications"
 workshop
 @CERN

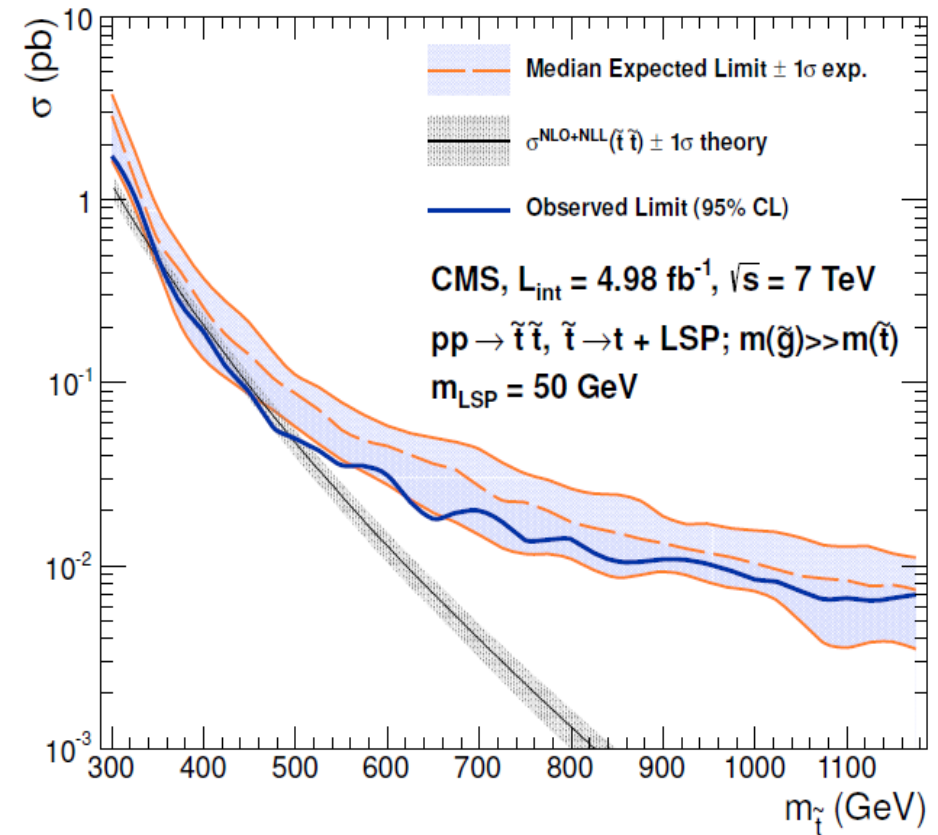
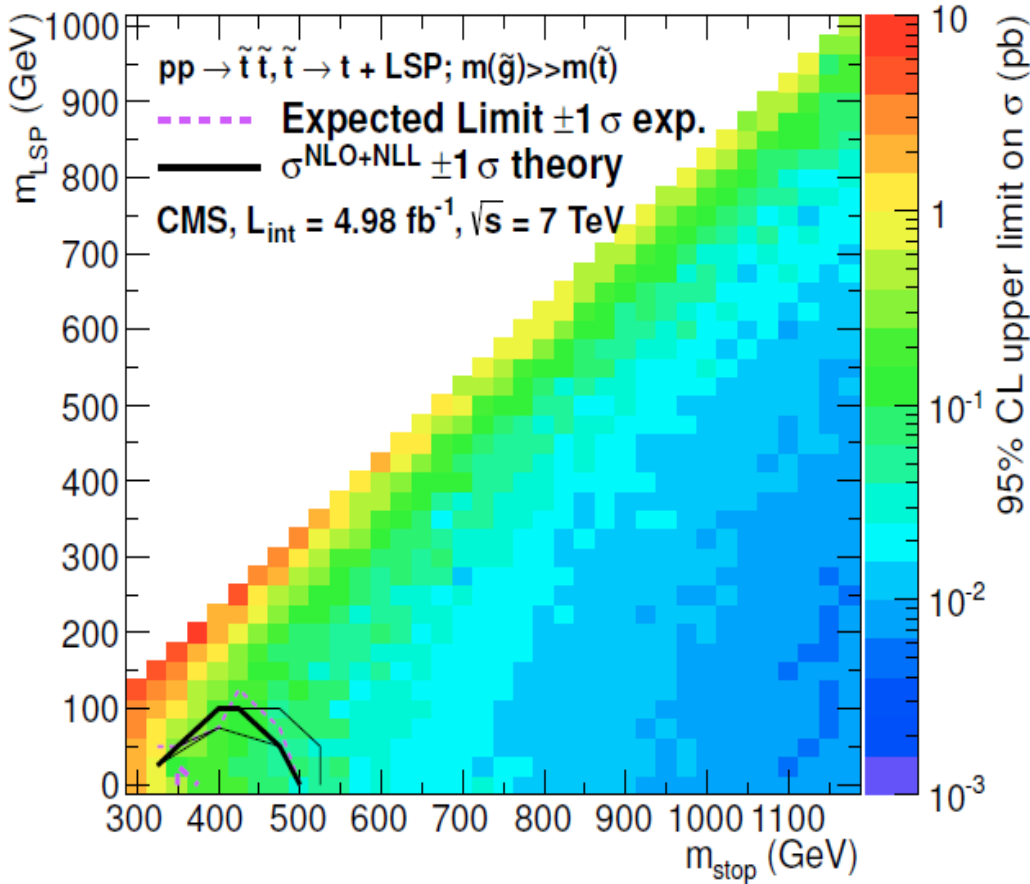
Impressive results from both collaborations

Bounds up to stop mass of $\sim 500 \text{ GeV}$

Third generation - stop searches

Direct stop production: $\tilde{t} \rightarrow t\tilde{\chi}_1^0$

CMS-SUS-11-022-003

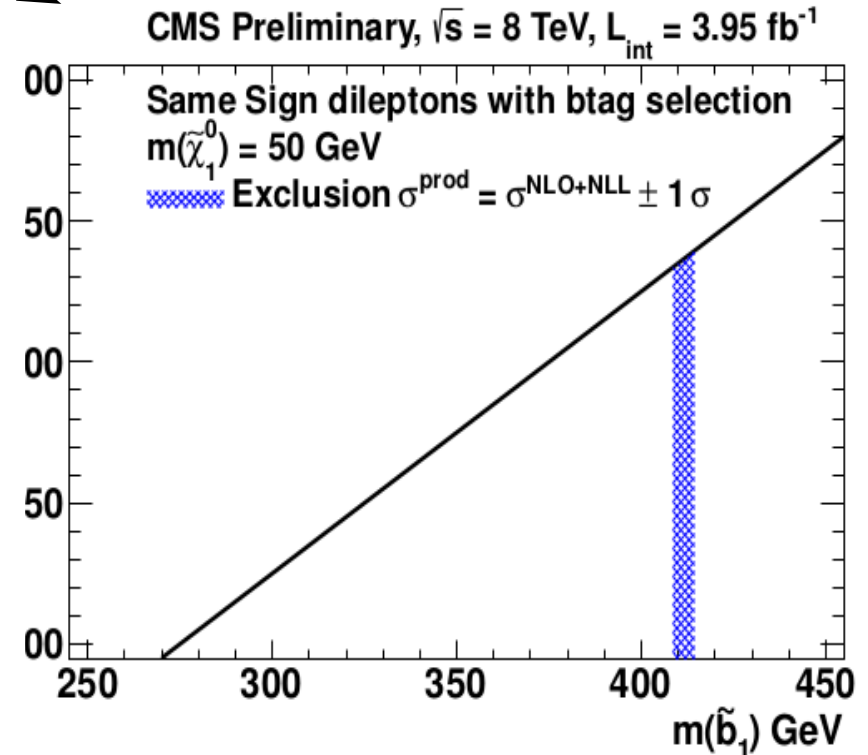
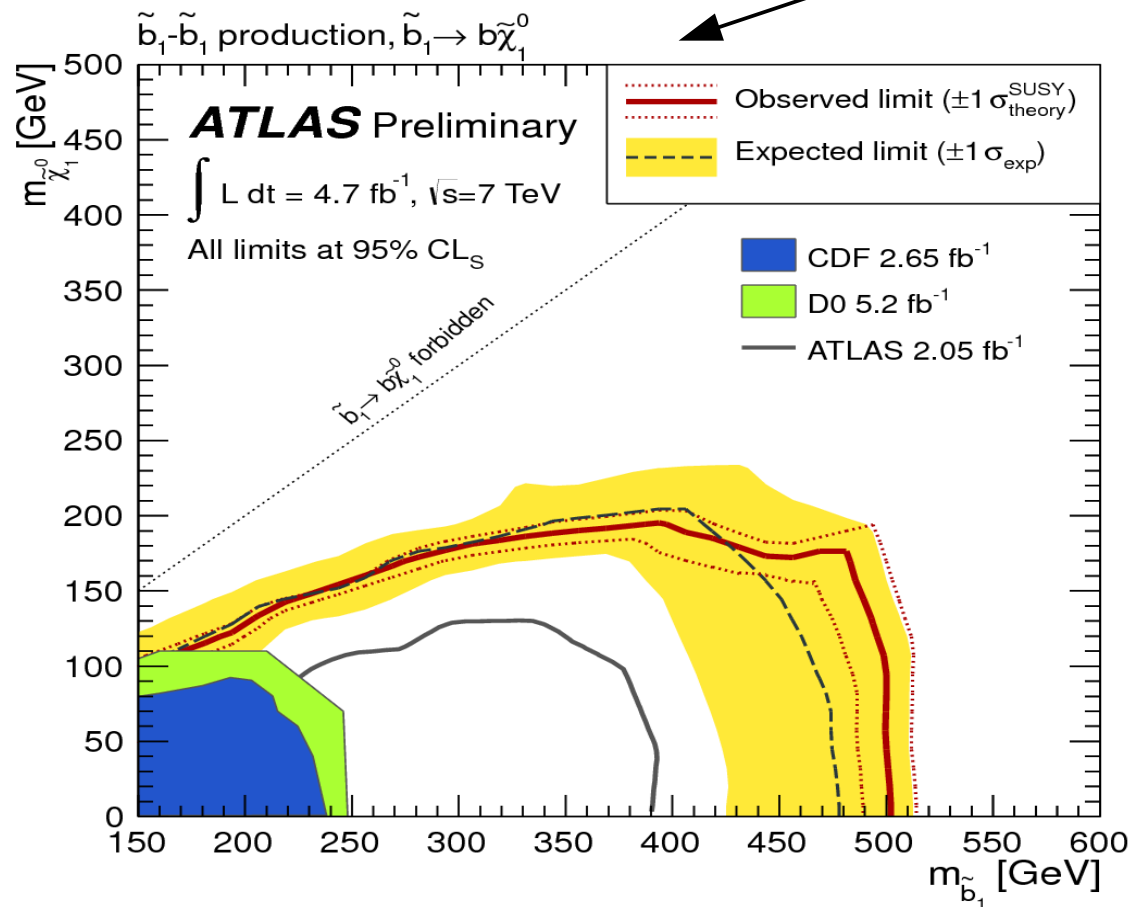


alpha_T based study with b-jets in the final state

For LSP mass = 50 GeV, Excluded stop mass between: 350 - 475 GeV

Third generation - sbottom searches

Sbottom production: $\tilde{b} \rightarrow b\tilde{\chi}_1^0; t\tilde{\chi}_1^-$



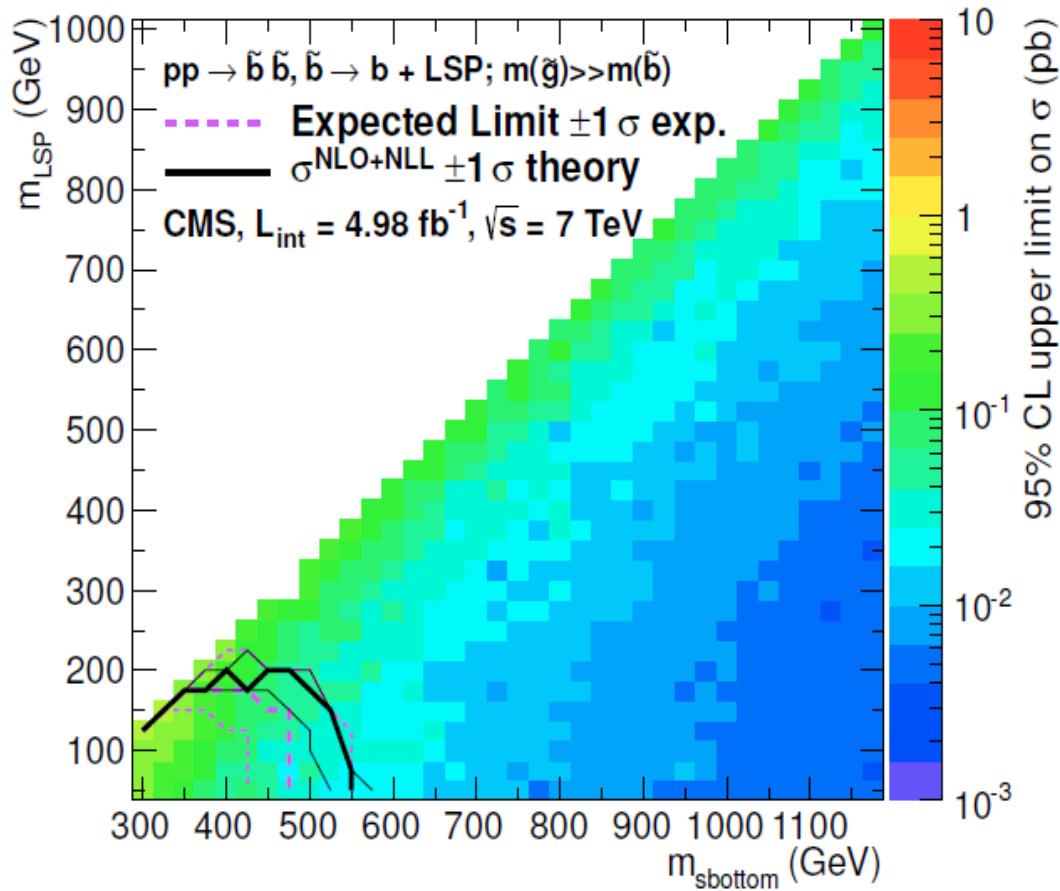
ATLAS hadronic btag jet (ATLAS-CONF-2012-106) searches provides

upper bound of $\sim 500 \text{ GeV}$ (LSP $\sim 50 \text{ GeV}$)

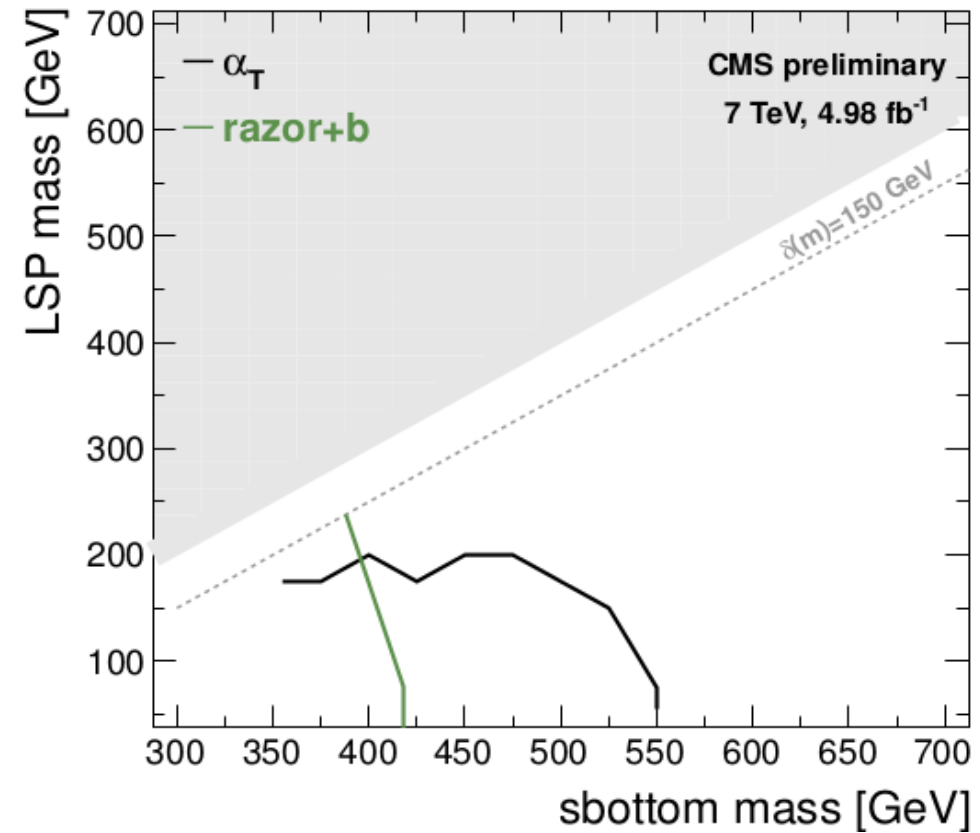
CMS Same sign dileptons with btag (SUS-12-017) $\sim 410 \text{ GeV}$

