



HEP 2013
Stockholm
18-24 July 2013
(info@eps-hep2013.eu)



Experimentally Viable Mass of the Fermionic Dark Matter

Based on Phys.Rev. D87 (2013) 115016

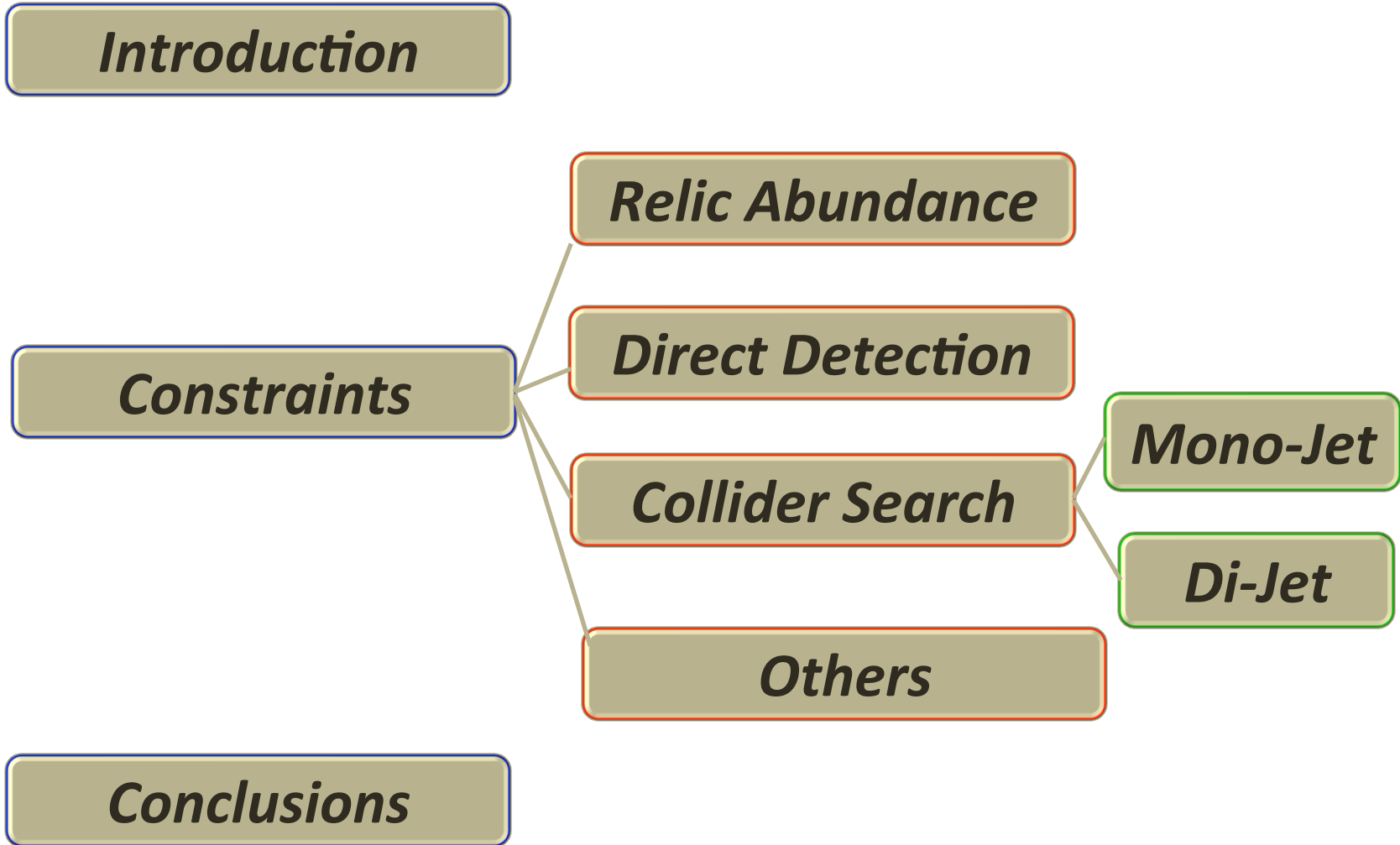
In coll. with Kwei-Chou Yang

Ho-Chin Tsai

CYCU Taiwan

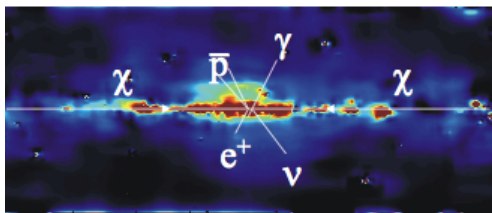
2013.07.18 (Thu) 18:15-18:30pm

Outline

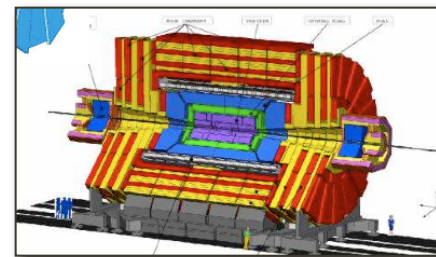
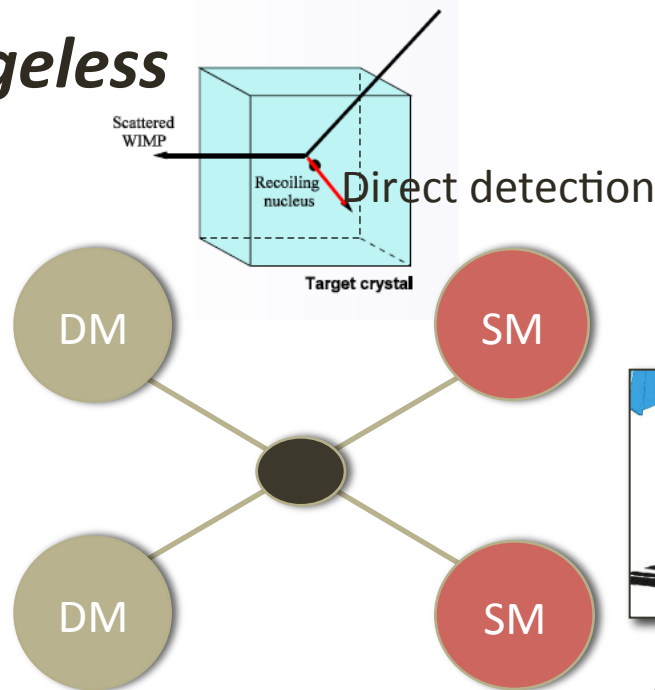


Dark Matter Introduction

- **CMB, Bullet Clusters, Galactic Rotational Curve,...**
