

# Double parton distributions - bounds and evolution

Tomas Kasemets

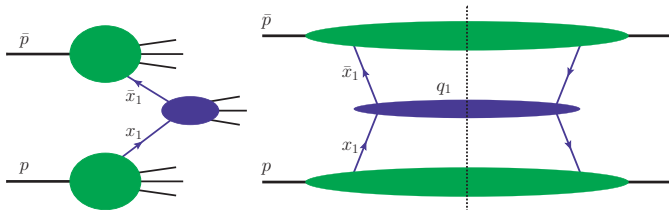


DESY  
Theory Group

*arXiv: 1210.5434, 1303.0842*

In collaboration with M. Diehl

# What is double parton scattering?

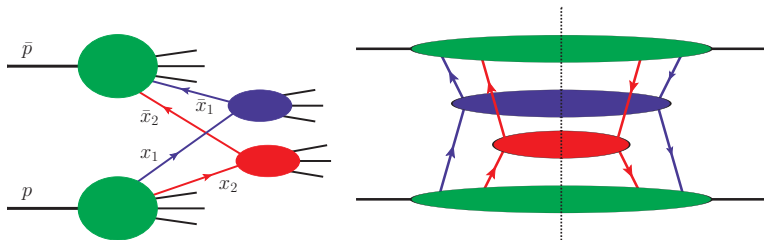


- Single parton scattering cross section

$$\sigma = \sum_{ij} \int dx_1 \int d\bar{x}_1 f_i(x_1) f_j(\bar{x}_1) \hat{\sigma}_{ij}$$

(Assuming factorization of hard scattering)

# What is double parton scattering?



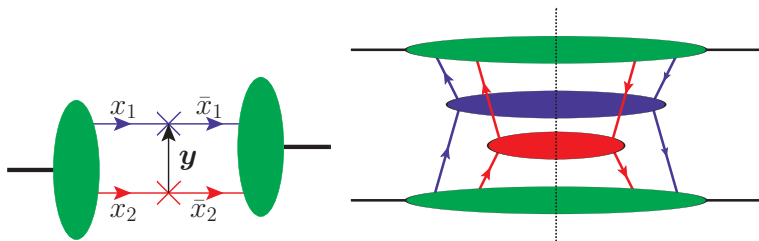
- Double parton scattering (DPS) cross section

$$\sigma_{DPS} \sim \sum_{ijkl} \int dx_1 dx_2 \int d\bar{x}_1 d\bar{x}_2 \int d^2\mathbf{y} f_{ik}(x_1, x_2, \mathbf{y}) f_{jl}(\bar{x}_1, \bar{x}_2, \mathbf{y}) \hat{\sigma}_{ij} \hat{\sigma}_{kl}$$

(Assuming factorization of hard scatters)

notice  $\mathbf{y}$  dependence

# What is double parton scattering?



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(Assuming factorization of hard scatters)

notice  $\mathbf{y}$  dependence

- Inclusive DPS cross section power suppressed compared to SPS cross section

$$\sigma_{DPS}/\sigma_{SPS} \sim \Lambda^2/Q^2$$

- Why study DPS?

- Inclusive DPS cross section power suppressed compared to SPS cross section

$$\sigma_{DPS}/\sigma_{SPS} \sim \Lambda^2/Q^2$$

- Why study DPS:

- Differential cross section (in  $\mathbf{q}_i$ ) not power suppressed
- $\sigma_{DPS}/\sigma_{SPS}$  expected to increase with collider energy
- Proton structure
- Background
- Multiparton interactions and Monte Carlo generators
- Experimental evidence found at ISR, SPS, Tevatron and LHC
- Many new effects, not present in single parton scattering
- No systematic theory

# Polarized double parton distributions

- Spin correlations, between partons in same proton

( M. Mekhfi, 1985; M. Diehl, A. Schäfer, 2011)

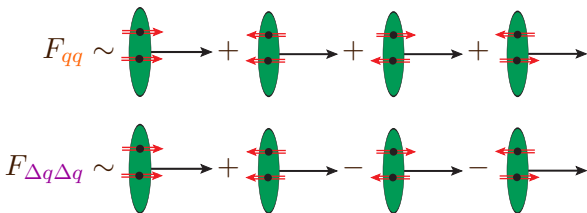
- Correlations described by parton distributions

$$F_{q_1 q_2} \sim \langle p | (\bar{q}_1 \Gamma_1 q_1) (\bar{q}_2 \Gamma_2 q_2) | p \rangle$$

- $\Gamma_{1/2} = \Gamma_q, \Gamma_{\Delta q}, \Gamma_{\delta q}^j$  projection operators
  - Unpolarized ( $q$ ), longitudinally- ( $\Delta q$ ) and transversely-polarized ( $\delta q$ ) quarks
  - Index  $j = 1, 2$  corresponds to the transverse spin-vector

