

# $b \rightarrow s$ transitions and Lattice QCD

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✓ **BIG News:** November 2012, first evidence of  $B_s \rightarrow \mu^+ \mu^-$  from LHCb

➡ **PANIC:** the measured  $\text{Br}^{\text{exp}}(B_s \rightarrow \mu\mu) = (3.2 \pm 1.5)10^{-9}$  is close

to the SM,  $\text{Br}^{\text{SM}}(B_s \rightarrow \mu\mu) = (3.3 \pm 0.3)10^{-9}$

✓ **NEWS:** information from  $B \rightarrow K\mu\mu$  and  $B \rightarrow K^*\mu\mu$  available BaBar & LHCb - 2012:

➡ **NO PANIC:**  $\text{Br}(B \rightarrow K^{(*)}\mu\mu)$  and  $\text{Br}(B_s \rightarrow \mu\mu)$  sensitive to different “ $b \rightarrow s$  couplings”

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$$Br(B_d \rightarrow X_s \gamma) \propto |C_7|$$

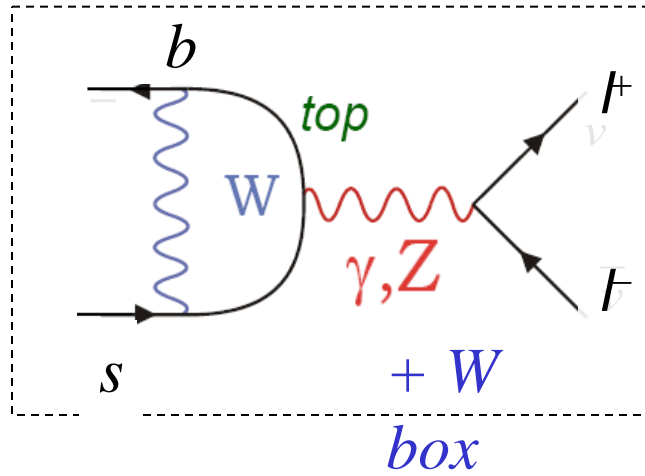
th (7%):  $(3.13 \pm 0.23) \times 10^{-4}$   
 exp (7%):  $(3.52 \pm 0.24) \times 10^{-4}$

Babar+Belle 1999-2007

$$Br(B_d \rightarrow X_s l^+ l^-) \propto C_7 C_9 + |C_9| + |C_{10}| + |C_7|$$

	exp (30%):	th (15-25%):
$[q^2 \in [0.04, 1.0] \text{ GeV}^2]$	$(0.6 \pm 0.5) \times 10^{-6}$	$(0.8 \pm 0.2) \times 10^{-6}$
$[q^2 \in [1.0, 6.0] \text{ GeV}^2]$	$(1.6 \pm 0.5) \times 10^{-6}$	$(1.6 \pm 0.1) \times 10^{-6}$
$[q^2 > 14.4 \text{ GeV}^2]$	$(4.4 \pm 1.3) \times 10^{-7}$	$(2.4 \pm 0.8) \times 10^{-7}$

Babar+Belle 2007

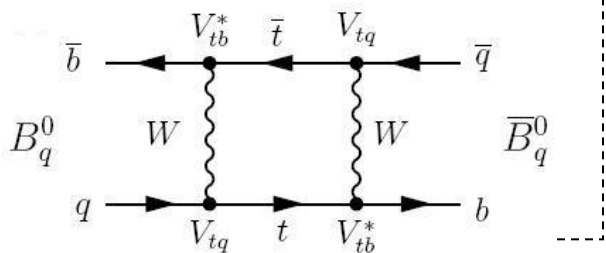


$$Heff = \left\{ \begin{array}{l} C_7 \times \bar{b}_R \sigma^{\mu\nu} s_L F_{\mu\nu} + \\ C_9 \times \bar{b}_L \gamma^\mu s_L \bar{l} \gamma^\mu l + C_{10} \times \bar{b}_L \gamma^\mu s_L \bar{l} \gamma^\mu \gamma_5 l \end{array} \right\} \text{ SM Ops.}$$

$$\left\{ \begin{array}{l} C_S \times \bar{b}_L s_R \bar{l} l + C_P \times \bar{b}_L s_R \bar{l} \gamma_5 l + \\ C_T \times \bar{b}_R \sigma^{\mu\nu} s_L \bar{l} \sigma^{\mu\nu} l + C_{T5} \times \bar{b}_R \sigma^{\mu\nu} s_L \bar{l} \sigma^{\mu\nu} \gamma_5 l \end{array} \right\} \text{ BSM Ops.}$$

+ L  $\Leftrightarrow$  R :  $(C'_9, C'_{10}, C'_S, C'_P, C'_{T5})$

$$\Delta M_S \propto C_{box} \times \bar{b}_L \gamma^\mu s_L \bar{b}_L \gamma^\mu s_L$$



Tevatron 2006

$$Br(B_s \rightarrow \mu^+ \mu^-) \propto C_{9,10,S,P}^{(b)}$$

LHCb 2012

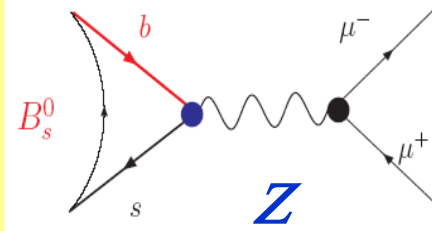
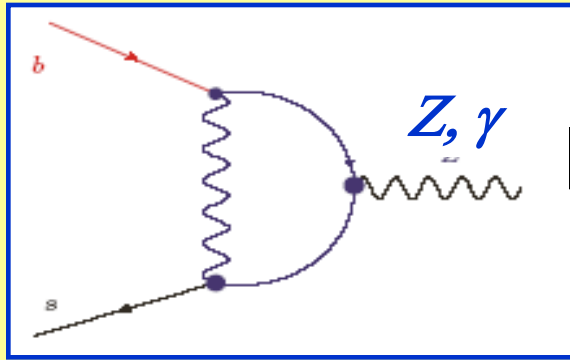
$$Br(B_d \rightarrow K^* \gamma) \propto |C_7^{(b)}|$$

$$Br(B_d \rightarrow K l^+ l^-) \propto C_{9,10,S,P,T}^{(b)}$$

$$B_d \rightarrow K^* l^+ l^- \propto C_{9,10,S,P,T}^{(b)}$$

☺ → here, we expect theory and exp. progress  
 ☺

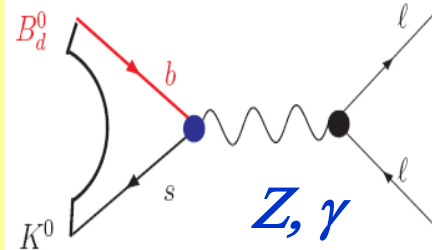
# Theory: Effective Lagrangian at $\mu \sim m_b$



$B_s \rightarrow \mu\mu$  (P=-1)

$$\langle 0 | \bar{b} \Gamma \gamma_5 s | B_s^0 \rangle \neq 0, \langle 0 | \bar{b} \Gamma s | B_s^0 \rangle = 0$$

$$\text{Br} \propto (m_l^2 (C_{10} - C'_{10}), C_{P,S} - C'_{P,S})$$



$B \rightarrow K ll$  (P=1)

$$\langle K | \bar{b} \gamma^\mu \gamma_5 s | B_s^0 \rangle = \langle K | \bar{b} \gamma_5 s | B_s^0 \rangle = 0$$

$$\text{Br} \propto (C_7 + C'_7, C_9 + C'_9, C_{10} + C'_{10}, C_{P,S} + C'_{P,S}, C_T, C_{T5})$$

+  $B \rightarrow X_s \gamma$  strongly constrains  $C_{7(8)}$

$$\begin{aligned} L_{\text{eff}} = & C_7 \bar{b} \sigma_L^{\mu\nu} s F_{\mu\nu} + C'_7 \bar{b} \sigma_R^{\mu\nu} s F_{\mu\nu} + C_9 (\bar{b} \gamma_L^\mu s) \bar{l} \gamma^\mu l + C'_9 (\bar{b} \gamma_R^\mu s) \bar{l} \gamma^\mu l \\ & + C_{10} (\bar{b} \gamma_L^\mu s) \bar{l} \gamma^\mu \gamma_5 l + C'_{10} (\bar{b} \gamma_R^\mu s) \bar{l} \gamma^\mu \gamma_5 l + C_S (\bar{b} L s) \bar{l} l + C'_S (\bar{b} R s) \bar{l} l \\ & + C_P (\bar{b} L s) \bar{l} \gamma_5 l + C'_P (\bar{b} R s) \bar{l} \gamma_5 l + C_T (\bar{b} \sigma_L^{\mu\nu} s) \bar{l} \sigma^{\mu\nu} l + C'_T (\bar{b} \sigma_R^{\mu\nu} s) \bar{l} \sigma^{\mu\nu} l \end{aligned}$$

