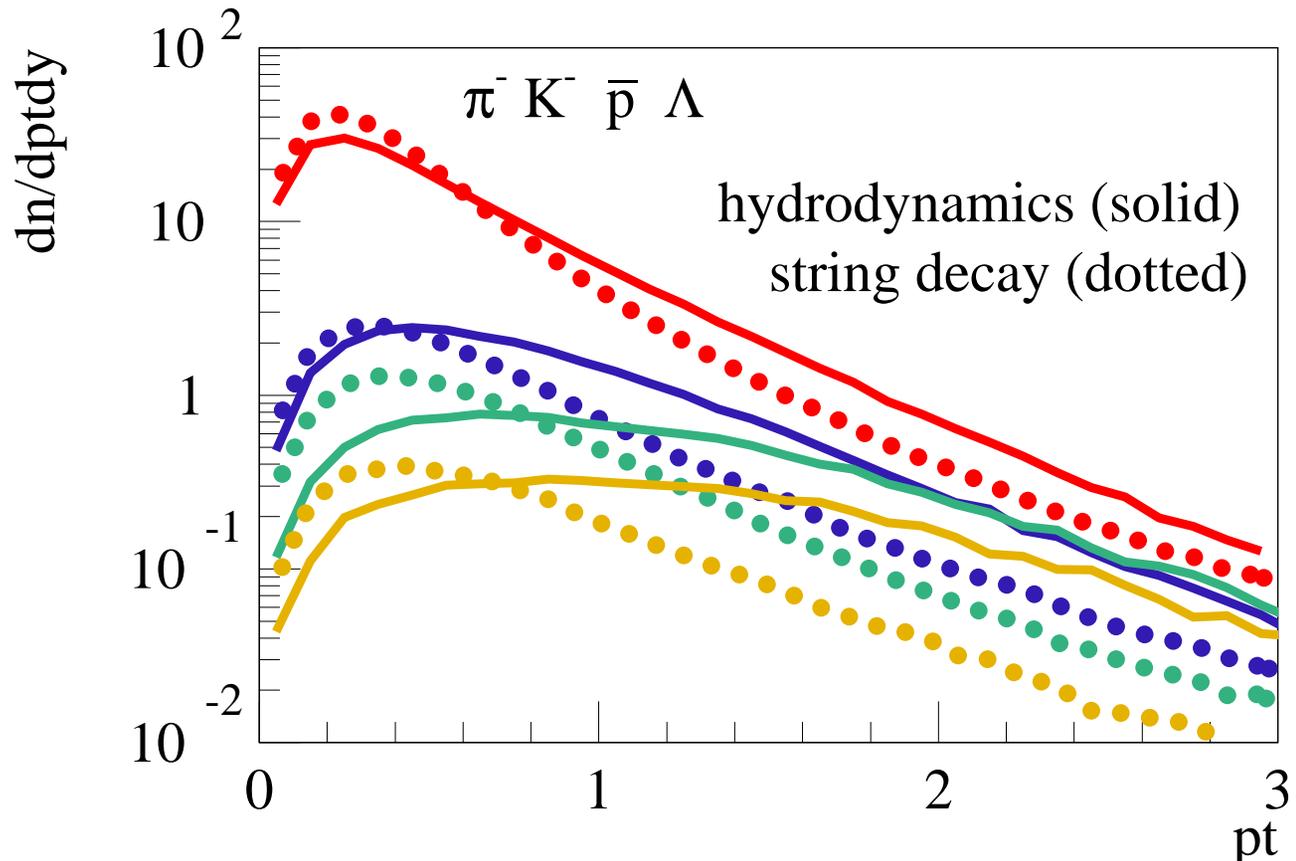


Flow in proton-nucleus collisions at 5 TeV

Manifestation of flow:

Particle spectra affected by radial flow



=> mass ordering of $\langle p_t \rangle$, lambda/K increase

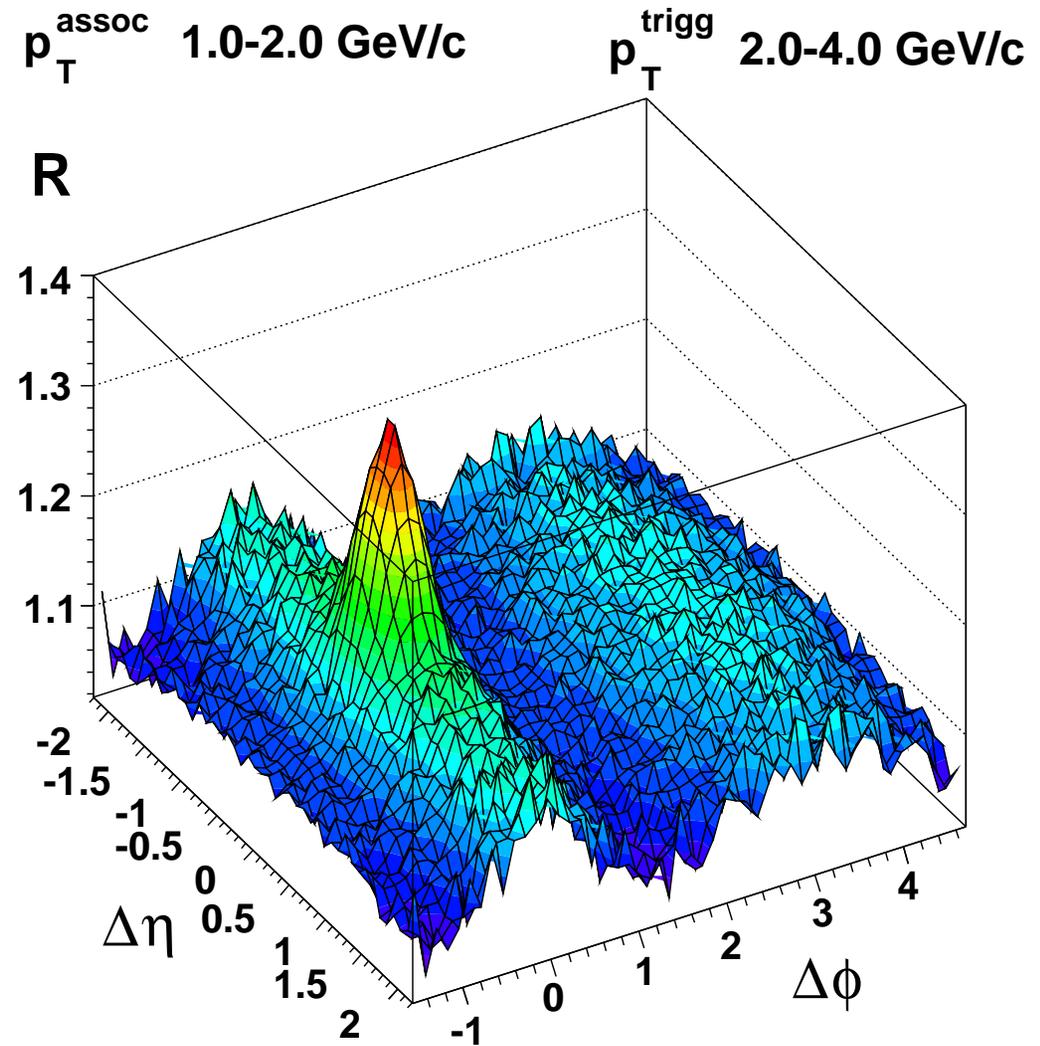
Ridges & flow harmonics

Ridges appear in

$$R = \frac{1}{N_{\text{trigg}}} \frac{dn}{d\Delta\phi\Delta\eta}$$

**due to initial
azimuthal
anisotropies**

(longitudinally
invariant)



EPOS3.074

pPb data, interpreted in terms of hydrodynamic flow

Models:

P. Bozek, W. Broniowski, arXiv:1304.3044

analysis of pPb@5TeV

- Glauber model (wounded nucleon model) initial conditions
- Viscous hydrodynamic expansion, $\eta/s = 0.08$ or 0.16
- Statistical hadronization using “Terminator”

A. Bzdak, B. Schenke, P. Tribedy, R. Venugopalan,

arXiv:1304.3403

- Theoretical study of flow in pp, pA, dA
- Glauber model or Color Glass Condensate initial conditions
- Viscous hydrodynamic expansion, $\eta/s = 0.08$

EPOS3, B. Guiot, Y. Karpenko, T. Pierog, K. Werner

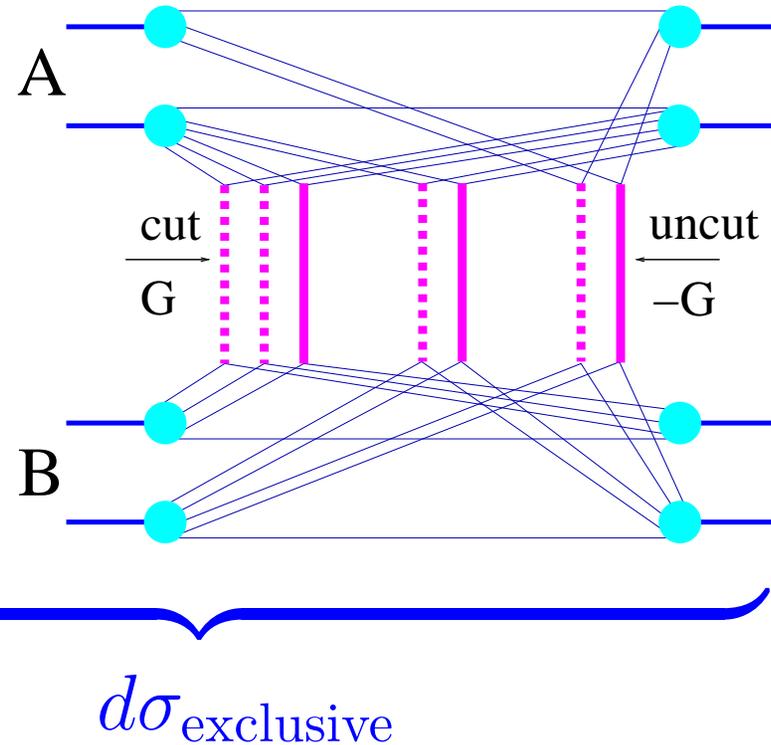
- Initial conditions:
Gribov-Regge multiple scattering approach,
elementary object = Pomeron = parton ladder,
using saturation scale $Q_s \propto N_{part} \hat{s}^\lambda$
- Core-corona approach
to separate fluid and jet hadrons
- Viscous hydrodynamic expansion, $\eta/s = 0.08$
- Statistical hadronization, final state hadronic cascade

EPOS3 will be used in the following

EPOS IC: Marriage pQCD+GRT+energy sharing

(Drescher, Hladik, Ostapchenko, Pierog, and Werner, Phys. Rept. 350, 2001)

$$\sigma^{\text{tot}} = \sum_{\text{cut P}} \int \sum_{\text{uncut P}} \int$$



$$\text{cut Pom} : G = \frac{1}{2\hat{s}} 2\text{Im} \{ \mathcal{FT} \{ T \} \} (\hat{s}, b), T = i\hat{s} \sigma_{\text{hard}}(\hat{s}) \exp(R_{\text{hard}}^2 t)$$

