

Dynamical Origin for the 125 GeV Higgs (a hybrid setup)

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Based on:

1. *“Hybrid dynamical EW symmetry breaking with heavy quarks and the 125 GeV Higgs”*,
M. Geller, SBS, A. Soni, arXiv:1302.2915 (PRD2013)
2. *“125 GeV Higgs in the context of four generations with two Higgs doublets”*,
M. Geller, SBS, G. Eilam, A. Soni, PRD86, 115008 (2012)

Need to account for the recently discovered 125 GeV Higgs-like ...

SM Higgs? So far looks like it but makes no sense:
where is the new physics? (hierarchy problem ...)

SUSY Higgs?
BUT: where is SUSY ?

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Light Higgs may be pseudo-Goldstone ?

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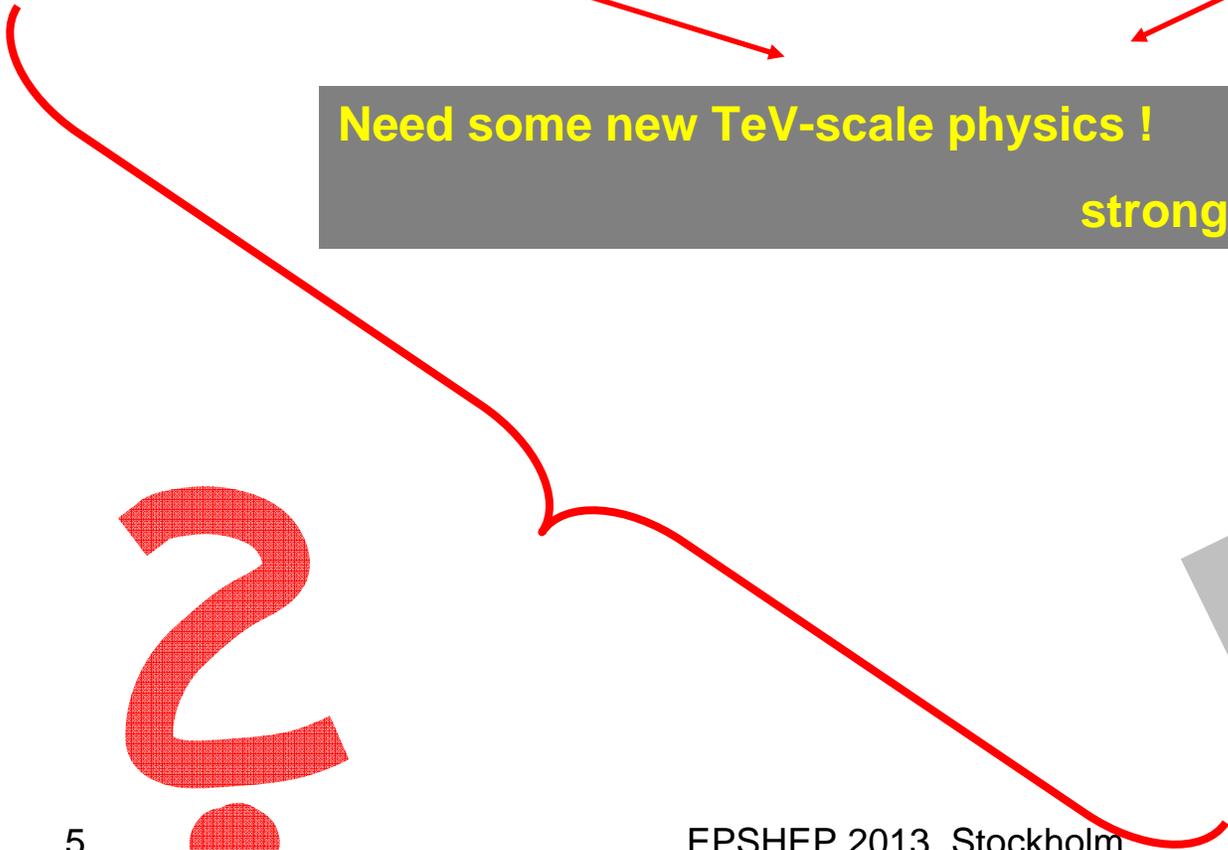
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An “old” idea: heavy fermions may be the agents for a Dynamical EWSB

Miransky, Tanabashi, Yamawaki, PLB221,177,1989 & PLA4,1043,1989;
Nambu, report#89-08,1989(unpublished); Bardeen, Hill, Lindner, PRD41,1647,1990

no need to introduce an elementary Higgs: $\langle \bar{q}' q' \rangle \neq 0 \dots$

e.g., 4th generation fermions/quarks can trigger
DEWSB through their condensations



