# Information and treatment of unknown correlations in BLUE combinations 

## Andrea Valassi (CERN IT-ES) <br> Roberto Chierici (IPN Lyon)

Top LHC Working Group, 29th November 2012

## Motivation

- Best Linear Unbiased Estimators (BLUE) are widely used for precision measurement combinations
- e.g. top mass @ Tevatron and @ LHC
- Larger event statistics have led (and will lead more and more) to combinations dominated by systematic errors
- And by the correlations between those systematics...
- In parallel, negative combination coefficients are appearing (and may appear more and more) in BLUE averages
- We should not confuse these coefficients ("weights" in BLUE weighted central values) with the statistical impact (another type of "weight") of a measurement
- Answer to typical questions
- How can we assess the "relative importance" of a measurement ("how much is my experiment/measurement contributing to a combination")?
- What is the scientific interpretation of negative BLUE coefficients ("are we correctly estimating the correlations used as inputs to the combination")?


## Main messages in this talk

1. Absolute values of (negative) BLUE coefficients should not be used to assess the "relative importance" of measurements

- We suggest an alternative procedure based on Fisher's information

2. Negative BLUE coefficients indicate a "high correlation" regime - We show that this is true in general for N measurements of one observable
3. Correlation estimates in this regime are not at all "conservative" and may need to be critically reassessed

- Overestimated correlations may lead to largely underestimated combined systematics and total errors
- We propose a few tools to help in this critical review of correlations (derivatives of information, information minimization, covariance "onionization", marginal inflow of information)
- But the ideal solution is still to measure correlations in data and MC
- Part 1 - information in BLUE
- "Relative importance" of the measurements in the presence of correlations
- Part 2 - interpretation of negative BLUE coefficients
- "Low-correlation" and "high-correlation" regimes
- Part 3 - "conservative" estimates of correlations?
- The risks in overestimating correlations and a few hints to try and avoid this
- Conclusions


## Part I: information in BLUE

- "Relative importance" of individual measurements with correlations
- BLUE primer
- What is the "contribution to knowledge" from various sources
- Absolute values of BLUE coefficients are used in the TOPLHCWG note
- The problem with absolute values of BLUE weights
- Fisher's information
- A proposed new way for presenting combinations: information weights
- First look at TOP LHC WG combination
- Simple case [Lyons] : measurements $\mathrm{y}_{\mathrm{i}}$ of a single observable Y
- A linear estimator $\quad \hat{Y}=\tilde{\lambda} \mathbf{y}=\sum_{i=1}^{n} \lambda_{i} y_{i} \quad$ is unbiased if $\tilde{\mathbf{U}} \boldsymbol{\lambda}=\sum_{i=1}^{n} \lambda_{i}=1$ [where U is a vector of all 1 's, i.e. contraction by U means summing all indexes]
- The square of the error on the combined estimate $\hat{\mathbf{Y}}$ is $\operatorname{var}(\hat{Y})=\sigma_{\hat{Y}}^{2}=\sum_{i=1}^{n} \sum_{j=1}^{n} \lambda_{i} \mathcal{M}_{i j} \lambda_{j}$ [where M is the $(\mathrm{NxN})$ input covariance of the N measurements] NB Assume Gaussian pdf with M known a priori and independent of unknown parameter
- The $\underline{\text { best linear unbiased estimate }}$ is given by $\lambda_{i}=\frac{\left(\mathcal{M}^{-1} \mathbf{U}\right)_{i}}{\left(\tilde{\mathbf{U}} \mathcal{M}^{-1} \mathbf{U}\right)}$
(i.e. the error is minimized for)
- The coefficients $\lambda_{i}$ in the BLUE linear combination can be negative [a.k.a. "weights" in the BLUE weighted average]
- The combined error on the BLUE is equal to: $\quad \operatorname{var}(\hat{Y})=\sigma_{\hat{Y}}^{2}=\frac{1}{\left(\tilde{\mathbf{U}} \mathcal{M}^{-1} \mathbf{U}\right)}$


## Negative weights (and absolute weighs) @LHC



BLUE Combination Coefficient [\%]


TOPLHC NOTE
ATLAS-CONF-2012-095
CMS PAS TOP-12-001
June 2012

Figure 1: (a): Input measurements and result of the LHC combination (see also Table 4); (b, c) : BLUE combination coefficients and relative importance of the input measurements. The relative importance of a measurement is defined as $\left|w_{i}\right| / \sum_{i}\left|w_{i}\right|$, where $\left|w_{i}\right|$ is the absolute value of the BLUE combination coefficient for the $i^{\text {th }}$ input measurement.

## The problem with absolute values

- Take A, B, C, with C uncorrelated to $\mathbf{A}$ or $\boldsymbol{B}: ~\left(\begin{array}{cc|c}\sigma_{A}^{2} & \rho_{A B} \sigma_{A} \sigma_{B} & 0 \\ \rho_{A B} \sigma_{A} \sigma_{B} & \sigma_{B}^{2} & 0 \\ \hline 0 & 0 & \sigma_{C}^{2}\end{array}\right)$
- Combining ( $\mathrm{A}, \mathrm{B}, \mathrm{C}$ ) in one go gives the same result for the combined error and for the linear coefficient of $C$ as first combining $(A, B)$ and then adding $C$ :

$$
\frac{1}{\left(\sigma_{\hat{Y}}^{2}\right)_{A B C}}=\frac{1}{\left(\sigma_{\hat{Y}}^{2}\right)_{A B}}+\frac{1}{\sigma_{C}^{2}}
$$

$$
\left(\lambda_{C}\right)_{A B C}=\frac{1 / \sigma_{C}^{2}}{1 /\left(\sigma_{\hat{Y}}^{2}\right)_{A B}+1 / \sigma_{C}^{2}}
$$

- The "relative importance" (RI) of C based on absolute values in combining $(A, B, C)$ in one go is

$$
\left(\mathrm{RI}_{C}\right)_{A, B, C}=\frac{\left|\left(\lambda_{C}\right)_{A B C}\right|}{\left|\left(\lambda_{A}\right)_{A B C}\right|+\left|\left(\lambda_{B}\right)_{A B C}\right|+\left|\left(\lambda_{C}\right)_{A B C}\right|}
$$

- It can easily be shown that there are two problems with this definition if $\lambda_{B}<0$ (which happens if $\rho_{A B}>\sigma_{A} / \sigma_{B}$ with $\sigma_{A}<\sigma_{B}-$ see later in this talk):
- first, $R I$ of $C$ is different if $C$ is combined with $A$ and $B$ separately(lower) or if it is combined with the combination of A and B (higher)
- second, RI of C underestimates the contribution of C with respect to its more intuitive (and "correct"?) assessment $\frac{1 / \sigma_{C}^{2}}{1 /\left(\sigma_{\hat{Y}}^{2}\right) A B+1 / \sigma_{C}^{2}}=\left(\lambda_{C}\right)_{A B C}$
$\rightarrow$ the "information weight" we propose for $C$ is exactly equal to $\left(\lambda_{C}\right)_{A B C}$



## Fisher's information primer

- Given a pdf $\mathrm{P}(\mathrm{y}, \mathrm{X})$ for the N measurements $\mathrm{y}_{\mathrm{i}}$ and the n parameters $\mathrm{X}_{\mathrm{a}}$, the $(n \times n)$ information matrix $I_{\alpha \beta}$ is defined as the covariance of $\partial \log P / \partial X_{\alpha}$
- Information has many nice properties [James]:
- Information is additive for independent measurements
- Information $I_{\alpha \beta}$ is defined with respect to the set of parameters $X_{\alpha}$ to be estimated
- Information is related to precision: the greater the available information, the more precise the estimates can be - the inverse of the information matrix is a lower bound (Cramer-Rao) for the covariance of any estimates of the parameters $X_{\alpha}$
- Information is positive semi-definite (for one observable it is a scalar $\geq 0$ )
- Assuming P is multivariate Gaussian, the covariance of the BLUE estimates is equal to the inverse of the information matrix
- For a single observable (e.g. $\mathrm{m}_{\text {top }}$ ), information is the inverse of the combined variance

$$
I^{(Y)}=\tilde{\mathbf{U}} \mathcal{M}^{-1} \mathbf{U}=\frac{1}{\operatorname{var}(\hat{Y})}=\frac{1}{\sigma_{\hat{Y}}^{2}}
$$

- For all practical purposes, in the following, "information" about one parameter is just a synonym for the inverse of its variance
- WE WILL ALWAYS BE TALKING ABOUT A SINGLE OBSERVABLE IN WHAT FOLLOWS


## Our proposal: Information Weights

- It is not only individual measurements, but also their correlation, that contribute to the information available in a combination
- The information contributed by correlations comes from the collective interplay of the measurements and cannot (in general) be attributed to any of them individually
- We propose to distinguish between the information weights from the measurements and the information weight from the ensemble of correlations
- Our definition of Information Weights (IW)
- Information weight for each measurement (always >0):
- In the absence of any correlations, this coincides with the BLUE linear coefficient $\lambda$

$$
\mathrm{IW}_{i}=\frac{1 / \sigma_{i}^{2}}{1 / \sigma_{\hat{Y}}^{2}}
$$

- Also true if one measurement is uncorrelated (e.g. C in previous $A, B, C$ example)
- Information weight for the correlations is all the rest (may be $<0$ ):
$\mathrm{IW}_{\text {corr }}=\frac{1 / \sigma_{\hat{Y}}^{2}-\sum_{i} 1 / \sigma_{i}^{2}}{1 / \sigma_{\hat{Y}}^{2}}$
- For N measurements of 1 observable: \# information weights are N (measurements) plus 1 (correlations)
- One issue: not easy/obvious how to subdivide the correlation weight
- e.g. split up contributions by error source and/or by pair of measurements?
- the interplay of correlations cannot be neglected, their effect is largely a collective one
- will come back to this only at the very end


## A practical example (more later): LHC m $_{\text {top }}$

| ATLAS + CMS Preliminary, $\sqrt{s}=7 \mathrm{TeV}$ |  |
| :---: | :---: |
|  |  |
| ATLAS 2010, I+jets $L_{\text {nt }}-35 \mathrm{pb}^{-1}, \text {, ( } \oplus \text { CR, UE syst.) }$ | -6.8 |
| ATLAS 2011, I+jets | 29.9 |
| ATLAS 2011, all jets $\mathrm{L}_{\mathrm{m}}=2 \mathrm{ff}^{-1}$, ( $\oplus$ CR, UE syst $)$ | -0.4 |
| CMS 2010, di-lepton $\mathrm{L}_{\mathrm{mt}}=36 \mathrm{pb}^{-1}$, ( $\oplus$ CR syst) | -1.9 |
| CMS 2010, I + jets $\mathrm{L}_{\mathrm{cta}}=36 \mathrm{pb} \mathrm{b}^{\mathbf{1}},(\oplus \mathrm{CR}$ syst) | -0.2 |
| CMS 2011, di-lepton $L_{\mathrm{mt}}=2.3 \mathrm{fb}$.', ( $\oplus$ CR, UE syst.) | -4.8 |
| CMS 2011, $\mu+j$ jets $\mathrm{L}_{\text {nt }}-4.9 \mathrm{fb}^{-1}$, , $\oplus \mathrm{CR}, \mathrm{UE}$ syst.) | 84.3 |
|  | LHC $m_{\text {too }}$ combination June 2012 |
| - | 1 |
| BLUE Combination Coefficient [\%] |  |



|  | $m_{\text {top }}$ <br> $[\mathrm{GeV}]$ | $\Delta I$ <br> $\left[1 / \mathrm{GeV}^{2}\right]$ | $\lambda$ | $\frac{\|\lambda\|}{\sum\|\lambda\|}$ | $\frac{\Delta I}{I}$ |
| :--- | :---: | ---: | ---: | ---: | ---: |
| CMS11 $\mu \mathrm{j}$ | $172.64 \pm 1.53$ | 0.425 | $84.3 \%$ | $65.7 \%$ | $83.8 \%$ |
| ATL11lj | $174.53 \pm 2.39$ | 0.175 | $29.9 \%$ | $23.3 \%$ | $34.5 \%$ |
| ATL10lj | $169.33 \pm 6.32$ | 0.025 | $-6.8 \%$ | $5.3 \%$ | $4.9 \%$ |
| CMS11ll | $173.30 \pm 2.96$ | 0.114 | $-4.8 \%$ | $3.7 \%$ | $22.5 \%$ |
| CMS10ll | $175.50 \pm 6.49$ | 0.024 | $-1.9 \%$ | $1.5 \%$ | $4.7 \%$ |
| ATL11aj | $174.90 \pm 4.44$ | 0.051 | $-0.4 \%$ | $0.3 \%$ | $10.0 \%$ |
| CMS10lj | $173.10 \pm 3.41$ | 0.086 | $-0.2 \%$ | $0.2 \%$ | $16.9 \%$ |
| Correlations | - | -0.392 | - | - | $-77.4 \%$ |
| Total | $173.33 \pm 1.40$ | 0.507 | $100.0 \%$ | $100.0 \%$ | $100.0 \%$ |

The real challenge (and the interesting part) is splitting up the correlation contribution and understanding where it comes from!