evidence in various scales.
- **Interaction weak to escape experimental searches**
but not too weaker to over produce DM relic
- **Invisible, Chargeless**
- **BSM fermion, scalar...**



Indirect detection
Relic \rightarrow



\leftarrow Collider



Dark Matter Searches

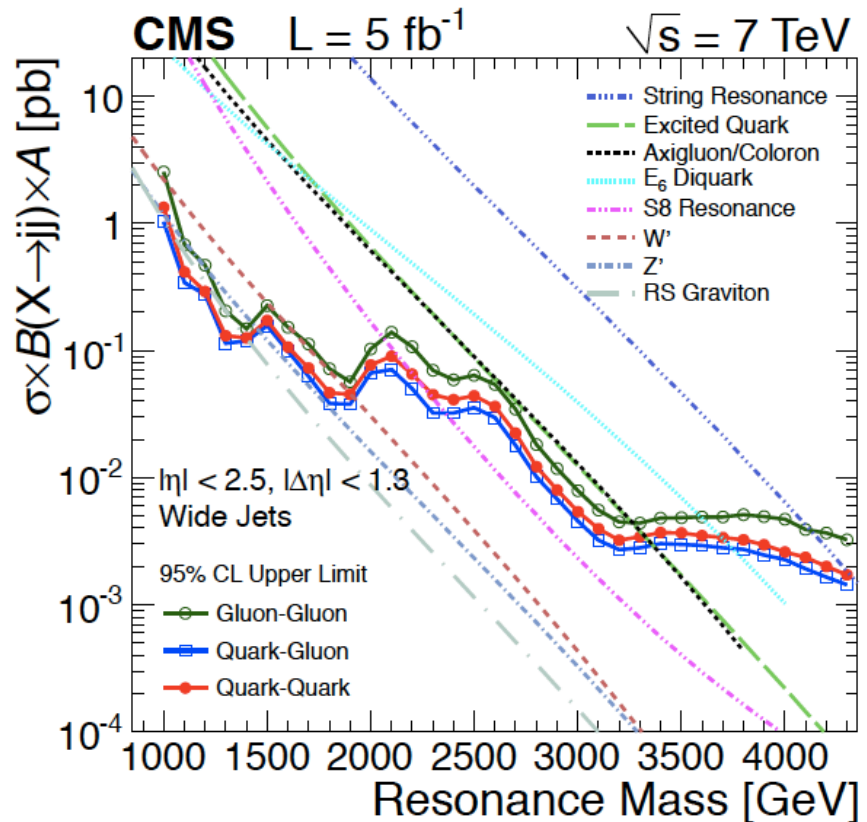
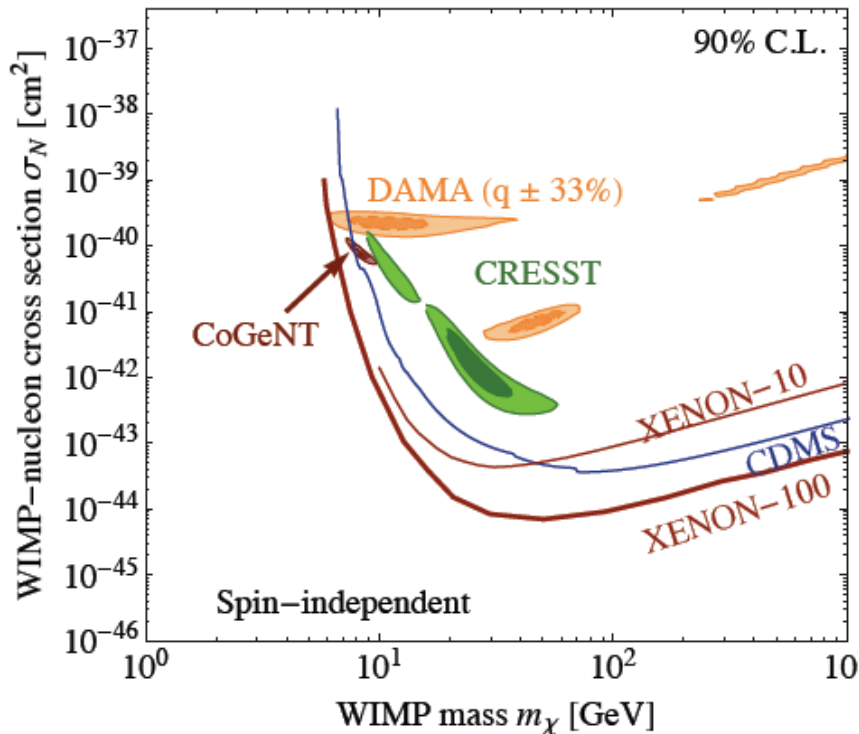
$$\text{Relic } \Omega_\chi h^2 \simeq \frac{0.1 \text{ pb} \cdot c}{\langle \sigma_{\text{ann}} v_{\text{rel}} \rangle} \sim 0.1$$

