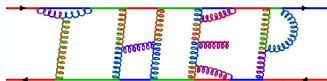


# QCD description of ATLAS jet veto data

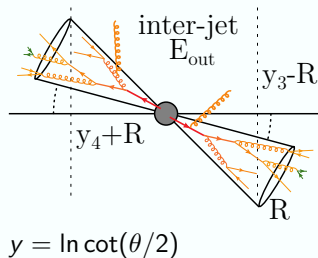
Based on: Y. Hatta, C. Marquet, C. Royon, G. Soyez, T. Ueda, DW, arxiv:1301.1910

EPS-HEP Stockholm, July 19th 2013  
Dominik Werder, Uppsala University

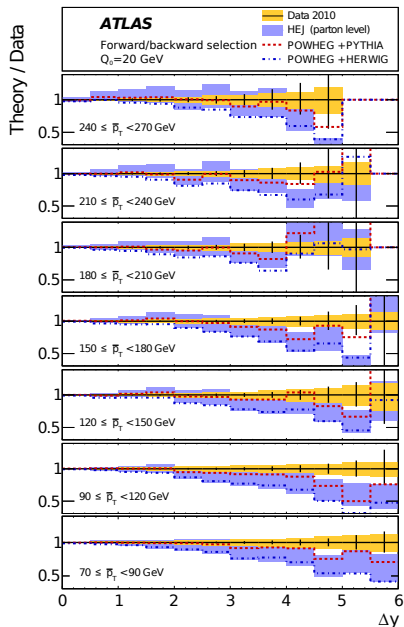


# Jets with veto on inter-jet radiation

- ▶ Di-jet events, high- $p_{\perp}$ ,  $\Delta y = y_3 - y_4$
- ▶ Veto on extra activity between these primary jets:  
Perturbative threshold e.g.  $E_{\text{out}} \simeq 20$  GeV
- ▶  $\mathcal{R}(\Delta y, p_{\perp}) = \frac{d\sigma^{\text{veto}}}{d\Delta y d^2p_{\perp}} / \frac{d\sigma^{\text{incl}}}{d\Delta y d^2p_{\perp}}$
- ▶ Contributing to additional activity between tagged jets:  
Wide angle emissions of gluons from primary, secondary ...



# Jets with veto on inter-jet radiation



Description very good for smaller  $\Delta y$ .

Theory below data for larger  $\Delta y$ .

Pythia with  $p_{\perp}$  ordered showers much closer.

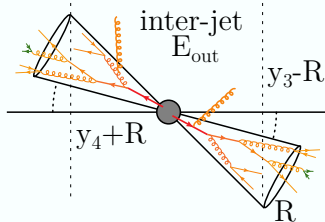
HEJ w/o showers very close as well.

Herwig: angular ordered showers seem to deviate most.

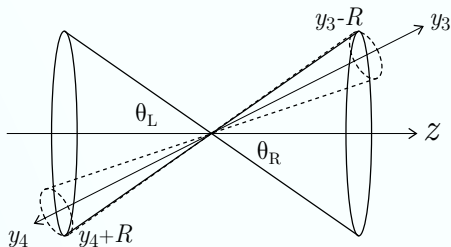
Soft emissions ordered in  $p_{\perp}$ .

# Wide angle emissions

- ▶ Applying veto: Cancellation spoiled  
 $(\alpha_s \Delta y \ln p_\perp / E_{\text{out}})^n$
- ▶ For  $p_\perp \gg E_{\text{out}}$  or large  $\Delta y$  sensitivity to resummation of large-angle gluon emissions
- ▶ Two types of logarithms appear:  
Sudakov (from primary partons)  
Non-global (2nd, 3rd ...)
- ▶ For wide angle gluon emissions from primary parton:  
Global logs  $\rightarrow e^{-c_0 \alpha_s \Delta y \ln(p_\perp / E_{\text{out}})}$
- ▶ Emissions from the secondary (and other) partons:  
More difficult, e.g. [hep-ph/0104277] ...  
We use: Banfi, Marchesini, Smye, 2002  
[hep-ph/0206076]



