

# Measurements of Quarkonium Production and Polarization at CMS

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on behalf of the  
**CMS Collaboration**

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**HEPHY**  
Institute of High Energy Physics

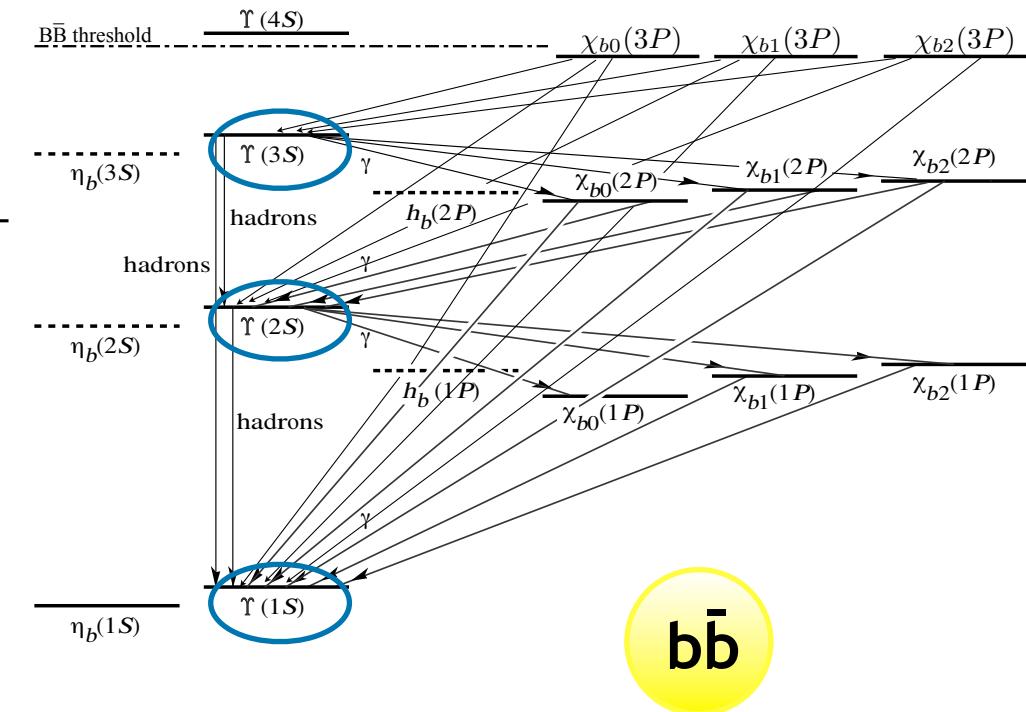
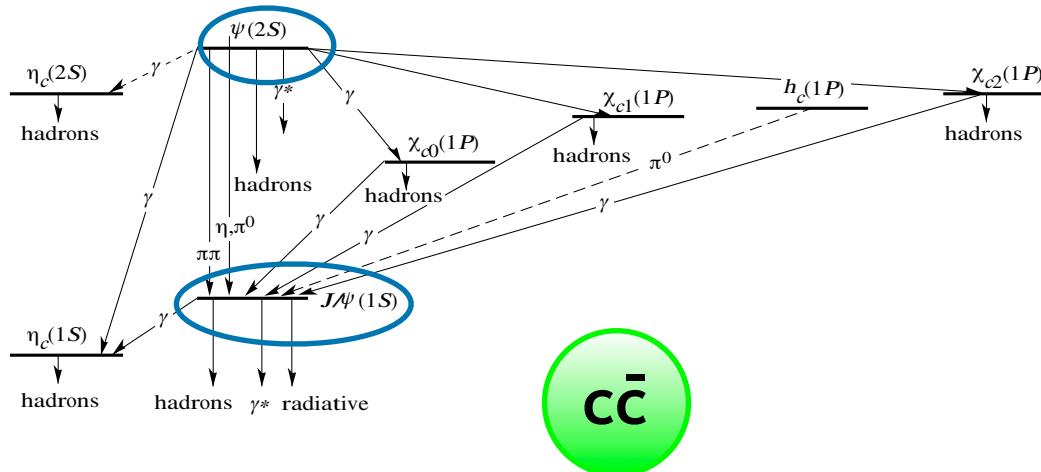


\* supported by Austrian Science Fund (FWF): P24167

# Outline

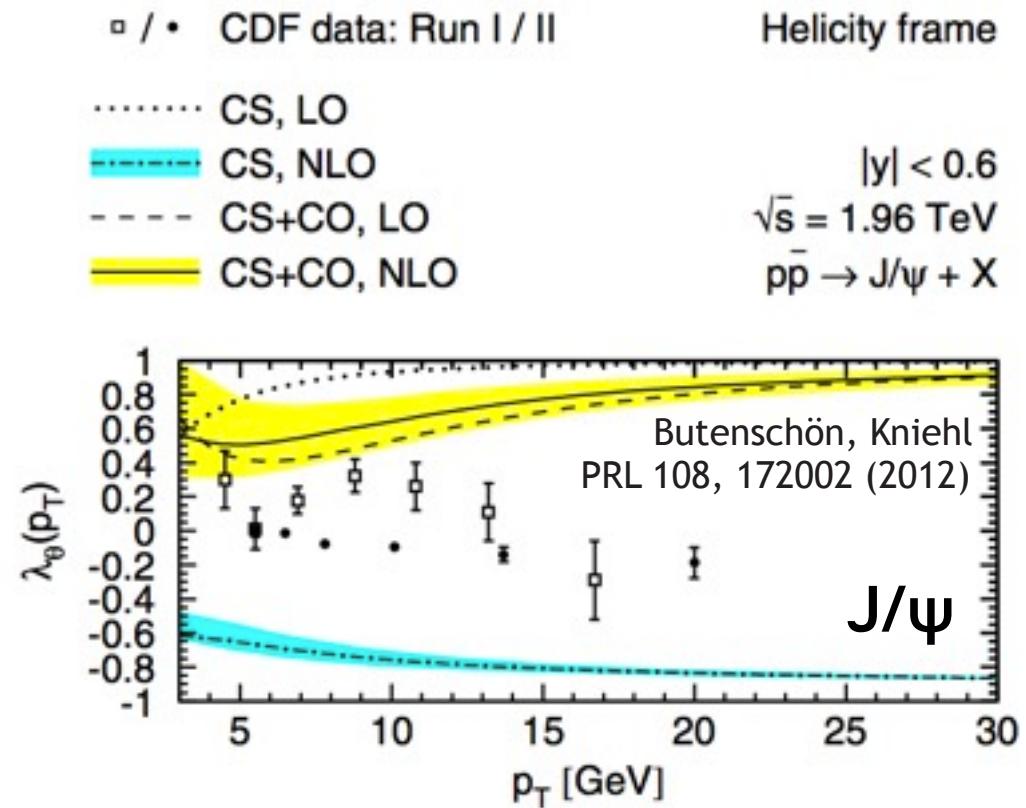
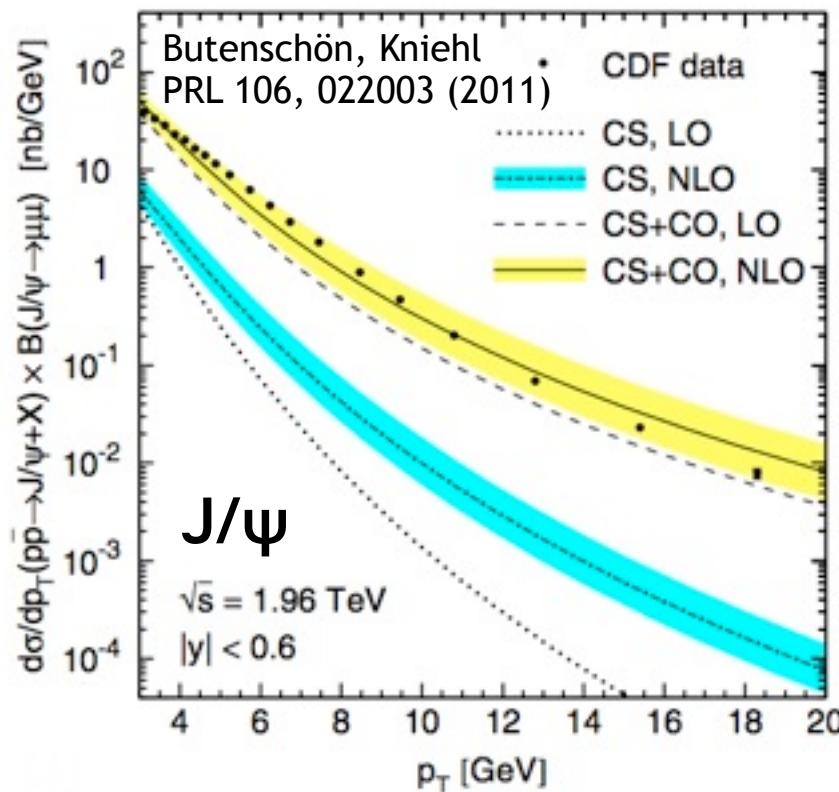
- $\Upsilon(nS)$  cross section
- $\Upsilon(nS)$  polarization
- Prompt  $J/\psi$  polarization *First time for CMS*
- Prompt  $\psi(2S)$  polarization

All results are based on a dimuon sample collected in pp collisions in 2011 at  $\sqrt{s} = 7$  TeV, corresponding to a total integrated luminosity of  $4.9 \text{ fb}^{-1}$



# Motivation

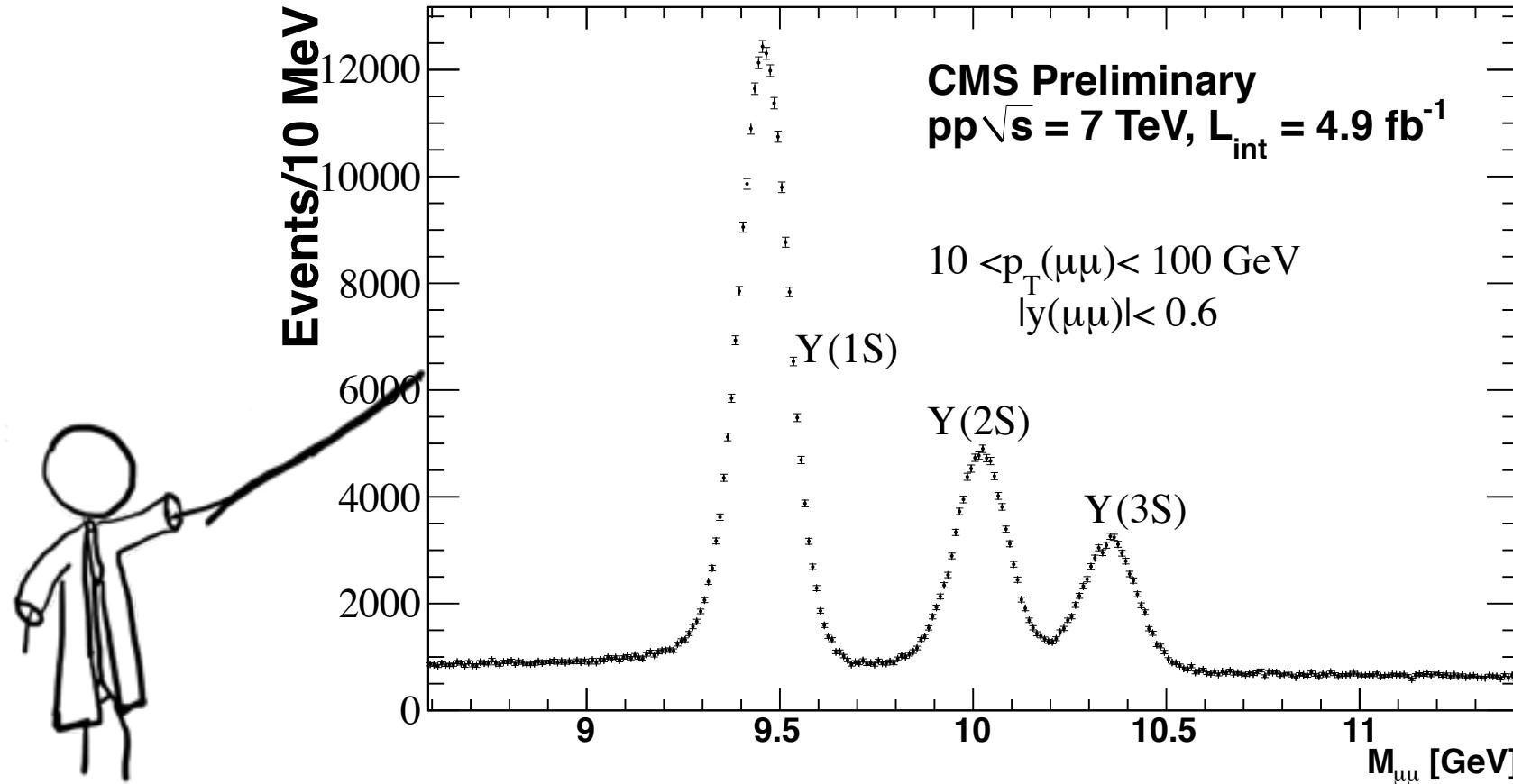
- No theory has simultaneously explained experimental measurements of both quarkonium production and polarization
- Most previous polarization measurements only determined one out of three polarization parameters
- CMS has a higher reach in transverse momentum,  $p_T$



# $\Upsilon(nS)$ Cross Section Analysis

- Uses results from CMS  $\Upsilon(nS)$  polarization measurement for the calculation of the acceptance
- Kinematic range of measurement:  $10 < p_T < 100 \text{ GeV}$ ,  $|y| < 0.6$

Details in CMS-PAS-BPH-12-006



# $\Upsilon(nS)$ Differential Cross Section

- $\Upsilon(nS)$  differential cross section  $\frac{d\sigma}{dp_T}$  times dimuon branching ratio  $\mathcal{B}$  integrated over  $|y| < 0.6$  in a given  $p_T$  bin of width  $\Delta p_T$

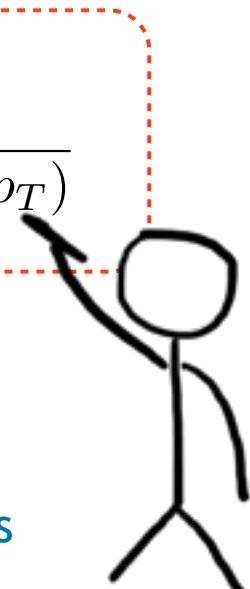
$$\frac{d\sigma(pp \rightarrow \Upsilon(nS))}{dp_T} \Big|_{y<0.6} \times \mathcal{B}(\Upsilon(nS) \rightarrow \mu^+ \mu^-) = \frac{N_{\Upsilon(nS)}^{fit}(p_T)}{L_{int} \cdot \Delta p_T \cdot \varepsilon(p_T) \cdot \mathcal{A}(p_T)}$$

$N_{\Upsilon(nS)}^{fit}(p_T)$  number of  $\Upsilon(nS)$  events in a  $p_T$  bin of width  $\Delta p_T$

$L_{int}$  integrated luminosity ( $4.9 \text{ fb}^{-1}$ )

