

Photoproduction of jets

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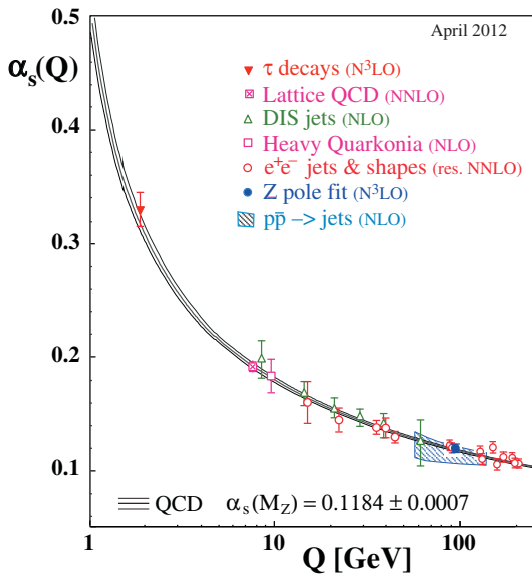


References

- MK
Theory of hard photoproduction
Rev. Mod. Phys. 74 (2002) 1221 [hep-ph/0206169]
- MK, G. Kramer, M. Michael
NNLO contributions to jet photoproduction and determination of α_s
Phys. Rev. D 89 (2014) 074032 [1310.1724]
- MK, G. Kramer
Dijet photoproduction of massless charm jets at next-to-leading order of QCD
Eur. Phys. J. C 71 (2011) 1774 [1104.0095]

Running of the strong coupling

PDG collaboration, Phys. Rev. D 86 (2012) 010001



Jet photoproduction in peripheral heavy-ion collisions

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Equivalent photon spectra

- Elastic scattering
- Inelastic scattering

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Differential $\gamma\gamma$ cross section:

$$d\sigma = \sum_{a,b} \int dx_a f_{a/\gamma}(x_a, \mu_a) \int dx_b f_{b/\gamma}(x_b, \mu_b) d\sigma_{ab}(\alpha_s, \mu, \mu_a, \mu_b)$$

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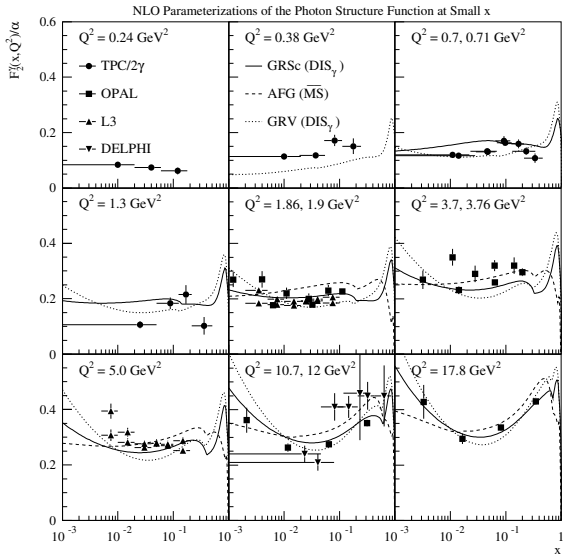
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Pointlike contribution to photon PDFs ($Q^2 \equiv \mu^2 \equiv \mu_{a,b}^2$):

$$f_{i/\gamma}^{\text{pl.}}(Q^2) = \frac{\alpha}{2\pi} \left[\frac{4\pi}{\alpha_s(Q^2)} a_i + b_i + \mathcal{O}(\alpha_s) \right]$$

