

Constraining dark matter capture and annihilation cross sections by searching for neutrino signature from the Earth core

Guey-Lin Lin

National Chiao-Tung University

Taiwan

Based upon work done with Fei-Fan Lee and Yue-Lin Sming Tsai

Outline:

- a. Dark matter capture and annihilation in the Earth core
- b. Calculation of neutrino signal and background event rates
- c. Three scenarios on σ^{SI}_p and their implications on constraining $\langle\sigma v\rangle$
- d. Summary

Capture and annihilation in the Earth core

$$\Gamma_A(t) = \frac{C_c}{2} \tanh^2\left(\frac{t}{\tau_A}\right)$$

$\Gamma_A(t)$: determines the neutrino flux

C_c : capture rate depends on σ^{SI}_p and

gravitational potential of specific chemical element

$\tau_A \equiv 1/\sqrt{C_c C_A}$ equilibrium time scale with the annihilation

coefficient

$$C_A = \frac{\langle\sigma v\rangle}{V_0} \left(\frac{m_\chi}{20\text{GeV}}\right)^{2/3}$$

Hence neutrino flux depends on both σ^{SI}_p and $\langle\sigma v\rangle$

DM induced neutrino flux from the Sun is more sensitive to σ^{SD}_p

General Background

- Probing $\chi\chi \rightarrow \nu\bar{\nu}$ from Earth core with IceCube were discussed in I. F. M. Albuquerque, L. J. B. e Silva, and C. P. de los Heros, Phys. Rev. D 85, 123539 (2012), for large DM mass (few hundred GeV to TeV).
- Probing $\chi\chi \rightarrow \tau^+\tau^-$ from Earth Core with IceCube were also discussed in C. Delaunay, P. J. Fox, and G. Perez, J. High Energy Phys.05 (2009) 099 for large DM mass as well.
- We consider light DM mass as well and include DeepCore capability in our discussions. All possible annihilation channels are discussed.

DM signal and atmospheric background

$$N_{\text{signal}} = \int_{E^{\text{th}}}^{m_\chi} \sum_{i,k} \frac{d\Phi_{\nu_i}^{\text{DM}}}{dE_\nu} A_{\text{eff}}^{i,k}(E_\nu) dE_\nu d\Omega$$

$$\frac{\text{DM signal}}{\sqrt{\text{ATM background}}} = 2.0$$

$$N_{\text{background}} = \int_{E^{\text{th}}}^{E_{\text{max}}} \sum_{i,k} \frac{d\Phi_{\nu_i}^{\text{ATM}}}{dE_\nu} A_{\text{eff}}^{i,k}(E_\nu) dE_\nu d\Omega$$