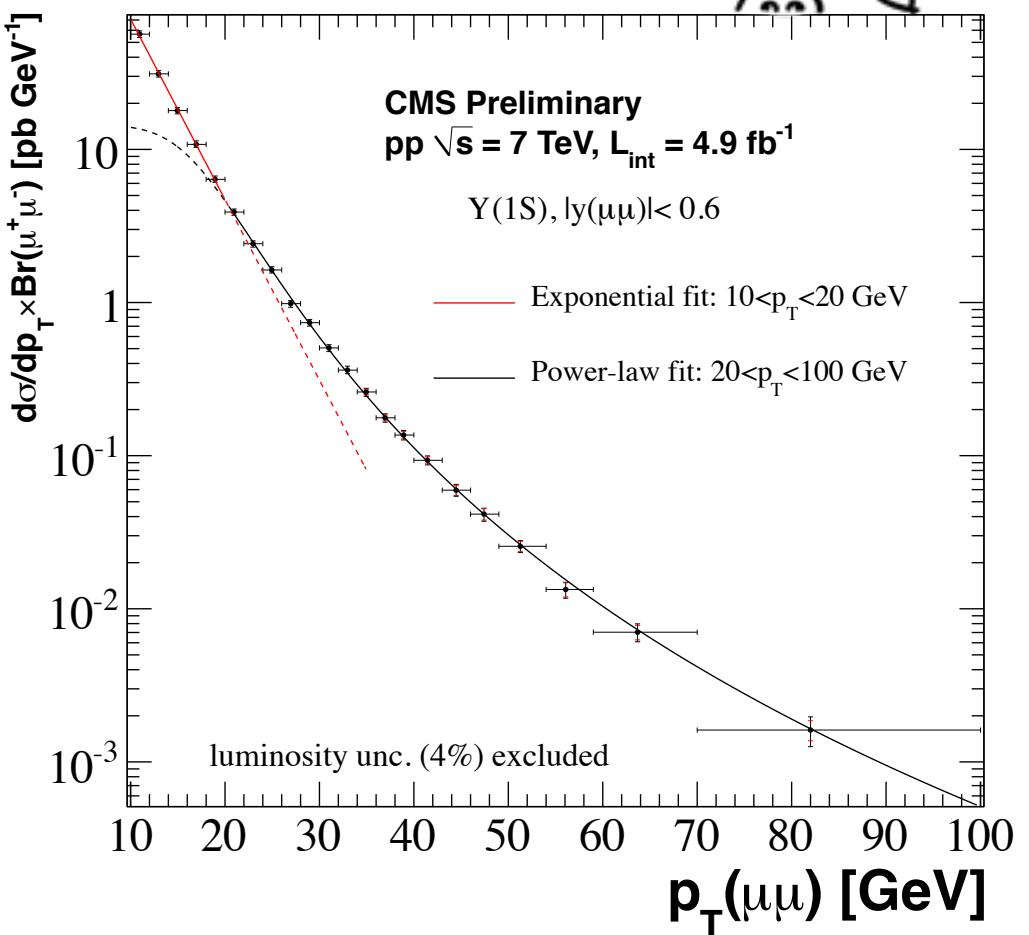
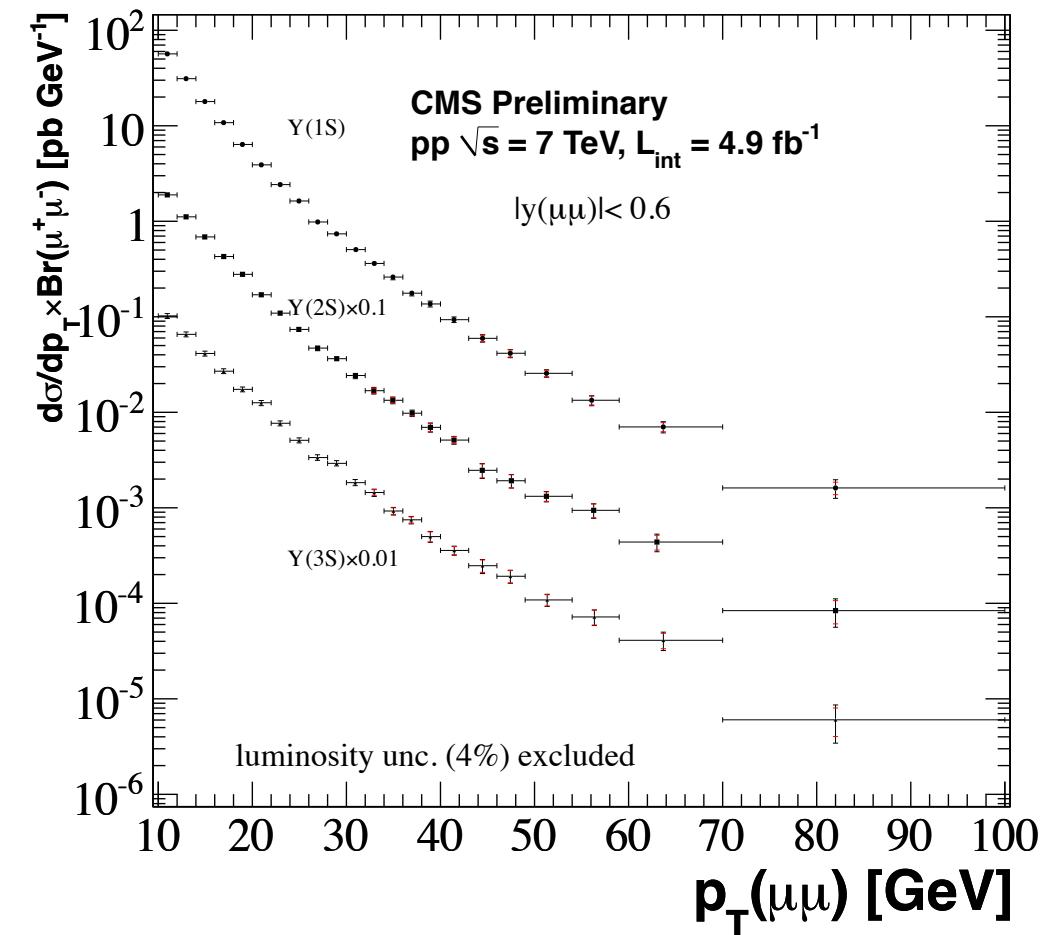
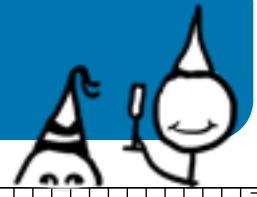
$\varepsilon(p_T)$  efficiency of trigger, reconstruction and analysis selections

$\mathcal{A}(p_T)$  acceptance calculated from Monte Carlo



- Acceptance is the polarization-weighted fraction of  $\Upsilon$  decays where the muons satisfy the kinematic requirements to the total of weighted events in a given  $p_T, y$  bin

# $\Upsilon(nS)$ Differential Cross Section



- Similar behaviour for all three  $\Upsilon$  states
- Change of slope for  $p_T > 20 \text{ GeV}$  suggests a change in the nature of the production process

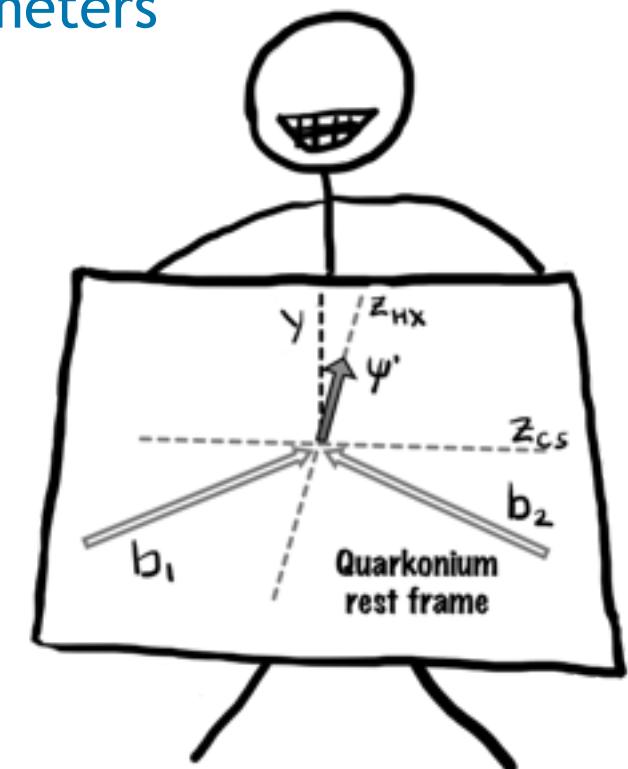
# Quarkonium Polarization

- Polarization is measured through the average angular decay distribution - for vector particles most generally written as

$$W(\cos \vartheta, \varphi | \vec{\lambda}) = \frac{3/(4\pi)}{(3 + \lambda_\vartheta)} (1 + \lambda_\vartheta \cos^2 \vartheta + \lambda_\varphi \sin^2 \vartheta \cos 2\varphi + \lambda_{\vartheta\varphi} \sin 2\vartheta \cos \varphi)$$

where  $\lambda_\vartheta, \lambda_\varphi, \lambda_{\vartheta\varphi}$  are the polarization parameters

- Angular decay distribution is measured with respect to a certain reference frame
  - center-of-mass helicity HX (polar axis  $z_{HX}$   
 $\approx$  direction of quarkonium momentum)
  - Collins-Soper CS ( $z_{CS} \approx$  direction of relative velocity of colliding particles)
  - perpendicular helicity PX ( $z_{PX} \perp z_{CS}$ )



# Need to Measure the Full Angular Distribution

- Two extreme angular decay distributions

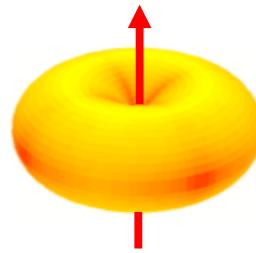
Longitudinal polarization

$$J_z = 0$$

$$\lambda_\vartheta = -1$$

$$\lambda_\varphi = 0$$

$$\lambda_{\vartheta\varphi} = 0$$



Transverse polarization

$$J_z = \pm 1$$

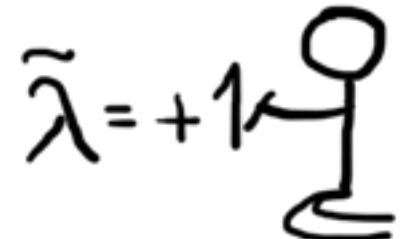
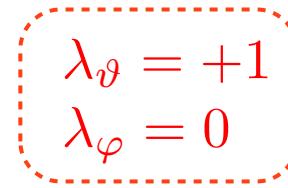
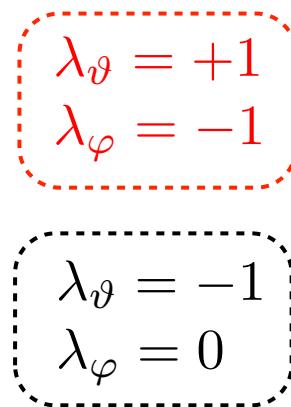
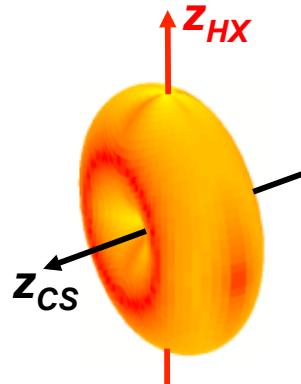
$$\lambda_\vartheta = +1$$

$$\lambda_\varphi = 0$$

$$\lambda_{\vartheta\varphi} = 0$$



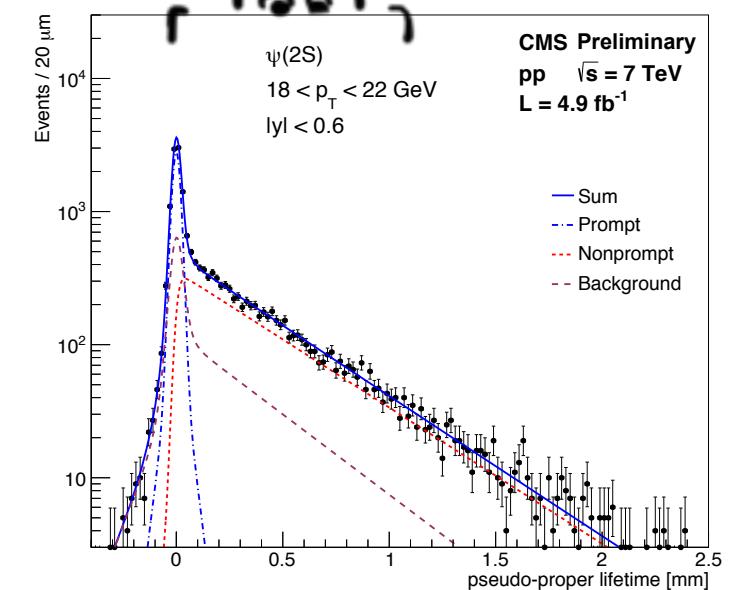
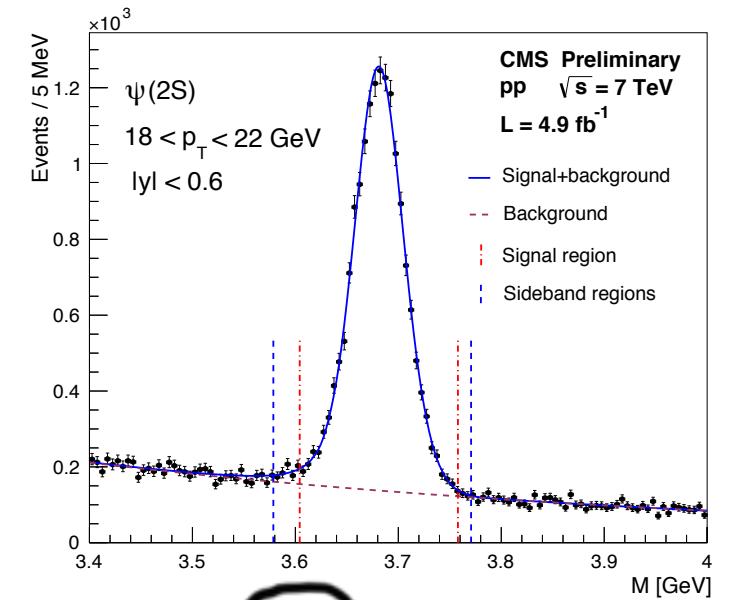
- The full angular distribution has to be measured. Otherwise two very different physical cases cannot be distinguished.
- The shape of the distribution is invariant and can be characterized by the frame invariant parameter  $\tilde{\lambda} = (\lambda_\vartheta + 3\lambda_\varphi)/(1 - \lambda_\varphi)$



# Quarkonium Polarization Measurements

- $\lambda_\theta, \lambda_\varphi, \lambda_{\theta\varphi}$  and  $\tilde{\lambda}$  are measured in three different reference frames (HX, CS, PX) for  $J/\psi$ ,  $\psi(2S)$ ,  $\Upsilon(1S)$ ,  $\Upsilon(2S)$  and  $\Upsilon(3S)$  mesons
- As a function of transverse momentum  $p_T$ 
  - $J/\psi$ :  $14 < p_T < 70$  GeV (10 bins)
  - $\psi(2S)$ :  $14 < p_T < 50$  GeV (4 bins)
  - $\Upsilon(nS)$ :  $10 < p_T < 50$  GeV (5 bins)
- And dimuon rapidity,  $|y|$ 
  - $J/\psi, \Upsilon(nS)$ :  $|y| < 1.2$  (2 bins)
  - $\psi(2S)$ :  $|y| < 1.5$  (3 bins)
- An additional non prompt component has to be taken into account for  $\psi(nS)$  states