# Theory: Hadronic Uncertainties

$$Br(B_s \rightarrow \mu^+ \mu^-) \propto C_{9,10,S,P}$$

$$Br(B_d \rightarrow K^{(*)} \ell \ell) \propto C_{9,10,S,P,T}$$

GOAL: calculate Matrix elements of 2-quark operators between hadrons (decay constants & Form factors)

*SM operators*

$$O_7 = \bar{b}_R \sigma^{\mu\nu} s_L F_{\mu\nu}$$

$$O_9 = (\bar{b} \gamma_L^\mu s) \bar{\ell} \gamma^\mu \ell$$

$$O_{10} = (\bar{b} \gamma_L^\mu s) \bar{\ell} \gamma^\mu \gamma_5 \ell$$

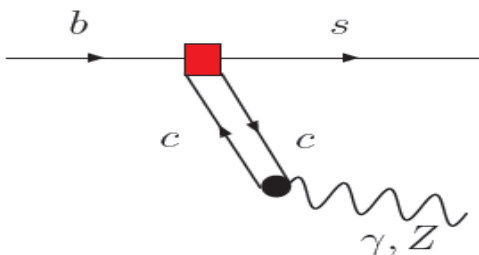
$$O_2 = (\bar{b} \gamma_L^\mu c) (\bar{c} \gamma_L^\mu s)$$

*BSM operators*

$$O_{S(P)} = (\bar{b}_R s_L) \bar{\ell} \ell_{S(P)}, O_T = (\bar{b}_R \sigma^{\mu\nu} s_L) \bar{\ell} \sigma^{\mu\nu} \ell$$

+L ↔ R

*Charm Loops*



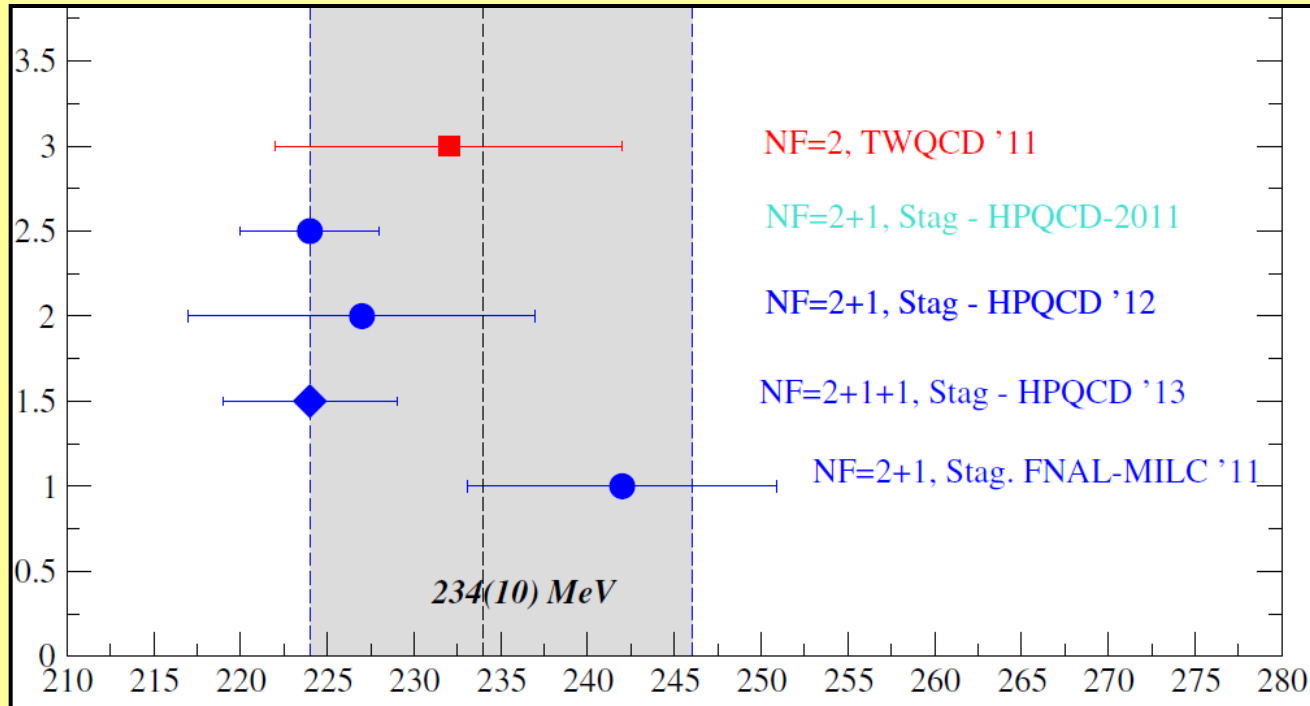
Under control (to some extent)  
at low and large  $q^2$ , out of resonance region

Khodjamirian's talk

# Theory: Hadronic Uncertainties

$B_s \rightarrow \mu\mu$

$f_{B_s}$



$B_s \rightarrow \mu\mu$

Only one hadronic parameter:  $f_{B_s}$

$$\langle 0 | \bar{b} \gamma^\mu \gamma_5 s | B_s^0 \rangle = i p^\mu f_{B_s}$$

$$f_{B_s} = (234 \pm 10) \text{ MeV}$$

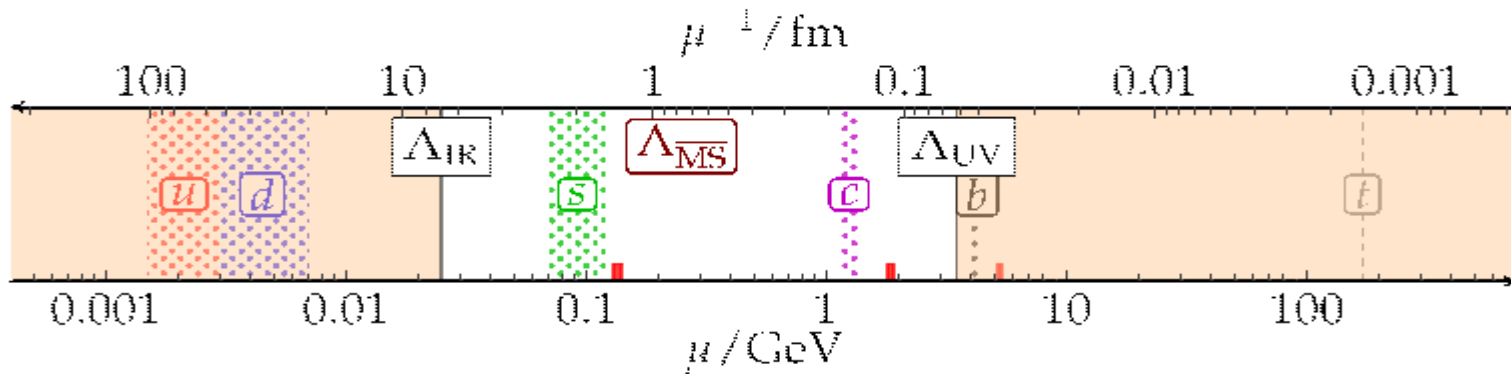
$$\langle 0 | \bar{b} \gamma_5 s | B_s^0 \rangle = -i f_{B_s} M_{B_s}^2 / m_b$$

4% hadronic uncertainty

Lattice: ETMC, MILC, HPQCD

$$\text{Br}(B_s \rightarrow \mu\mu)^{\text{SM}} = (3.3 \pm 0.3) \times 10^{-9} \text{ (6.5\%)}$$

# Challenge of B-physics: the multi scale-problem of QCD



hierarchy of disparate physical scales to be covered:

$$\Lambda_{\text{IR}} = L^{-1} \ll m_{\pi}, \dots, m_D, m_B \ll a^{-1} = \Lambda_{\text{UV}}$$

$$\left\{ O(e^{-Lm_{\pi}}) \Rightarrow L \gtrsim \frac{4}{m_{\pi}} \sim 6 \text{ fm} \right\} \rightsquigarrow L/a \gtrsim 120 \rightsquigarrow \left\{ am_D \lesssim \frac{1}{2} \Rightarrow a \approx 0.05 \text{ fm} \right\}$$