Nonlinear effects considered via saturation scale $Q_s \propto N_{\text{part}} \hat{s}^\lambda$

$$\begin{aligned}
 \sigma^{\text{tot}} = & \int d^2b \int \prod_{i=1}^A d^2b_i^A dz_i^A \rho_A(\sqrt{(b_i^A)^2 + (z_i^A)^2}) \\
 & \prod_{j=1}^B d^2b_j^B dz_j^B \rho_B(\sqrt{(b_j^B)^2 + (z_j^B)^2}) \\
 & \sum_{m_1 l_1} \dots \sum_{m_{AB} l_{AB}} (1 - \delta_{0 \Sigma m_k}) \int \prod_{k=1}^{AB} \left(\prod_{\mu=1}^{m_k} dx_{k,\mu}^+ dx_{k,\mu}^- \prod_{\lambda=1}^{l_k} d\tilde{x}_{k,\lambda}^+ d\tilde{x}_{k,\lambda}^- \right) \left\{ \right. \\
 & \prod_{k=1}^{AB} \left(\frac{1}{m_k!} \frac{1}{l_k!} \prod_{\mu=1}^{m_k} G(x_{k,\mu}^+, x_{k,\mu}^-, s, |\vec{b} + \vec{b}_{\pi(k)}^A - \vec{b}_{\tau(k)}^B|) \right. \\
 & \left. \left. \prod_{\lambda=1}^{l_k} -G(\tilde{x}_{k,\lambda}^+, \tilde{x}_{k,\lambda}^-, s, |\vec{b} + \vec{b}_{\pi(k)}^A - \vec{b}_{\tau(k)}^B|) \right) \right\} \\
 & \prod_{i=1}^A \left(1 - \sum_{\pi(k)=i} x_{k,\mu}^+ - \sum_{\pi(k)=i} \tilde{x}_{k,\lambda}^+ \right)^\alpha \prod_{j=1}^B \left(1 - \sum_{\tau(k)=j} x_{k,\mu}^- - \sum_{\tau(k)=j} \tilde{x}_{k,\lambda}^- \right)^\alpha \left. \right\}
 \end{aligned}$$

The hydrodynamic equations (Israel-Stewart formulation) in arbitrary coordinate system (implemented/solved by Yuri Karpenko), always $\eta/S = 0.08$, $\zeta/S = 0$

$$\partial_{;\nu} T^{\mu\nu} = \partial_\nu T^{\mu\nu} + \Gamma_{\nu\lambda}^\mu T^{\nu\lambda} + \Gamma_{\nu\lambda}^\nu T^{\mu\lambda} = 0$$

$$\gamma (\partial_t + v_i \partial_i) \pi^{\mu\nu} = -\frac{\pi^{\mu\nu} - \pi_{\text{NS}}^{\mu\nu}}{\tau_\pi} + I_\pi^{\mu\nu}$$

$$\gamma (\partial_t + v_i \partial_i) \Pi = -\frac{\Pi - \Pi_{\text{NS}}}{\tau_\Pi} + I_\Pi$$

$T^{\mu\nu} = \epsilon u^\mu u^\nu - (p + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$,

$\partial_{;\nu}$ denotes a covariant derivative,

$\Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$ is the projector orthogonal to u^μ ,

$\pi^{\mu\nu}$ and Π are the shear stress tensor and bulk pressure, respectively.

$\pi_{\text{NS}}^{\mu\nu} = \eta(\Delta^{\mu\lambda} \partial_{;\lambda} u^\nu + \Delta^{\nu\lambda} \partial_{;\lambda} u^\mu) - \frac{2}{3} \eta \Delta^{\mu\nu} \partial_{;\lambda} u^\lambda$

$\Pi_{\text{NS}} = -\zeta \partial_{;\lambda} u^\lambda$

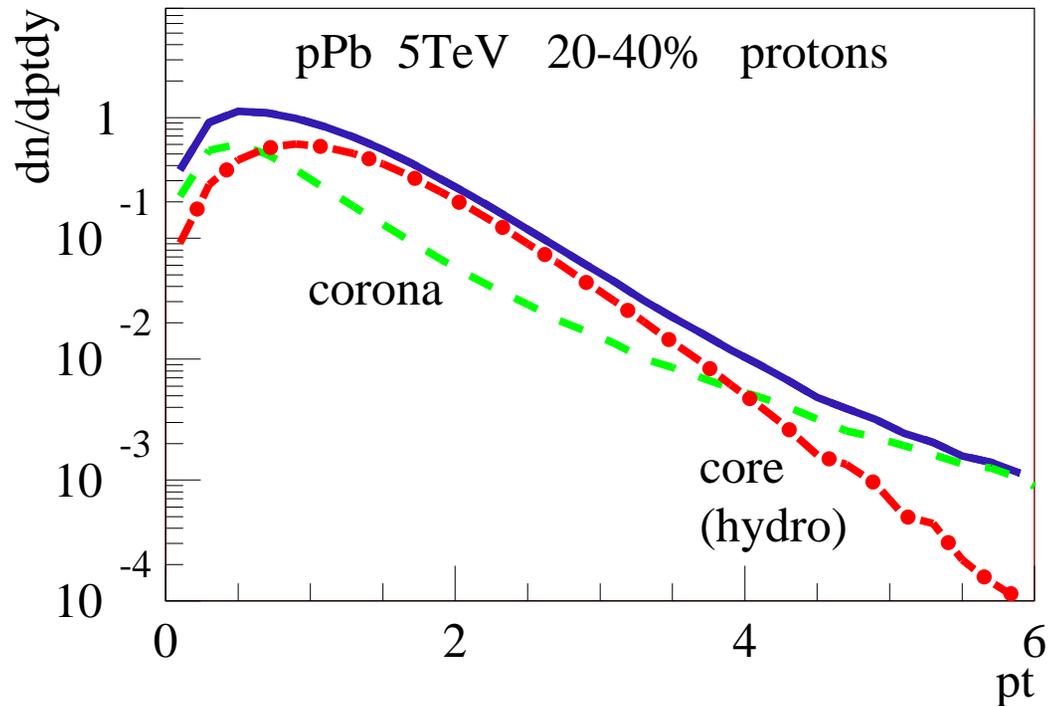
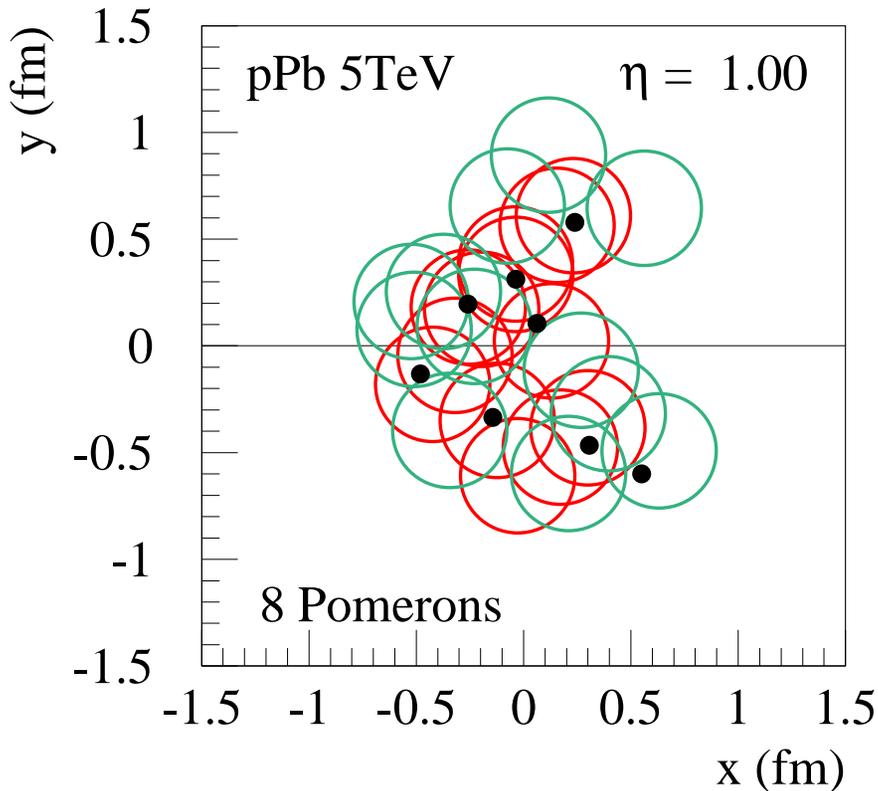
$I_\pi^{\mu\nu} = -\frac{4}{3} \pi^{\mu\nu} \partial_{;\gamma} u^\gamma - [u^\nu \pi^{\mu\beta} + u^\mu \pi^{\nu\beta}] u^\lambda \partial_{;\lambda} u_\beta$

$I_\Pi = -\frac{4}{3} \Pi \partial_{;\gamma} u^\gamma$

EPOS3:

Pomeron => parton ladder => flux tube (kinky string)

String segments with high p_t escape => **corona**,
 the others form the **core** = initial condition for hydro
 depending on the local string density



CMS: Multiplicity dependence of pion, kaon, proton pt spectra

CMS, arXiv:1307.3442

We plot 4 “centrality” classes:

$\langle N_{\text{tracks}} \rangle = 8, 84, 160, 235$ (in $|\eta| < 2.4$)

