Transverse Momentum Resummation in Top Quark Pair Production

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Outline

- Introduction: top quark pair production
- Introduction: **Q**_T resummation
- Q_T resummation in top quark pair production
- Summary

Top quark pair production



- Main production mechanism of top quarks at hadron colliders
- NNLO results for total cross section emerging
 Bärnreuther, Czakon, Mitov: 1204.5201
 Czakon, Mitov: 1207.0236 & 1210.6832
 waiting for the gg channel...
- Pair differential distributions interesting
 - Searching for new resonances in the invariant mass distribution
 - Forward-backward asymmetry shows intriguing dependence on the invariant mass and the transverse momentum of the pair

Invariant mass distribution Ahrens, Ferroglia, Neubert, Pecjak, LLY: 1003.5827

NLO+NNLL threshold resummed prediction



Invariant mass distribution in the boosted regime

- For *M* much larger than *m_t*, may simultaneously resum threshold logs and mass logs
- Framework set up Ferroglia, Pecjak, LLY: 1205.3662

$$C_{ij}(z, M, m_t, \cos \theta, \mu_f) = C_D^2(m_t, \mu_f) \operatorname{Tr} \left[\boldsymbol{H}_{ij}(M, t_1, \mu_f) \, \boldsymbol{S}_{ij}(\sqrt{\hat{s}}(1-z), t_1, \mu_f) \right]$$

$$\otimes C_{ff}^{ij}(z, m_t, \mu_f) \otimes C_{t/t}(z, m_t, \mu_f) \otimes C_{t/t}(z, m_t, \mu_f)$$

cross section

$$\otimes S_D(m_t(1-z), \mu_f) \otimes S_D(m_t(1-z), \mu_f) + \mathcal{O}(1-z) + \mathcal{O}\left(\frac{m_t}{M}\right)$$

- NNLO massless soft function calculated allows soft+virtual at NNLO Ferroglia, Pecjak, LLY: 1207.4798
- Phenomenological studies in progress

partonic

Small Q7 region

- Suppress the gluon channel enhance the forward-backward asymmetry
- Theoretically challenging fixed-order divergent, resummation necessary
- Existing parton shower tools need to be validated in this region by analytic calculations



D0 Collaboration: 1107.4995

Q_T resummation

• Originally developed by Collins, Soper, Sterman (CSS)

Collins, Soper, Sterman: Nucl. Phys. B 250, 199 (1985)

• Refinement by Catani, de Florian, Grazzini

Catani, de Florian, Grazzini: hep-ph/0008014

- New SCET based approaches
 - Becher, Neubert: 1007.4005
 - Chiu, Jain, Neill, Rothstein: 1104.0881
 - Echevarria, Idilbi, Scimemi: 1111.4996

Applied to color-neutral final states so far!

Q_T resummation in Higgs production

Becher, Neubert, Wilhelm: 1212.2621

$$d\sigma = \sigma_{0}(\mu) C_{t}^{2}(m_{t}^{2},\mu) \left| C_{S}(-m_{H}^{2},\mu) \right|^{2} \frac{m_{H}^{2}}{\tau s} dy \frac{d^{2}q_{\perp}}{(2\pi)^{2}} \int d^{2}x_{\perp} e^{-iq_{\perp} \cdot x_{\perp}}$$

$$\times 2\mathcal{B}_{c}^{\mu\nu}(\xi_{1},x_{\perp},\mu) \mathcal{B}_{\bar{c}}_{\mu\nu}(\xi_{2},x_{\perp},\mu) \mathcal{S}(x_{\perp},\mu),$$
transverse PDFs
transverse soft function
trivial with analytic regulator

$$\bar{C}_{gg\leftarrow ij}(z_1, z_2, q_T^2, m_H^2, \mu) = \frac{1}{2} \int_0^\infty dx_T \, x_T \, J_0(x_T q_T) \, \exp\left[g_A(\eta, L_\perp, a_s)\right] \\ \times \sum_{n=1,2} \, \bar{I}_{g\leftarrow i}^{(n)}(z_1, L_\perp, a_s) \, \bar{I}_{g\leftarrow j}^{(n)}(z_2, L_\perp, a_s) \,,$$