Third generation searches

Sbottom production: $\tilde{b} \rightarrow b\tilde{\chi}_1^0$



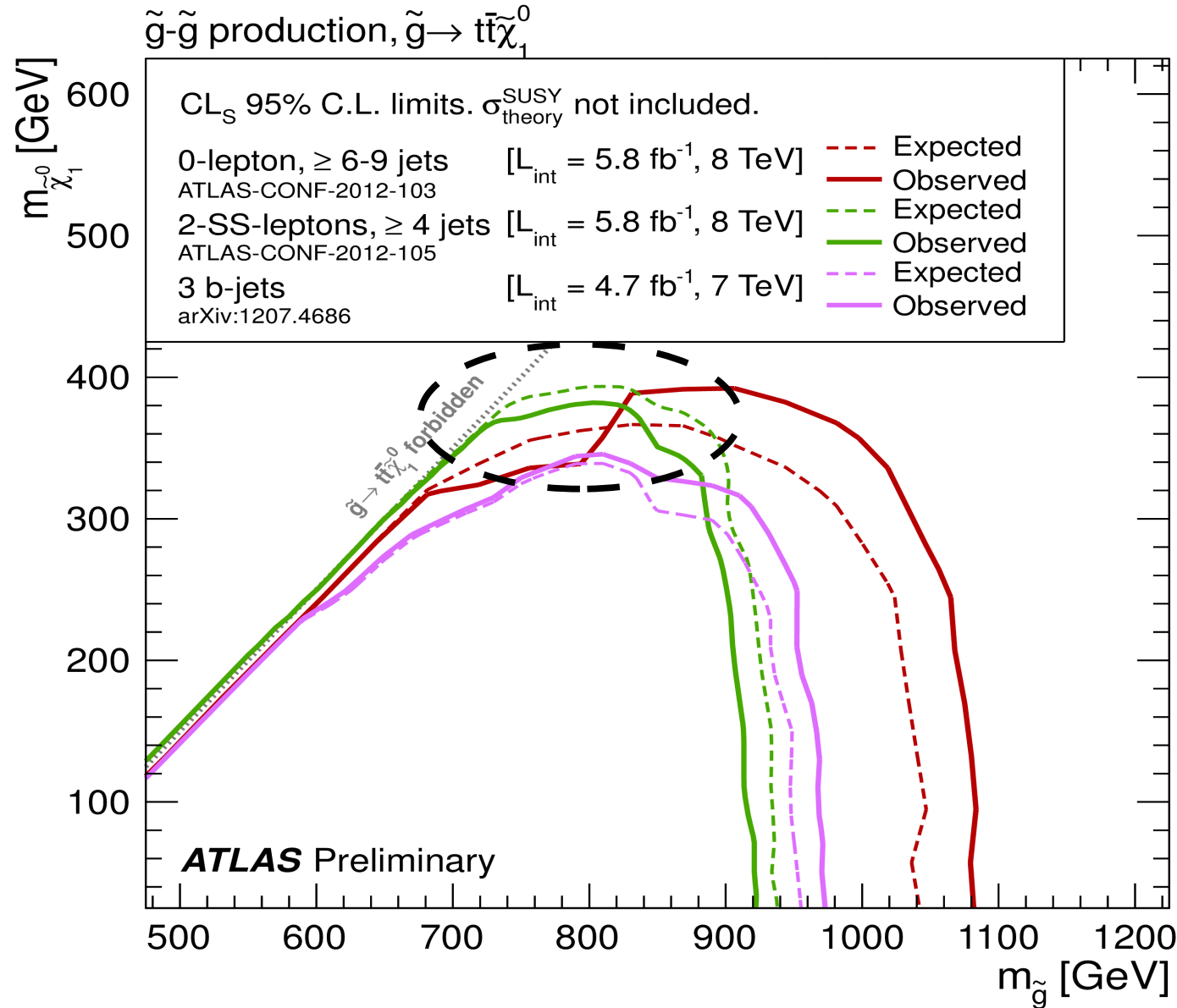
95% exclusion limits for $\tilde{b} \rightarrow b\tilde{\chi}_1^0$; $m(\tilde{g}, \tilde{q}) \gg m(\tilde{b})$



CMS hadronic studies exclude masses up to ~ 550 GeV (LSP = 50 GeV)

Third generation squark productions via gluinos

ATLAS-CONF-2012-103

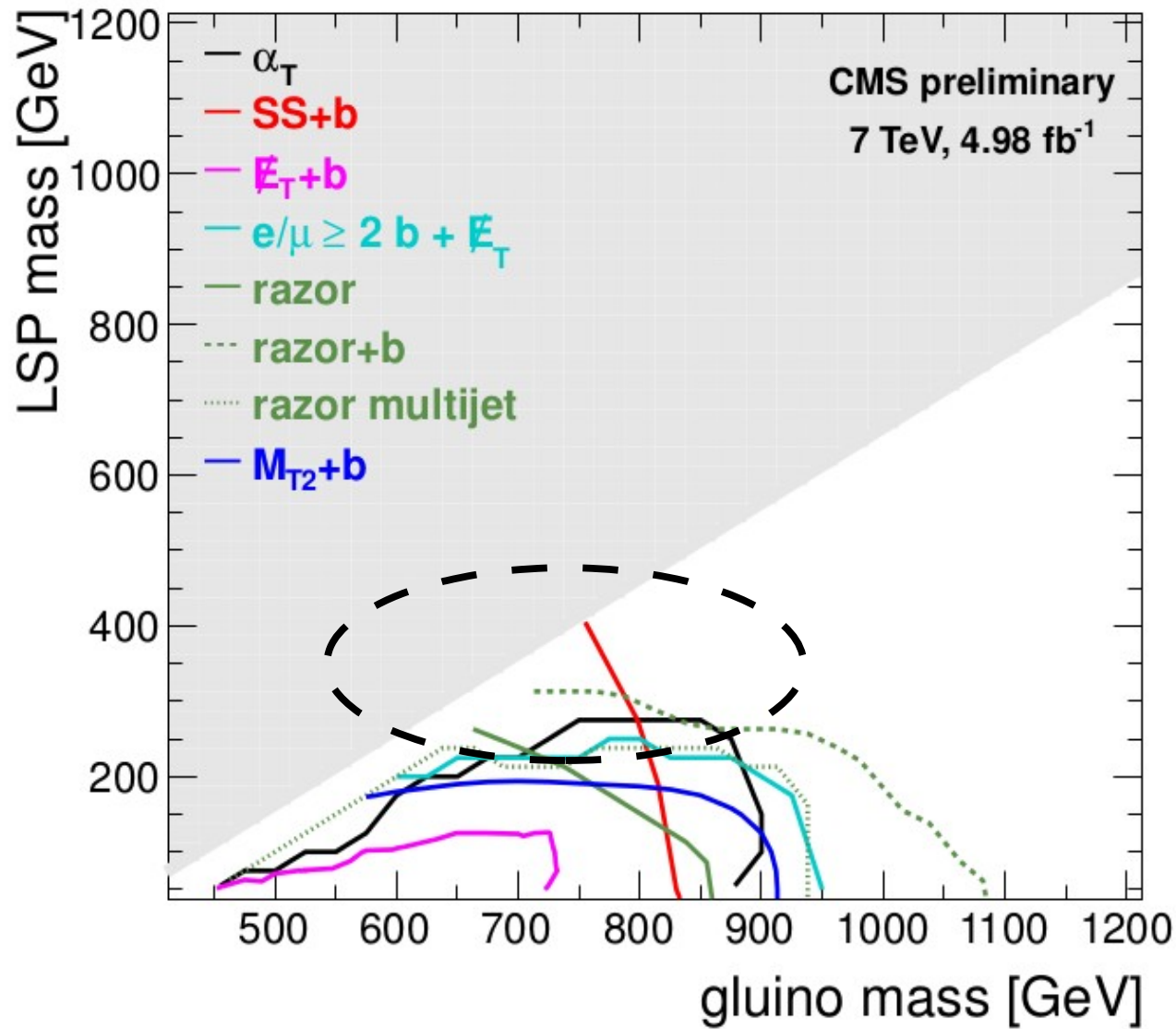


Impressive summary results - excludes gluino mass ~ 1.1 TeV (LSP mass ~ 50 GeV)

Same sign dilepton study probes into the “compressed” area

Third generation squark productions via gluinos

95% exclusion limits for $\tilde{g} \rightarrow t t \tilde{\chi}^0$; $m(\tilde{q}) \gg m(\tilde{g})$



Impressive summary results - excludes gluino mass ~ 1.1 TeV (LSP mass ~ 50 GeV)

Summary on SUSY Colored sector

Impressive search results (with assumptions) from both ATLAS and CMS Collab.

- For large 1st & 2nd generation squark masses, gluinos ~ 1.1 TeV excluded
- For large gluino mass, 1st and 2nd generation squarks ~ 1.5 TeV excluded
- With LSP mass ~ 50 GeV and decoupled gluinos
 - Direct stop production leads to mass limit ~ 500 GeV (ATLAS)
 - Direct stop production leads to mass limit $\sim 350 - 475$ GeV (CMS)
 - Direct sbottom production leads to mass limit ~ 550 GeV (CMS)
- Gluino production with third generation squarks ~ 1.1 TeV excluded

All of the above satisfies the naturalness conditions in SUSY

Searches for SUSY weak production

Experimental constraints from LEP

Chargino ($\tilde{\chi}_i^\pm; i = 1, 2$) and Neutralino ($\tilde{\chi}_i^0; i = 1 - 4$) productions at LEP:

$$e^+e^- \rightarrow \tilde{\chi}^+ \tilde{\chi}^- \rightarrow W^+W^- \tilde{\chi}_1^0 \tilde{\chi}_1^0$$

$$e^+e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^0 \rightarrow l^+l^- \tilde{\chi}_1^0 \tilde{\chi}_1^0$$

Neutralino pairs via s-channel Z or t-channel with slepton exchange

Using mSUGRA or CMSSM framework

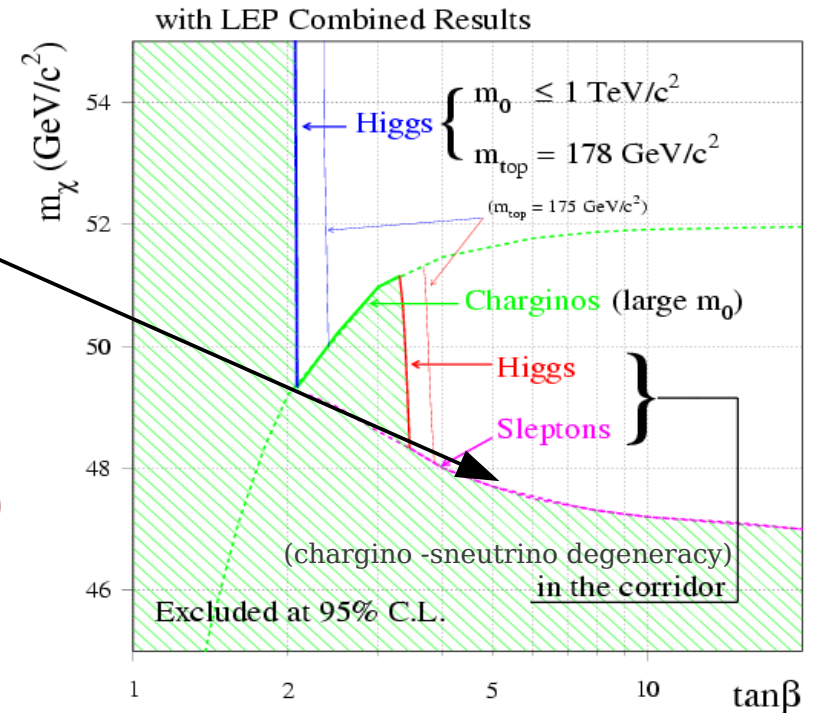
(assuming mixing in stau sector is small)

LSP mass below 47/50 GeV is excluded

However several assumptions are involved:

- mSUGRA / CMSSM
- gaugino mass unification
- $\tan\beta < 3.3$ limits at large M_0 (+higgs, chargino)
(M_0 - common sfermion mass at GUT)
- $\tan\beta > 3.3$ the limit is using small M_0

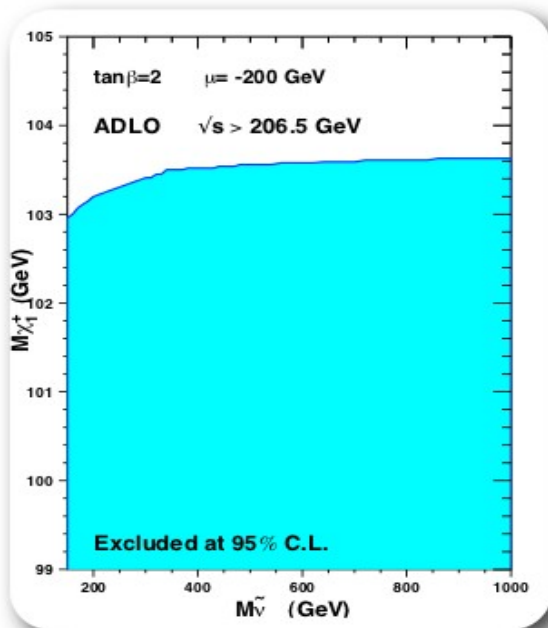
No mass limit in general outside these assumptions



Experimental constraints from LEP

Charginos via: s-channel γ/Z or t-channel with sneutrino exchange

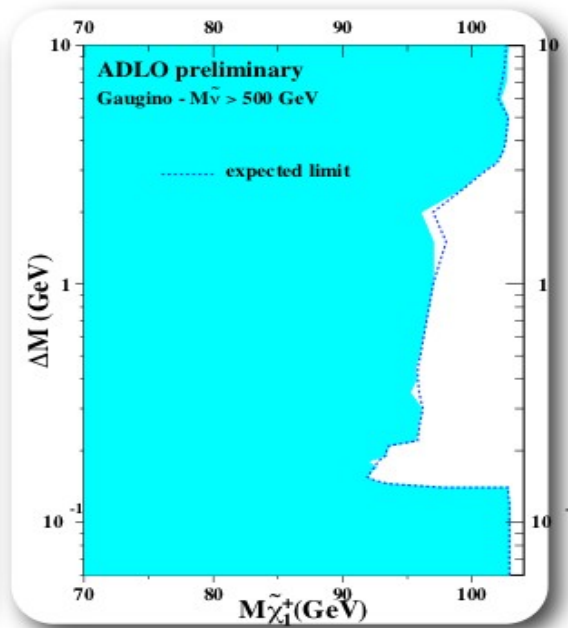
canonical case



**$m_{\tilde{\chi}_1^\pm} > 103.5$ GeV
for $m_{\text{sneue}} > 300$ GeV**

LEPSUSYWG/01-03.1

degenerate case



**$m_{\tilde{\chi}_1^\pm} > 91.9 / 92.4$
GeV**

LEPSUSYWG/02-04.1

Unification of gaugino masses at GUT scale is assumed.

$$- M1 = (5/3)\tan^2(\theta_W) M2 \sim 0.5 M2$$

Canonical case:

- With $M(\text{sneutrino}) > 300$ GeV

Degenerate case:

- $M1$ and $M2$ nearly degenerate

- Large $M0$ ($m(\text{snu}) \sim 500$ GeV)

In general Charginos up to ~ 100 GeV in mass are excluded by the LEP experiments

Experimental constraints from Tevatron

D0 Collaboration: $p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$

Three leptons + MET signature

- e, μ , and τ

4 Channels (eel, $\mu\mu l$, $e\mu\tau l$, $\mu\tau l$)

Dominant bkg: WZ, ZZ in MET tails

Within the context of MSUGRA

Assuming:

$$m_{\tilde{\chi}_1^\pm} \sim m_{\tilde{\chi}_2^0} \sim 2m_{\tilde{\chi}_1^0}$$

- and neglecting the slepton mixing

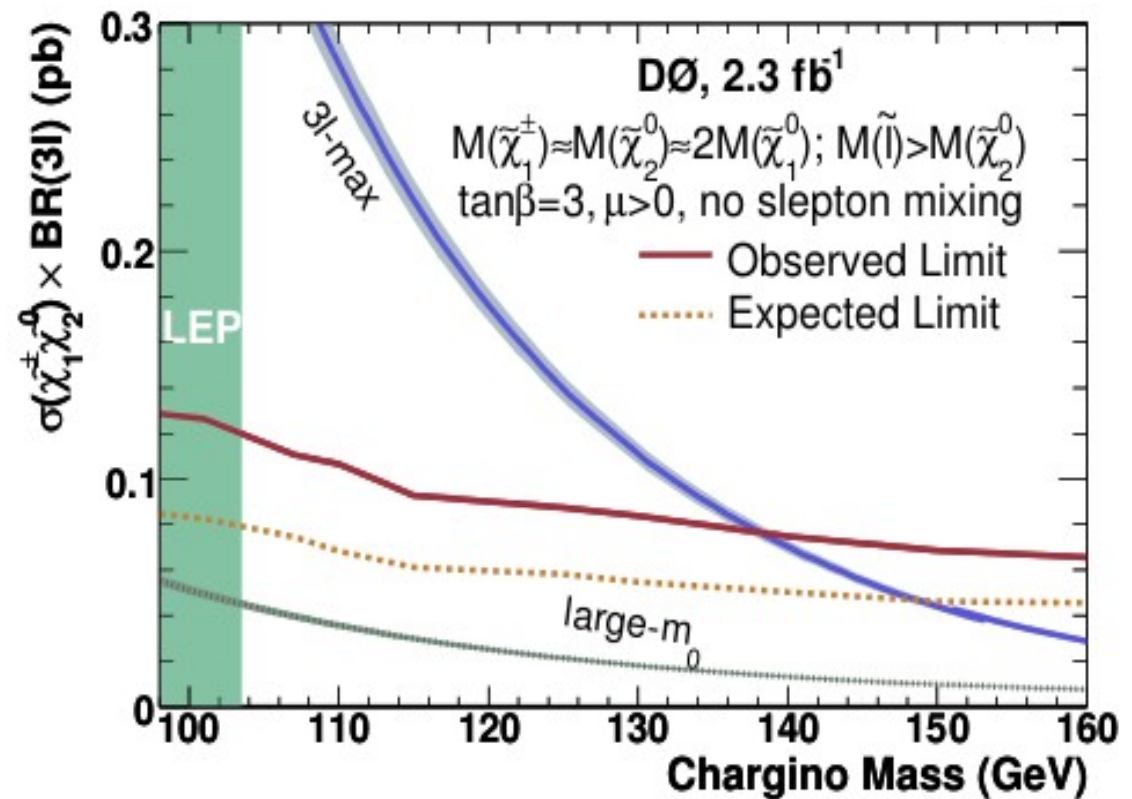
- sleptons and sneutrinos heavier than lightest charginos and next lightest neutralino

In the limit of heavy sleptons (large m_0 scenario):

- the slepton mass is just above mass of $\tilde{\chi}_2^0$, leptonic BR is maximized (3l max case)

Chargino mass < 138 GeV is excluded by this study

D0 Collab. Phys. Lett. B680, 34 (2009)



Experimental constraints from Tevatron

CDF Collaboration: $p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$

CDF Note: 10636

CDF Run II Preliminary (L=5.8 fb⁻¹)

Three leptons + MET signature

Several SRs in the plane - MET & M_{ll}

Modes:

- eel, μμl ; l = e, μ, τ (or single track)

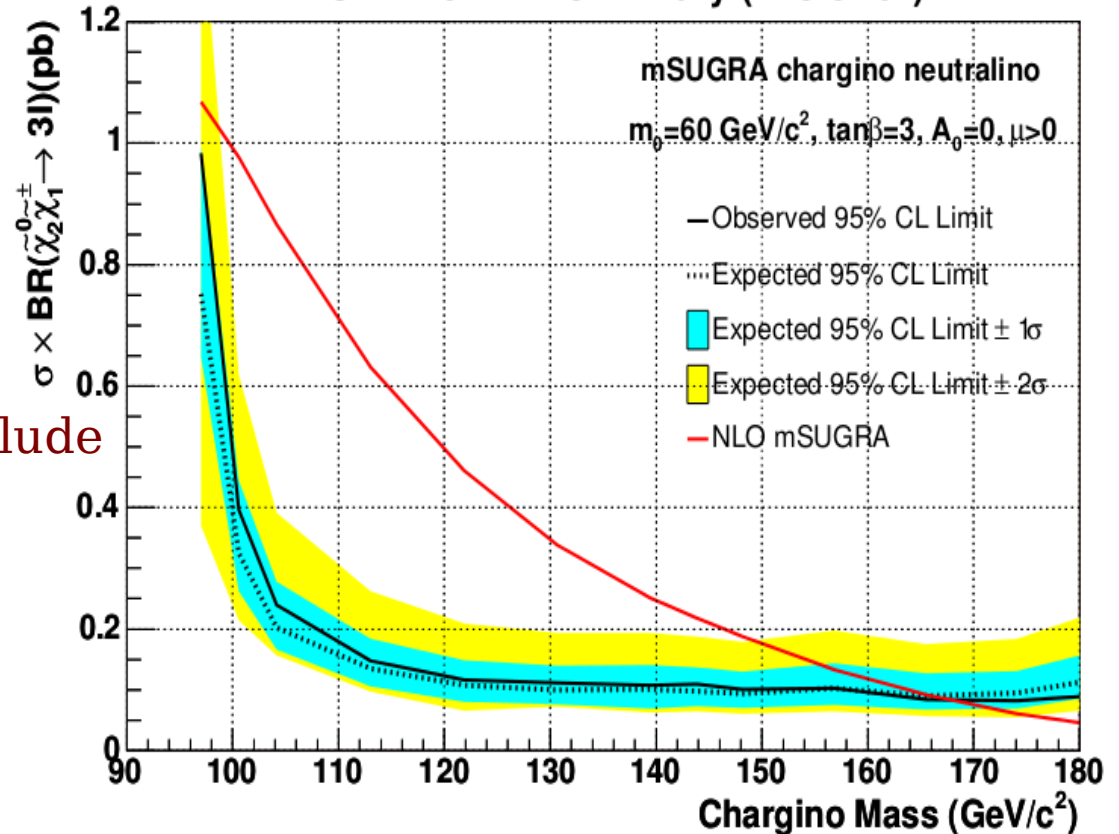
- Expanded the acceptance & also include

low p_T leptons ~ 5 GeV threshold.