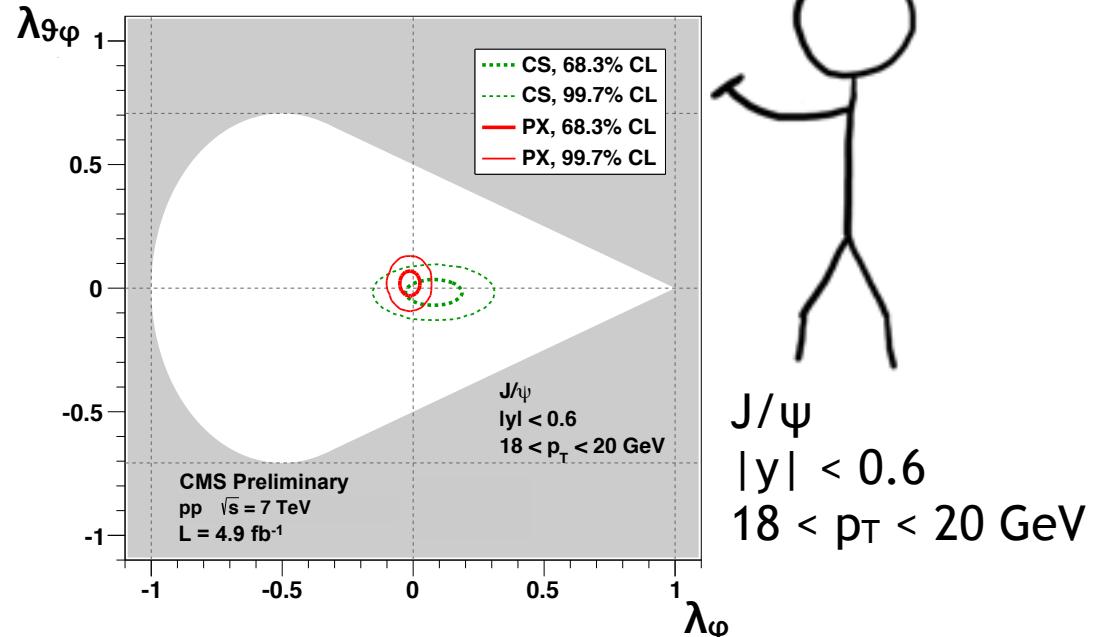
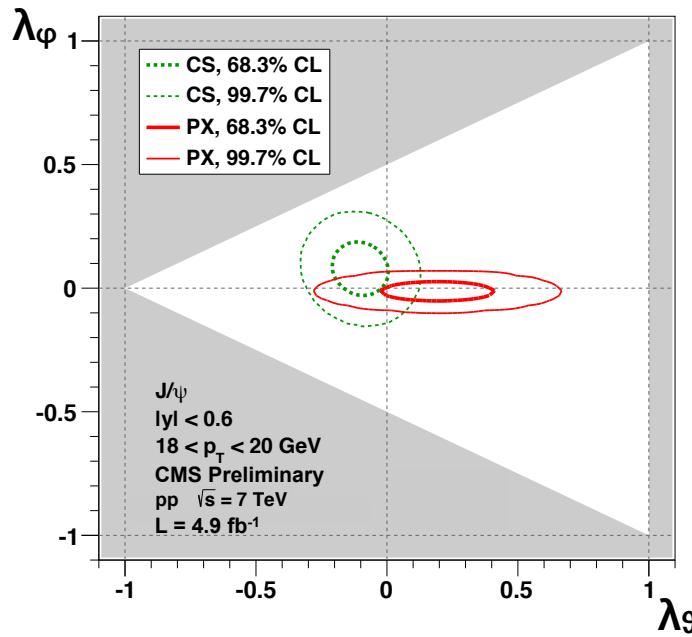
Details in PRL 110, 081802 (2013) and  
CMS-PAS-BPH-13-003



# Obtaining Polarization Parameters

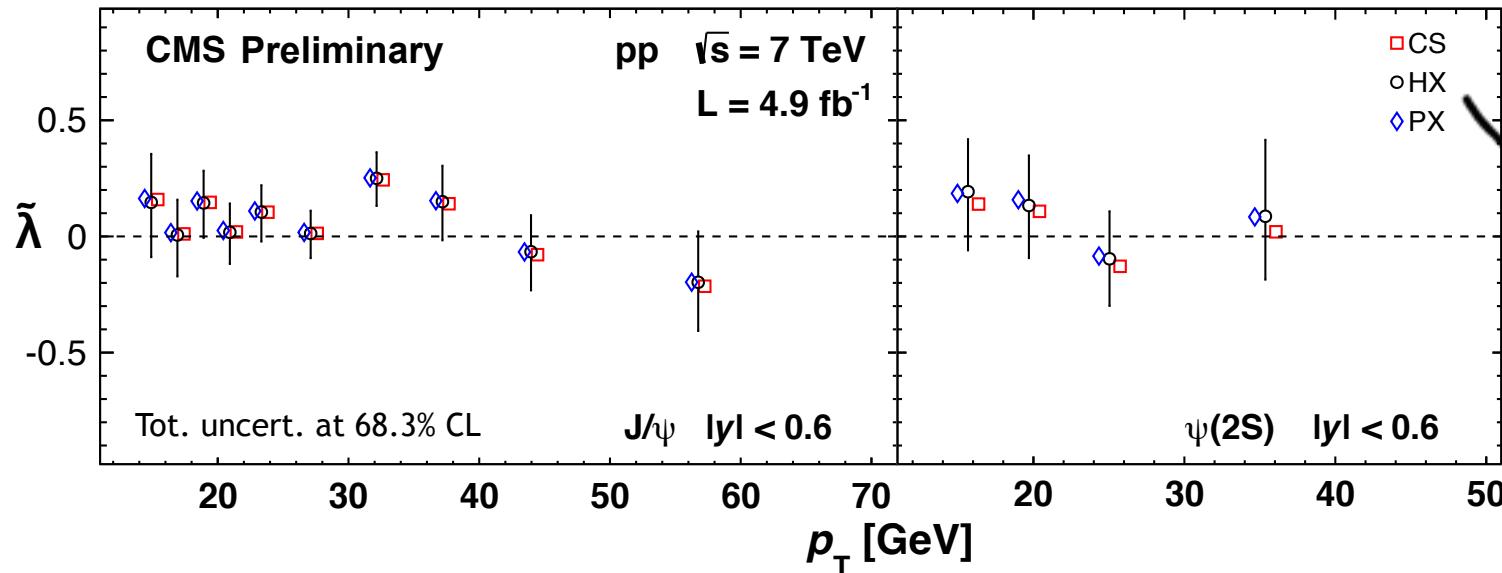
- Full and direct calculation of the Posterior Probability Distribution (PPD) of the polarization parameters  $\lambda_\theta$ ,  $\lambda_\varphi$ ,  $\lambda_{\theta\varphi}$

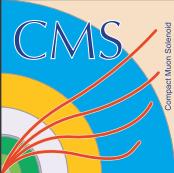
  1. Events distributed as in the background model are subtracted from the data sample until the previously determined background fraction is reached
  2. Definition of the PPD from the remaining signal-like events
  3. Numerical results and graphical representations are determined from 1D and 2D projections of the PPD



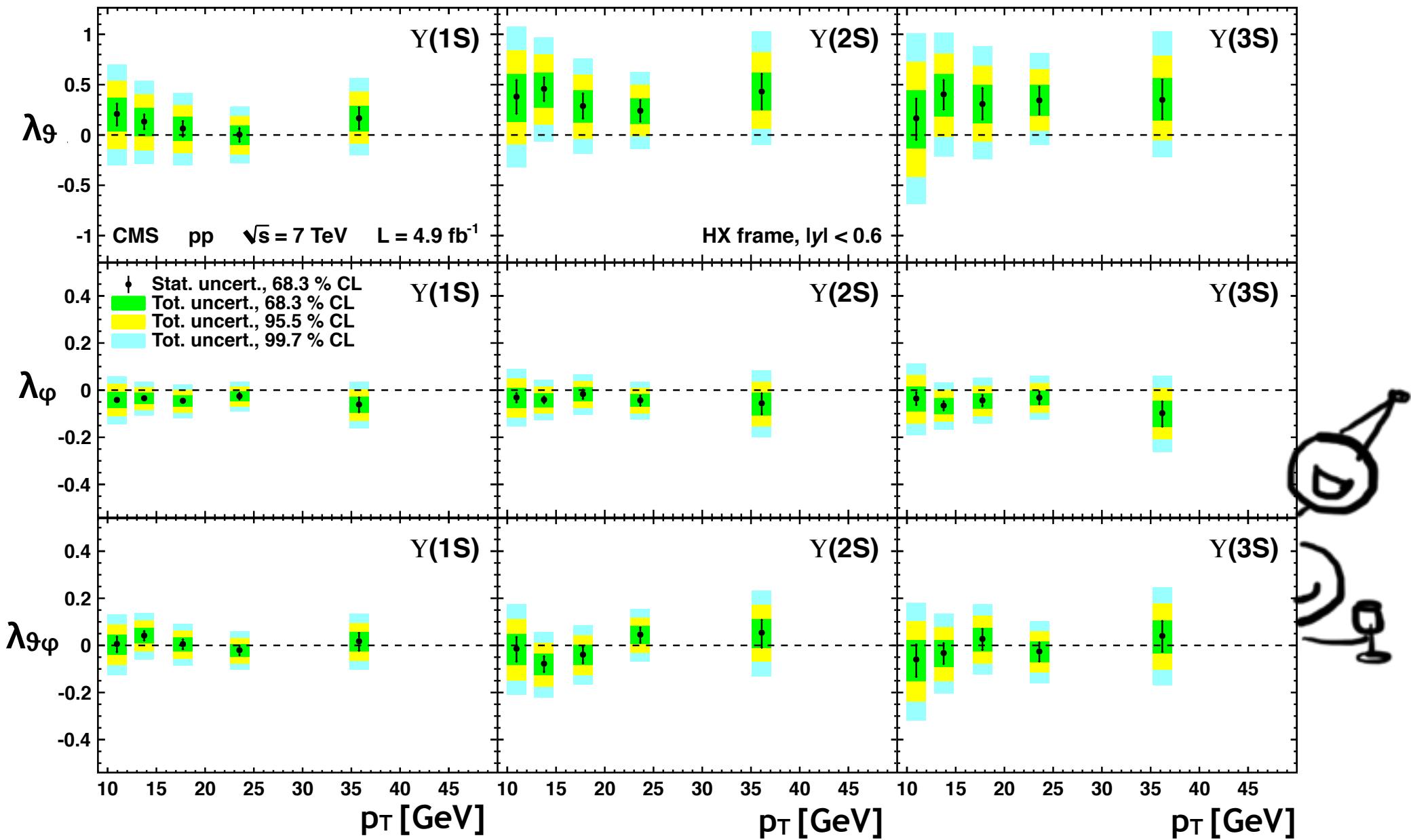
# Systematic Effects

- Sources of systematic effects:
  - Extraction of polarization parameters
  - Background model
  - Muon efficiencies
- Systematic uncertainties are propagated to the PPD
- Good agreement between the  $\tilde{\lambda}$  parameters in the three reference frames shows no indication of unaccounted systematic uncertainties

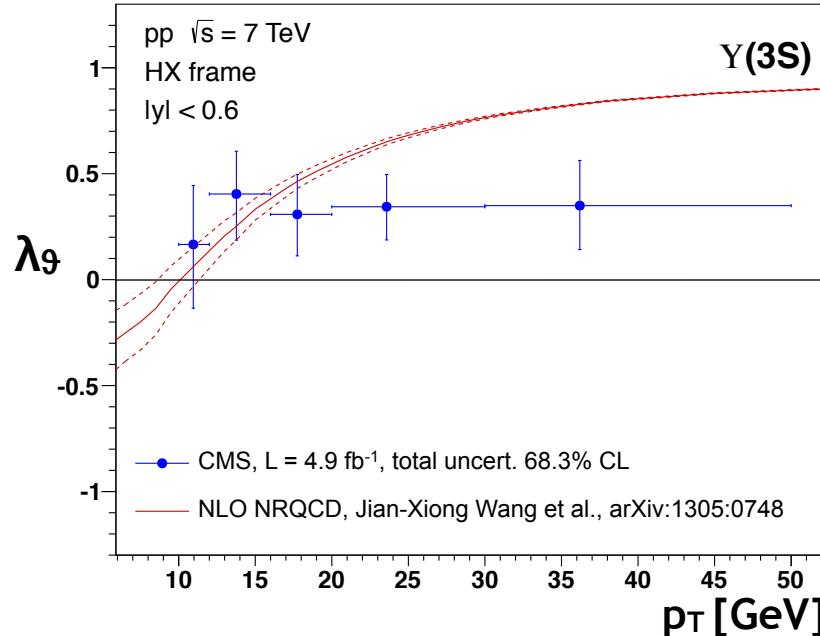
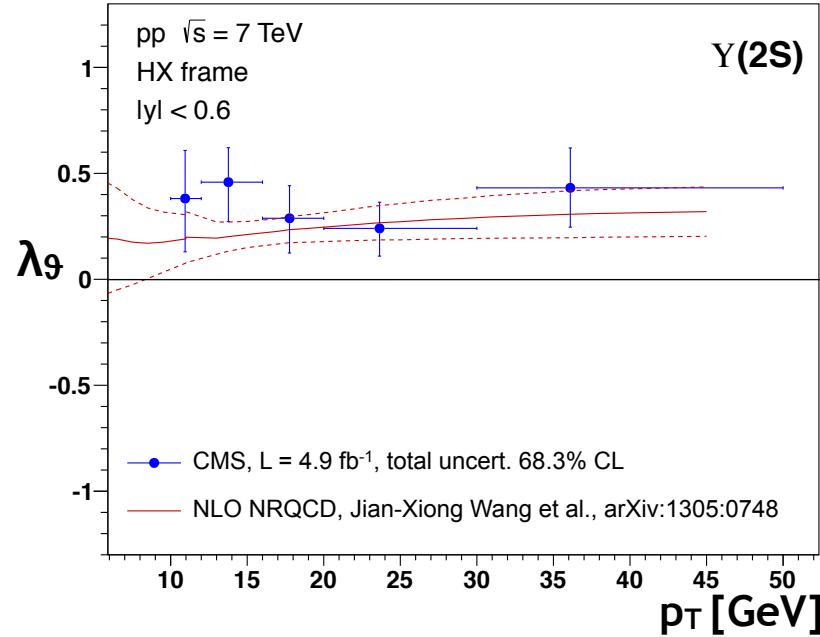
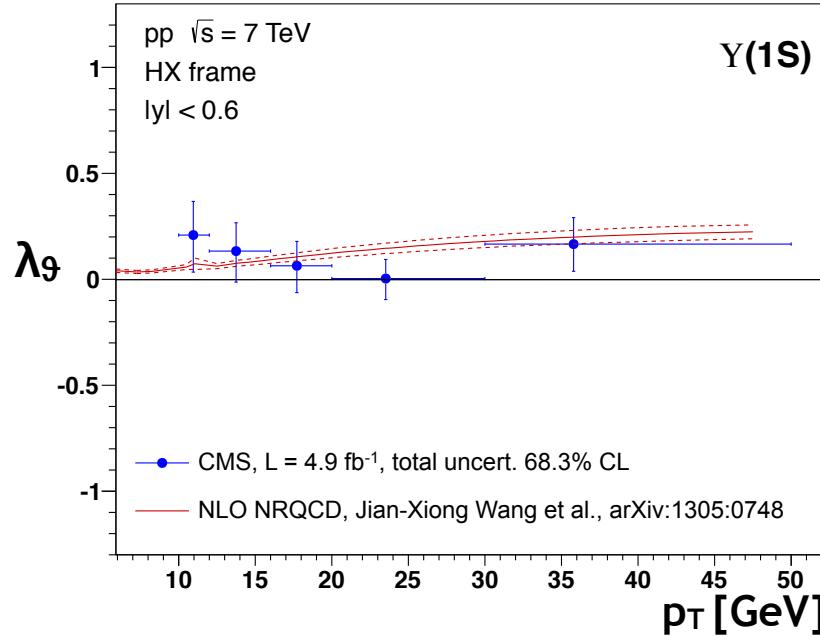