Currently  $a^{-1} < 4 \text{ GeV}$ , **b** quarks cannot be directly simulate at their physical mass due to large discretization errors ( $a m_b \ll 1$ )

❑ *effective theories: like NRQCD action*

❑ *simulate heavy quark in the charm region and extrapolate to the B + HQET.*

## ☹ Comments:

### - Discretized NRQCD action

#### ➤ Quite Sophisticated procedure!

⇒ larger set of  $1/(am_Q)$  corrections on the lattice w.r.t the continuum

➤  $O[\alpha_s^n/(am_Q)]$  divergences to be subtracted to get the continuum limit

➤ On the other hand, large experience from MILC/FNAL/HPQCD

➤ ☺ Successful strategy for  $f_B$  when comparing with unquenched results from strategy

# Theory: Hadronic Uncertainties

$$B \rightarrow K^* \gamma, B \rightarrow K^* l \bar{l}$$

$$\langle V(p', \varepsilon) | \bar{q} \hat{\gamma}^\mu b | B(p) \rangle = \frac{2iV(q^2)}{m_B + m_V} \varepsilon^{\mu\nu\rho\sigma} \varepsilon_\nu^* p'_\rho p_\sigma$$

$$\begin{aligned} \langle V(p', \varepsilon) | \bar{q} \hat{\gamma}^\mu \hat{\gamma}^5 b | B(p) \rangle &= 2m_V A_0(q^2) \frac{\varepsilon^* \cdot q}{q^2} q^\mu \\ &+ (m_B + m_V) A_1(q^2) \left( \varepsilon^{*\mu} - \frac{\varepsilon^* \cdot q}{q^2} q^\mu \right) \\ &- A_2(q^2) \frac{\varepsilon^* \cdot q}{m_B + m_V} \left( (p + p')^\mu - \frac{m_B^2 - m_V^2}{q^2} q^\mu \right) \end{aligned}$$

$$q^\nu \langle V(p', \varepsilon) | \bar{q} \hat{\sigma}_{\mu\nu} b | B(p) \rangle = 2T_1(q^2) \varepsilon_{\mu\rho\tau\sigma} \varepsilon^{*\rho} p^\tau p'^\sigma \longrightarrow \text{Br}(B \rightarrow K^* \gamma) \text{ one ff. at } q^2=0$$

$$\begin{aligned} q^\nu \langle V(p', \varepsilon) | \bar{q} \hat{\sigma}_{\mu\nu} \hat{\gamma}^5 b | B(p) \rangle &= iT_2(q^2) [\varepsilon_\mu^* (m_B^2 - m_V^2) - (\varepsilon^* \cdot q)(p + p')_\mu] \\ &+ iT_3(q^2) (\varepsilon^* \cdot q) \left[ q_\mu - \frac{q^2}{m_B^2 - m_V^2} (p + p')_\mu \right] \end{aligned}$$

$\text{Br}(B \rightarrow K^* l \bar{l})$ : 7 form factors in QCD

! LATTICE QCD: only to compute the full ff basis at large  $q^2$ :

☺ no  $O(\Lambda/m_b)$  uncertainty from Isgur-Wise relation at LO!



# Theory: Hadronic Uncertainties

## $B \rightarrow K \ell \ell$

### $B \rightarrow K \ell \ell$

Dominant uncertainties come from the form 3 factors:  $f_+(q^2)$ ,  $f_0(q^2)$ ,  $f_T(q^2)$

$$\langle B(p) | \bar{b} \gamma^\mu s | K(k) \rangle = (p^\mu + k^\mu - \frac{m_B^2 - m_K^2}{q^2} q^\mu) f_+(q^2) + \frac{m_B^2 - m_K^2}{q^2} q^\mu f_0(q^2)$$

$$\langle B(p) | \bar{b} \sigma^{\mu\nu} s | K(k) \rangle = \frac{if_T}{m_B + m_K} [(p^\mu + k^\mu) q^\nu - (p^\nu + k^\nu) q^\mu]$$

$$\diamond C_{9,10}^{(\prime)} \rightarrow f_+(q^2), f_0(q^2), C_{S,P}^{(\prime)} \rightarrow f_0/m_b \quad C_7^{(\prime)} \rightarrow f_T,$$

❖ Wide range of  $q^2 = [0, (m_B - m_K)^2]$  -> **Opportunities for different nonperturbative techniques: Lattice QCD-QCSR** – relative th. error 30% -> large room for improvement

*! LATTICE QCD: only to compute the full ff basis at large  $q^2$ :  
☺ no  $O(\Lambda/m_b)$  uncertainty from Isgur-Wise relation at LO!*

# Studies of form-factor calculations on the Lattice:

$B \rightarrow K l l$

$N_F=0$ : Quenched lattice QCD: **relativistic fermions**

❖ D. Becirevic, N. Kosnik, F. M., E. Schneider, 2012

$N_F=2+1$  staggered fermions: NRQCD

❖ FNAL/MILC, 2012      ❖ HPQCD, 2012  
 ❖ Cambridge (*prelims*), 2012

$f_+(q^2), f_-(q^2)$   
 $f_T(q^2)$

$B \rightarrow K^* l l$

$N_F=0$ : Quenched lattice QCD: **relativistic fermions**

❖ D. Becirevic, V. Lubicz & F. M. 2007

$N_F=2+1$  staggered fermions: **NRQCD**

❖ Cambridge (*prelims*), 2012

$T_{12}(q^2)$

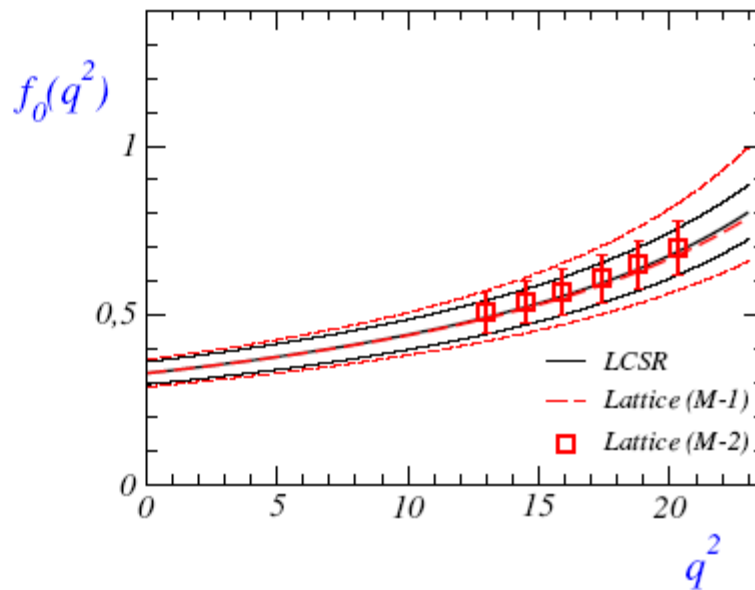
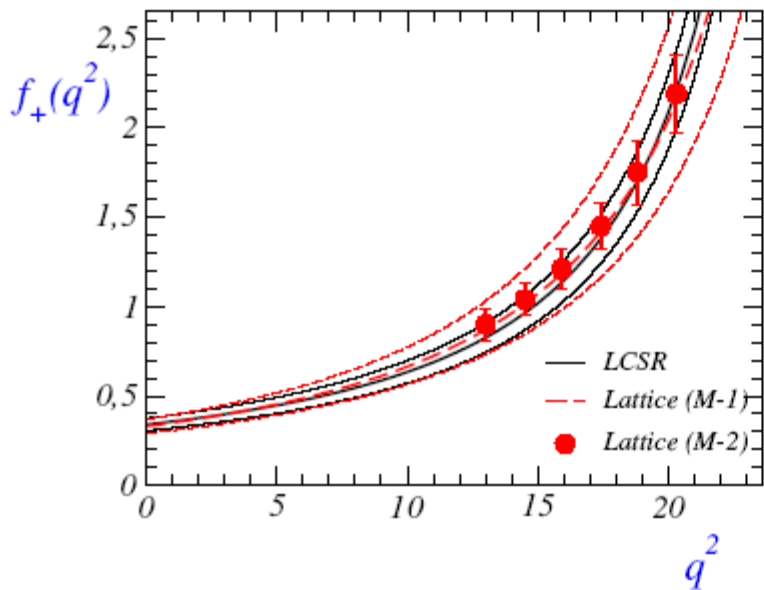
$T_{12}(q^2)$   
 $V(q^2), A_{012}(q^2)$