## Part II: negative BLUE coefficients

"Low-correlation" and "High-correlation" regimes

- The simple case of two measurements of one observable [Lyons]
- Dependencies of BLUE coefficients and combined error and IW as the correlation varies
- One linear coefficient $\lambda$ goes to 0 and flips sign where the combined error is minimum
- Interpretation of information in the two regimes $\lambda>0$ and $\lambda<0$
- Generalization to the case of N measurements of one observable
- Information inflow - adding the $\mathrm{N}^{\text {th }}$ measurement with its $\mathrm{N}-1$ correlations
- Information derivatives - analyzing $\mathrm{N}_{\mathrm{x}}(\mathrm{N}-1) / 2$ correlations independently
- Low-correlation regime (all $\lambda>0$ ) and high-correlation regime (one or more $\lambda<0$ )

WE WILL ONLY DISCUSS POSITIVE CORRELATIONS IN THE FOLLOWING

## 2 measurements of 1 observable

- For two measurements A and B with errors $\sigma_{A}<\sigma_{B}$ and correlation $\rho$ [Lyons etall:

$$
\begin{aligned}
\hat{Y}=\lambda_{A} y_{A}+\lambda_{B} y_{B} & \lambda_{A} & =\frac{\sigma_{B}^{2}-\rho \sigma_{A} \sigma_{B}}{\sigma_{A}^{2}+\sigma_{B}^{2}-2 \rho \sigma_{A} \sigma_{B}} \\
\sigma_{\hat{Y}}^{2}=\frac{\sigma_{A}^{2} \sigma_{B}^{2}\left(1-\rho^{2}\right)}{\sigma_{A}^{2}+\sigma_{B}^{2}-2 \rho \sigma_{A} \sigma_{B}}=\frac{1}{I} & \lambda_{B} & =\frac{\sigma_{A}^{2}-\rho \sigma_{A} \sigma_{B}}{\sigma_{A}^{2}+\sigma_{B}^{2}-2 \rho \sigma_{A} \sigma_{B}}
\end{aligned}
$$

- Varying $\rho$, the combined error is maximum (equal to $\sigma_{A}$ ) for $\rho=\sigma_{A} / \sigma_{B}$, where $\lambda_{B}$ flips sign
- Low correlation region - $\rho<\sigma_{A} / \sigma_{B}$
- Both $\lambda_{A}$ and $\lambda_{B}$ are $>0$
- Info decreases (error increases) as $\rho$ increases
- High correlation region - $\rho>\sigma_{A} / \sigma_{B}$
- One coefficient $\lambda_{\mathrm{B}}$ is $<0$
- Info increases (error decreases) as $\rho$ increases
- Boundary (most conservative corr.) - $\rho=\sigma_{A} / \sigma_{B}$
- One coefficient $\lambda_{B}$ is $=0$
- Error is maximum $=\sigma_{A}$ (marginal info from $B$ is 0$)$
- Three information weights (A, B, correlation)
- Sum is equal to 1 by construction
- IW from $\rho$ is minimum (largest negative) for $\lambda_{B}=0$



## Two measurements - interpretation

- Low correlation regime, $\lambda_{B}>0$ (i.e. $\sigma_{A}{ }^{2}-\rho \sigma_{A} \sigma_{B}>0$ )
- The covariance can be seen as the sum of a common error and an uncorrelated error

- You can combine based on the uncorrelated error (compute statistical weights from basic error propagation) and add the common error only at the end
- Remember: this happens if the off-diagonal covariance $\rho \sigma_{A} \sigma_{B}$ is smaller than the smaller of the two variances $\sigma_{A}{ }^{2}$ (this will be used later in our "onionization" prescription proposal)
- Adding the less precise measurement B to the combination brings in additional information because B contributes independent (uncorrelated) knowledge about the unknown parameter
- High correlation regime, $\lambda_{B}<0$ (i.e. $\sigma_{A}{ }^{2}-\rho \sigma_{A} \sigma_{B}<0$ )
- The covariance can NOT be seen as the sum of a common error and an uncorrelated error
- Adding the less precise measurement $B$ to the combination brings in additional information because $B$ helps to constrain a correlated error on A to which it has a different sensitivity
- $\rho$ may be overestimated or this may be perfectly legitimate: you may leverage on the different sensitivity of two measurements to a common background or a common MC parameter to reduce the uncertainty on them
- The additional information only comes from the interplay of $B$ with $A$ - neither can claim it as their exclusive merit
- Boundary is $\lambda_{\mathrm{B}}=0-$ where $B$ brings no additional information
- Adding the less precise measurement B to the combination brings in NO additional information because its error has a part common with A plus an additional "uncorrelated" component


## Generalization to N measurements

- The previous discussion can be generalized from 2 to $\mathbf{N}$ measurements
- Fixing the variances on the measurements and varying only the correlations between them, any BLUE combination can fall in one of two regimes:
- Either Increasing any correlation decreases information and increases the combined error ("low" correlation regime because it includes the case where correlations are all 0)
- Or Increasing some correlations increases information and decreases the combined error ("high" correlation regime because you must increase a correlation beyond a threshold to enter it)
- It can be shown that these regimes are defined by the BLUE coefficients:
- Low correlations (info decreases if any correlation increases) $\leftrightarrow$ All BLUE coefficients are $\geq 0$
- High correlations (info increases if a correlation increases) $\leftrightarrow$ Some BLUE coefficients are <0
- In other words: if any BLUE coefficients are $<0$, it is because some correlations are high
- Marginal info from $\mathbf{N}^{\text {th }}$ measurement: fix all variances, vary $\mathbf{N}$-1 correlations ' $\mathbf{c}$ ' $\mathcal{M}_{N \times N}=\left(\begin{array}{cc}\mathcal{D}_{(N-1) \times(N-1)} & \mathbf{c} \\ \tilde{\mathbf{c}} & d\end{array}\right) \quad \Delta I_{N} \propto\left(\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)-1\right)^{2} \geq 0 \quad \lambda_{N} \propto-\left(\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)-1\right)$
- General case: fix variances, vary $\mathbf{N x}(\mathrm{N}-1) / 2$ correlations and use derivatives
- Information is minimum when some coefficients are 0
- If it is impossible to properly estimate correlations, the most "conservative" option is to exclude some measurements...

$$
\frac{\partial I}{\partial \rho_{i j}}=-2 I^{2} \lambda_{i} \lambda_{j} \sigma_{i} \sigma_{j}
$$

## Summary of part II

- If any BLUE coefficients are $<0$, it is because some correlations are high
- The BLUE combination is in a "high correlation" regime
- The contribution to information from measurements with BLUE coefficients <0 comes only from their interplay with other measurements through correlations
- They help constrain systematic uncertainties on other measurements
- If any BLUE coefficients are $<0$, the total BLUE error is smaller than that which would be obtained if some correlations were smaller
- Are correlations really as high or are they being overestimated?

IF any BLUE coefficients are <0
AND IF correlations were "conservatively" estimated to be 100\%, THEN correlation estimates should be reassessed

See part III for a few tools and suggestions

## Part III :

## "conservative" estimates of correlations

- The risks in overestimating correlations and a few hints to try and avoid this - with practical examples from $\mathrm{m}_{\text {top }}$ @ LHC
- " $100 \%$ " is not conservative when statistical errors are low
- Using information derivatives to identify the most relevant correlations
- Minimization and "onionization" procedures
- Splitting up the correlation contribution to information using marginal inflow
- Practical recommendations


## Is $\rho=100 \%$ "conservative"?

- Two measurements AB, each with a statistical and a systematic error
- Statistical errors are uncorrelated, systematic errors are correlated with a correlation $\rho_{\text {cor }}$
- If $\rho_{\text {cor }}$ is unknown, we want to look for the most "conservative" of $\rho_{\text {cor }}$ i.e. that which makes the combined BLUE error as large as possible
- The most conservative $\rho_{\text {cor }}$ is that which makes the total correlation $\rho$ as large as possible, as long as $\rho$ remains in the "region low-correlation" $\rho<\sigma_{A} / \sigma_{B}$ (assuming $\sigma_{A}<\sigma_{B}$ )
- If measurements are statistically dominated, the most conservative $\rho_{\text {cor }}$ is $\mathbf{1 0 0 \%}$
- But the most conservative $\rho_{\text {cor }}$ is $<\mathbf{1 0 0 \%}$ if systematic errors are not negligible


With N measurements, finding the most conservative correlations is not as easy!


## Prioritizing correlations to review

- So, I think I overestimated my correlations: and now, where do I start?


## - 1-Which correlations should I consider reassessing first?

- Use information derivatives to analyze which correlations (between which 2 measurements and for which error sources) have the largest effect
- Rescale the nominal covariance using a different rescaling factor (between 0 and 1 - do not increase correlations or flip their sign) for each off-diagonal element and error source

$$
\mathcal{M}_{i j}^{[s]} \rightarrow\left(\mathcal{M}^{\prime}\right)_{i j}^{[s]}=\left\{\begin{array}{ll}
f_{i j}^{[s]} \mathcal{M}_{i j}^{[s]} & \text { if } i \neq j, \\
\mathcal{M}_{i i}^{[s]} & \text { if } i=j,
\end{array} \quad 0 \leq f_{i j}^{[s]} \leq 1\right.
$$

- Will show in the following the derivatives (at nominal correlations) with respect to these scale factors

$$
\frac{\partial I^{\prime}}{\partial f_{i j}^{[s]}}=-2\left(I^{\prime}\right)^{2} \lambda_{i}^{\prime} \lambda_{j}^{\prime} \mathcal{M}_{i j}^{[s]}
$$

- Will normalize derivatives to total information at nominal correlations to make them a-dimensional
- For the top measurement at LHC there would be $21 \times 16$ derivatives
- 21 combinations (of 7 measurements) and 16 correlated error sources
- They should be studied individually, but in the following will show them summed by off-diagonal element (21) and by error source (16); the grand-total is always the derivative by a global scale factor


# LHC $m_{\text {top }}$ information derivatives 



- The first correlations that should be reviewed (if needed) are:
- Those between ATL10lj/ATL111j, ATL10lj/CMS11mj, CMS11mj/CMS11II
- Those for the bJES, dJES and RAD error sources
- Next slides: look at dJES and bJES for ATL101j/ATL111j and CMS11mj/CMS11II


