A natural framework for baryonic R-parity violation in GUT

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CP³ - Origins

Particle Physics & Origin of Mass

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Baryonic R-parity violation and GUT

 To appear soon, in collaboration with Luca Di Luzio (KIT, Karlsruhe) Andrea Romanino (SISSA, Trieste)

Outline

- •SUSY with RPV
- •GUT & Baryonic RPV: the problem
- •A simple model in SO(10)
- Conclusions

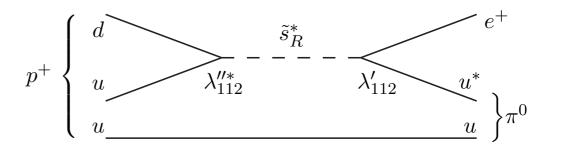
MSSM with R-parity

Most generic superpotential

$$\begin{split} W_{RPC} &= y_U^{ij} \hat{q}_i \hat{u}_j^c \hat{h}_u + y_D^{ij} \hat{h}_d \hat{q}_i \hat{d}_j^c + y_E^{ij} \hat{h}_d \hat{\ell}_i \hat{e}_j^c + \mu \, \hat{h}_u \hat{h}_d \,, \\ W_{RPV} &= \mu^i \hat{h}_u \hat{\ell}_i + \frac{1}{2} \lambda^{ijk} \hat{\ell}_i \hat{\ell}_j \hat{e}_k^c + (\lambda')^{ijk} \hat{\ell}_i \hat{q}_j \hat{d}_k^c + \frac{1}{2} (\lambda'')^{ijk} \hat{u}_i^c \hat{d}_j^c \hat{d}_k^c \,, \end{split}$$

- B and L numbers violating couplings
- Dangerous contribution to the proton decay

$$\Gamma_{p \to e^+ \pi^0} \sim m_{\text{proton}}^5 \sum_{i=2,3} |\lambda'^{11i} \lambda''^{11i}|^2 / m_{\widetilde{d}_i}^4,$$

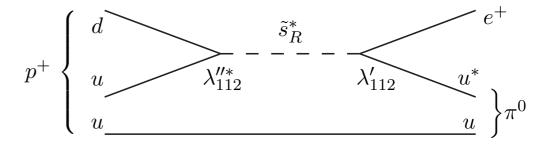


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- Possible solution: MSSM + R-parity $P_R = (-1)^{3(B-L)+2s}$ $P_M = (-1)^{3(B-L)}$
- Consequences:

(i) Lightest SUSY particle stable, possible DM candidate

(ii) SUSY particle pair produced at colliders

(iii) Each S-particle decay chain ends with an odd numbers of LSP (missing energy @ LHC)



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$$W_5 \supset \frac{\hat{q}\hat{q}\hat{q}\hat{L}}{\Lambda} + \frac{\hat{u}^c\hat{u}^c\hat{d}^c\hat{e}^c}{\Lambda}$$

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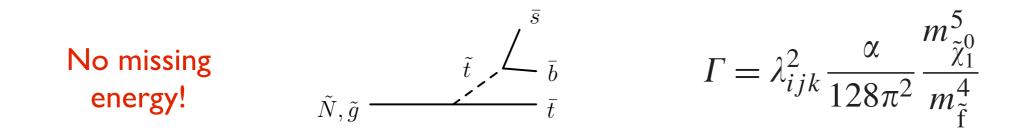
3) R-Parity violation is welcome

R-parity conservation leads to the prototypical SUSY signal: missing energy.

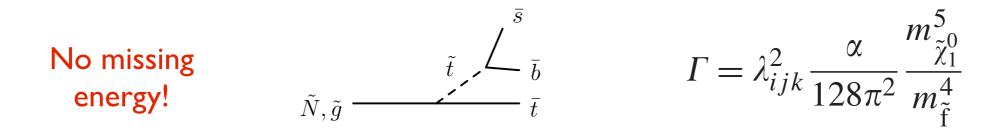
$$m_{\tilde{q}}, m_{\tilde{g}} \gtrsim 1 \,\,\mathrm{TeV}$$

Bounds can be relaxed in the presence of RPV

•RPV violating decays:



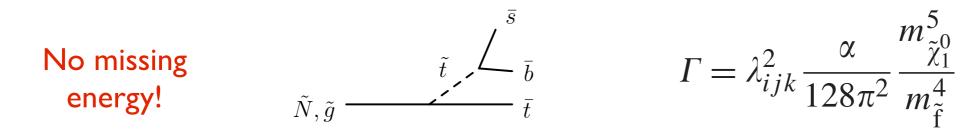
•RPV violating decays:



•Requiring a prompt decay (ex. Neutralino LSP)

$$L = 100 \,\mu\mathrm{m} \left(\beta\gamma\right) \left(\frac{m_{\tilde{f}}}{500 \,\,\mathrm{GeV}}\right)^4 \left(\frac{100 \,\,\mathrm{GeV}}{m_{\chi}}\right)^5 \left(\frac{2 \cdot 10^{-3}}{\lambda}\right)^2$$