☹ *very preliminary unquenched activities:*

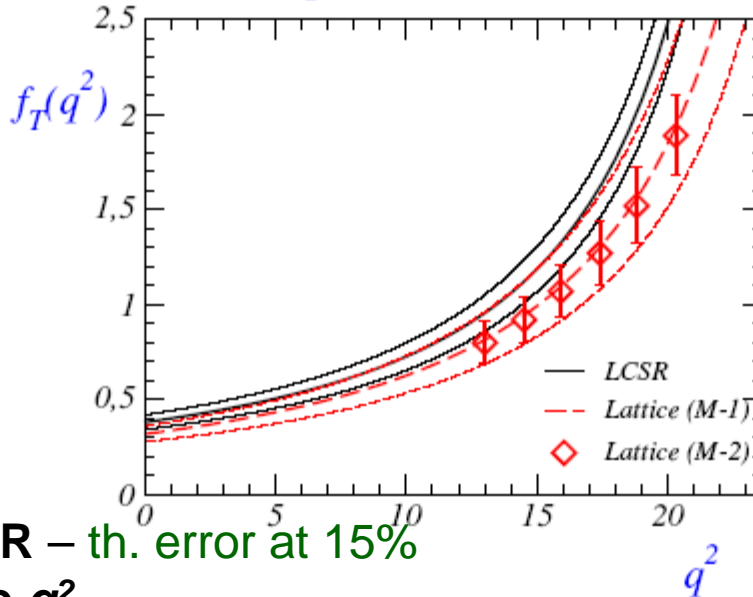
*overall agreement between Quenched and LCSR*

☹  $q^2$  dependence: *further complication with respect to  $f_B$  or  $B_B$*

# $B \rightarrow K_{ll}$ - STRATEGY 1: QCD + extrapolating from charm region:



**Lattice QCD Quenched:**  
 $f_+, f_0, f_T \rightarrow$  D. Becirevic et al. 2012



**Light cone QCD  
 sum rules [Ball'05,  
 Khodjamirian'07,  
 '10]**

- 1) Lattice QCD and LCSR – th. error at 15%
- 2) Lattice points at large  $q^2$
- 3) Agreement with LQSR

$B \rightarrow K \ell \ell$

Dominant uncertainties come from the 3 form factors:  $f_+(q^2)$ ,  $f_0(q^2)$ ,  $f_T(q^2)$

$$\langle K | \bar{b} \gamma^\mu \gamma_5 s | B \rangle \Leftrightarrow f_{+,0}(q^2) \quad \langle K | \bar{b} \sigma^{\mu\nu} s | B \rangle \Leftrightarrow f_T(q^2)$$

$$\diamond C_{9,10}^{(\prime)} \rightarrow f_+(q^2), f_0(q^2), C_{S,P}^{(\prime)} \rightarrow f_0/m_b \quad C_7^{(\prime)} \rightarrow f_T,$$

$\diamond$  Wide range of  $q^2 = [0, (m_B - m_K)^2]$  -> Opportunities for different nonperturbative techniques: Lattice QCD and LCSR – **relative error 30%**

$$\text{Br}(B \rightarrow K \ell^+ \ell^-)_{\text{SM}} = \begin{cases} (7.5 \pm 1.4) \times 10^{-7} & \text{LQCD,} \\ (6.8 \pm 1.6) \times 10^{-7} & \text{LCSR.} \end{cases}, \quad \text{Our average} \quad \text{Br}(B \rightarrow K \ell^+ \ell^-)_{\text{SM}} = (7.0 \pm 1.8) \times 10^{-7}$$

**still th. error large 30%**

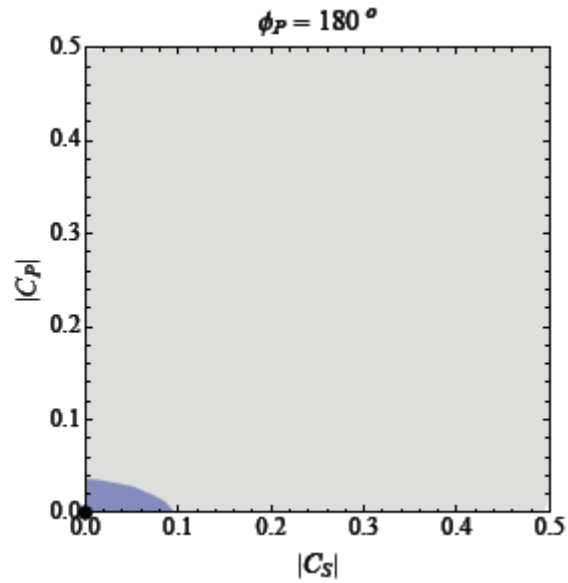
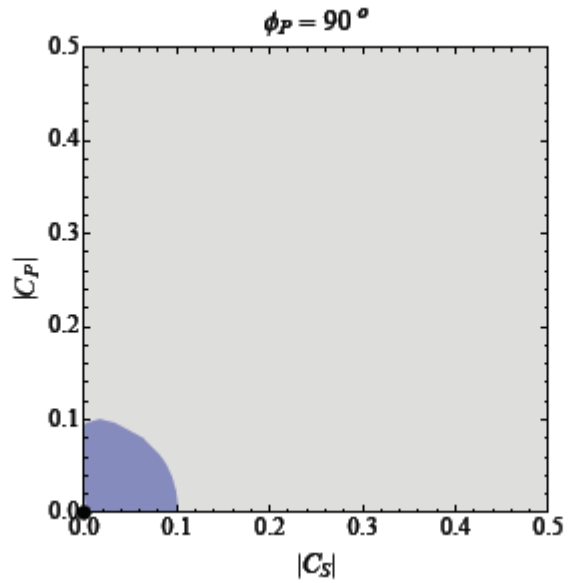
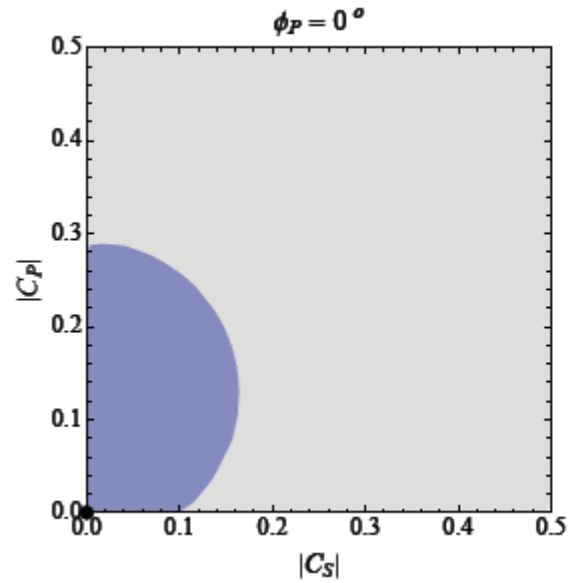
**BaBar'12**