## LHC $m_{\text {top }}$ systematics

| $\begin{aligned} & \text { N} \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & + \\ & \hline \end{aligned}$ | Uncertainty Categories |  |  | Size [GeV ] |  |  |  |  |  |  | Correlation |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Tevatron | ATLAS | CMS | ATLAS |  |  | CMS |  |  |  | $\rho_{\text {exp }} \rho_{\text {LHC }}$ |  |
|  |  |  |  | $\begin{gathered} 2010 \\ l+\text { jets } \end{gathered}$ | $\begin{gathered} 2011 \\ l+\text { jets } \end{gathered}$ | $2011$ <br> all jets | $\begin{gathered} 2010 \\ \text { di-l } \end{gathered}$ | $\begin{gathered} 2010 \\ l+\text { jets } \end{gathered}$ | $\begin{gathered} 2011 \\ \text { di-l } \end{gathered}$ | $\begin{gathered} 2011 \\ \mu+\text { jets } \end{gathered}$ |  |  |
|  | Statistics |  |  | 4.0 | 0.6 | 2.1 | 4.6 | 2.1 | 1.2 | 0.4 | 0 | 0 |
|  | iJES | Jet Scale Factor | Jet Scale Factor |  | 0.4 |  |  |  |  | 0.4 | 0 | 0 |
| ర | aJES |  |  |  |  |  |  |  |  |  |  |  |
| 응 | bJES | $J E S_{b-j e t}$ | $J E S_{b-j e t}$ | 2.5 | 1.6 | 1.4 | 0.9 | 0.9 | 1.1 | 0.7 | 1 | 0.5 |
| U | cJES |  |  |  |  |  |  |  |  |  |  |  |
| $\pm$ | dJES | $J E S_{\text {light-jet }}$ | $J E S_{\text {light-jet }}$ | 2.1 | 0.7 | 2.1 | 2.1 | 2.1 | 2.0 | 0.2 | 1 | 0 |
| ¢ | rJES |  | residual-JES |  |  |  | 3.3 |  |  |  | 0 | 0 |
| O | LepPt |  | Lepton $p_{\mathrm{T}}$ Scale |  |  |  | 0.3 |  | 0.2 |  | 1 | 0 |
| $\stackrel{C}{4}$ | MC | MC Generator Hadronisation | MC Generator | $\begin{aligned} & \hline 0.7 \\ & 0.7 \end{aligned}$ | $\begin{aligned} & \hline 0.3 \\ & 0.2 \end{aligned}$ | $\begin{aligned} & 0.5 \\ & (\cdot) \end{aligned}$ | 0.4 |  | 0.1 |  |  |  |
| - |  | Sum | Sum | 1.0 | 0.4 | 0.5 | 0.4 |  | 0.1 |  | 1 | 0.5 |
| $\begin{aligned} & 0 \\ & \frac{0}{0} \\ & \frac{0}{0} \end{aligned}$ | Rad | ISR/FSR | ISR/FSR Q-Scale Jet-Parton Scale | 2.5 | 1.0 | 1.7 | $\begin{aligned} & \hline 0.2 \\ & 0.6 \\ & 0.7 \\ & \hline \end{aligned}$ | $\begin{aligned} & 0.2 \\ & 1.1 \\ & 0.4 \end{aligned}$ | $\begin{aligned} & 0.4 \\ & 0.7 \end{aligned}$ | $\begin{aligned} & 0.8 \\ & 0.3 \end{aligned}$ |  |  |
| O |  | Sum | Sum | 2.5 | 1.0 | 1.7 | 0.9 | 1.2 | 0.8 | 0.8 | 1 | 0.5 |
| $\stackrel{+}{\square}$ | CR | Colour Recon. |  | 0.6 | 0.6 | 0.6 | 0.5 | 0.5 | 0.5 | 0.5 | 1 | 1 |
| 안 | PDF | Proton PDF | Proton PDF | 0.5 | 0.1 | 0.6 | 0.5 | 0.1 | 0.4 | 0.1 | 1 | 1 |
| \#1 |  |  | Jet Energy Res. $\begin{array}{r}b \text {-tagging } \\ E_{\mathrm{T}}^{\text {miss }} \\ \hline\end{array}$ | $\begin{aligned} & \hline 0.9 \\ & 0.5 \\ & 0.5 \end{aligned}$ | $\begin{gathered} 0.1 \\ <0.05 \\ 0.3 \\ 0.1 \\ \hline \end{gathered}$ | $\begin{aligned} & \hline 0.3 \\ & 0.2 \\ & 0.3 \end{aligned}$ | $\begin{aligned} & \hline 0.5 \\ & 0.4 \\ & 0.1 \\ & \hline \end{aligned}$ | $\begin{aligned} & 0.1 \\ & 0.1 \\ & 0.4 \\ & \hline \end{aligned}$ | $\begin{aligned} & \hline 0.3 \\ & \\ & 0.5 \\ & 0.4 \\ & \hline \end{aligned}$ | $\begin{aligned} & 0.2 \\ & \\ & 0.2 \\ & 0.1 \\ & \hline \end{aligned}$ |  |  |
|  | DetMod | Sum | Sum | 1.2 | 0.3 | 0.5 | 0.7 | 0.4 | 0.7 | 0.3 | 1 | 0 |
|  | UE | Underlying Event | Underlying Event | 0.6 | 0.6 | 0.6 | 1.4 | 0.2 | 0.6 | 0.6 | 1 | 0 |
|  |  | W+jet Norm. $W+$ jet Shape | background | $\begin{aligned} & 1.6 \\ & 0.8 \end{aligned}$ | 0.1 |  | 0.1 | 0.2 |  | 0.1 |  |  |
|  | BGMC | Sum | Sum | 1.8 | 0.1 |  | 0.1 | 0.2 |  | 0.1 | 1 | 1 |
|  | BGData | W+jet Norm. QCD Norm. QCD Shape Sum | QCD Norm. | $\begin{aligned} & 0.5 \\ & 0.4 \\ & 0.6 \end{aligned}$ | $\begin{aligned} & \hline 0.4 \\ & 0.2 \\ & 0.3 \\ & 0.5 \end{aligned}$ | $\begin{aligned} & 1.9 \\ & 1.9 \end{aligned}$ |  | $\begin{aligned} & 0.4 \\ & 0.4 \end{aligned}$ | $\begin{aligned} & 0.4 \\ & 0.4 \end{aligned}$ |  | 0 | 0 |
|  | Method | Method Calib. | Method Calib. | 0.4 | 0.1 | 1.0 | 0.3 | 0.1 | 0.4 | 0.2 | 0 | 0 |
|  | MHI | Pile-up | Pile-up | 0.7 | <0.05 |  | 1.0 | 0.1 | 0.2 | 0.4 | 1 | 1 |

Example:

- bJES and dJES are 100\% correlated between ATL11lj and ATL10lj, or CMS11mj and CMS11II, but have very different variances in the 2 measurements in each exp.
- The larger sensitivities of ATL10lj and CMS11Il to bJES and dJES with respect to ATL11lj and CMS11mj makes the former (overall less precise) measurements contribute information and reduce the combined error by constraining the dJES and bJES systematics in the latter (overall more precise) measurements


## LHC $m_{\text {top }}$ marginal combinations

|  | Comb. $m_{\text {top }}$ | Stat | iJES | a.JES | bJES | cJES | dJES | r.JES | Lept | MC | Rad | CR | PDF | DTMO | UE | BGMC | BGDT | Meth | MHI | $\chi^{2} /$ ndof |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CMS11 $\mu \mathrm{j}$ | $172.64 \pm 1.53$ | 0.37 | 0.43 | 0.00 | 0.66 | 0.00 | 0.23 | 0.00 | 0.00 | 0.00 | 0.80 | 0.54 | 0.05 | 0.28 | 0.59 | 0.09 | 0.00 | 0.15 | 0.38 |  |
| $\oplus$ ATL11lj | $173.02 \pm 1.46$ | 0.32 | 0.36 | 0.00 | 0.74 | 0.00 | 0.23 | 0.00 | 0.00 | 0.07 | 0.76 | 0.54 | 0.06 | 0.24 | 0.49 | 0.10 | 0.10 | 0.12 | 0.31 | 0.64/1 |
| $\oplus$ ATL10lj | $173.44 \pm 1.42$ | 0.44 | 0.36 | 0.00 | 0.71 | 0.00 | 0.18 | 0.00 | 0.00 | 0.04 | 0.69 | 0.54 | 0.03 | 0.22 | 0.48 | 0.02 | 0.16 | 0.12 | 0.25 | 2.06/2 |
| $\oplus$ CMS11ll | $173.40 \pm 1.41$ | 0.46 | 0.38 | 0.00 | 0.68 | 0.00 | 0.09 | 0.00 | 0.01 | 0.04 | 0.69 | 0.55 | 0.01 | 0.19 | 0.48 | 0.01 | 0.16 | 0.13 | 0.26 | 2.14/3 |
| $\oplus$ CMS10ll | $173.34 \pm 1.40$ | 0.47 | 0.38 | 0.00 | 0.68 | 0.00 | 0.08 | 0.06 | 0.01 | 0.04 | 0.69 | 0.55 | 0.01 | 0.19 | 0.46 | 0.01 | 0.16 | 0.13 | 0.25 | 2.40/4 |
| $\oplus$ ATL11aj | $173.33 \pm 1.40$ | 0.47 | 0.38 | 0.00 | 0.68 | 0.00 | 0.07 | 0.06 | 0.01 | 0.04 | 0.69 | 0.55 | 0.01 | 0.19 | 0.46 | 0.01 | 0.16 | 0.13 | 0.25 | 2.49/5 |
| $\oplus \mathrm{CMS10} 1 \mathrm{j}$ | $173.33 \pm 1.40$ | 0.47 | 0.38 | 0.00 | 0.68 | 0.00 | 0.07 | 0.06 | 0.01 | 0.04 | 0.69 | 0.55 | 0.01 | 0.19 | 0.47 | 0.01 | 0.16 | 0.13 | 0.25 | 2.49/6 |

Table 8: LHC $m_{\text {top }}$ BLUE combinations with the marginal addition of only one measurement at a time. The measurements are added in the following order: first those with $\lambda>0$, ordered by decreasing $\lambda$; then those with $\lambda<0$, ordered by decreasing $|\lambda|$.

- The large effect of bJES and dJES correlations can also be seen by adding measurements one by one (order by decreasing $\lambda>0$ and then decreasing $|\lambda|$ )
- The first 2 measurements ( $\lambda>0$ ) bring uncorrelated information (reduce statistical errors)
- The next $2(\lambda<0)$ reduce common systematics because of high correlations
- The third measurement $(\lambda=-7 \%)$ has a small effect on the combined error, but it reduces systematics (bJES, dJES, RAD...) and pushes the statistical error much higher! It also shifts the central value significantly. Is this reasonable or are correlations overestimated?
- The last 3 ( $\lambda<0,|\lambda|<2 \%)$ have virtually no effect on the combination
- Note that (in general) all measurements contribute to $\chi^{2}$ even if $\lambda \leq 0$
- It is not true that a measurement with $\lambda=0$ is effectively ignored
- Why does $X^{2}$ not change here for a growing number of degrees of freedom?


## Making correlations more conservative?

- So, I think I overestimated my correlations: and now, where do I start?
- 2 - Is there a way to (blindly) make correlations more conservative?
- We tested four procedures to reduce the nominal correlations
- Minimize information by varying one scale factor per error source (16 for $m_{\text {top }}$ )
- Minimize information by varying one scale factor per measurement pair (21 for $m_{\text {top }}$ )
- This may lead into non-physical space (non positive definite covariances) - give up in that case
- Minimize information by varying one global scale factor (1 for $m_{\text {top }}$ )
- "Onionize" covariance matrices as described in the next slide


## But carefully measuring correlations is much better than any rule of thumb!

## The "onionization" prescription

- Generalize to $\mathbf{N}$ measurements an observation made for two: keep each offdiagonal covariance $\rho_{i j} \sigma_{i} \sigma_{j}$ smaller than both variances $\sigma_{i}{ }^{2}$ and $\sigma_{i}{ }^{2}$
- We do this for each error source of uncertainty [s]

| $\rho_{\mathrm{cor}}^{[s]} \sigma_{A, \mathrm{cor}}^{[s]} \sigma_{B, \mathrm{cor}}^{[s]}$ | $\leq\left(\sigma_{A, \mathrm{cor}}^{[s]}\right)^{2}$ | $\forall s$ |
| :--- | :--- | :--- |
| $\rho_{\mathrm{cor}}^{[s]} \sigma_{A, \mathrm{cor}}^{[s]} \sigma_{B, \mathrm{cor}}^{[s]}$ | $\leq\left(\sigma_{B, \mathrm{cor}}^{[s]}\right)^{2}$ | $\forall s$ |


| $\left(\sigma_{A, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{A, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{A, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{A, \text { cor }}^{[s]}\right)^{2}$ |
| :---: | :---: | :---: | :---: |
| $\left(\sigma_{A, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{B, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{B, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{B, \text { cor }}^{[s]}\right)^{2}$ |
| $\left(\sigma_{A, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{B, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{C, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{C, \text { cor }}^{[s]}\right)^{2}$ |
| $\left(\sigma_{A, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{B, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{C, \text { cor }}^{[s]}\right)^{2}$ | $\left(\sigma_{D, \text { cor }}^{[s]}\right)^{2}$ |
| ( |  |  |  |
| $\sigma_{\text {A,cor }}$ | $\sigma_{\mathrm{B}, \mathrm{cor}}$ | $\sigma_{\mathrm{C}, \text { cor }}$ | $\leq \sigma_{\mathrm{D}, \mathrm{cor}}$ |

- This may lead to underestimated $\rho$ and is not even enough to ensure "conservativeness"
- But it may help avoid strange effects with highly correlated errors of very different sizes

PRELIMINARY! PRELIN
WORK IN PROGRESS!

- Example: covariance for the RAD uncertainty
- e.g. reduce the covariance for ATL10lj and ATL111j (very different variances: 2.5 and 1.0)

| CMS11ll | ( 0.59 | 0.61 | 0.73 | 0.39 | 0.91 | 0.65 | 0.94 | CMS111l | ( 0.59 | 0.59 |  | 0.59 | 0.59 | 0.59 | 0.59 | CMS11ll | ( 0.59 | 0.59 | 0.59 |  |  | 0.59 | 0.59 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CMS11 $\mu \mathrm{j}$ | 0.61 | 0.64 | 0.75 | 0.40 | 0.95 | 0.68 | 0.98 | CMS11 $\mu \mathrm{j}$ | 0.59 | 0.64 | 0.64 | 0.64 | 0.64 | 0.64 | 0.64 | CMS11 $\mu$ j | 0.59 | 0.64 | 0.64 | 0.40 | 0.64 | 0.64 | 0.64 |
| CMS10ll | 0.73 | 0.75 | 0.89 | 0.48 | 1.12 | 0.80 | 1.15 | CMS10ll | 0.59 | 0.64 | 0.89 | 0.89 | 0.89 | 0.89 | 0.89 | CMS10ll | 0.59 | 0.64 | 0.89 | 0.48 | 0.89 | 0.80 | 0.89 |
| ATL11lj | 0.39 | 0.40 | 0.48 | 1.02 | 0.60 | 1.72 | 2.47 | ATL11lj | 0.59 | 0.64 | 0.89 | 1.02 | 1.02 | 1.02 | 1.02 | ATL11lj | 0.39 | 0.40 | 0.48 | 1.02 | 0.60 | 1.02 | 1.02 |
| CMS10lj | 0.91 | 0.95 | 1.12 | 0.60 | 1.41 | 1.01 | 1.45 | CMS10lj | 0.59 | 0.64 | 0.89 | 1.02 |  | 1.41 | 1.41 | CMS10lj | 0.59 | 0.64 | 0.89 | 0.60 | 1.41 | 1.01 | 1.41 |
| ATL11aj | 0.65 | 0.68 | 0.80 | 1.72 | 1.01 | 2.89 | 4.16 | ATL11aj | 0.59 | 0.64 | 0.89 | 1.02 | 1.41 |  | 2.89 | ATL11aj | 0.59 | 0.64 | 0.80 | 1.02 | 1.01 | 2.89 | 2.89 |
| ATL10lj | ( 0.94 | 0.98 | 1.15 | 2.47 | 1.45 | 4.16 | 5.98 | ATL10lj |  | 0.64 | 0.89 | 1.02 | 1.41 | 2.89 | 5.98 | ATL10lj | ( 0.59 | 0.64 | 0.89 | 1.02 | 1.41 | 2.89 | 5.98 |
| Table 13: Onionization of the RAD covariance. Nominal RAD. |  |  |  |  |  |  |  | Table 14: Onionization of the RAD covariance. Onion upper bound. |  |  |  |  |  |  |  | Table 15: | Onionizatio | on of the | he RAD | covarian | nce. Or | nionize | RAD. |
| OLD |  |  |  |  |  |  |  | MAX |  |  |  |  |  |  |  |  |  |  |  | NE |  |  |  |
| BLUE@ TOPLHCWG - 29th November 2012 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  | A. Valassi, R. Chierici - 24 |  |  |  |  |  |  |  |