Holdom, PRL57, 2496, 1986 & JHEP,0608,076,2006; King, PLB234,108,1990; Luty, PRD41,2893,1990;
Hill, Luty, Paschos, PRD43,3011, 1991; Hung, Isidori, PLB402,122,1997; Burdman, Da Rold,
JHEP,0712,086,2007; Hung, Xiong, NPB848,288,2011 & PLB694,430,2011; Mimura, Hou, Kohyama,
arXiv:1206.6063; Hashimoto, Miransky, arXiv:0912.4453

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no need to introduce an elementary Higgs: $\langle \bar{q}' q' \rangle \neq 0 \dots$

* top-condensation $\langle \bar{t} t \rangle \neq 0$ doesn't work
- m_t (dynamical) ~ 230 GeV
- $\Lambda \sim 10^{17}$ GeV \Rightarrow fine tuning ...

Bardeen, Hill, Lindner, PRD41,1647,1990

Adopt a simple logic ...

Be modest (no explicit model building ...):

- **ASSUME** a strongly interacting fermionic sector above Λ

and

- **TRADE** ignorance for effective couplings + boundary conditions @ Λ



NJL mechanism

At the compositeness scale, Λ , physics is described by an effective (attractive) 4-Fermi interaction of the strongly coupled fermions:

$$L_{NJL} = G_{\psi} (\bar{\psi}_L \psi_R)(\bar{\psi}_R \psi_L)$$

Solve the Gap equation (keeping only fermion loops – bubble diagrams):

If G is greater than some critical value, $G_{\psi} > G_{\text{critical}}$, then:

- **EW symmetry can be broken**
- **the field ψ acquires a mass**
- **the low-energy theory contains a scalar bound state: $S \sim \langle \psi \psi \rangle$!**

Bardeen, Hill, Lindner, PRD41,1647,1990

DEWSB with the NJL mechanism & “matching”

To get a realistic framework;

use an auxiliary field H which reproduces the 4-Fermi interaction when it is integrated out:

$$L_{aux} = g_0^\psi (\bar{\psi}_L \psi_R H + H.c) - m_0^2 H^+ H \xrightarrow{\text{integrate } H} L_{NJL} = \frac{(g_0^\psi)^2}{m_0^2} (\bar{\psi}_L \psi_R)(\bar{\psi}_R \psi_L)$$

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Then, below the cutoff Λ , H develops kinetic terms and quatric interactions (from the fermion loops) in the effective action

- H becomes a dynamical field
- the theory containing H is exactly equivalent to the theory written in terms of fermions, i.e., in terms of L_{NJL} with $G_\psi = \frac{(g_0^\psi)^2}{m_0^2}$

⇒ H is interpreted as a bound state: $H \sim \langle \psi \psi \rangle \dots$

More generally (& naturally):

should expect multiple bound states

e.g., $\langle t' t' \rangle$, $\langle b' b' \rangle$...

\Rightarrow multiple Higgs bosons below the compositeness scale

**Luty,PRD41,2893,1990; Hill,Luty,Paschos,PRD43,3011,1991; Burdman,Da Rold, JHEP,0712,086,2007;
Hung,Xiong,NPB848,288,2011 & NPB847,160,2011 & PLB694,430,2011;
Ho,Hung,Kepart,JHEP1206,45,2012; Burdman,de Lima,Matheus,PRD83,035012,2011;
Burdman,Haluch,JHEP1112,038,2011; Holdom,arXiv:1301.0329;Smetana,arXiv:1301.1554; Burdman,Da
Rold,Eboli,Matheus,PRD79,075026,2009; Hashimoto,Miransky,PRD81,055014,2010;
Hernandez,Dib,Neill,Zewekh,JHEP1202,132,2012;**

