Combining Z' and W' searches at the LHC

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Outlook

- Introduction: leptonic signatures
- Singlets
- Triplets
- Mixture
- Combined limits

Z' and W' searches at LHC

Cleanest signatures:

Dilepton — Neutral resonance
 Lepton + MET — Charged resonance

Effective Theory Field content: SM + new resonances Symmetries: Lorentz, SU(3)xSU(2)xU(1) Leading order: Local 4D operators

Vector resonances in leptonic final states:

- Single production
- Renormalizable interactions with SM fields
- Couple to quarks and leptons

Direct couplings

 $\mathcal{B} \in (1,1)_0$ $\mathcal{W} \in (1,3)_0$ Couplings via mixing with $Z \mbox{ and } W$

 $\mathcal{B}^1 \in (1,1)_1$ $\mathcal{W}^1 \in (1,3)_1$

Vector resonances in leptonic final states:

EWPT with Higgs at 126 GeV



Small Mixing ϕ

Direct couplings

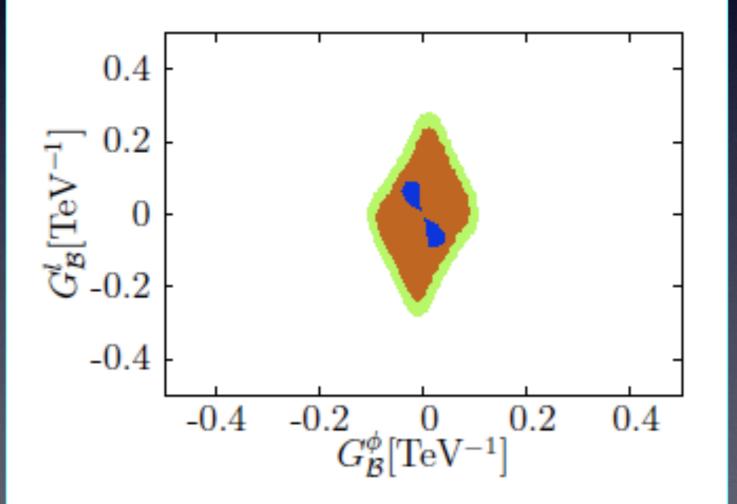
 $\mathcal{B} \in (1,1)_0$ $\mathcal{W} \in (1,3)_0$ Couplings via mixing with $Z \mbox{ and } W$

 $\mathcal{B}^1 \in (1,1)_1^{\varphi} \triangleleft 0.00_1$

 $1, 3]_1$ $\phi < 0.003$

Singlets $\mathcal{B} \in (1,1)_0 \rightarrow Z'$ (Universal) Couplings g_q, g_u, g_d, g_l, g_e

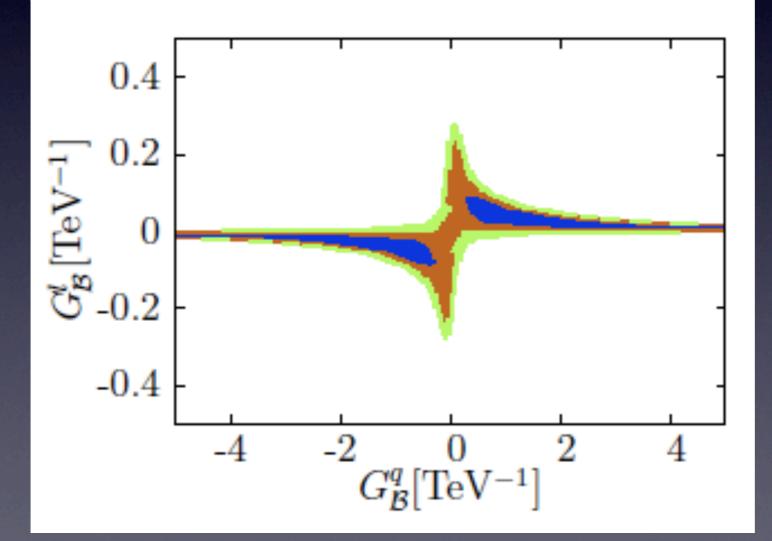
EWPT with Higgs at 126 GeV



$\begin{array}{l} \text{Mixing} \\ \phi \lesssim 0.002 \end{array}$

Singlets $\mathcal{B} \in (1,1)_0 \rightarrow Z'$ (Universal) Couplings g_q, g_u, g_d, g_l, g_e