# $\Upsilon(nS)$ Polarization in the HX Frame, $|y| < 0.6$

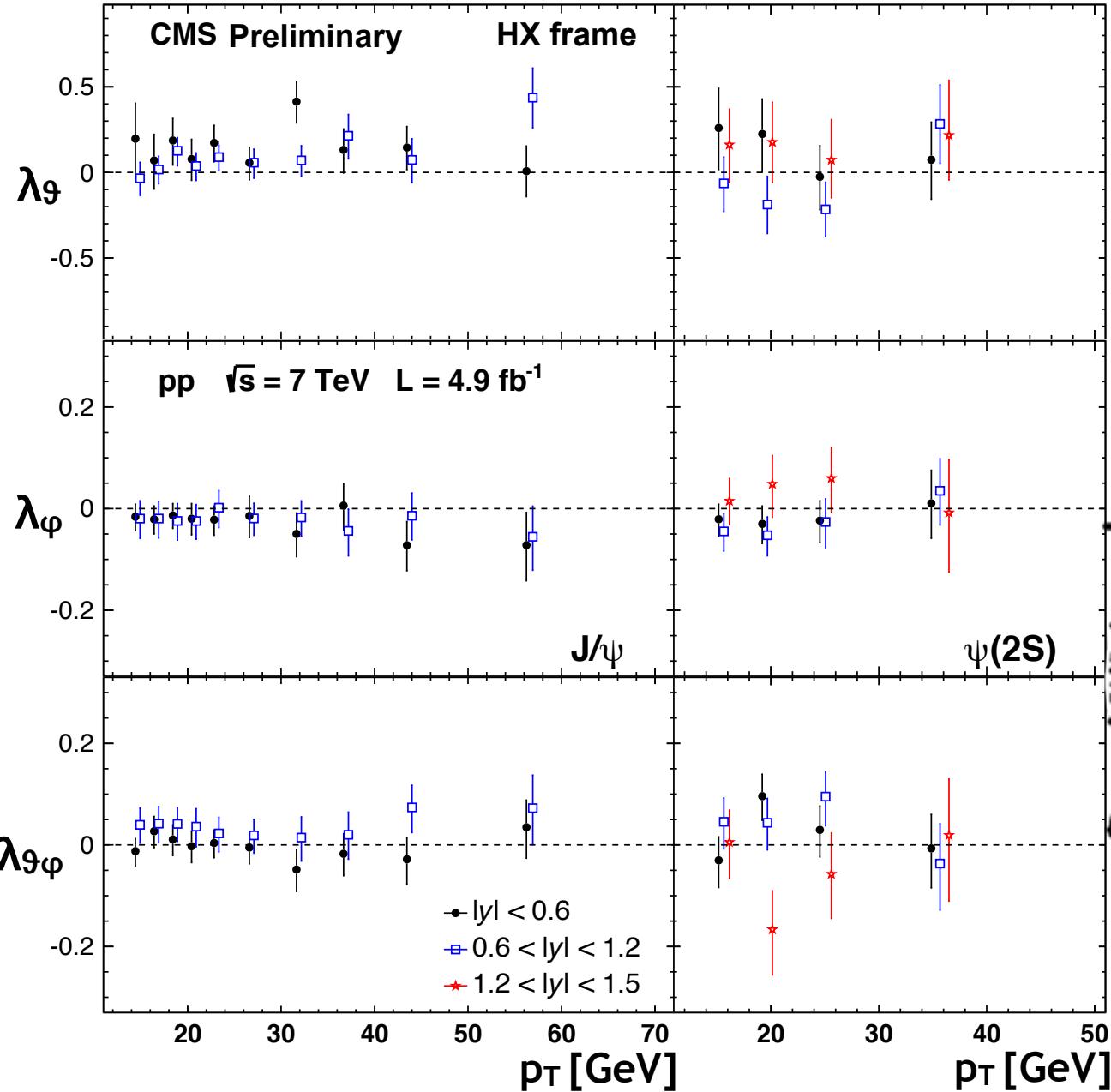


# $\Upsilon(nS)$ : Comparison to NLO NRQCD



- Theory calculation accounts for feed-down contributions to  $\Upsilon(1S)$  and  $\Upsilon(2S)$  states
- Prediction for  $\Upsilon(3S)$  may change when including feed-down from  $\chi_b$  (3P) states
- Color octet matrix elements are fit to hadroproduction data only

# Prompt $\psi(nS)$ Polarization in the HX Frame



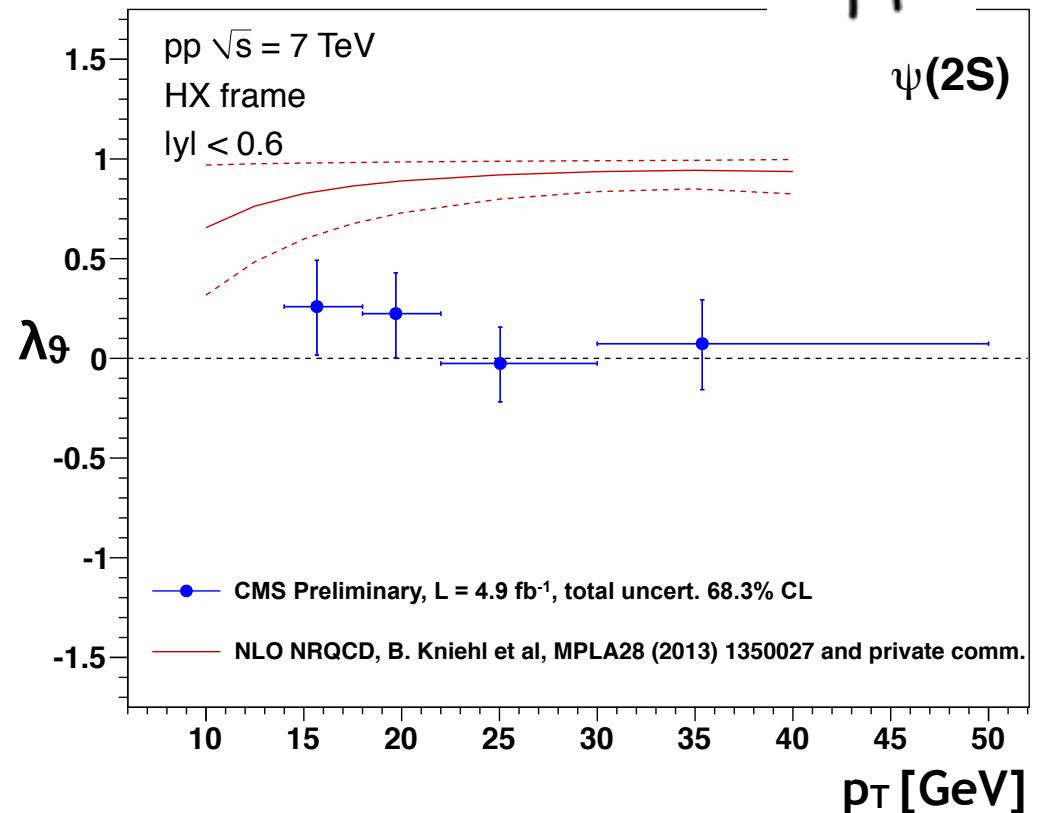
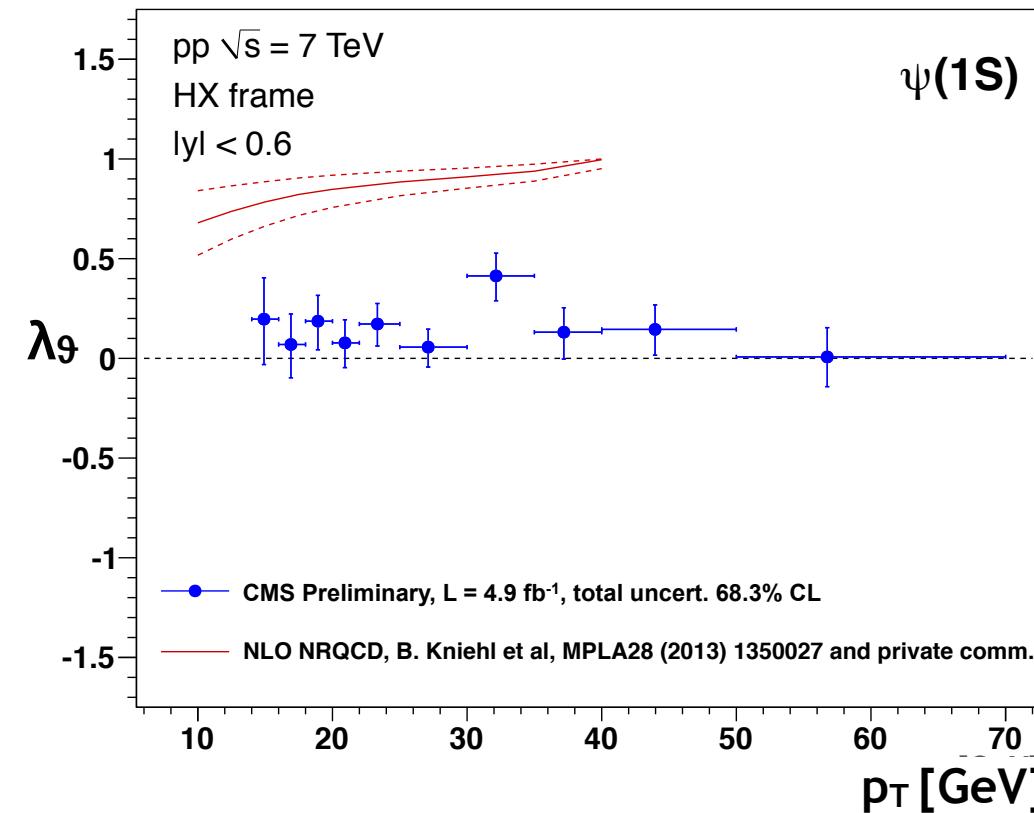
- No sign of strong polarization
- $\psi(2S)$  is not affected by feed-down decays from higher states



Error bars show total uncertainties at 68.3% CL

# $\psi(nS)$ : Comparison to NLO NRQCD

- CMS results disagree with the NLO NRQCD calculations
- Calculations use a global fit of color octet matrix elements to photo- as well as hadroproduction data
- Theory predicts polarization only for directly produced J/ $\psi$ 's



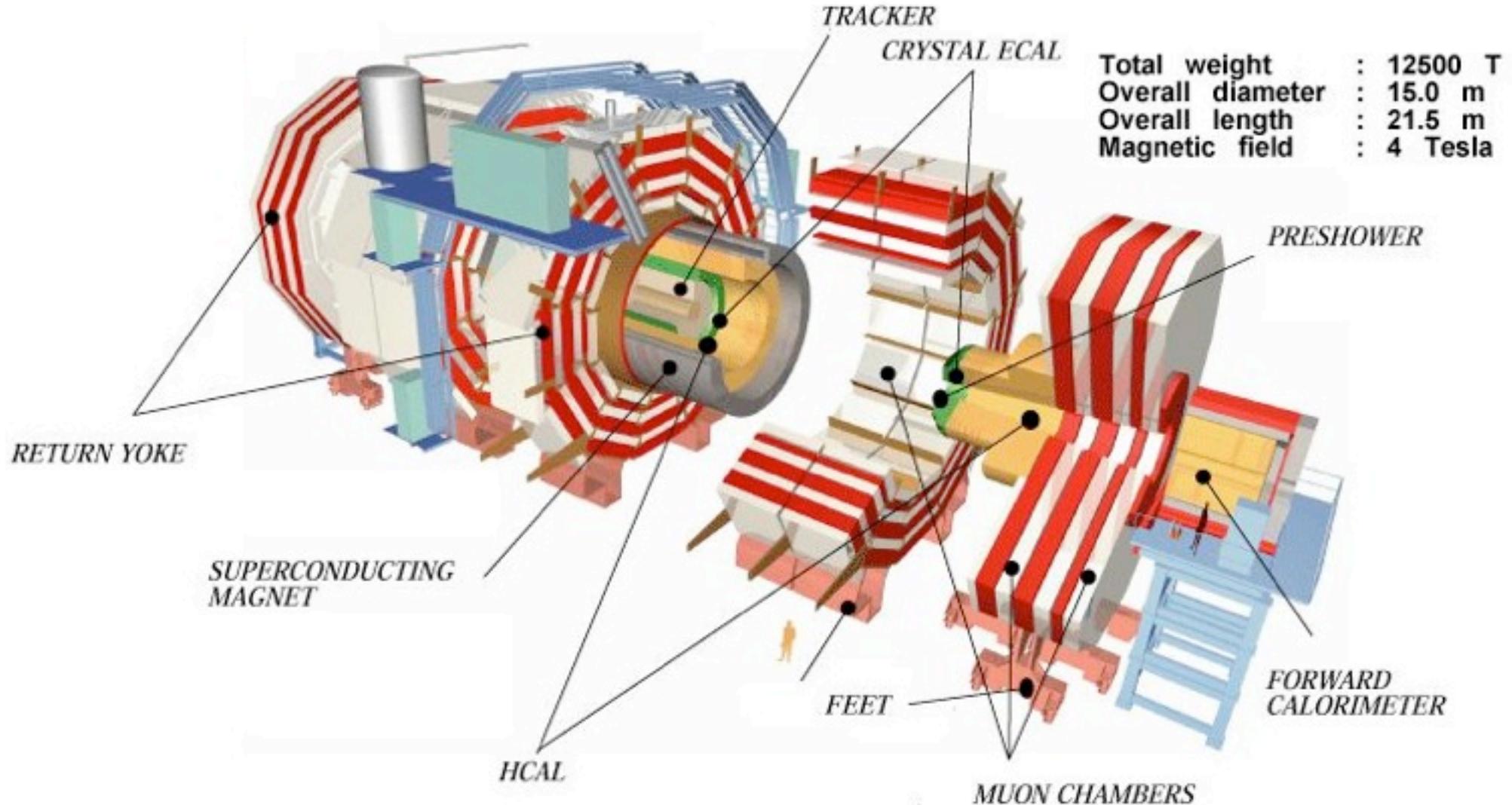
# Summary and Conclusions

- 🍩  $\Upsilon(1S)$ ,  $\Upsilon(2S)$  and  $\Upsilon(3S)$  differential cross sections for  $10 < p_T < 100$  GeV are measured using CMS polarization results
- 🍩 A change in the slope of the differential cross sections from an exponential to power-law is observed
- 🥜 Frame dependent polarization parameters  $\lambda_\theta$ ,  $\lambda_\phi$ ,  $\lambda_{\theta\phi}$  and the frame invariant parameter  $\tilde{\lambda}$  are measured in three different frames (CS, HX, PX) for the  $\Upsilon(1S)$ ,  $\Upsilon(2S)$ ,  $\Upsilon(3S)$  and prompt  $J/\psi$  and  $\psi(2S)$  mesons
- 🥜  $J/\psi$  results are shown for the first time
- 🥜 No evidence of strong longitudinal or transverse polarizations has been observed
- 🥜  $J/\psi$  and  $\psi(2S)$  measurements are in disagreement with current theoretical predictions