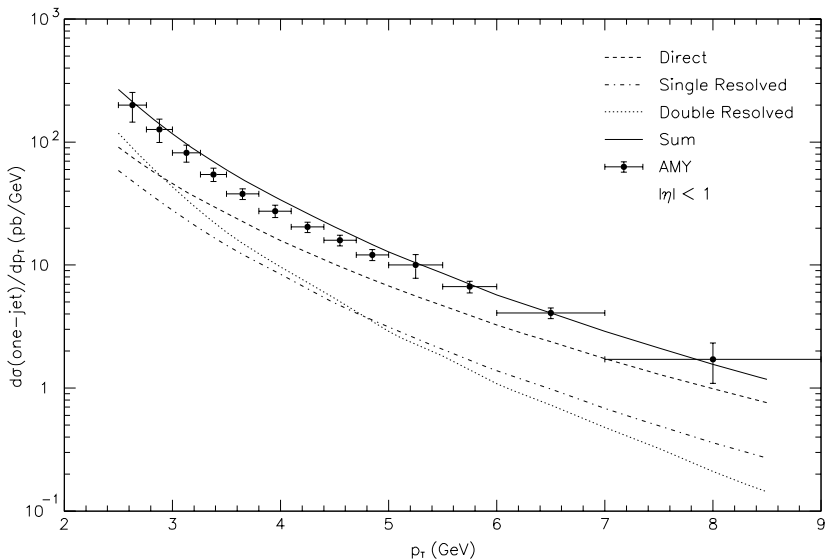
Photon structure function

MK, Rev. Mod. Phys. 74 (2002) 1221 [hep-ph/0206169]



Jets in $\gamma\gamma$ collisions at TRISTAN ($\sqrt{s} = 58$ GeV)

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Experimental conditions at TRISTAN ($\sqrt{s} = 58$ GeV)

AMY Coll., Phys. Lett. B 325 (1994) 248

Event selection:

- ≥ 4 charged tracks w/ $|p| > 0.75$ GeV, ≥ 1 w/ $p_T > 1$ GeV
- Net $p_T \leq 5$ GeV, net charge $\sum_i q_i \leq 2$
- $E_{\max}^{\text{had.}} \leq E_{\text{beam}}/4$, 4 GeV $\leq W_{\text{vis.}} \leq 20$ GeV (assigning m_π)

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- Cone algorithm with $R = 1$
- $p_T^{\text{jet}} \geq 2.5$ GeV
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LHC conditions:

- Rapidity gaps \rightarrow lower p_T threshold than for incl. production
- Jets interact also strongly \rightarrow SD contribution with $S = 0.1$
- No further suppression from SD to DD

NNLO corrections from threshold resummation (1)

N. Kidonakis, Int. J. Mod. Phys. A 19 (2004) 1793 [hep-ph/0303186]

Universal threshold logarithms spoil perturbative expansion:

$$D_I = \left[\frac{\ln^I(1-z)}{1-z} \right]_+ \quad \text{with} \quad z = \frac{M_{jj}^2}{s}$$

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Resummation to all orders in α_s :

$$\sigma^{\text{res.}} = E_a(N) E_b(N) H_{ab} \exp[G_{ab}(N)]$$

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Perturbative re-expansion (approximate NLO, NNLO, ...):

$$\sigma^{\text{res.}} = \sigma^{(0)} + \sigma^{(1)} + \sigma^{(2)} + \dots$$

NNLO corrections from threshold resummation (2)

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NLO partonic cross section:

$$\begin{aligned}\sigma^{(1)} &= \sigma^{(0)} \frac{\alpha_s}{\pi} [c_3 D_1(z) + c_2 D_0(z) + c_1 \delta(1-z)] \\ &+ \frac{\alpha_s^{d_{\alpha_s}+1}}{\pi} [A^c D_0(z) + T_1^c \delta(1-z)]\end{aligned}$$

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NNLO partonic cross section:

$$\sigma^{(2)} = \sigma^{(0)} \frac{\alpha_s^2}{\pi^2} \left[\frac{1}{2} c_3^2 D_3(z) + \left(\frac{3}{2} c_3 c_2 - \frac{\beta_0}{4} \right) D_2(z) + \dots \right]$$

NNLO corrections from threshold resummation (3)

MK, G. Kramer, M. Michael, Phys. Rev. D 89 (2014) 074032 [1310.1724]

Coefficients for $\gamma q \rightarrow qg$:

$$\begin{aligned}c_3 &= C_F - N_C, \\c_2 &= C_F \left[-\ln \left(\frac{\mu_p^2}{s} \right) - \frac{3}{4} + 2 \ln \left(\frac{-u}{s} \right) \right] + N_C \ln \left(\frac{t}{u} \right) - \frac{\beta_0}{4}, \\c_1^\mu &= -\frac{3C_F}{4} \ln \left(\frac{\mu_p^2}{s} \right) + \frac{\beta_0}{4} \ln \left(\frac{\mu^2}{s} \right).\end{aligned}$$