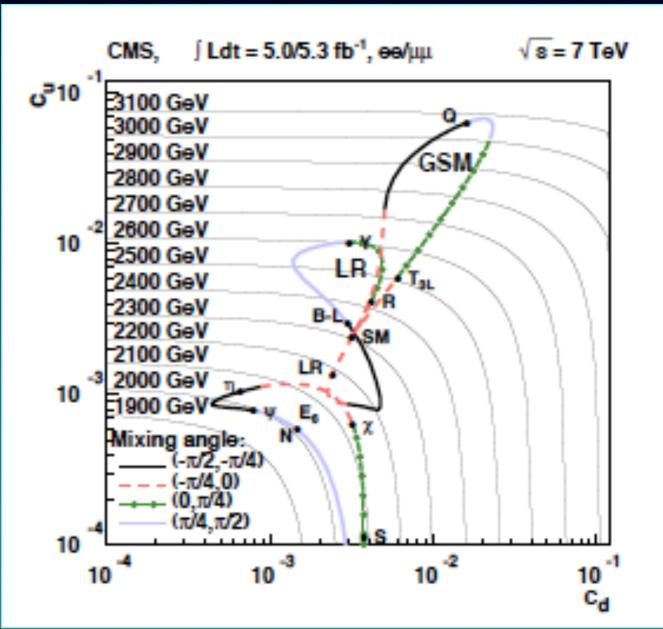
LHC searches

Narrow Width Approximation

$$\sigma(pp \to l^+l^-) = \frac{\pi}{6s} \left[c_u \omega_u(s, m_{Z'}^2) + c_d \omega_d(s, m_{Z'}^2) \right]$$

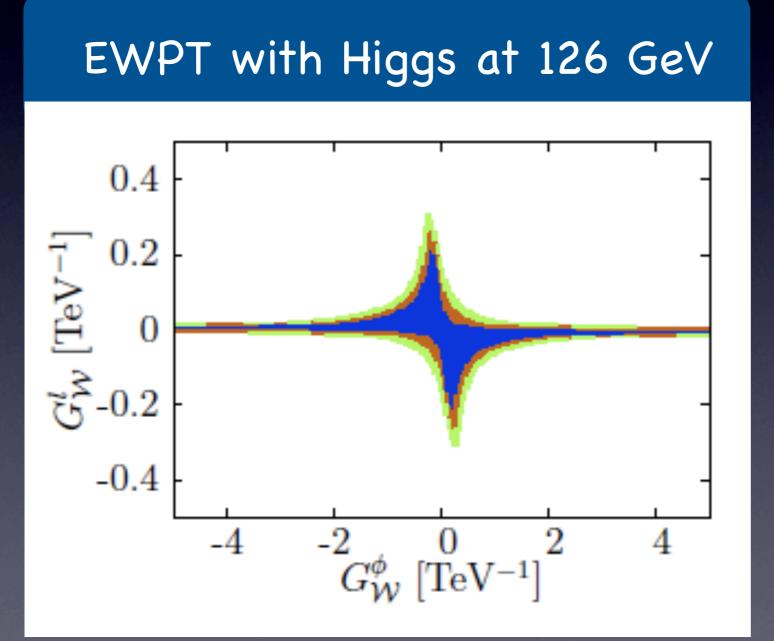
[Carena, Daleo, Dobrescu, Tait '04]

$$\begin{aligned} \mathbf{c}_{u,d} &= \left(g_q^2 + g_{u,d}^2\right) \\ &\times \frac{g_l^2 + g_e^2}{2g_l^2 + g_e^2 + 6g_q^2 + 3g_u^2 + 3g_d^2} \end{aligned}$$