See Becher's talk for details

Q₇ subtraction

- Closely related to Q_T resummation
- Exploits behavior at small Q_T to subtract IR divergences
- Works in principle for any process where Q_T resummation is available
- One of the first subtraction methods applied successfully at NNLO

$$d\sigma^{F}_{(N)NLO} = \mathcal{H}^{F}_{(N)NLO} \otimes d\sigma^{F}_{LO} + \left[d\sigma^{F+\text{jets}}_{(N)LO} - d\sigma^{CT}_{(N)LO} \right]$$

Q_T resummation in $t\bar{t}$ production

- Early attempts naively supplementing the CSS formalism with final state interactions
 Berger, Meng: hep-ph/9310341 Mrenna, Yuan: hep-ph/9606361
 - Results don't agree between the two works!
 - (Partial) NLL only. No initial-final interference.
- New SCET based systematic factorization/resummation framework
 Li, Li, Shao, LLY, Zhu: 1208.5774
 - Works in principle to all orders

Rest of this talk!

NNLL numerical results

Factorization formula

• An all-order factorization formula is proposed

$$\frac{d^{4}\sigma}{dQ_{T}^{2}dydMd\cos\theta} = \frac{\pi\beta_{t}}{sM} \int \frac{d^{2}x_{\perp}}{4\pi} e^{-iQ_{\perp}\cdot x_{\perp}} \\ \times \operatorname{Tr} \left[\boldsymbol{H}(\boldsymbol{M},\cos\theta,\boldsymbol{m}_{t},\boldsymbol{\mu}) \, \boldsymbol{S}(\boldsymbol{L}_{\perp},\boldsymbol{M},\cos\theta,\boldsymbol{m}_{t},\boldsymbol{\mu}) \right] \\ \times \mathcal{B}(\xi_{1},\boldsymbol{L}_{\perp},\boldsymbol{\mu}) \, \bar{\mathcal{B}}(\xi_{2},\boldsymbol{L}_{\perp},\boldsymbol{\mu}) \\ \text{same as in} \\ \text{threshold resummation} \\ \text{Ahrens, Ferroglia, Neubert,} \\ \text{Pecjak, LLY: 1003.5827} \\ \text{same as in} \\ \text{DY and Higgs} \\ \text{Becher, Neubert, Wilhelm:} \\ 1109.6027 \& 1212.2621 \end{aligned}$$

Possible factorization breaking effects

Mitov, Sterman: 1209.5798



• Due to interactions among top quarks and spectators

See also Collins, Qiu: 0705.2141 Rogers, Mulders: 1001.2977

- Numerical impacts unclear
- We focus on the partonic cross section (neglecting spectators)

The transverse soft function

$$\begin{aligned} \mathbf{S}(L_{\perp}, M, \cos\theta, m_t, \mu) &= \frac{1}{d_R} \sum_{X_s} \int \frac{d\phi_t}{2\pi} d^2 Q_{\perp} e^{iQ_{\perp} \cdot x_{\perp}} \\ &\times \langle 0 | Y_n^{\dagger} Y_{\bar{n}}^{\dagger} Y_t^{\dagger} Y_{\bar{t}}^{\dagger} | X_s \rangle \, \delta^{(2)}(Q_{\perp} + \hat{P}_{\perp}) \, \langle X_s | Y_n Y_{\bar{n}} Y_t Y_{\bar{t}} | 0 \rangle \end{aligned}$$

- Definition similar to the threshold one, but calculation much more complicated due to the presence of the transverse vector
- We had to integrate over the azimuthal angle calculation for the more exclusive case failed ideas?

NLO soft function

$$\begin{aligned} \mathbf{S}_{i\bar{i}}^{(1)} &= 4L_{\perp} \left(2\mathbf{w}_{i\bar{i}}^{13} \ln \frac{-t_1}{m_t M} + 2\mathbf{w}_{i\bar{i}}^{23} \ln \frac{-u_1}{m_t M} + \mathbf{w}_{i\bar{i}}^{33} \right) \\ &- 4 \left(\mathbf{w}_{i\bar{i}}^{13} + \mathbf{w}_{i\bar{i}}^{23} \right) \operatorname{Li}_2 \left(1 - \frac{t_1 u_1}{m_t^2 M^2} \right) + 4\mathbf{w}_{i\bar{i}}^{33} \ln \frac{t_1 u_1}{m_t^2 M^2} \\ &- 4 \left(\mathbf{w}_{i\bar{i}}^{13} + \mathbf{w}_{i\bar{i}}^{23} \right) \operatorname{Li}_2 \left(1 - \frac{t_1 u_1}{m_t^2 M^2} \right) + 4\mathbf{w}_{i\bar{i}}^{33} \ln \frac{t_1 u_1}{m_t^2 M^2} \\ &+ 4 \ln x_s \ln \cos \frac{\theta}{2} \,. \end{aligned}$$