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Coefficients for $\gamma g \rightarrow q\bar{q}$:

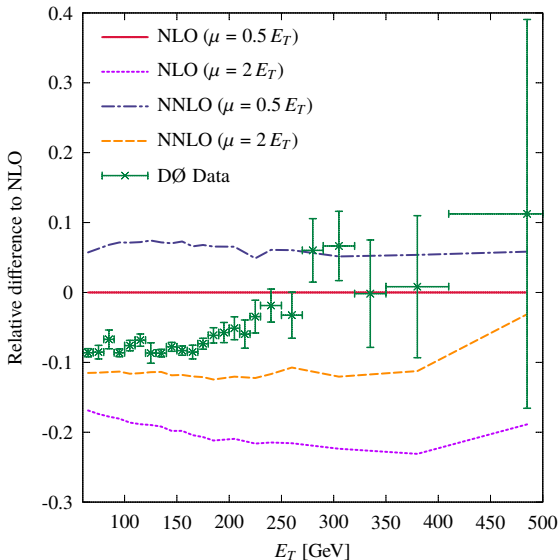
$$c_3 = 2(N_C - C_F),$$

$$c_2 = -\frac{3C_F}{2} + N_C \left[-\ln \left(\frac{\mu_p^2}{s} \right) + \ln \left(\frac{tu}{s^2} \right) \right],$$

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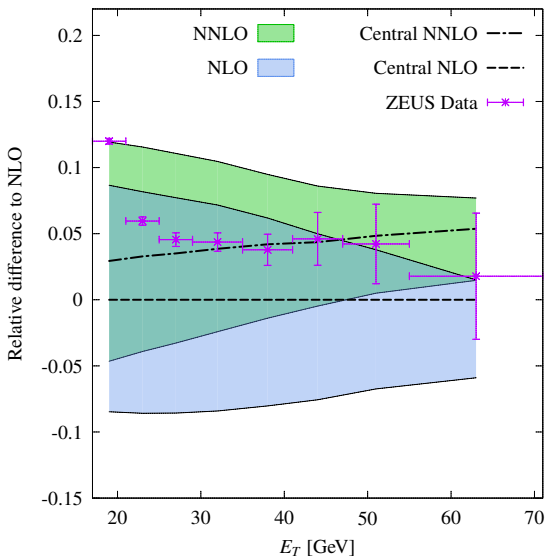
Jet hadroproduction at the Tevatron

MK, G. Kramer, M. Michael, Phys. Rev. D 89 (2014) 074032 [1310.1724]



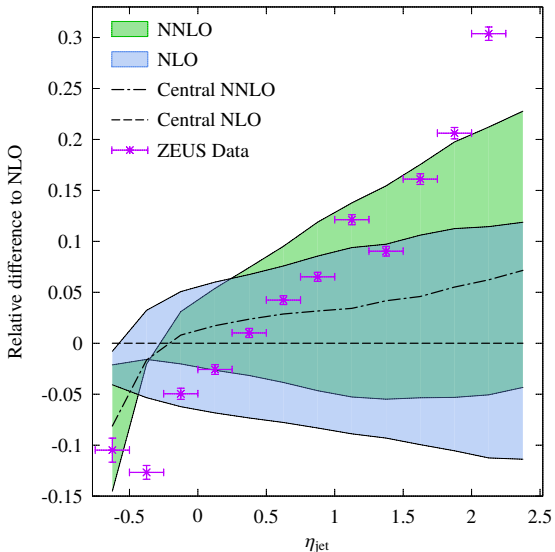
Jet photoproduction at HERA

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ZEUS determination of $\alpha_s(M_Z)$ at NLO:

$$\alpha_s(M_Z) = 0.121 \pm 0.002 \text{ (exp.)} \pm 0.004 \text{ (th.)}$$

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Our determination of $\alpha_s(M_Z)$ up to **NNLO**:

$$\alpha_s(M_Z) = 0.120 \pm 0.002 \text{ (exp.)} \pm 0.003 \text{ (th.)}$$

Charm contribution

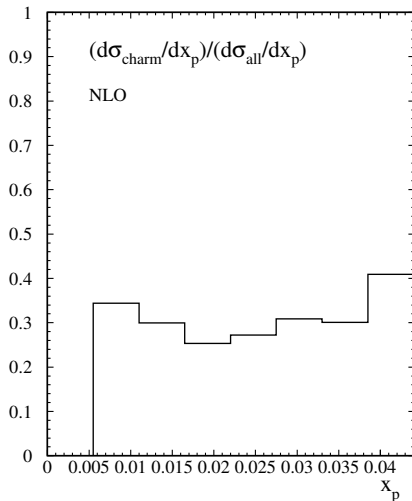
MK, G. Kramer, Eur. Phys. J. C 71 (2011) 1774 [1104.0095]

Direct photoproduction: $d\sigma_{\text{charm}}/d\sigma_{\text{all}} \sim e_c^2 / \sum_q e_q^2 = 2/5$

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Theoretical approaches to heavy-quark cross sections

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Massive calculations:

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ZEUS measurement [Phys. Lett. B 565 (2003) 87]:

- $E_T^{\text{jet}} > 5 \text{ GeV}$, $|\eta^{\text{jet}}| < 2.4$
- $p_T^{D^*} > 3 \text{ GeV}$, $|\eta^{D^*}| < 1.5$

Comparison to the ZEUS measurement

MK, G. Kramer, Eur. Phys. J. C 71 (2011) 1774 [1104.0095]

Observed momentum fractions in the photon / proton:

$$x_{\gamma}^{\text{obs.}} = \frac{\sum_{\text{jets}} E_T^{\text{jet}} e^{-\eta^{\text{jet}}}}{2yE_e} \quad x_p^{\text{obs.}} = \frac{\sum_{\text{jets}} E_T^{\text{jet}} e^{+\eta^{\text{jet}}}}{2E_p}$$

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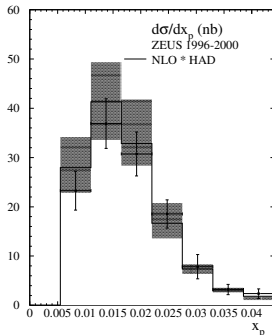
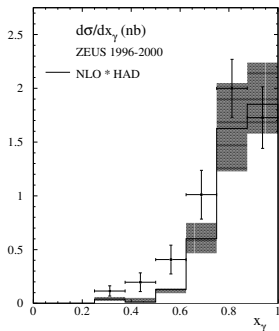
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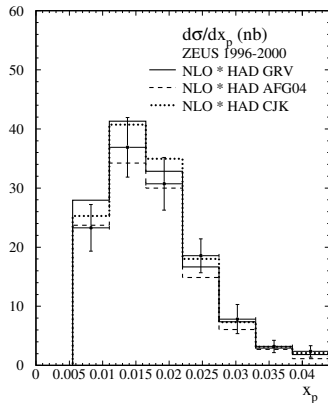
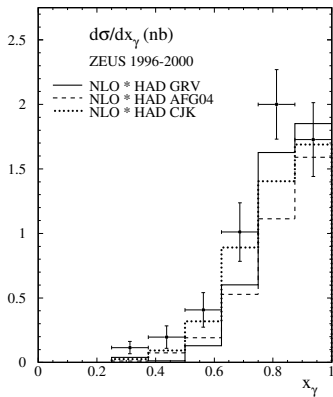
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Differential cross sections:



Sensitivity to the charm content in the photon

MK, G. Kramer, Eur. Phys. J. C 71 (2011) 1774 [1104.0095]



Conclusion

Applicability of perturbative QCD:

- Clean environment crucial $\rightarrow e^+e^-$ or ultraperipheral collisions
- pQCD works in photoproduction and jets down to 2.5 GeV

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Physical applications:

- Determination of α_s , photon and proton PDFs
- Diffractive, nuclear PDFs and form factors