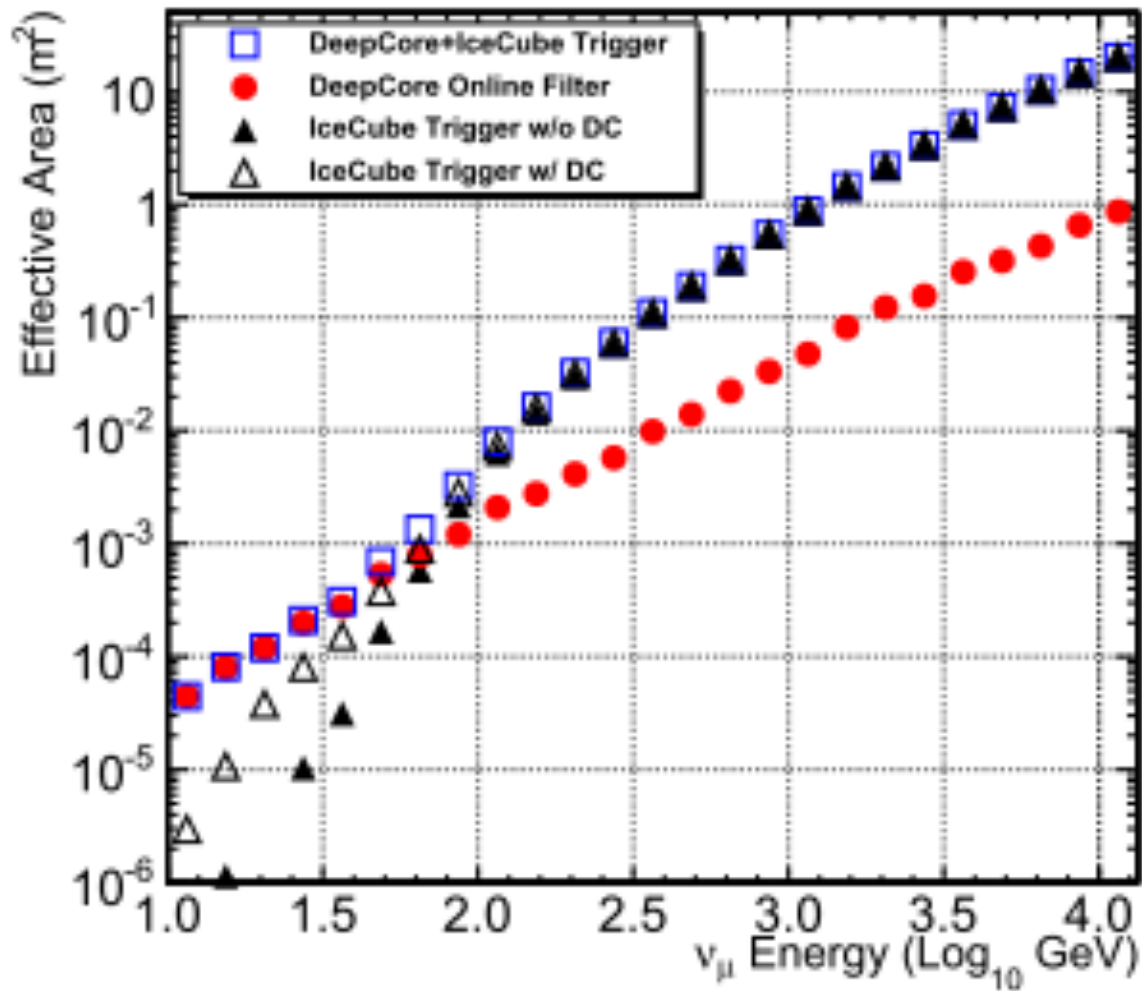
2σ detection
Significance in
5 years

i: neutrino flavor, k: CC or NC interaction

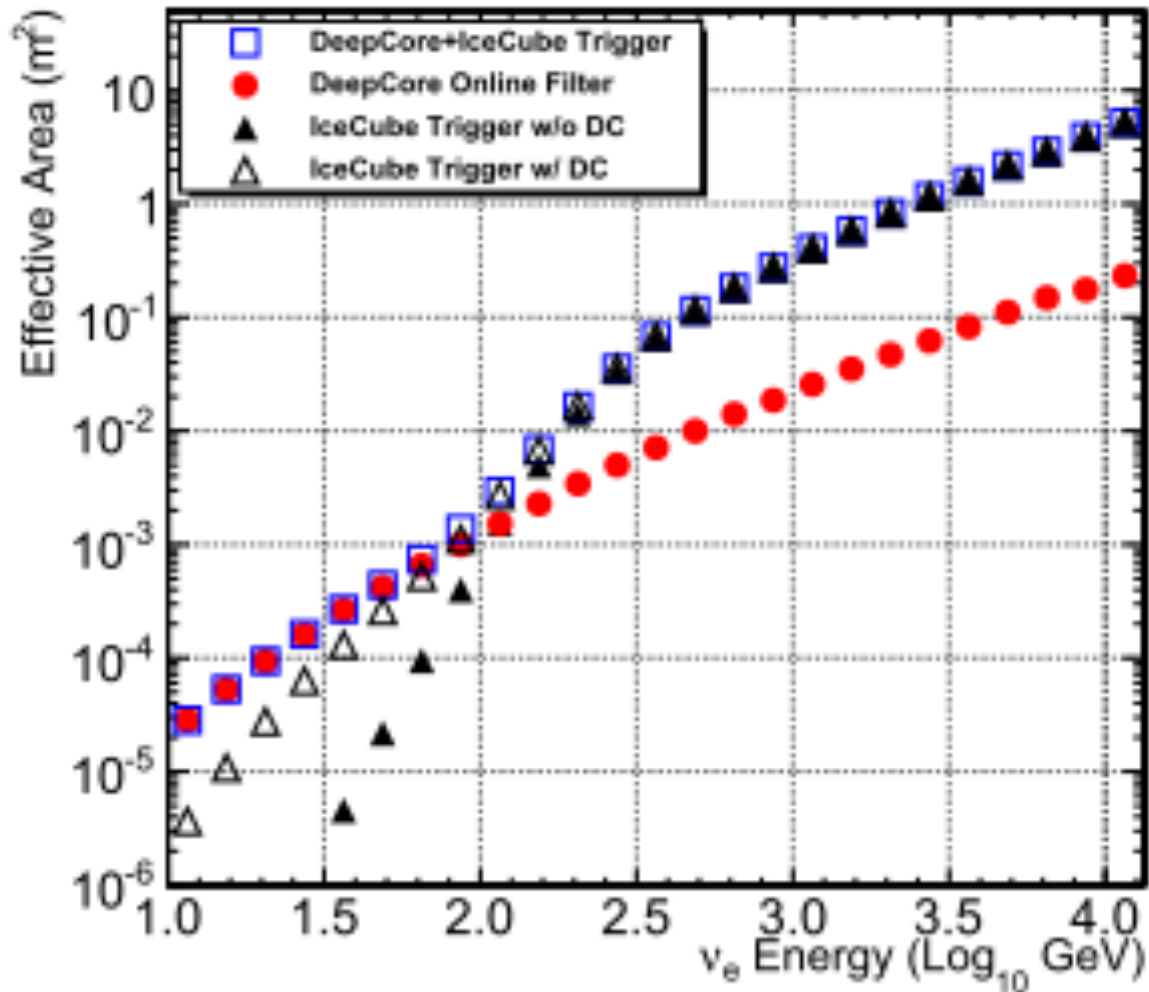
Calculation done by WIMPSIM

M. Blennow, J. Edsjo and T. Ohlsson, JCAP 0801, 021 (2008)

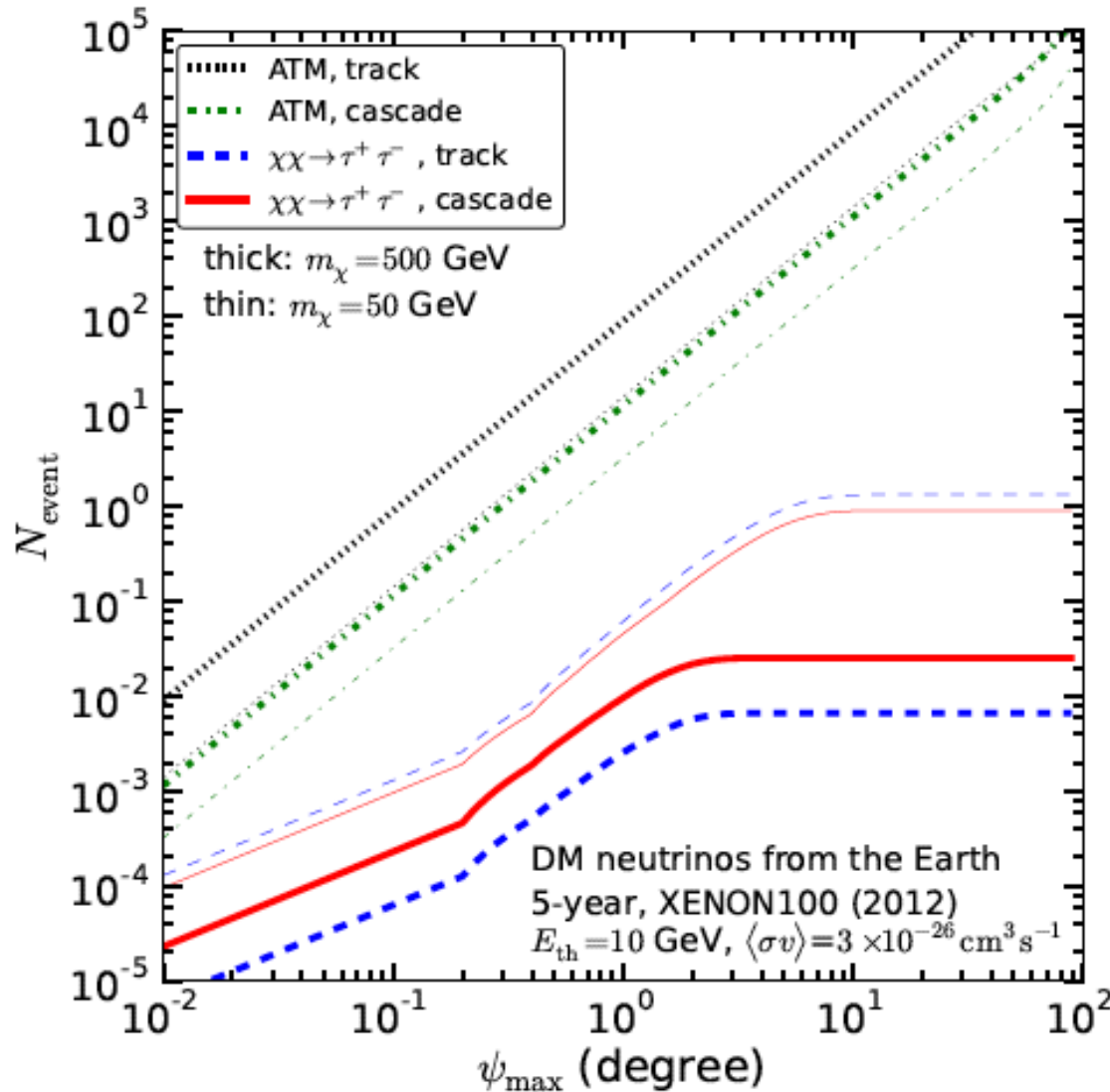