$$\text{Br}(B \rightarrow K \ell \ell) = (4.7 \pm 0.6) \times 10^{-7}$$

**LHCb'12**

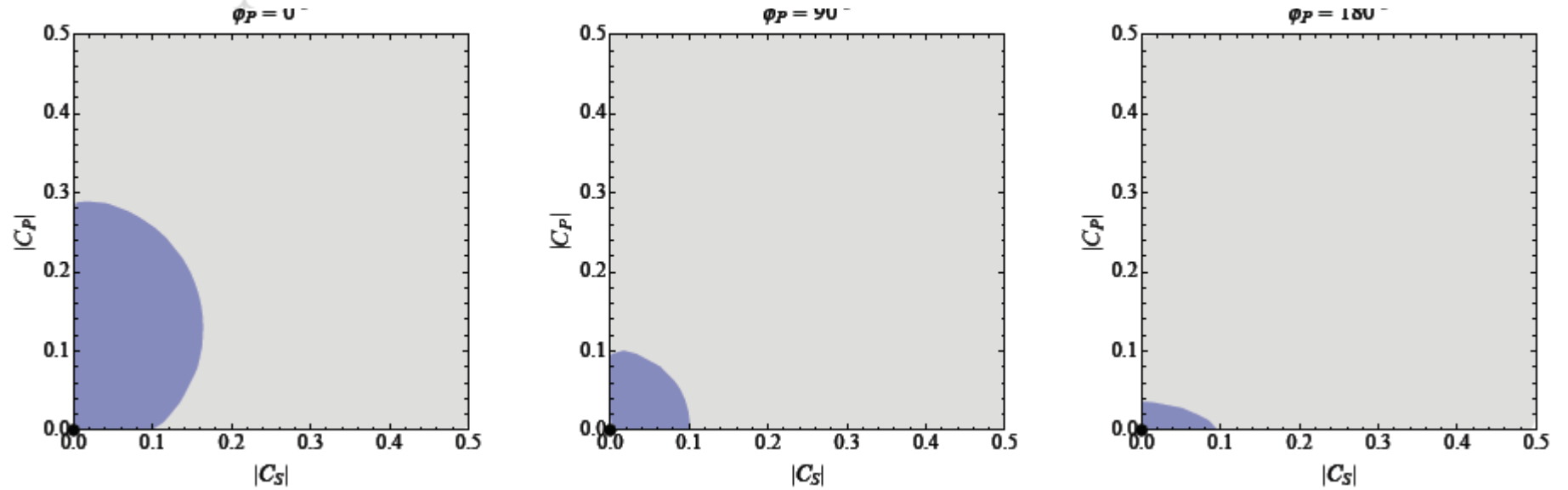
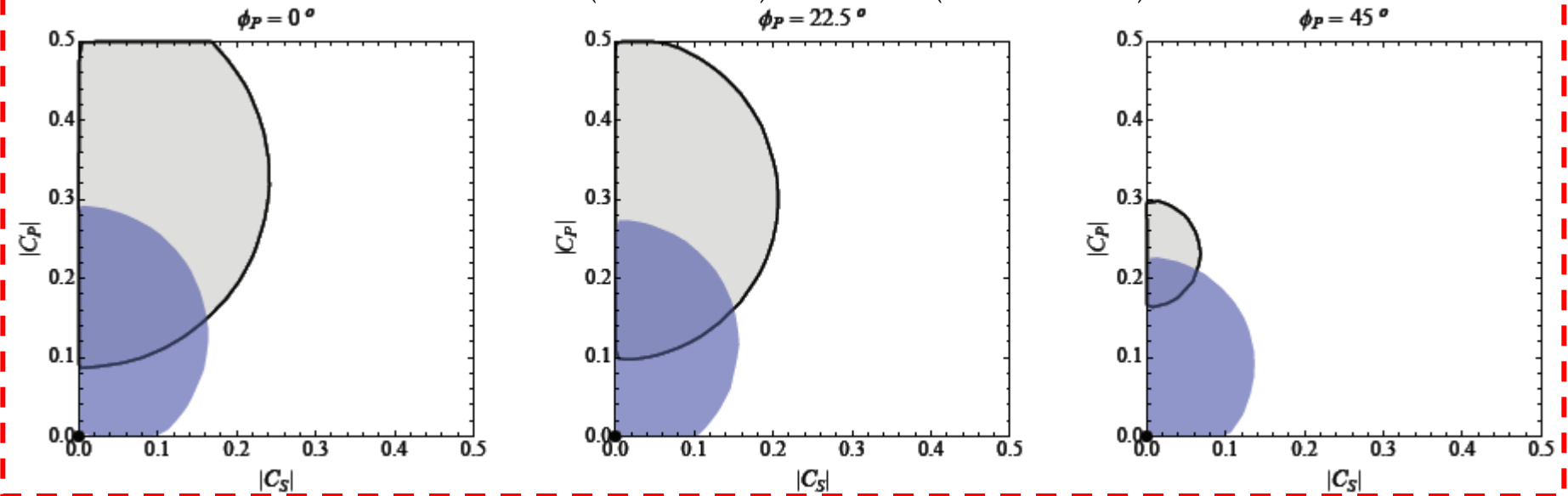
$$\text{Br}(B^+ \rightarrow K^+ \mu \mu) = (3.1 \pm 0.7) \times 10^{-7}$$

New Physics: **SM** +  $C_S (\bar{b}(1-\gamma_5)s) \bar{\ell}\gamma_5\ell$  +  $C_P (\bar{b}\gamma_5(1-\gamma_5)s) \bar{\ell}\gamma_5\ell$



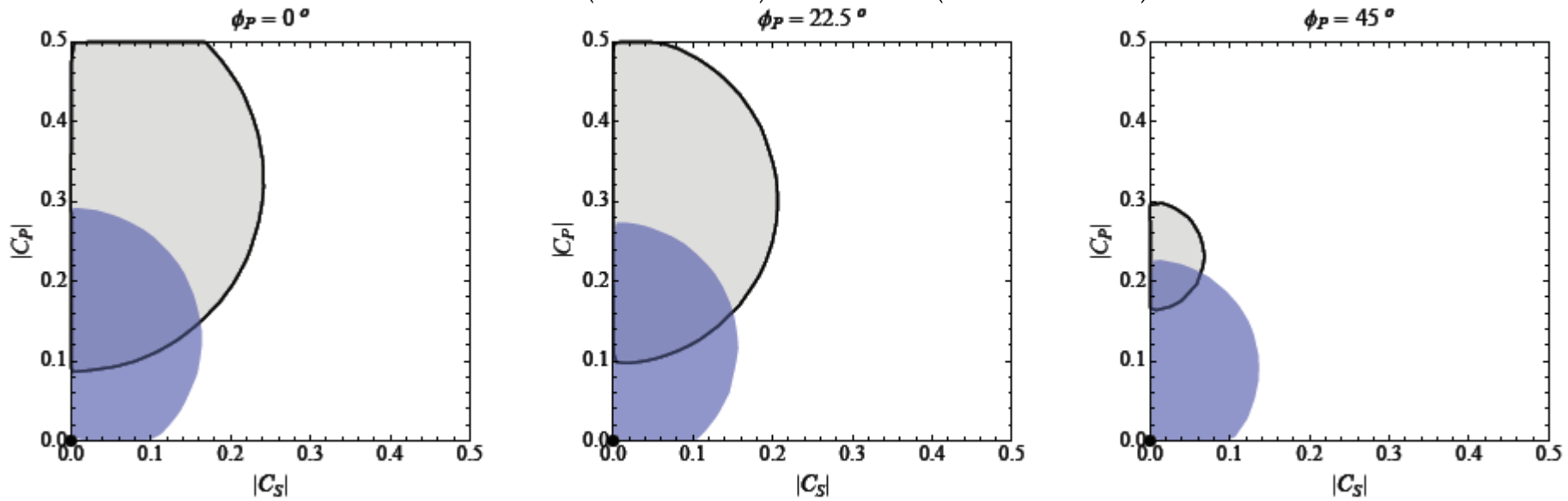
Lowering th. error on  $B \rightarrow Kll$ : 20% smaller than now

**New Physics: SM** +  $C_S (\bar{b}(1-\gamma_5)s) \bar{l}\gamma_5 l$  +  $C_P (\bar{b}\gamma_5(1-\gamma_5)s) \bar{l}\gamma_5 l$



Lowering th. error on  $B \rightarrow Kll$ : 20% smaller than now

**New Physics: SM** +  $C_S (\bar{b}(1-\gamma_5)s) \bar{\ell}\gamma_5\ell + C_P (\bar{b}\gamma_5(1-\gamma_5)s) \bar{\ell}\gamma_5\ell$