Multiplicity = centrality measure

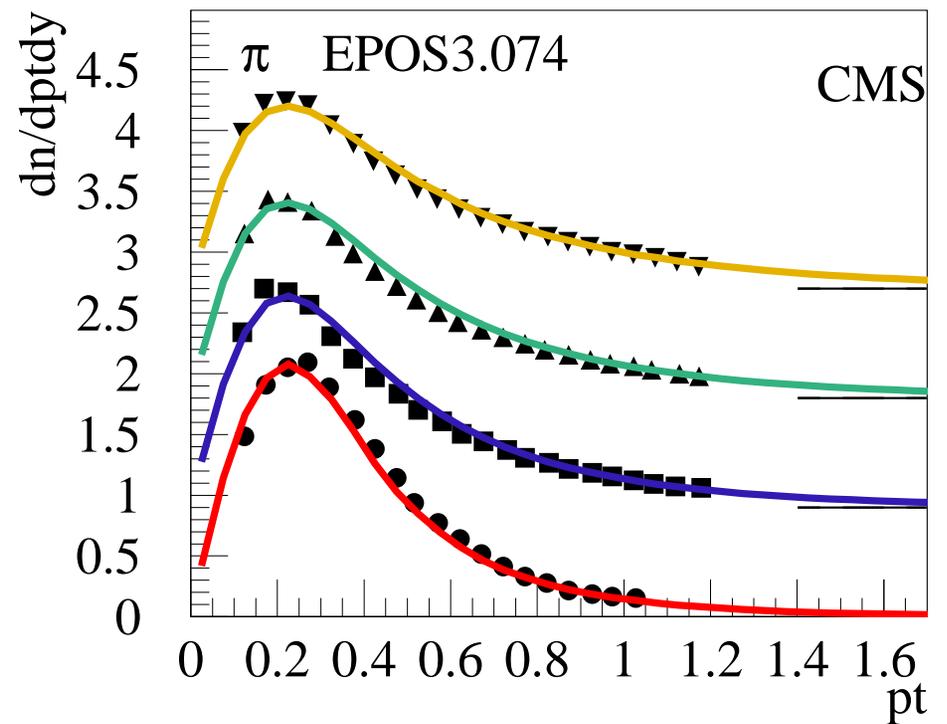
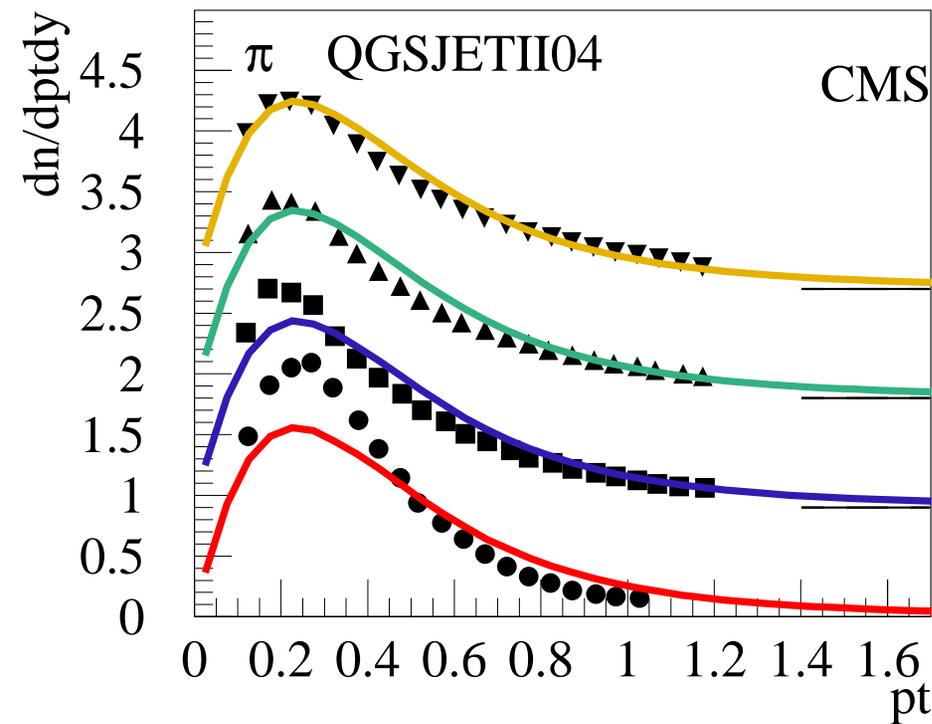
in EPOS: high multiplicity = many Pomerons

Data compared to

- EPOS3 (hydrodynamic expansion, flow)**
- QGSJETII (no flow effects, only string decay)**

Pions

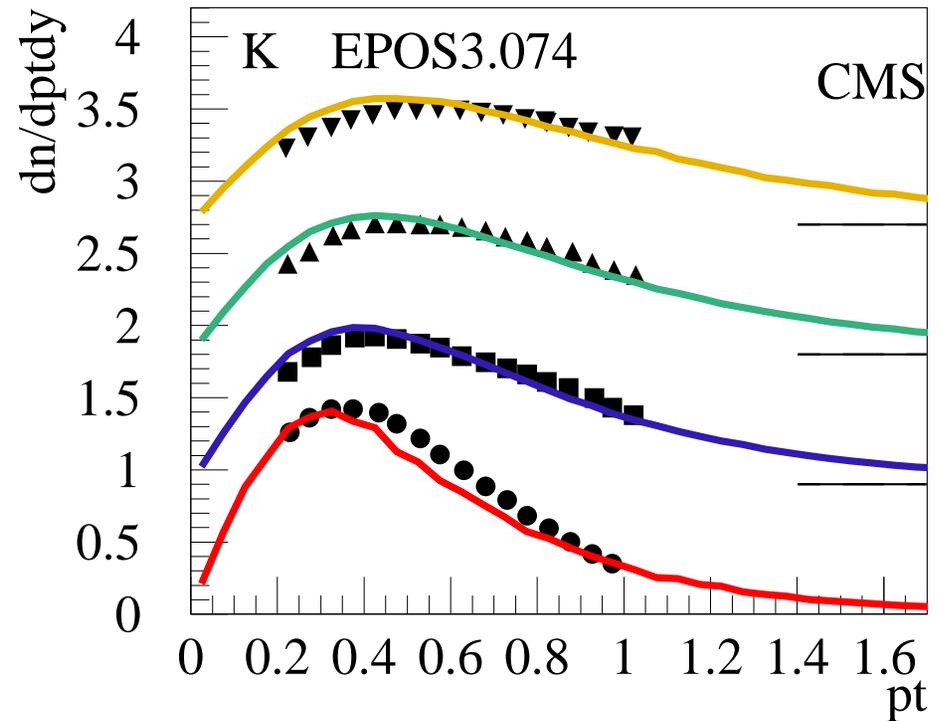
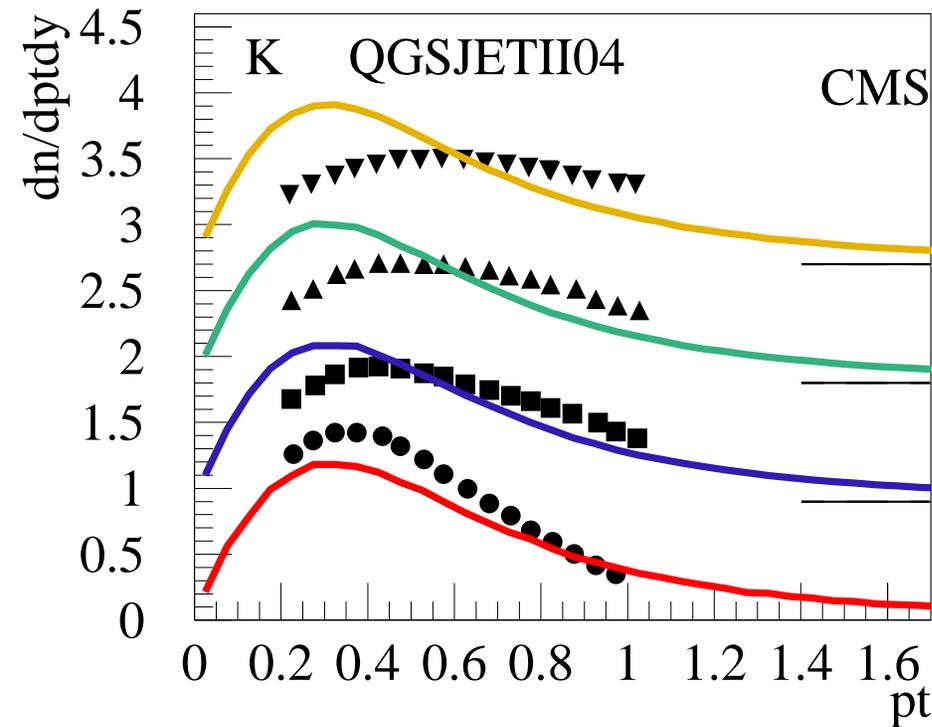
$\langle N_{\text{tracks}} \rangle = 8, 84, 160, 235$, from bottom to top, curves shifted by 0.9
 spectra normalized to unity, lines = theory



Little change with $\langle N_{\text{tracks}} \rangle$ for pions

Kaons

$\langle N_{\text{tracks}} \rangle = 8, 84, 160, 235$, from bottom to top, curves shifted by 0.9
 spectra normalized to unity, lines = theory

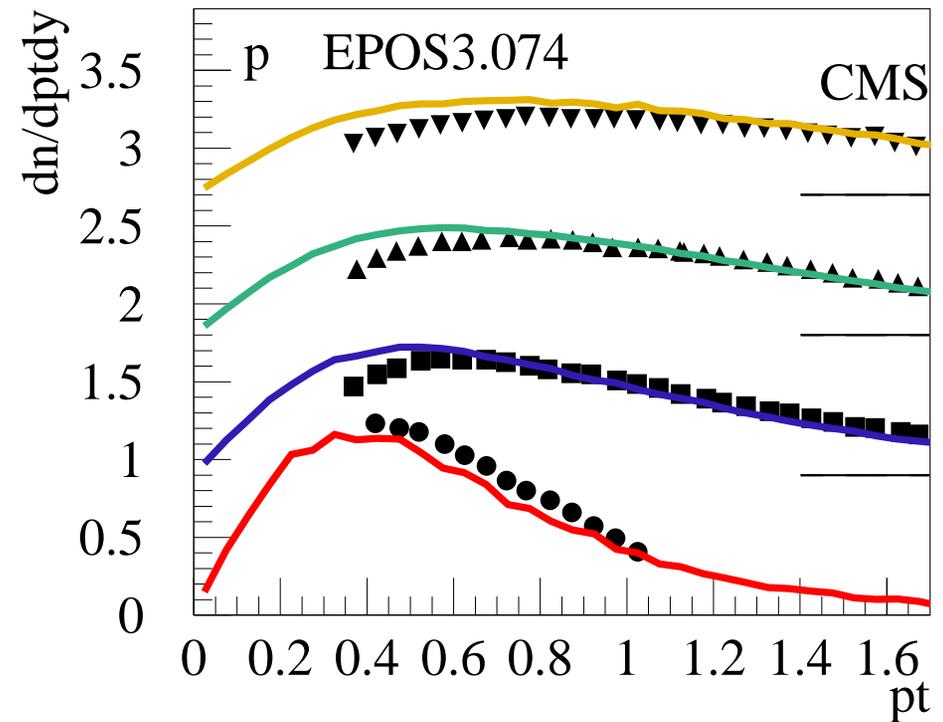
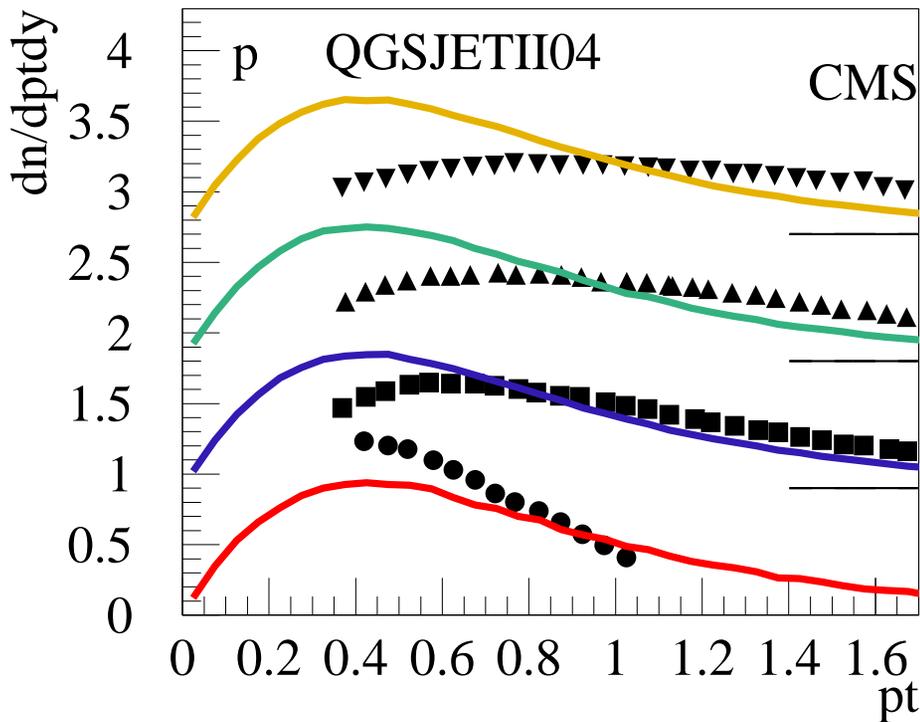