Monojet + missing transverse energy

DM Direct Detection: $10^{-9} \sim 10^{-10}$ pb

Di-Jet Resonance Search

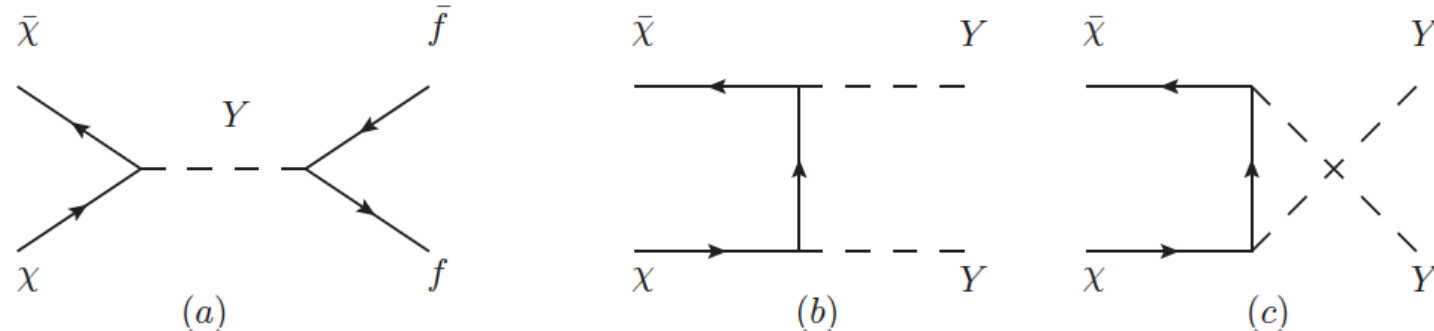
DM Indirect Detection



SM + Fermionic DM + Scalar Y

- Effective $Y \bar{\chi} (\lambda_s^{\chi} + i\lambda_p^{\chi} \gamma^5) \chi + Y \bar{f} (\lambda_s^f + i\lambda_p^f \gamma^5) f$

- DM annihilation



- Similar formula for SS, SP, PS, PP

$$\langle \sigma v_{\text{rel}} \rangle^{\underline{SS}} = N_c \frac{\lambda_s^{\chi^2} \lambda_s^{f^2}}{2\pi} \frac{p_{\chi}^2 p_f^3}{E^3 [(4E^2 - m_Y^2)^2 + m_Y^2 \Gamma_{SS,Y}^2]} + \frac{\lambda_s^{\chi^4}}{2\pi} \frac{E m_{\chi}^2 p_{\chi}^2 p_Y \theta(E - m_Y)}{[4(E^2 - m_Y^2) m_{\chi}^2 + m_Y^4]^2}$$

- DM relic $\Omega_{\chi} h^2 \simeq \frac{0.1 \text{ pb} \cdot c}{\langle \sigma_{\text{ann}} v_{\text{rel}} \rangle} \sim 0.1$ fix coupling to $\lambda_{\chi}^4, \lambda_f^4(m_{\chi}, m_Y)$