# BACKUP

# CMS Detector

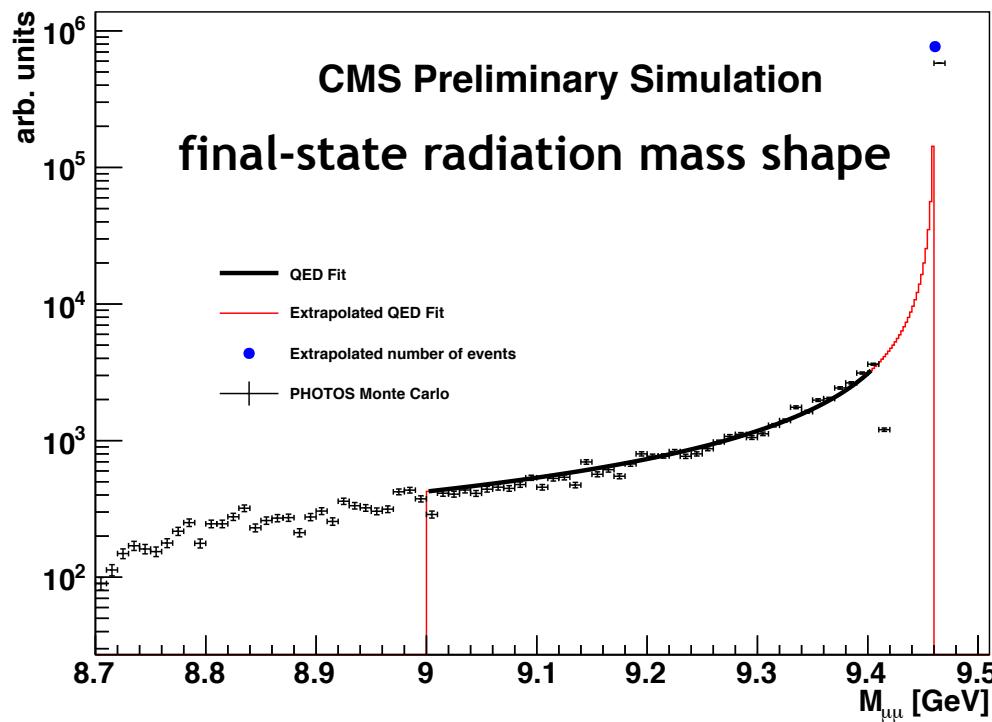


# Quarkonium Production

# Line Shape Determination

- Mass PDF was defined with the help of radiative line shape determined from Monte Carlo and the mass resolution

$$\mathcal{F}(M_{\mu\mu}; c_w, \delta m) = \frac{1}{N} \sum_{i=1}^N \frac{1}{\sqrt{2\pi} c_w \zeta_i} e^{-(M_{\mu\mu} - m_i - \delta m)^2 / 2c_w^2 \zeta_i^2}$$



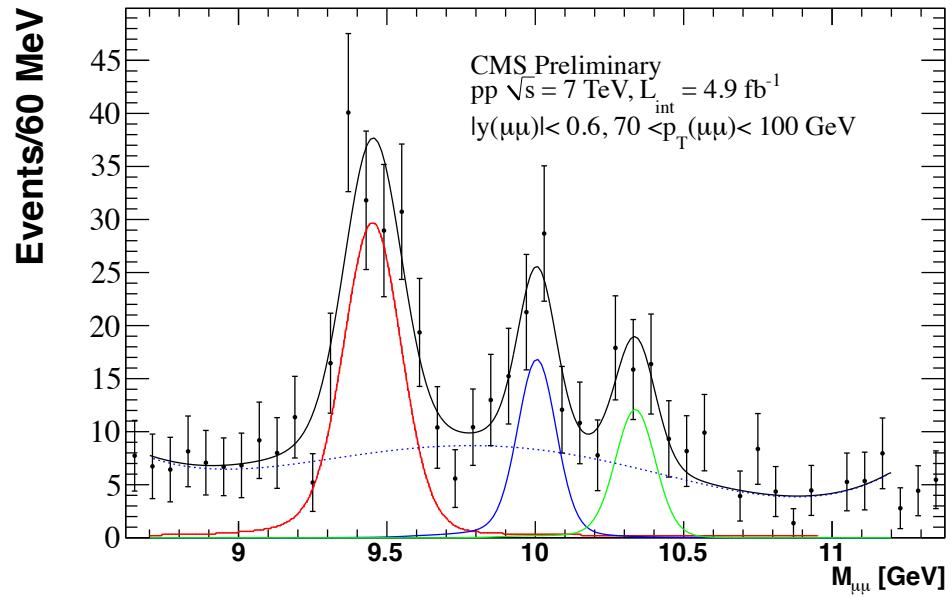
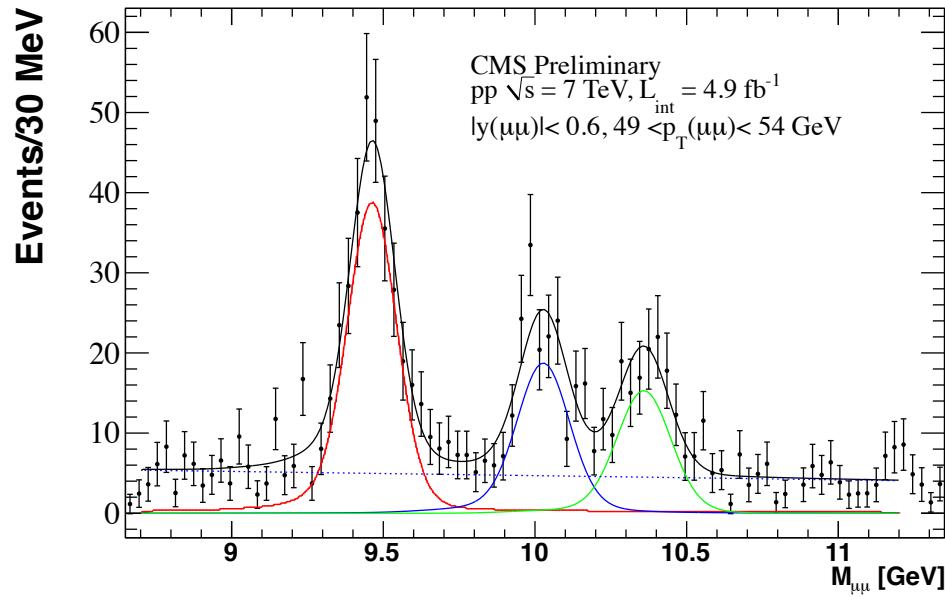
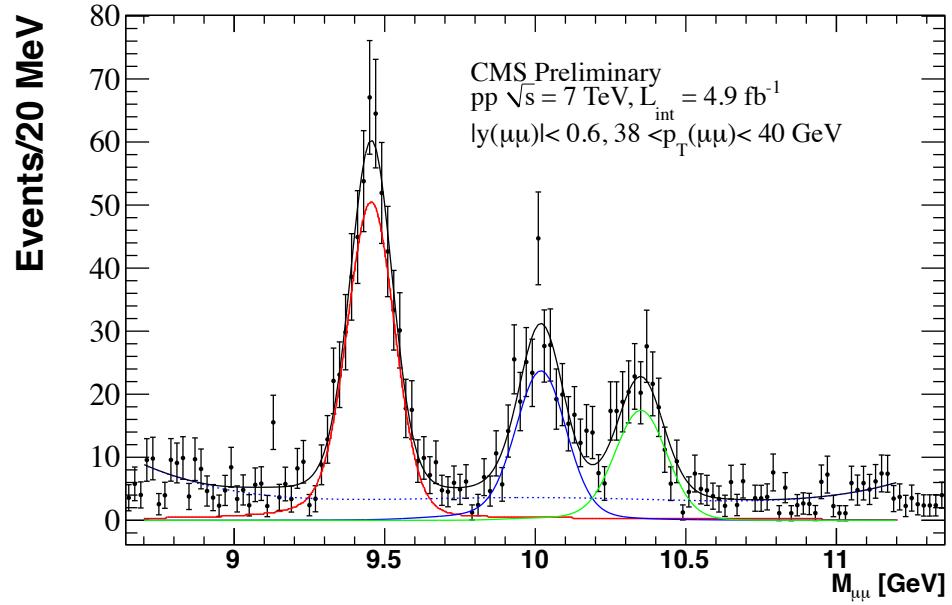
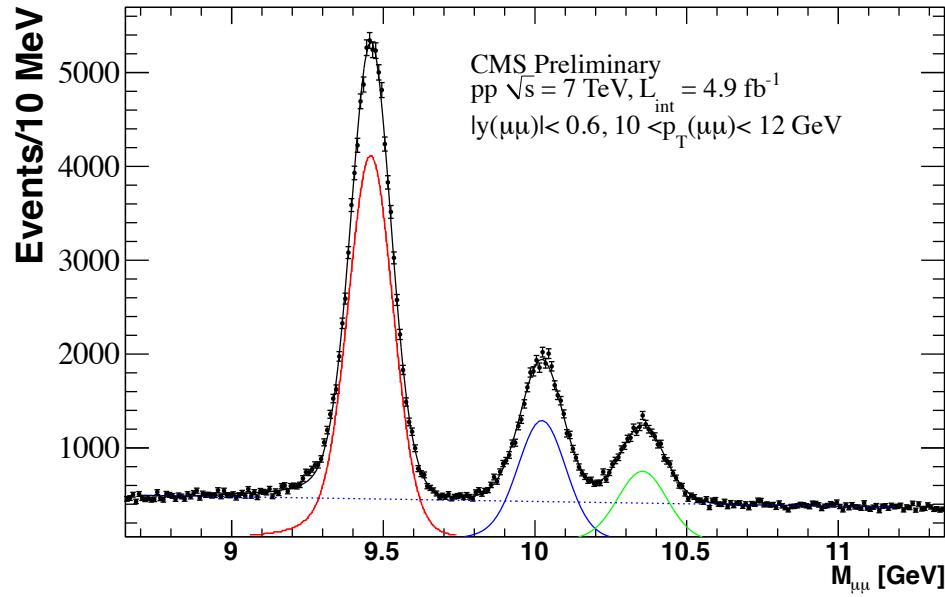
$c_w$  width scale factor

$\delta m$  mass shift scale factor

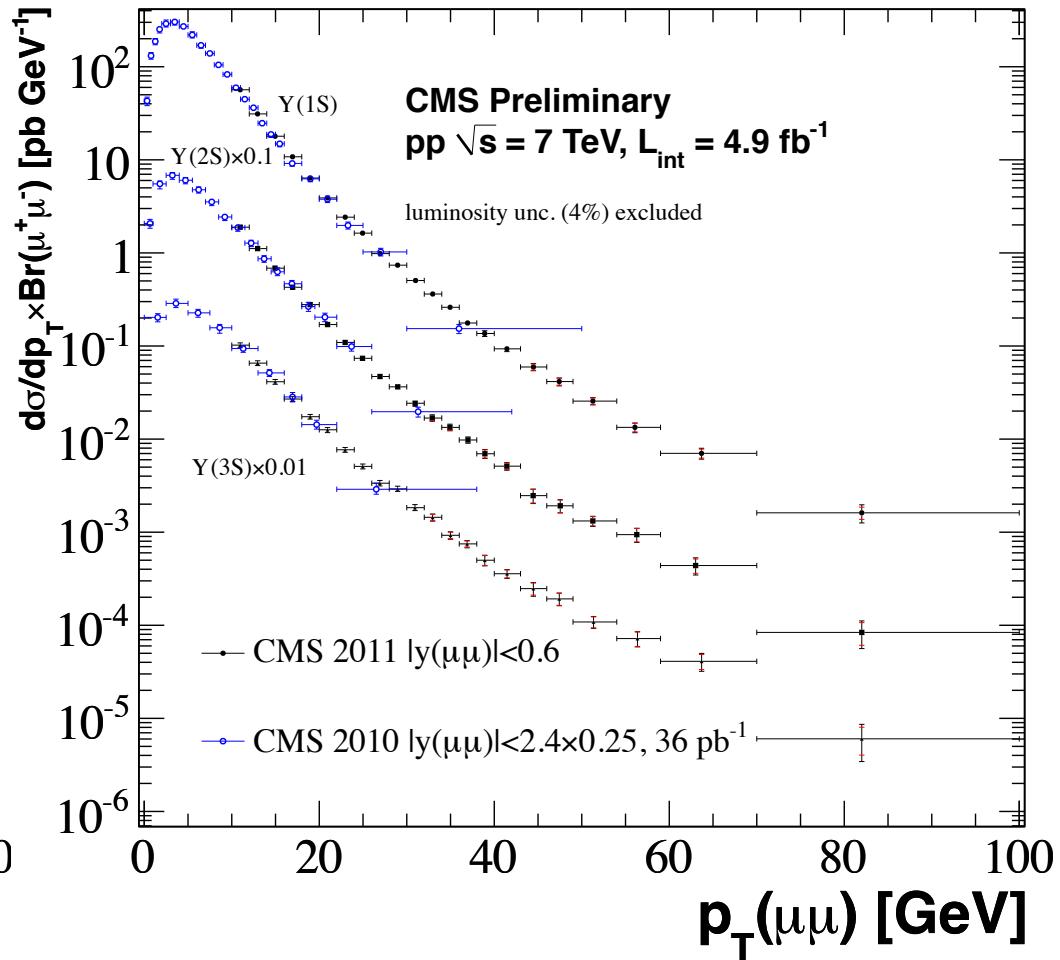
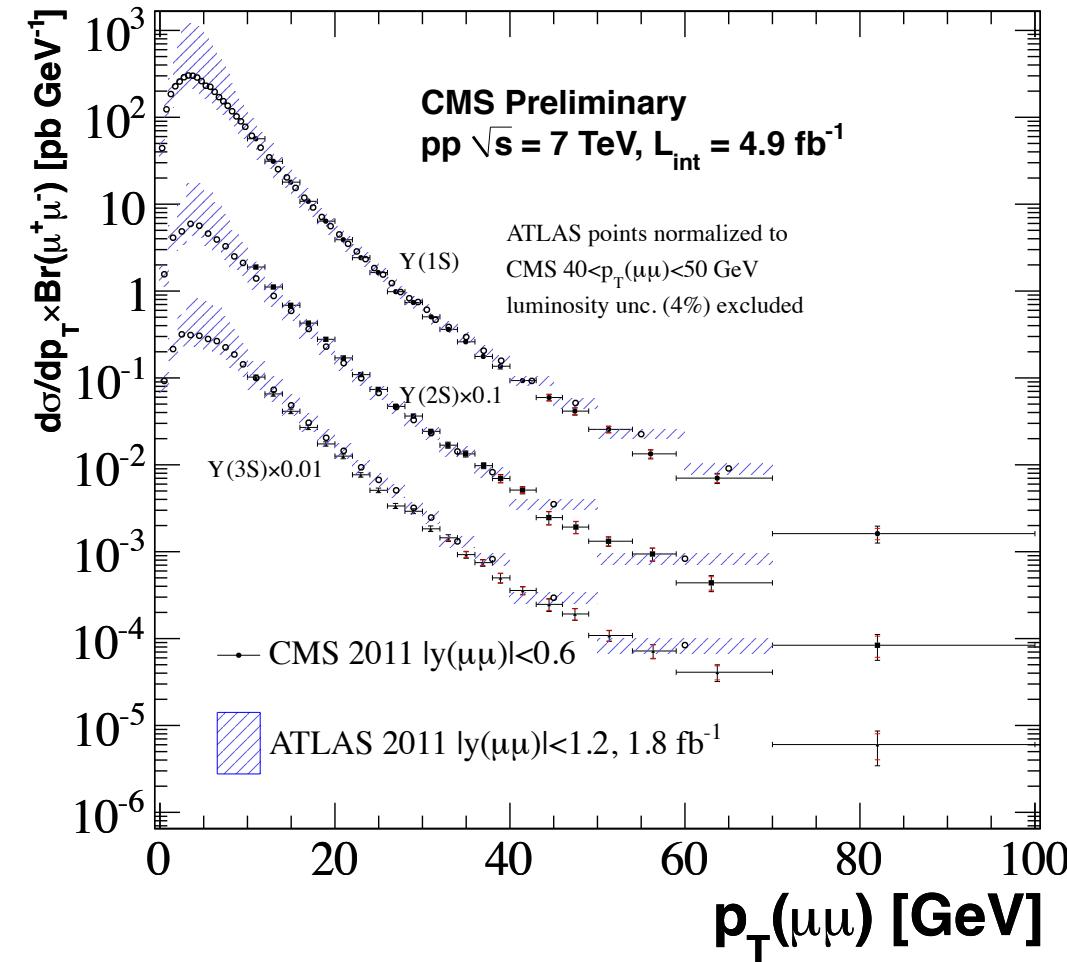
$m_i$  was sampled from the final-state radiation mass shape