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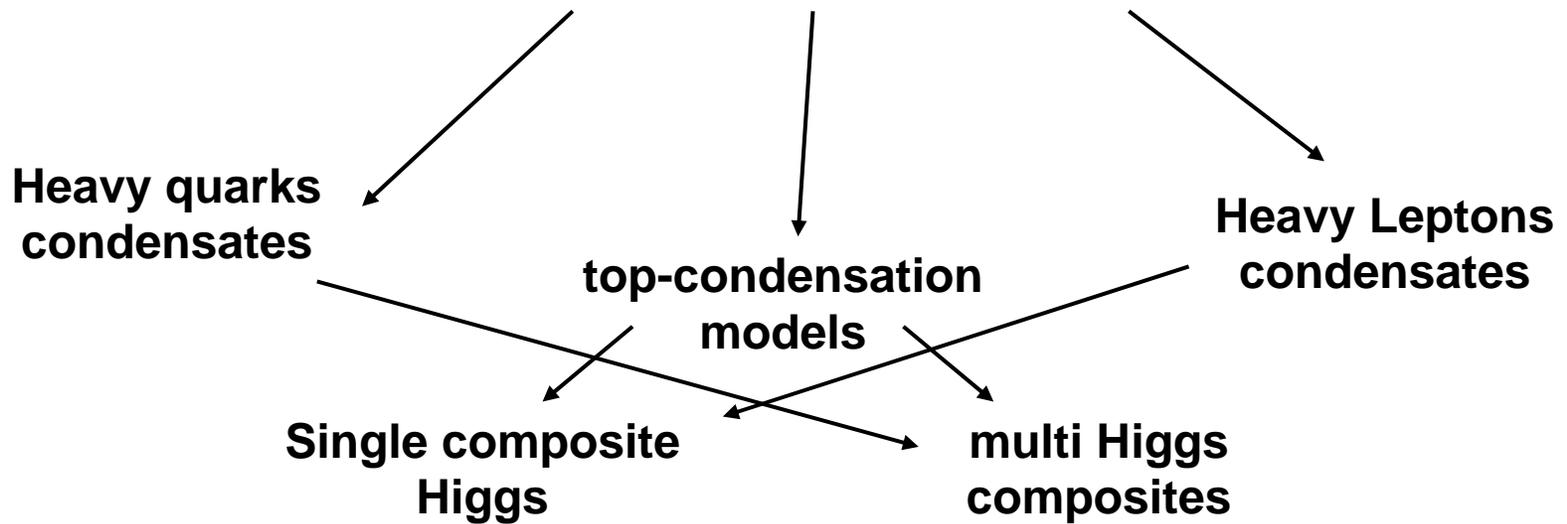
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Heavy fermions + multiple Higgs also more favorable when confronted with the data ...

M. Geller, SBS, G. Eilam, A. Soni, PRD86, 115008 (2012)
SBS, S. Nandi, A. Soni, PRD84, 035009 (2011)

SBS, M. Geller, S. Nandi, A. Soni, arXiv:1208.3195, in: "Advances in High Energy Physics" V2013

The DEWSB/NJL “spectrum”



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Heavy quarks
condensates

top-condensation
models

Heavy Leptons
condensates

Single composite
Higgs

multi Higgs
composites

$\Lambda \sim O(10^{17}) \text{ GeV}$

$m_f \sim v_{EW}$
e.g., top-condensation

$\Lambda \sim O(1) \text{ TeV}$

$m_f \sim 500 \text{ GeV}$
e.g., “4th gen” ...

A major caveat:

In the “conventional” NJL models for DEWSB the typical mass of the composite H $\sim \langle \psi\psi \rangle$ tends to lie in the range:

$$m_\psi < m_{\langle \bar{\psi}\psi \rangle} < 2m_\psi$$



Too heavy to account for the recently observed 125 GeV Higgs-like particle

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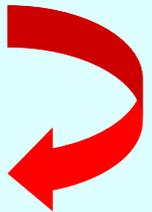


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Our proposed HYBRID “solution” leads to $m_H \sim O(m_W)$

Hybrid setup with heavy quarks [& $\Lambda \sim O(1)$ TeV] :

add an unconstrained scalar (“fundamental”, ϕ_l) at the compositeness scale, where the super-critical 4-Fermi operators form an additional heavy composite ϕ_h



Hybrid DEWSB with heavy quarks and the 125 GeV Higgs

M. Geller, SBS, A. Soni, arXiv:1302.2915

As a “toy framework”: use a chiral 4th gen. doublet, assumed to be charged under some new strong interaction that dynamically break EW symmetry:

At the compositeness scale Λ : *The bare SM Lag. at Λ (with ϕ_1)*

$$\mathcal{L} = \mathcal{L}_{SM}(\Lambda) + G_{t'} \bar{Q}'_L t'_R \bar{t}'_R Q'_L + G_{b'} \bar{Q}'_L b'_R \bar{b}'_R Q'_L + G_{t'b'} (\bar{Q}'_L b'_R \bar{t}'_R i\tau_2 Q'^e_L + h.c.)$$

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The above Lag/theory can be reproduced by introducing at Λ an auxiliary Higgs doublet ϕ_h which couples **ONLY** to the 4th gen. quarks:

$$\mathcal{L}_{q'}(\Lambda) = g_{b'}^0 (\bar{Q}'_L \Phi_h b'_R + h.c.) + g_{t'}^0 (\bar{Q}'_L \tilde{\Phi}_h t'_R + h.c.) - (\mu_h^0)^2 \Phi_h^\dagger \Phi_h$$

The full theory at Λ is then described by: $\mathcal{L}(\Lambda) = \mathcal{L}_{SM}(\Lambda) + \mathcal{L}_{q'}(\Lambda) + (\mu_{he}^0)^2 (\Phi_h^\dagger \Phi_e + h.c.)$