Major backgrounds:

- WZ, ZZ, dileptons + fakes

Within the context of MSUGRA



Exclude at 95% CL $\sigma(\tilde{\chi}_1^\pm \tilde{\chi}_2^0) \times BR(lll)$ above 0.1 fb

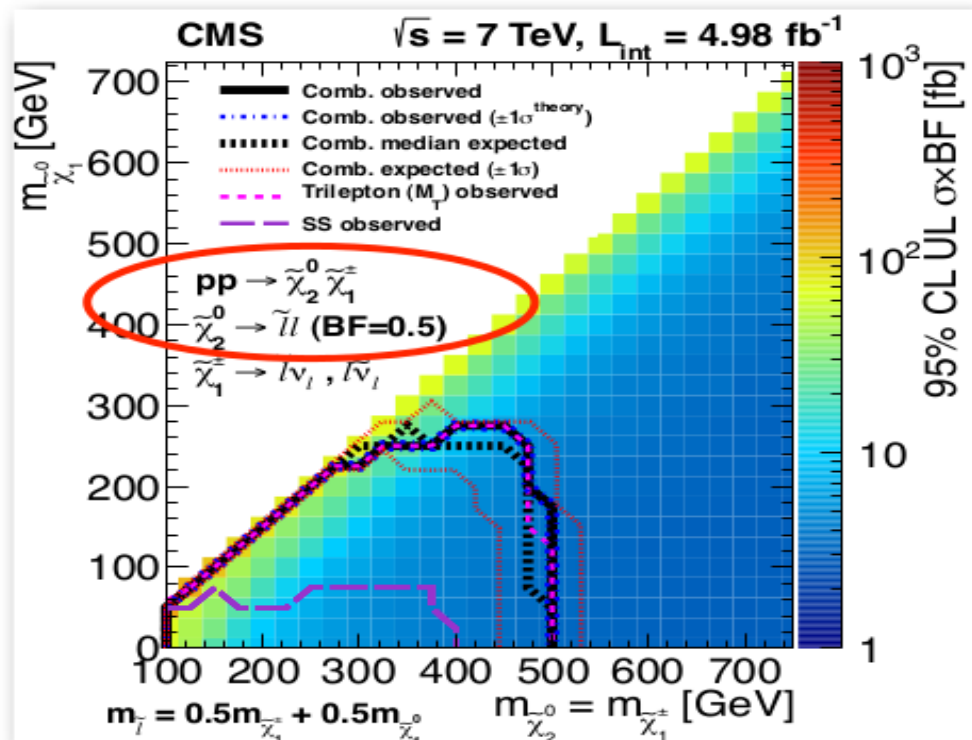
Chargino mass below 168 GeV is excluded by this study

Direct electroweak production at the LHC

Dominant modes $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$

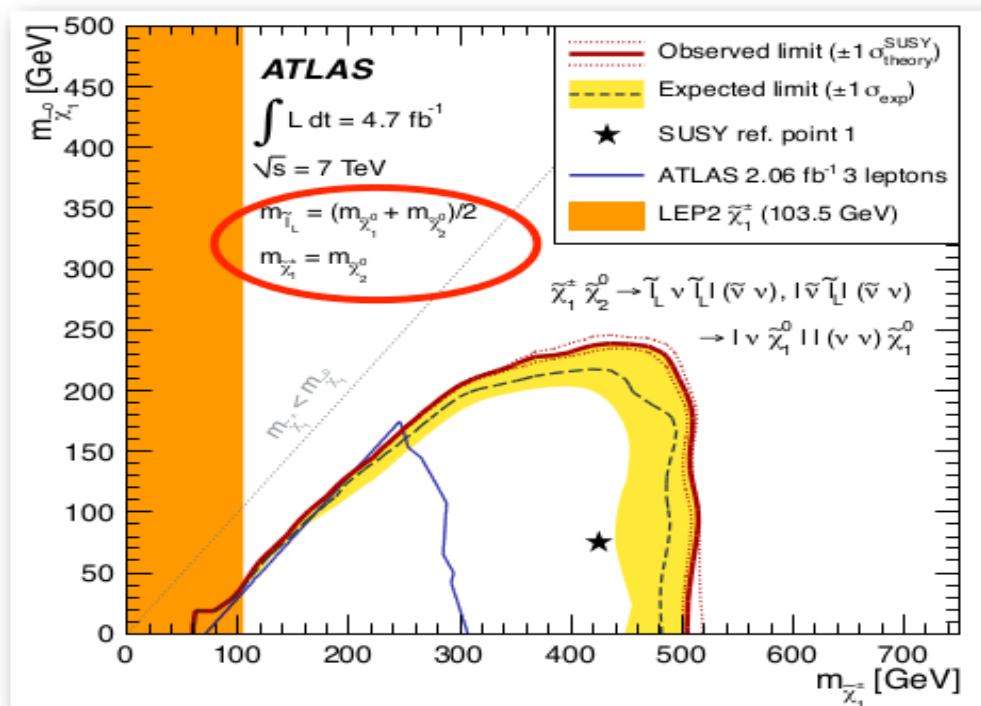
dilepton or trilepton + MET

CMS 1209.6620



trilepton + MET

ATLAS 1208.3144



60 GeV < m_{\tilde{\chi}_{1\pm}} < 500 GeV

Limits are weaker for - heavy slepton_L

- $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$ being Higgsinos

- small mass difference

Lepton rich final state to enhance reach

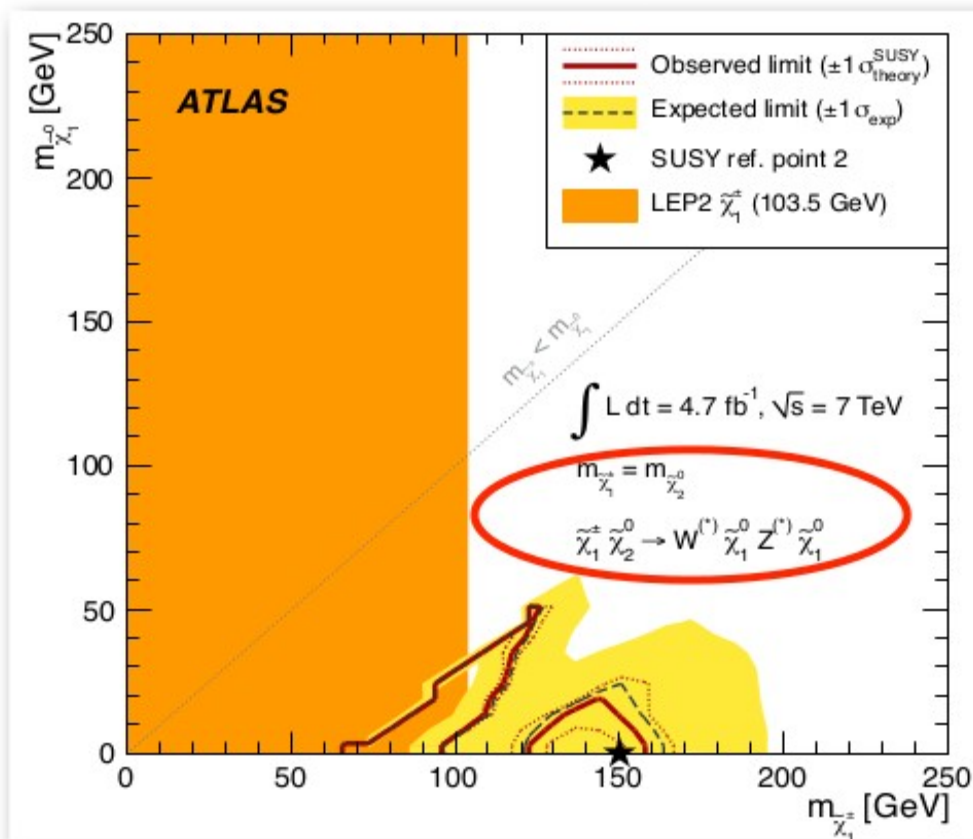
Direct electroweak production at the LHC

Dominant modes $\tilde{\chi}_2^0 \tilde{\chi}_1^+$

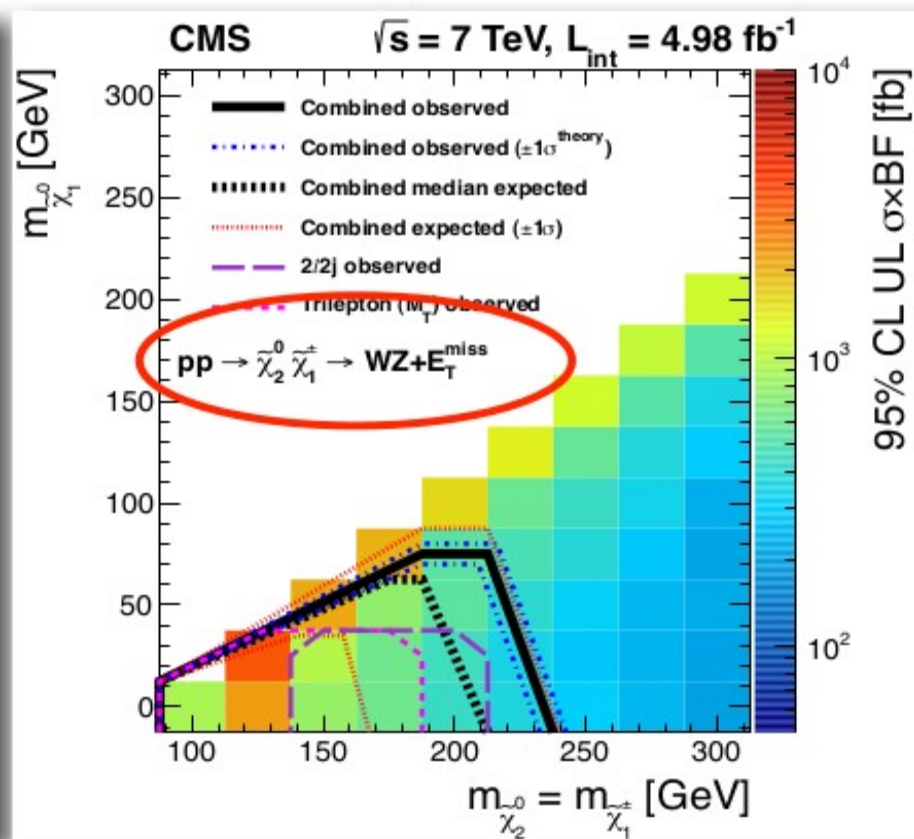
trilepton + MET

Much weaker limits!

dilepton/trilepton + MET



ATLAS 1208.3144



CMS 1209.6620

Excludes $\sim 200 \text{ GeV}$ for LSP mass of 50 GeV

100% BR - Usually not realized

Direct electroweak production

The remaining dominant contribution in EWKino (light Wino) sector is from: $\tilde{\chi}^+ \tilde{\chi}^-$

I think none of the LHC experiments have bounds in this mode

- *Re-interpretation/exclusion by M. Lisanti & N. Weiner (arXiv:1112.483)*
- Use $H \rightarrow WW$ (and $H \rightarrow ZZ$) results from ATLAS and CMS
- Simplified topology

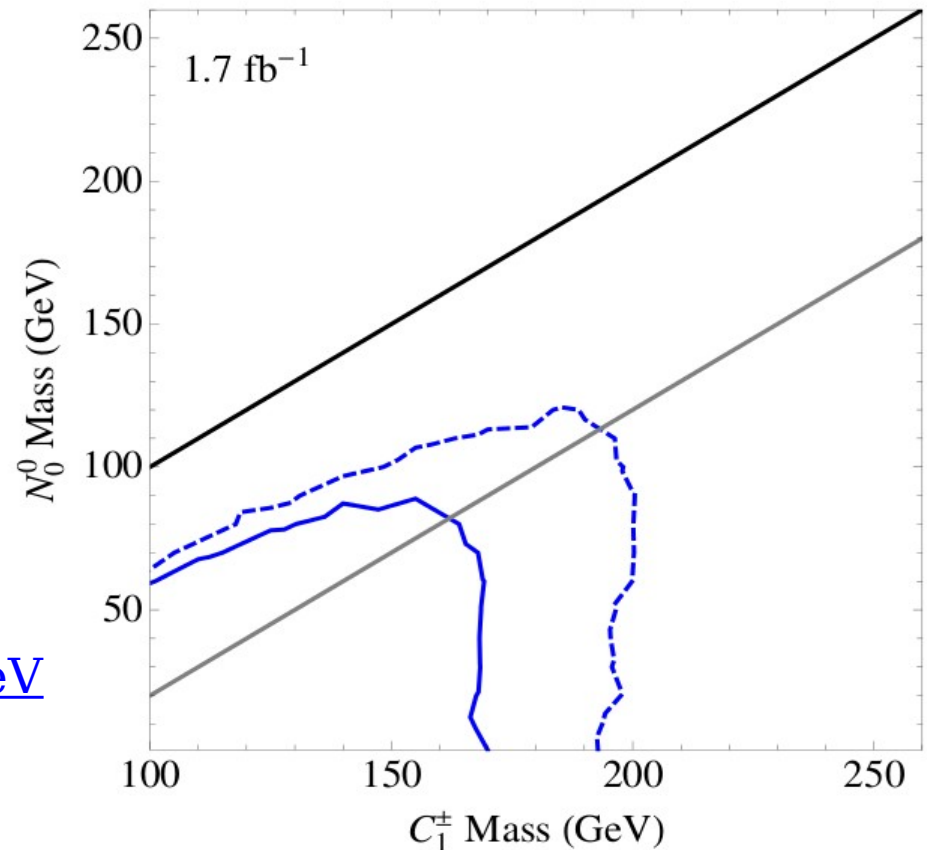
Dashed and solid lines indicate

- 1 and 2σ sensitivity

The sensitivity is defined as:

$$N_{\text{signal}} \geq 2 \times \sqrt{B_{\text{stat}}^2 + B_{\text{sys}}^2}.$$

[Charginos are most likely above ~200 GeV based on this study](#)



Charginos and Neutralinos in the light of the Higgs boson

In collaboration with T. Han (Pittsburg) and S. Su (Arizona)

Natural SUSY Charginos/Neutralinos

Assuming Higgs connection

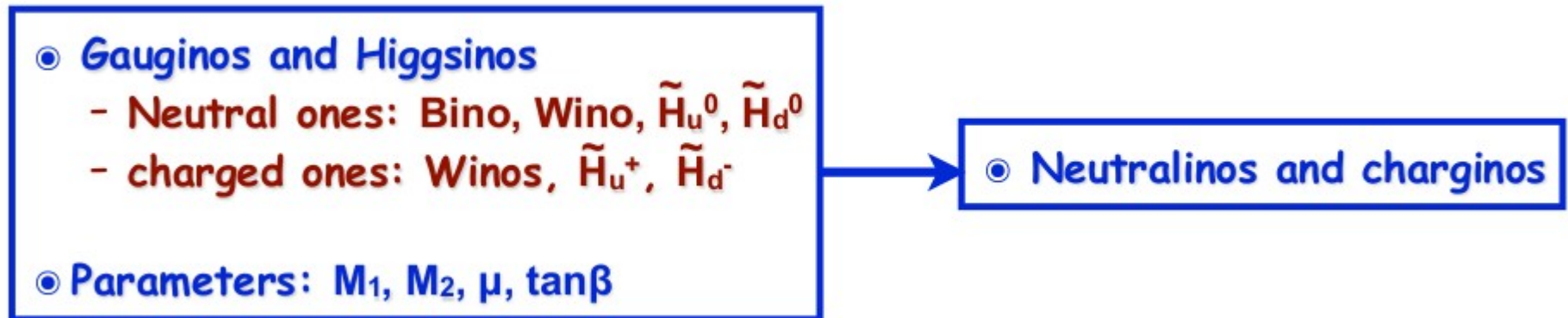
- Natural SUSY \rightarrow Light gauginos and Higgsinos

Colored superparticles might be heavy (See previous slides)

- Electroweak sector + stops/sbottoms might be the only accessible particles
- no indication from current LHC searches, $m_{sq}, m_{\text{gluino}} > 1 \text{ TeV}$

Connection to lepton collider

In MSSM :



Natural SUSY Charginos/Neutralinos

$$\frac{1}{2}M_Z^2 = \frac{(m_{H_d}^2 + \Sigma_d) - (m_{H_u}^2 + \Sigma_u) \tan^2 \beta}{(\tan^2 \beta - 1)} - \mu^2.$$