IceCube Effective Areas-Track Events



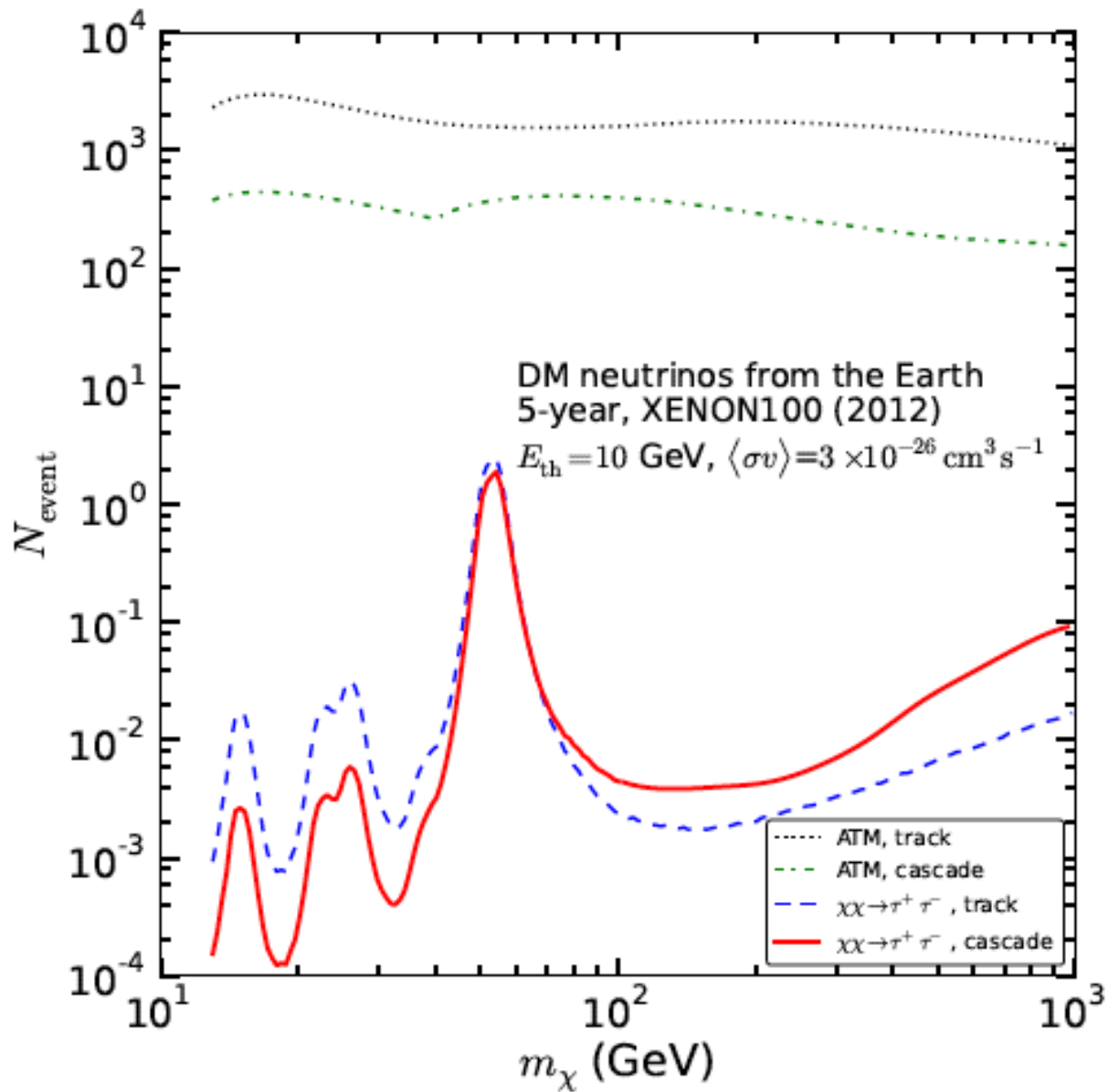
IceCube Effective Areas-Cascade Events



ψ_{\max} : observation open angle toward the Earth core

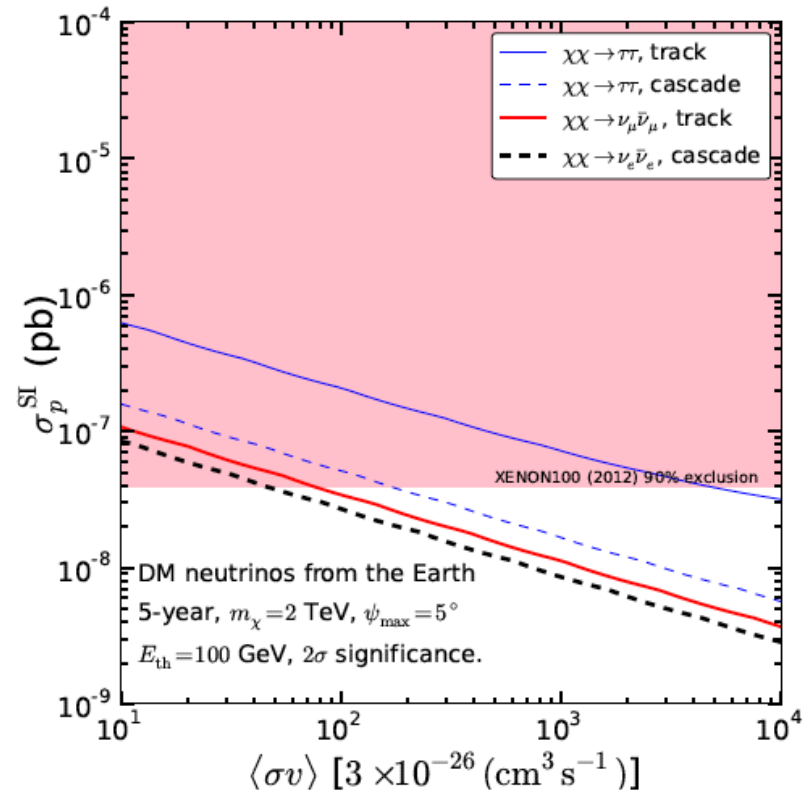
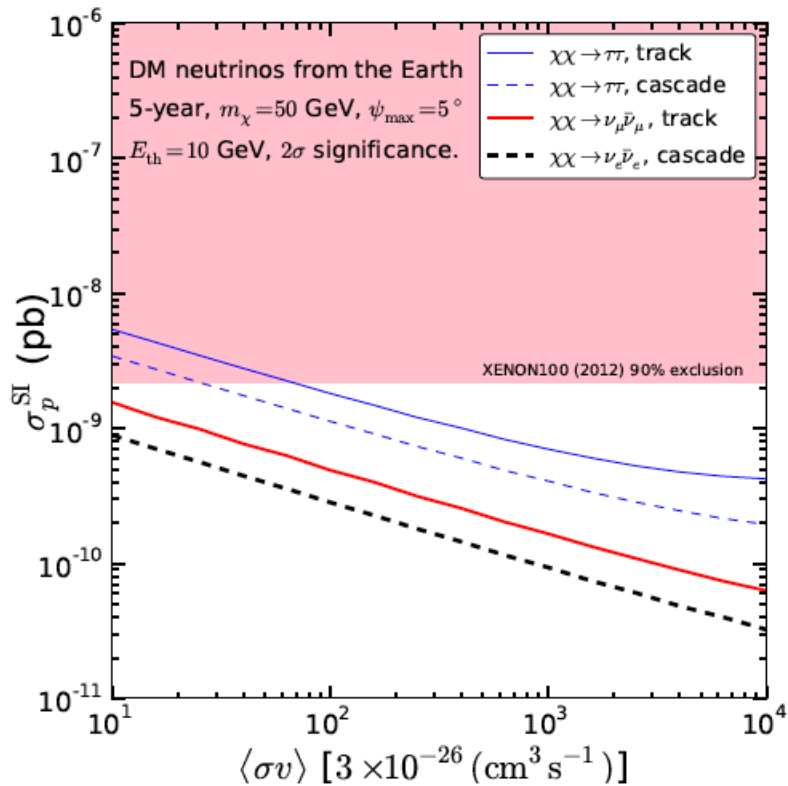