# LHC $m_{\text {top }}$ - conservative correlations? ${ }^{\text {CERN } T T_{\text {Dep }}}$ 

|  | Comb. $m_{\text {top }}$ | Stat | iJES | bJES | dJES | rJES | Lept | MC | Rad | CR | PDF | DTMO | UE | BGMC | BGDT | Meth | MHI | $\chi^{2} / \mathrm{ndof}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| JUNE 2012 TOP WG | $173.33 \pm 1.40$ | 0.47 | 0.38 | 0.68 | 0.07 | 0.06 | 0.01 | 0.04 | 0.69 | 0.55 | 0.01 | 0.19 | 0.47 | 0.01 | 0.16 | 0.13 | 0.25 | 2.49/6 |
| MinimizeGlobalFactor1 | $173.29 \pm 1.41$ | 0.40 | 0.36 | 0.71 | 0.19 | 0.05 | 0.01 | 0.06 | 0.71 | 0.53 | 0.04 | 0.22 | 0.47 | 0.04 | 0.14 | 0.12 | 0.26 | 2.11/6 |
| MinimizeErrorSrc18 | $173.22 \pm 1.44$ | 0.35 | 0.35 | 0.72 | 0.25 | 0.02 | 0.00 | 0.09 | 0.73 | 0.54 | 0.05 | 0.24 | 0.48 | 0.10 | 0.12 | 0.12 | 0.29 | 1.97/6 |
| Onionize | $173.08 \pm 1.44$ | 0.31 | 0.33 | 0.72 | 0.26 | 0.01 | 0.00 | 0.07 | 0.76 | 0.54 | 0.06 | 0.24 | 0.47 | 0.09 | 0.10 | 0.11 | 0.28 | 1.67/6 |
| MinimizeOffDiagonal21 | $173.02 \pm 1.46$ | 0.32 | 0.36 | 0.74 | 0.23 | 0.00 | 0.00 | 0.07 | 0.76 | 0.54 | 0.06 | 0.24 | 0.49 | 0.10 | 0.10 | 0.12 | 0.31 | 1.83/6 |
| CMS11 $\mu \mathrm{j} \oplus$ ATL11lj | $173.02 \pm 1.46$ | 0.32 | 0.36 | 0.74 | 0.23 | 0.00 | 0.00 | 0.07 | 0.76 | 0.54 | 0.06 | 0.24 | 0.49 | 0.10 | 0.10 | 0.12 | 0.31 | 0.64/1 |
| NO CORRELATIONS | $173.25 \pm 1.05$ | 0.39 | 0.22 | 0.48 | 0.39 | 0.09 | 0.02 | 0.08 | 0.47 | 0.29 | 0.07 | 0.19 | 0.32 | 0.07 | 0.16 | 0.11 | 0.18 | 1.09/6 |

- Largest combined error (most "conservative" result) is found by minimizing information when rescaling correlations independently for the 21 pairs
- Result is identical (but for the $x^{2}$ ) with the combination of CMS11 $\mu \mathrm{j}$ and ATL11jI alone
- Checked that 20 information derivatives are 0 (for all covariances but that of CMS11 $\mu \mathrm{j}$ and ATL11jl)
- Two BLUE coefficients for CMS11 $\mu \mathrm{j}$ and ATL11jl are $>0$; five others are 0 (remember dl/dM $\mathrm{M}_{\mathrm{ij}} \sim-\lambda_{i} \lambda_{\mathrm{j}}$ )
- The most conservative option consists in not using the results with $\lambda<0$ !
- Some negative BLUE coefficients remain for the other two minimizations and (fewer) for the onionization procedure
- All of these procedures change the sys/stat balance, not only the total error!
- Nominal correlations: lower total error and systematics, higher statistical error
- Modified correlations: higher total error and systematics, lower statistical error


## Recommendations?

- Assume you have chosen your best estimate of correlations ("nominal")
- If all $\lambda>0$ at "nominal" correlations, you are in a low-correlation regime
- Even if correlations are high, their estimates are "conservative" enough: they do not increase the available information, they decrease it - nothing to worry about! :)
- If some $\boldsymbol{\lambda}$ are negative, you are in a high-correlation regime
- You should be aware that these high correlations are adding information!
- Effectively, the measurements with $\lambda<0$ are only adding information through these correlations
- We advise that all correlations should then be reviewed on a case by case basis
- There are legitimate cases where high correlations exist and may help you constrain common systematics (e.g. you varied same MC parameter or depend on a common $E_{\text {beam }}$ measurement)
- But if you "conservatively" estimated correlations to be $100 \%$, then you are definitely wrong!
- If correlations are overestimated, you are certainly underestimating the combined systematics, even more than you are underestimating the total combined error - and your central values are shifted, too
- What can you do to provide a more realistic (or "conservative") estimate?
- Prioritize the correlations to analyze, using information derivatives
- Then measure these correlations with your data and MC if this is possible
- For correlations you cannot measure precisely, consider some of the tools we presented
- You may try to minimize information in those correlations alone, or try the onionization procedure (if $\sigma_{S}$ for $A$ and $B$ is correlated, and $\sigma_{S, B}>\sigma_{S, A}$, strip out $\sigma_{S^{\prime}, B}=\sigma^{2}{ }_{S, B} \sigma^{2}{ }_{\mathrm{S}, \mathrm{A}}$ as uncorrelated)
- If you have no clue, omitting the measurements with $\lambda<0$ is the most conservative option


## Summary:

"relative importance" revisited and conclusions

## LHC $m_{\text {top }}$ marginal information weights

|  | $m_{\text {top }}$ <br> $[\mathrm{GeV}]$ | $\Delta I$ <br> $\left[1 / \mathrm{GeV}^{2}\right]$ | $\lambda$ | $\frac{\|\lambda\|}{\sum\|\lambda\|}$ | $\frac{\Delta I}{I}$ |
| :--- | :---: | ---: | :---: | :---: | :---: |
| CMS11 $\mu \mathrm{j}$ | $172.64 \pm 1.53$ | 0.425 | $84.3 \%$ | $65.7 \%$ | $83.8 \%$ |
| ATL11lj | $174.53 \pm 2.39$ | 0.175 | $29.9 \%$ | $23.3 \%$ | $34.5 \%$ |
| ATL10lj | $169.33 \pm 6.32$ | 0.025 | $-6.8 \%$ | $5.3 \%$ | $4.9 \%$ |
| CMS111ll | $173.30 \pm 2.96$ | 0.114 | $-4.8 \%$ | $3.7 \%$ | $22.5 \%$ |
| CMS10ll | $175.50 \pm 6.49$ | 0.024 | $-1.9 \%$ | $1.5 \%$ | $4.7 \%$ |
| ATL11aj | $174.90 \pm 4.44$ | 0.051 | $-0.4 \%$ | $0.3 \%$ | $10.0 \%$ |
| CMS10lj | $173.10 \pm 3.41$ | 0.086 | $-0.2 \%$ | $0.2 \%$ | $16.9 \%$ |
| Correlations | - | -0.392 | - | - | $-77.4 \%$ |
| Total | $173.33 \pm 1.40$ | 0.507 | $100.0 \%$ | $100.0 \%$ | $100.0 \%$ |

## "Split \#0": information weights (IW) from Part I of this talk

- One IW per measurement and one for all correlations together
"Split \#1": one possible way to split up IW for correlations
- Split it up completely amongst all measurements
- Use "marginal information inflow"

|  |  |  |  | $\mathrm{I}_{\text {tot }}=0.507$ |  |  | plit \#1 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $I_{\text {corr }}=I-I_{\mathrm{A}}-I_{\mathrm{B}}=\frac{1}{\sigma_{\hat{Y}}^{2}}-\frac{1}{\sigma_{A}^{2}}-\frac{1}{\sigma_{B}^{2}}$ |  | $\begin{gathered} m_{\mathrm{top}} \\ {[\mathrm{GeV}]} \end{gathered}$ | $\begin{gathered} \Delta I \\ {\left[1 / \mathrm{GeV}^{2}\right]} \end{gathered}$ | $\lambda$ | $\frac{\|\lambda\|}{\sum\|\lambda\|}$ | $\Delta I$ $I$ | $\sum \frac{\Delta I}{I}$ |
|  | CMS11 $\mu \mathrm{j}$ | $172.64 \pm 1.53$ | 0.425 | 84.3\% | 65.7\% | 83.8\% | 83.8\% |
|  | Low sorr. | - | - | - | - |  |  |
| Split low and high correlations by rescaling c as$\mathbf{c}^{\prime}=\frac{1}{\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)} \mathbf{c}$ | ATL11lj | $174.53 \pm 2.39$ | 0.175 | 29.9\% | 23.3\% | 34.5\% | \% |
|  | Low corr. | - | -0.131 | - | - | -25.9\% | 8.6\% |
|  | ATL10lj | $169.33 \pm 6.32$ | 0.025 | -6.8\% | 5.3\% | 4.9\% |  |
|  | ow corr | - | -0.025 | - | - | -4.9\% | 5.8\% |
|  | High corr. | - | 0.030 | - | - | 5.8\% |  |
| For adding ATL10lj, this factor is 0.543 | CMSIIIl | $173.30 \pm 2.96$ | 0.114 | -4.8\% | 3.7\% | 22.5\% |  |
|  | Low corr. | - | -0.114 | - | - | -22.5\% | 1.1\% |
|  | High corr. | - | 0.006 | - | - | 1.1\% |  |
|  | CMS10ll | $175.50 \pm 6.49$ | 0.024 | -1.9\% | 1.5\% | 4.7\% |  |
|  | Low corr. | - | -0.024 | - | - | -4.7\% | 0.7\% |
|  | High corr. | - | 0.003 | - | - | 0.7\% |  |
|  | ATL11aj | $174.90 \pm 4.44$ | 0.051 | -0.4\% | 0.3\% | 10.0\% |  |
| Part I of this talk relations together | Low corr. | - | -0.051 | - | - | -10.0\% | 0.0\% |
|  | High corr. | - | 0.000 | - | - | 0.0\% |  |
|  | CMS10lj | $173.10 \pm 3.41$ | 0.086 | -0.2\% | 0.2\% | 16.9\% |  |
| W for correlations | Low corr. | - | -0.086 | - | - | -16.9\% | 0.0\% |
|  | High corr. | - | 0.000 | - | - | 0.0\% |  |
| - | Total | $173.33 \pm 1.40$ | 0.507 | 100.0\% | 100.0\% | 100.0\% | 100.0\% |

- This requires a ranking by BLUE coefficient
- Two measurements contribute $0.0 \%$ information if they are added last

More complex ways to split this up may be more satisfactory

- It would be nice to treat differently negative/positive information contributions in low/high correlation regimes


## Conclusions

- We propose to use information contributions (instead of absolute BLUE coefficients) to assess the "relative importance" of measurements
- One contribution to information comes from the correlations between the measurements and cannot generally be attributed to any one of them individually
- Marginal information inflow is a possible alternative but implies a ranking
- Negative BLUE coefficients indicate a "high correlation" regime
- It is crucial to properly assess correlations in this regime: overestimated correlations may lead to largely underestimated combined systematics
- We propose a few tools to help a critical review of correlation estimates
- Derivatives of information may help to prioritize the most sensitive correlations
- Information minimization and covariance onionization may help be more conservative
- Understanding the marginal inflow of information from the addition of one measurement to the combination may help decide what to do with it
- Of course the ideal solution is to measure correlations in the data and MC

High correlations have a large effect on the combination: as much effort should go into properly estimating correlations than individual errors

## Some useful references

- A. C. Aitken, On Least Squares and Linear Combinations of Observations, Proc. Roy. Soc. Edinburgh 55 (1935), 42
- The first published description of the BLUE technique (AFAIK)
- L. Lyons, D. Gibaut, P. Clifford, How to combine correlated estimates of a single physical quantity, Nucl. Instr. Meth. A270 (1988) 110
- Extensive discussion of correlations and negative weights for two measurements of a single observable! A very useful read! ()
- A. Valassi, Combining correlated measurements of several different physical quantities, Nucl. Instr. Meth. A500 (2003) 391
- Generalization of Lyons formulas and computation of individual error contributions for many observables (e.g. LEPEWWG)
- F. James, Statistical Methods in Experimental Physics (2 ${ }^{\text {nd }}$ Edition), World Scientific (2006)
- A very useful textbook covering information and estimation
- A. van den Bos, Parameter Estimation for Scientists and Engineers, WileyInterscience (2007)
- Another very useful textbook covering information and estimation


## RESERVE SLIDES

## Negative weights (and absoulte weighs) @Tevatron



http://arxiv.org/pdf/1107.5255.pdf<br>FERMILAB-TM-2504-E<br>TEVEWWG/top 2011/xx<br>CDF Note 10549<br>D $\varnothing$ Note 6222<br>July 2011

|  | Run I published |  |  |  |  | Run II published |  |  |  |  | Run II preliminary CDF |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | CDF |  |  | Dø |  | CDF |  |  | Dø |  |  |  |
|  | $1+\mathrm{jt}$ | di-1 | allh | 1+jt | di-1 | $1+\mathrm{jt}$ | di-1 | Lxy | $1+\mathrm{jt}$ | di-1 | allh | Met |
| Pull | $+0.40$ | -0.51 | +1.12 | +1.33 | -0.37 | -0.23 | -0.81 | -0.67 | 1.52 | 0.27 | +0.40 | -0.36 |
| Weight [\%] | -4.7 | -1.0 | -0.8 | -0.0 | -0.2 | +56.6 | +1.4 | +0.3 | +27.2 | +1.5 | +14.0 | +6.7 |

The weights of some of the measurements are negative. In general, this situation can occur if the correlation between two measurements is larger than the ratio of their total uncertainties. This is indeed the case here. In these instances the less precise measurement will usually acquire a negative weight. While a weight of zero means that a particular input is effectively ignored in the combination, a negative weight means that it affects the resulting $M_{\mathrm{t}}$ central value and helps reduce the total uncertainty. To visualize the weight each measurement carries in the combination, Fig. 2 shows the absolute values of the weight of each measurement divided by the sum of the absolute values of the weights of all input measurements.

