Onset of deconfinement and search for the critical point of strongly interacting matter at CERN SPS energies

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(for the NA49 and NA61/SHINE Collaborations)





NA61/SHINE at the CERN SPS

SHINE - SPS Heavy Ion and Neutrino Experiment

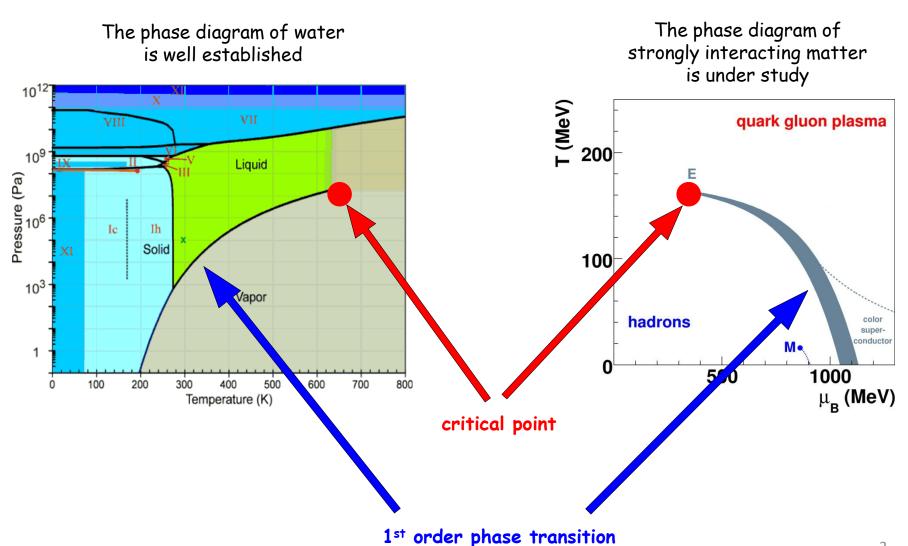


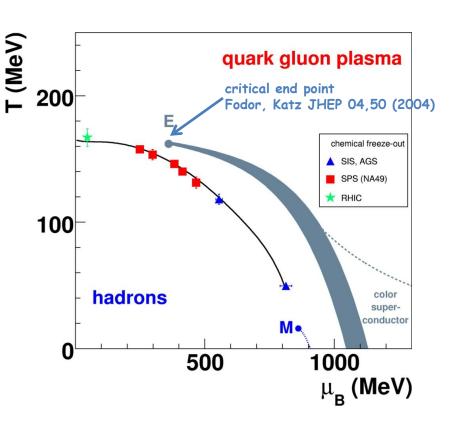
- > Fixed target experiment in the north area of the CERN SPS
- > Succesor of NA49
 - √Pb+Pb at 20A 158A GeV/c
 - √p+p, C+C, Si+Si at 158A GeV/c
- > Started in 2007
- > Beams:
 - √Ions (secondary: Be, primary: Ar and Xe)

at 13A - 158A GeV/c

√Hadrons (secondary): p at 13 - 158 GeV/c,

 π^- at 158 and 350 GeV/c, K- at 158 GeV/c





QCD considerations suggest a 1st order phase boundary ending in a critical point

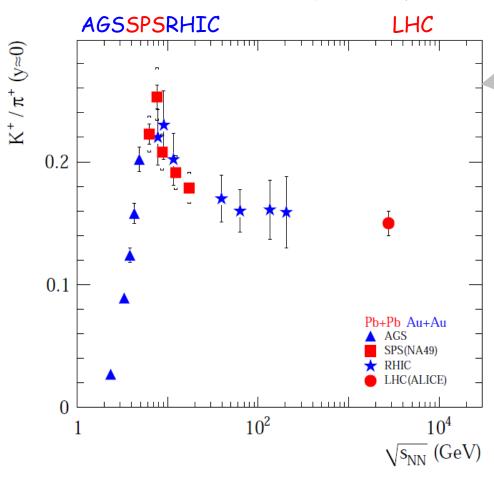
hadro-chemical freeze-out points are obtained from statistical model fits to measured particle yields