Kaon spectra change significantly with $\langle N_{\text{tracks}} \rangle$

in EPOS3: more and more flow contribution

Protons

$\langle N_{\text{tracks}} \rangle = 8, 84, 160, 235$, from bottom to top, curves shifted by 0.9
 spectra normalized to unity, lines = theory



Strong variation of proton spectra

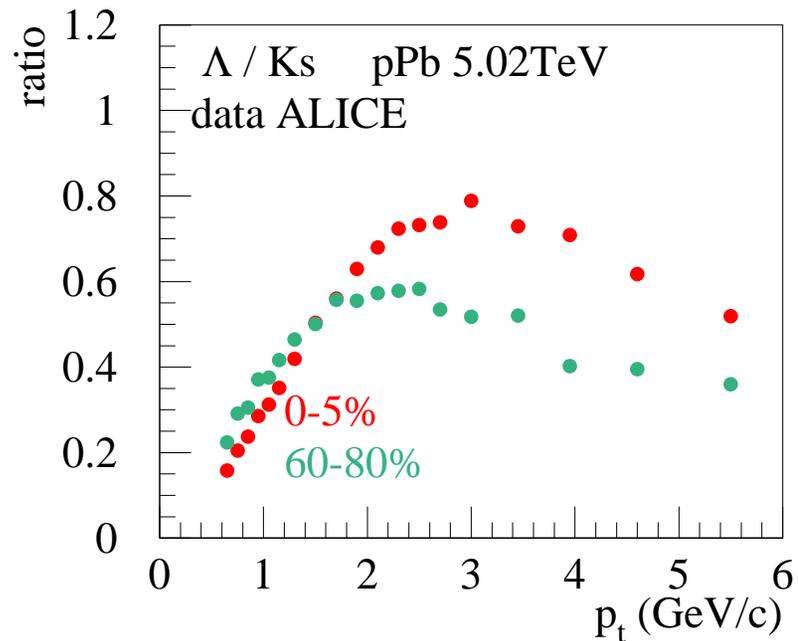
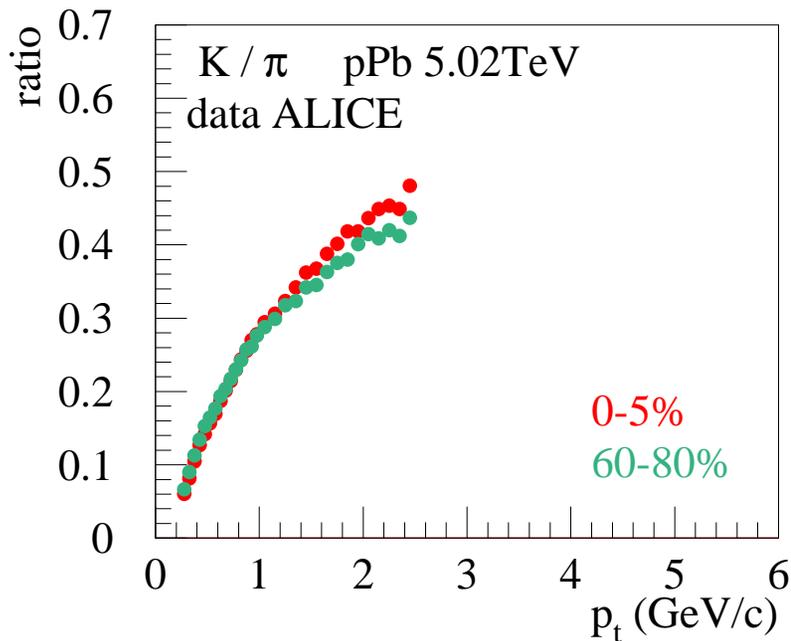
=> flow helps

ALICE: compare pt spectra for identified particles in different multiplicity classes: 0-5%,...,60-80%

(in $2.8 < \eta_{lab} < 5.1$)

R. Preghenella, ALICE, talk Trento workshop 2013

Useful : ratios (K/pi, p/pi...)

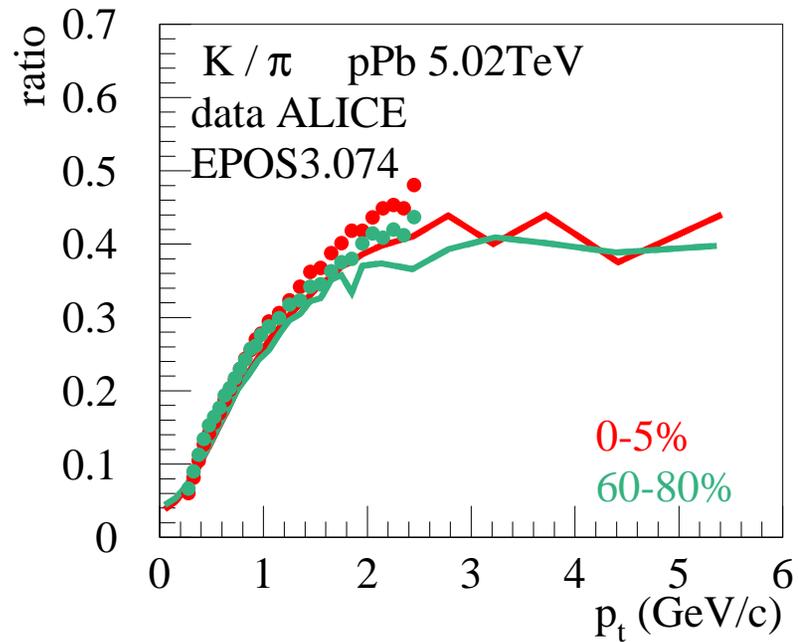
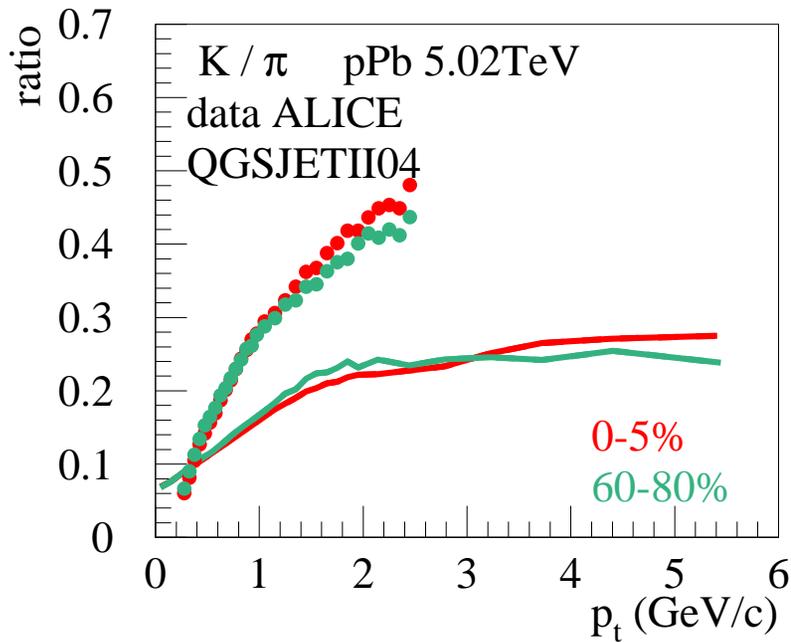


Significant variation of lambda/K – like in PbPb

$$K/\pi$$

High multiplicity (0-5%, red),
low multiplicity (60-80%, green)

lines = theory
 points = data



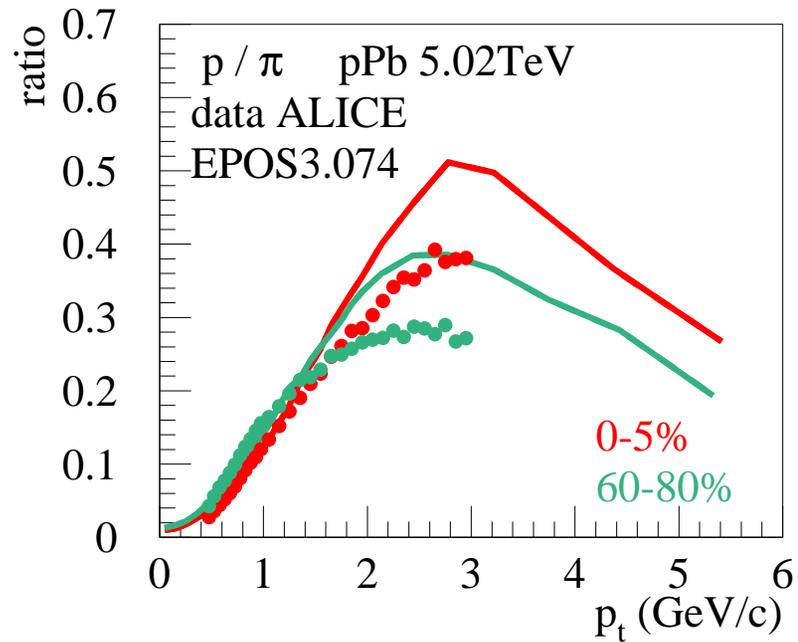
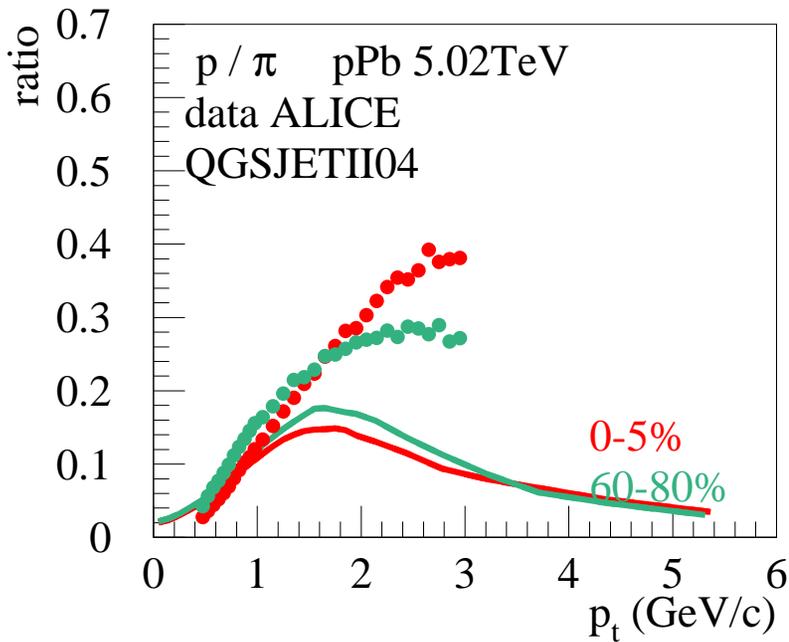
No multiplicity dependence

not trivial to get the peripheral right !!