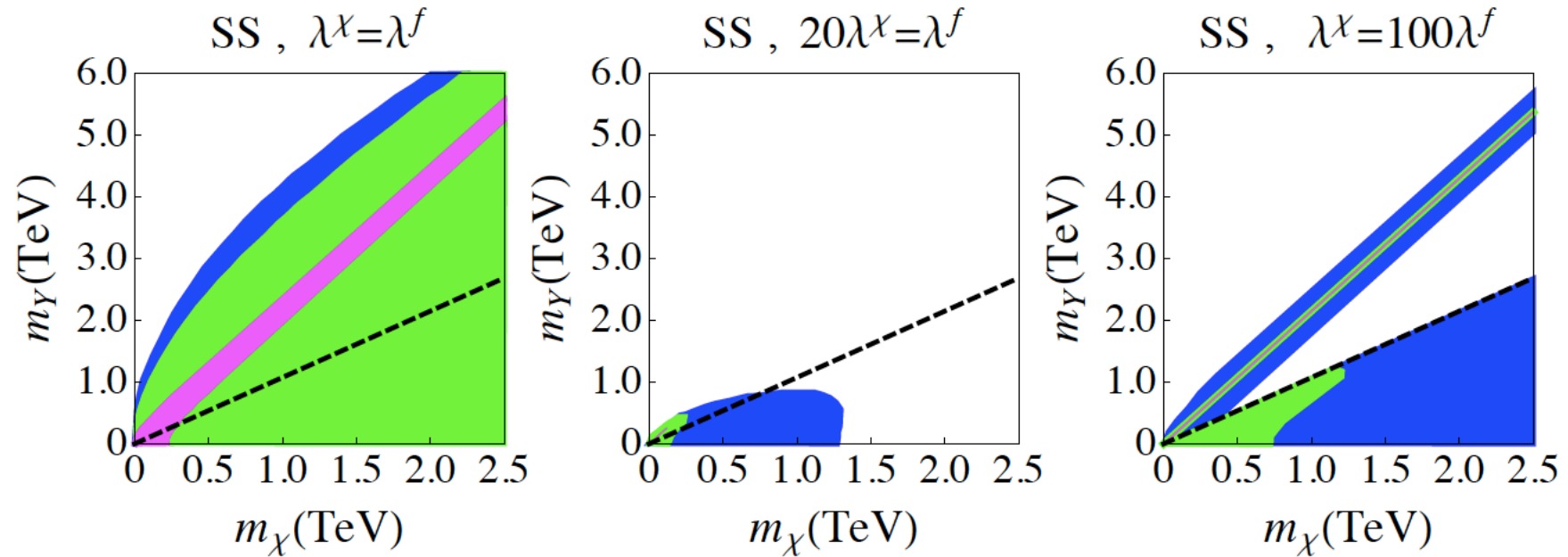


Plot m_χ v.s. m_Y with Right Relic

$$\text{Max}\{\lambda_\chi, \lambda_f\}$$

$$m_Y \sim E \sim m_\chi + \frac{3}{40} m_\chi$$

Color code: $0 < \text{Magenta} < 0.3 < \text{Green} < 1 < \text{Blue} < 3 \sim \text{Perturbativity} < \sqrt{4\pi}$



Similar plots for SP, PS and PP cases



χ -N Scattering Direct Detection

$$\sigma_{el}^{SS}(\chi N \rightarrow \chi N) = \frac{\lambda_s^{\chi^2} \lambda_s^{f^2}}{\pi} \frac{m_\chi^2 m_N^2}{(m_\chi + m_N)^2 m_Y^4} f_N^2$$

SP

$$p^2 / 2m_N^2$$

PS

$$p^2 / 2m_\chi^2$$

PP

$$p^4 / 3m_N^2 m_\chi^2$$

Pseudo-scalar velocity

Suppression

1

$$p = \mu v$$

10^{-6}

$$v \sim 10^{-3} c$$

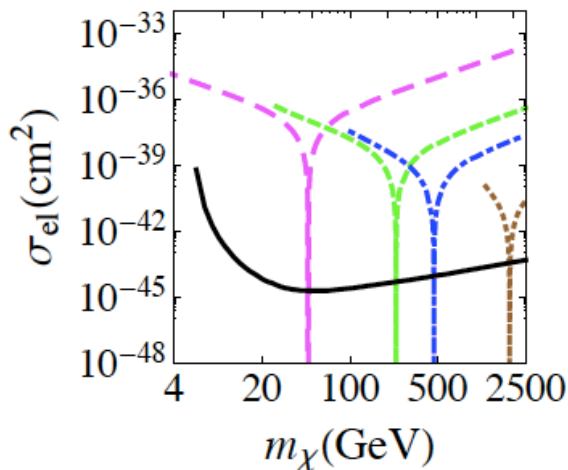
10^{-10}

10^{-12}

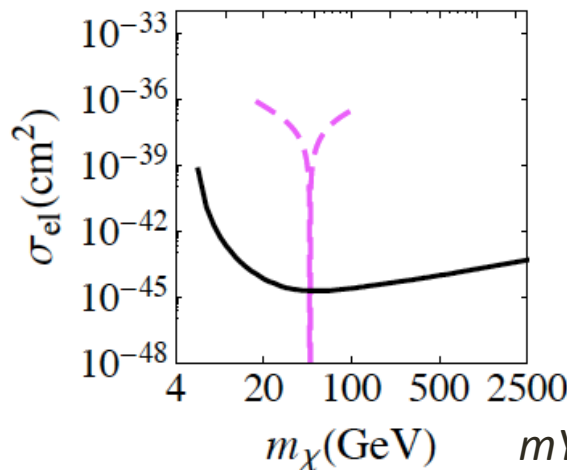
PS, PP are very rarely constrained by **Xenon100**