ψ_{\max}^c :
critical value of ψ_{\max}
when N_{event} no longer
increases



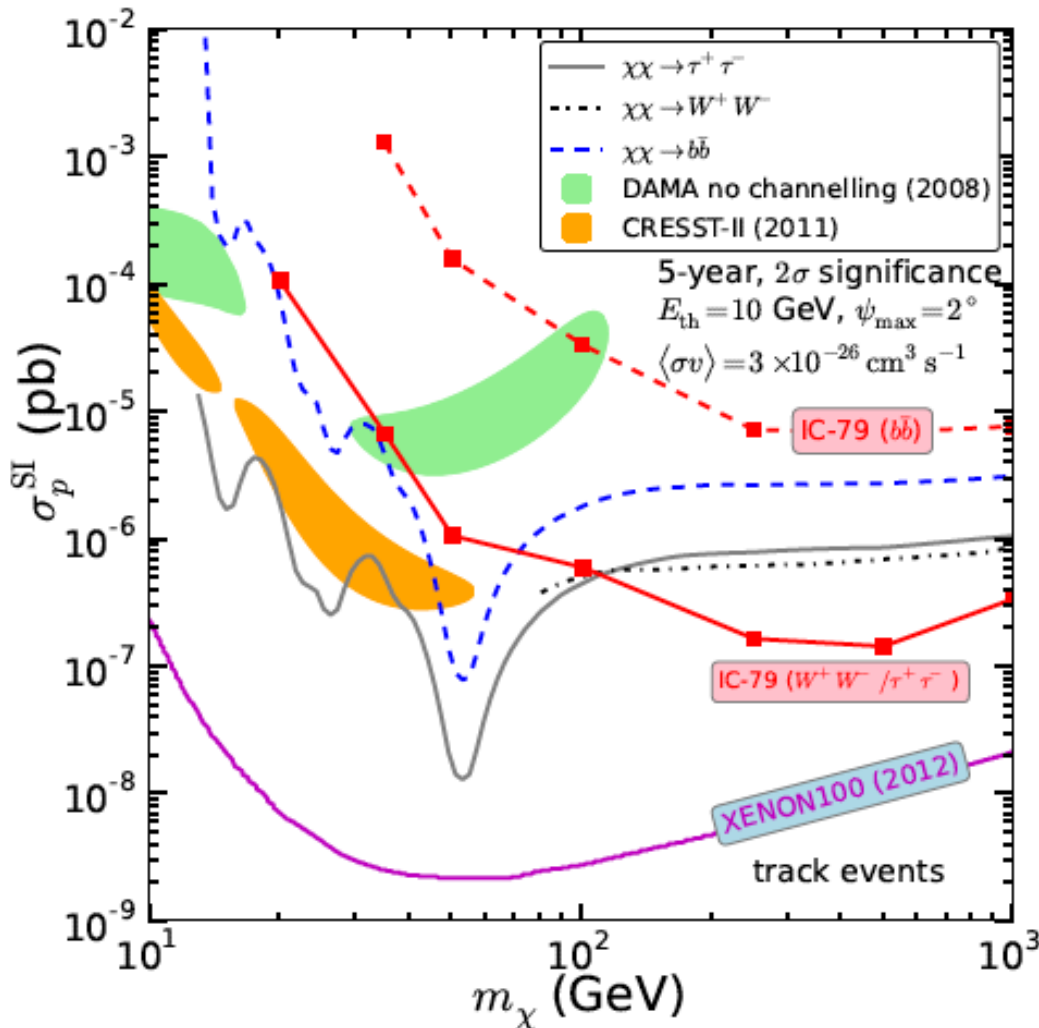
Several resonance peaks for N_{event}

Comparison of small and large m_χ



- For larger m_χ , 2σ significance requires larger $\langle\sigma v\rangle$
- Sensitivity for each channel is derived by assuming dominant branching fraction for that particular channel.

Scenario A: Neutrino observation implied by DAMA and CRESST-II



- IC-79: IceCube 79 string result from the search for DM annihilation in the Sun

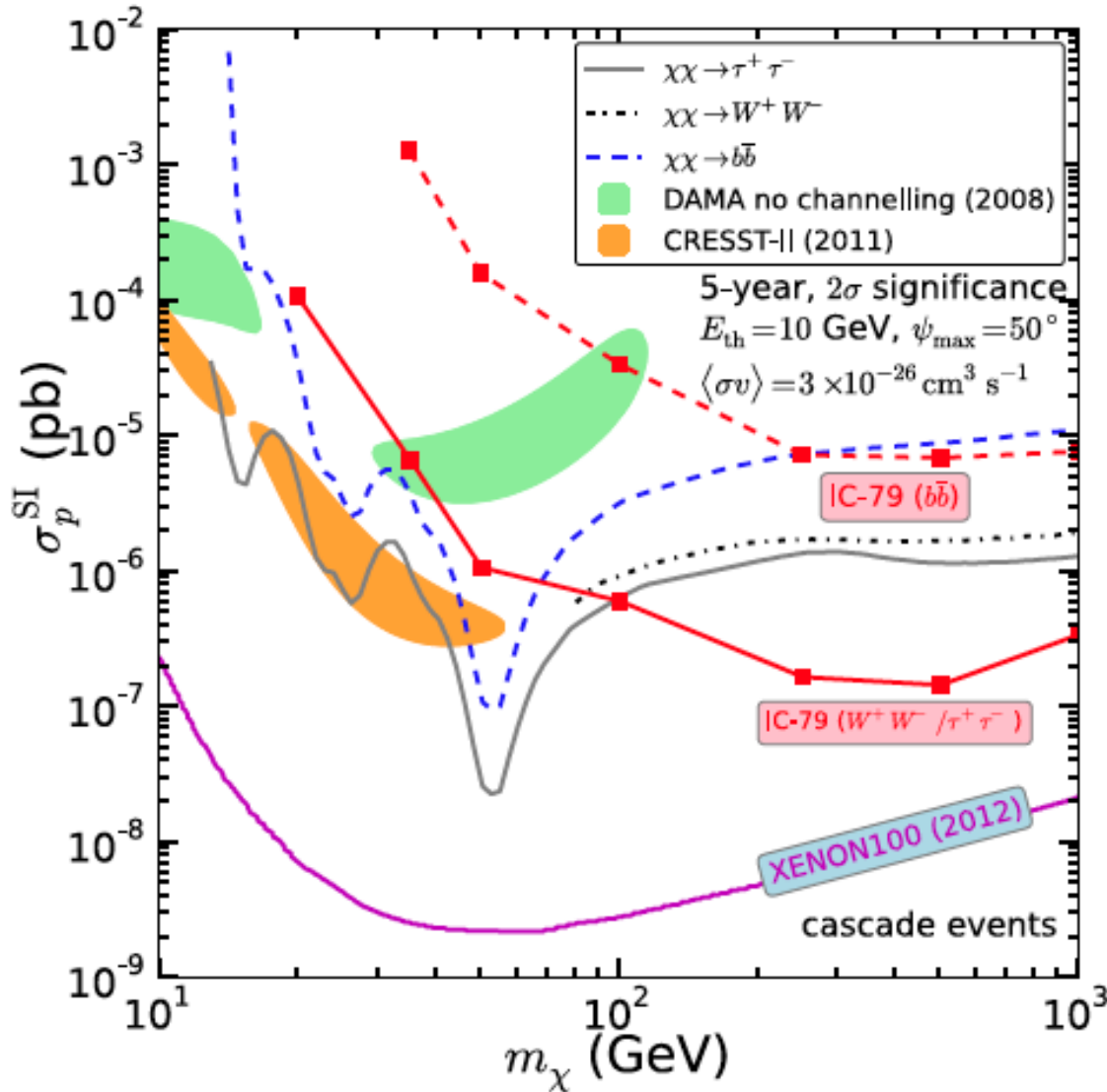
- M. G. Aartsen et al. [IceCube Collaboration], Phys. Rev. Lett. 110, 131302 (2013)