## Information

- General definition for many observables

$$
\begin{aligned}
\mathcal{I}_{\alpha \beta}^{(\mathbf{X})} & =E\left[\frac{\partial \log p(\mathbf{y} ; \mathbf{X})}{\partial X_{\alpha}} \frac{\partial \log p(\mathbf{y} ; \mathbf{X})}{\partial X_{\beta}}\right] \\
& =\int \frac{\partial \log p(\mathbf{y} ; \mathbf{X})}{\partial X_{\alpha}} \frac{\partial \log p(\mathbf{y} ; \mathbf{X})}{\partial X_{\beta}} p(\mathbf{y} ; \mathbf{X}) d y_{1} \ldots d y_{n}
\end{aligned}
$$

- For multivariate Gaussians (as in BLUE assumption) everything becomes much easier

$$
p(\mathbf{y} ; \mathbf{X})=\frac{1}{(2 \pi)^{n / 2}(\operatorname{det} \mathcal{M})^{n / 2}} \exp \left(-\frac{1}{2} \sum_{i=1}^{n} \sum_{j=1}^{n}(y-\mathcal{U} X)_{i} \mathcal{M}_{i j}^{-1}(y-\mathcal{U} X)_{j}\right)
$$

- Information becomes

$$
\mathcal{I}_{\alpha \beta}^{(\mathrm{X})}=\left(\tilde{\mathcal{U}} \mathcal{M}^{-1} \mathcal{U}\right)_{\alpha \beta}
$$

related to the BLUE covariance $\operatorname{cov}\left(\hat{x}_{\alpha}, \hat{x}_{\beta}\right)=\sum_{i=1}^{n} \sum_{j=1}^{n} \lambda_{\alpha i} \mathcal{M}_{i j} \lambda_{\beta j}=\left(\tilde{\mathcal{U}} \mathcal{M}^{-1} \mathcal{U}\right)_{\alpha \beta}^{-1}$

$$
\text { by } \operatorname{cov}\left(\hat{x}_{\alpha}, \hat{x}_{\beta}\right)=\left(\mathcal{I}^{(\mathbf{X})}\right)_{\alpha \beta}^{-1}
$$

i.e. BLUE is efficient (satisfies the Cramer Rao lower bound)

## What if correlations are fixed a priori? ${ }^{\text {CERNT }}$

Fixing $\sigma_{A}$ and $\rho$ and varying $\sigma_{B}$, the result is somewhat surprising, but completely in line with the previous comments:


- Once in the high-correlation regime ( $\left.\sigma_{\mathrm{B}}>\sigma_{\mathrm{A}} / \mathrm{\rho}\right)$, the higher is $\sigma_{\mathrm{B}}$, the higher the total information and the lower the combined error!
- The better you measure $B$ (as long as $\left.\sigma_{B}>\sigma_{A} / \rho\right)$, the worse the combination!
- The derivative of total information with respect to $\left(1 / \sigma_{B}\right)^{2}$ is negative - it is actually $\lambda_{B}$ !
- We also considered using integrals of the dl/df information derivatives to split up the information weight of correlations
- The differential of I for a step in the multi dimensional space of rescaling factors can be easily written as the sum of contributions from each off-diagonal element in each error source, individually
- Integrate this by defining a "path" to transform 0 correlations (all scale factors $=0$ ) to nominal correlations (all scale factors $=1$ )
- Use two segments of straight lines, from 0 to a minimum (e.g. minimum by offdiagonal element) and then up to nominal values, to further split up two separate contributions in the low-correlation and high-correlation regimes
- Some interesting results but in the end why bother?
- Derivatives alone provide useful tool to prioritize correlations to re-analyze
- Marginal information analysis seems to be a more natural and better fit to split up "contributions to knowledge" if really needed
- series converges quite fast in high-correlation regime because we start from minimum (only $\lambda>0$ ) and we then add measurements in a carefully chosen order


## $\lambda=0$ - information vs. $\mathrm{X}^{2}$

- A measurement with a zero weight $\lambda_{i}=0$ in a BLUE combination does not contribute any information...
- i.e. it does not reduce the variance of the combined estimate
- ... but note that it does contribute to the $\chi^{2}$ calculation
- the $x^{2}$ is essentially a weighted sum of the products of the differences to the estimate, with weights $\left(\mathrm{M}^{-1}\right)_{\mathrm{ij}}$ given by the inverse of the covariance matrix
- the three sets of "weights", BLUE coefficients, information weights and $\mathrm{x}^{2}$ weights are all equal to $1 / \sigma_{\mathrm{i}}^{2}$ if there are no correlations
- Useful to assess if correlations are correctly estimated
- similarly, measurements with $\lambda_{i}<0$ also contribute to the $\chi^{2}$


## The problem with absolute values

- Take A, B, C, with C uncorrelated to A or B:
- Combining A, B, C gives the same result as first combining $A, B$ and then adding $C$ :

$$
\left(\lambda_{A}\right)_{A B C}=\frac{\left(\lambda_{A}\right)_{A B} /\left(\sigma_{\hat{Y}}^{2}\right)_{A B}}{1 /\left(\sigma_{\hat{Y}}^{2}\right)_{A B}+1 / \sigma_{C}^{2}}
$$

$$
\frac{1}{\left(\sigma_{\hat{Y}}^{2}\right)_{A B C}}=\frac{1}{\left(\sigma_{\hat{Y}}^{2}\right)_{A B}}+\frac{1}{\sigma_{C}^{2}}
$$

$$
\begin{aligned}
\left(\lambda_{B}\right)_{A B C} & =\frac{\left(\lambda_{B}\right)_{A B} /\left(\sigma_{\hat{Y}}^{2}\right)_{A B}}{1 /\left(\sigma_{\hat{Y}}^{2}\right)_{A B}+1 / \sigma_{C}^{2}} \\
\left(\lambda_{C}\right)_{A B C} & =\frac{1 / \sigma_{C}^{2}}{1 /\left(\sigma_{\hat{Y}}^{2}\right)_{A B}+1 / \sigma_{C}^{2}}
\end{aligned}
$$

- Intuitively, the relative contribution of C is its "information weight"

$$
\left(\mathrm{IW}_{C}\right)_{A, B, C}=\frac{I_{C}}{I_{C}+I_{A B}}=\frac{1 / \sigma_{C}^{2}}{1 /\left(\sigma_{\hat{Y}}^{2}\right)_{A B}+1 / \sigma_{C}^{2}}=\left(\lambda_{C}\right)_{A B C}=\frac{\left(\lambda_{C}\right)_{A B C}}{\left(\lambda_{A}\right)_{A B C}+\left(\lambda_{B}\right)_{A B C}+\left(\lambda_{C}\right)_{A B C}}
$$

while the "relative importance" based on absolute values is

$$
>1 \quad\left(\mathrm{RI}_{C}\right)_{A, B, C}=\frac{\left|\left(\lambda_{C}\right)_{A B C}\right|}{\left|\left(\lambda_{A}\right)_{A B C}\right|+\left|\left(\lambda_{B}\right)_{A B C}\right|+\mid\left(\lambda_{C}\right)_{A B C}}
$$

- Two issues if $\rho_{A B}>\sigma_{A} / \sigma_{B}$ : first, RI underestimates the contribution of $C$; second, RI gives different results for ( $A, B, C$ ) and for ( $A B, C$ )!

$$
\left(\mathrm{RI}_{C}\right)_{A, B, C}<\left(\mathrm{RI}_{C}\right)_{A B, C}=\left(\mathrm{IW}_{C}\right)_{A B, C}=\left(\mathrm{IW}_{C}\right)_{A, B, C} \quad \text { if } \rho>\sigma_{A} / \sigma_{B}
$$

## Two measurements of 1 observable

- For two measurements $A, B$ with errors $\sigma_{A}<\sigma_{B}$ and correlation $\rho$ :

$$
\begin{aligned}
\hat{Y}=\lambda_{A} y_{A}+\lambda_{B} y_{B} & \lambda_{A} & =\frac{\sigma_{B}^{2}-\rho \sigma_{A} \sigma_{B}}{\sigma_{A}^{2}+\sigma_{B}^{2}-2 \rho \sigma_{A} \sigma_{B}} \\
\sigma_{\hat{Y}}^{2}=\frac{\sigma_{A}^{2} \sigma_{B}^{2}\left(1-\rho^{2}\right)}{\sigma_{A}^{2}+\sigma_{B}^{2}-2 \rho \sigma_{A} \sigma_{B}}=\frac{1}{I} & \lambda_{B} & =\frac{\sigma_{A}^{2}-\rho \sigma_{A} \sigma_{B}}{\sigma_{A}^{2}+\sigma_{B}^{2}-2 \rho \sigma_{A} \sigma_{B}}
\end{aligned}
$$

- The effect of correlations was extensively discussed by Lyons et al.
- Fixing $\sigma_{A}$ and $\sigma_{B}$ and varying $\rho$, the combined error has a maximum (equal to $\sigma_{A}$ ) for $\rho=\sigma_{A} / \sigma_{B}$, which is where $\lambda_{B}$ flips sign from $>0$ to $<0$



## IW's in two measurements of 1 observable

## By construction the sum of all information weights ( N meas. +1 correlation) is $=1$



Most "conservative" correlation: largest combined error, minimum (largest negative) ínformation weight from correlations

## Two measurements - interpretation

- Low correlation regime, $\lambda_{B}>0$ (i.e. $\sigma_{A}{ }^{2}-\rho \sigma_{A} \sigma_{B}>0$ )
$>0$ (positive definite matrix)
- The covariance can be seen as the sum of a common error and an uncorrelated error $\rightarrow$
$\left(\begin{array}{cc}\sigma_{A}^{2} & \rho \sigma_{A} \sigma_{B} \\ \rho \sigma_{A} \sigma_{B} & \sigma_{B}^{2}\end{array}\right)=\left(\begin{array}{cc}\rho \sigma_{A} \sigma_{B} & \rho \sigma_{A} \sigma_{B} \\ \rho \sigma_{A} \sigma_{B} & \rho \sigma_{A} \sigma_{B}\end{array}\right)_{\text {com }}+\left(\begin{array}{cc}\sigma_{A}^{2}-\rho \sigma_{A} \sigma_{B} & 0 \\ 0 & \sigma_{B}^{2}-\rho \sigma_{A} \sigma_{B}\end{array}\right)_{\text {unc }}$
- You can combine based on the uncorrelated error (compute statistical weights from basic error propagation) and add the common error only at the end
- Remember: this happens if the off-diagonal covariance $\rho \sigma_{A} \sigma_{B}$ is smaller than the smaller of the two variances $\sigma_{A}{ }^{2}$ (this will be used later in our "onionization" prescription proposal)
- Adding the less precise measurement B to the combination brings in additional information because B contributes independent (uncorrelated) knowledge about the unknown parameter
- High correlation regime, $\lambda_{B}<0$ (i.e. $\sigma_{A}{ }^{2}-\rho \sigma_{A} \sigma_{B}<0$ )
- The covariance can NOT be seen as the sum of a common error and an uncorrelated error
- Subtracting a common error matrix $\rho \sigma_{A} \sigma_{B}$ from the covariance would give a non-positive-definite matrix
- The covariance can instead be seen as the sum of a common error and of a $100 \%$ positively correlated effect
- Adding the less precise measurement B to the combination brings in additional information because B helps to constrain a correlated error on A to which it has a different sensitivity
- There are cases where this is perfectly legitimate: you may leverage on the different sensitivity of two measurements to a common background or a common MC parameter to reduce the uncertainty on them
- But this may effectively (and incorrectly) happen also if correlations are being overestimated
- Boundary is $\lambda_{B}=0$ - where $B$ brings no additional information
- Adding the less precise measurement $B$ to the combination brings in NO additional information because its error has a part common with A plus an additional "uncorrelated" component