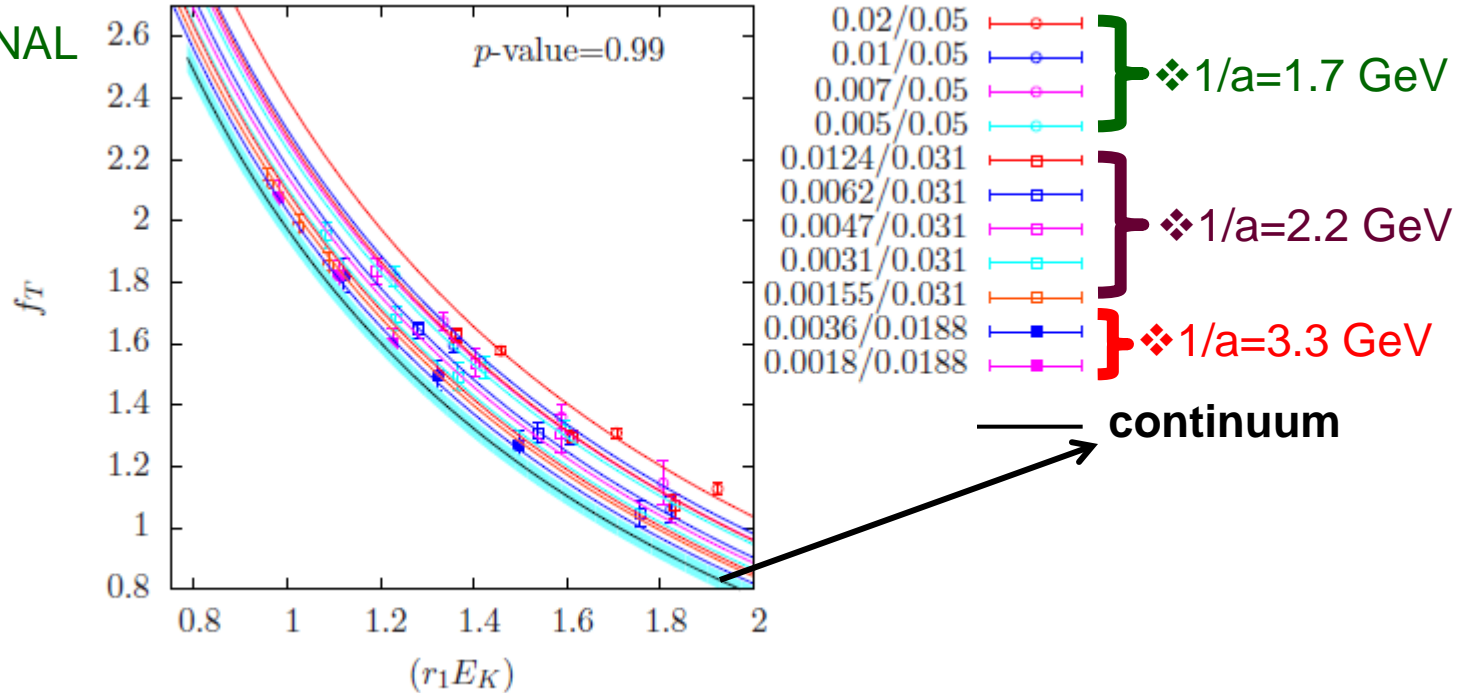
This "toy-scenario" would prefer nonzero  $C_P$ .

Lowering th. error on  $B \rightarrow Kll$ : 20% smaller than now

# $B \rightarrow K_{ll}$ - STRATEGY 2: (NRQCD) effective action for the $b$ quark, $v \ll c$

❖ sea quarks

❖ MILC-FNAL



❖ fit in  $q^2$ ,  $m_{sea}$ , and lattice spacing ( $a$ )

$$f_{\parallel} = \frac{C_{\parallel}^{(0)}}{f_{\pi}} \left[ 1 + \log s + C_{\parallel}^{(1)} m_l + C_{\parallel}^{(2)} (2m_l + m_s) + C_{\parallel}^{(3)} E_K + C_{\parallel}^{(4)} E_K^2 + C_{\parallel}^{(5)} a^2 \right],$$

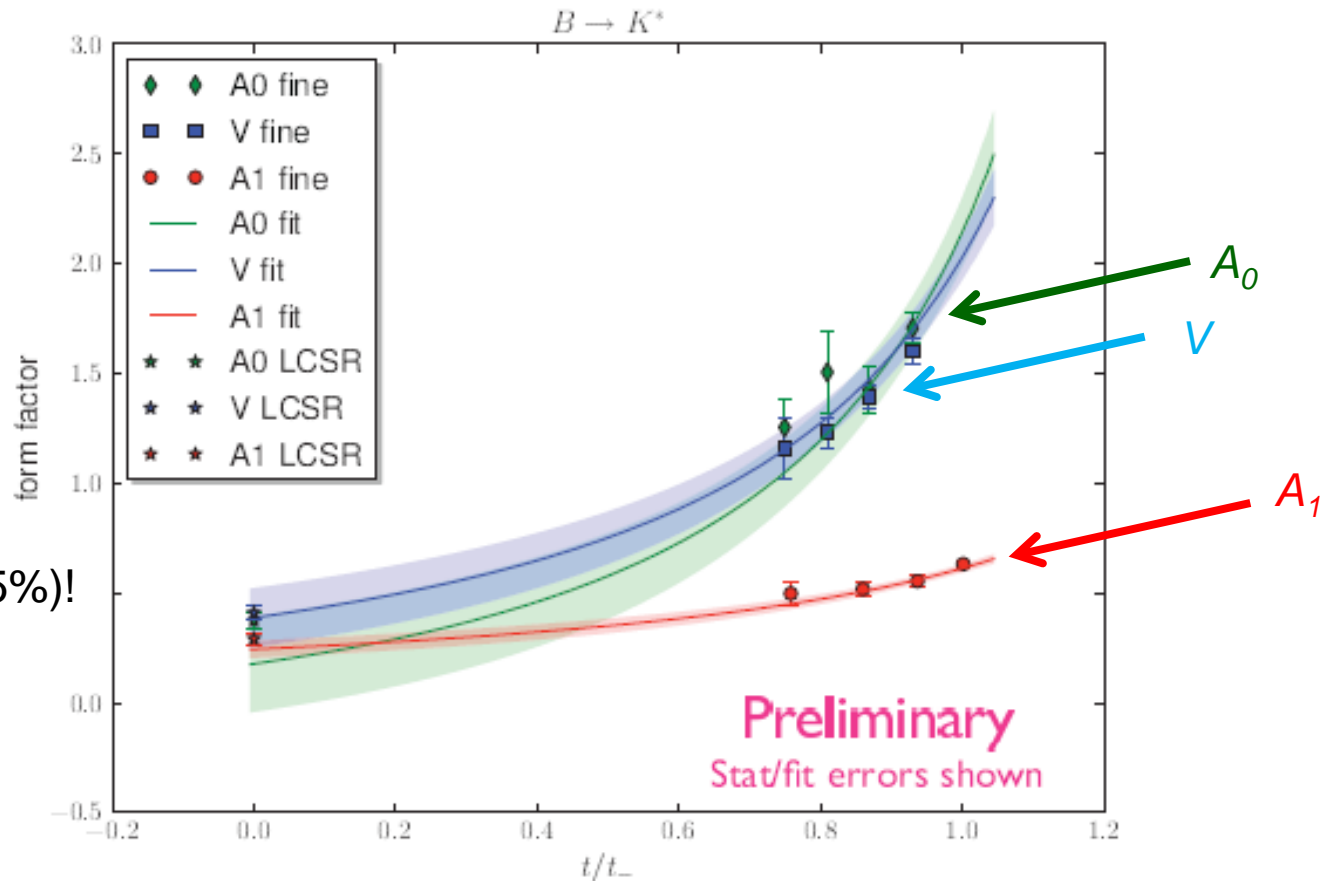
❖ preliminary unquenched results (NF=2+1)