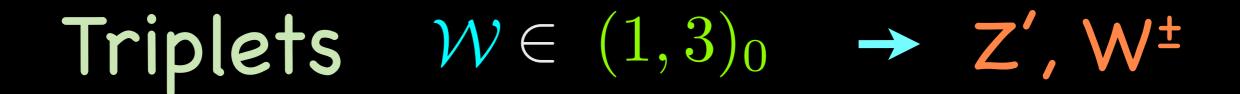


Triplets $\mathcal{W} \in (1,3)_0 \rightarrow Z', W^{\pm}$

(Universal) Couplings $g_q, \overline{g_l}$

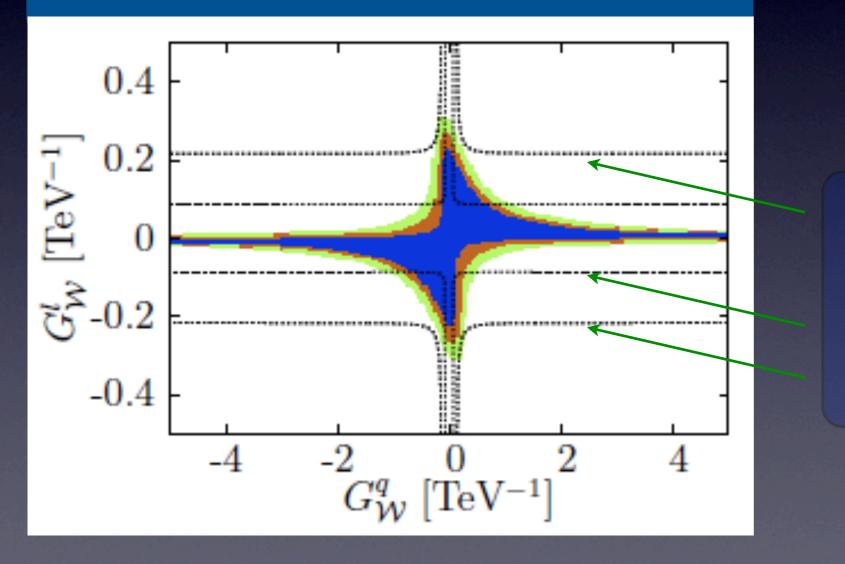


Preserves custodial symmetry (p=1), but mixings constraint for significant lepton couplings, $\varphi \lesssim 0.01$ and also by perturbativity, $\varphi \lesssim 0.1$



(Universal) Couplings g_q, g_l





Constant

$$\tilde{g} = \frac{2}{g_q g_l} \sqrt{3g_q^2 + g_l^2}$$

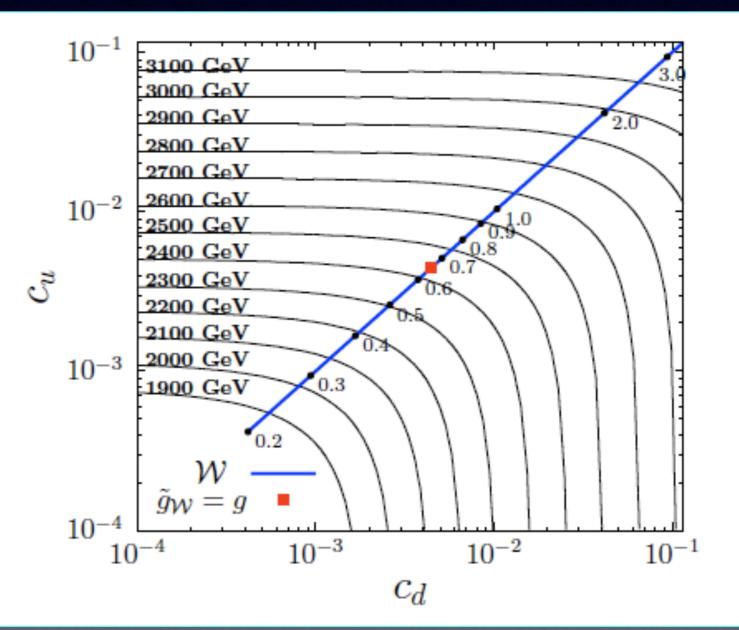
LHC searches

$$\sigma(pp \to l^+ l^-) = \frac{\pi}{6s} \begin{bmatrix} c_u \omega_u(s, m_{Z'}^2) + c_d \omega_d(s, m_{Z'}^2) \end{bmatrix}$$

$$\sigma(pp \to l^\pm \nu) = \frac{\pi}{6s} c_c \omega_c(s, m_{W'}^2) \qquad \text{NWA}$$