Track events

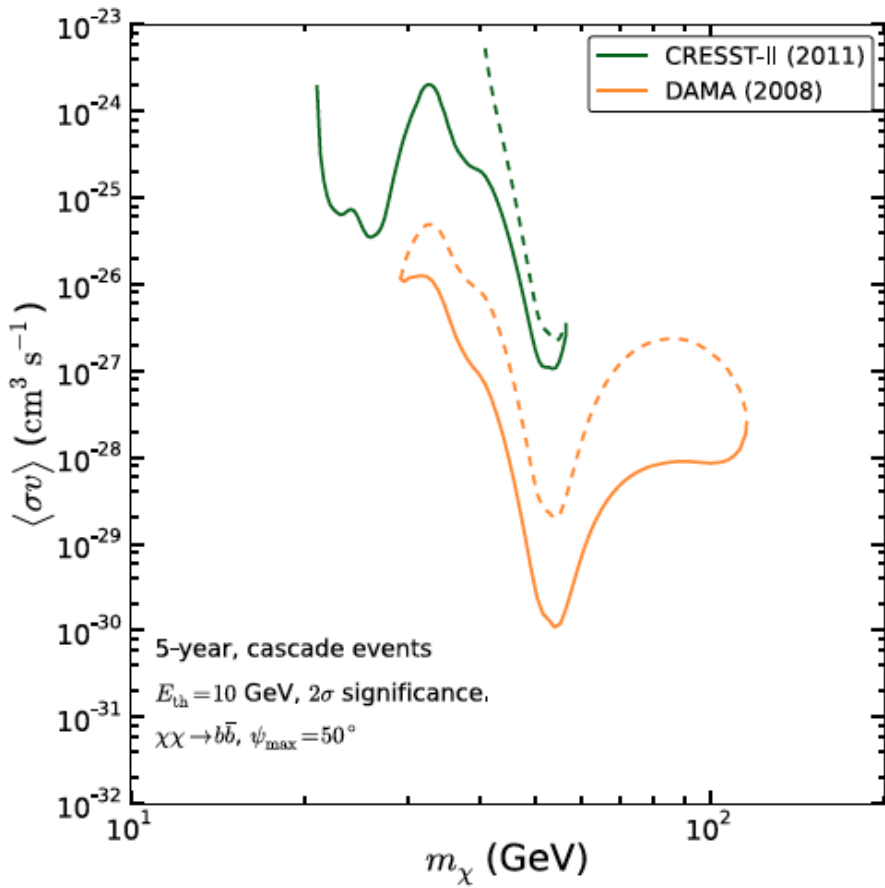
$$\Psi_{\text{max}} = 2^\circ$$

$$\langle\sigma v\rangle = 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

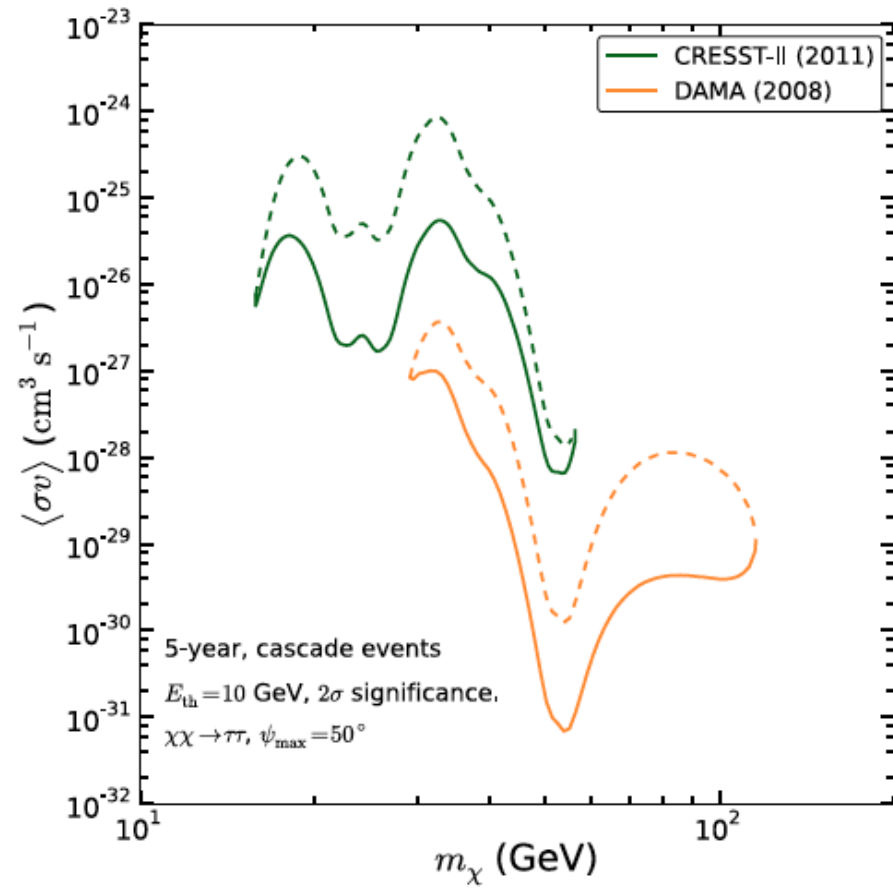
Cascade events $\Psi_{\max}=50^\circ$



One can probe $\langle \sigma v \rangle$ to values much smaller than $3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$ if σ_p^{SI} is given by DAMA or CRESST-II