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$$\phi_h \sim g_{t'} \langle \bar{Q}'_L t'_R \rangle + g_{b'} \langle \bar{Q}'_L b'_R \rangle$$

Integrating out ϕ_h , recovers the 4-Fermi Lag with:

$$G_{t'} = \frac{(g_{t'}^0)^2}{(\mu_h^0)^2}, \quad G_{b'} = \frac{(g_{b'}^0)^2}{(\mu_h^0)^2}, \quad G_{t'b'} = -\frac{g_{t'}^0 g_{b'}^0}{(\mu_h^0)^2}$$

Hybrid DEWSB with heavy quarks and the 125 GeV Higgs

ϕ_h - couples to new heavy quarks & leptons – **composite**

⇒ Q' have purely dynamical masses ...

ϕ_l - mass of SM fermions & CKM flavor structure – **“fundamental”/unconstrained**

**“fundamental” = unconstrained at Λ
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At energies below Λ :

ϕ_h acquires a kinetic term as well as LARGE self interactions from the heavy fermion loops (ϕ_l quartic do not receive such large corrections !)



\Rightarrow the theory behaves like a 2HDM below Λ

(specifically: like the 4G2HDM of SBS, Nandi, Soni PRD2011)

The low-energy Hybrid 4G2HDM (h4G2HDM)

$$\begin{aligned}
 V_{h4G2HDM}(\Phi_h, \Phi_\ell) = & \mu_\ell^2 \Phi_\ell^\dagger \Phi_\ell + \mu_h^2 \Phi_h^\dagger \Phi_h - \mu_{h\ell}^2 (\Phi_h^\dagger \Phi_\ell + h.c.) \\
 & + \frac{1}{2} \lambda_\ell (\Phi_\ell^\dagger \Phi_\ell)^2 + \frac{1}{2} \lambda_h (\Phi_h^\dagger \Phi_h)^2 + \lambda_3 (\Phi_h^\dagger \Phi_h) (\Phi_\ell^\dagger \Phi_\ell) + \lambda_4 (\Phi_h^\dagger \Phi_\ell) (\Phi_\ell^\dagger \Phi_h)
 \end{aligned}$$

how dynamical:

$$Z_\phi(\Lambda) \rightarrow 0$$

from the f' -loops

$$Z_\phi |D\phi|^2$$

Matching the 4-Fermi theory to the h4G2HDM at Λ :

⇒ Solve RGE of the h4G2HDM with the compositeness boundary conditions :

$$g_{q'}(\Lambda) \rightarrow \infty, \quad \lambda_{h,3,4}(\Lambda) \rightarrow \infty, \quad \frac{\lambda_h(\Lambda)}{g_{q'}^4(\Lambda)} \rightarrow 0, \quad \frac{\lambda_{3,4}(\Lambda)}{g_{q'}^2(\Lambda)} \rightarrow 0$$

i.e., the composite theory is effectively a strongly coupled

Higgs-Yukawa and Higgs-quartic systems ...

While $\lambda_\ell(\mu \rightarrow \Lambda) \rightarrow \lambda_\ell^{(0)}$, where $\lambda_\ell^{(0)}$ is a free parameter

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\hookrightarrow how dynamical: $Z_\phi | \langle \phi \rangle |^2$
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 \hookrightarrow from the f' -loops

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Higgs mass parameters (minimizing the potential):

$$\mu_\ell^2 \simeq \frac{\mu_{h\ell}^2}{t_\beta} - \frac{v^2}{2} c_\beta^2 \lambda_\ell, \quad \mu_h^2 \simeq t_\beta \mu_{h\ell}^2 - \frac{v^2}{2} s_\beta^2 \lambda_h$$

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how dynamical: $Z_\phi \rightarrow 0$
 $Z_\phi \sim |\phi|^2$
 from the t' -loops

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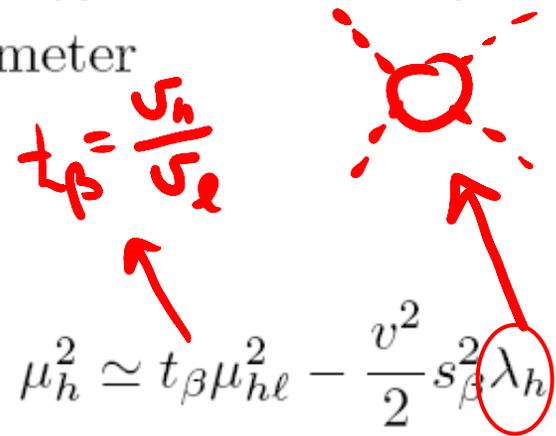
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RGE & spectrum