 $c_{u,d} = \frac{\tilde{g}^2}{96}$ $c_c = \frac{\tilde{g}^2}{24}$

$$\tilde{g} = \frac{2}{g_q g_l} \sqrt{3g_q^2 + g_l^2}$$



Singlet
$$\mathcal{W} \in (1,3)_0$$
 \mathcal{W}^{\pm} ,+Triplet $\mathcal{B} \in (1,1)_0$ \mathcal{W}^{\pm} ,

$$\mathcal{L} \supset g_{\mathcal{WB}} \mathcal{W}^{a}_{\mu} \mathcal{B}^{\mu} \phi^{\dagger} \frac{\sigma_{a}}{2} \phi$$

Neutral mass-eigenstates: $Z'_{1} = \cos \theta \mathcal{W}^{3} - \sin \theta \mathcal{B}$ $Z'_{2} = \sin \theta \mathcal{W}^{3} + \cos \theta \mathcal{B}$

Eigenmasses:

$$M_{Z_1'}^2 = M_{\mathcal{W}}^2 - \Delta \sin^2 \theta$$
$$M_{Z_2'}^2 = M_{\mathcal{W}}^2 + \Delta \cos^2 \theta$$

$$\operatorname{an} 2\theta = \frac{g_{\mathcal{W}\mathcal{B}} v^2}{2\left(M_{\mathcal{B}}^2 - M_{\mathcal{W}}^2\right)} \qquad |\Delta| = \sqrt{\left(M_{\mathcal{B}}^2 - M_{\mathcal{W}}^2\right)^2 + g_{\mathcal{W}\mathcal{B}}^2 \frac{v^4}{4}}$$

Singlet $\mathcal{W} \in (1,3)_0$ \mathcal{W}^{\pm} ,+Triplet $\mathcal{B} \in (1,1)_0$ \mathcal{X}'_1 , \mathcal{Z}'_2

Scenario 1

 $|M_{\mathcal{B}}^2 - M_{\mathcal{W}}^2| \gg v^2$

$Z'_2 \simeq Z_{\mathcal{B}}$

W'

 $Z_1' \simeq Z_{\mathcal{W}}'$

Scenario 2 $|M_{\mathcal{B}}^2 - M_{\mathcal{W}}^2| \sim v^2$

$$Z_2^\prime W^\prime \ Z_1^\prime$$



LHC searches Scenario 2

- Charged sector as in triplet

standard Breit-Wigner not valid [Cacciapaglia, Deandrea,
 non-diagonal imaginary terms De Curtis, '09]



Scenario 2

Generalized NWA taking interference into account

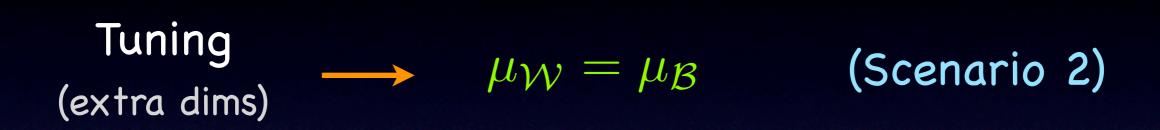
$$c_q = \frac{1}{6\pi \text{Tr}\Sigma} \left\{ \text{Tr}\left[G_l G_q\right] + \frac{\text{Tr}\left[G_l \tilde{\Sigma} G_q \tilde{\Sigma}\right]}{\det \Sigma} \right\}$$

$$G_{ij}^{f} = \frac{1}{2} \left[(g_{L}^{f})_{i} (g_{L}^{f})_{j} + (g_{R}^{f})_{i} (g_{R}^{f})_{j} \right]$$

$$\Sigma_{ij} = \frac{1}{12\pi} \sum_{f=u,d,e,\nu} G^f_{ij}$$

$$\tilde{\Sigma}_{ij} = \epsilon_{jm} \epsilon_{in} \Sigma_{mn}$$

Example: Generalized Sequential Model $[SU(2) \times U(1)]^2 \rightarrow SU(2) \times U(1)$