$P=10^{-3}$ for $m_\chi = 100 m_N$

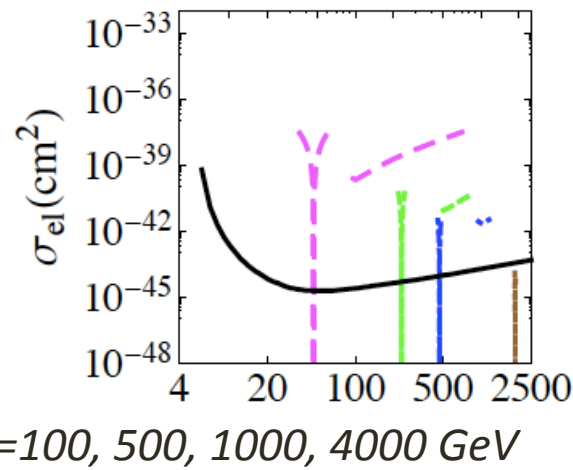
SS, $\lambda^\chi = \lambda^f$



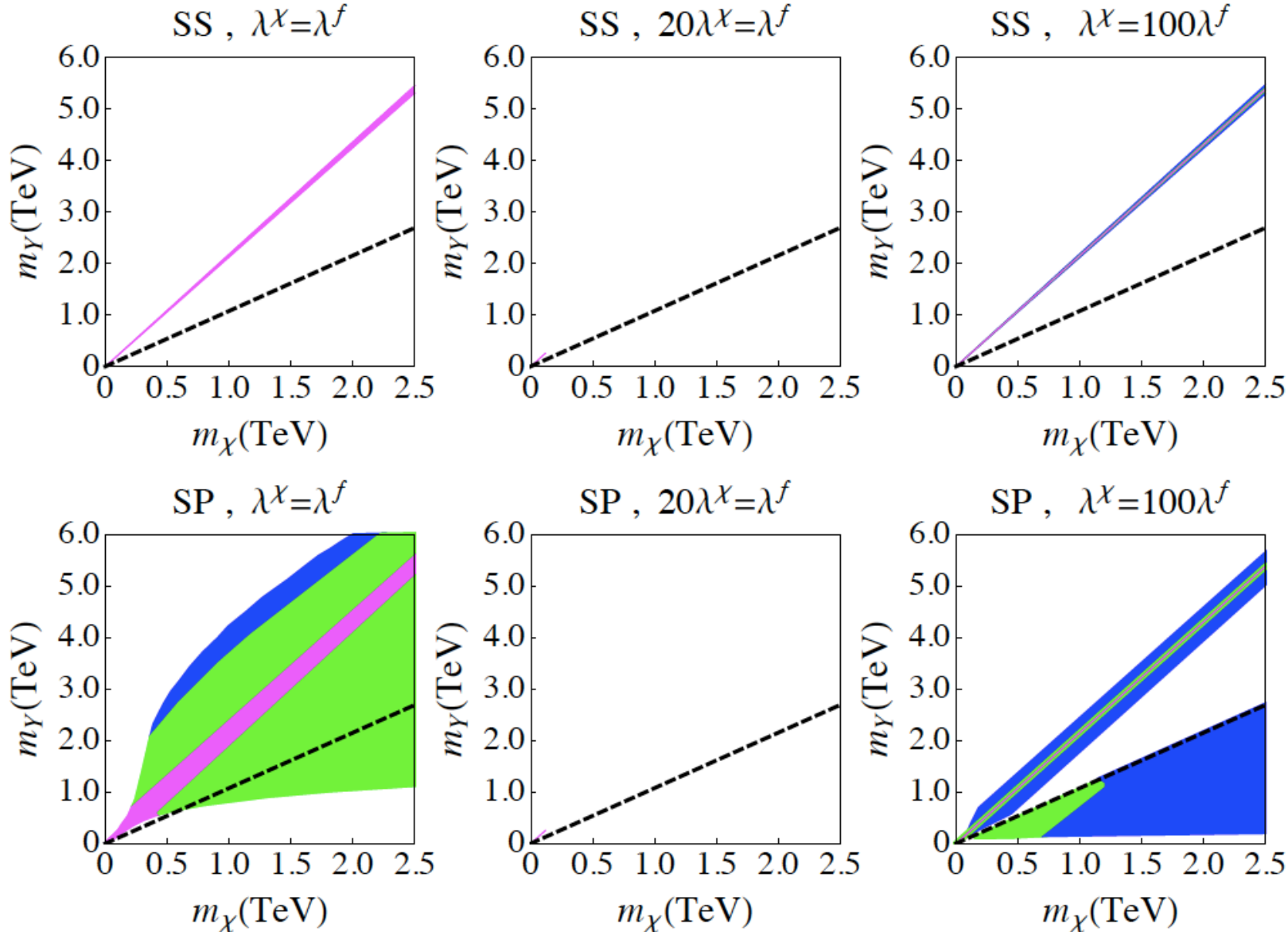
SS, $20\lambda^\chi = \lambda^f$



SS, $\lambda^\chi = 100\lambda^f$

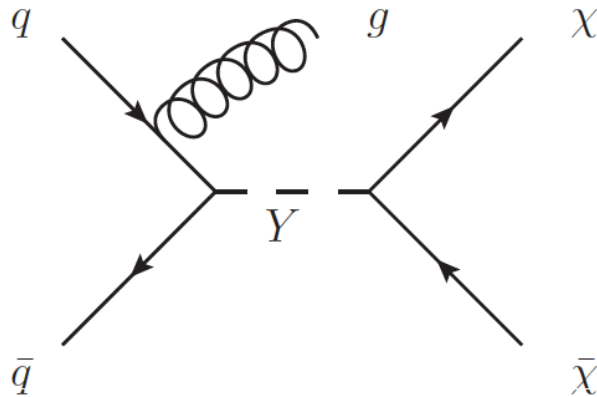


Relic+Direct Detections Results



DM in Mono-jet + Missing E_T

Invisible Y



For light DMs