Assume LSP based on SUSY breaking mass parameters M_1 , M_2 and μ

- Decouple the SUSY colored sector

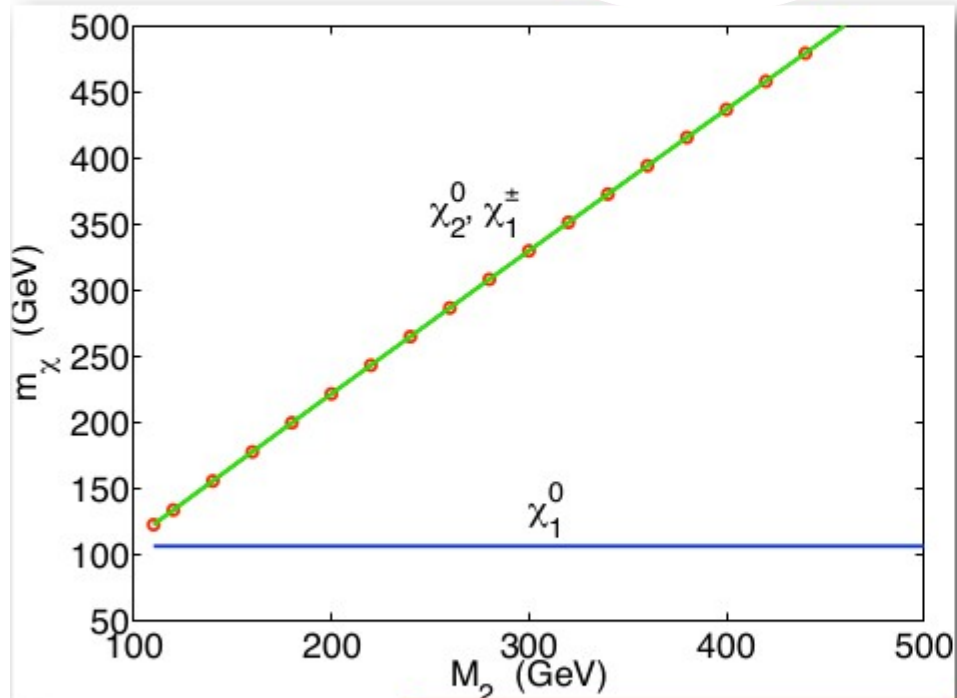
There can be three cases:

- a) Bino LSP ($M_1 < M_2, \mu$)
- b) Wino LSP ($M_2 < M_1, \mu$)
- c) Higgsino LSP ($\mu < M_1, M_2$)

Masses: Bino LSP

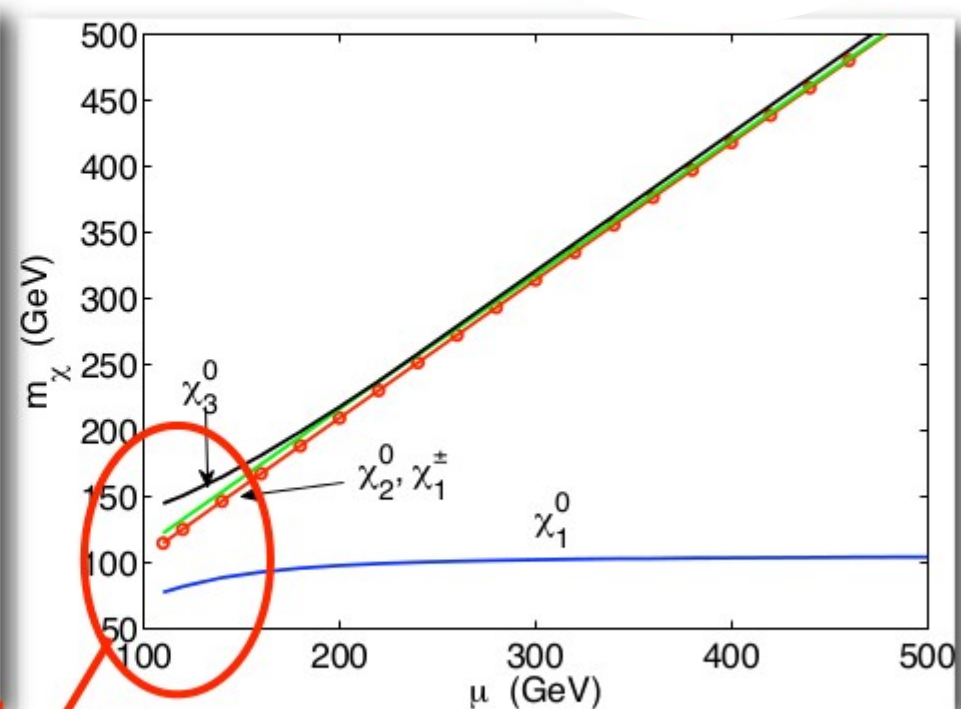
Case AI:
 $M_1 < M_2 < \mu$

$\mu = 1 \text{ TeV}$



Case AII:
 $M_1 < \mu < M_2$

$M_2 = 1 \text{ TeV}$



**large mixing, natural
 compressed spectrum**

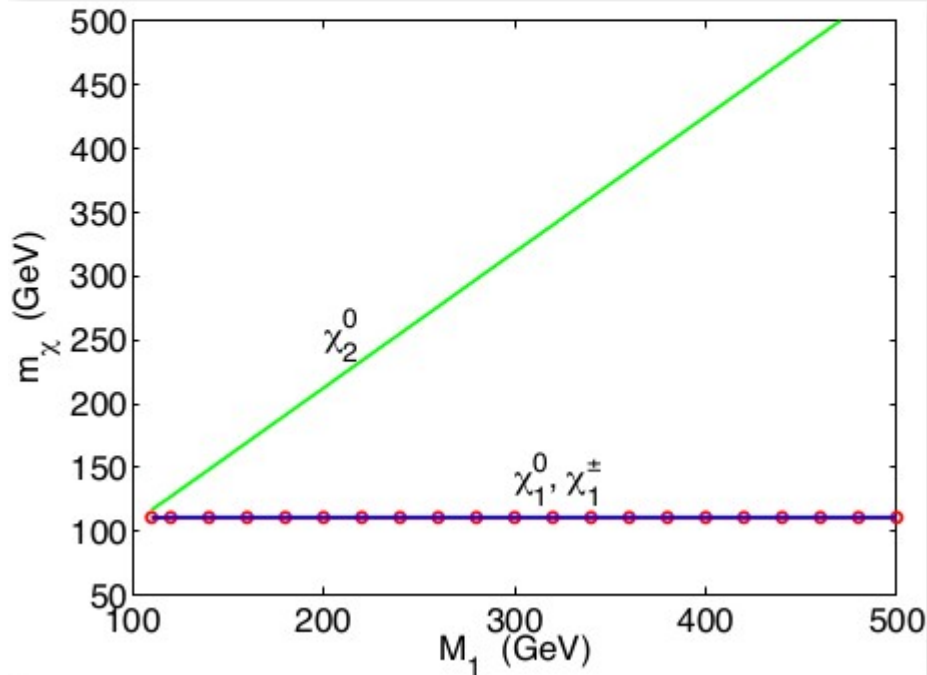
Case AI : $M_2 < \mu$, χ_1^\pm, χ_2^0 are Wino – like; $\chi_2^\pm, \chi_{3,4}^0$ are Higgsino – like;

Case AII : $\mu < M_2$, $\chi_1^\pm, \chi_{2,3}^0$ are Higgsino – like, χ_2^\pm, χ_4^0 are Wino – like.

Masses: Wino LSP

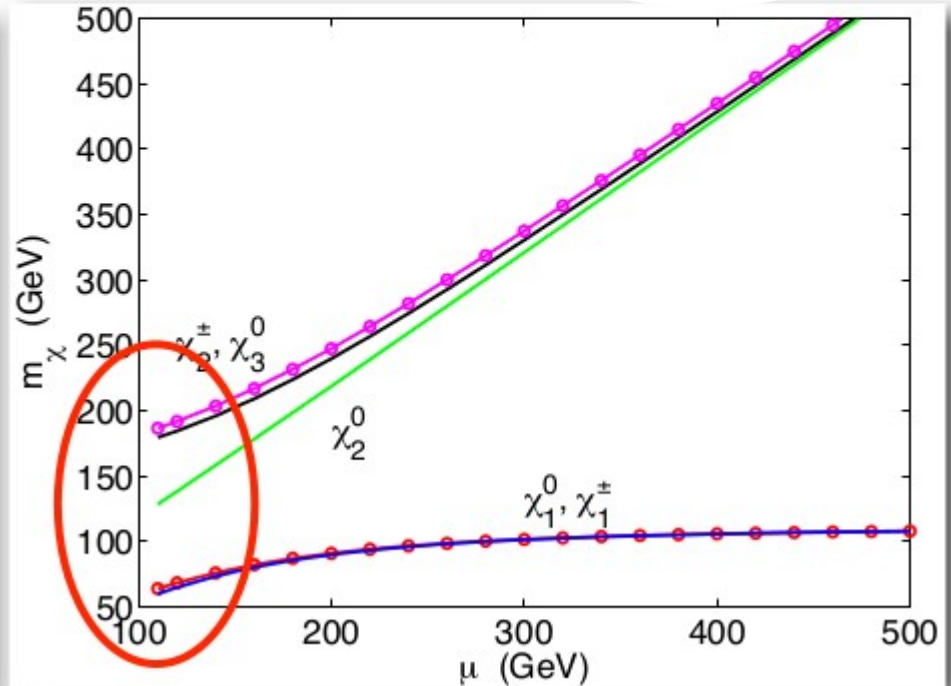
Case BI:
 $M_2 < M_1 < \mu$

$\mu = 1 \text{ TeV}$



Case BII:
 $M_2 < \mu < M_1$

$M_1 = 1 \text{ TeV}$



With wino LSP:

Case BI : $M_1 < \mu$, χ_2^0 Bino – like; $\chi_2^\pm, \chi_{3,4}^0$ Higgsino – like;

Case BII : $\mu < M_1$, $\chi_{2,3}^\pm, \chi_{2,3}^0$ Higgsino – like; χ_4^0 Bino – like.

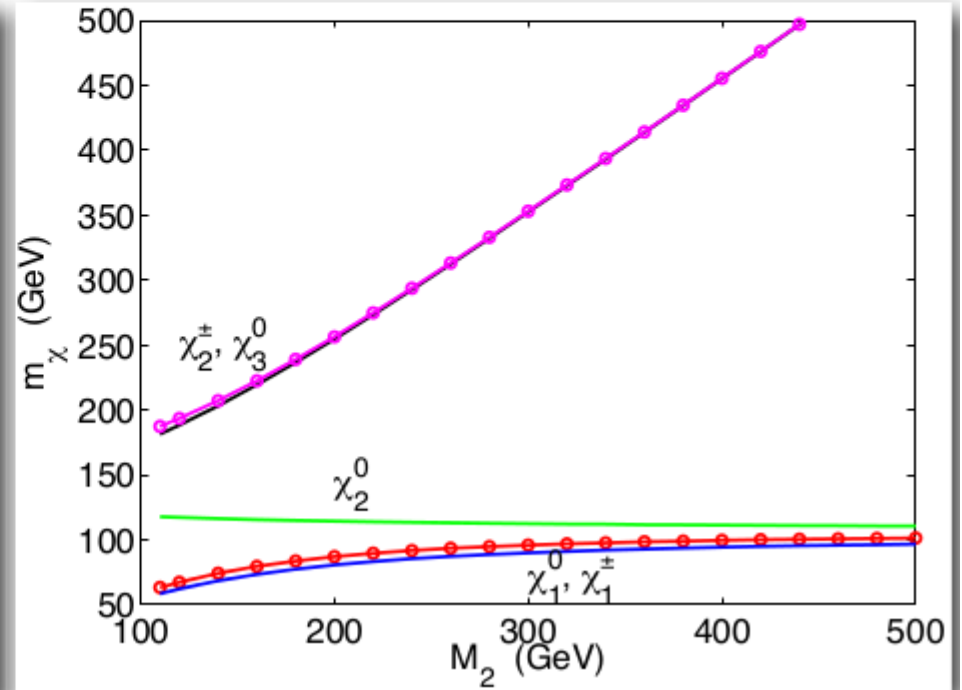
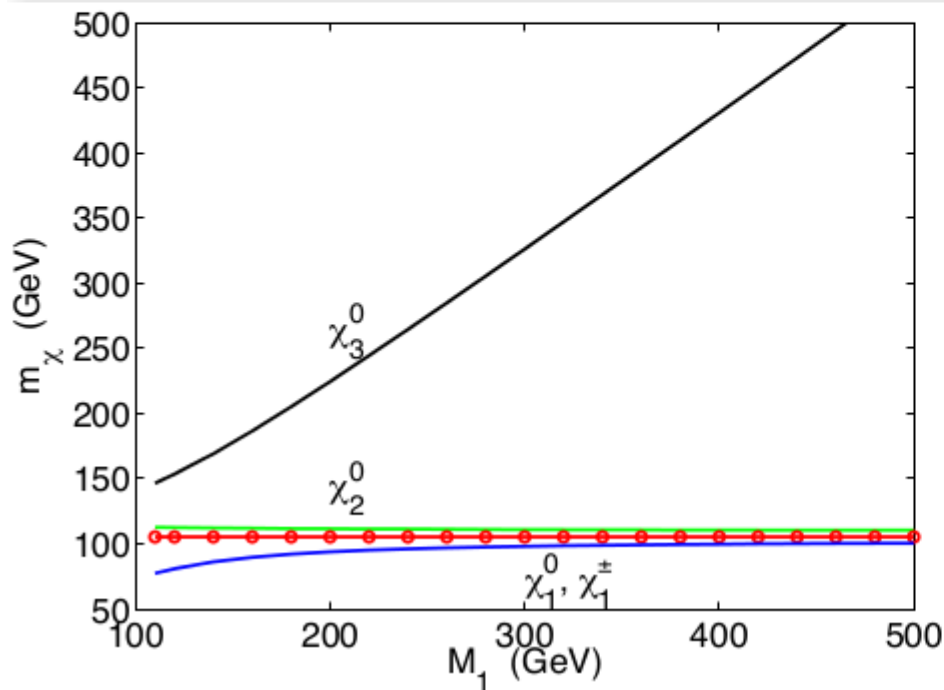
Masses: Higgsino LSP

Case CI:
 $\mu < M_1 < M_2$

Case CII:
 $\mu < M_2 < M_1$

$M_2 = 1 \text{ TeV}$

$M_1 = 1 \text{ TeV}$

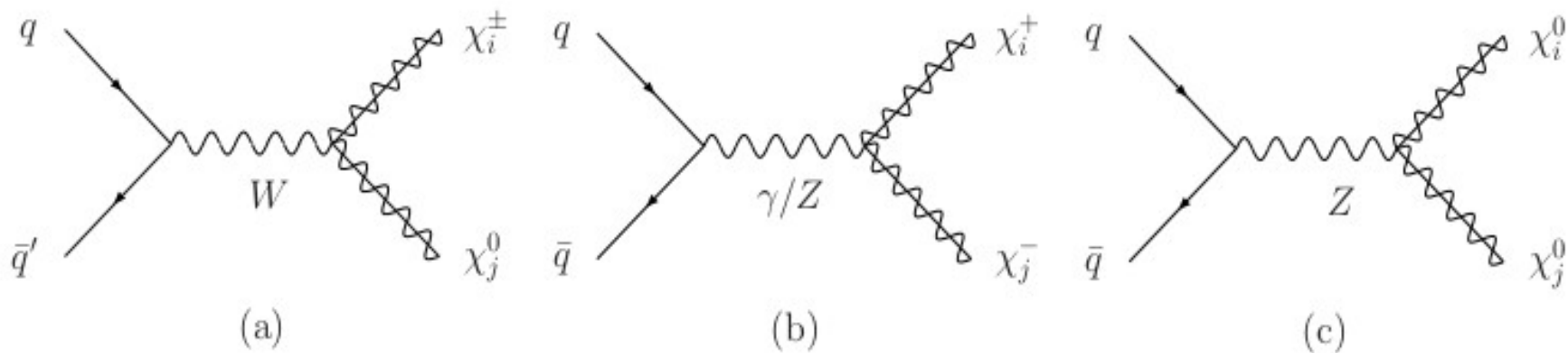


With higgsino LSP:

Case CI : $M_1 < M_2$, χ_3^0 Bino – like; χ_2^\pm, χ_4^0 Wino – like;

Case CII : $M_2 < M_1$, χ_2^\pm, χ_3^0 Wino – like; χ_4^0 Bino – like.

Productions of SUSY weak sector

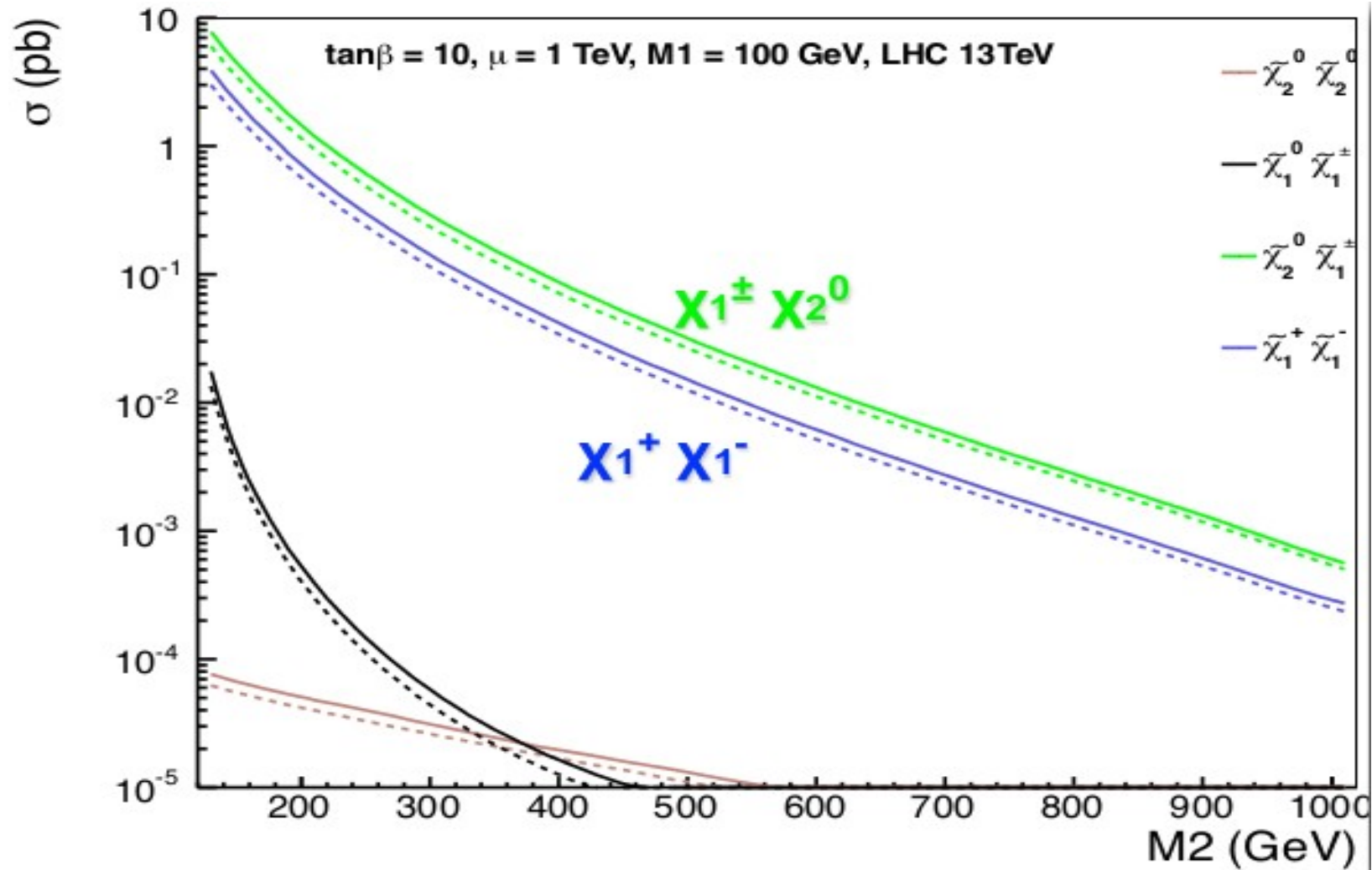


Dominant production:

- Wino pair production: $\chi_i^+ \chi_j^-, \chi_i^\pm \chi_j^0$
- Higgsino pair production: $\chi_i^+ \chi_j^-, \chi_i^\pm \chi_j^0, \chi_i^0 \chi_j^0$

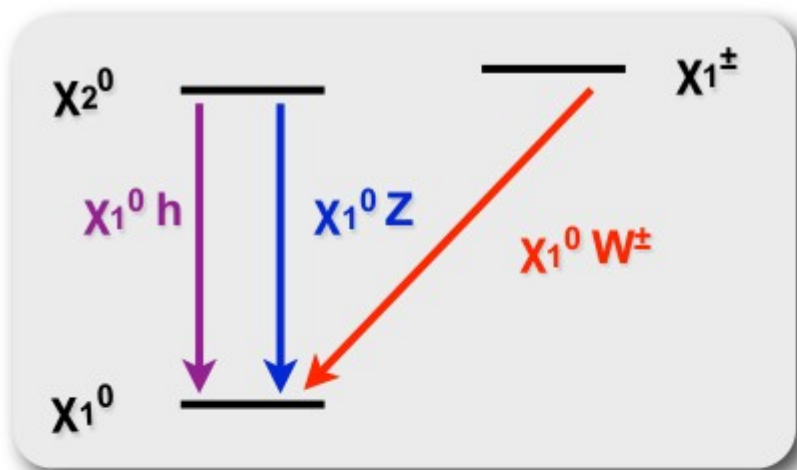
Productions of Bino LSP, Wino NLSP

Case A1: $M_1 < M_2 < \mu$



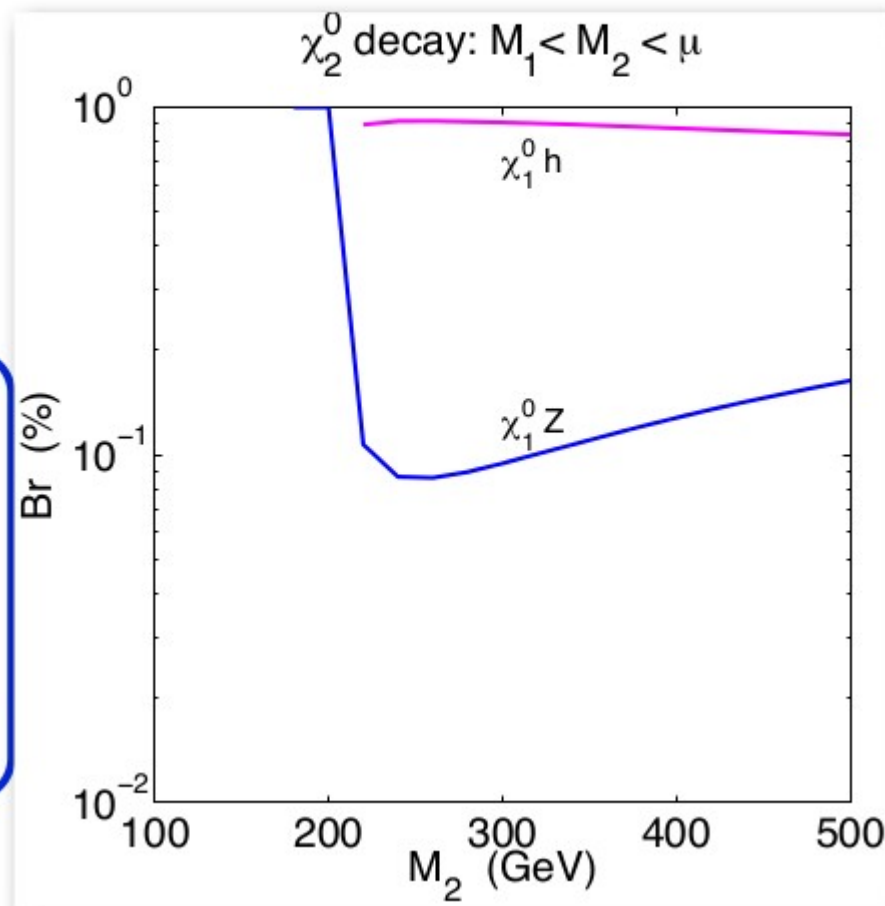
Dominant contributions are from: $pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0 X, \tilde{\chi}_1^+ \tilde{\chi}_1^- X$

Decays with Bino LSP, Wino NLSP



χ_{1^\pm} decay 100% via on/off-shell W

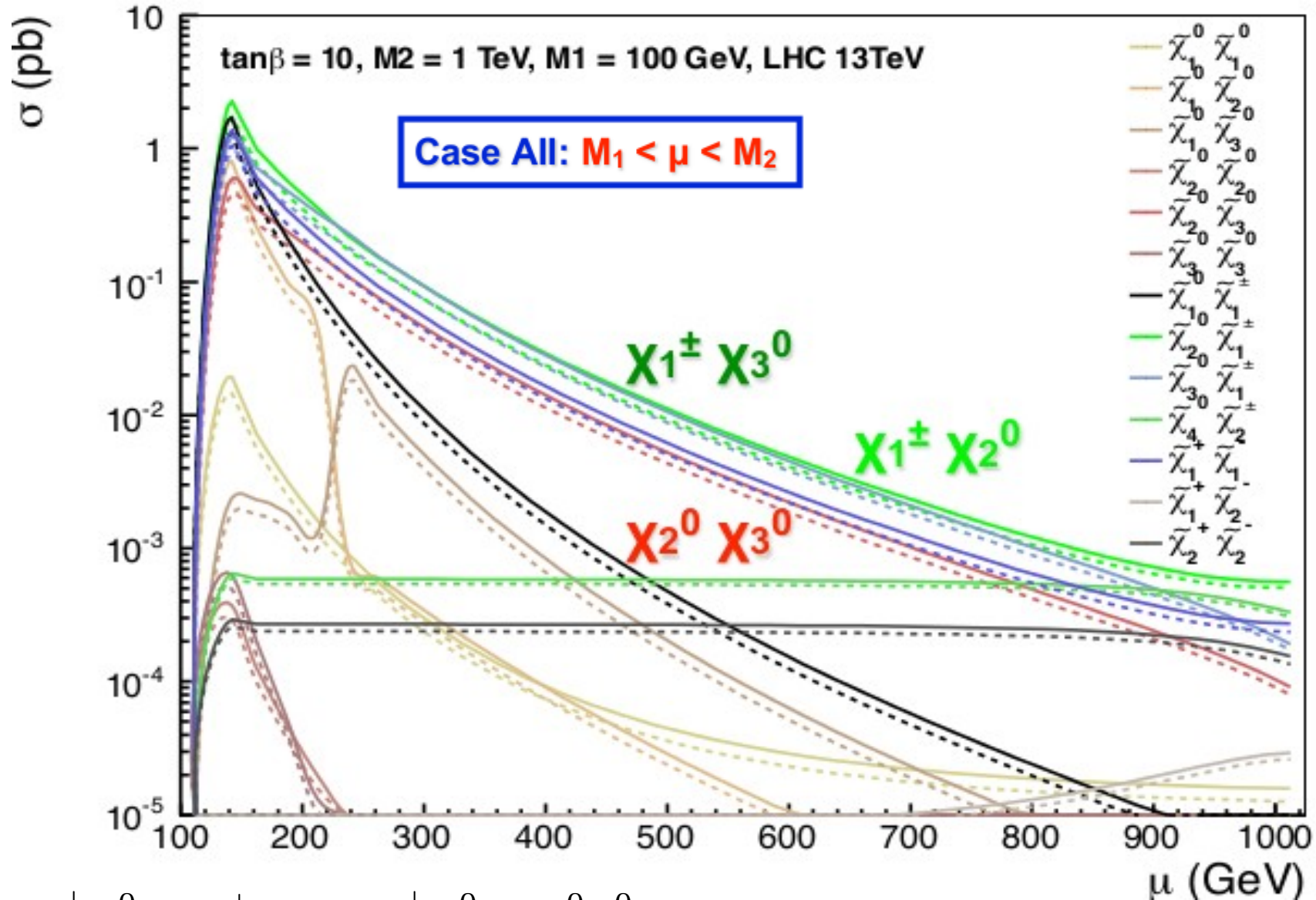
- below h threshold, decay via on/off-shell Z
- χ_2^0 on-shell decay to h dominate over on-shell Z for $\mu > 0$
- χ_2^0 decay to h and Z flipped for $\mu < 0$



Dominant contributions are from:

$$pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0 X, \tilde{\chi}_1^+ \tilde{\chi}_1^- X; \chi_{1^\pm} \rightarrow W^\pm \chi_1^0, \chi_2^0 \rightarrow (h/Z) \chi_1^0$$

Productions of Bino LSP, Higgsino NLSP



$$pp \rightarrow \chi_{1^\pm} \chi_{2^0} X, \chi_{1^+} \chi_{1^-} X, \chi_{1^\pm} \chi_{3^0} X, \chi_{2^0} \chi_{3^0} X$$

$$\chi_{1^\pm} \rightarrow W^\pm \chi_1^0; \chi_{2^0} \rightarrow (h/Z) \chi_1^0, \chi_{3^0} \rightarrow (Z/h) \chi_1^0$$

For details on other LSP (M₂, μ)
 → See the upcoming paper

SUSY weak productions

Final states that can be explored:

WW, WZ, Wh, Zh, ZZ, hh

- BR(WZ) < 100% in most cases, sometimes highly suppressed
- Wh complementary to WZ channel : a new discovery mode
- Zh/hh could also be explored.

Experimentally challenging depending on “compression” between the mass states:
(e.g: Also the depends on the choice of the LSP)

- If the mass difference is in MeV: $\chi_2^0 - \chi_1^0$ or $\chi_1^\pm - \chi_1^0$
 - Expect “appearing tracks” within few cms if the associated particle is neutral
 - Expect highly ionizing tracks (dE/dx) associated with charged particle
- If the mass difference is in GeV → prompt decays

SUSY weak productions

In terms of searches:

1. If both parents are un-compressed:

- Standard analysis, trigger on any or both of the visible decay products

2. If one of the parents is compressed e.g: $\chi_2^0 \chi_1^\pm$; $M(\chi_1^\pm) \approx M(\chi_1^0)$

- Use trigger based on one visible decay product

3. If both parents are compressed

- e.g: $\chi_1^+ (\rightarrow W \chi_1^0) \chi_1^- (\rightarrow W \chi_1^0)$; $M(\chi_1^\pm) \approx M(\chi_1^0)$

- Use mono-jet kind of analysis with trigger on ISR jets (Parked data?)

Possible future LHC searches with Higgs in the final state

● Wh channel: 1l+jets + MET

- Isolated $e(\mu)$, $P_t > 30(20)$ GeV, $|\eta| < 2.5$
- Veto any additional e/μ with $P_t > 10$ GeV, $|\eta| < 2.5$
- Veto any Taus or isolated Tracks
- 2 Jets $P_t > 30$ GeV, $|\eta| < 2.5$
- Veto 3rd Jet with $P_t > 20$ GeV
- 2 bjets with $P_t > 30$ GeV, $|\eta| < 2.5$
- 2 bjets in one hemi-sphere ←
- Invariant mass of two bjets $100 < M_{bb}$ (GeV) < 140
- M_T (MET and the Higgs) > 200 GeV
- $MET > 50$ GeV

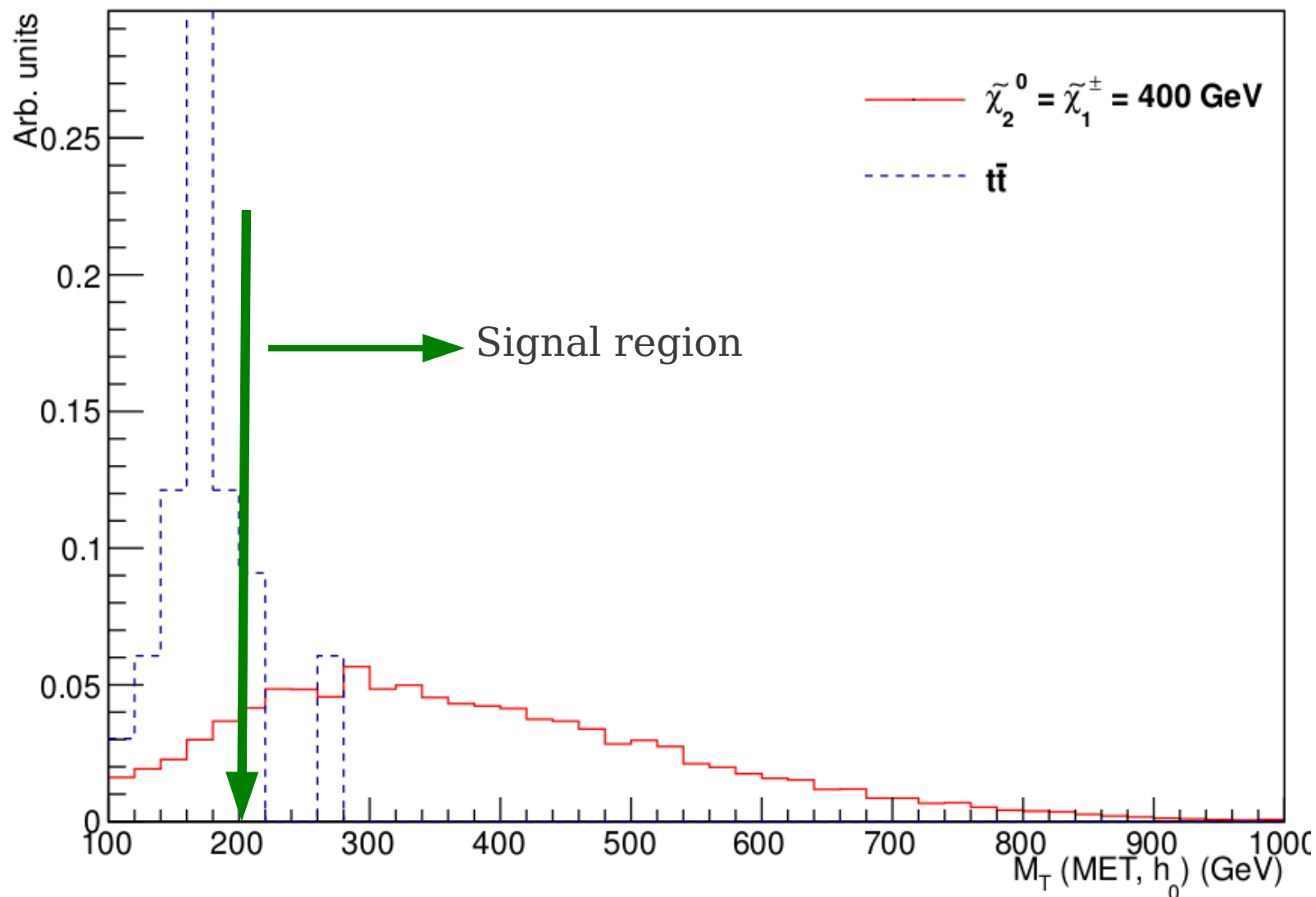
Signal regions:

$(M_T, MET) > (200, 50), (600, 50), (200, 100), (600, 100)$ GeV

10 fb⁻¹

| Processes | MET > 50, MT > 200 (Baseline) | MET > 50, 200 < MT < 400 | MET > 50, 400 < MT < 600 | MET > 50, MT > 600 | MET > 100, MT > 200 | MET > 100, 200 < MT < 400 | MET > 100, 400 < MT < 600 | MET > 100, MT > 600 |
|-----------|-------------------------------|--------------------------|--------------------------|--------------------|---------------------|---------------------------|---------------------------|---------------------|
| Total bg | 46.15 ± 12.01 | 43.27 ± 11.96 | 2.40 ± 1.02 | 0.48 ± 0.48 | 33.63 ± 10.69 | 30.75 ± 10.63 | 2.40 ± 1.02 | 0.49 ± 0.48 |

Possible LHC searches with Higgs in the final state

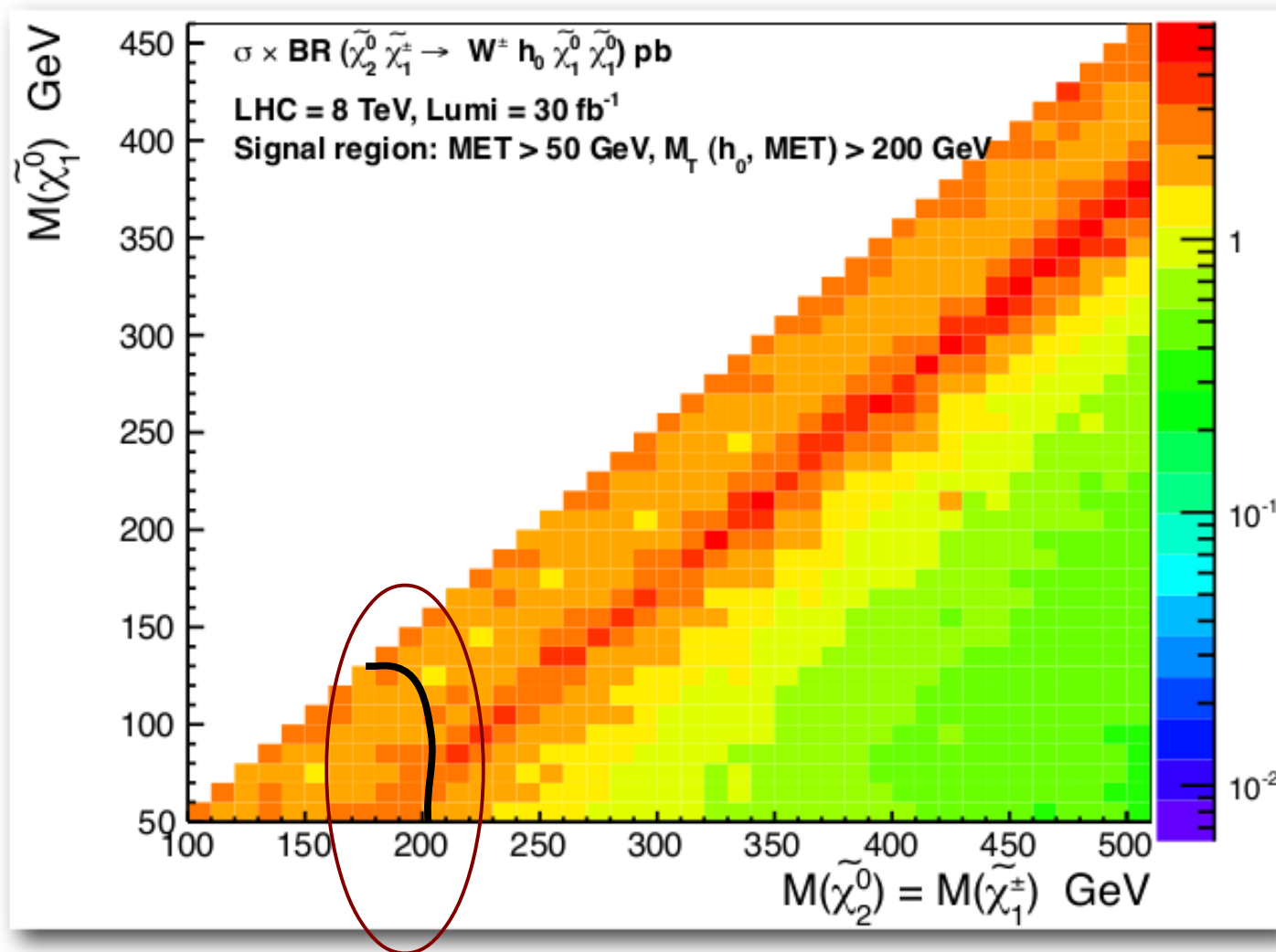


- Higgs Transverse Mass (Event simulation using Delphes)
- Background dominated by $t\bar{t}$ events

Possible LHC searches with Higgs in the final state

Wh: 1l+jets + MET

95% C.L. upper limit on signal cross section



With background only hypothesis, one can be sensitive to ~ 200 GeV in mass

Summary and Conclusion

Naturalness in SUSY can be valuable guiding principle for current/future searches

SUSY results from ATLAS and CMS show the breath of physics analyses

SUSY searches from the LHC

- constraints both squarks and gluinos up to TeV scale (with assumptions)
- the direct stop/sbottom limits up to ~ 500 GeV in mass
- the constraints on direct electroweak productions are soft.