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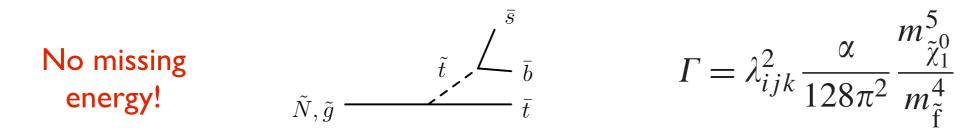


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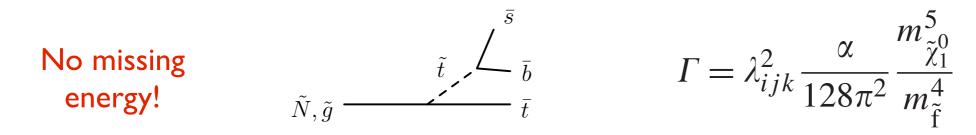
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•We improve on naturalness respect to RPC, but now, what about GUT?

•Is it possible to obtain the obtain the BNV operator without generating the LNV ones in GUT?

$$W_{\text{eff}}^{\text{ren}} = W_{\text{MSSM}} + \lambda_{ijk}^{\prime\prime} u_i^c d_j^c d_k^c,$$

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•Decomposed in terms of MSSM superfields, we obtain

$$\Lambda_{ijk} e_i^c \ell_j \ell_k + 2\Lambda_{ijk} q_i d_j^c \ell_k + \Lambda_{ijk} u_i^c d_j^c d_k^c$$

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- •The problem is due to the fact that quarks and leptons share same irreducible GUT representations
- Previous attempts:
 I. Smirnov, Vissani (1996) SU(5) + fine tuning
 2. Tamvakis(1996) SU(5) + large representations
 3. Giudice-Rattazzi (1997) Flipped SU(5)
 4. Bhattacherjee, et al. (2013) SU(5)xSU(3)

•Consider SO(10) with matter in the $16_i = q_i \oplus u_i^c \oplus d_i^c \oplus \ell_i \oplus e_i^c \oplus \nu_i^c$

•To get the BNV operator we need a trilinear term in the 16, for example

$$\frac{16_H 16_i 16_j 16_k}{\Lambda} \xrightarrow{\langle 16_H \rangle} e^c \ell \ell + q d^c \ell + u^c d^c d^c$$

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- •A possible way to project out the SU(2) doublets from the 16 is trough $\langle 45_R \rangle = V_R T_{3R} \quad (\langle 45_R \rangle 16_i)_{16} = V_R \left(-\frac{1}{2}u_i^c + \frac{1}{2}d_i^c + \frac{1}{2}e_i^c - \frac{1}{2}\nu_i^c \right)$

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•Finally we get
$$16_H (45_R 16_i)_{16} (45_R 16_j)_{16} 16_k$$

 Λ^3 $\langle 16_H \rangle \langle 45_R \rangle \quad u^c d^c d^c$

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$$\underbrace{16_H(45_R16_i)_{16}(45_R16_j)_{16}16_k}_{\Lambda^3}$$
 $\underbrace{16_H}\langle 45_R\rangle$ $u^c d^c d^c$
•This can be obtained in renormalizable theory

(16)

$W_{\Delta B=1} = \mu_{16} 16 \,\overline{16} + \mu_{10} 10 \,10 + \alpha_i 16_i 45_R \overline{16} + \beta_i 16_i 16_H 10 + \eta \,16 \,16 \,10$

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$$\begin{split} W_{\text{non-pure}}^{16} &= \mu_{16} 16 \,\overline{16} + \alpha_i 16_i 45_R \overline{16} \\ &\supset \mu_{16} \left(q_{16} \overline{q}_{\overline{16}} + l_{16} \overline{l}_{\overline{16}} \right) \\ &+ \left(\mu_{16} \, d_{16}^c + V_R \, \alpha_i d_{16_i}^c \right) \overline{d}_{\overline{16}}^c \\ &+ \left(\mu_{16} \, u_{16}^c - V_R \, \alpha_i u_{16_i}^c \right) \overline{u}_{\overline{16}}^c \\ &+ \left(\mu_{16} \, e_{16}^c + V_R \, \alpha_i e_{16_i}^c \right) \overline{e}_{\overline{16}}^c , \end{split}$$

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• $\eta 16 16 10$ generates only baryonic RPV!

Yukawa sector

•MSSM chiral superfields are in non pure embedding of $16_i \oplus 16 \oplus \overline{16} \oplus 10$

•Assuming that the Higgs up and down are in 10_v , we get at the ren. level

 $W_Y = y_{ij} 16_i 16_j 10_v + z_i 16_i 16 10_v$

•Remember that light MSSM chiral superfields are "selected" by

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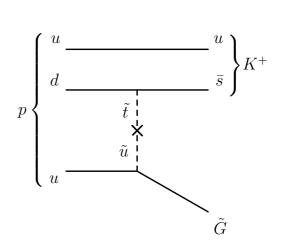
•Size of baryonic RPV couplings is related to structure of the flavor sector

$$\eta 16 \ 16 \ 10 \longrightarrow \begin{cases} 16 \to a_i u_i^c \\ 16 \to b_i d_i^c \\ 10 \to c_i d_i^c \end{cases} \longrightarrow \lambda_{ijk}^{\prime\prime} \propto a_i \left(b_j c_k - b_k c_j \right)$$