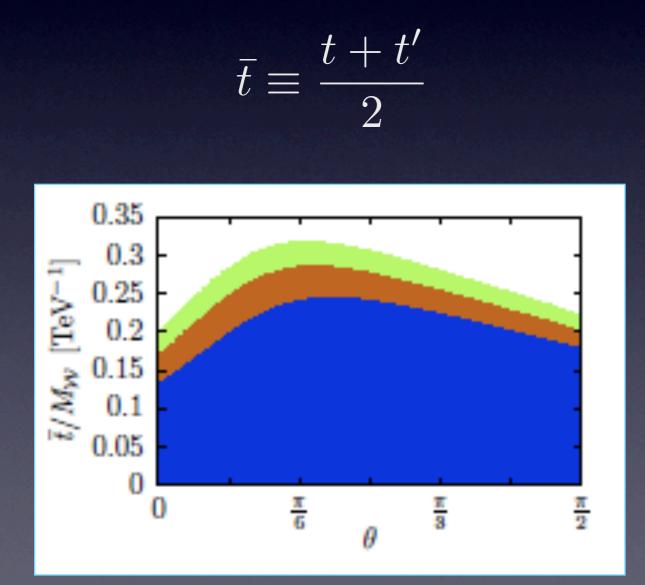


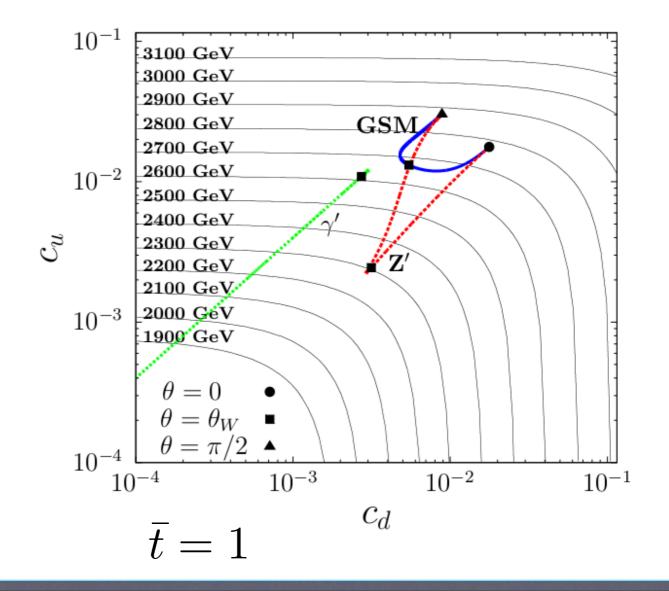
Mass matrix: shifted version of SM

• Couplings of (W', Z'_1, Z'_2) equal to the ones of (W, Z, γ) with the replacements $g \to gt, \quad \theta_W \to \theta \in [0, \pi/2] \quad (g' \to g't')$

Example: Generalized Sequential Model $\left[SU(2) imes U(1) ight]^2 ightarrow SU(2) imes U(1)$

$g \to gt, \quad \theta_W \to \theta \in [0, \pi/2] \quad (g' \to g't')$





 Vectorial I+MET resonance comes with vectorial di-lepton resonance(s)