Discovery of Higgs is just a starting point to move into a new territory

Search for new physics with Higgs in the final state

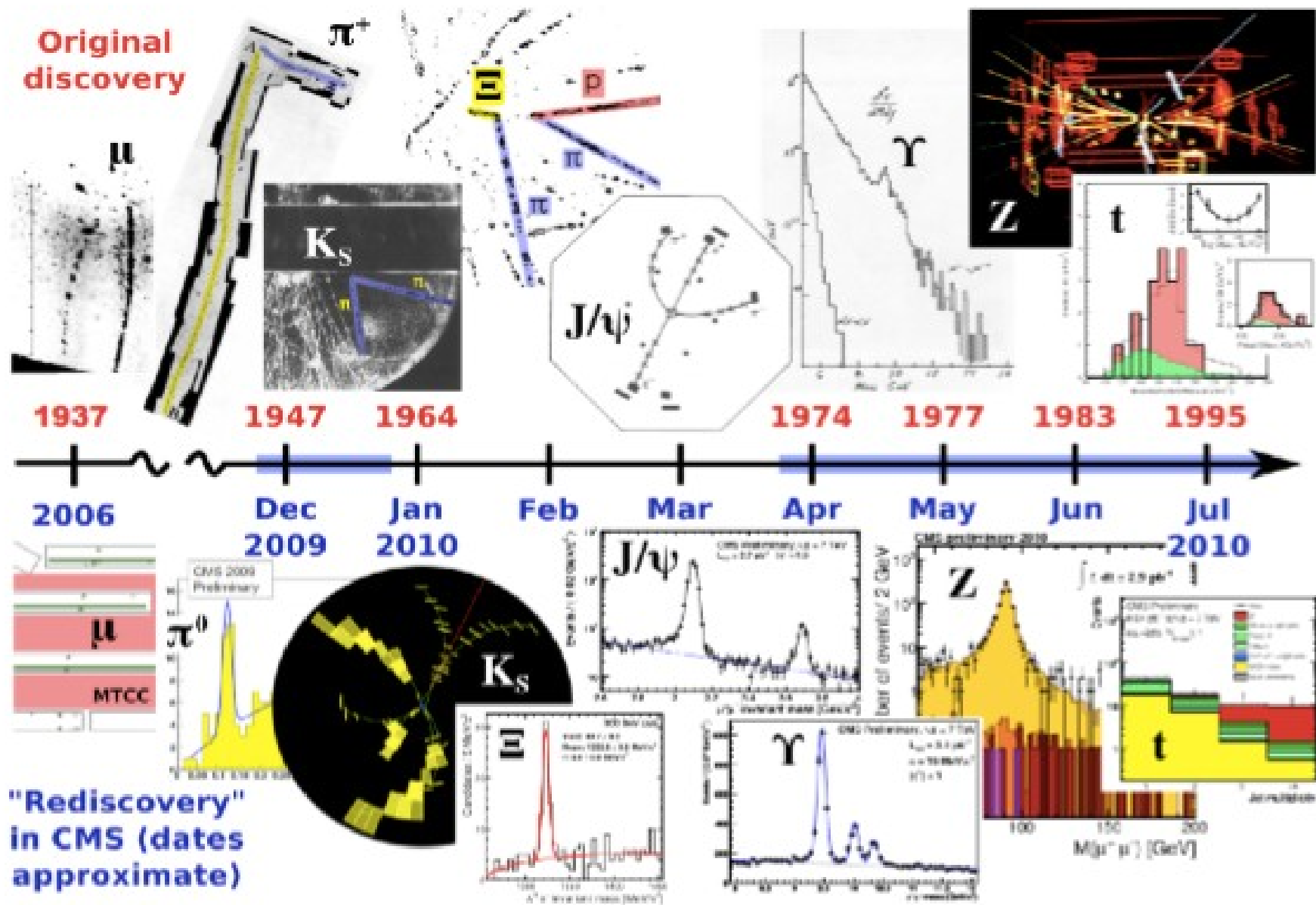
- Can be the next discovery mode ...

Studies towards naturally compressed spectra

→ Essential for next phase of LHC studies

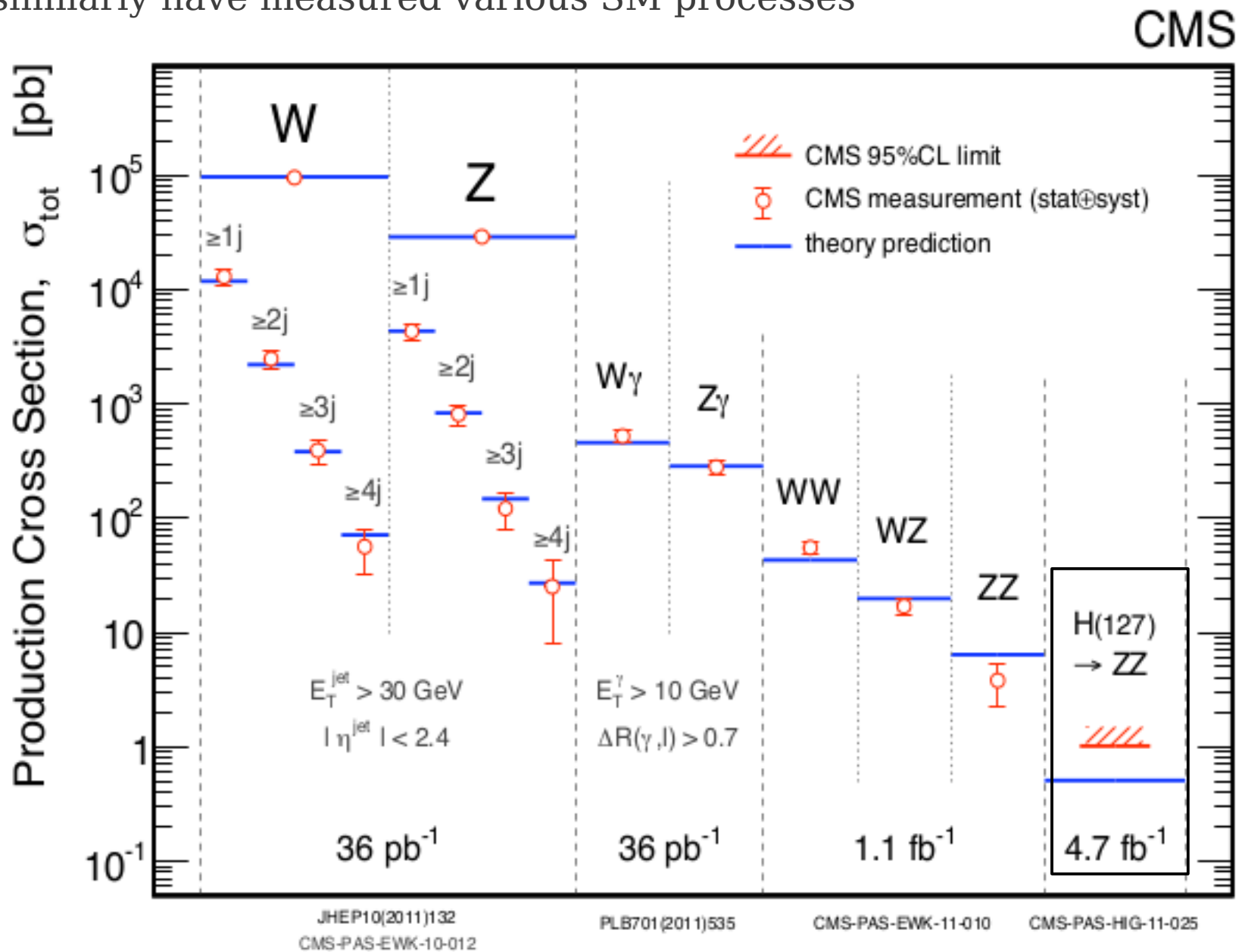
Backup slides

Re-Discovery of the Standard Model



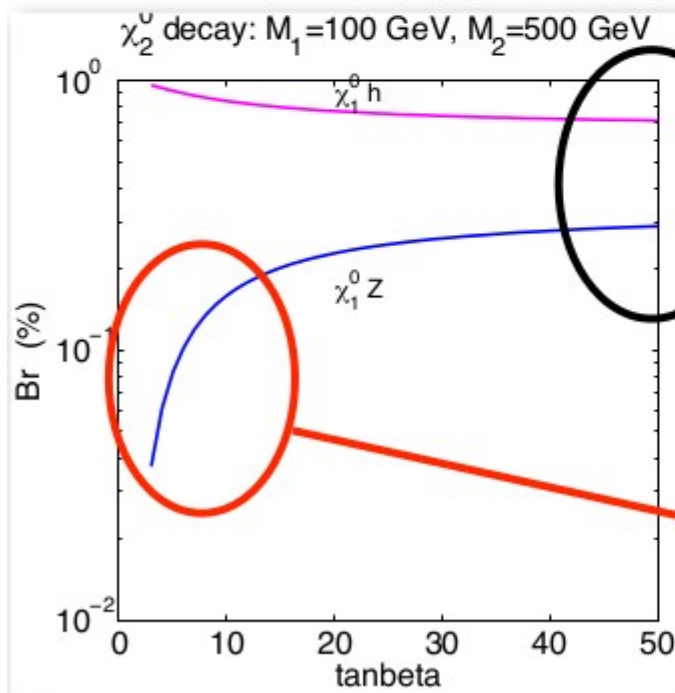
Standard Model Measurements

ATLAS similarly have measured various SM processes



Tanbeta dependency

- decay occur via mixing through Higgsino
- $M_2 \gg M_1$, $\chi_2^0 \rightarrow \chi_1^0 Z$ dominated by the decay via Z_L (goldstone mode G^0)
- h, G^0 as mixture of H_u^0 and H_d^0



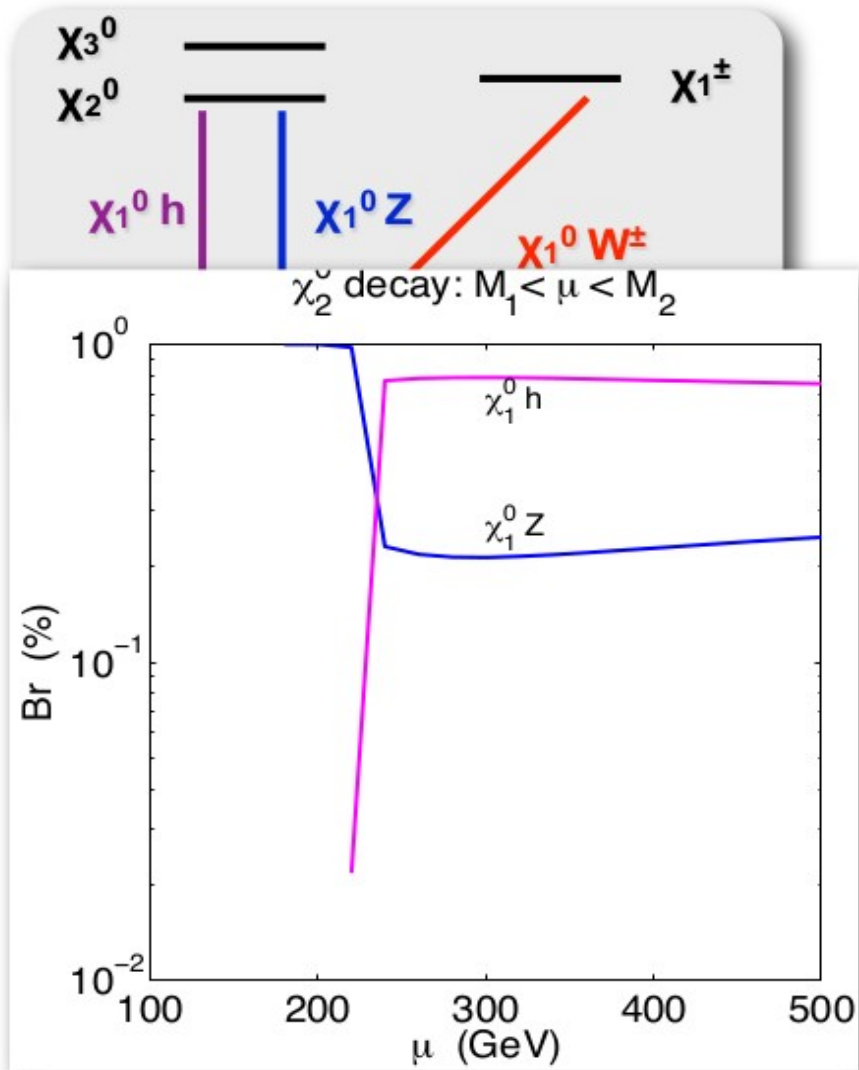
$$\Gamma(\chi_2^0 \rightarrow \chi_1^0 h) \propto \left(2s_{2\beta} + \frac{M_2}{\mu} \right)^2 \left[(M_2 + M_1)^2 - m_h^2 \right],$$

$$\Gamma(\chi_2^0 \rightarrow \chi_1^0 Z) \propto \left(c_{2\beta} \frac{M_2}{\mu} \right)^2 \left[(M_2 - M_1)^2 - m_Z^2 \right].$$

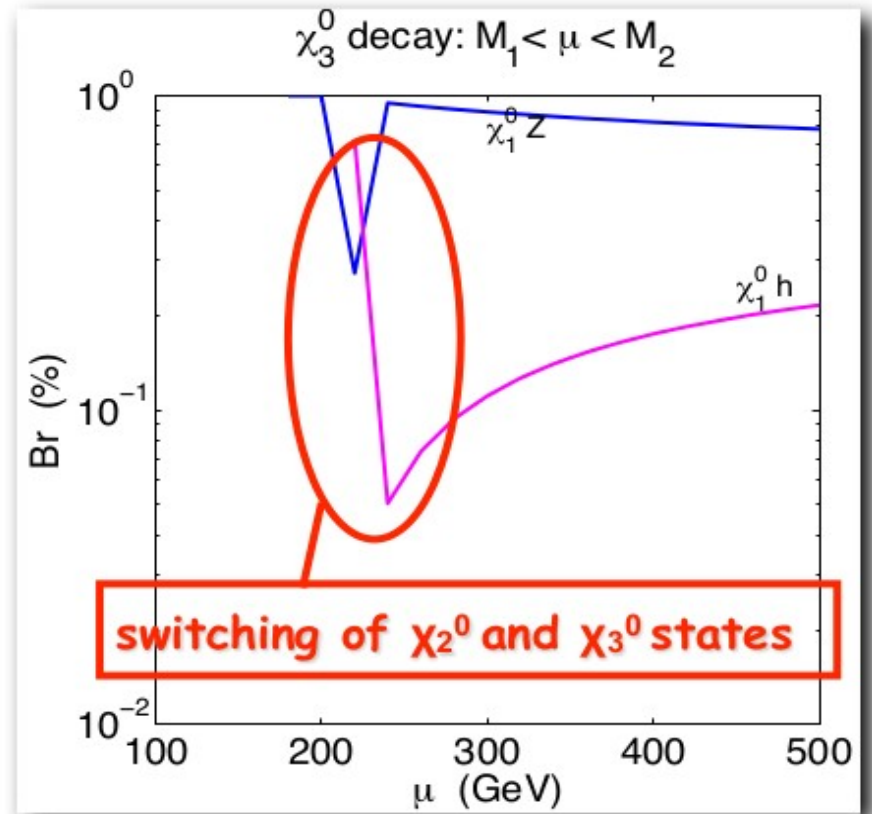
large $\tan\beta$, $[(M_2 + M_1)^2 - m_h^2] / [(M_2 - M_1)^2 - m_h^2]$

small $\tan\beta$, Z channel relatively suppressed

Case AII: Bino LSP-Higgsino NLSP



χ_1^\pm decay 100% via on/off-shell W



Neutralinos

• Neutralinos

$$\psi^0 = (\tilde{B}, \tilde{W}^0, \tilde{H}_d^0, \tilde{H}_u^0)$$

$$M_{\tilde{N}} = \begin{pmatrix} M_1 & 0 & -c_\beta s_W m_Z & s_\beta s_W m_Z \\ 0 & M_2 & c_\beta c_W m_Z & -s_\beta c_W m_Z \\ -c_\beta s_W m_Z & c_\beta c_W m_Z & 0 & -\mu \\ s_\beta s_W m_Z & -s_\beta c_W m_Z & -\mu & 0 \end{pmatrix},$$

| | | |
|---------|-----------------|---|
| M_1 | Bino | $\begin{pmatrix} \chi_1^0 \\ \chi_2^0 \\ \chi_3^0 \\ \chi_4^0 \end{pmatrix} = \begin{pmatrix} 1 & \mathcal{O}\left(\frac{m_Z}{M}\right) & \mathcal{O}\left(\frac{m_Z}{M'}\right) & \mathcal{O}\left(\frac{m_Z}{M}\right) \\ \mathcal{O}\left(\frac{m_Z}{M}\right) & 1 & \mathcal{O}\left(\frac{m_Z}{M}\right) & \mathcal{O}\left(\frac{m_Z}{M}\right) \\ \mathcal{O}\left(\frac{m_Z}{M}\right) & \mathcal{O}\left(\frac{m_Z}{M}\right) & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \mathcal{O}\left(\frac{m_Z}{M}\right) & \mathcal{O}\left(\frac{m_Z}{M}\right) & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} \tilde{B} \\ \tilde{W}^0 \\ \tilde{H}_d^0 \\ \tilde{H}_u^0 \end{pmatrix}$ |
| M_2 | Wino | |
| $ \mu $ | Higgsino | |
| $ \mu $ | Higgsino | |