## Two measurements - interpretation

- Low correlation regime, $\lambda_{B}>0$ (i.e. $\sigma_{A}{ }^{2}-\rho \sigma_{A} \sigma_{B}>0$ )
$\geq 0$ (positive definite matrix)
- The covariance can be seen as the sum of a common error and an uncorrelated error $\rightarrow$

$$
\left(\begin{array}{cc}
\sigma_{A}^{2} & \rho \sigma_{A} \sigma_{B} \\
\rho \sigma_{A} \sigma_{B} & \sigma_{B}^{2}
\end{array}\right)=\left(\begin{array}{cc}
\rho \sigma_{A} \sigma_{B} & \rho \sigma_{A} \sigma_{B} \\
\rho \sigma_{A} \sigma_{B} & \rho \sigma_{A} \sigma_{B}
\end{array}\right)_{\mathrm{com}}+\left(\begin{array}{cc}
\sigma_{A}^{2}-\rho \sigma_{A} \sigma_{B} & 0 \\
0 & \sigma_{B}^{2}-\rho \sigma_{A} \sigma_{B}
\end{array}\right)_{\mathrm{unc}}
$$

- You can combine based on the uncorrelated error (compute statistical weights from basic error propagation) and add the common error only at the end
- Remember: this happens if the off-diagonal covariance $\rho \sigma_{A} \sigma_{B}$ is smaller than the smaller of the two variances $\sigma_{A}{ }^{2}$ (this will be used later in our "onionization" prescription proposal)
- Adding the less precise measurement B to the combination brings in additional information because $B$ contributes independent (uncorrelated) knowledge about the unknown parameter
- High correlation regime, $\boldsymbol{\lambda}_{\mathbf{B}}<\mathbf{0}$ (i.e. $\left.\boldsymbol{\sigma}_{\mathbf{A}}{ }^{2}-\boldsymbol{\rho} \boldsymbol{\sigma}_{\mathbf{A}} \boldsymbol{\sigma}_{\mathbf{B}}<\mathbf{0}\right) \quad \begin{aligned} & y_{A}=Y_{\text {true }}+\delta_{\text {com }}-\lambda_{B} \delta_{\text {cor }} \\ & y_{B}=Y_{\text {true }}+\delta_{\text {com }}+\lambda_{A} \delta_{\text {cor }}\end{aligned}$
- The covariance can NOT be seen as the sum of a common error and an uncorrelated error
- Subtracting a common error matrix $\rho \sigma_{A} \sigma_{\mathrm{B}}$ from the covariance would give a non-positive-definite matrix
- The covariance can be seen as the sum of a common error and of a $100 \%$ positively correlated effect $\left(\begin{array}{cc}\sigma_{A}^{2} & \rho \sigma_{A} \sigma_{B} \\ \rho \sigma_{A} \sigma_{B} & \sigma_{B}^{2}\end{array}\right)=\left(\begin{array}{cc}\sigma_{\hat{Y}}^{2} & \sigma_{\hat{Y}}^{2} \\ \sigma_{\hat{Y}}^{2} & \sigma_{\hat{Y}}^{2}\end{array}\right)_{\mathrm{com}}+\left(\sigma_{A}^{2}+\sigma_{B}^{2}-2 \rho \sigma_{A} \sigma_{B}\right)\left(\begin{array}{cc}\lambda_{B}^{2} & -\lambda_{A} \lambda_{B} \\ -\lambda_{A} \lambda_{B} & \lambda_{A}^{2}\end{array}\right)_{\text {cor }} \ggg 0$
- Adding the less precise measurement $\mathbf{B}$ to the combination brings in additional information because $B$ helps to constrain a correlated error on $A$ to which it has a different sensitivity
- Boundary is $\boldsymbol{\lambda}_{\mathrm{B}}=\mathbf{0}$ - where B brings no additional information
- Adding the less precise measurement $B$ to the combination brings in NO additional information because its error has a part common with A plus an additional "uncorrelated" component


## Marginal information inflow

- How much info does the $\mathbf{N}^{\text {th }}$ measurement add to a combination with $\mathrm{N}-1$ ?
- Call $M$ the $N_{x} N$ covariance and $D$ the ( $\mathrm{N}-1$ ) $\times(\mathrm{N}-1)$ covariance
- Call c the ( $\mathrm{N}-1$ ) dimensional vector of covariances of the $\mathrm{N}^{\text {th }}$ measurement $\quad \mathcal{M}=\left(\begin{array}{cc}\mathcal{D} & \mathbf{c} \\ \tilde{\mathbf{c}} & d\end{array}\right)$

$$
\mathcal{M}=\left(\begin{array}{cc}
\mathcal{D} & \mathbf{c} \\
\tilde{\mathbf{c}} & d
\end{array}\right)
$$

- Define $u$ as an ( $\mathrm{N}-1$ ) dimensional vector of 1 's - contraction by this vector is a sum on $\mathrm{N}-1$ indices
- The marginal information from the $\mathbf{N}^{\text {th }}$ measurement is always $\geq \mathbf{0}$ [van den Bos]
- The information inflow $\Delta l$ from the $\mathrm{N}^{\text {th }}$ measurement is minimum for (i.e. the $\mathrm{N}^{\text {th }}$ measurement brings no marginal information)

$$
\Delta I=0 \Longleftrightarrow\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)-1=0
$$

- This is a hyperplane in the $N$ - 1 dimensional space of correlations of the $N^{\text {th }}$ to all other measurements
- But the BLUE coefficient for the $\mathbf{N}^{\text {th }}$ measurement is $\lambda_{N} \propto-\left(\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)-1\right)$
- In other words, $\Delta I=0$ implies $\lambda_{N}=0$ and viceversa!
- The hyperplane $\lambda_{N}=0$ separates the space of correlations into two half-spaces:
- Low correlation regime (including $\mathrm{C}=0$ ), where $\lambda_{N}>0$
- Here $\Delta l>0$ because the $\mathrm{N}^{\text {th }}$ measurement adds independent (uncorrelated) information
- High correlation regime (not including $\mathrm{C}=0$ ), where $\lambda_{\underline{N}} \leq 0$
- Here $\Delta l>0$ because the high correlations of the $\mathrm{N}^{\text {th }}$ measurement help reduce common systematics
- Note that rescaling all covariances downwards by a common scale factor to $c^{\prime}=\frac{1}{\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)} \mathbf{c}$
would effectively reduce both $\lambda_{N}$ and $\Delta l$ to 0

$$
\mathbf{c}^{\prime}=\frac{1}{\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)} \mathbf{c}
$$



- How much additional information does the $\mathrm{n}^{\text {th }}$ measurement add to the previous combination with only n -1?
- The inverse of the covariance matrix $\quad \mathcal{M}=\left(\begin{array}{cc}\mathcal{D} & \mathbf{c} \\ \tilde{\mathbf{c}} & d\end{array}\right)$ is $\quad \mathcal{M}^{-1}=\left(\begin{array}{cc}\mathcal{D}^{-1}+\frac{\left(\mathcal{D}^{-1} \mathrm{c}\right)\left(\tilde{\mathcal{c}^{-1}}\right)}{d\left(-\bar{c} \mathcal{D}^{-1} \mathbf{c}\right)} & \frac{\left(\mathcal{D}^{-1} \mathrm{c}\right)}{d-\left(\tilde{c} \mathcal{D}^{-1} \mathbf{c}\right)} \\ \frac{-\left(\tilde{\mathrm{c}} \mathcal{D}^{-1}\right)}{d-\left(\tilde{\mathfrak{c}} \mathcal{D}^{-1} \mathrm{c}\right)} & \frac{1}{d-\left(\tilde{\tilde{c}} \mathcal{D}^{-1} \mathrm{c}\right)}\end{array}\right)$
- Defining ũ as a vector of 1 's, it can be shown [van den Bos] that the information contributed by the $\mathrm{n}^{\text {th }}$ measurement is

$$
\Delta I=\frac{\left(\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)-1\right)^{2}}{d-\left(\tilde{\mathbf{c}} \mathcal{D}^{-1} \mathbf{c}\right)} \geq 0
$$

i.e. the inflow of information is always non-negative

- It is simply equal to $1 / d$ (inverse of variance of $n^{\text {th }}$ measurement) if all correlations are 0


## Information inflow and correlations

- The inflow $\boldsymbol{\Delta I}$ is zero for $\Delta I=0 \Longleftrightarrow\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)-1=0$
- This defines a hyperplane in the $n-1$ dimensional space of covariances $c$ between the $\mathrm{n}^{\text {th }}$ and all other measurements
- But $\lambda=\frac{\tilde{\mathbf{U}} \mathcal{M}^{-1}}{I}$ means that $\lambda_{n}=-\frac{1}{I} \times \frac{\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)-1}{d-\left(\tilde{\mathbf{c}} \mathcal{D}^{-1} \mathbf{c}\right)}$
- In other words, $\Delta I=0$ implies $\lambda_{n}=0$ and viceversa!
- The hyperplane $\lambda_{n}=0(\Delta \mathrm{I}=0)$ separates the space of covariances $c$ into two half-spaces: $\quad \lambda_{N}$
- Low correlation regime (including $\mathrm{C}=0$ ), where $\lambda_{n} \geq 0$
- Here $\Delta \mathrm{l}>0$ because the $\mathrm{n}^{\text {th }}$ measurement adds independent information
- High correlation regime (not including $\mathrm{c}=0$ ), where $\lambda_{n} \leq 0$

$$
\lambda_{n} \leq 0 \Longleftrightarrow\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)-1 \geq 0
$$

- Here $\Delta l>0$ because correlations help reduce common systematics!
- If one suspects that correlations are overestimated, rescaling downwards the covariance by $\mathrm{c}^{\prime}=\frac{1}{\left(\tilde{\mathbf{u}}^{-1} \mathrm{c}\right)} \mathrm{c}$ effectively reduces both $\lambda$ and $\Delta \mathrm{l}$ to 0
- This generalizes to n measurements the discussion of correlations for 2 measurements in the Lyons et al. paper


## Information sensitivity to $\rho_{\mathrm{ij}}$

- What is the effect on Fisher's information (i.e. on the combined BLUE error) of a change in the correlation $\rho_{\mathrm{ij}}$ between two measurements $y_{i}$ and $y_{j}$ ?
- Keep the $N$ variances $\sigma_{i}$ fixed and vary only the $N x(N-1) / 2$ correlations $\rho_{i j}$
- Using matrix derivatives [van den Bos], it is easy to show that the derivative of the information with respect to the correlation $\rho_{\mathrm{ij}}$ is

$$
\frac{\partial I}{\partial \rho_{i j}}=-2 I^{2} \lambda_{i} \lambda_{j} \sqrt{\sigma_{i}} \sqrt{\sigma_{j}}
$$

- The $\mathbf{N}_{\mathrm{x}}(\mathrm{N}-1) / 2$ dimensional space of correlations is split into two regimes:
- Low correlation regime (including " $\rho_{\mathrm{ij}}=0$ for all ij "), where all $\lambda>0$
- All derivatives are negative: increasing any correlation decreases the information (i.e. increasing any correlation increases the combined error - correlations "conservative")
- High correlation regime (not including " $\rho_{i j}=0$ for all ij "), where some $\lambda<0$
- Some derivatives are positive: increasing those correlations increases the information (i.e. increasing those correlations decreases the combined error - correlations not "conservative")
- The boundary between the two regimes is a hypersurface where some of the $\lambda$ are $=0$
- Information is locally minimized (derivatives are 0 ) with respect to some correlations
- We checked numerically in some examples (see next slides) that the global minimum of the information (the "most conservative choice of correlations") requires that all $\lambda$ are either $>0$ or $=0$ : in other words, a conservative option consists in removing the measurements with $\lambda<0$


## Information sensitivity to $\mathrm{M}_{\mathrm{ij}}$

- Keeping in mind that $\frac{\partial \mathcal{A}^{-1}}{\partial v_{p}}=-\mathcal{A}^{-1} \frac{\partial \mathcal{A}}{\partial v_{p}} \mathcal{A}^{-1}$ [van den Bos] and $I=\tilde{\mathbf{U}} \mathcal{M}^{-1} \mathbf{U}$, the partial derivatives of information with respect to a vector $v_{p}$ are:

$$
\frac{\partial I}{\partial v_{p}}=-\tilde{\mathbf{U}} \mathcal{M}^{-1} \frac{\partial \mathcal{M}}{\partial v_{p}} \mathcal{M}^{-1} \mathbf{U}=-I^{2} \tilde{\lambda} \frac{\partial \mathcal{M}}{\partial v_{p}} \lambda
$$

- Deriving with respect to the off-diagonal element $\mathrm{M}_{\mathrm{ij}}=\mathrm{M}_{\mathrm{ii}}(\mathrm{i} \neq \mathrm{j})$ yields

$$
\frac{\partial I}{\partial \mathcal{M}_{i j}}=-2 I^{2} \lambda_{i} \lambda_{j} \text { because }\left(\frac{\partial \mathcal{M}}{\partial \mathcal{M}_{i j}}\right)_{k l}=\delta_{i l} \delta_{j k}+\delta_{i k} \delta_{j l}=\left\{\begin{array}{ll}
1 & \text { if } i=k \text { and } j=l, \\
1 & \text { if } i=l \text { and } j=k, \\
0 & \text { otherwise },
\end{array},\right.
$$