Relevant new RGE (rest is SM-like ...)
(& $\lambda_{3,4} \ll \lambda_h, \lambda_1$)

$$D \equiv 16\pi^2 \mu \frac{d}{d\mu}$$
$$Dg_{q'} \approx 6g_{q'}^3,$$
$$D\lambda_h \approx 4\lambda_h (3\lambda_h + 6g_{q'}^2) - 24g_{q'}^4$$

$$g_{q'}(\mu) = \sqrt{\frac{4\pi^2}{3\ln\frac{\Lambda}{\mu}}}, \quad \lambda_h(\mu) = \frac{4\pi^2}{3\ln\frac{\Lambda}{\mu}}$$

RGE & spectrum

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The compositeness scale ($m_{q'} = v_h g_{q'}(\mu = m_{q'}) / \sqrt{2}$):

$$\Lambda \approx m_{q'} \cdot \exp\left(\frac{2\pi^2 (s_\beta v)^2}{3m_{q'}^2}\right) \xrightarrow{m_{q'} \sim O(500 \text{ GeV}), t_\beta \sim O(1)} \Lambda \sim 1 - 1.5 \text{ TeV}$$

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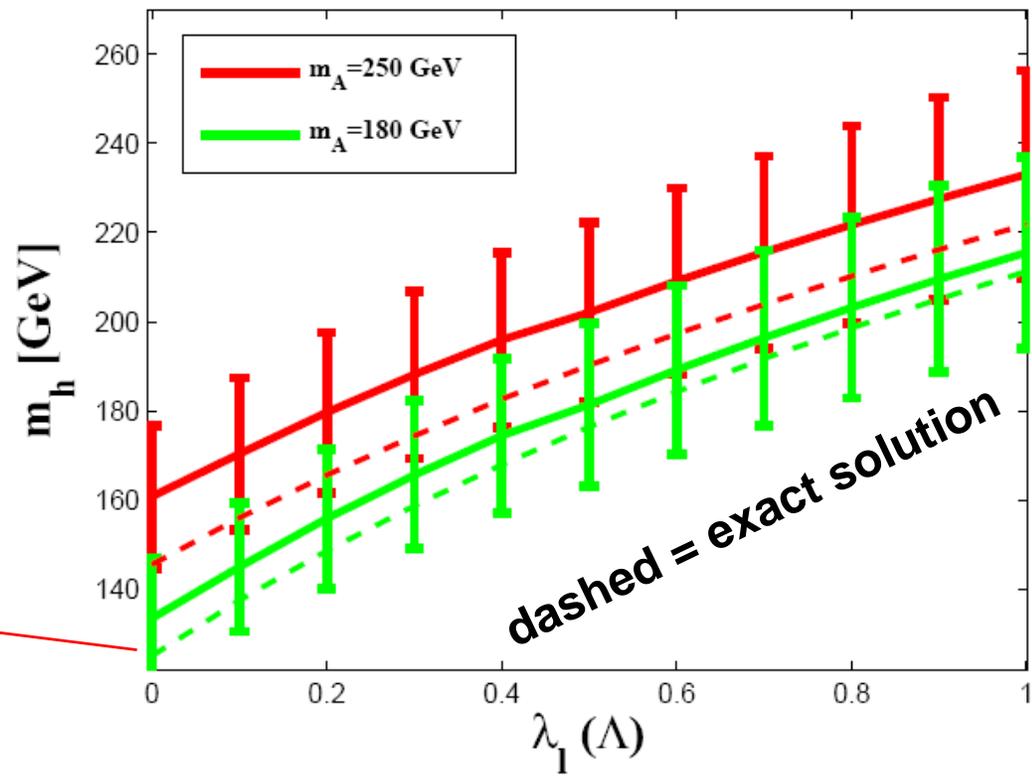
Recall: the top-condensate scenario had a severe fine-tuning problem:

$$\Lambda \sim m_t \cdot \exp\left(\frac{16\pi^2 v^2}{9m_t^2}\right) \sim 10^{17} \text{ GeV} \quad \text{!!!}$$

Higgs masses

A light Higgs **requires** $\lambda_1(\Lambda) \rightarrow 0$:

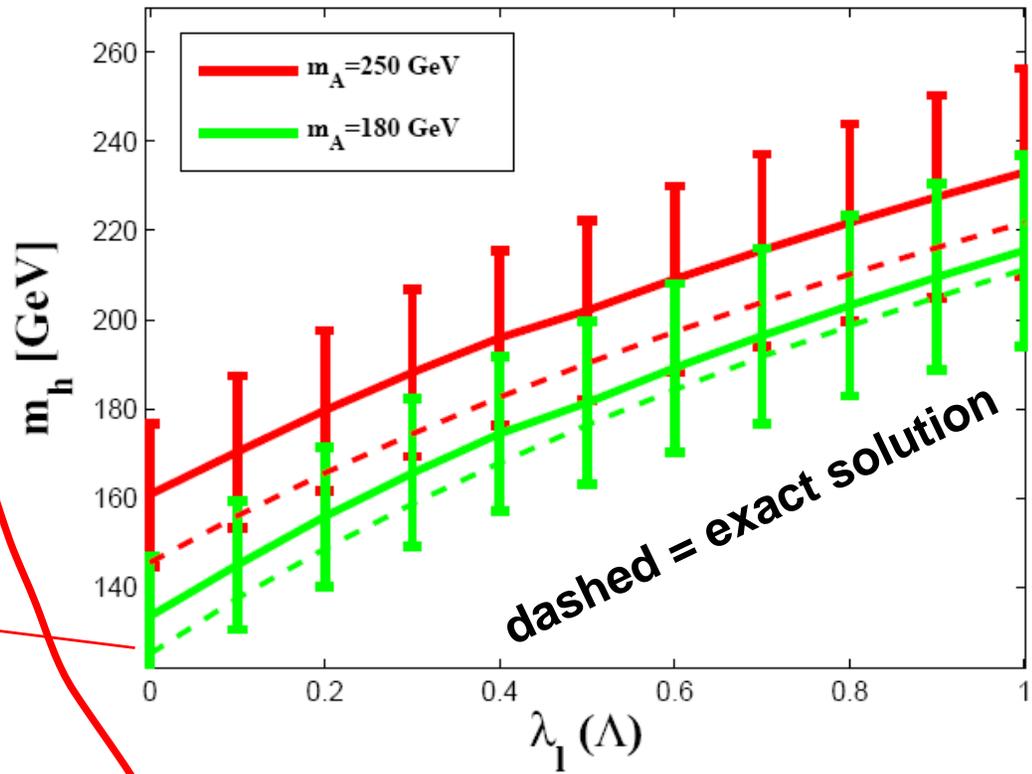
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\rightarrow pseudo-Goldstone!!

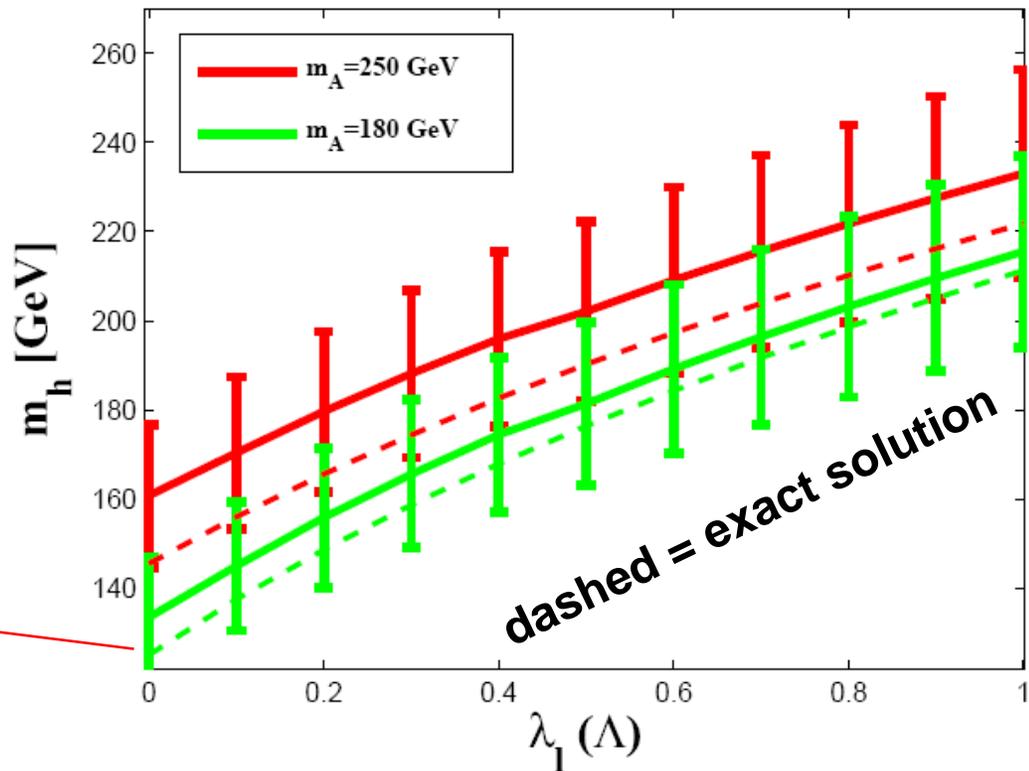
$\rightarrow \lambda_2(\Lambda) \rightarrow 0$