T and μ_B approach phase boundary and estimated critical point at SPS

evidence of onset of deconfinement from rapid changes of hadron production properties

search for indications of the critical point as a maximum in fluctuations

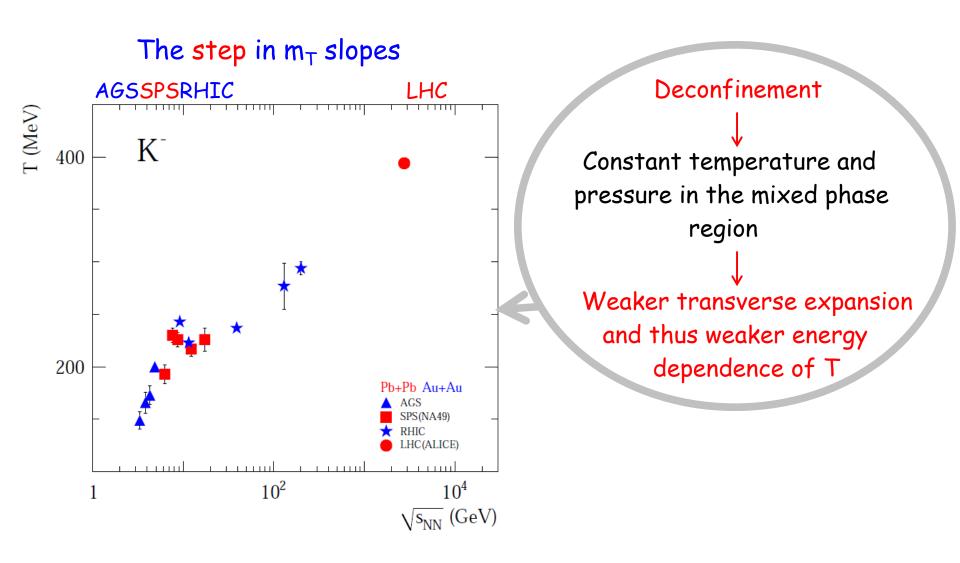




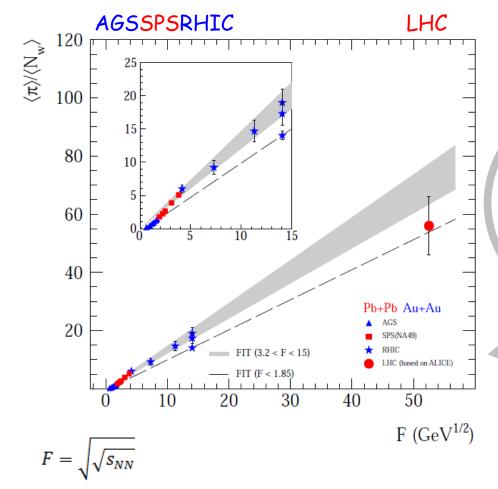
Deconfinement

Decrease of masses of strangeness carriers and the number ratio of strange to non-strange degrees of freedom