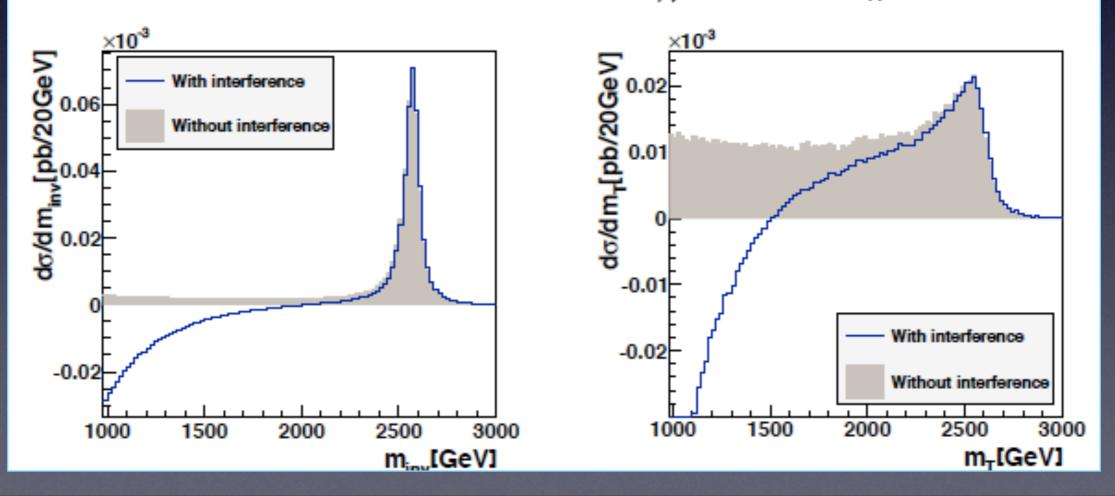
Nearly-degenerate masses

Related couplings

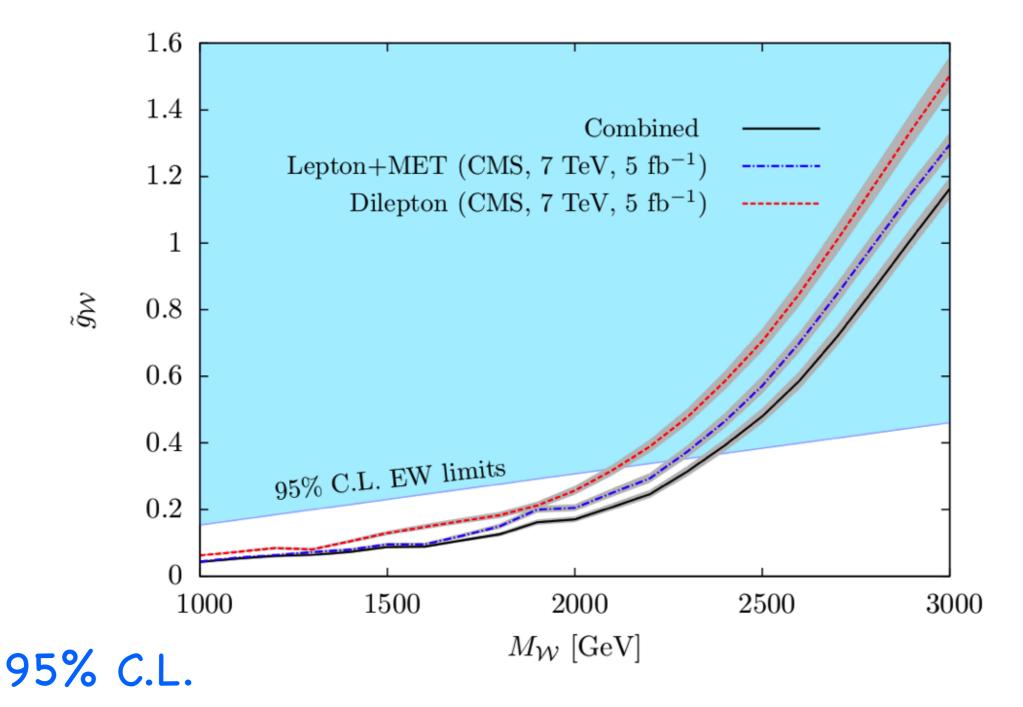
Combine I+MET and di-lepton data in a single likelihood

- Data from CMS, 5 fb⁻¹, 7 TeV
- Montecarlo based on Madgraph-Pythia-PGS
- Interference of nearby heavy resonances
- Interference with SM amplitudes