$$p/\pi$$

High multiplicity (0-5%, red),
low multiplicity (60-80%, green)

lines = theory
 points = data



Significant multiplicity dependence

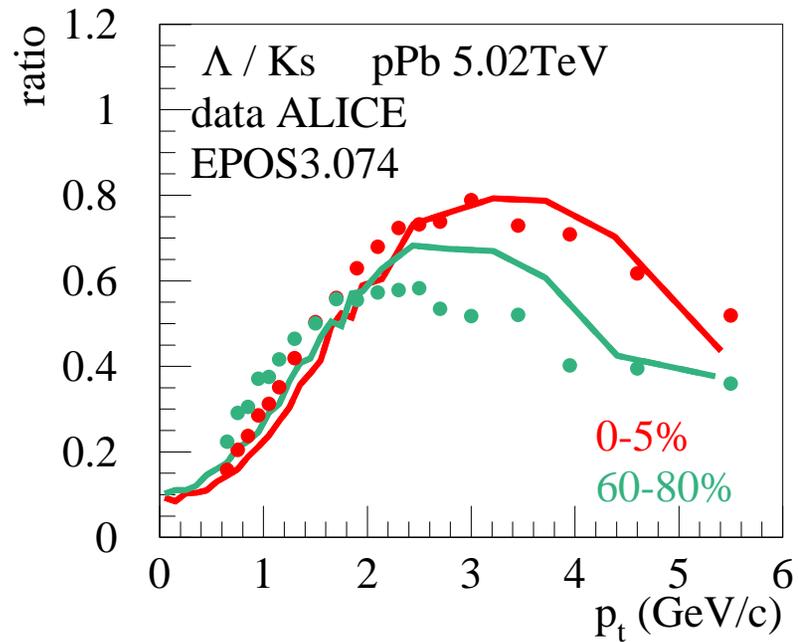
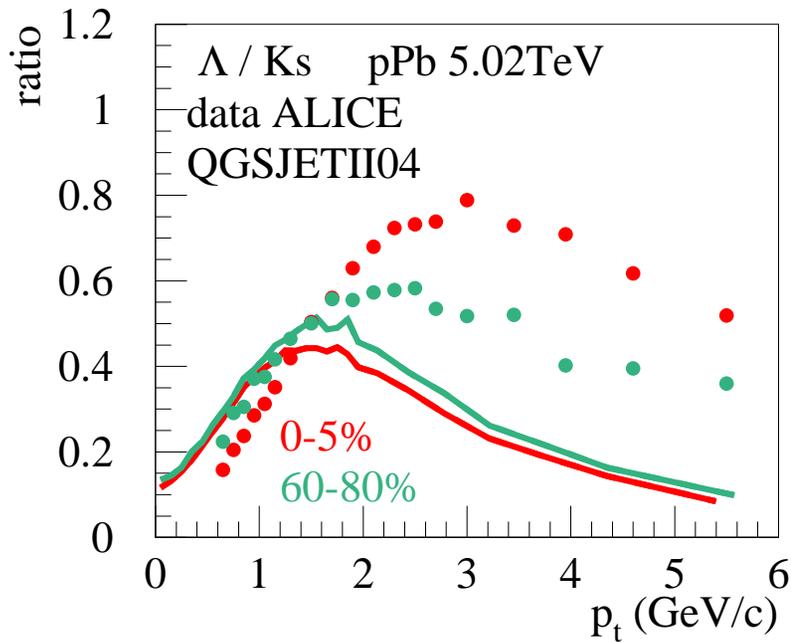
in EPOS, flow already affects the low multiplicity case

”flow peak” around 2-3 GeV/c,
beyond 5GeV/c corona (minjets) dominate

$$\Lambda / K_s$$

High multiplicity (0-5%, red),
low multiplicity (60-80%, green)

lines = theory
 points = data



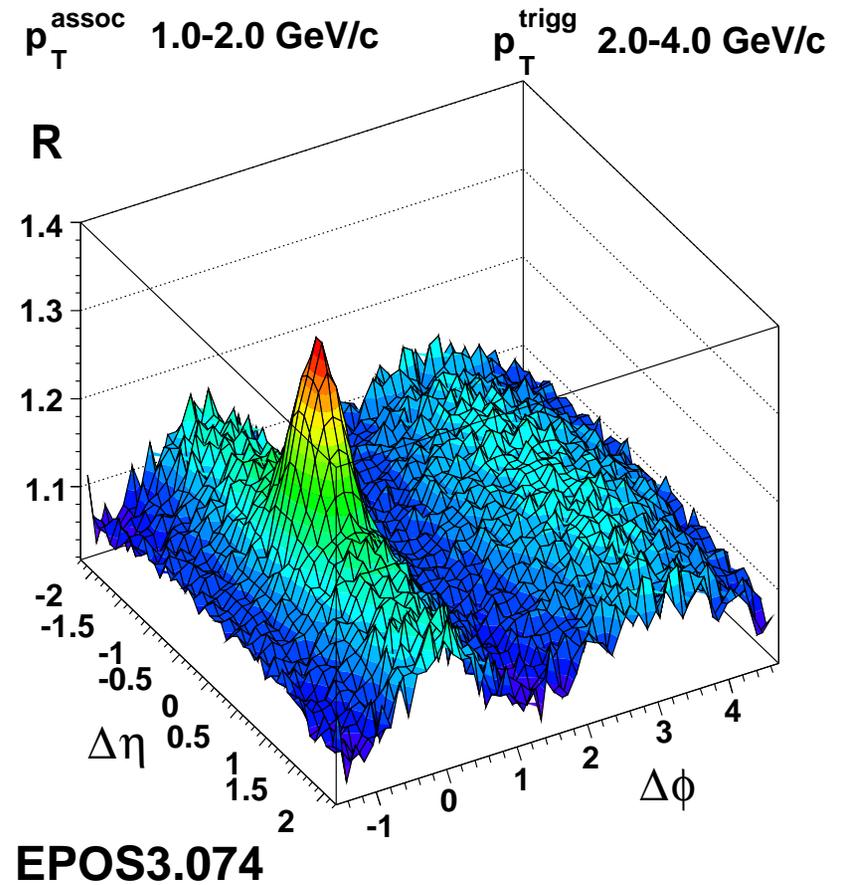
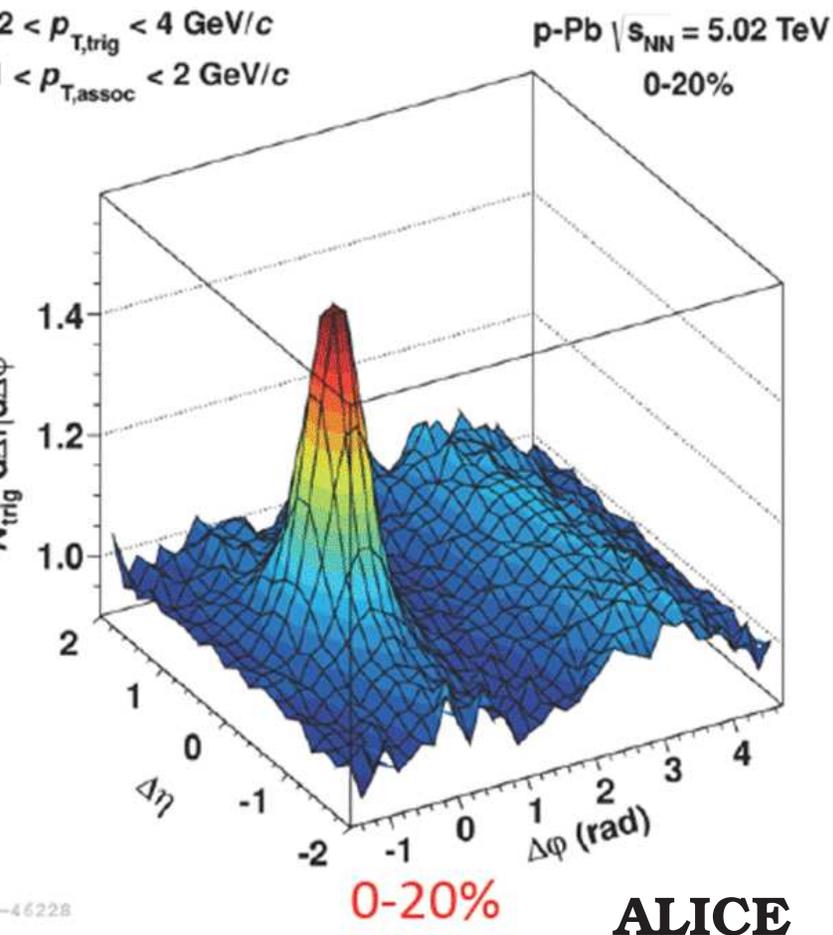
Significant multiplicity dependence

again, flow already needed for low multiplicity (even in pp!)