A sharp maximum in the strangeness to pion ratio



The kink in pion multiplicity



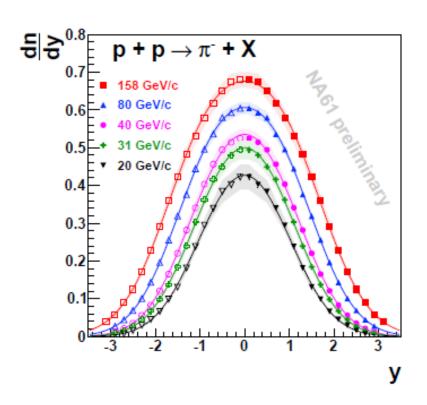
Deconfinement

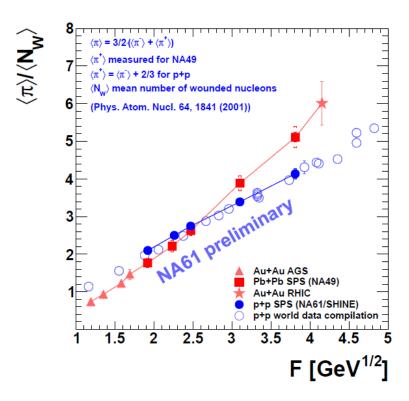
Increased entropy production

Steepening of the increase of pion production

 $\langle \pi \rangle$ - total pion multiplicity

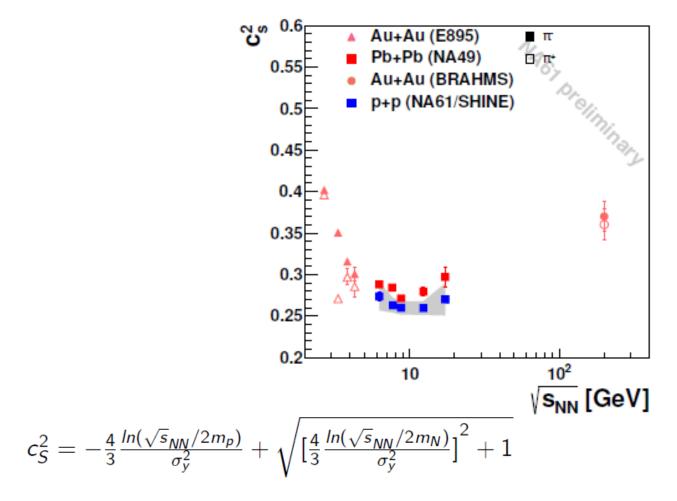
 $\langle N_W \rangle$ - number of interacting nucleons





 $F \approx s_{NN}^{1/4}$

- ✓NA61 results agree with the previous measurements
- $\checkmark\pi$ multiplicity about 200 times higher in central Pb+Pb than in p+p
- $\checkmark \pi$ multiplicity at the SPS energies increases faster in central Pb+Pb than in p+p ("kink")
- √The two dependences cross each other at about 40A GeV/c

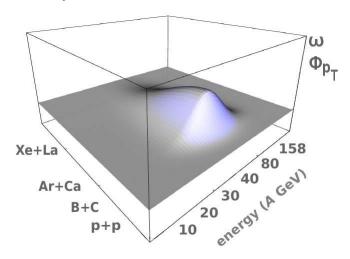


 $\sqrt{c_s^2}$ dependence is similar in p+p and Pb+Pb

Pb+Pb: Phys. Rev. C 66, 054902 (2002), Phys. Rev. C 77, 024903 (2008) c_s²: E. Shuryak, Yad. Fiz. **16**, 707 (1972) [Sov. J. Nucl. Phys. **16**, 395 (1973)], LV. Bravina et al. Phys. Rev. C 60, 024904

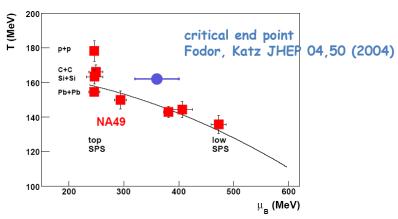
search strategy: 2-dimensional (T, μ_B) scan of phase diagram

expected "hill" of fluctuations



freeze-out points from stat. model

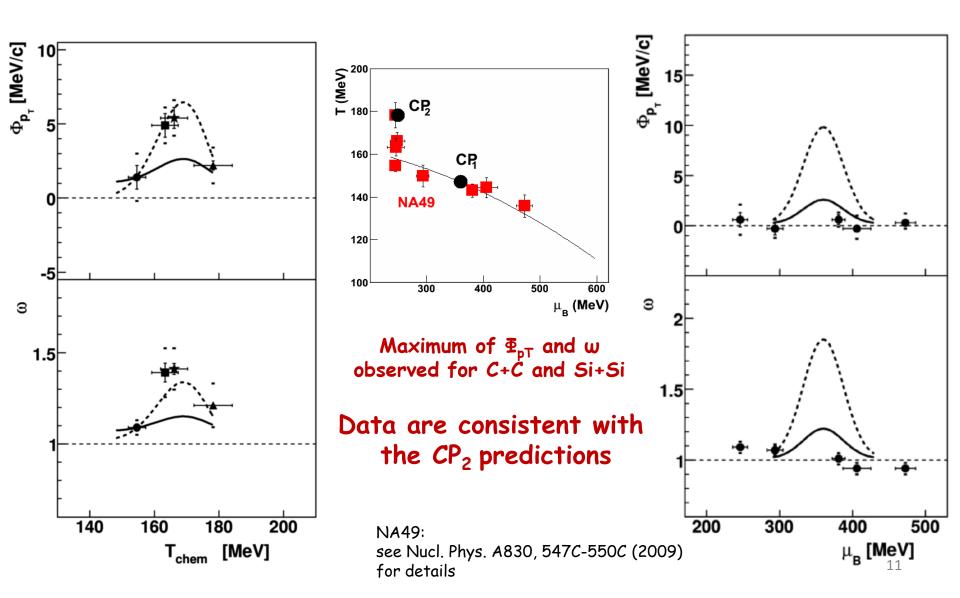
Becattini et al, PRC73, 044905 (2006)