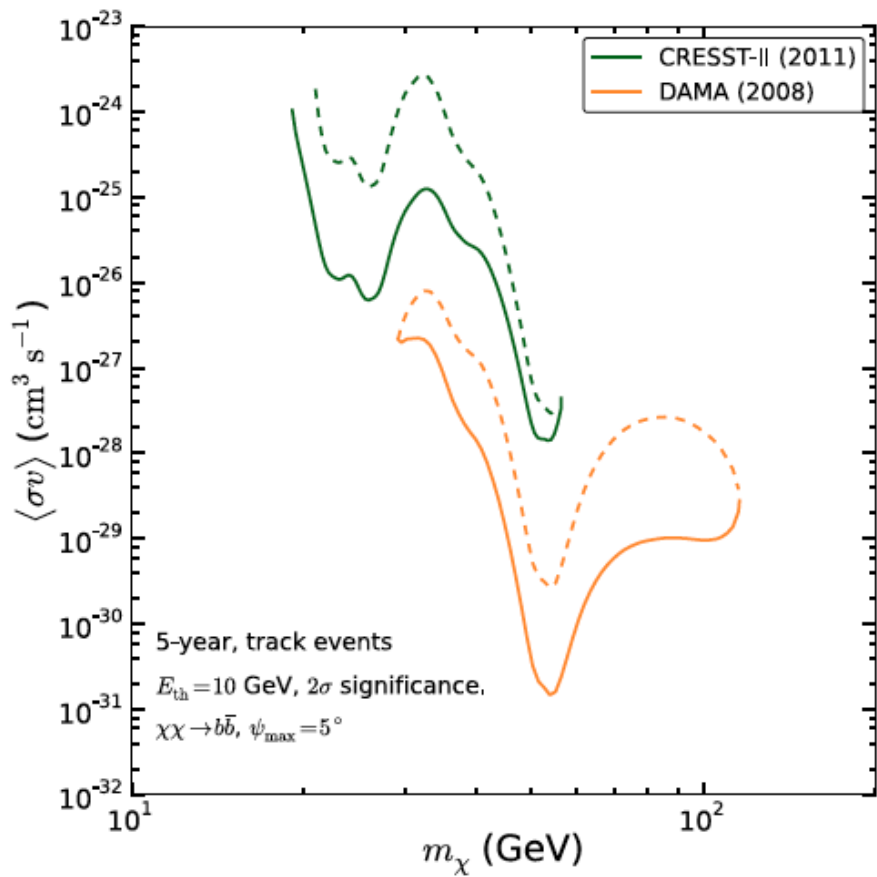


$b\bar{b}$ final state



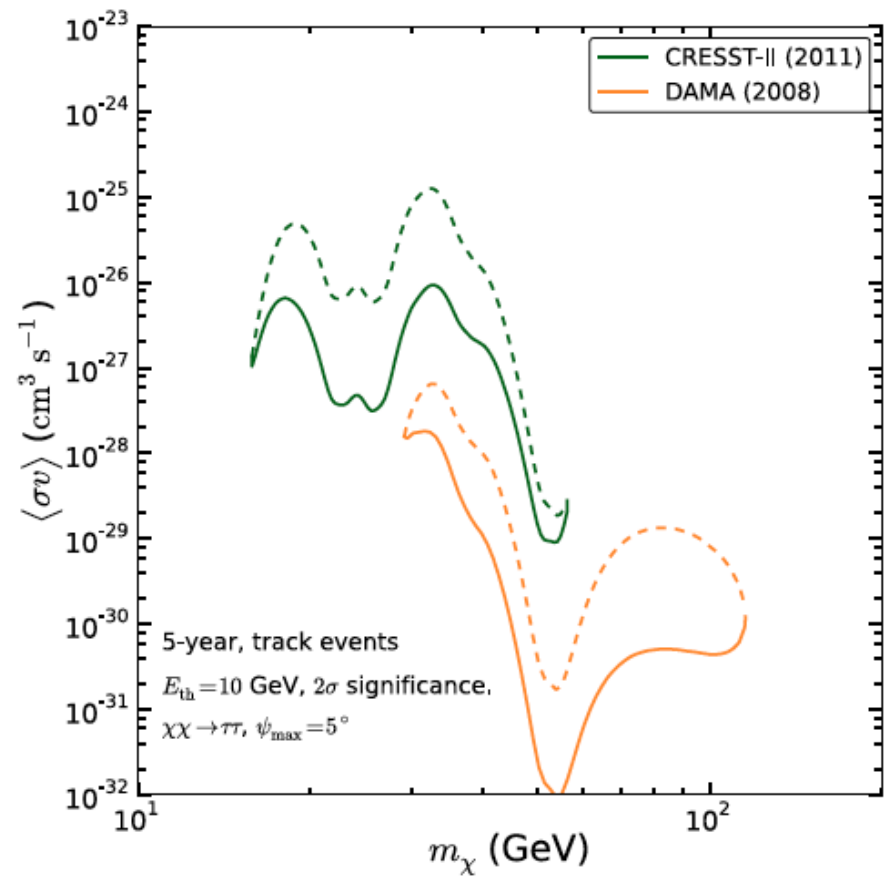
$\tau^+\tau^-$ final state

- Cascade events with $\psi_{\text{max}}=50^\circ$
- Upper limits of $\langle\sigma v\rangle$ with $\sigma_{\text{Si}_p}^{\text{Si}}$ ranges given by DAMA or CRESST-II



$b\bar{b}$ final state

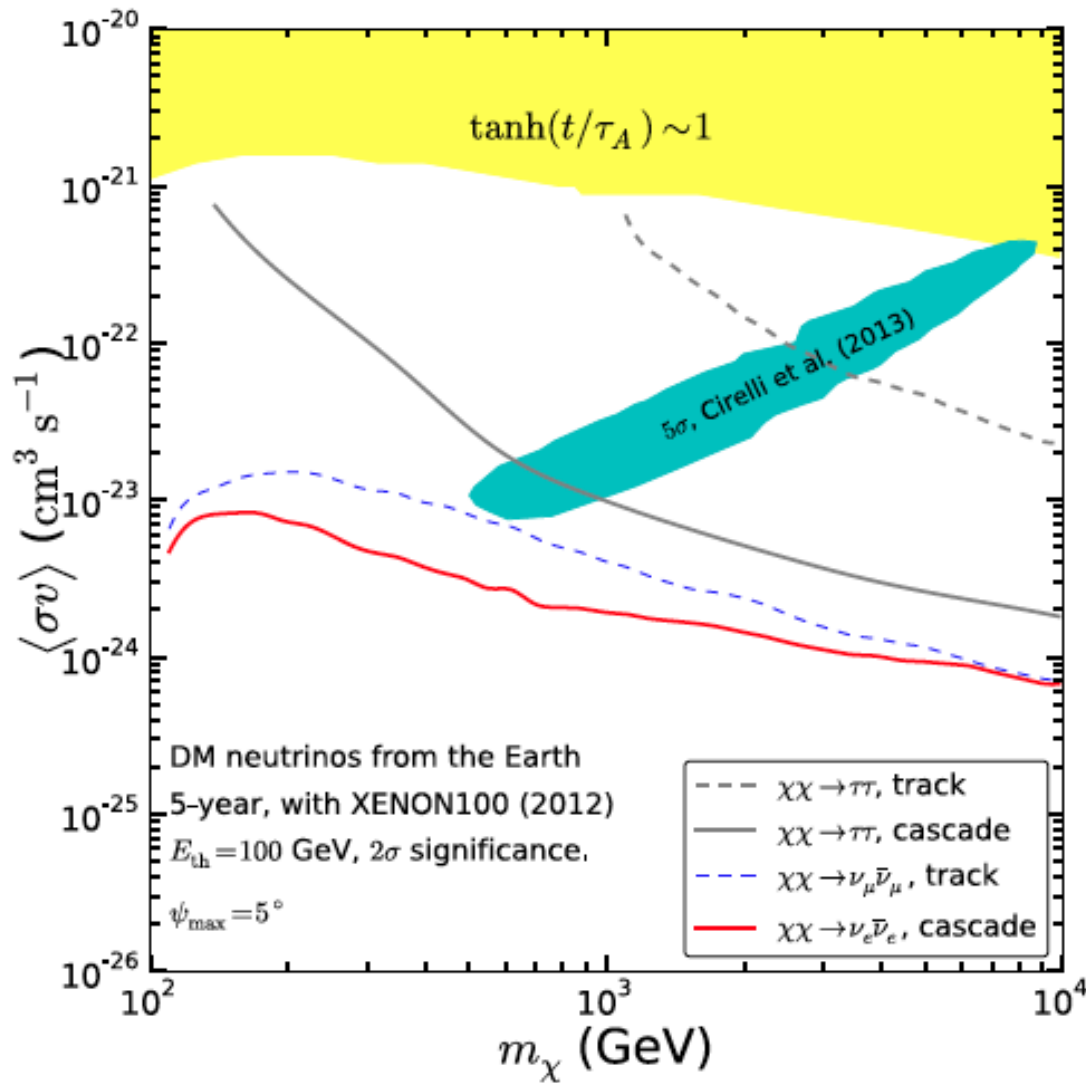
- Track events with $\psi_{\text{max}} = 5^\circ$