$\hookrightarrow \lambda_2(m_w) \ll \lambda_h(m_w)$

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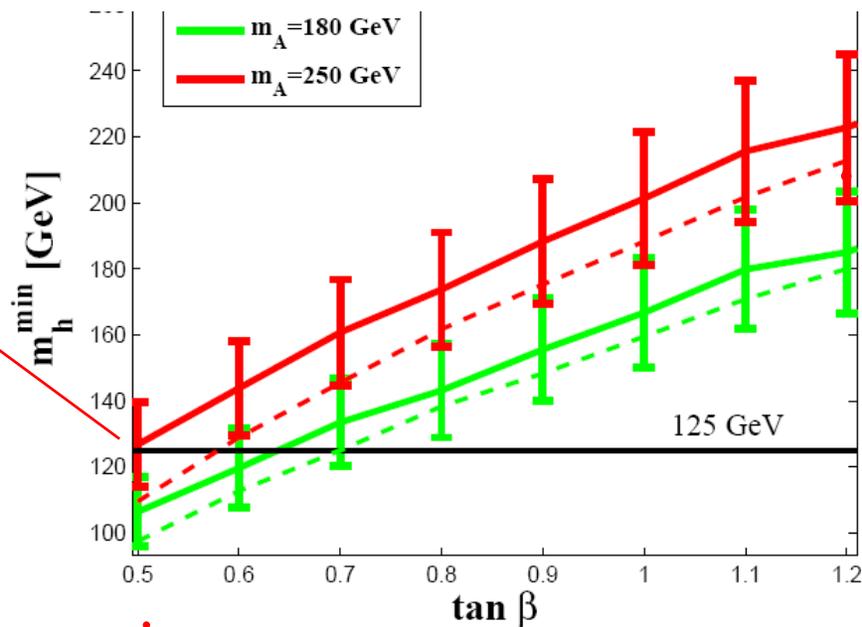
$$m_h \xrightarrow{\lambda_1(\Lambda) \rightarrow 0 \text{ \& \; } \tan \beta \sim O(1)} \frac{m_A}{\sqrt{2}}$$



$$m_h \sim 125 \text{ GeV } (\pm 10\%)$$

Requires:

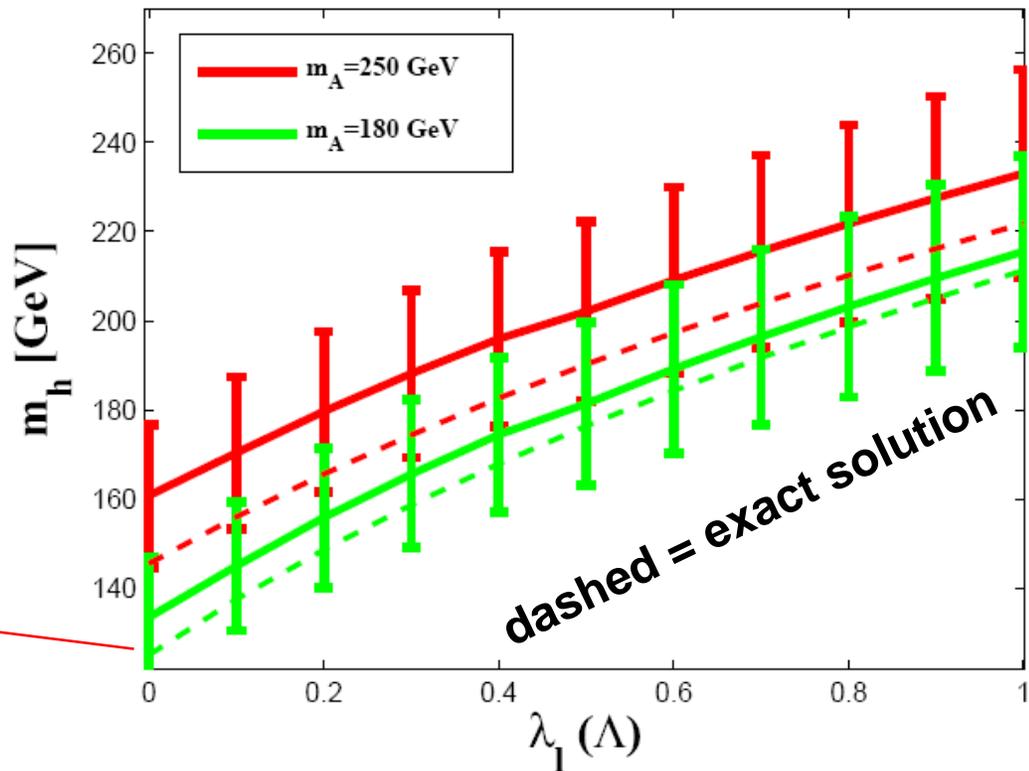
$$\tan \beta \lesssim 0.7 \text{ and } m_A \lesssim 250 \text{ GeV}$$



