

Jet quenching in a strongly interacting plasma

A lattice approach

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- 1 Motivation
- 2 Theoretical approach
- 3 Soft physics contribution from a Euclidean setup
- 4 Lattice implementation
- 5 Results
- 6 Discussion and conclusions

Based on:

- ★ M. P., K. Rummukainen and A. Schäfer, *A lattice study of the jet quenching parameter*, [arXiv:1307.xxxx](https://arxiv.org/abs/1307.xxxx)

Related works:

- S. Caron-Huot, *$O(g)$ plasma effects in jet quenching*, [arXiv:0811.1603](https://arxiv.org/abs/0811.1603)
- M. Benzke, N. Brambilla, M. A. Escobedo and A. Vairo *Gauge invariant definition of the jet quenching parameter*, [arXiv:1208.4253](https://arxiv.org/abs/1208.4253)
- M. Laine, *A non-perturbative contribution to jet quenching*, [arXiv:1208.5707](https://arxiv.org/abs/1208.5707)
- J. Ghiglieri, J. Hong, A. Kurkela, E. Lu, G. D. Moore and D. Teaney, *Next-to-leading order thermal photon production in a weakly coupled quark-gluon plasma*, [arXiv:1302.5970](https://arxiv.org/abs/1302.5970)
- M. Laine and A. Rothkopf, *Light-cone Wilson loop in classical lattice gauge theory*, [arXiv:1304.4443](https://arxiv.org/abs/1304.4443)

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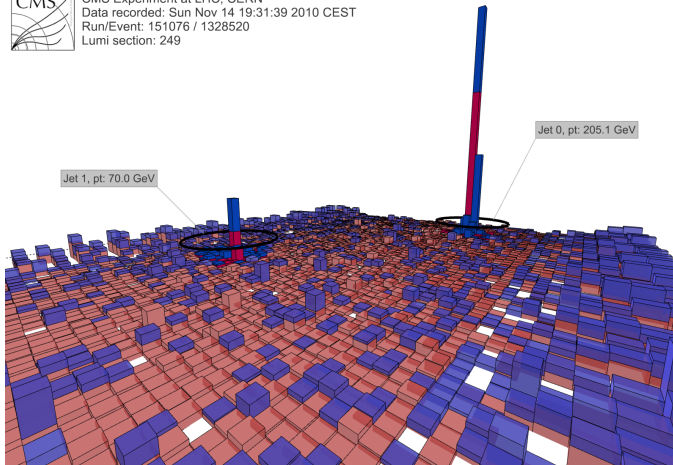
5 *Results*

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Jet quenching is the suppression of high- p_T particles and back-to-back correlations in nuclear collisions

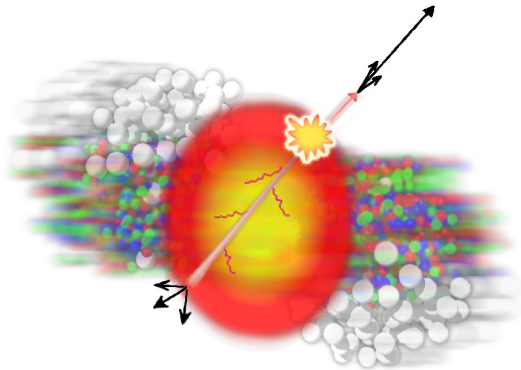


CMS Experiment at LHC, CERN
Data recorded: Sun Nov 14 19:31:39 2010 CEST
Run/Event: 151076 / 1328520
Lumi section: 249



Jet quenching is the suppression of high- p_T particles and back-to-back correlations in nuclear collisions

Provides important experimental evidence for the quark-gluon plasma (QGP) existence [Bjorken, 1982](#)



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Jet quenching belongs to the class of *hard probes* to heavy-ion collisions, involving a large energy scale Q (see [Casalderrey-Solana and Salgado, 2007](#) and [Apolinário's talk at EPS HEP 2013](#))

QCD factorization theorems:

$$\sigma_{(M+N \rightarrow \text{hadron})} = f_M(x_1, Q^2) \otimes f_N(x_2, Q^2) \otimes \sigma(x_1, x_2, Q^2) \otimes D_{\text{parton} \rightarrow \text{hadron}}(z, Q^2)$$