Charginos

● Charginos

$$\psi^\pm = (\tilde{W}^+, \tilde{H}_u^+, \tilde{W}^-, \tilde{H}_d^-)$$

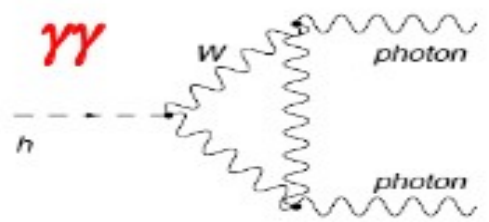
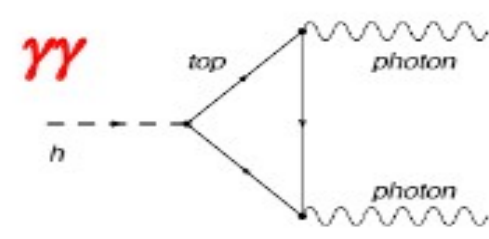
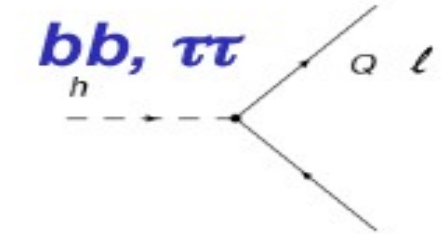
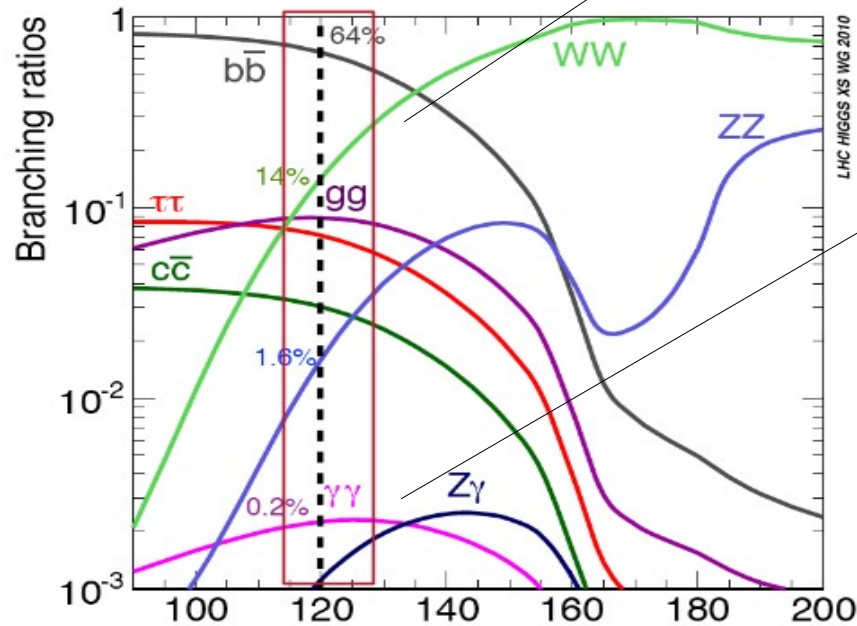
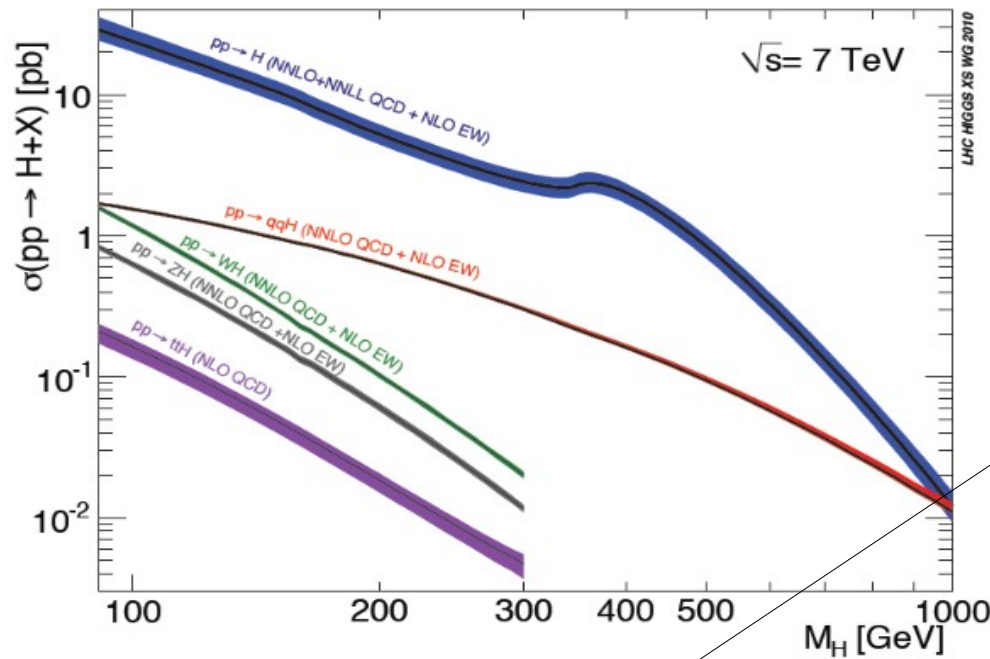
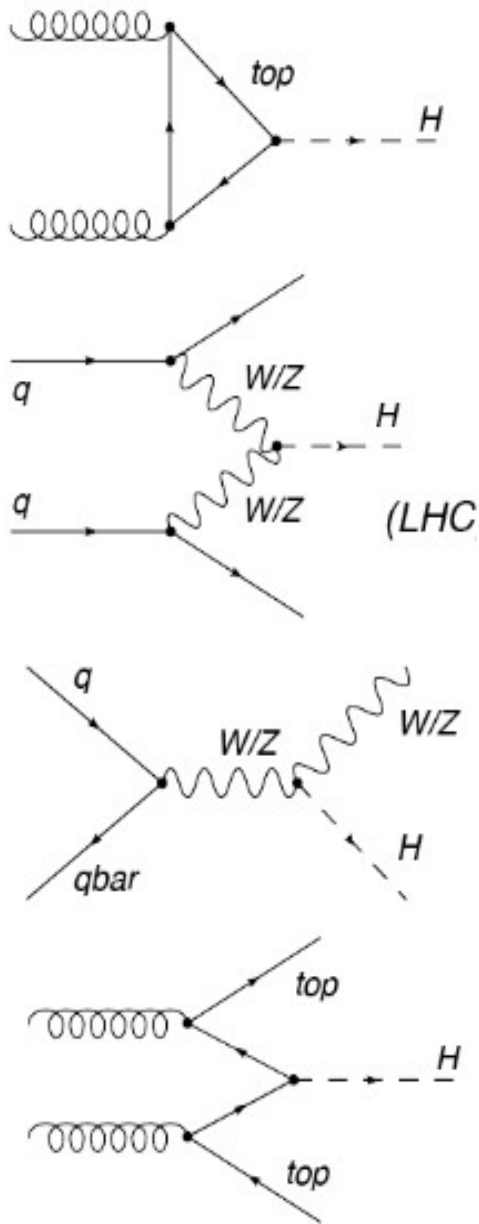
$$M_{\tilde{C}} = \begin{pmatrix} 0_{2 \times 2} & X_{2 \times 2}^T \\ X_{2 \times 2} & 0_{2 \times 2} \end{pmatrix}, \quad \text{with } X_{2 \times 2} = \begin{pmatrix} M_2 & \sqrt{2}s_\beta m_W \\ \sqrt{2}c_\beta m_W & \mu \end{pmatrix}$$

M_2 Wino
 $|\mu|$ Higgsino

$$\begin{pmatrix} \chi_1^+ \\ \chi_2^+ \end{pmatrix} = \begin{pmatrix} 1 & \mathcal{O}\left(\frac{m_Z}{M}\right) \\ \mathcal{O}\left(\frac{m_Z}{M}\right) & 1 \end{pmatrix} \begin{pmatrix} \tilde{W}^+ \\ \tilde{H}_u^+ \end{pmatrix}$$

$$\begin{pmatrix} \chi_1^- \\ \chi_2^- \end{pmatrix} = \begin{pmatrix} 1 & \mathcal{O}\left(\frac{m_Z}{M}\right) \\ \mathcal{O}\left(\frac{m_Z}{M}\right) & 1 \end{pmatrix} \begin{pmatrix} \tilde{W}^- \\ \tilde{H}_d^- \end{pmatrix}$$

Higgs production and decay at the LHC



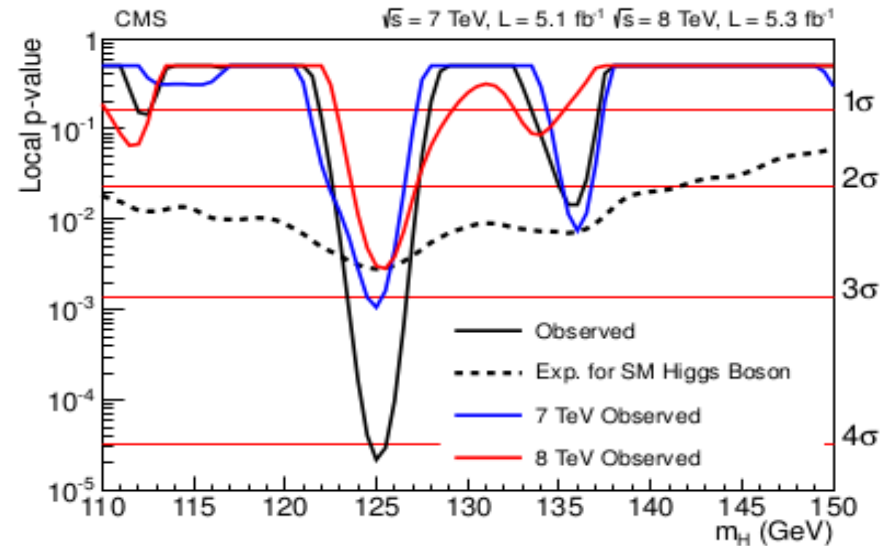
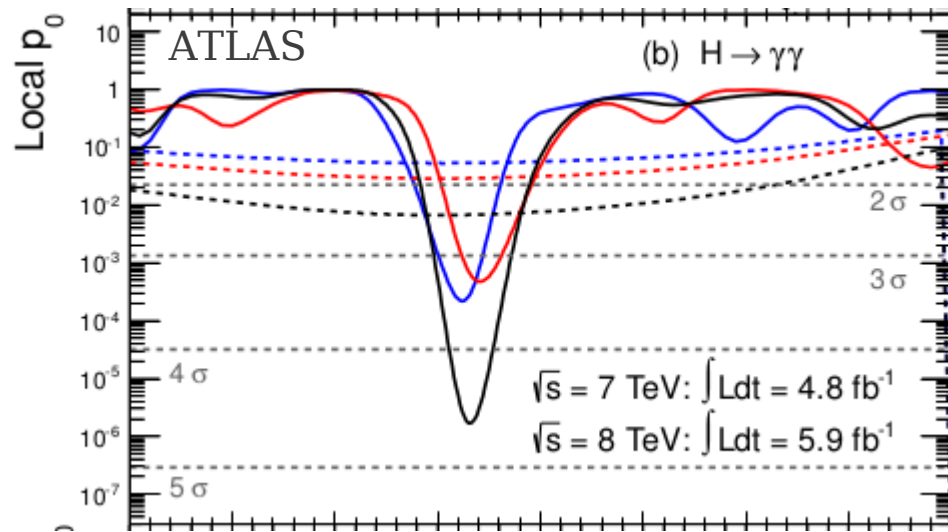
Expect:

$\sim 400 \gamma\gamma$ events

$\sim 10 ZZ$ events

Discovery of Higgs boson at the LHC

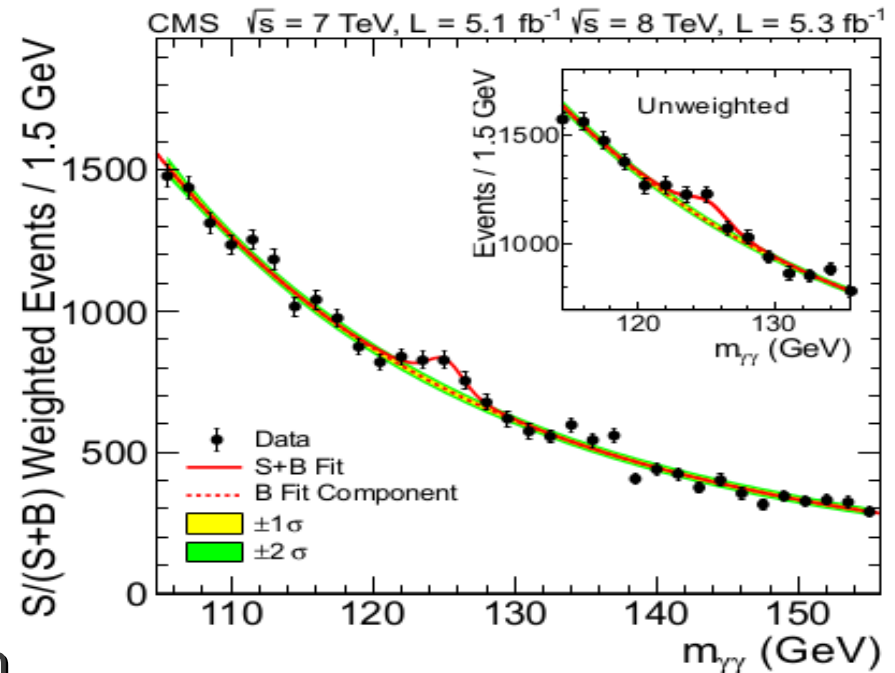
ATLAS (hep-ex: 1207.7214) and CMS (hep-ex: 1207.7235) Higgs $\rightarrow \gamma\gamma$ results



p-value: probability that background fluctuates to give an excess as large as the (average) signal size expected for a SM Higgs.

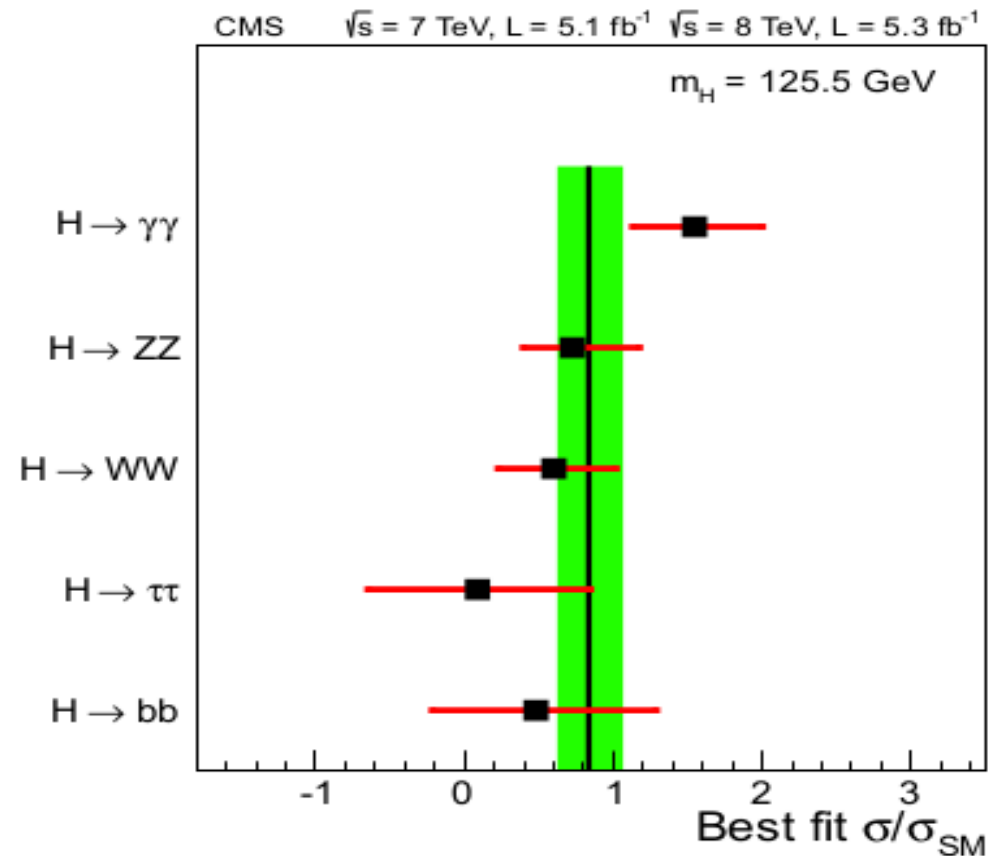
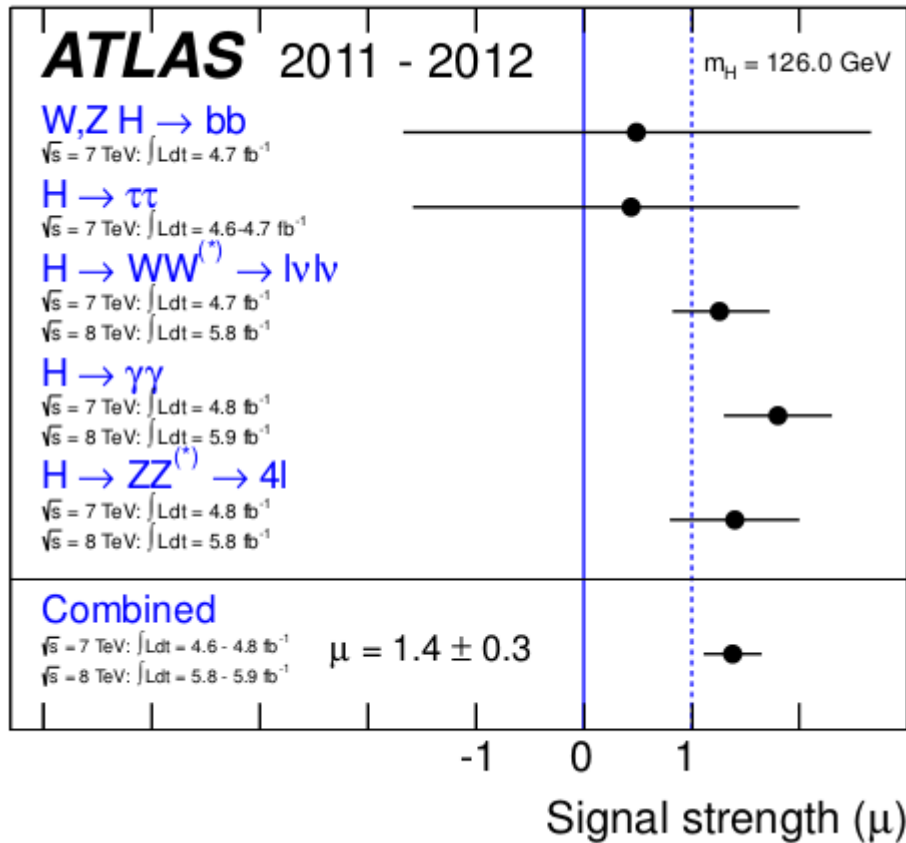
ATLAS: Observed a peak at 126.5 GeV
 Local signif.: 4.5σ (obs), 2.4σ (expected)