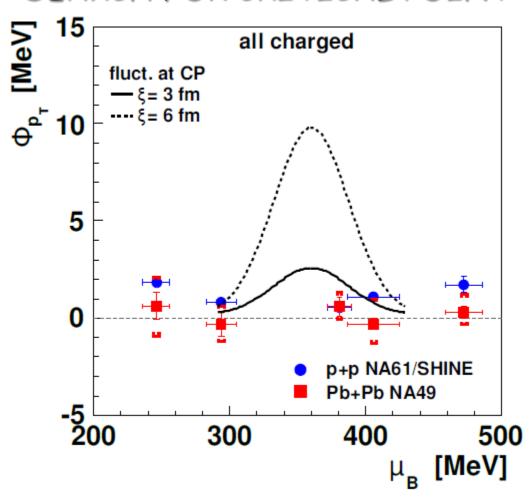


- ✓ Deconfinement necessary for observing CP effect (above 30A GeV)
- \checkmark Expected size of fluctuation signals (\sim ξ²) limited by short lifetime and size of collision system (correlation lengths ξ \sim 3-6 fm for Pb+Pb)

 M.Stephanov, K.Rajagopal, E.Shuryak, PRD60,114028(1999)

Average p_t and multiplicity fluctuations: Dependence on phase diagram coordinates





T. Czopowicz and B. Maksiak,, CPOD 2013

Pion and proton intermittency analysis

Predictions of critical QCD

- 1. Net baryon density at midrapidity is an order parameter for the QCD critical point.
- 2. At the critical point the density-density correlation function in transverse momentum space of net baryons at midrapidity obeys a power-law:

$$\langle n_B(\overrightarrow{p_T})n_B(0)\rangle \sim |\overrightarrow{p_T}|^{2\phi_{2,c}}$$

- 3. For the 3D Ising universality class $\phi_{2,c} = 5/6$
- 4. The critical power-law behaviour of the net baryon density-density correlation is transferred also to the proton density-density correlation.

Methodology

- •Such a power-law distribution can be observed through proton intermittency analysis in transverse momentum space.
- •We have to calculate the second factorial moment of the proton transverse momentum distribution $F_2(M)$ as a function of M (M^2 = number of transverse momentum bins).
- •For protons originating from a critical state (without background) we expect:

$$F_2(M) \sim M^{5/3}$$

In real data background is always present and has to be removed.

Pion and proton intermittency analysis

Intermittency analysis was performed in the following systems:

- C+A with A=C, N (50000 events)
- Si+A with A = AI, Si, P (100000 events)
- · Pb+Pb (1500000 events)

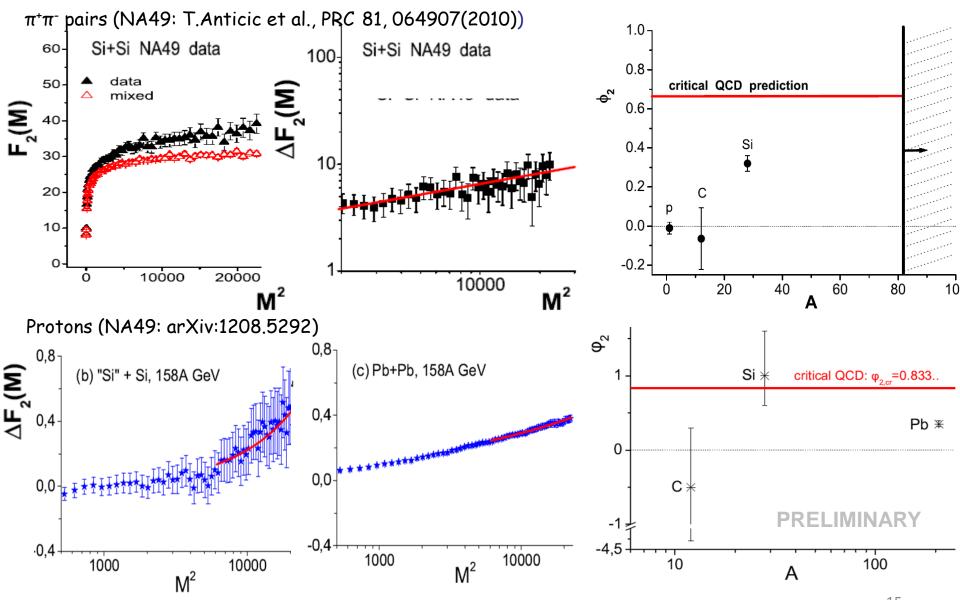
Event and track selection criteriae:

- Events corresponding to central collisions (centrality 0-12%)
- Particles with center of mass rapidity in the interval [-0.75, 0.75]
- Tracks corresponding to identified protons with at least 80% purity

Background is removed by subtracting the moments of constructed mixed events from those of the data:

$$\Delta F_2(M) = F_2^{(data)}(M) - F_2^{(mixed)}(M)$$

We look for a power-law behaviour $\Delta F_2(M) \sim M^{2\phi_2}$ (exactly at the critical point $\phi_2 = \phi_{2,c} = 5/6$)



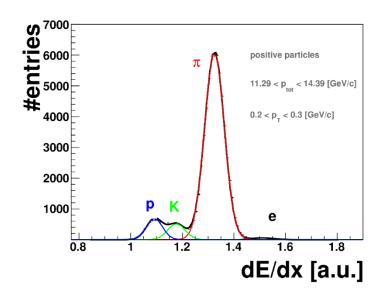
Critical point close to freeze-out points of Si+Si and Pb+Pb system at 158A GeV? 15

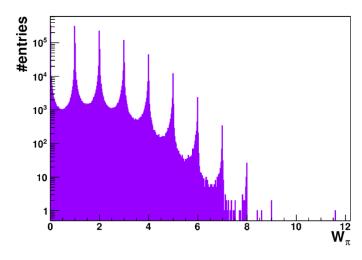
FLUCTUATIONS OF IDENTIFIED PARTICLES

dE/dx information doesn't allow to identify each particle uniquely as the dE/dx distributions overlap

The identity method* allows to obtain second and third moments (pure and mixed) of identified particle multiplicity distributions corrected for misidentification effect

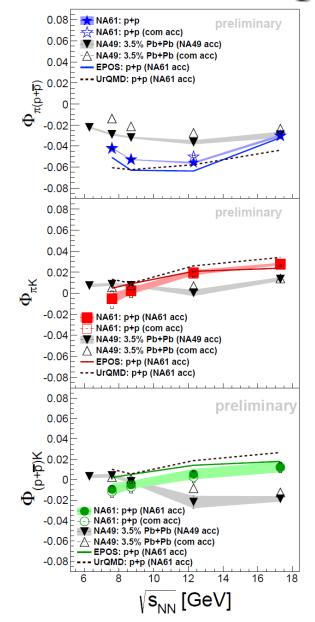
First fluctuation measurements of p+p - unique opportunity to compare Pb+Pb and p+p results at the SPS energies





^{*}Phys. Rev. C83, 054907 (2011), Phys. Rev C84, 024902 (2011) arXiv:nucl-th/1204.6632

♠ MEASURE OF THE FLUCTUATIONS BETWEEN DIFFERENT PARTICLES



$$\begin{split} & \varPhi_{ij} = \frac{\sqrt{\langle N_i \rangle \langle N_j \rangle}}{\langle N_i + N_j \rangle} \cdot \left(\sqrt{\sum_{ij}} - 1 \right) \\ & \text{where } \sum_{ij} = \left[\langle N_i \rangle \omega_j + \langle N_j \rangle \omega_i - 2 \left(\langle N_i N_j \rangle - \langle N_i \rangle \langle N_j \rangle \right) \right] / \langle N_i + N_j \rangle \end{split}$$

 Φ =0 for independent particle production

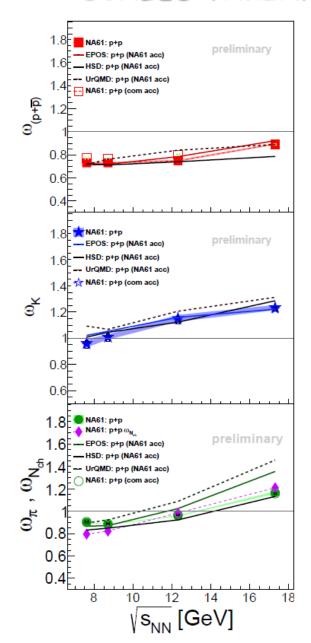
Fluctuations of πp affected by conservation laws and resonance decays (PRC70, 064903).