$f_A(x, Q^2)$: parton distribution functions

$\sigma(x, y, Q^2)$: short-distance cross-section

$D_{\text{parton} \rightarrow \text{hadron}}(z, Q^2)$: fragmentation function

Here: Focus on propagation of a high-energy parton in QGP medium

Multiple soft-scattering description, in the *eikonal approximation*

Leading effect: *transverse momentum broadening*, described by the jet quenching parameter:

$$\hat{q} = \frac{\langle p_{\perp}^2 \rangle}{L}$$

Can be evaluated in terms of a *collision kernel* $C(p_{\perp})$ (differential parton-plasma constituents collision rate)

$$\hat{q} = \int \frac{d^2 p_{\perp}}{(2\pi)^2} p_{\perp}^2 C(p_{\perp})$$

$C(p_{\perp})$ can be related to a two-point correlator of *light-cone Wilson lines*

What tools are available?

- Perturbation theory (PT) expansions
 - ✓ Based on first principles
 - ✓ Well established technology
 - ✓ Problems with infrared divergences are well understood
 - ✗ May not be reliable at RHIC or LHC temperatures
- Holographic computations
 - ✓ Mathematically *beautiful*
 - ✓ Ideally suited for strong coupling
 - ✗ Not directly applicable to real-world QCD
- Lattice simulations
 - ✓ Based on first principles
 - ✓ Well established (computer) technology
 - ✓ Do not rely on weak- or strong-coupling assumptions
 - ✗ Euclidean setup, so *generally* unsuitable for real-time phenomena

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Energy scale hierarchy in high-temperature, perturbative QCD:

$$g^2 T / \pi \text{ (ultrasoft)} \ll gT \text{ (soft)} \ll \pi T \text{ (hard)}$$

IR divergences accounted for by 3D effective theories [Braaten and Nieto, 1995](#):

- electrostatic QCD (3D Yang-Mills + adjoint scalar field) for soft scale
- magnetostatic QCD (3D pure Yang-Mills) for ultrasoft scale

Large NLO corrections hindering PT due to *soft*, essentially *classical* fields

Observation: Soft contributions to physics of light-cone partons *insensitive* to parton velocity → Can turn the problem Euclidean!

Spatially separated ($|t| < |z|$) light-like Wilson lines:

$$\begin{aligned}
 G^<(t, x_{\perp}, z) &= \int d\omega d^2 p_{\perp} dp^z \tilde{G}^<(\omega, p_{\perp}, p^z) e^{-i(\omega t - x_{\perp} \cdot p_{\perp} - z p^z)} \\
 &= \int d\omega d^2 p_{\perp} dp^z \left[\frac{1}{2} + n_B(\omega) \right] \left[\tilde{G}_R(\omega, p_{\perp}, p^z) - \tilde{G}_A(\omega, p_{\perp}, p^z) \right] e^{-i(\omega t - x_{\perp} \cdot p_{\perp} - z p^z)}
 \end{aligned}$$

Shift $p'^z = p^z - \omega t/z$, integrate over frequencies by analytical continuation into upper (lower) half-plane for retarded (advanced) contribution \rightarrow sum over Matsubara frequencies

$$G^<(t, x_{\perp}, z) = T \sum_{n \in \mathbb{Z}} \int d^2 p_{\perp} dp^z \tilde{G}_E(2\pi n T, p_{\perp}, p^z) e^{i(x_{\perp} \cdot p_{\perp} + z p^z)}$$

- $n \neq 0$ contributions: exponentially suppressed at large separations
- Soft contribution: from $n = 0$ mode. Time-independent: evaluate in EQCD

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Super-renormalizable EQCD Lagrangian

$$\mathcal{L} = \frac{1}{4} F_{ij}^a F_{ij}^a + \text{Tr}((D_i A_0)^2) + m_E^2 \text{Tr}(A_0^2) + \lambda_3 (\text{Tr}(A_0^2))^2$$

Parameters chosen (by matching) to reproduce soft physics of high- T QCD

- 3D gauge coupling: $g_E^2 = g^2 T$
- Debye mass parameter: $m_E^2 = (1 + \frac{n_f}{6}) g^2 T$
- 3D quartic coupling: $\lambda_3 = \frac{9-n_f}{24\pi^2} g^4 T$

Standard Wilson lattice regularization

Our setup: QCD with $n_f = 2$ light flavors, two temperature ensembles:

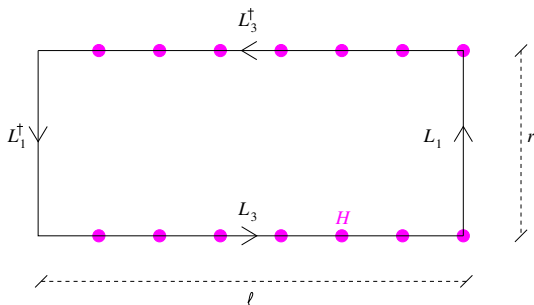
- $T \simeq 398$ MeV
- $T \simeq 2$ GeV

Light-cone Wilson line correlator

$$\langle W(\ell, r) \rangle = \left\langle \text{Tr} \left(L_3 L_1 L_3^\dagger L_1^\dagger \right) \right\rangle \sim \exp[-\ell V(r)]$$

with

$$L_3 = \prod U_3 H \quad L_1 = \prod U_1 \quad H = \exp(-ag_E^2 A_0)$$



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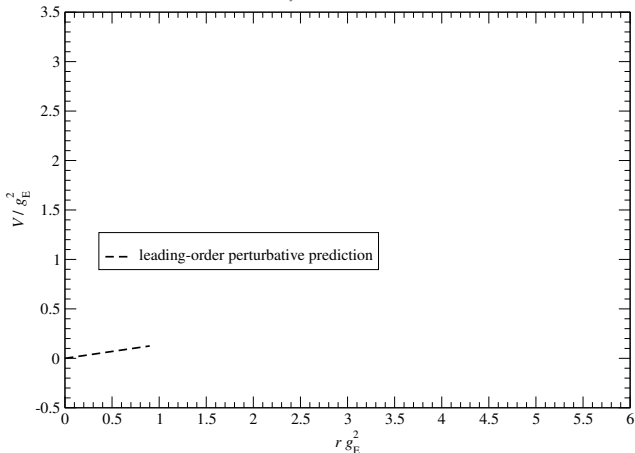
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Non-perturbative soft contribution to \hat{q} from short-distance behavior of $V(r)$

Potential from the decorated loop operator in EQCD

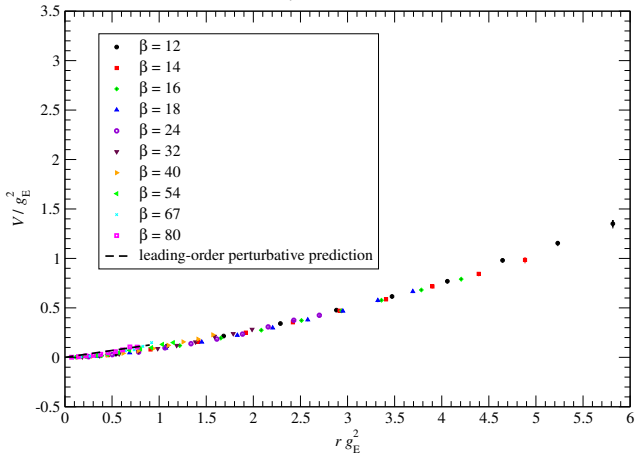
($n_f = 2, T \simeq 398$ MeV)



Non-perturbative soft contribution to \hat{q} from short-distance behavior of $V(r)$

Potential from the decorated loop operator in EQCD

($n_f = 2, T \simeq 398 \text{ MeV}$)



Soft NLO contribution to \hat{q} is quite *large*:

$$\hat{q}_{\text{EQCD}} \simeq \begin{cases} 0.55(5)g_E^6 & \text{for } T \simeq 398 \text{ MeV} \\ 0.45(5)g_E^6 & \text{for } T \simeq 2 \text{ GeV} \end{cases}$$

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These results lead to an *approximate* estimate $\hat{q} \sim 6 \text{ GeV}^2/\text{fm}$ at RHIC temperatures

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- Lattice approach *possible* for certain real-time problems
- Here: focus on soft physics in thermal QCD
- Outlined approach is *systematic*
- Tentative estimate of jet quenching parameter
- Results in ballpark of
 - holographic computations [Liu, Rajagopal and Wiedemann, 2006](#) ✓
 - estimates from experiments [Eskola et al., 2004](#) ✓
- Further work in progress