SM BG $p p \rightarrow j Z \rightarrow j \nu \nu$

ALTAS 1202.0158 MadGraph5

VeryHighPT: 7 TeV, 1fb-1 missing $E_T > 300\text{GeV}$

$\sigma_{1j} < 0.045\text{pb}$

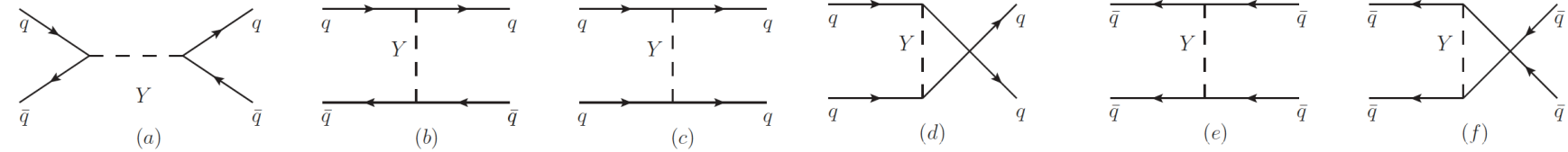


Dijet Resonances Search for Y

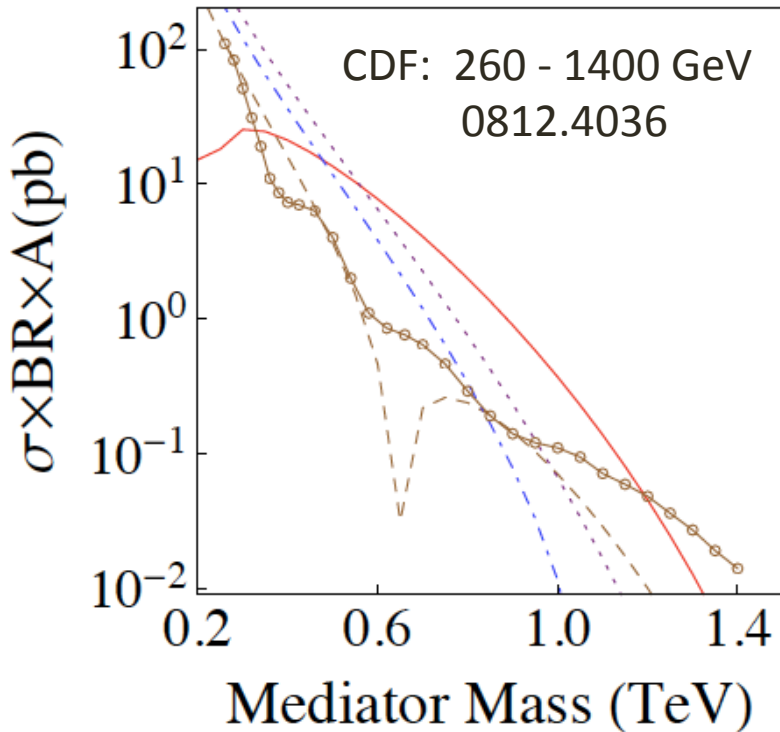


visible Y

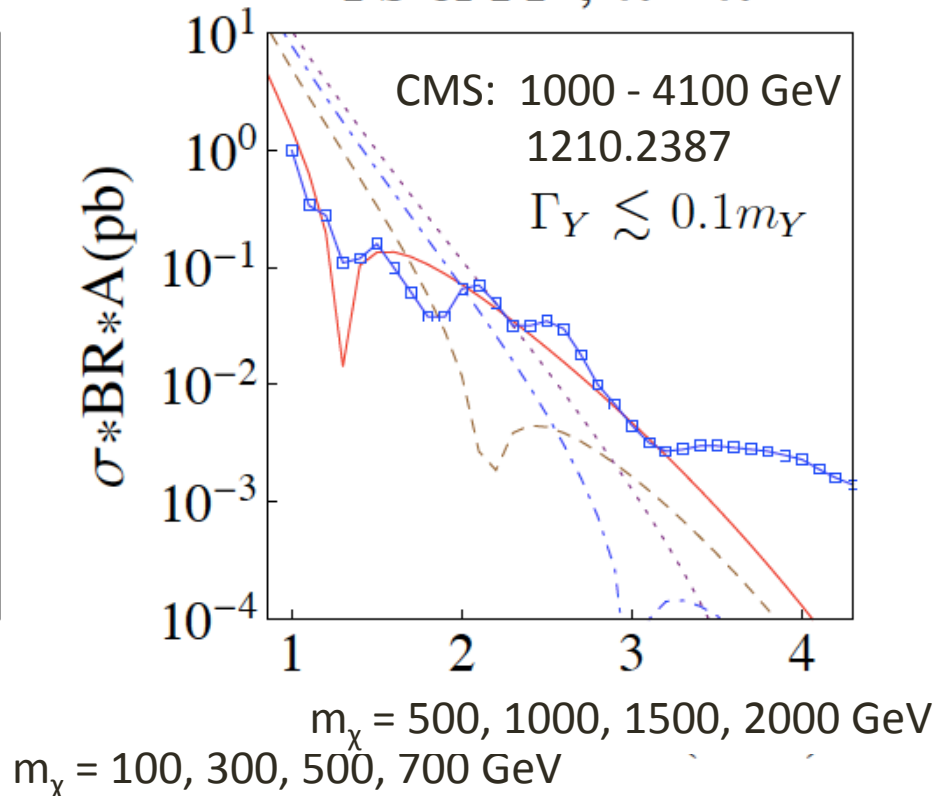
Review 1110.5302 by R.M. Harris & K. Kousouris