Phenomenological consequences

•Proton decay... depends on the gravitino mass

 $m_{\tilde{G}} > 1 \text{ GeV}$



,

•Delta B=2 processes

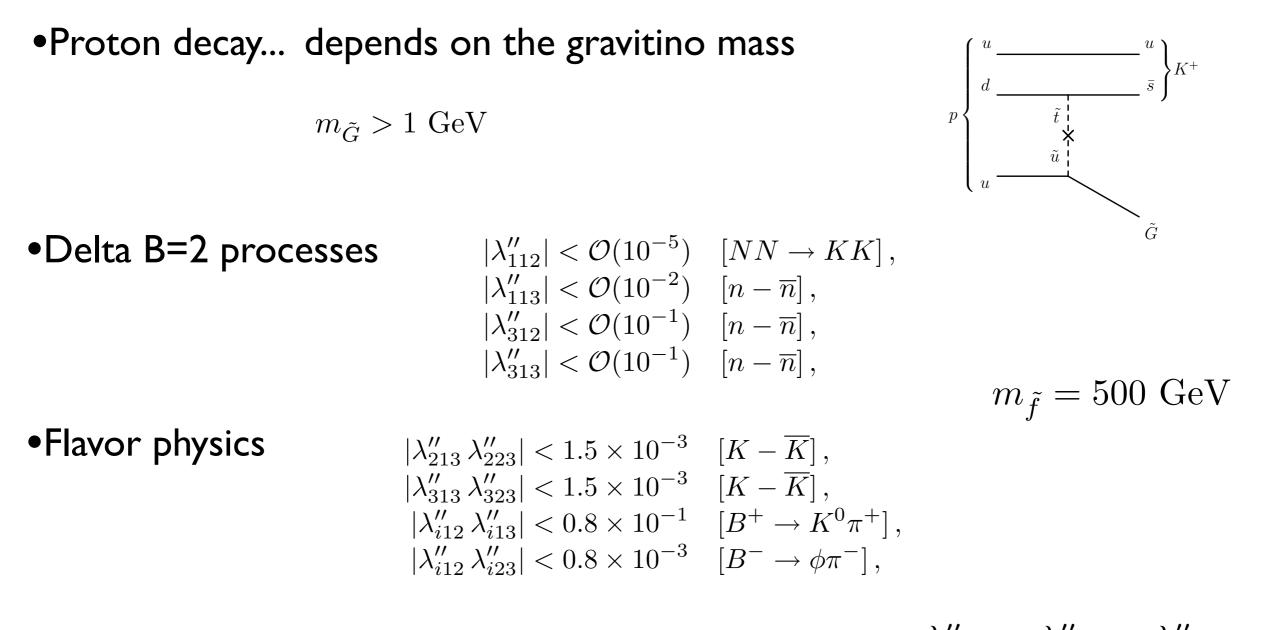
$$\begin{aligned} |\lambda_{112}''| &< \mathcal{O}(10^{-5}) \quad [NN \to KK] \\ |\lambda_{113}''| &< \mathcal{O}(10^{-2}) \quad [n - \overline{n}] , \\ |\lambda_{312}''| &< \mathcal{O}(10^{-1}) \quad [n - \overline{n}] , \\ |\lambda_{313}''| &< \mathcal{O}(10^{-1}) \quad [n - \overline{n}] , \end{aligned}$$

$$m_{\tilde{f}} = 500 \text{ GeV}$$

•Flavor physics

$$\begin{aligned} |\lambda_{213}'' \lambda_{223}''| &< 1.5 \times 10^{-3} \quad [K - \overline{K}], \\ |\lambda_{313}'' \lambda_{323}''| &< 1.5 \times 10^{-3} \quad [K - \overline{K}], \\ |\lambda_{i12}'' \lambda_{i13}''| &< 0.8 \times 10^{-1} \quad [B^+ \to K^0 \pi^+], \\ |\lambda_{i12}'' \lambda_{i23}''| &< 0.8 \times 10^{-3} \quad [B^- \to \phi \pi^-], \end{aligned}$$

Phenomenological consequences



•GUT correlation $\lambda_{ijk}^{\prime\prime} \propto a_i \left(b_j c_k - b_k c_j \right) \longrightarrow \frac{\lambda_{i12}^{\prime\prime}}{\lambda_{j12}^{\prime\prime}} = \frac{\lambda_{i13}^{\prime\prime}}{\lambda_{j13}^{\prime\prime}} = \frac{\lambda_{i23}^{\prime\prime}}{\lambda_{j23}^{\prime\prime}}$

•Natural to expect larger couplings with the heavier families like in MFV or Partial Compositeness [in progress]



- •Baryonic RPV and GUT are compatible
- •General framework: vector-like families + splitting mechanism
- •Simple models can be constructed
- •Flavor structure of the BNV couplings is non generic