-> stats + (systs) errors at 5%

3 coarse lattice spacings -. Continuum limit



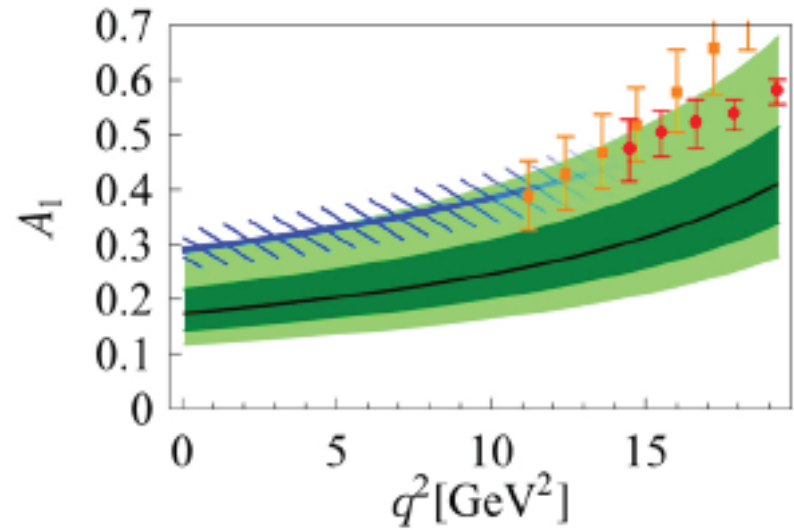
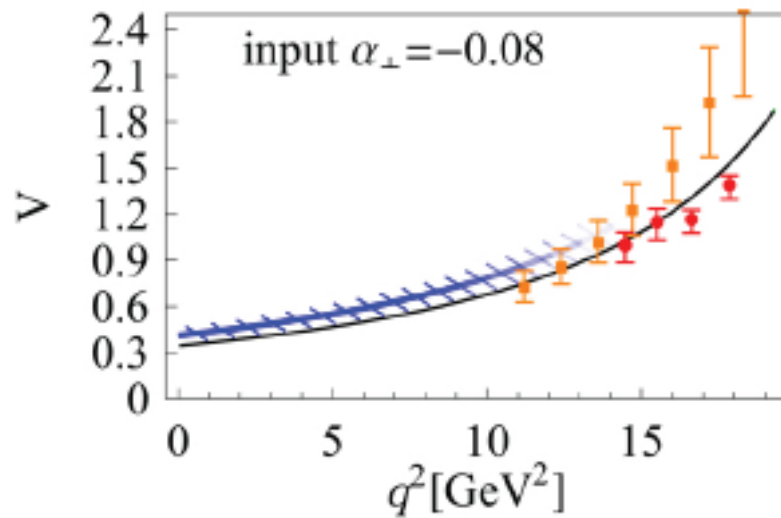
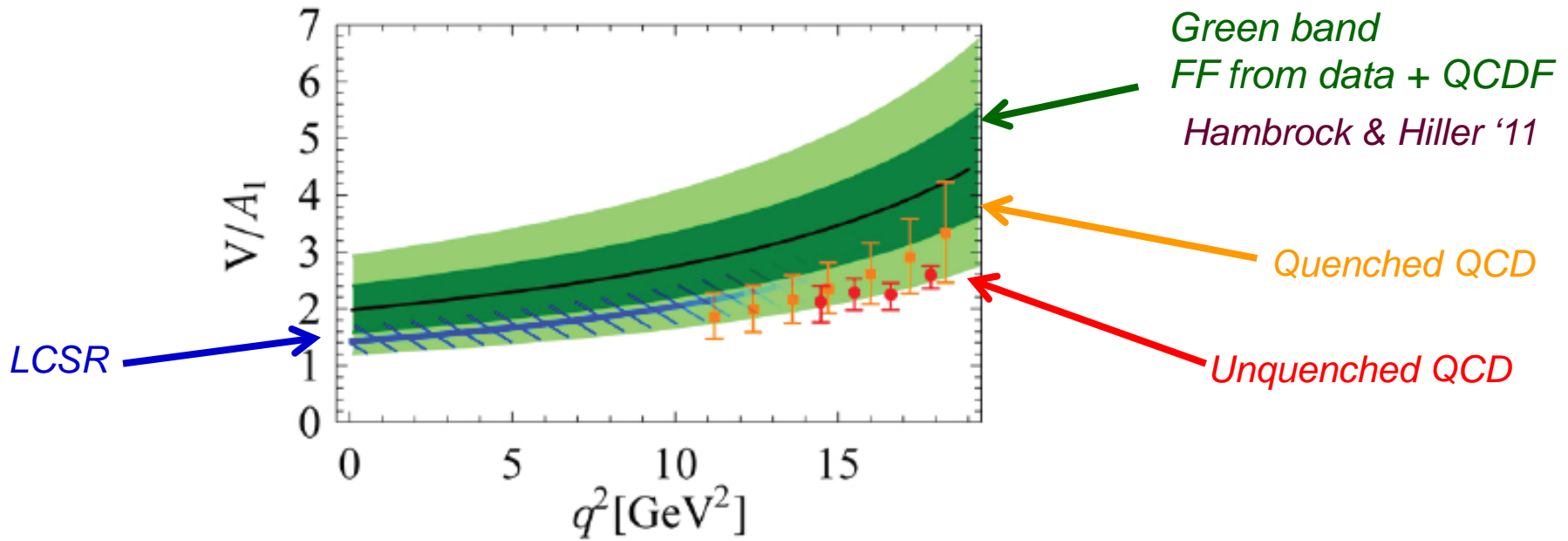
# $B \rightarrow K^* \parallel$ form factors from Cambridge/W&M/Edinburgh.



- Only stats errors (at 5%)!
- Promising study

Preliminary results on  $B \rightarrow K^* \parallel V, A_0,$  and  $A_1$  vs.  $q^2/q_{\max}^2$ . (by M. Wingate at lattice 2012)

# Comparison of $B \rightarrow K^* l l$ form factor calculations



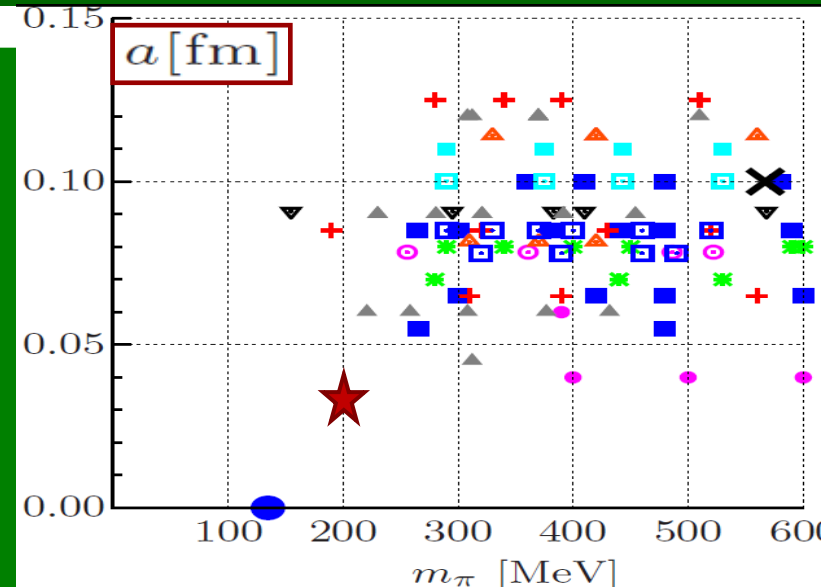
# Conclusions:

## LATTICE QCD -> touchable progress in recent years:

- ➔ reliable unquenched simulations with pions close to the physical point =>  $m_\pi=156$  MeV (PACS-CS),  $m_\pi=190$  MeV (BMW)
- ➔  $f_K/f_\pi$  &  $f_B$  paradigm of present lattice progress!
- ➔ promising studies at percent level on the way for B Physics ffs

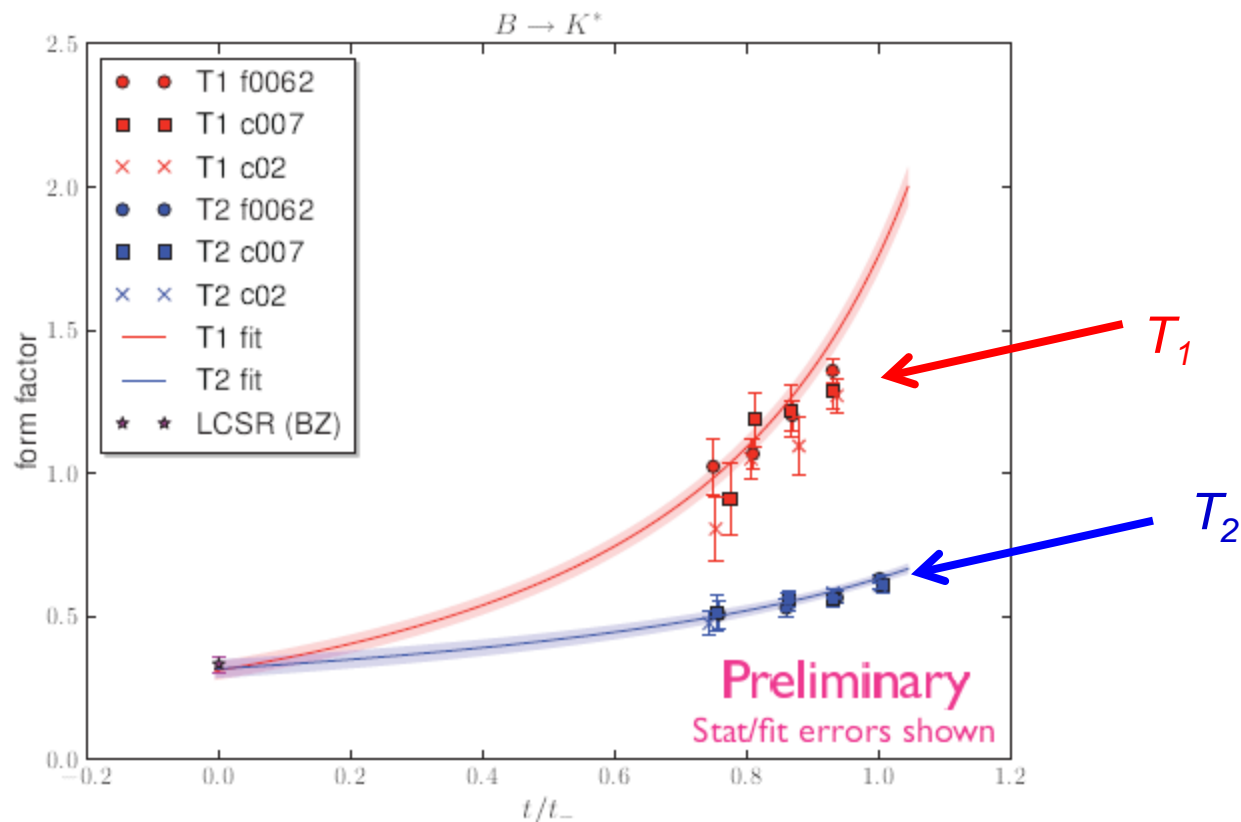
## Still a long work to assess 1%-precision needed for B physics