$\zeta$  approximates the detector effects on the mass error

# Dimuon Mass Fits



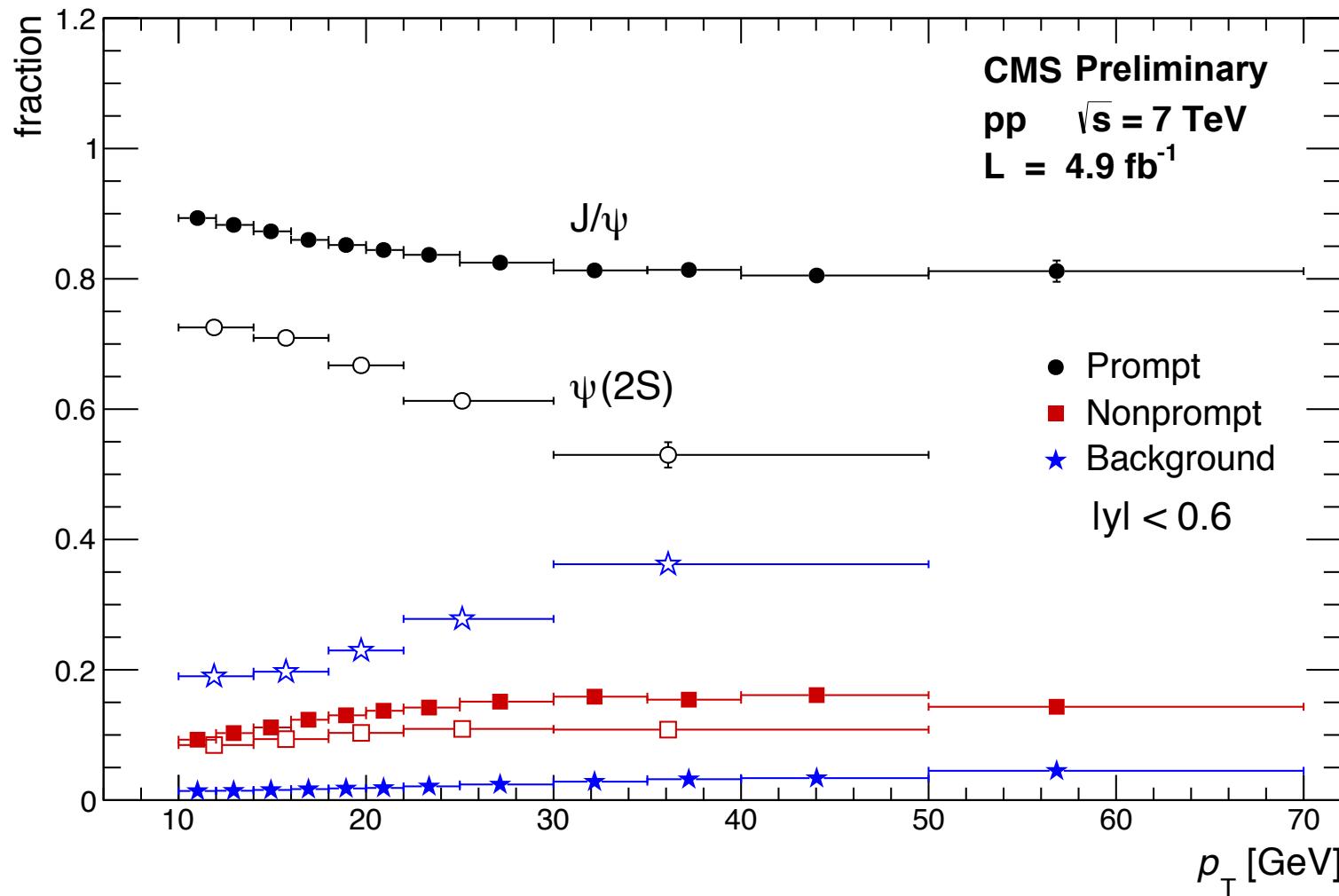
# Comparison to ATLAS and Previous CMS Measurements



# Quarkonium Polarization

# Contributions to the $\psi(nS)$ Prompt Signal Region

- The prompt-signal region is defined as a 2D window of  $\pm 3\sigma$  widths in dimuon mass and (pseudo-proper) lifetime



# Definition of the PPD

$$\mathcal{P}(\vec{\lambda}) \propto \prod_i \frac{1}{\mathcal{N}(\vec{\lambda})} W(\cos \theta^{(i)}, \phi^{(i)} | \vec{\lambda}) \varepsilon(p_1^{(i)} p_2^{(i)})$$

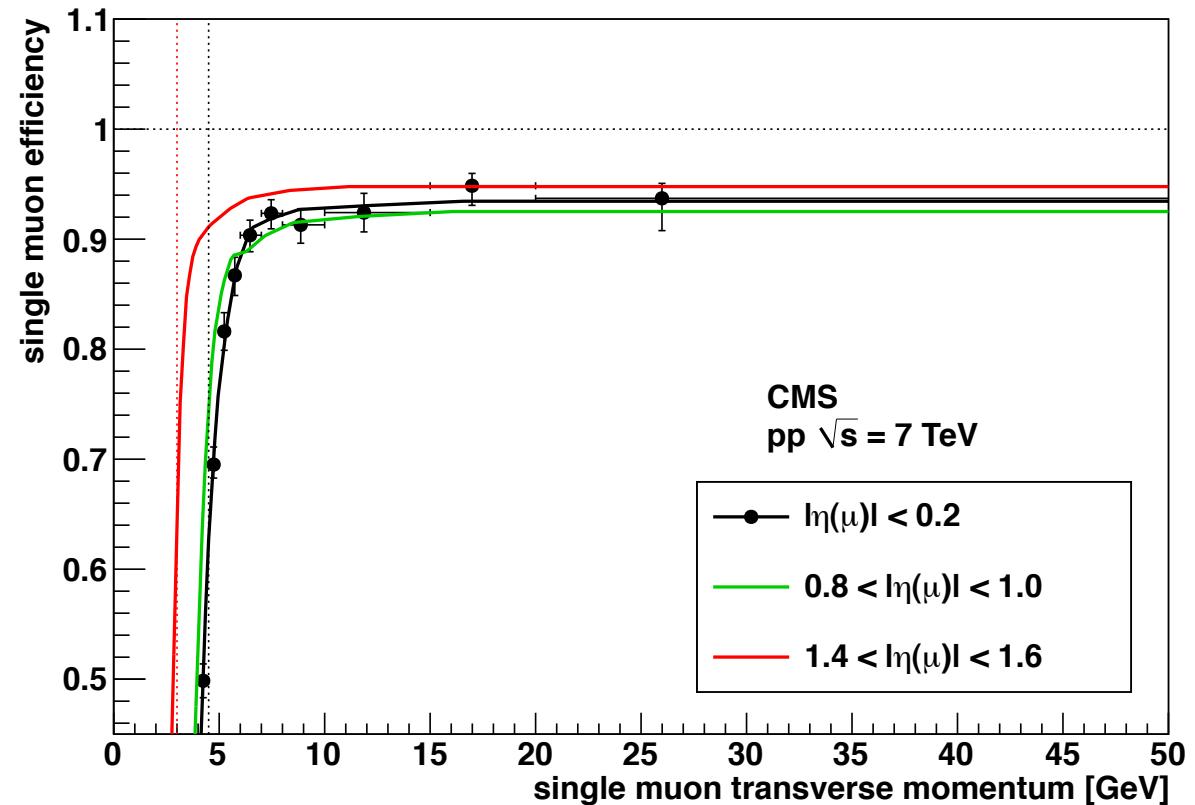
$\mathcal{N}$ : normalization

$W$ : general angular distribution

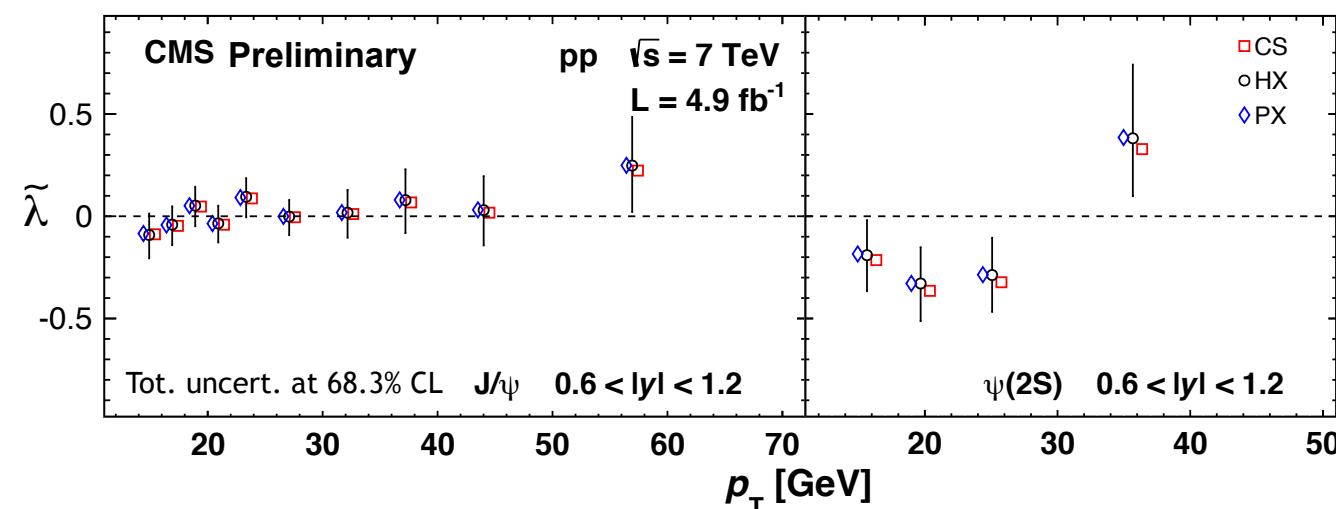
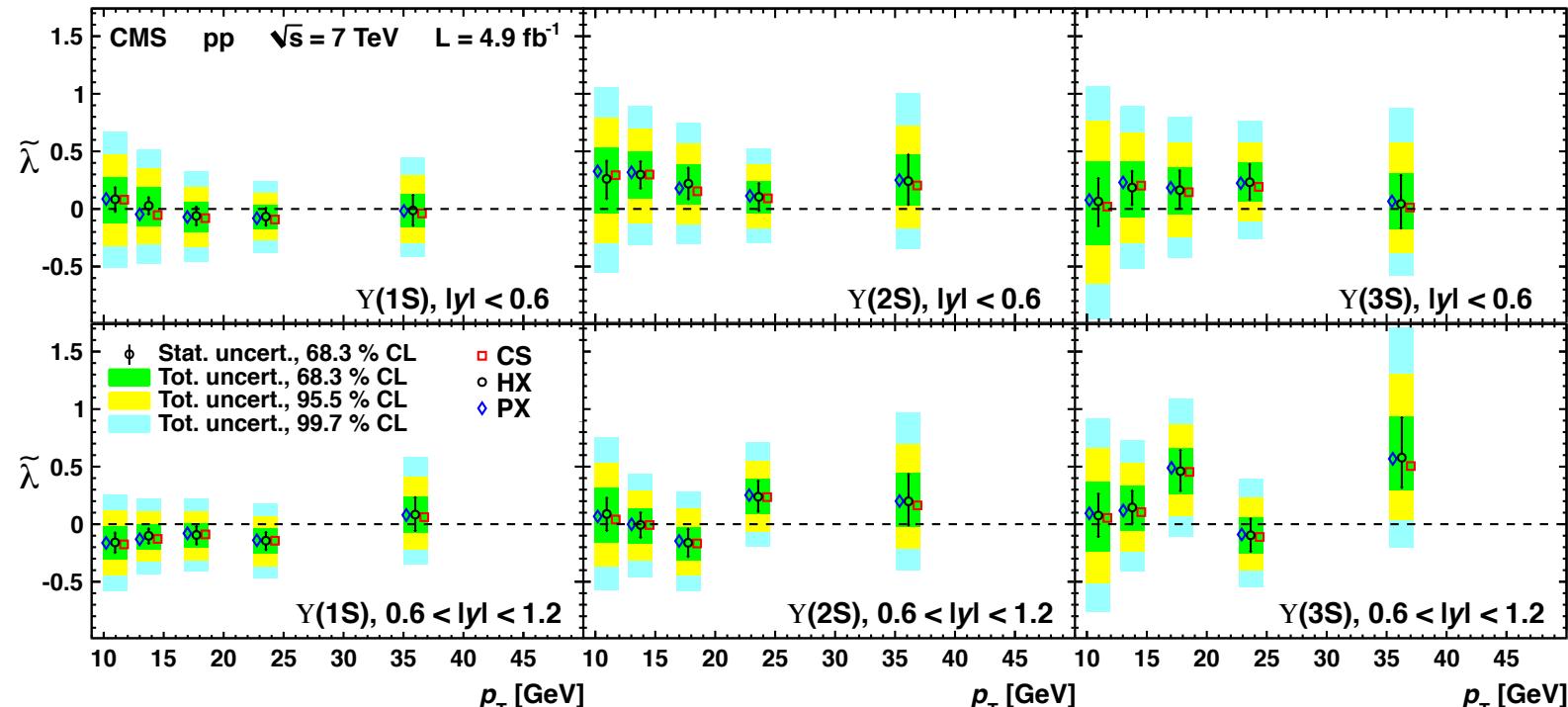
$\varepsilon$ : dimuon efficiency as a function of the muon momenta