# BMS approach for non-global resummation



- ▶ LO matrix element
- ▶ Primary jets, given by  $(y_3, \varphi)$ ,  $(y_4, \varphi)$
- ▶ Cones enclose tagged jets: Veto region between the cones
- ▶ Kinematical cuts:  
 $p_{\perp} = (p_{3\perp} + p_{4\perp})/2 \geq 70 \text{ GeV}$   
 $|y| \leq 4.4$   
Jet reconstruction: anti- $k_{\perp}$ ,  $R = 0.6$
- ▶ For  $\sigma^{\text{veto}}$  require  $E_{\perp}^{\text{inter-jet}} \leq E_{\text{out}} \sim 20 \text{ GeV}$

# BMS approach for non-global resummation

- ▶ Probability  $P$  for  $E_{\perp}^{\text{inter-jet}} < E_{\text{out}}$

boost invariance,  $y_R = y_3 - R$ ,  $y_L = y_4 + R$

$$P_{\tau}(y_R, y_L, y_3, y_4) = P_{\tau}\left(\frac{y_R - y_L}{2}, -\frac{y_R - y_L}{2}, y_3 - \frac{y_R + y_L}{2}, y_4 - \frac{y_R + y_L}{2}\right)$$

- ▶ Evolution variable, for running  $\alpha_s$

$$\tau = \int_{E_{\text{out}}}^{p_{\perp}} \frac{dk_{\perp}}{k_{\perp}} \frac{\alpha_s(k_{\perp}) N_c}{\pi} = \frac{1}{2b} \ln(\alpha_s(E_{\text{out}})/\alpha_s(p_{\perp})) = \frac{1}{2b} \ln \frac{\ln(p_{\perp}/\Lambda_{\text{QCD}})}{\ln(E_{\text{out}}/\Lambda_{\text{QCD}})}$$

- ▶ Evolution equation, large- $N_c$

$$\begin{aligned} \partial_{\tau} P_{\tau}(\Omega_{\alpha}, \Omega_{\beta}) &= - \int_{C_{\text{out}}} \frac{d^2 \Omega_{\gamma}}{4\pi} \frac{1 - \cos \theta_{\alpha\beta}}{(1 - \cos \theta_{\alpha\gamma})(1 - \cos \theta_{\gamma\beta})} P_{\tau}(\Omega_{\alpha}, \Omega_{\beta}) \\ &+ \int_{C_{\text{in}}} \frac{d^2 \Omega_{\gamma}}{4\pi} \frac{1 - \cos \theta_{\alpha\beta}}{(1 - \cos \theta_{\alpha\gamma})(1 - \cos \theta_{\gamma\beta})} (P_{\tau}(\Omega_{\alpha}, \Omega_{\gamma}) P_{\tau}(\Omega_{\gamma}, \Omega_{\beta}) - P_{\tau}(\Omega_{\alpha}, \Omega_{\beta})) \end{aligned}$$

- ▶ Numerically solved (Y. Hatta, T. Ueda, 2009)

# LO ME + single leading log resummed

- ▶ Contributing processes: Leading order QCD  $2 \rightarrow 2$

$qq' \rightarrow qq', qq \rightarrow qq, q\bar{q} \rightarrow q\bar{q}, q\bar{q} \rightarrow q'\bar{q}', q\bar{q} \rightarrow gg \dots$

$$\frac{d\sigma}{dp_{\perp} d\Delta Y} = \sum_{ij} \int_Y dY x_1 f_i(x_1) x_2 f_j(x_2) \frac{1}{\pi} \frac{d\sigma_{ij}}{d\hat{t}}$$

- ▶ Weighted QCD cross section with BMS veto factor
- ▶ Standard proton PDF
- ▶ Running  $\alpha_s$  with  $\alpha_s(m_Z = 91.2 \text{ GeV}) = 0.12$
- ▶ Take ratio with inclusive cross section ( $P = 1$ ):

$$\frac{\text{Diagram 1}}{\text{Diagram 2} + \text{Diagram 3}}$$

The diagram shows a ratio of two terms. The numerator is a single diagram consisting of two vertices connected by two lines, forming a diamond shape. The denominator is the sum of two diagrams: the first is the same diamond shape as the numerator, and the second is a diagram where the two vertices are connected by two lines, but the lines are crossed in the middle, forming a figure-eight shape.