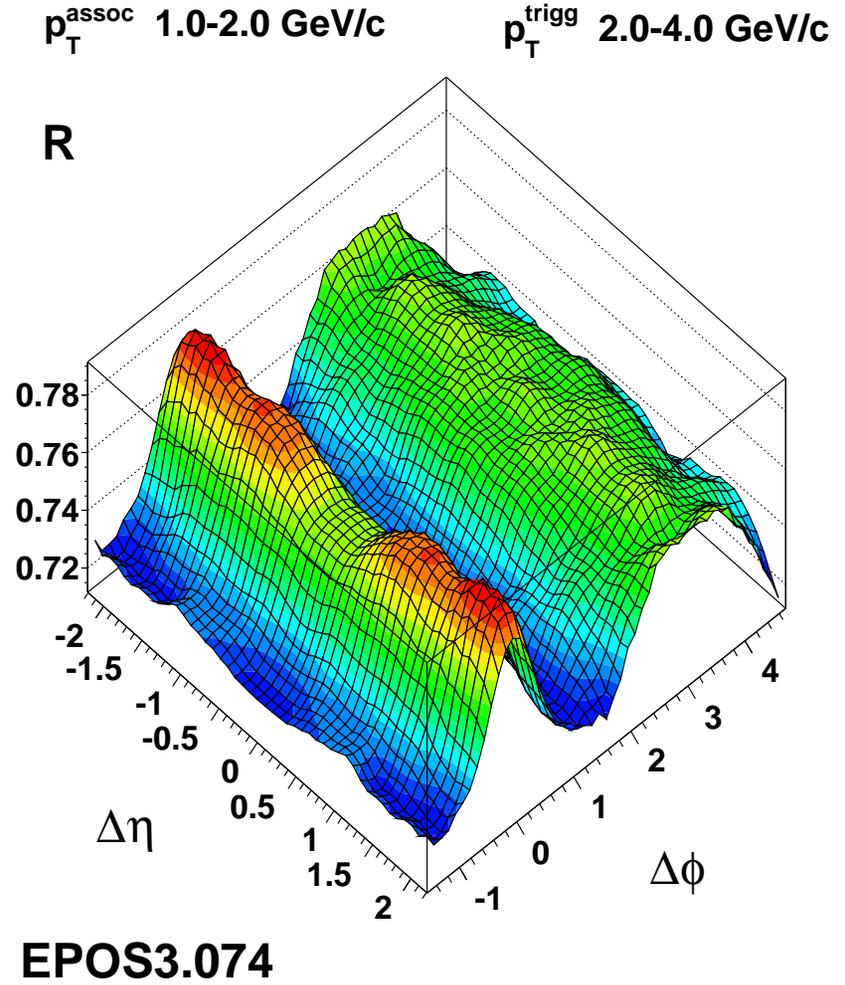
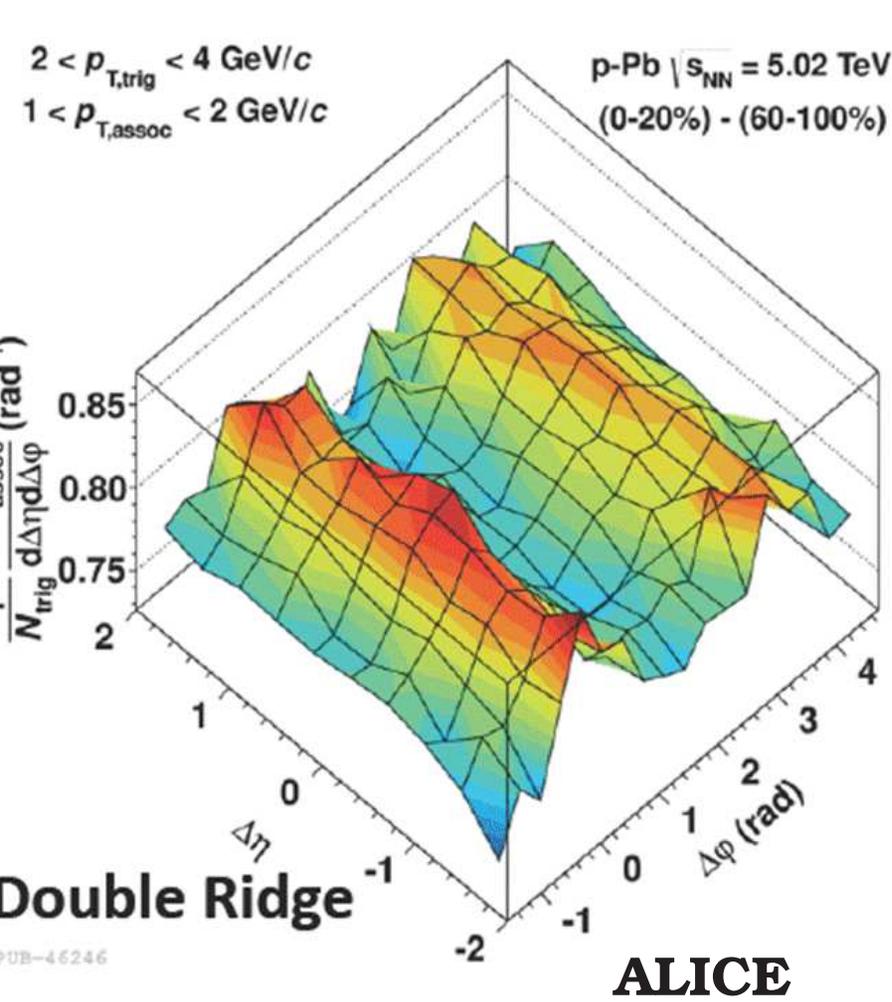
flow dominates 2-5 GeV/c

“Ridges” in pA

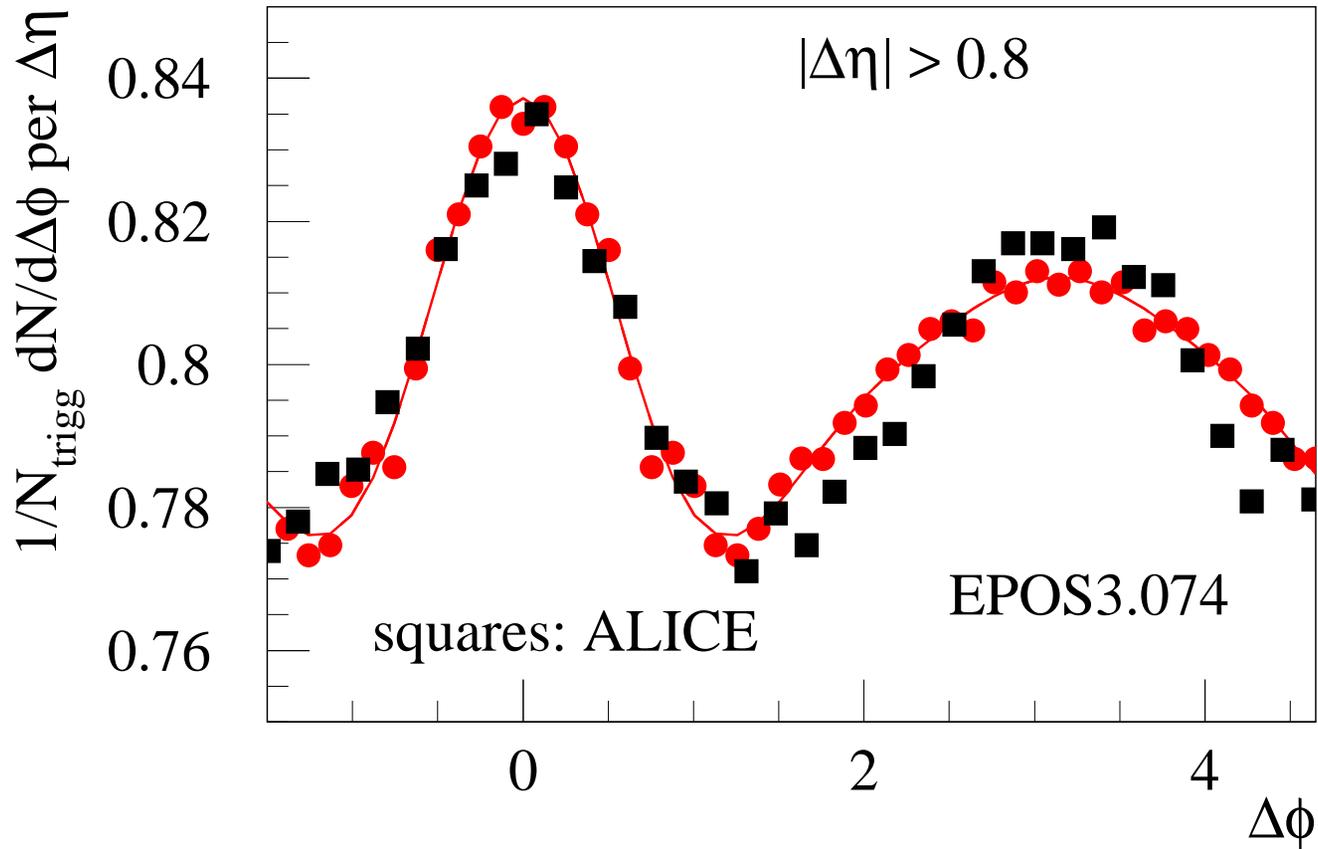
ALICE, arXiv:1212.2001, arXiv:1307.3237



Central - peripheral (to get rid of jets)



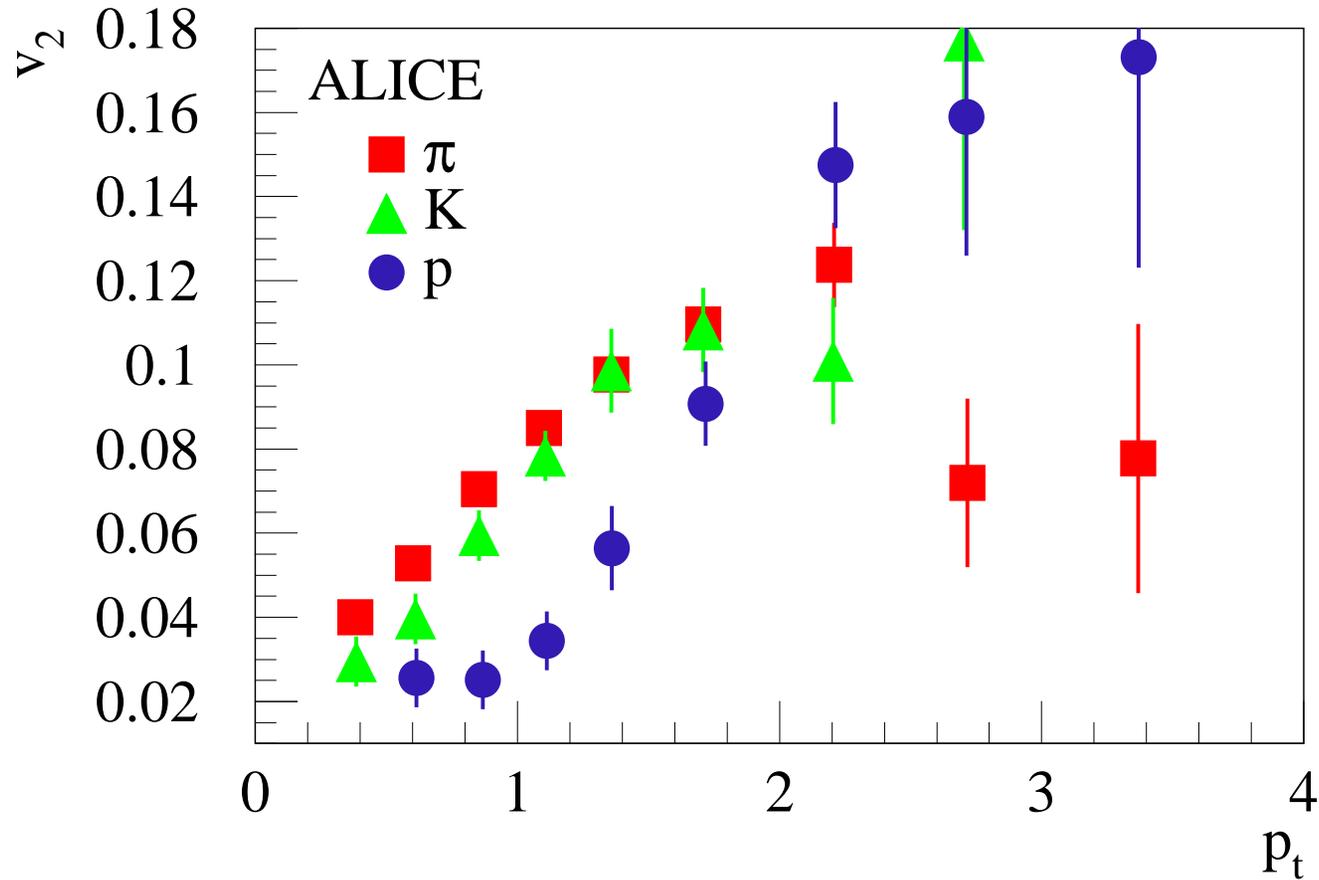
Projection



red line : $\sum 2a_n \cos(n\Delta\phi);$

$$\implies v_n = \sqrt{\frac{a_n}{b}}$$

Identified particle v_2



mass splitting, as in PbPb !!!