# Efficiencies

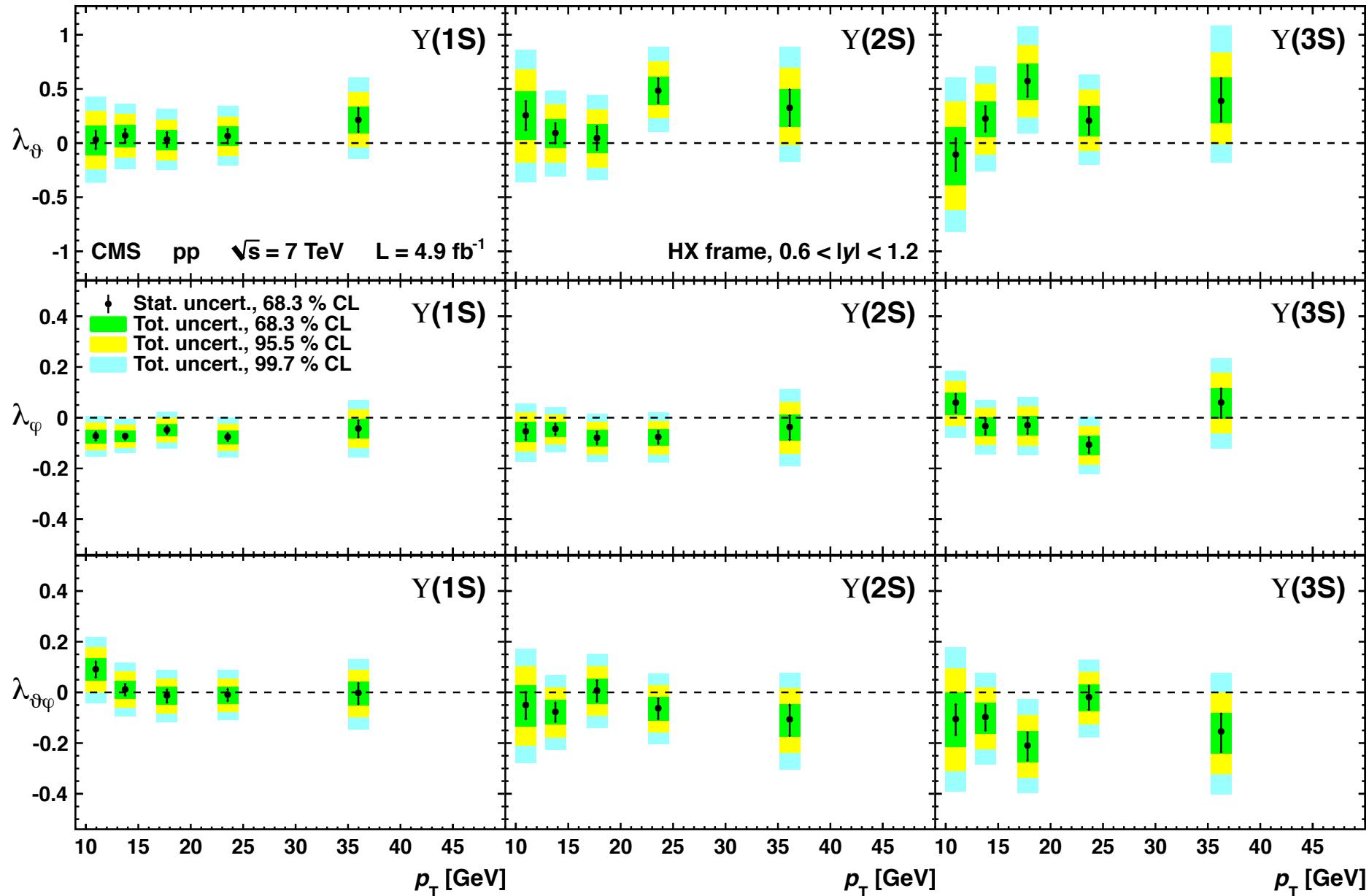
- Data-driven single muon efficiencies measured with the *Tag&Probe* method
- Precise knowledge of efficiencies needed to avoid introducing artificial polarization
- Dimuon efficiencies are calculated as the product of single muon efficiencies
- Correlations between muons are negligible as seen in detailed MC studies
- Efficiencies are accounted for on an event-by-event basis



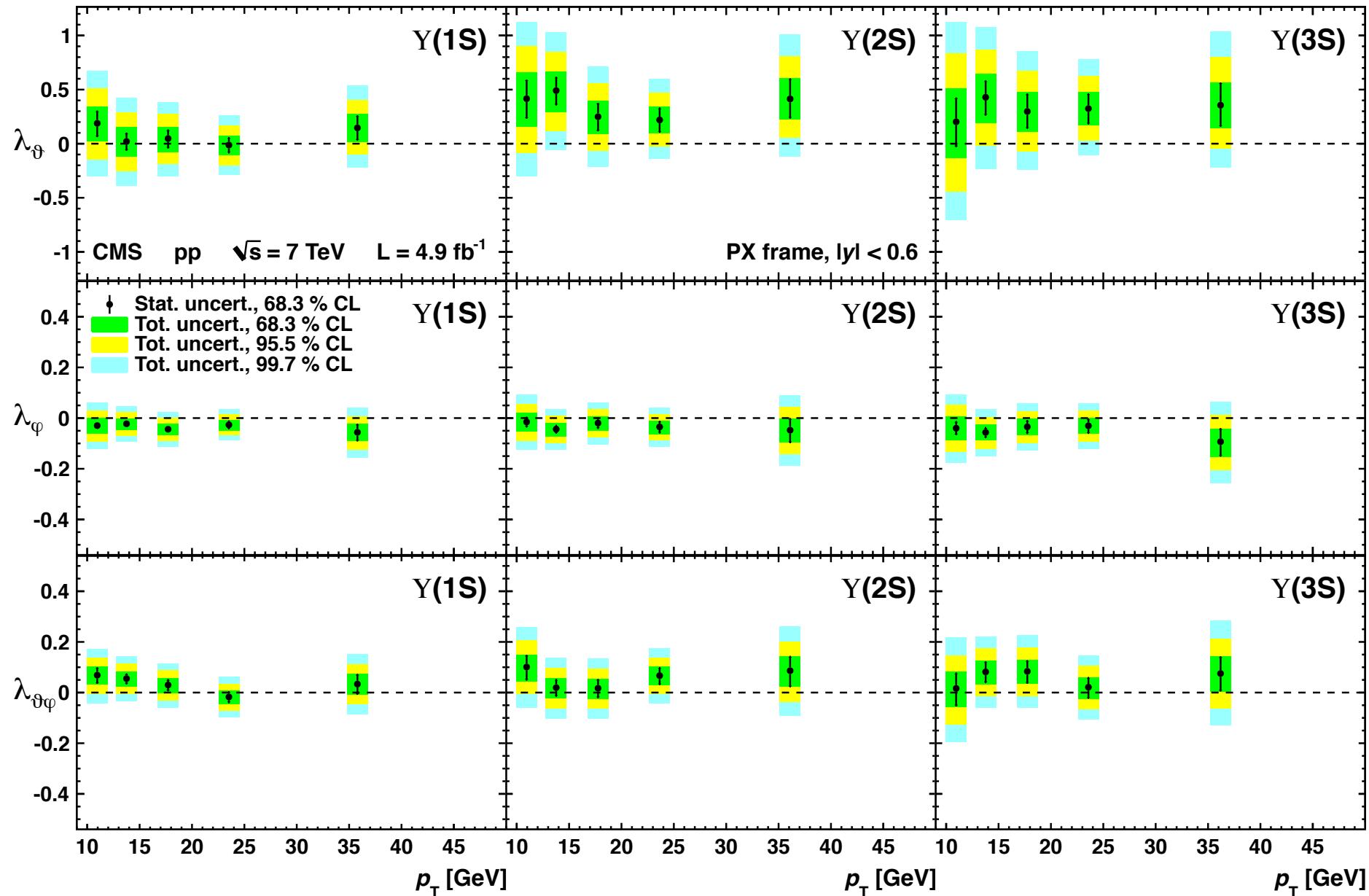
# Frame Invariant Parameter $\tilde{\lambda}$



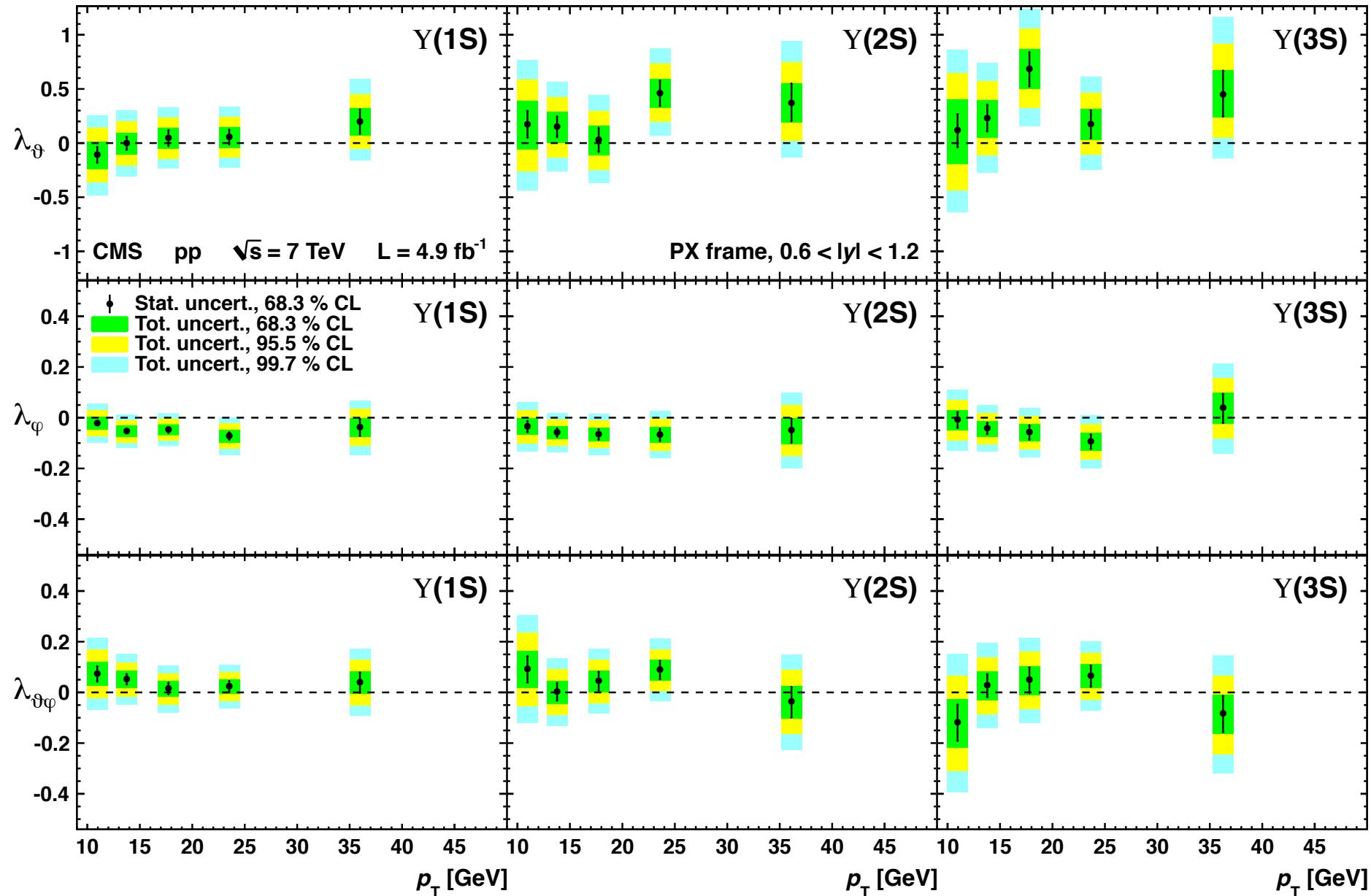
# $\Upsilon(nS)$ Polarization in the HX Frame, $0.6 < |y| < 1.2$



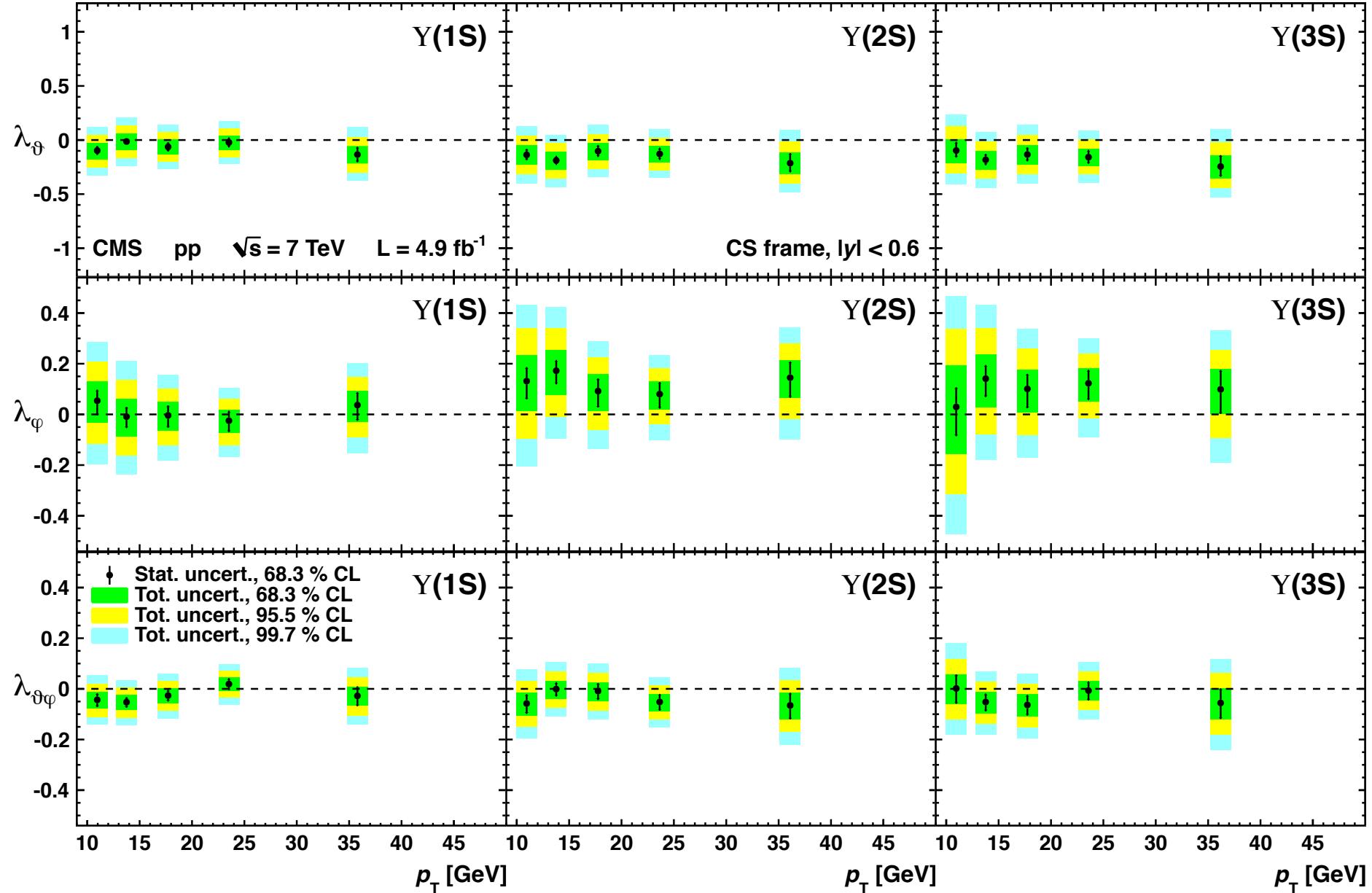
# $\Upsilon(nS)$ Polarization in the PX Frame, $|y| < 0.6$