Remember:  $q$ ,  $\Delta q$  and  $\delta q$

- Unpolarized and longitudinally polarized quarks



- Transverse polarization ( $F_{q\delta q}^i, F_{\delta q q}^i, F_{\delta q \delta q}^{ij}$ ) - helicity interference



- Polarized double parton distributions
  - Unpolarized and longitudinally polarized quarks

$$F_{qq} = f_{qq}(x_1, x_2, \mathbf{y})$$

$$F_{\Delta q \Delta q} = f_{\Delta q \Delta q}$$

- Singly transversely polarized quarks ( $\tilde{y}^i = y^j \epsilon^{ij}$ )

$$F_{q\delta q}^i = M \tilde{y}^i f_{q\delta q}$$

$$F_{\delta q q}^i = M \tilde{y}^i f_{\delta q q}$$

- Doubly transversely polarized quarks

$$F_{\delta q \delta q}^{ij} = \delta^{ij} f_{\delta q \delta q} + (2y^i y^j - y^2 \delta^{ij}) M^2 f_{\delta q \delta q}^t$$

6 different polarized distributions

- Double Drell-Yan cross section ( $W^\pm, Z, \gamma^*$ )
  - Longitudinal polarization change both magnitude and angular modulation
  - Azimuthal correlations

$$d\sigma^{(2)} \sim \sin^2 \theta_1 \sin^2 \theta_2 \left[ A \cos 2(\varphi_1 - \varphi_2) - B \sin 2(\varphi_1 - \varphi_2) \right] \\ \times \int d^2 \mathbf{y} f_{\delta q_1 \delta q_2}(x_1, x_2, \mathbf{y}) \bar{f}_{\delta \bar{q}_1 \delta \bar{q}_2}(\bar{x}_1, \bar{x}_2, \mathbf{y})$$

(transversely polarized quarks in both interactions)

(TK, M. Diehl, 2012 )

- Similar features in other cross sections (such as double dijet)
- Double gluon and mixed quark-gluon distributions
- For gluon  $g$  unpolarized,  $\Delta g$  longitudinally polarized and  $\delta g$  linearly polarized.

polarized distributions have profound effect on cross section

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how big can the polarized distributions be?

## Positivity bounds

- Many distributions, poorly known  $f_{ab}(x_1, x_2, \mathbf{y})$
- Derive constraints on their sizes
- Constraints from probabilistic interpretation
- Can derive positivity bounds, similar to single parton distributions  
( A. Bacchetta, M. Boglione, P.J. Mulders, 1999; M. Diehl, Ph. Hägler, 2005)
- Positivity bounds on DPD's:

$$f_{ab} + h_{\delta a \delta b} - h_{\delta a \delta b}^t \pm \sqrt{(h_{\delta a b} + h_{a \delta b})^2 + (f_{\Delta a \Delta b} - h_{\delta a \delta b} - h_{\delta a \delta b}^t)^2} \geq 0$$

$$f_{ab} - h_{\delta a \delta b} + h_{\delta a \delta b}^t \pm \sqrt{(h_{\delta a b} - h_{a \delta b})^2 + (f_{\Delta a \Delta b} + h_{\delta a \delta b} + h_{\delta a \delta b}^t)^2} \geq 0$$

( M. Diehl, TK, 2013)

- Bounds can aid construction of double parton distributions
- Sets an upper limit on the size of spin correlations

- How are the bounds affected by evolution?
- For single parton in polarized proton (Soffer bound), first argued to be broken by evolution but later shown to hold

(C. Bourrely, J. Soffer and O. Teryaev, 1998)

- Evolution of unpolarized double parton distribution ( $x_1$  parton)

$$\frac{\partial}{\partial \tau_1} \left( \text{green oval with } x_1 \text{ and } x_2 \right) = \left( \text{green oval with } x_1 \text{ and } x_2 \text{ and loop on } x_1 \right) + \left( \text{green oval with } x_1 \text{ and } x_2 \text{ and loop on } x_2 \right)$$

$$\frac{\partial f_{qq}(x_1, x_2, \mathbf{y}; \mu_1, \mu_2)}{\partial \log \mu_1^2} \sim P_{qq} \otimes_1 f_{qq} + P_{qg} \otimes_1 f_{gq},$$

where the  $P_{ab}$ 's are the usual splitting kernels (without inhomogeneous term)

- Analogous equations for polarized partons
- Bounds stable under leading order evolution

(M.Diehl, TK, 2013)

# Numerical solution of evolution equations

- Numerically study the evolution
- Evolution code:
  - Modified version of evolution code written by J. Gaunt and J. Stirling
- Generally evolution washes out spin correlations

(J. Gaunt, J. Stirling 2009)