pPb in EPOS3

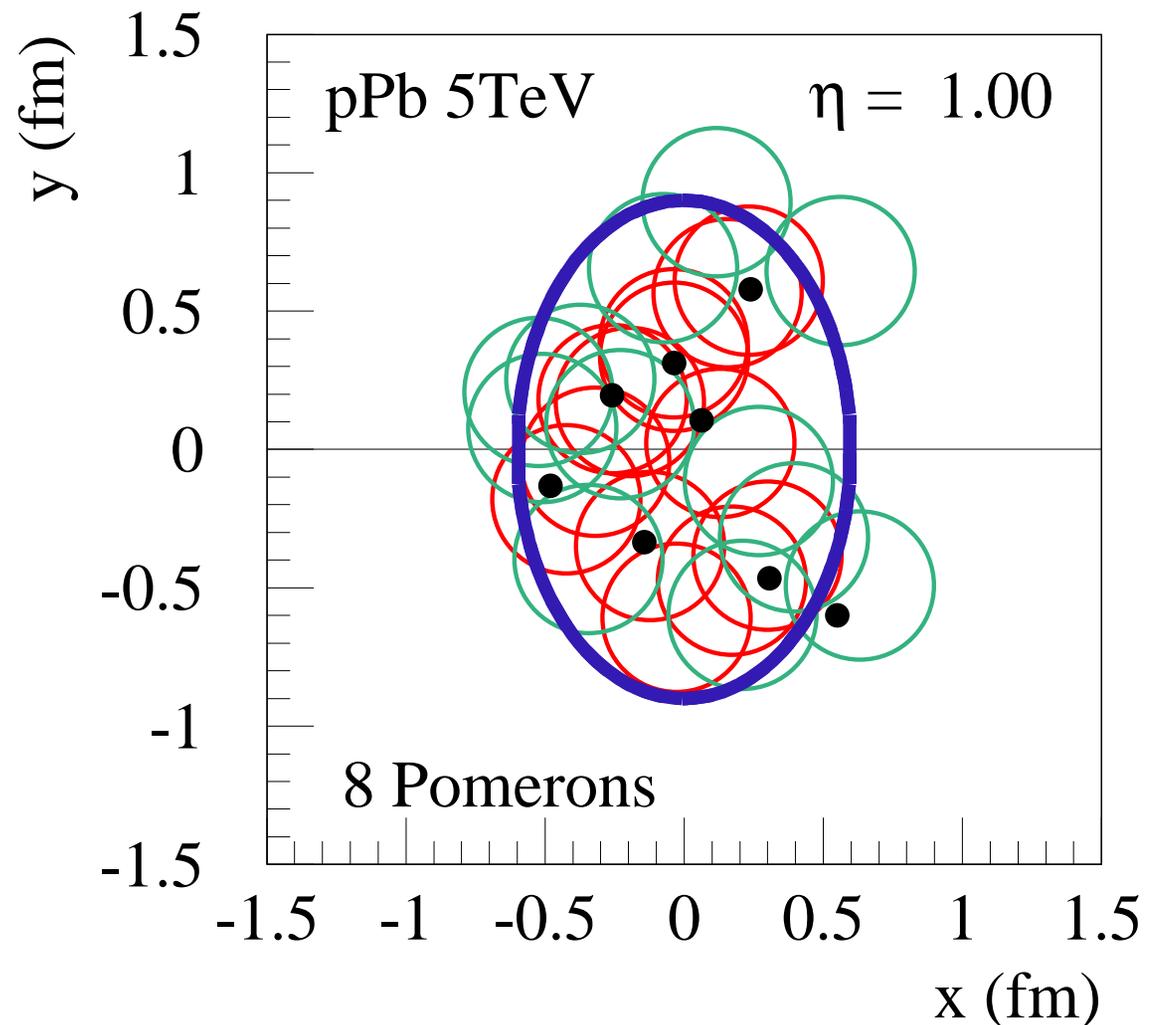
Pomerons (number and positions) characterize geometry (P. number \propto multiplicity)