$\tau^+\tau^-$ final state

- Muon neutrinos in this mode comes from tau lepton decays or oscillations from tau neutrinos

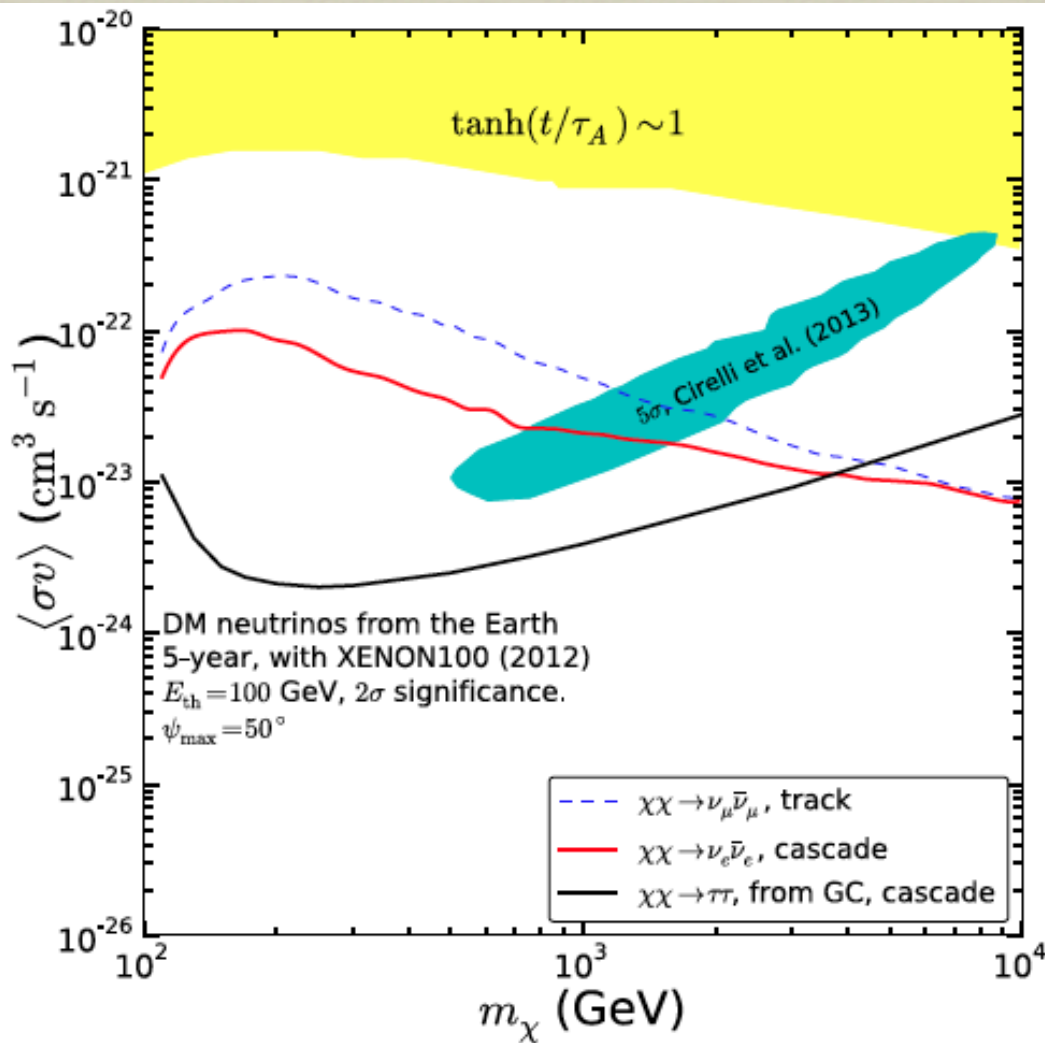
Scenario B: Probing $\langle\sigma v\rangle$ with XENON100 bound on σ^{SI}_p as input



$$\psi_{\text{max}}=5^\circ$$

- Challenging for cascade Events; however, see R. Auer, Nucl. Instrum. Methods Phys. Res., Sect. A 602,84 (2009).

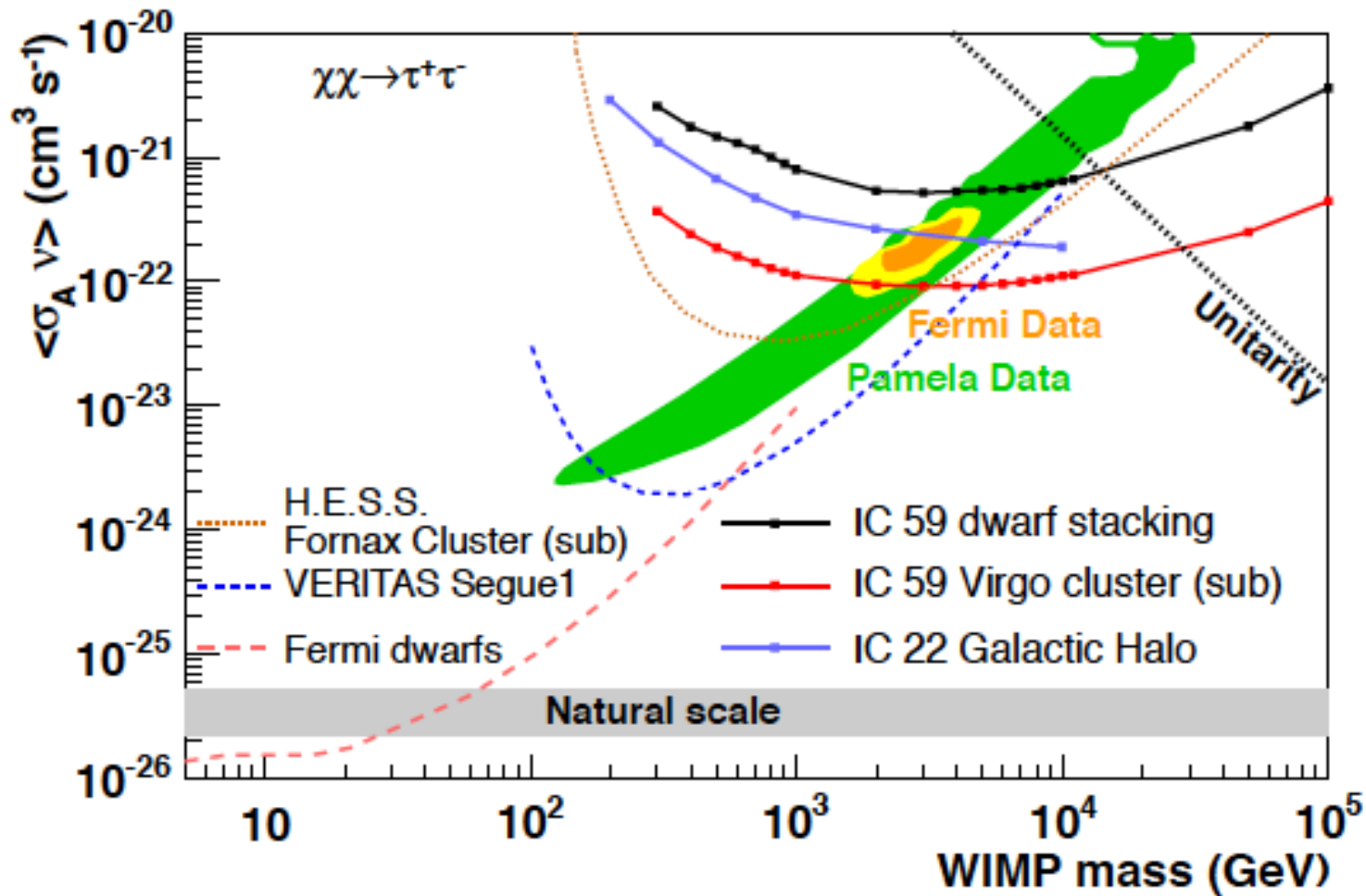
- Green region: AMS02+ PAMELA preferred region