# Numerical solution of evolution equations

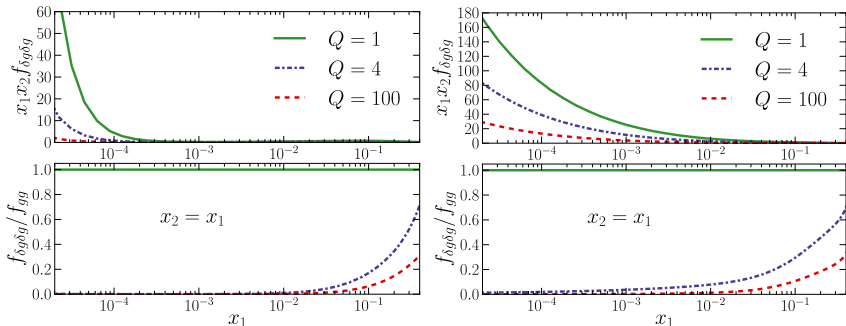
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## Need ansatz for DPDs at starting scale

- Build on products of single parton distributions,  $f_{ab} \sim f_a f_b$
- Ansatz saturating the positivity bounds:
  - polarized distributions equal the unpolarized at starting scale

- Linearly polarized gluons ( $f_{\delta g \delta g}$ )
- Starting from leading order distributions from MSTW (left) or GJR (right)

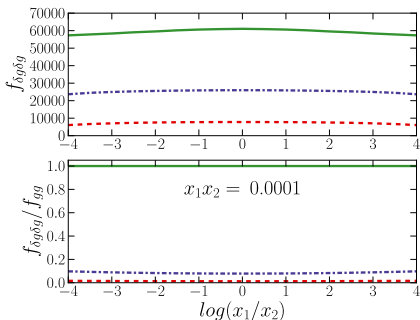
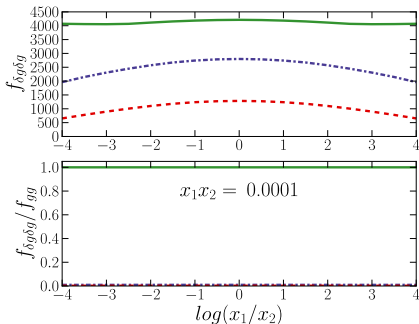
(A. Martin, W. Stirling, R. Thorne, G. Watt 2009, M. Gluck, P. Jimenez-Delgado, E. Reya 2007)



- Quickly washed out by evolution

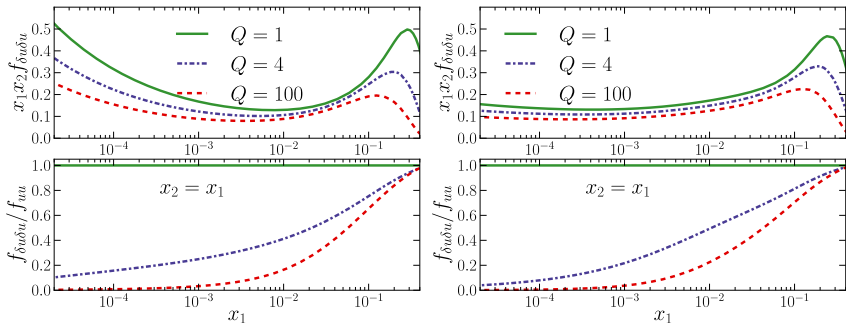


- $f_{\delta g \delta g}$  at fixed  $x_1 x_2$



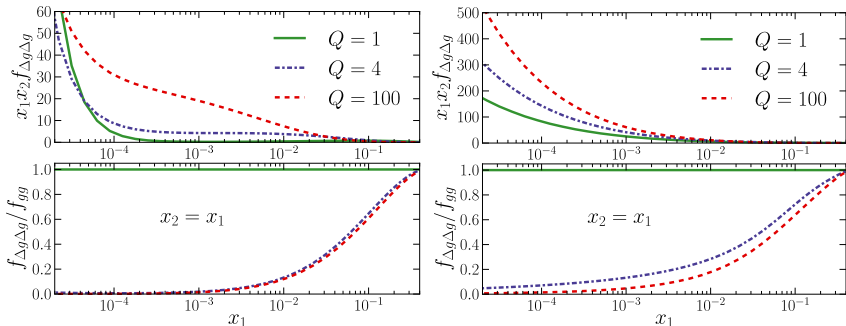
- $\log x_1/x_2 \sim$  rapidity difference between interactions
- Quickly washed out  $\Rightarrow$  negligible already at medium scales
- Differences depending on starting single pdfs

- Transversely polarized up-quarks ( $f_{\delta u \delta u}$ )