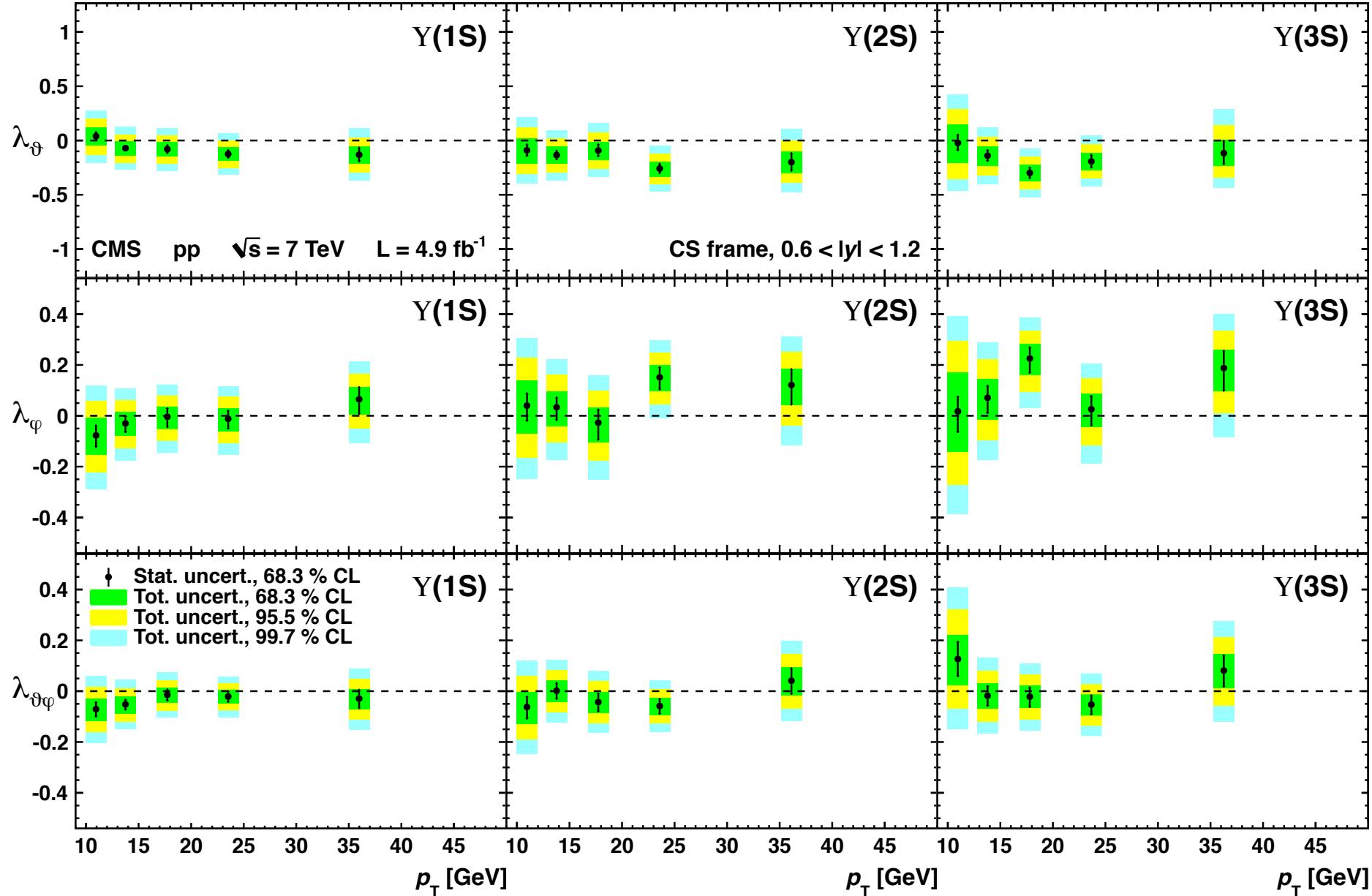
# $\Upsilon(nS)$ Polarization in the PX Frame, $0.6 < |y| < 1.2$



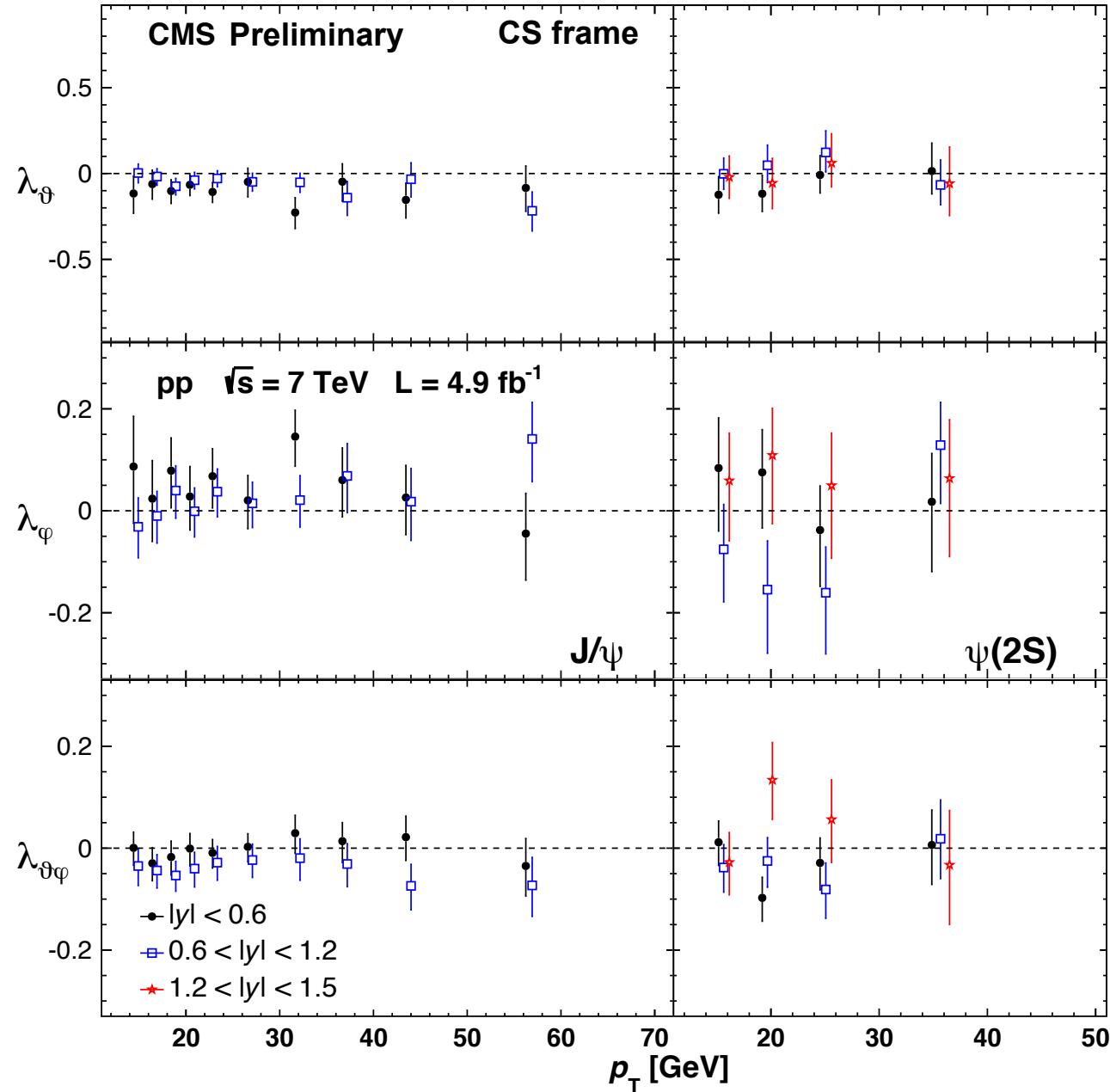
# $\Upsilon(nS)$ Polarization in the CS Frame, $|y| < 0.6$



# $\Upsilon(nS)$ Polarization in the CS Frame, $0.6 < |y| < 1.2$

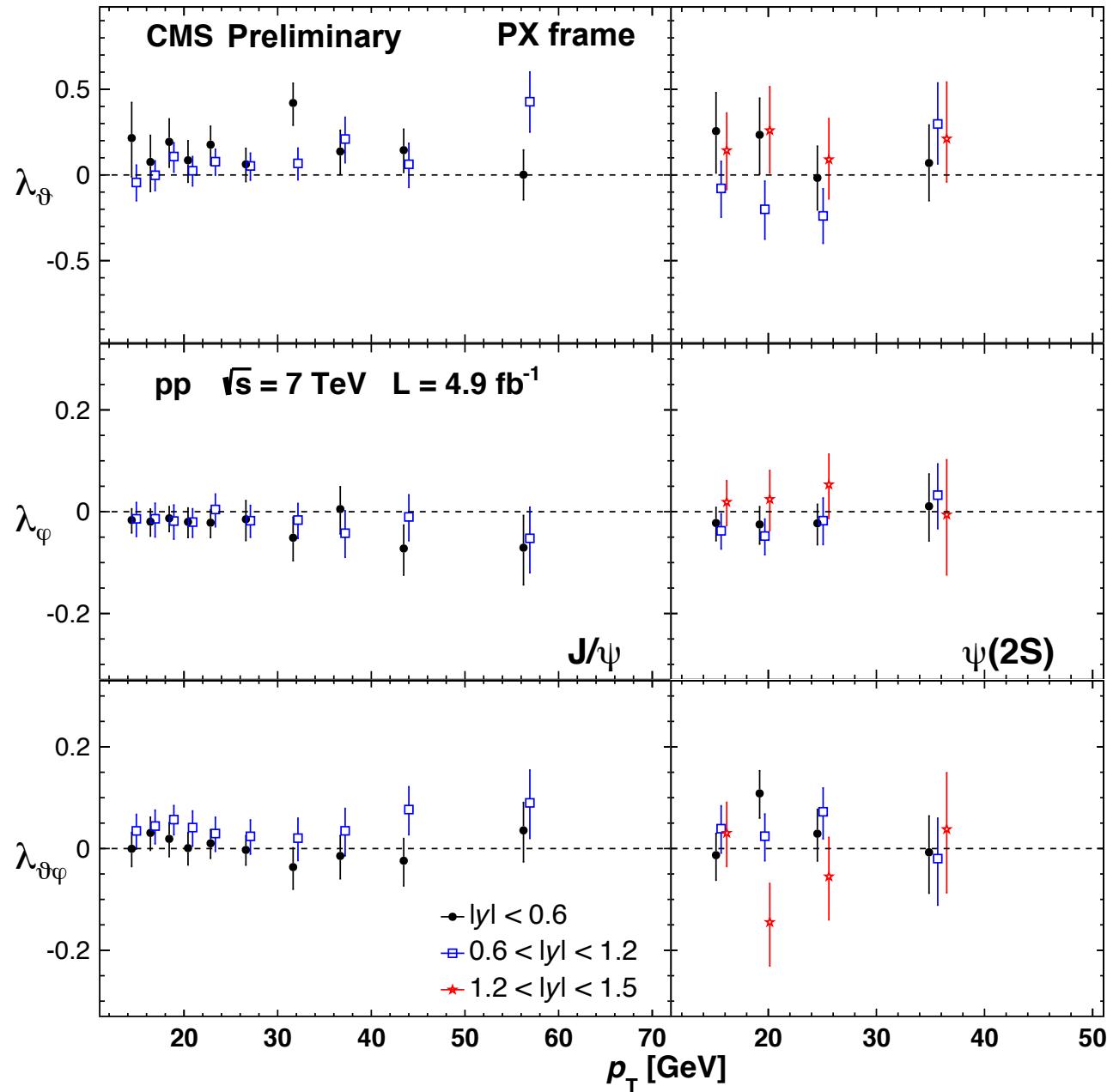


# Prompt $\psi(nS)$ Polarization in the CS frame



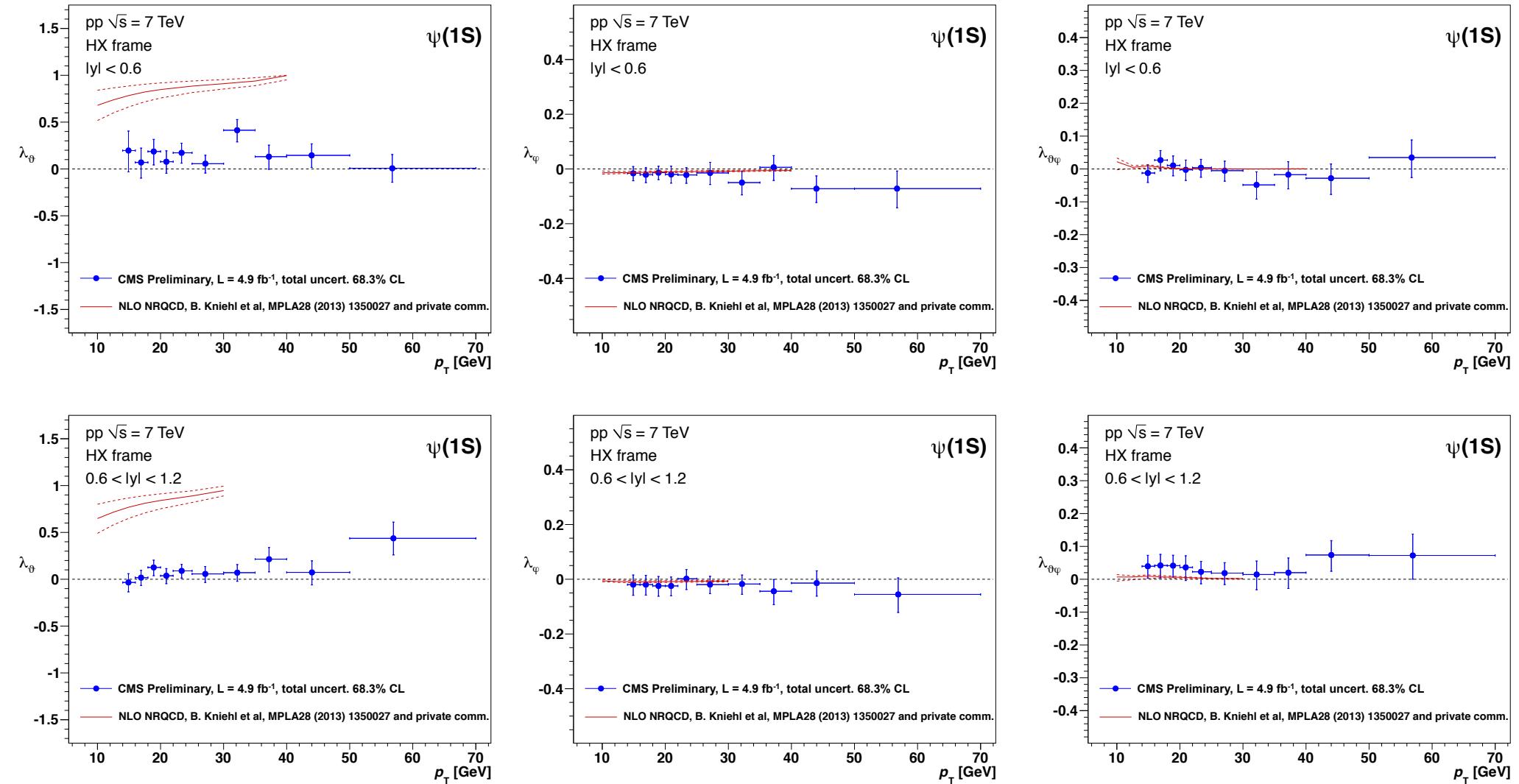
Error bars show total uncertainties at 68.3% CL

# Prompt $\psi(nS)$ Polarization in the PX frame



Error bars show total  
uncertainties at  
68.3% CL

# $\psi(1S)$ : Comparison to NLO NRQCD



# $\psi(2S)$ : Comparison to NLO NRQCD