PS or PP , $\lambda^X = \lambda^f$



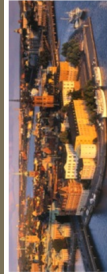
PS or PP , $\lambda^X = \lambda^f$



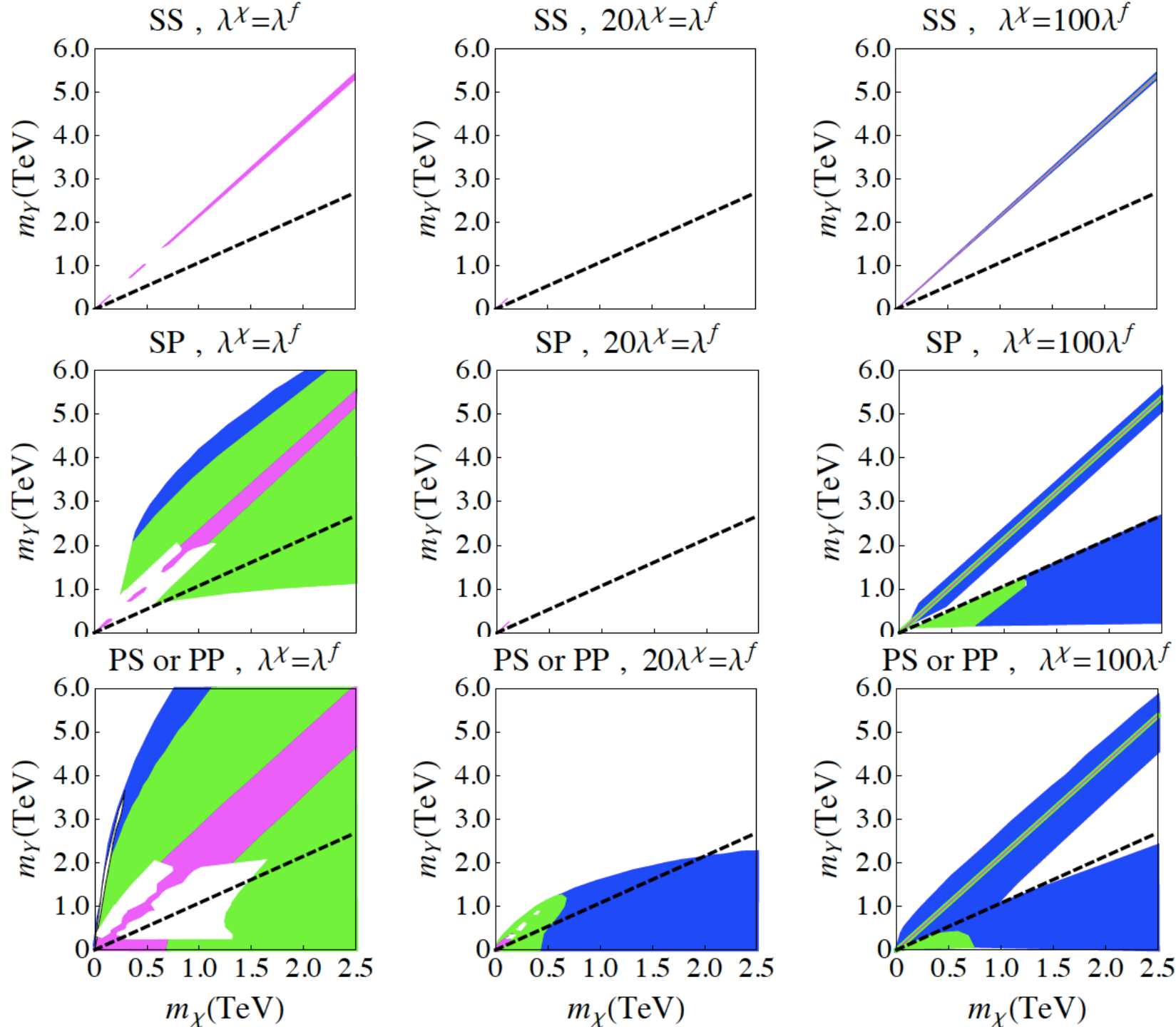
Other Constraints

- Indirect detection: No constraint on SS and SP cases due to v^2 suppression, but $m_\chi < 30$ GeV of PS and PP cases is disfavored by Fermi-LAT.
- Di-photon search: This loop effect is too small in our uniform coupling case.
- Y is fat when $\lambda_\chi, \lambda_f > 1 \Rightarrow \Gamma_Y > m_Y$
- Narrow width $\lambda_\chi, \lambda_f < 0.3 \Rightarrow \Gamma_Y < 0.1m_Y$





總結



Conclusions

- Fermionic DM couple universally to the SM quarks through scalar particle Y with scalar or pseudo-scalar coupling.
- Viable mass domain (m_χ, m_Y) is small for SS, SP cases but is large for PS, and PP cases.
- Direct detection favors DM pseudo-scalar coupling case due to velocity suppression.
- Mono-jet further constrain PS and PP further. Di-jet further constrain all cases.
- DM may be near if Y found.



Thank you very much.



Di-Jet Visible Y Search

SEARCHES FOR DIJET RESONANCES AT HADRON COLLIDERS *arXiv: 1110.5302*

ROBERT M. HARRIS (Fermilab) and KONSTANTINOS KOUSOURIS (CERN)

$$\sigma = \sum_{ij} \int dx_1 dx_2 f_i(x_1, \mu_F^2) f_j(x_2, \mu_F^2) \hat{\sigma}_{ij} \left(\alpha_s(\mu_R^2), \frac{Q^2}{\mu_F^2}, \frac{Q^2}{\mu_R^2} \right)$$

$$\frac{dL_{ij}}{d\tau} = \int_0^1 \int_0^1 dx_1 dx_2 f_i(x_1) f_j(x_2) \delta(x_1 x_2 - \tau)$$

$$\tau = x_1 x_2 = \frac{\hat{s}}{s}$$

$$\hat{\sigma}(m) = \frac{16\pi \times \mathcal{N} \times \mathcal{A}_{\cos\theta^*} \times BR \times \Gamma_R^2}{(m^2 - m_R^2)^2 + m_R^2 \Gamma_R^2}$$

$$\frac{1}{(m^2 - m_R^2)^2 + m_R^2 \Gamma_R^2} \approx \frac{\pi}{m_R \Gamma_R} \delta(m^2 - m_R^2)$$

$$\sigma_{had}(m_R) = 16\pi^2 \times \mathcal{N} \times \mathcal{A}_{\cos\theta^*} \times BR \times \left[\frac{1}{s} \frac{dL(\bar{y}_{min}, \bar{y}_{max})}{d\tau} \right]_{\tau=m_R^2/s} \times \frac{\Gamma_R}{m_R}$$