random azimuthal asymmetry

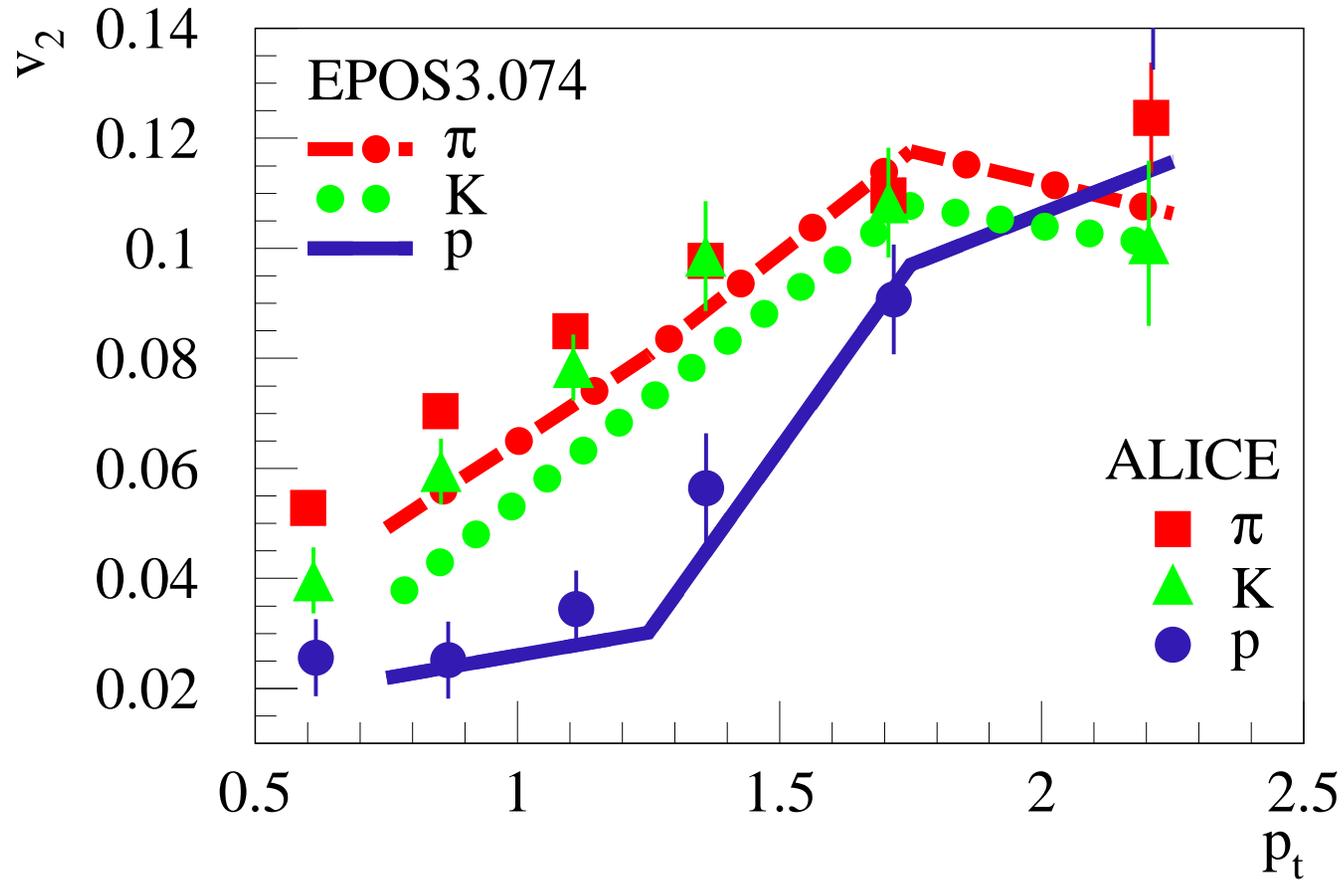
=>

asymmetric flow seen at higher pt for heavier ptls

Robust results

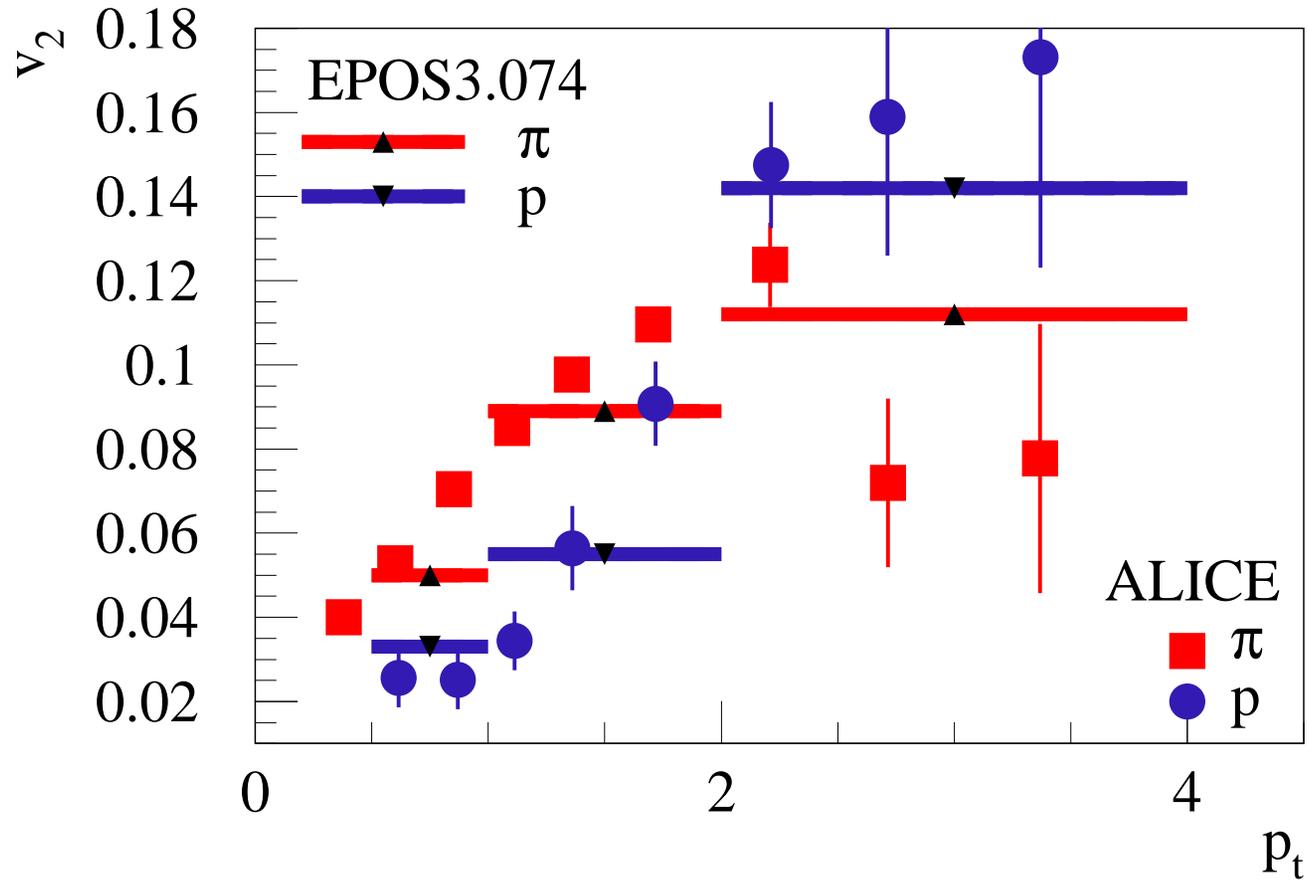


v_2 for π , K, p clearly differ



mass splitting, due to flow

different binning:

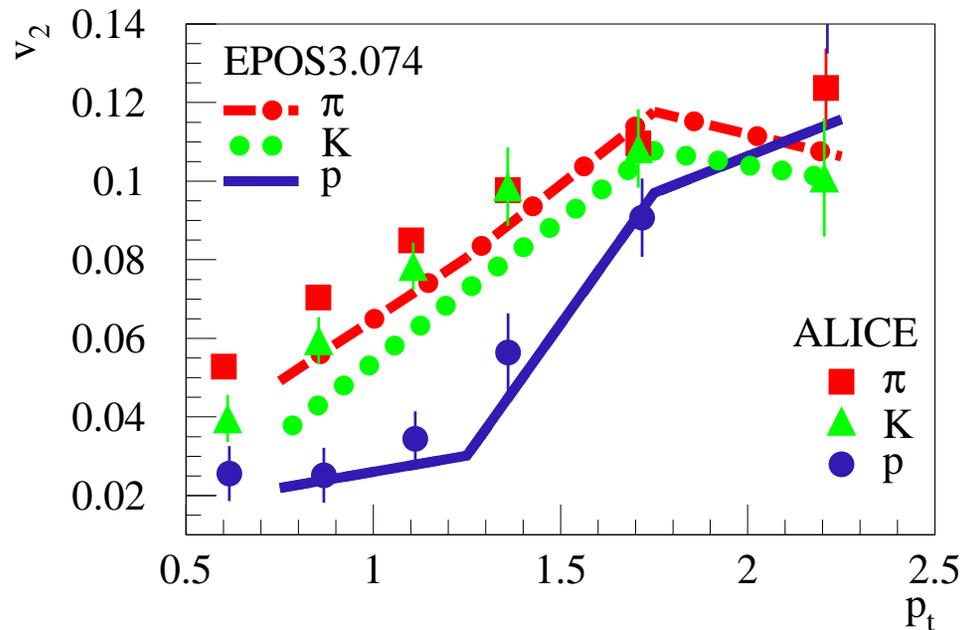


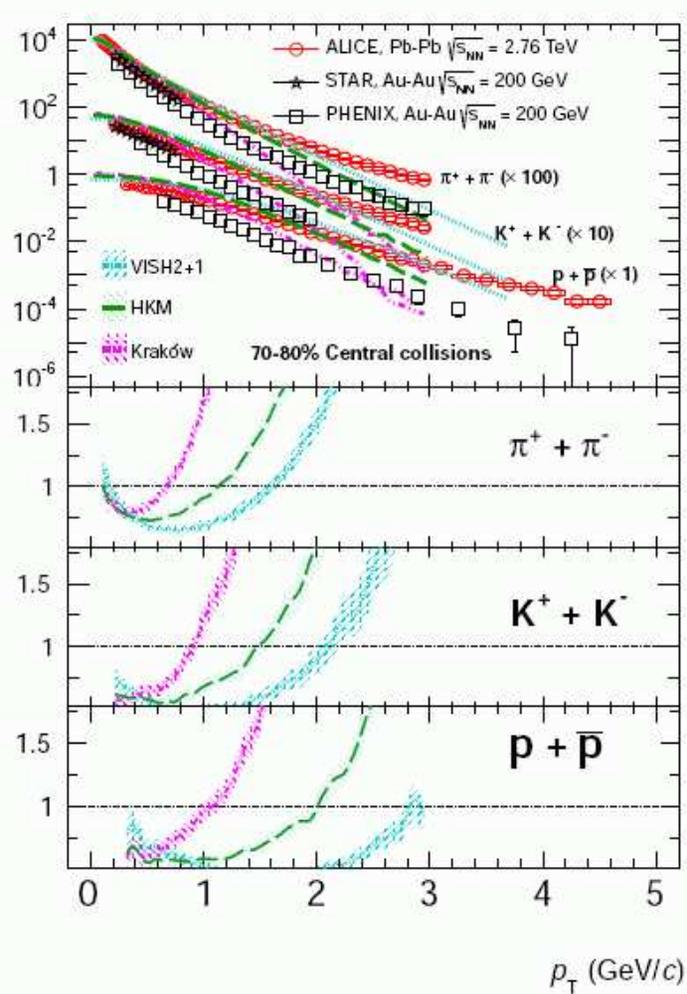
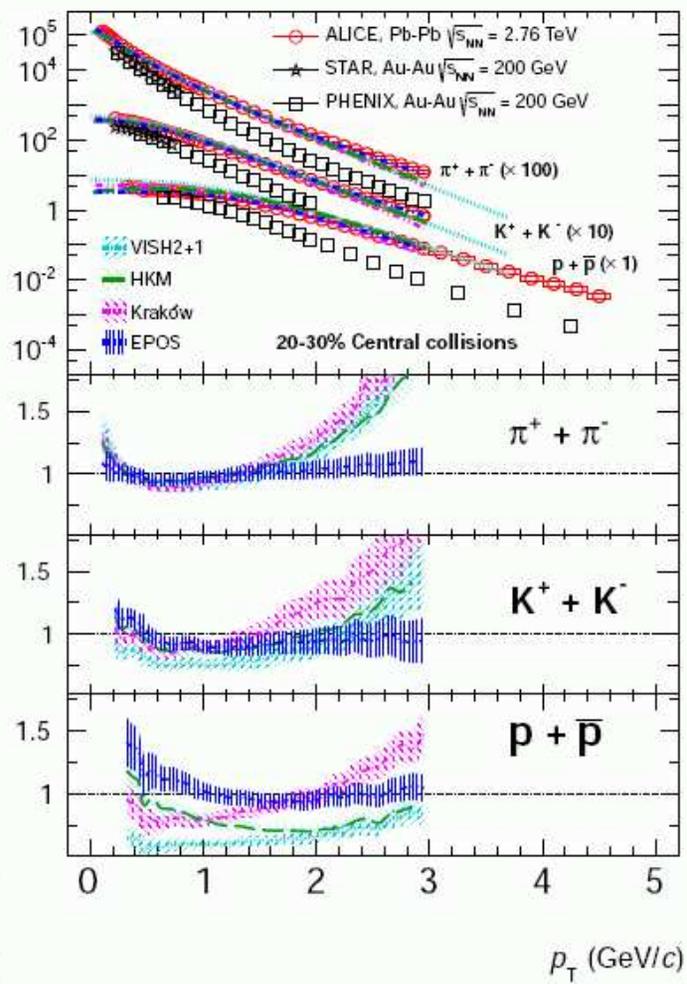
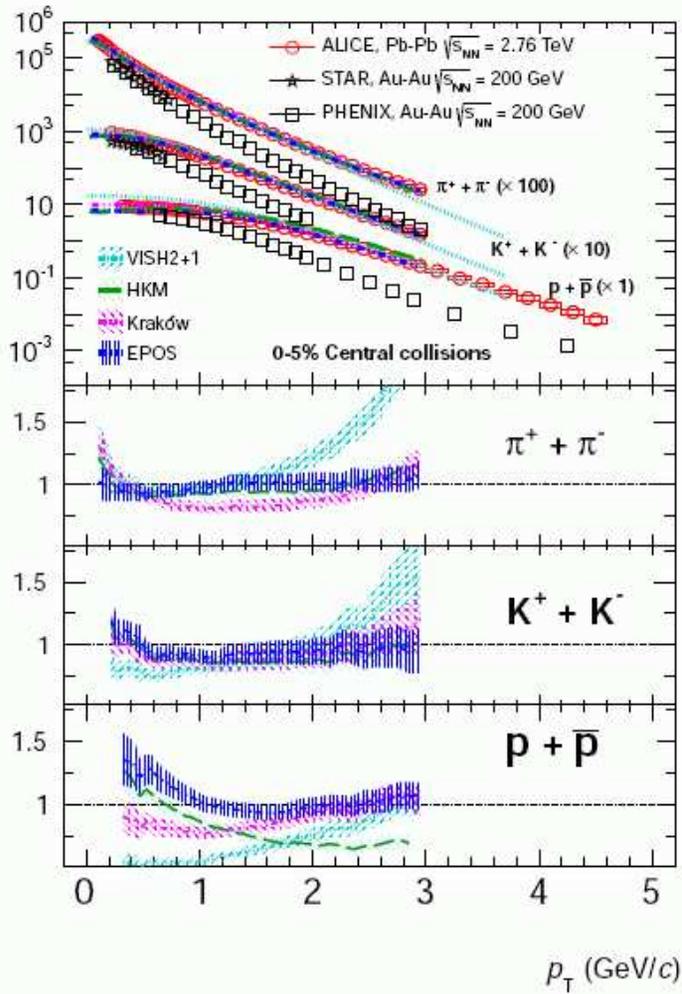
$v_2(\text{protons}) > v_2(\text{pions})$ beyond 2GeV

Summary

Analyzing p_t -spectra, ratios, and dihadrons correlations for identified hadrons:

- **pPb looks very much like a hydrodynamically expanding system**
(more clean than PbPb, where hydro and minijets heavily interact, as well as the final hadrons among themselves)





ALICE arXiv:1303.0737