- This shows again that the $n(n-1) / 2$ dimensional space of correlations is split into two regimes:
- Low correlation regime (including " $\mathrm{M}_{\mathrm{ij}}=0$ for all ij "), where all $\lambda>0$
- All derivatives are negative: increasing any correlation will decrease information
- High correlation regime (not including " $\mathrm{M}_{\mathrm{i}}=0$ for all ij "), where some $\lambda<0$
- Some derivatives are positive: increasing those correlations increases information
- The boundary between the two regimes, where information is minimized as $\mathrm{dl} / \mathrm{dM}=0$ ("most conservative" correlations), requires that some $\lambda=0$
- In doubt, a conservative option consists in removing the measurements with $\lambda<0$


## Is $\rho=100 \%$ "conservative"?

- Simple case of two measurements. We want:

$$
\begin{gathered}
\rho_{\text {cor }} \sigma_{A, \mathrm{cor}} \sigma_{B, \mathrm{cor}} \leq\left(\sigma_{A, \mathrm{unc}}^{2}+\sigma_{A, \mathrm{cor}}^{2}\right) \\
\rho_{\mathrm{cor}} \sigma_{A, \mathrm{cor}} \sigma_{B, \mathrm{cor}} \leq\left(\sigma_{B, \mathrm{unc}}^{2}+\sigma_{B, \mathrm{cor}}^{2}\right)
\end{gathered}
$$

- Then:

$$
\sigma_{A, \mathrm{cor}} \leq \sigma_{B, \mathrm{cor}}
$$

$$
\rho_{\mathrm{cor}} \leq \frac{\sigma_{A, \mathrm{unc}}^{2}+\sigma_{A, \mathrm{cor}}^{2}}{\sigma_{A, \mathrm{cor}} \sigma_{B, \mathrm{cor}}}=\frac{1}{\left(\frac{\sigma_{A, \mathrm{cor}}}{\sigma_{A}}\right)^{2}} \times \frac{\sigma_{A, \mathrm{cor}}}{\sigma_{B, \mathrm{cor}}}
$$



- Always true for statistically dominated

$$
\sigma_{A, \mathrm{unc}} \gg \sigma_{A, \mathrm{cor}} \text { and } \sigma_{A, \mathrm{unc}} \gg \sigma_{B, \text { cor }}
$$

- If systematic dominated, rho<1 only if:

$$
\left(\sigma_{B, \text { cor }} / \sigma_{A, \text { cor }}\right) \geq\left(\sigma_{A} / \sigma_{A, \text { cor }}\right)^{2}
$$

- For N measurements (rather: unknown correlations) the procedure for finding the maximum error is much less defined



## Derivatives w.r.t correlation re-scale factors

- Can use derivatives to analyze which correlations (between which 2 measurements) and error sources have largest effect on information
- Rescale $\mathcal{M}_{i j}^{[s]} \rightarrow\left(\mathcal{M}^{\prime}\right)_{i j}^{[s]}= \begin{cases}f_{i j}^{[s]} \mathcal{M}_{i j}^{[s]} & \text { if } i \neq j, \text { with } 0 \leq f_{i j}^{[s]} \leq 1 \\ \mathcal{M}_{i i}^{[s]} & \text { if } i=j,\end{cases}$
using a different rescaling factor for each off-diagonal element and error source
- From $\mathcal{M}_{i j}^{\prime}=\sum_{s=1}^{S}\left(\mathcal{M}^{\prime}\right)_{i j}^{[s]}=\sum_{s=1}^{S} f_{i j}^{[s]} \mathcal{M}_{i j}^{[s]}$ you get $\frac{\partial \mathcal{M}_{i j}^{\prime}}{\partial f_{i j}^{[s]}}=\mathcal{M}_{i j}^{[s]}$ that implies $\frac{\partial I^{\prime}}{\partial f_{i j}^{[s]}}=-2\left(I^{\prime}\right)^{2} \lambda_{i}^{\prime} \lambda_{j}^{\prime} \mathcal{M}_{i j}^{[s]}$
- In the following we will aggregate these derivatives ( $21 \times 16$ for $\mathrm{m}_{\text {top }}$ ) to reduce complexity, by off-diagonal element (21 for $\left.\mathbf{m}_{\text {top }}\right) \frac{\partial I^{\prime}}{\partial f_{i j}}=\sum_{s}^{S} \frac{\partial I^{\prime}}{\partial f_{i j}^{[s]}}$
and error source (16 for $\mathbf{m}_{\text {top }}$ ) $\frac{\partial I^{\prime}}{\partial f^{[s]}}=\sum_{i<j}^{n} \frac{\partial I^{\prime}}{\partial f_{i j}^{[s]}}$
- Total is the derivative by a global scale factor $\frac{\partial I^{\prime}}{\partial f}=\sum_{s}^{S} \sum_{i<j}^{n} \frac{\partial I^{\prime}}{\partial f_{i j}^{[s]}}=\sum_{s}^{S} \frac{\partial I^{\prime}}{\partial f^{[s]}}=\sum_{i<j}^{n} \frac{\partial I^{\prime}}{\partial f_{i j}}$
- We will normalize the derivatives to the total information at nominal correlations, so that the derivatives become a-dimensional numbers


## Minimizations

- We studied two procedures to modify the nominal covariance and reduce correlations (e.g. because they have been overestimated)
- Rescaling correlations as described before and minimizing information
- A procedure called "onionization" that we describe later on
- We minimize information (i.e. maximize the error to be as "conservative" as possible) with respect to the three sets of rescaling factors described in a previous slide
- Vary one rescaling factor per error source (16 for $m_{\text {top }}$ )
- For a given error source, use the same factor for all off-diagonal elements (i.e. for the whole partial covariance for that error source)
- Vary one rescaling factor per off-diagonal element ( 21 for $\mathrm{m}_{\text {top }}$ )
- For a given off-diagonal element, use the same factor for all error sources
- This may lead into non-physical space (non positive definite covariances) - give up in that case
- Vary one global rescaling factor ( 1 for $\mathrm{m}_{\text {top }}$ )


## The "onionization" prescription

- How can we ensure that $\sum_{s=1}^{S} \rho_{\mathrm{cor}}^{[s]} \sigma_{\mathrm{A}, \operatorname{cor}}^{[s]} \sigma_{\mathrm{B}, \mathrm{cor}}^{[s]} \leq\left(\sigma_{\mathrm{A}, \text { inc }}^{2}+\sum_{s=1}^{S}\left(\sigma_{\mathrm{A}, \mathrm{cor}}^{[s]}\right)^{2}\right) \quad \boldsymbol{?}$ (*)

$$
\begin{equation*}
\sum_{s=1}^{S} \rho_{\mathrm{cor}}^{[s]} \sigma_{\mathrm{A}, \text { cor }}^{[[] s} \sigma_{\mathrm{B}, \text { cor }}^{[s]} \leq\left(\sigma_{\mathrm{B}, \mathrm{unc}}^{2}+\sum_{s=1}^{S}\left(\sigma_{\mathrm{B}, \mathrm{cor}}^{[s]}\right)^{2}\right) \tag{*}
\end{equation*}
$$

We considered: $\rho_{\text {cor }}^{[s]} \sigma_{\mathrm{A}, \text { cor }}^{[s]} \sigma_{\mathrm{B}, \text { cor }}^{[s]} \leq\left(\sigma_{\mathrm{A}, \text { una }}^{2}+\sum_{s^{\prime}}^{S}\left(\sigma_{\mathrm{A}, \text { orr }}^{\left[s^{1}\right]}\right)^{2}\right) \quad \forall$.

$$
\rho_{\operatorname{cor}}^{[s]} \sigma_{\mathrm{A}, \text { cor }}^{[s]} \sigma_{\mathrm{B}, \text { cor }}^{[s]} \leq\left(\sigma_{\mathrm{B}, \text { inc }}^{2}+\sum_{s^{\prime}=1}^{S}\left(\sigma_{\mathrm{B}, \text { ocr }}^{\left[s^{[j]}\right]}\right)^{2}\right) \quad \forall s
$$

- But we prefer the "onionization" of every source of uncertainty: remain $\leq$
i.e.: : $\begin{array}{lll}\rho_{\mathrm{cor}}^{[s]} \sigma_{A, \mathrm{cor}}^{[s]} \sigma_{B}^{[s]}[\mathrm{cor} & \leq\left(\sigma_{A, \operatorname{cor}}^{[s]}\right)^{2} & \forall s \\ \rho_{\mathrm{cor}}^{[s]} \sigma_{A, \mathrm{cor}}^{[s]} \sigma_{B, \mathrm{cor}}^{[s]} \leq\left(\sigma_{B, \mathrm{cor}}^{[s]}\right)^{2} & \forall s\end{array}$
$\left(\begin{array}{ccccc}\left(\sigma_{A, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{A, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{A, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{A, \mathrm{cor}}^{[s]}\right)^{2} & \ldots \\ \left(\sigma_{A, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{B, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{B, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{B, \mathrm{cor}}^{[s]}\right)^{2} & \ldots \\ \left(\sigma_{A, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{B, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{C, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{C, \mathrm{cor}}^{[s]}\right)^{2} & \ldots \\ \left(\sigma_{A, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{B, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{C, \mathrm{cor}}^{[s]}\right)^{2} & \left(\sigma_{D, \mathrm{cor}}^{[s]}\right)^{2} & \ldots \\ \ldots & \ldots & \ldots & \ldots & \ldots\end{array}\right)$
- However this may transform overestimates of $\rho$ into underestimates
- It means assuming that the difference in sensitivity between measurements to a given systematic effect cannot be used to constrain or determine the actual size of that effect
- Which is generally wrong - there are legitimate cases for exploiting this leverage effect
- And with $>2$ measurements this condition ( ${ }^{*}$ ) is not even enough to ensure $\lambda>0$ for each $\lambda$
- Of course the only correct way forward is
- measure the correlations between measurements (from the data, the MC parameters...)
- if not possible, assess case by case whether $100 \%$ (or e.g. onionizing) is appropriate



## Onionization example - RAD @ LHC miop

| CMS11ll | ( 0.59 | 0.61 | 0.73 | 0.39 | 0.91 | 0.65 | 0.94 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CMS11 $\mu \mathrm{j}$ | 0.61 | 0.64 | 0.75 | 0.40 | 0.95 | 0.68 | 0.98 |
| CMS10ll | 0.73 | 0.75 | 0.89 | 0.48 | 1.12 | 0.80 | 1.15 |
| ATL11lj | 0.39 | 0.40 | 0.48 | 1.02 | 0.60 | 1.72 | 2.47 |
| CMS10lj | 0.91 | 0.95 | 1.12 | 0.60 | 1.41 | 1.01 | 1.45 |
| ATL11aj | 0.65 | 0.68 | 0.80 | 1.72 | 1.01 | 2.89 | 4.16 |
| ATL10lj | ( 0.94 | 0.98 | 1.15 | 2.47 | 1.45 | 4.16 | 5.98 |

Table 13: Onionization of the RAD covariance. Nominal RAD.

| CMS11ll | ( 0.59 | 0.59 | 0.59 | 0.59 | 0.59 | 0.59 | 0.59 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CMS11 $\mu \mathrm{j}$ | 0.59 | 0.64 | 0.64 | 0.64 | 0.64 | 0.64 | 0.64 |
| CMS10ll | 0.59 | 0.64 | 0.89 | 0.89 | 0.89 | 0.89 | 0.89 |
| ATL11lj | 0.59 | 0.64 | 0.89 | 1.02 | 1.02 | 1.02 | 1.02 |
| CMS10lj | 0.59 | 0.64 | 0.89 | 1.02 | 1.41 | 1.41 | 1.41 |
| ATL11aj | 0.59 | 0.64 | 0.89 | 1.02 | 1.41 | 2.89 | 2.89 |
| ATL10lj | ( 0.59 | 0.64 | 0.89 | 1.02 | 1.41 | 2.89 | 5.98 |

Table 14: Onionization of the RAD covariance. Onion upper bound.

| CMS11ll | ( 0.59 | 0.59 | 0.59 | 0.39 | 0.59 | 0.59 | 0.59 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CMS11 $\mu \mathrm{j}$ | 0.59 | 0.64 | 0.64 | 0.40 | 0.64 | 0.64 | 0.64 |
| CMS10ll | 0.59 | 0.64 | 0.89 | 0.48 | 0.89 | 0.80 | 0.89 |
| ATL11lj | 0.39 | 0.40 | 0.48 | 1.02 | 0.60 | 1.02 | 1.02 |
| CMS10lj | 0.59 | 0.64 | 0.89 | 0.60 | 1.41 | 1.01 | 1.41 |
| ATL11aj | 0.59 | 0.64 | 0.80 | 1.02 | 1.01 | 2.89 | 2.89 |
| ATL10lj | ( 0.59 | 0.64 | 0.89 | 1.02 | 1.41 | 2.89 | 5.98 |

PRELIMINARY!
WORK IN PROGRESS!

Table 15: Onionization of the RAD covariance. Onionized RAD.