Small increase of πK fluctuations with increasing energy.

 $\Phi \approx 0$ indicates weak if any correlations in particle production.

UrQMD and EPOS closely reproduce data

SCALED VARIANCE OF MULTIPLICITY DISTRIBUTION



$$\omega = \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle} = \frac{\text{Var}(N)}{\langle N \rangle}$$

w=1 for Poisson distribution

Fluctuations can not be corrected for the limited acceptance (results are presented in NA61 acceptance*)

 ω_{p+p} is below 1 for all SPS energies possibly due to baryon number conservation (B = 2).

 w_K is above 1 for all SPS energies possibly due to strangeness conservation (S = 0).

 ω_π increases with increasing energy reflecting increase of ω_{Nch} measured in full phase-space acceptance.

HSD and EPOS reproduce measured scaled variances.

^{*}https://edms.cern.ch/document/1237791/1

SUMMARY

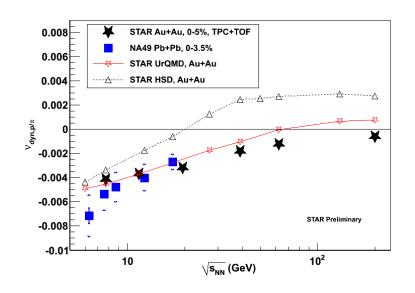
Onset of deconfinement: discovery confirmed, systematic studies in progress:

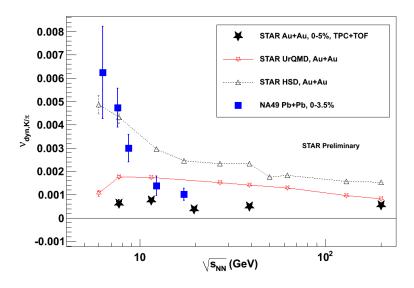
- results from RHIC agree with the relevant NA49 data,
- first LHC data confirm the interpretation,
- NA61/SHINE studies system size dependence of the observed signals.

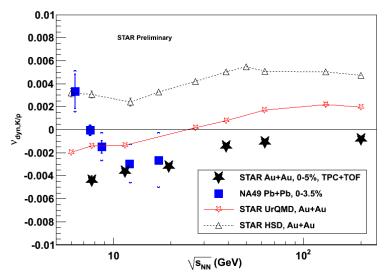
Search for the critical point:

- NA49 hints of a maximum of fluctuations in Si+Si at 158A GeV
- NA61/SHINE performs systematic search via system size energy scan,
- conservation laws seem to play important role in p+p,
- EPOS model reproduces chemical fluctuations in p+p collisions.

ADDITIONAL SLIDES







[K,p] and [K, π] results from NA49 and STAR are significantly different at low energies.

NA49 results use identification via the identity method

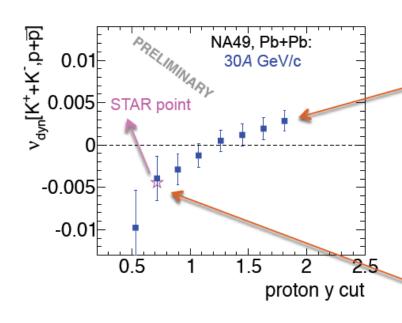
What is the reason for this difference?