Y width

$$\Gamma(Y \rightarrow \bar{\chi}\chi)_S = \frac{1}{8\pi} m_Y \lambda_s^{\chi^2} \beta_\chi^3 \theta(m_Y - 2m_\chi),$$

$$\Gamma(Y \rightarrow \bar{\chi}\chi)_P = \frac{1}{8\pi} m_Y \lambda_p^{\chi^2} \beta_\chi \theta(m_Y - 2m_\chi),$$

$$\Gamma(Y \rightarrow \bar{f}f)_S = \frac{N_c}{8\pi} m_Y \lambda_s^{f^2} \beta_f^3,$$

$$\Gamma(Y \rightarrow \bar{f}f)_P = \frac{N_c}{8\pi} m_Y \lambda_p^{f^2} \beta_f,$$

$$E \sim m_\chi + \frac{3}{40} m_\chi$$

$$\beta_\chi = \sqrt{1 - 4m_\chi^2/m_Y^2}$$

$$\beta_f = \sqrt{1 - 4m_f^2/m_Y^2}$$

$$\Omega_D h^2 \sim 0.1 \left(\frac{3 \times 10^{-26} \text{ cm}^3 / \text{sec}}{\langle \sigma_{ann} v_{rel} \rangle} \right)$$

$$\Gamma_{SS,Y} = \Gamma(Y \rightarrow \bar{\chi}\chi)_S + \Gamma(Y \rightarrow \bar{f}f)_S,$$

$$\Gamma_{SP,Y} = \Gamma(Y \rightarrow \bar{\chi}\chi)_S + \Gamma(Y \rightarrow \bar{f}f)_P,$$

$$\Gamma_{PS,Y} = \Gamma(Y \rightarrow \bar{\chi}\chi)_P + \Gamma(Y \rightarrow \bar{f}f)_S,$$

$$\Gamma_{PP,Y} = \Gamma(Y \rightarrow \bar{\chi}\chi)_P + \Gamma(Y \rightarrow \bar{f}f)_P,$$

$$\lambda_\chi, \lambda_f > 1 \Rightarrow \Gamma_Y > m_Y$$

$$\lambda_\chi, \lambda_f < 0.3 \Rightarrow \Gamma_Y < 0.1 m_Y$$

Narrow width



χ - N Elastic Cross-section

$$\sigma_{el}^{SS}(\chi N \rightarrow \chi N) = \frac{\lambda_s^2 \lambda_s^2}{\pi} \frac{m_\chi^2 m_N^2}{(m_\chi + m_N)^2 m_Y^4} f_N^2,$$

$$\mu \equiv m_\chi m_N / (m_\chi + m_N)$$

$$\sigma_{el}^{SP}(\chi N \rightarrow \chi N) = \frac{\lambda_s^2 \lambda_p^2}{\pi} \frac{m_\chi^2 p^2}{2(m_\chi + m_N)^2 m_Y^4} g_N^2,$$

$$p = \mu v$$

$$\sigma_{el}^{PS}(\chi N \rightarrow \chi N) = \frac{\lambda_p^2 \lambda_s^2}{\pi} \frac{p^2 m_N^2}{2(m_\chi + m_N)^2 m_Y^4} f_N^2,$$

$$v \sim 10^{-3} c$$

$$\sigma_{el}^{PP}(\chi N \rightarrow \chi N) = \frac{\lambda_p^2 \lambda_p^2}{\pi} \frac{p^4}{3(m_\chi + m_N)^2 m_Y^4} g_N^2,$$

$$\langle N | m_q \bar{q} q | N \rangle = m_N f_{T_q}^{(N)} \bar{u}_N u_N$$

with $m_D = 100$ mN

$$f_N = \sum_{q=u,d,s} f_{T_q}^{(N)} \frac{m_N}{m_q} + \frac{2}{27} f_{T_G}^{(N)} \sum_{Q=c,b,t} \frac{m_N}{m_Q}$$

$$\begin{aligned} f_{T_u}^{(p)} &= 0.023, & f_{T_d}^{(p)} &= 0.034, & f_{T_s}^{(p)} &= 0.14, & f_{T_G}^{(p)} &= 0.803, \\ f_{T_u}^{(n)} &= 0.019, & f_{T_d}^{(n)} &= 0.041, & f_{T_s}^{(n)} &= 0.14, & f_{T_G}^{(n)} &= 0.8. \end{aligned}$$

Dark SUSY

$$g_N = (1 - \eta) \Delta u^{(N)} \frac{m_N}{m_u} + (1 - \eta z) \Delta d^{(N)} \frac{m_N}{m_d} + (1 - \eta w) \Delta s^{(N)} \frac{m_N}{m_s}$$

$$\Delta u^{(p)} = \Delta d^{(n)} = 0.77,$$

$$\Delta d^{(p)} = \Delta u^{(n)} = -0.40,$$

$$\Delta s^{(p)} = \Delta s^{(n)} = -0.12.$$

$$\langle N' | \bar{q} i \gamma_5 q | N \rangle \simeq (1 - \eta \delta_q) \Delta q^{(N)} \frac{m_N}{m_q} \bar{u}'_N i \gamma_5 u_N$$

where $\delta_u = 1$, $\delta_d = z$, and $\delta_s = w$ with $\eta = (1 + z + w)^{-1}$, $z = m_u/m_d$, and $w = m_u/m_s$.