$\psi_{\text{max}}=50^\circ$
 Only monochromatic
 channels remain

$\chi\chi \rightarrow \tau^+ \tau^-$
 from galactic center
 produces
 cascade events—
 sensitive to AMS02+
 PAMELA region

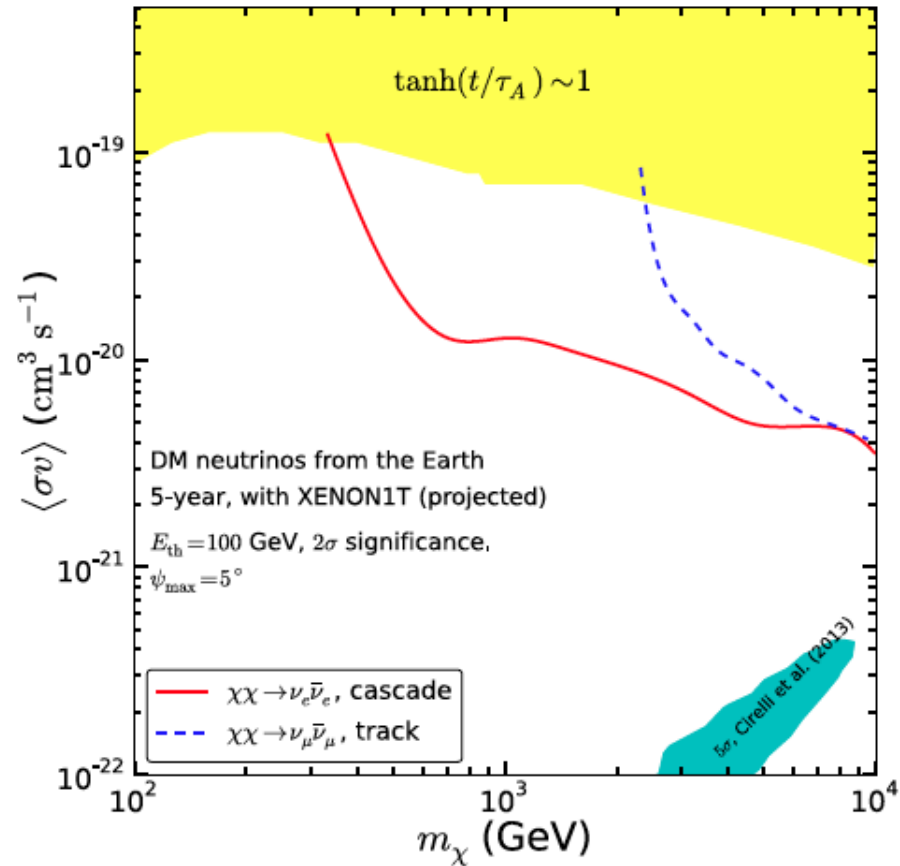
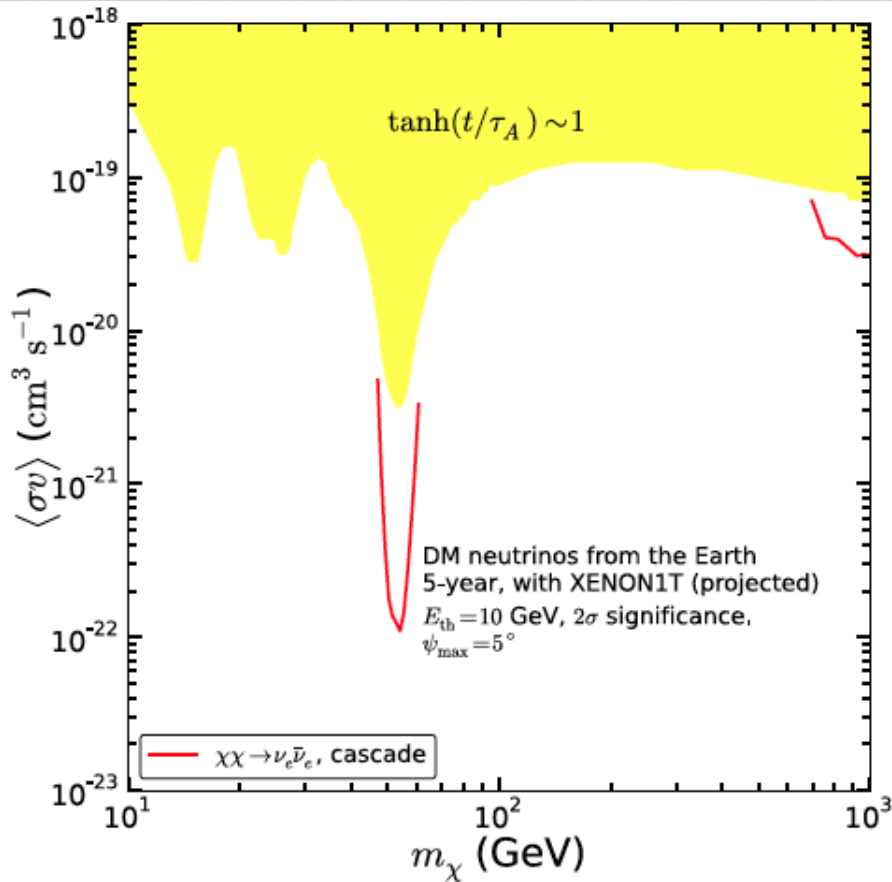
S. K. Mandal et al.,
 Phys. Rev. D81,
 043508 (2010)



IceCube new results on DM search

M. G. Aartsen *et al.*, arxiv: 1307.3473 [astro-ph.HE]

Scenario C: assuming non-detection by XENON1T



$E_{\text{th}} = 10$ GeV

$\chi\chi \rightarrow \nu_\tau \bar{\nu}_\tau$ not shown

$E_{\text{th}} = 100$ GeV

Only monochromatic channels remain

Summary:

- With $\langle\sigma v\rangle=3\times 10^{-26}\text{ cm}^3\text{s}^{-1}$, search for neutrinos produced by DM annihilation in the Earth core gives better limit on σ^{SI}_p than the current IC-79 search for DM induced neutrinos from the Sun.
- Taking DAMA and CRESST-II preferred regions as inputs for σ^{SI}_p , one can probe $\langle\sigma(\chi\chi\rightarrow\tau^+\tau^-)v\rangle$ to $10^{-32}\text{ cm}^3\text{s}^{-1}$ and $\langle\sigma(\chi\chi\rightarrow b\bar{b})v\rangle$ to $10^{-31}\text{ cm}^3\text{s}^{-1}$ for m_χ at Iron resonance.
- Taking XENON100(2012) upper limit as input for σ^{SI}_p , one can probe $\langle\sigma(\chi\chi\rightarrow\tau^+\tau^-)v\rangle$ into the region favored by PAMELA and AMS02, but the limit is not as strong as the galactic search.

Summary:

- Taking XENON1T upper limit as input for $\sigma_{\text{p}}^{\text{SI}}$, one can only probe monochromatic channel $\langle \sigma(\chi\chi \rightarrow \nu\bar{\nu})\nu \rangle$