# LHC $m_{\text {top }}$ - conservative correlations? ${ }^{\text {CERNT }}$ 

|  | Comb. $m_{\text {top }}$ | Stat | iJES | bJES | dJES | rJES | Lept | MC | Rad | CR | PDF | DTMO | UE | BGMC | BGDT | Meth | MHI | $\chi^{2} / \mathrm{ndof}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| PUBLISHED | $173.33 \pm 1.40$ | 0.47 | 0.38 | 0.68 | 0.07 | 0.06 | 0.01 | 0.04 | 0.69 | 0.55 | 0.01 | 0.19 | 0.47 | 0.01 | 0.16 | 0.13 | 0.25 | 2.49/6 |
| MinimizeGlobalFactor1 | $173.29 \pm 1.41$ | 0.40 | 0.36 | 0.71 | 0.19 | 0.05 | 0.01 | 0.06 | 0.71 | 0.53 | 0.04 | 0.22 | 0.47 | 0.04 | 0.14 | 0.12 | 0.26 | 2.11/6 |
| MinimizeErrorSrc18 | $173.22 \pm 1.44$ | 0.35 | 0.35 | 0.72 | 0.25 | 0.02 | 0.00 | 0.09 | 0.73 | 0.54 | 0.05 | 0.24 | 0.48 | 0.10 | 0.12 | 0.12 | 0.29 | 1.97/6 |
| Onionize | $173.08 \pm 1.44$ | 0.31 | 0.33 | 0.72 | 0.26 | 0.01 | 0.00 | 0.07 | 0.76 | 0.54 | 0.06 | 0.24 | 0.47 | 0.09 | 0.10 | 0.11 | 0.28 | 1.67/6 |
| MinimizeOffDiagonal21 | $173.02 \pm 1.46$ | 0.32 | 0.36 | 0.74 | 0.23 | 0.00 | 0.00 | 0.07 | 0.76 | 0.54 | 0.06 | 0.24 | 0.49 | 0.10 | 0.10 | 0.12 | 0.31 | 1.83/6 |
| CMS11 $\mu$ j $\oplus$ ATL11lj | $173.02 \pm 1.46$ | 0.32 | 0.36 | 0.74 | 0.23 | 0.00 | 0.00 | 0.07 | 0.76 | 0.54 | 0.06 | 0.24 | 0.49 | 0.10 | 0.10 | 0.12 | 0.31 | 0.64/1 |
| NO CORRELATIONS | $173.25 \pm 1.05$ | 0.39 | 0.22 | 0.48 | 0.39 | 0.09 | 0.02 | 0.08 | 0.47 | 0.29 | 0.07 | 0.19 | 0.32 | 0.07 | 0.16 | 0.11 | 0.18 | 1.09/6 |

Table 9: LHC $m_{\text {top }}$ BLUE combinations with nominal and modified correlations. From top to bottom: published results with all 7 measurements and nominal correlations; all correlations rescaled by 1 global factor, maximizing the error; correlations rescaled by 18 scale factors, one per error source, maximizing the error; onionization of the covariance matrix in each error source; correlations rescaled by 21 scale factors, one per off-diagonal element (same factor for all error sources), maximizing the error; results with only the two measurements with positive $\lambda$ and nominal correlations; all correlations set to 0 everywhere. From top to bottom, the combination essentially moves from a high-correlation regime to a low-correlation regime, possibly passing through the absolute minimum at the minimization by off-diagonal element, which seems to coincide with the combination with two measurements.

- Absolute minimum is most likely at "MinimizeOffDiagonal21"
- Checked that 20 information derivatives are 0 (for all covariances but that of CMS11 $\mu \mathrm{j}$ and ATL11j)
- Indeed (remember dl/dM $\mathrm{M}_{\mathrm{ij}} \sim-\lambda_{\mathrm{i}} \lambda_{\mathrm{j}}$ !) all 5 lambdas become 0 except for CMS11 $\mu \mathrm{j}$ and ATL11jI
- Not surprisingly results coincide with the combination of CMS11 $\mu \mathrm{j}$ and ATL11jl alone, but for the $\chi^{2}$
- A very conservative option consists in not using the results with $\lambda<0$ !
- Some negative BLUE coefficients remain for the other two minimizations and (less) for the onionization procedure
- Note again that all of these procedures change the sys/stat balance, not only the total error!
- Some of these procedures may also not be completely physical (covariances not positive definite?)
- Carefully measuring correlations is much better than any rule of thumb!


## Recommendations?

- Assume you have chosen your best estimate of correlations ("nominal")
- If all $\lambda>0$ at "nominal" correlations, you are in a low-correlation regime
- Even if correlations are high, their estimates are "conservative" enough: they do not increase the available information, they decrease it - nothing to worry about! :)
- If some $\boldsymbol{\lambda}$ are negative, you are in a high-correlation regime
- You should be aware that these high correlations are adding information!
- Effectively, the measurements with $\lambda<0$ are only adding information through these correlations
- We advise that all correlations should then be reviewed on a case by case basis
- There are legitimate cases where high correlations exist and may help you constrain common systematics (e.g. you varied same MC parameter or depend on a common $E_{\text {beam }}$ measurement)
- But if you "conservatively" estimated correlations to be $100 \%$, then you are definitely wrong!
- If correlations are overestimated, you are certainly underestimating the combined systematics, even more than you are underestimating the total combined error - and your central values are shifted, too
- What can you do to provide a more realistic (or "conservative") estimate?
- Prioritize the correlations to analyze, using information derivatives
- Then measure these correlations with your data and MC if this is possible
- For correlations you cannot measure precisely, consider some of the tools we presented
- You may try to minimize information in those correlations alone, or try the onionization procedure (if $\sigma_{S}$ for $A$ and $B$ is correlated, and $\sigma_{S, B}>\sigma_{S, A}$, strip out $\sigma_{S^{\prime}, B}=\sigma^{2}{ }_{S, B} \sigma^{2}{ }_{\mathrm{S}, \mathrm{A}}$ as uncorrelated)
- If you have no clue, omitting the measurements with $\lambda<0$ is the most conservative option


# $I_{\text {corr }}=I-I_{A}-I_{B}=\frac{1}{\sigma_{\hat{Y}}^{2}}-\frac{1}{\sigma_{A}^{2}}-\frac{1}{\sigma_{B}^{2}}$ <br> <br> LHC miop marginal information weights 

 <br> <br> LHC miop marginal information weights}

|  | $\begin{gathered} m_{\text {top }} \\ {[\mathrm{GeV}]} \end{gathered}$ | $\begin{gathered} \Delta I \\ {\left[1 / \mathrm{GeV}^{2}\right]} \end{gathered}$ | $\lambda$ | $\frac{\|\lambda\|}{\sum\|\lambda\|}$ | ( I | $\sum \frac{\Delta I}{I}$ | $\begin{gathered} \Delta I^{\prime} \\ {\left[1 / \mathrm{GeV}^{2}\right]} \\ \hline \end{gathered}$ | $\sum \frac{\Delta I^{\prime}}{I}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CMS11 $\mu \mathrm{j}$ <br> Low corr. | $172.64 \pm 1.53$ | 0.425 | $84.3 \%$ | $65.7 \%$ | 83.8\% | 83.8\% | $\begin{array}{r} \hline 0.425 \\ -0.049 \end{array}$ | 74.1\% |
| ATL11lj <br> Low corr. | $174.53 \pm 2.39$ | $\begin{array}{r} \hline 0.175 \\ -0.131 \end{array}$ | $\begin{array}{r} 29.9 \% \\ \hline \end{array}$ | $23.3 \%$ | $\begin{array}{r} 34.5 \% \\ -25.9 \% \end{array}$ | 8.6\% | $\begin{array}{r} 0.175 \\ \hline-0.082 \end{array}$ | 18.4\% |
| ATL10lj <br> Low corr. <br> High corr | $169.33 \pm 6.32$ | $\begin{array}{r} 0.025 \\ \hline-0.025 \\ 0.030 \end{array}$ | $-6.8 \%$ |  | $\begin{array}{r} 4.9 \% \\ \hline-4.9 \% \\ 5.8 \% \end{array}$ | 5.8\% | 0.025 | 5.8\% |
| Civishll <br> Low corr. <br> High corr. | $173.30 \pm 2.96$ | $\begin{array}{r} \hline 0.114 \\ -0.114 \\ 0.006 \\ \hline \end{array}$ | $-4.8 \%$ $\qquad$ | $\begin{array}{r} 3.7 \% \\ - \\ - \end{array}$ | $\begin{array}{r} \hline 22.5 \% \\ \hline-22.5 \% \\ 1.1 \% \\ \hline \end{array}$ | 1.1\% | 0.006 | 1.1\% |
| CMS10ll <br> Low corr. <br> High corr. | $175.50 \pm 6.49$ $-$ | $\begin{array}{r} \hline 0.024 \\ -0.024 \\ 0.003 \end{array}$ | $-1.9 \%$ | $\begin{gathered} 1.5 \% \\ - \\ \hline \end{gathered}$ | $\begin{gathered} 4.7 \% \\ \hline-4.7 \% \\ 0.7 \% \end{gathered}$ | 0.7\% | 0.003 | 0.7\% |
| ATL11aj Low corr. High corr. | $174.90 \pm 4.44$ | $\begin{array}{r} 0.051 \\ -0.051 \\ 0.000 \end{array}$ | $-0 / 4 \%$ - - - | $\begin{gathered} 0.3 \% \\ \hline \end{gathered}$ | $\begin{array}{r} \hline 10.0 \% \\ \hline-10.0 \% \\ 0.0 \% \end{array}$ | 0.0\% | 0.000 | 0.0\% |
| CMS10lj <br> Low corr. <br> High corr. | $173.10 \pm 3.41$ | $\begin{array}{r}0.086 \\ -0.086 \\ 0.009 \\ \hline\end{array}$ | $\begin{array}{r} -0.2 \% \\ - \\ - \end{array}$ | $0.2 \%$ | $\begin{array}{r} 16.9 \% \\ \hline-16.9 \% \\ 0.0 \% \\ \hline \end{array}$ | 0.0\% | 0.000 | 0.0\% |
| Total | $173.33 \pm 1.40$ | 0.507 | 100.0\% | 100.0\% | 100.0\% | 100.0\% | 0.507 | 100.0\% |

Table 7: LHC $m_{\text {top }}$ marginal information weights. The measurements are listed in the following order: first those with $\lambda>0$, ordered by decreasing $\lambda$; then those with $\lambda<0$, ordered by decreasing $|\lambda|$. The information inflows $\Delta I$ represent the marginal contributions of adding to the combination the $n^{\text {th }}$ measurement and its correlations to the previous $n-1$ measurements. The "effective" information inflows $\Delta I^{\prime}$ for the measurements are computed as follows: for those with $\lambda<0$, they are are equal to $\Delta I$; for those with $\lambda \geq 0$ (the first two only), a BLUE combination of those measurements alone is performed, and the total information contribution from the correlations between them is split up so that the totals for each measurement and its correlations are proportional to the BLUE coefficients ( $80.1 \%$ and $19.9 \%$ ) in that combination. As a result, the two columns with the information weights computed as

Split low and high correlations by
rescaling cas

$$
\mathbf{c}^{\prime}=\frac{1}{\left(\tilde{\mathbf{u}} \mathcal{D}^{-1} \mathbf{c}\right)} \mathbf{c}
$$

For adding ATL10lj, this factor is 0.543

|  | $m_{\text {top }}$ <br> $[\mathrm{GeV}]$ | $\Delta I$ <br> $\left[1 / \mathrm{GeV}^{2}\right]$ | $\lambda$ | $\frac{\|\lambda\|}{\sum\|\lambda\|}$ | $\frac{\Delta I}{I}$ |
| :--- | :---: | ---: | :---: | :---: | :---: |
| CMS11 $\mu \mathrm{j}$ | $172.64 \pm 1.53$ | 0.425 | $84.3 \%$ | $65.7 \%$ | $83.8 \%$ |
| ATL11lj | $174.53 \pm 2.39$ | 0.175 | $29.9 \%$ | $23.3 \%$ | $34.5 \%$ |
| ATL10lj | $169.33 \pm 6.32$ | 0.025 | $-6.8 \%$ | $5.3 \%$ | $4.9 \%$ |
| CMS11ll | $173.30 \pm 2.96$ | 0.114 | $-4.8 \%$ | $3.7 \%$ | $22.5 \%$ |
| CMS10ll | $175.50 \pm 6.49$ | 0.024 | $-1.9 \%$ | $1.5 \%$ | $4.7 \%$ |
| ATL11aj | $174.90 \pm 4.44$ | 0.051 | $-0.4 \%$ | $0.3 \%$ | $10.0 \%$ |
| CMS10lj | $173.10 \pm 3.41$ | 0.086 | $-0.2 \%$ | $0.2 \%$ | $16.9 \%$ |
| Correlations | - | -0.392 | - | - | $-77.4 \%$ |
| Total | $173.33 \pm 1.40$ | 0.507 | $100.0 \%$ | $100.0 \%$ | $100.0 \%$ |

These are two possible ways to split up the info contribution from the correlations, based on "marginal" contributions

Note that two measurements contribute $0.0 \%$ information
$\sum \frac{\Delta I^{\prime}}{I}$ and $\sum \frac{\Delta I^{\prime}}{I}$ differ only for $\lambda \geq 0$.