# Estimation of uncertainties

- ▶ Uncertainties: Variation of  $\alpha_s$ -scale ( $\times 2$ )

- ▶ in ME
- ▶ in  $\tau$

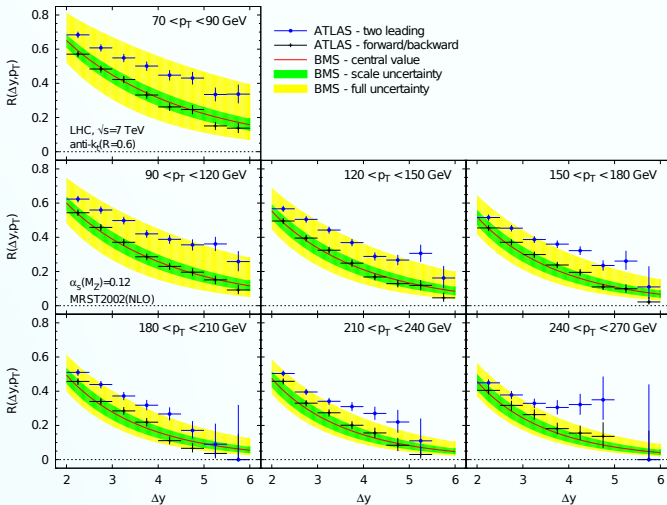
Factor  $\alpha_s$  in ME cancels in ratio

- ▶ Variation of scale in PDF ( $\times 2$ )
- ▶ Cross-check PDFs (CT10, MRST2002, ...)
- ▶ Subleading logs not accounted for in  $(\alpha_s \ln(p_\perp/E_{\text{out}}))^n$   
Vary boundary in  $\tau$  integral

$$\begin{aligned}\tau &= \int_{E_{\text{out}}}^{p_\perp} \frac{dk_\perp}{k_\perp} \frac{\alpha_s(k_\perp) N_c}{\pi} \\ &= \frac{1}{2b} \ln(\alpha_s(E_{\text{out}})/\alpha_s(p_\perp))\end{aligned}$$



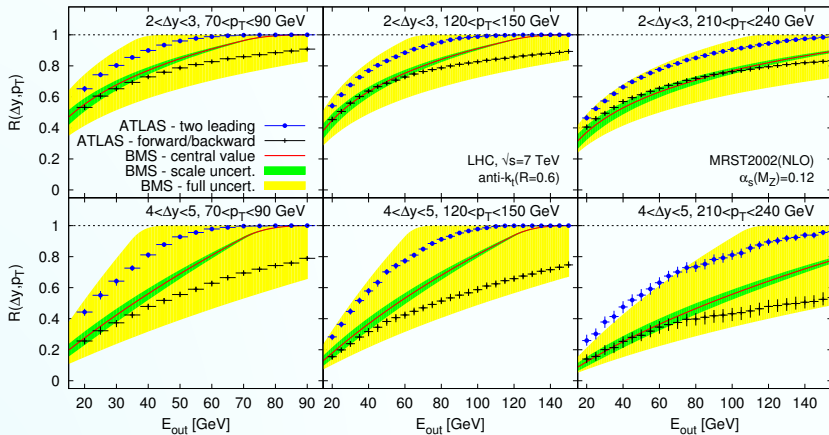
# Results



Y. Hatta, C. Marquet, C. Royon, G. Soyez, T. Ueda, DW, arxiv:1301.1910

- ▶ Very nice agreement with forward/backward jet selection
- ▶ Larger uncertainties from estimation of subleading logs

# Results



- ▶ Best description for large scale difference  
Large scale needed for formalism

# Conclusions

- ▶ Jets with veto on inter-jet region
- ▶ Emissions from primaries (Sudakov) vs. secondaries
- ▶ Resummation possible for first case, numerical solution via BMS for second
- ▶ Compared to ATLAS jet veto data:  
Very good agreement with the forward/backward dijet selection

## Comment on BFKL

- ▶ Choice of veto  $E_{\text{out}}$  too large for sensitivity to BFKL
- ▶ Would increase  $P$   
Scattered quarks form dipoles most dominantly with their remnants  
⇒ Less radiative activity between the boundary jets
- ▶ Need for BFKL not evident for the choice of  $E_{\text{out}}$



Gap fraction