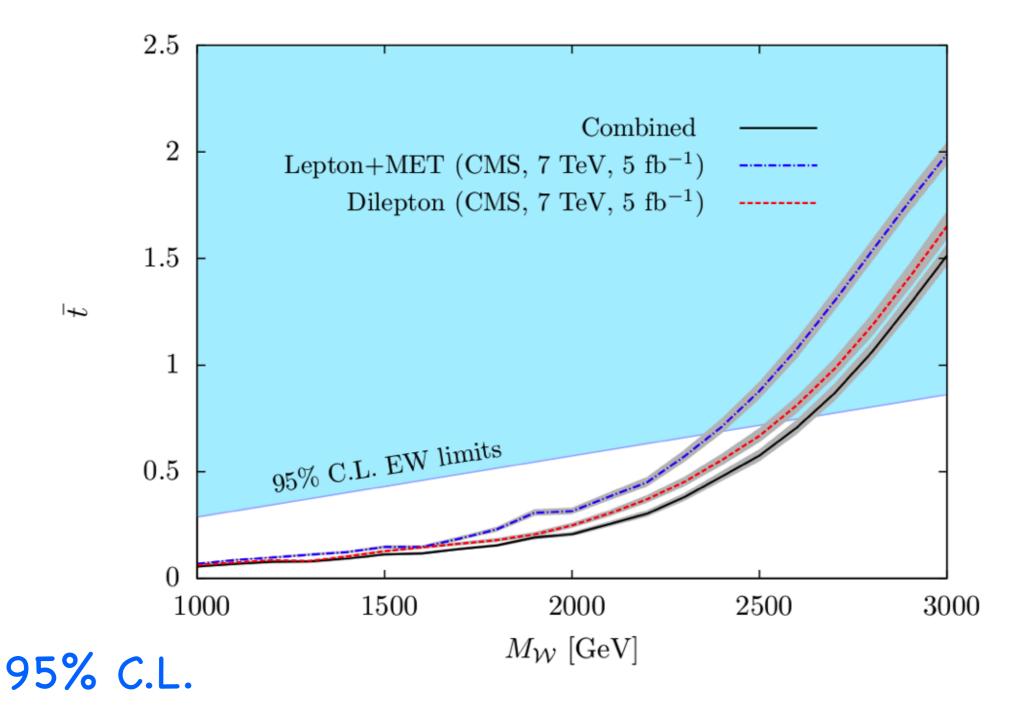
Simulation at CMS at $\sqrt{s} = 7$ TeV, for a GSM ($M_W = 2.6$ TeV, $\theta = \theta_W, \bar{t}=1$):



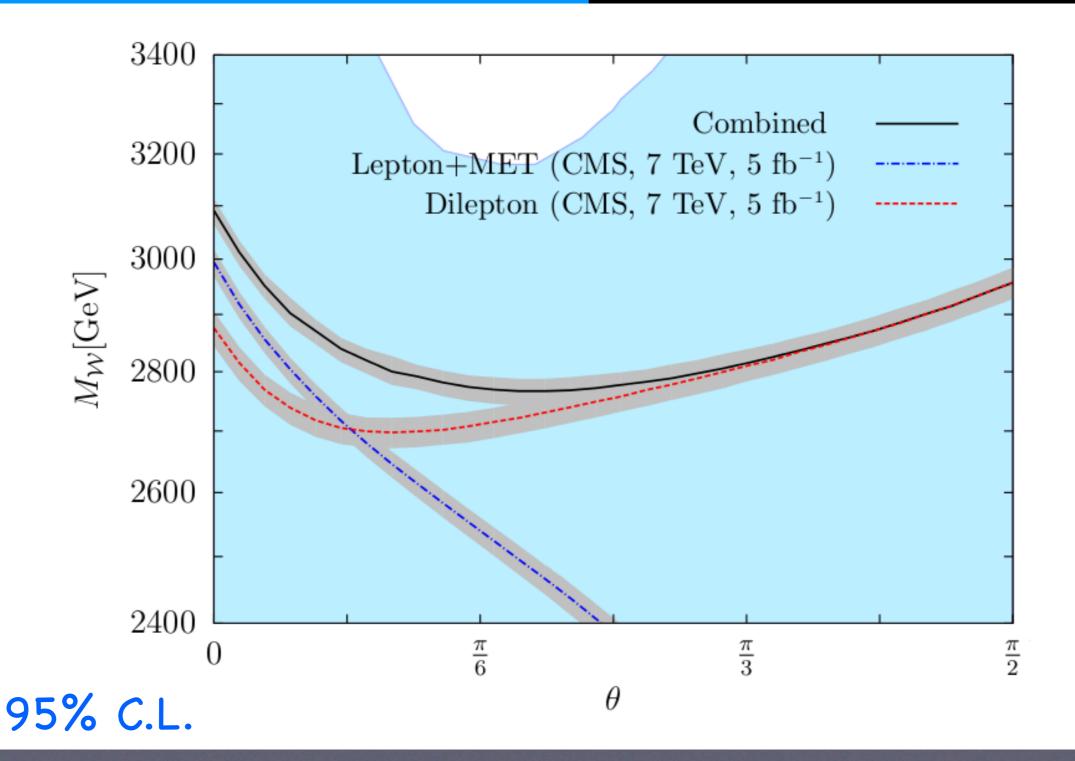












Conclusions

 W' decaying into charged lepton-neutrino is accompanied by Z' of similar mass decaying into two charged leptons

 Combination of neutral and charged channels improves individual limits

EWPT competitive with LHC for large couplings /masses