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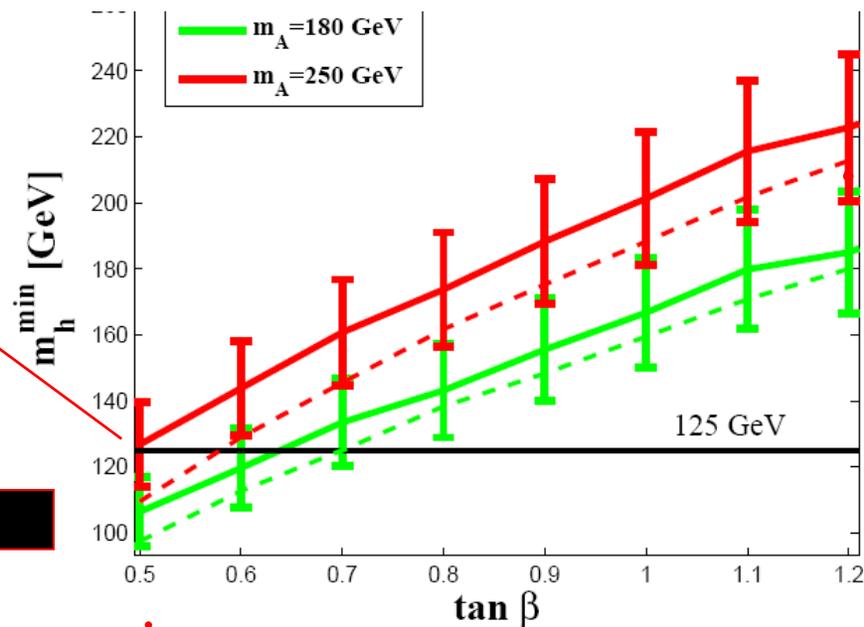
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Reproduces all observed Higgs signals, see M. Geller's poster

Higgs masses cont.

Heavy CP-even Higgs is **mostly composite**:

$$m_H \approx v \cdot \sqrt{\frac{\lambda_h}{2}} \xrightarrow{m_{q'} \sim O(500 \text{ GeV})} \sim 500 \pm 100 \text{ GeV}$$



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Higgs mixing angle $\sim O(10^0)$

$$h = \cos \alpha \cdot \text{Re}(\Phi_\ell^0) - \sin \alpha \cdot \text{Re}(\Phi_h^0)$$

$$H = \cos \alpha \cdot \text{Re}(\Phi_h^0) + \sin \alpha \cdot \text{Re}(\Phi_\ell^0)$$

Light Higgs mostly "fundamental"

$$\tan 2\alpha \simeq \left(\cot 2\beta - \frac{v^2}{2\mu_{h\ell}^2} \cdot (s_\beta^2 \lambda_h - c_\beta^2 \lambda_\ell) \right)^{-1}$$

Summary & concluding remarks

SM Higgs – makes no sense:
where is the new physics? (hierarchy problem ...)

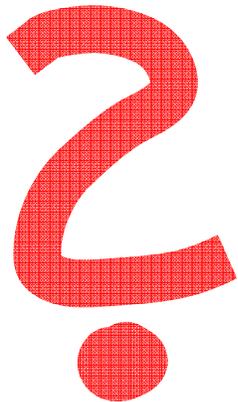
SUSY Higgs? - where is it ?

Need some new TeV-scale physics !

strong dynamics?

Difficult to realize with a
light composite Higgs ...

Light Higgs may be
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SM Higgs – makes no sense:
where is the new physics? (hierarchy problem ...)

SUSY Higgs? - where is it ?

Need some new TeV-scale physics !

strong dynamics?

Hybrid setup:
"fundamental" + composite scalars at TeV-scale



Needs heavy fermions !

Difficult to realize with a
light composite Higgs ...

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The $h4G2HDM$

- New heavy quarks with $y_{q'}$, $(\Lambda \sim \text{TeV}) \rightarrow \infty$
 \Rightarrow New physics scale: $\Lambda \sim \mathcal{O}(\text{TeV})$
 - $q' \in$ strong dynamics \Rightarrow composite ϕ_h
 - $m_{q'} \in$ dynamical origin
 - SM fermions + flavor \in “fundamental” ϕ_l (@TeV)
-
- GET: a viable 125 GeV Higgs candidate - mostly fundamental
(protected from large q' loops)
 - consistent with all currently measured 125 Higgs signals
 - Other Higgs states: $m_{H^\pm} \sim m_A < 250 \text{ GeV}$, $m_H \sim 500 \text{ GeV}$