CMS: Observed a peak at 125.3 GeV
 Local signif.: 4.1σ (obs), 2.8σ (expected)



Must be a spin 0 or 2 boson

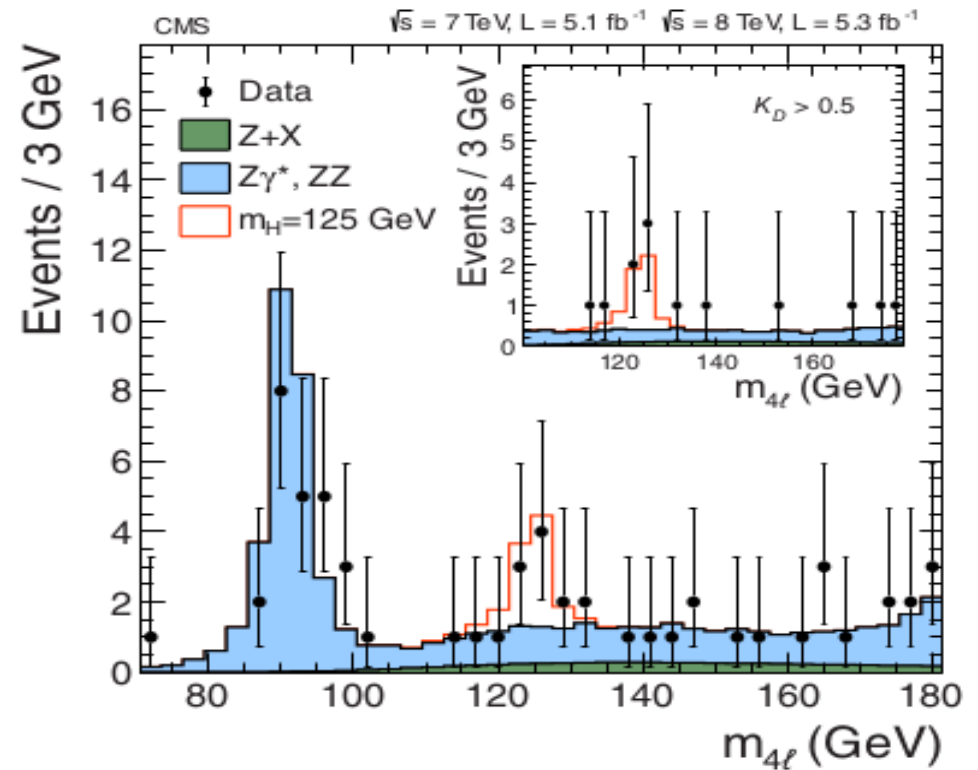
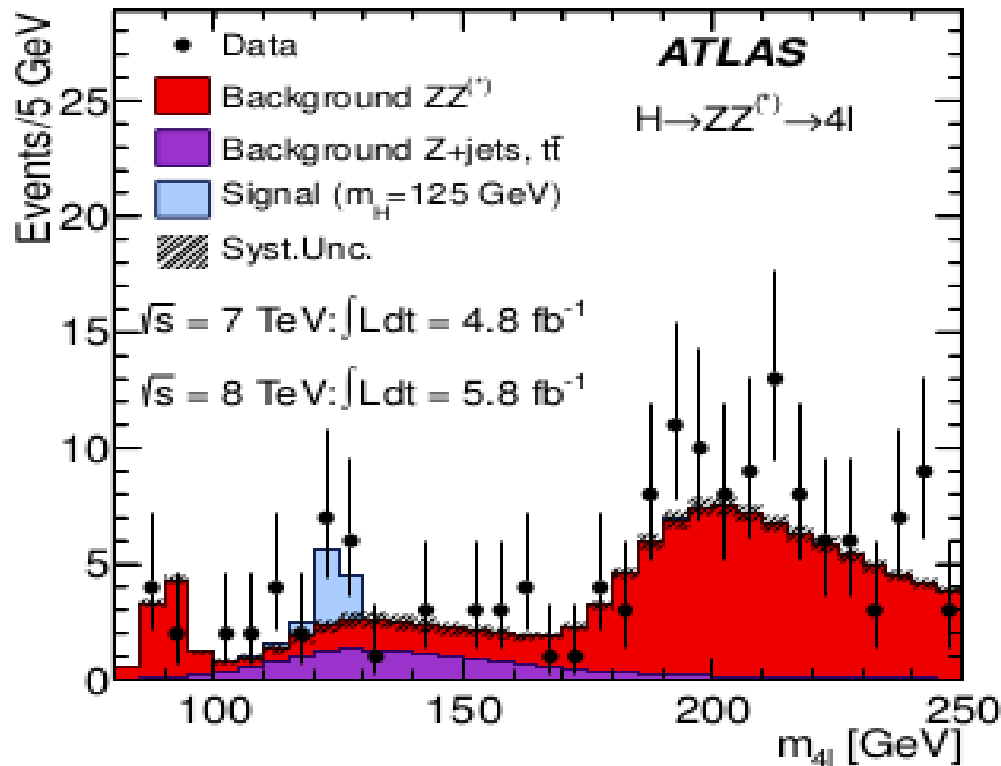
Compatibility with SM Higgs boson



Results are self consistent within errors .. so far

Discovery of Higgs boson at the LHC

ATLAS (hep-ex: 1207.7214) and CMS (hep-ex: 1207.7235) Higgs \rightarrow ZZ results

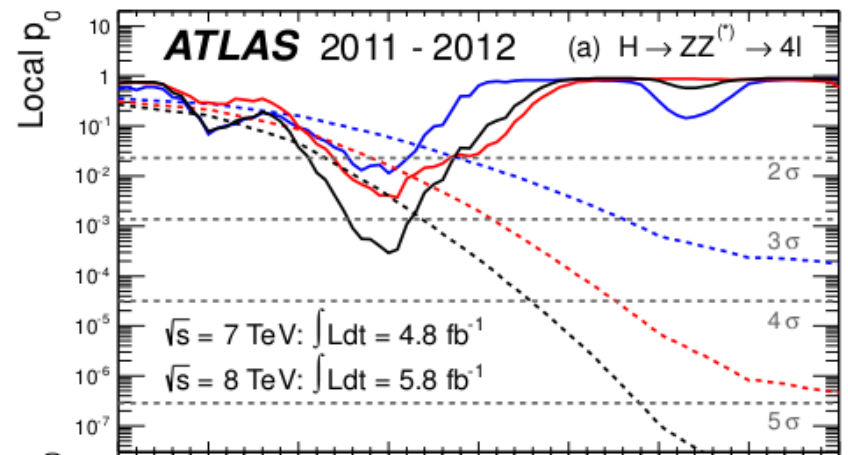


ATLAS: Observed a peak at 125 GeV

Local signif.: 3.6σ (obs), 2.7σ (expected)

CMS: Observed a peak at 125.5 GeV

Local signif.: 3.2σ (obs), 3.8σ (expected)



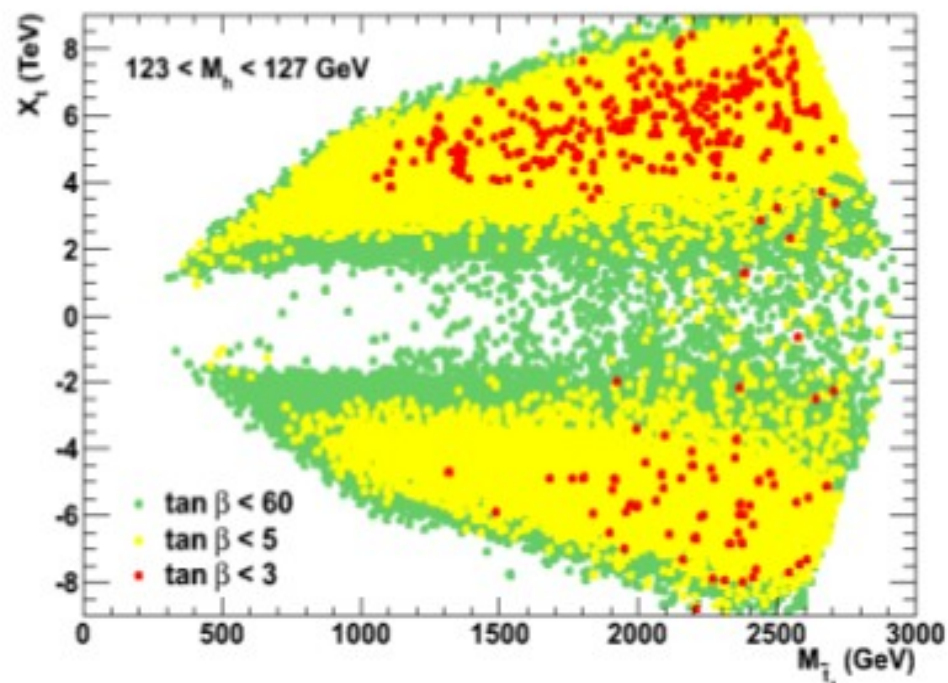
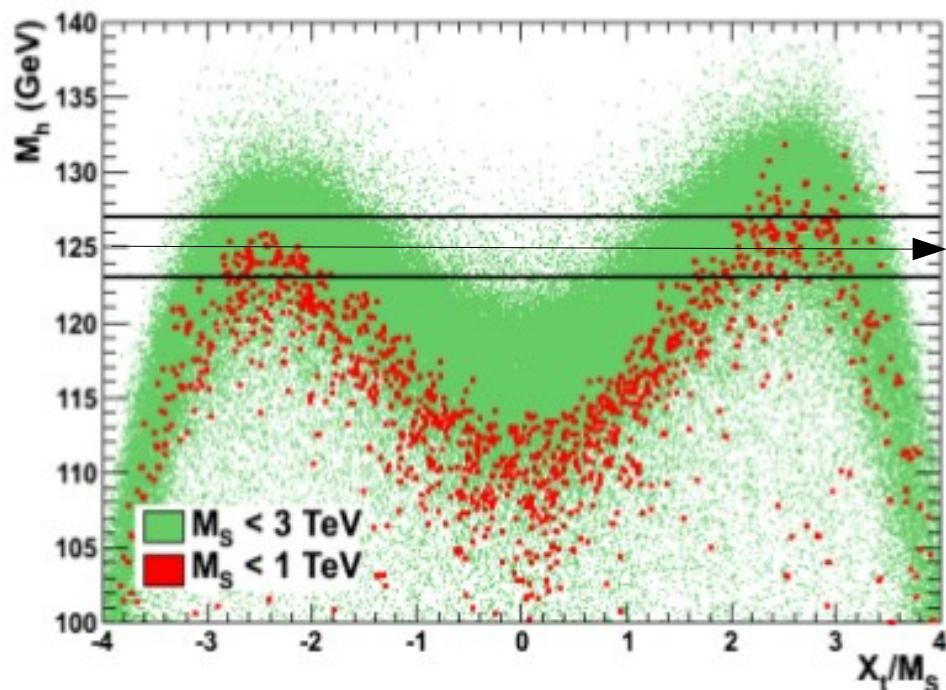
Impact of Higgs discovery on new physics

Consider Higgs to be $123 < M_h \text{ (GeV)} < 127$

$$M_h^2 \stackrel{M_A \gg M_Z}{\approx} M_Z^2 \cos^2 2\beta + \frac{3m_t^4}{2\pi^2 v^2} \left[\log \frac{M_S^2}{m_t^2} + \frac{X_t^2}{M_S^2} \left(1 - \frac{X_t^2}{12M_S^2} \right) \right]$$

The consequences of this in pMSSM (19 parameters)

A. Arbey, M. Battaglia, A. Djouadi, F.M., J. Quevillon, Phys.Lett. B708 (2012) 162



A large part of the pMSSM still survives

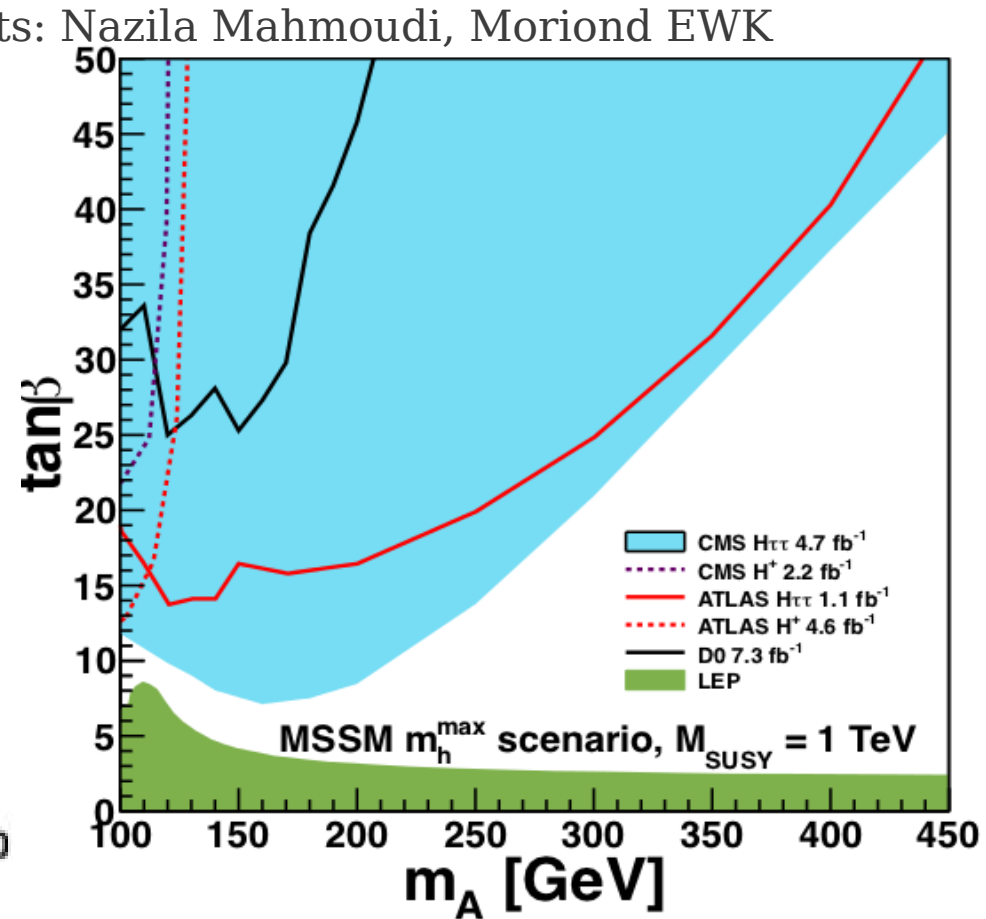
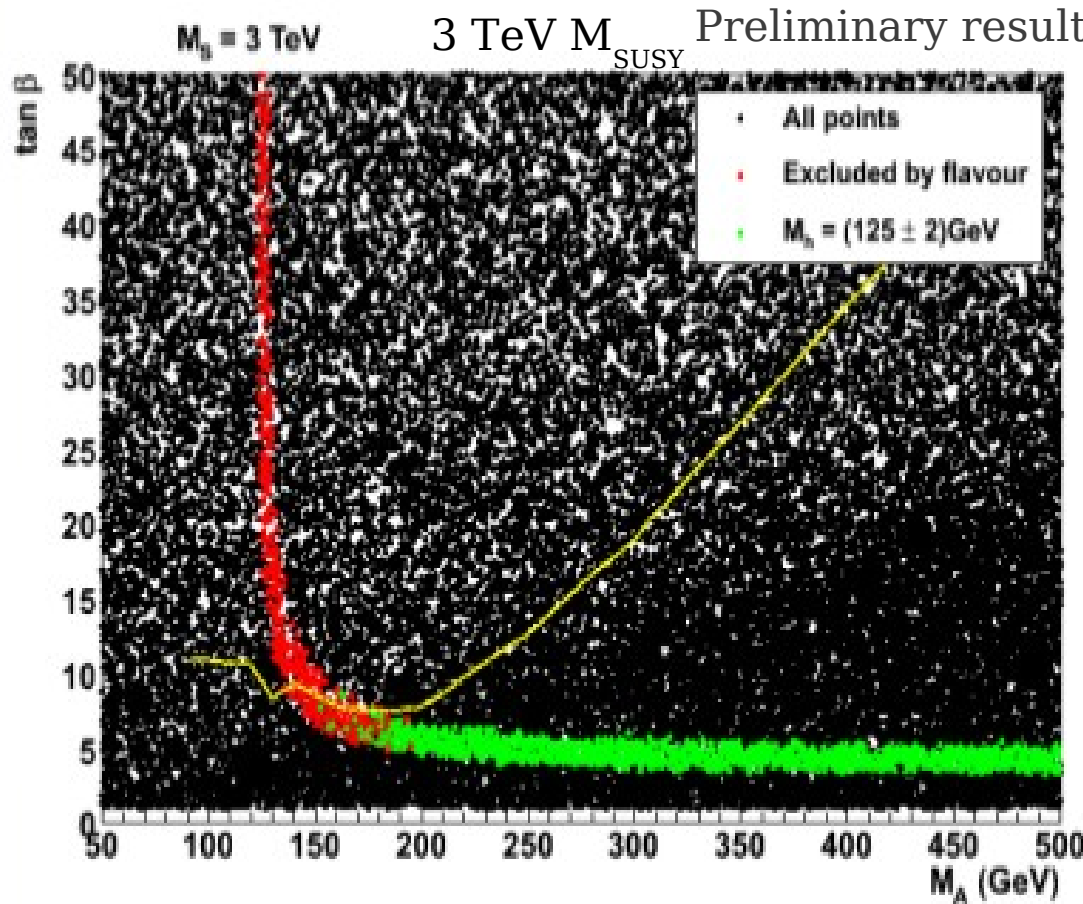
No mixing cases ($X_t \sim 0$) is excluded for $M_S < 1 \text{ TeV}$

- Even at $M_S < 3 \text{ TeV}$, chances are narrow

Small stop masses are still allowed

Impact of Higgs discovery on new physics

In the maximal mixing case ($X_t = \sqrt{6} M_S$) $123 < M_h$ (GeV) < 127



Very strong constraint on the neutral Higgs searches!

Flavour constraints: $b \rightarrow s\gamma$, $B \rightarrow \tau\nu$ and new LHCb B_s results