- ① discretization errors:  $a*m_B \ll 1$   
=>  $a \sim 0.033$  fm (6 GeV): ( $a \geq 0.07$  fm)
- ② finite volume effects:  $L*m_\pi \gg 1$   
=>  $L \geq 4.5$  fm ( $L \leq 3$  fm)
- ③ chiral regime:  $200 \leq m_\pi \leq 300$  MeV



courtesy of G. Herdoiza

# $B \rightarrow K^* \ell \ell$ form factors from Cambridge/W&M/Edinburgh.



: Preliminary results on  $B \rightarrow K^* \ell \ell$   $T_1$  and  $T_2$  vs.  $q^2/q_{\max}^2$  (by M. Wingate at lattice 2012)

- Only stats errors at 5%!
- Promising study

## Conclusions

✧  $Br(B_s \rightarrow \mu\mu)$  is genuinely sensitive to (pseudo)scalar operators

$$O'_S = (\bar{b}P_{R,L}S) \bar{\ell}\ell, \text{ and } O_P = (\bar{b}P_{R,L}S) \bar{\ell}\gamma_5\ell$$

➡ Only one hadronic parameter enters,  $f_{B_s} \rightarrow$  **small th. error**

✧  $Br(B \rightarrow K\ell\ell)$  is sensitive to (pseudo)scalar + vector operators (+ tensors)

➡ 3 hadronic parameters,  $f_{0,+T}$  form factors  $\rightarrow$  **large th. error**

➡ With respect to  $B_s \rightarrow \mu\mu$  it probes the effective Hamiltonian in the “orthogonal” direction!

➡ Improvement in form factors calculation would make the two observables a high resolution probe of scalar operators

➡ with tensor operators tested by  $A_{FB}(B \rightarrow K^{(*)}\ell\ell)$

➡ with vector ones by  $B \rightarrow X_s\ell\ell$  spectrum and transverse asymmetries in  $B \rightarrow K^*\ell\ell$

# New Physics: tensor contributions

- tensor contributions to  $B \rightarrow K \mu^+ \mu^-$

$$\mathcal{O}_T = \frac{e^c}{16\pi^2} (\bar{s} \sigma^{\mu\nu} b) (\bar{\ell} \sigma_{\mu\nu} \ell) \quad \mathcal{O}_{T5} = \frac{e^2}{16\pi^2} (\bar{s} \sigma^{\mu\nu} b) (\bar{\ell} \sigma_{\mu\nu} \gamma_5 \ell)$$

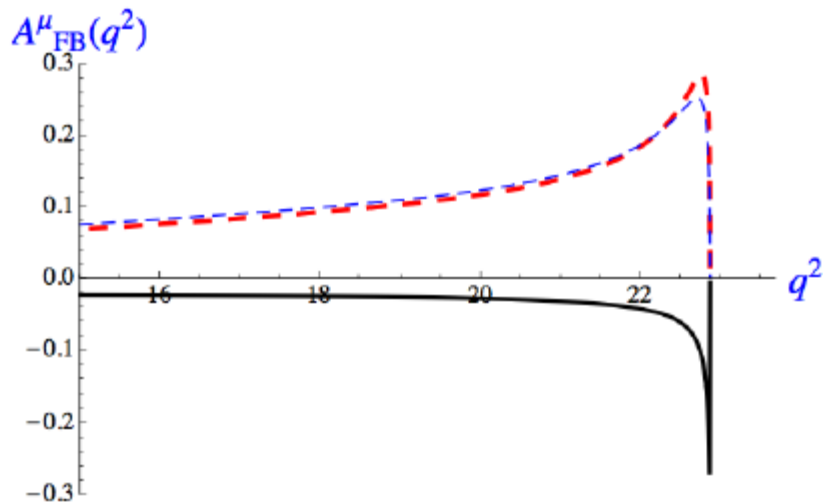
Assumption testable in

- $\mathcal{B}(B \rightarrow X_s \mu^+ \mu^-)$

$$\int_{1 \text{ GeV}^2}^{6 \text{ GeV}^2} \frac{d\mathcal{B}(B \rightarrow X_s \mu^+ \mu^-)}{dq^2} dq^2 = 1.59(17) \times 10^{-6} [1 + 0.59(2)(|C_T|^2 + |C_{T5}|^2)]$$

- Forward-backward asymmetry of  $B \rightarrow K \mu^+ \mu^-$

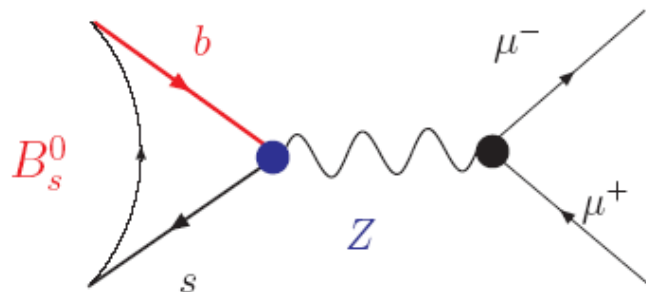
$$A_{FB}^\ell(q^2) = \frac{2 \mathcal{C}(q^2)}{\Gamma_\ell} \frac{m_B - m_K}{m_b} \sqrt{\lambda(q^2)} f_0(q^2) \left\{ (C_S C_T + C_P C_{T5}) q^2 f_T(q^2) \right. \\ \left. + m_\ell [C_S C_9 (m_B + m_K) f_P(q^2) + 2m_b (C_S C_7 + 2C_{T5} C_{10}) f_T(q^2)] + \mathcal{O}(m_\ell^2) \right\}$$



- $C_T = C_{T5} = 1.6$
- Thick  $\rightarrow C_S = C_P = 0$
- Dashed  $\rightarrow \{C_S, C_P\} = (1, 0)$

# 2013 Experimental situation

$B_s \rightarrow \mu\mu$



LHCb'12

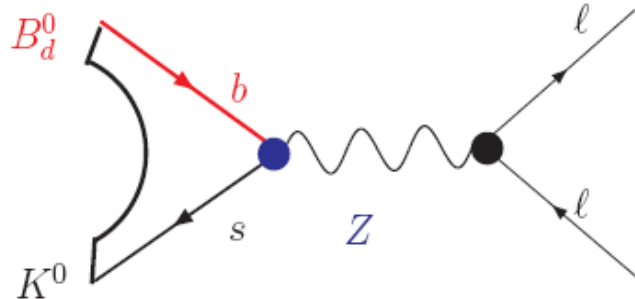
$$\text{Br}^{\text{exp}}(B_s \rightarrow \mu\mu) = (3.2 \pm 1.5)10^{-9}$$

exp. error 50%

$$\text{Br}(B_s \rightarrow \mu\mu)^{\text{SM}} = (3.3 \pm 0.3) \times 10^{-9}$$

th. error 6.5%

$B_d^0 \rightarrow K^0 \ell\ell$



$B_u^+ \rightarrow K^+ \ell\ell$

BaBar'12-04

$$\text{Br}(B \rightarrow K\ell\ell) = (4.7 \pm 0.6) \times 10^{-7}$$

- averaging  $B^+/B^0$  and  $e/\mu$

LHCb'12-05

$$\text{Br}(B^0 \rightarrow K^0 \mu\mu) = (3.1 \pm 0.7) \times 10^{-7}$$

- puzzle  $B^+$  and  $B^0$  asymmetries?



$$\text{(our) Br}(B \rightarrow K\ell\ell)^{\text{SM}} = (7.0 \pm 1.8) \times 10^{-7}$$