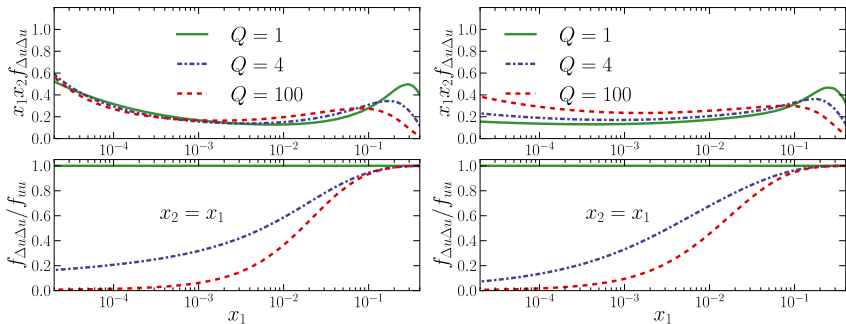
- Can be significant up to larger scales
- Smaller differences between MSTW and GJR

- Longitudinally polarized gluons ( $f_{\Delta g \Delta g}$ )



- Stays sizable up to larger scales
- Differences depending on starting single pdfs

- Longitudinally polarized up-quarks ( $f_{\Delta u \Delta u}$ )



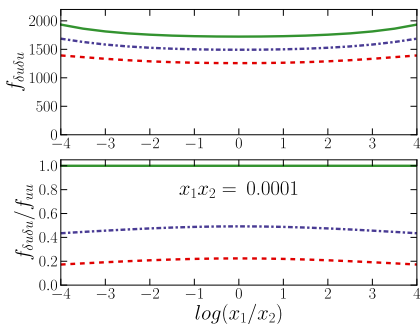
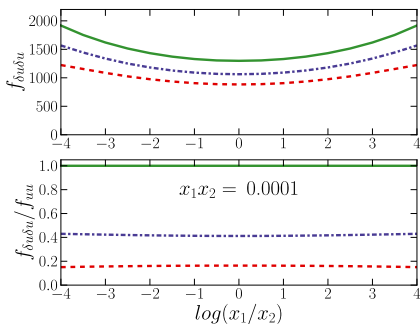
- Remains large up to larger scales

## Summary

- Many unexplored effects in double parton scatterings
- **Spin correlations** have profound effects on cross sections
- We derived **positivity bounds** on distributions
- Bounds are **stable under** leading order **evolution**
- Bounds set **upper limit on spin correlations** between partons in the proton, at any scale
- Linear gluon polarization quickly washed out by evolution
- **Transverse quark polarization** significant up to **high scales**
- **Longitudinal quark and gluon** polarization significant up to **high scales**

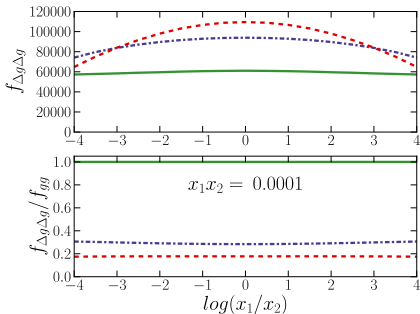
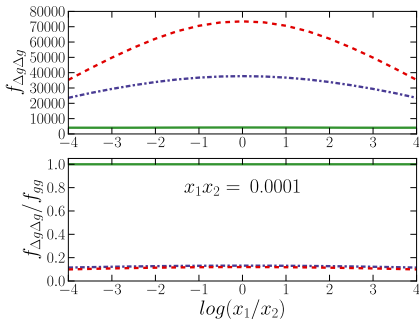
# Backup Slides

- $f_{\delta u \delta u}$  at fixed  $x_1 x_2$



- Can be significant up to larger scales
- Small differences between MSTW and GJR

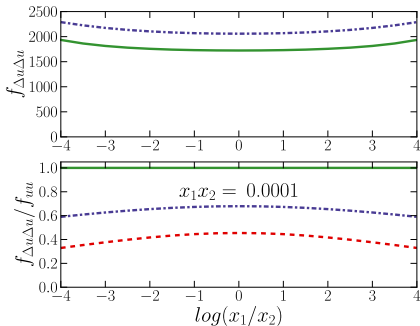
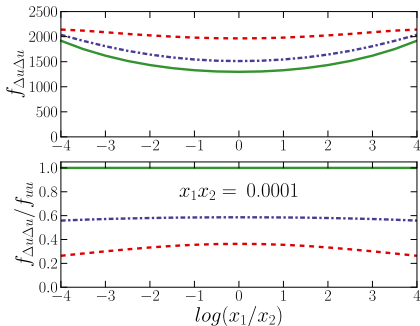
- $f_{\Delta g \Delta g}$  at fixed  $x_1 x_2$



- Stays sizable up to larger scales
- Differences depending on starting single pdfs



- $f_{\Delta u \Delta u}$  at fixed  $x_1 x_2$



- Remains large up to larger scales

## Angles and coordinate systems