- Rapidity divergence in intermediate steps regularized using method in Becher, Bell: 1112.3907
- No rapidity divergence in the final soft function (as expected)
- Dependence on the scattering angle cannot be written as simple functions of the invariants *t*₁ and *u*₁ (unlike the threshold soft function)

Consistency checks

- NLO expansion compared with exact result at small **Q**_T
- Soft function important for the nice agreement
- Constant term checked by reproducing the NLO total cross section
- Would be great to compare the NNLO expansion with the NLO *t*t̄+jet code



Resummation formula

$$C_{i\bar{i}\leftarrow ab}(z_1, z_2, q_T, M, \cos\theta, m_t, \mu) = \frac{1}{2} \int_0^\infty db \, b \, J_0(bq_T)$$

$$\times \exp\left[g_i(\eta_i, L_\perp, \alpha_s)\right] \left[\bar{I}_{i/a}(z_1, L_\perp, \alpha_s) \, \bar{I}_{\bar{i}/b}(z_2, L_\perp, \alpha_s) + \delta_{gi} \, \bar{I}'_{g/a}(z_1, L_\perp, \alpha_s) \, \bar{I}'_{g/b}(z_2, L_\perp, \alpha_s)\right]$$

$$\times \operatorname{Tr}\left[\boldsymbol{H}_{i\bar{i}}(M, \cos\theta, m_t, \mu_h, \mu) \, \boldsymbol{S}_{i\bar{i}}(L_\perp, M, \cos\theta, m_t, \mu)\right]$$

Scale choices follow Becher, Neubert, Wilhelm for DY and Higgs

 $\mu_i = q_i^* + Q_T$

• Much weaker dependence on non-perturbative physics due to larger mass — protected all the way down to $Q_T = 0$

$$q_q^* \gtrsim 3.0 \text{GeV} \qquad q_g^* \gtrsim 14.0 \text{GeV}$$

• Evolution of the hard function given in

Ferroglia, Neubert, Pecjak, LLY: 0907.4791 & 0908.3676 Ahrens, Ferroglia, Neubert, Pecjak, LLY: 1003.5827

Resummed prediction

- Good perturbative convergence, significantly reduced scale dependence from NLL to NNLL
- Consistent with NLO parton shower results
- Agree well with existing data.
 More data required to check the shape in the low *Q_T* region!



Forward-backward asymmetry

• LO parton shower can generate non-zero asymmetry due to color coherence

Skands, Webber, Winter: 1205.1466

- In our framework, color coherence is built into the soft function
- Our prediction agrees well with NLO parton shower



NNLO Q_T subtraction in $t\bar{t}$ production?

$$\frac{d^{4}\sigma}{dQ_{T}^{2}dydMd\cos\theta} = \frac{\pi\beta_{t}}{sM} \int \frac{d^{2}x_{\perp}}{4\pi} e^{-iQ_{\perp}\cdot x_{\perp}}$$

$$\times \operatorname{Tr}\left[\boldsymbol{H}(M,\cos\theta,m_{t},\mu) \boldsymbol{S}(L_{\perp},M,\cos\theta,m_{t},\mu)\right]$$

$$\times \mathcal{B}(\xi_{1},L_{\perp},\mu) \,\bar{\mathcal{B}}(\xi_{2},L_{\perp},\mu)$$
???

Gehrmann, Lübbert, LLY 1209.0682 and work in progress

Summary

- An all-order Q₇ resummation framework for top quark pair production is proposed for the first time
- Numerical results are given at NLO+NNLL accuracy
- Consistent with NLO parton shower tools and with experimental data
- May provide an alternative NNLO subtraction method for $t\bar{t}$ production
- Towards matching at NNLO

Thank you!