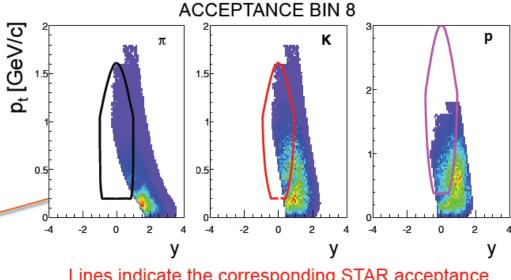
- √bias in the used methods
- ✓acceptance effects

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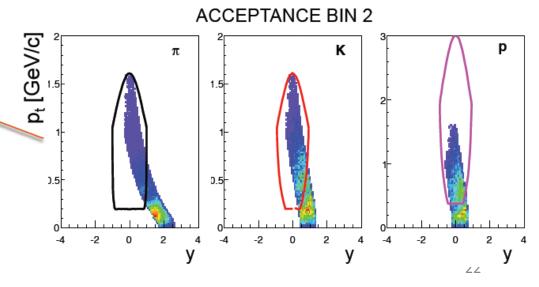
Dependence on acceptance

NA49 central Pb+Pb data at 30A GeV/c

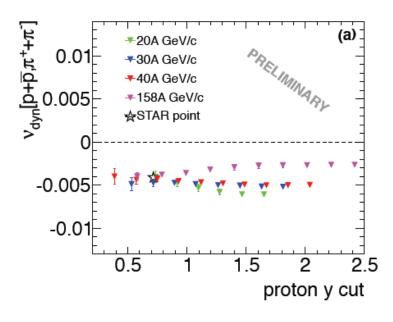




Lines indicate the corresponding STAR acceptance



Dependence on acceptance



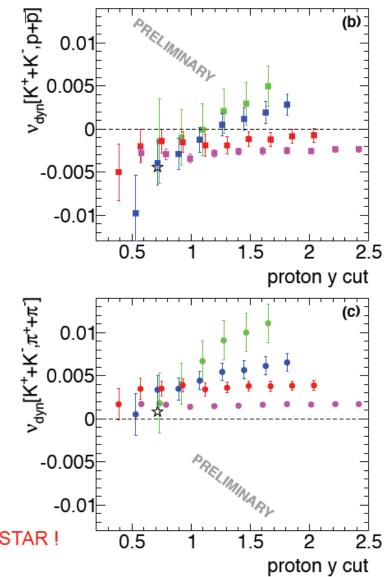
[p,π]: no strong acceptance dependence

[K, π]: at 20A and 30A GeV/c there is a strong

acceptance dependence

[K,p]: at 20A and 30A GeV/c there is a strong

acceptance dependence



acceptance coverage appears to explain the difference with STAR!

 $v_{dyn}[A,B]$ reexamined

$$V_{dyn}[A,B] = \frac{\left\langle A^2 \right\rangle}{\left\langle A \right\rangle^2} + \frac{\left\langle B^2 \right\rangle}{\left\langle B \right\rangle^2} - 2\frac{\left\langle AB \right\rangle}{\left\langle A \right\rangle \left\langle B \right\rangle} - \left(\frac{1}{\left\langle A \right\rangle} + \frac{1}{\left\langle B \right\rangle}\right)$$

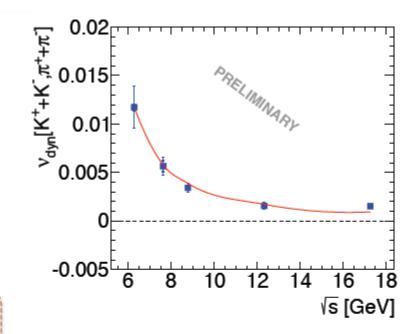
Poisson scaling:

$$v_{dyn}(\sqrt{s}) = v_{dyn}(6.3 \,\text{GeV}) \frac{\left[\frac{1}{\langle \pi \rangle} + \frac{1}{\langle K \rangle}\right]_{\sqrt{s}}}{\left[\frac{1}{\langle \pi \rangle} + \frac{1}{\langle K \rangle}\right]_{6.3 \,\text{GeV}}}$$

After little algebra:

$$\frac{v_{dyn}[A,B]}{\frac{1}{\langle A \rangle} + \frac{1}{\langle B \rangle}} = \Sigma^{AB} - 1$$

$$\Sigma^{AB} = \frac{\langle B \rangle \omega_A + \langle A \rangle \omega_B - 2(\langle AB \rangle - \langle A \rangle \langle B \rangle)}{\langle A + B \rangle}$$



V. Koch, T. Schuster, PRC 81, 034910 (2010) M. I. Gorenstein, M. Gazdzicki et al., PRC 84, 014904 (2011)

Poisson scaled $v_{\mathit{dyn}}[A,B]$ is nothing else but the shifted \sum^{AB}