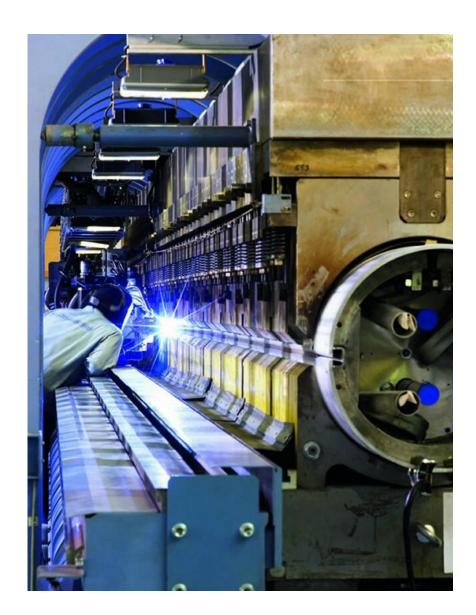
### Lecture 5: Practical matters

#### Plan

- LHC quench protection
- current leads
- accelerator coil winding and curing
- forces and clamping
- magnet assembly, collars and iron
- installation
- some superconducting accelerators



Martin Wilson Lecture 5 slide1 JUAS February 2013

# LHC dipole protection: practical implementation

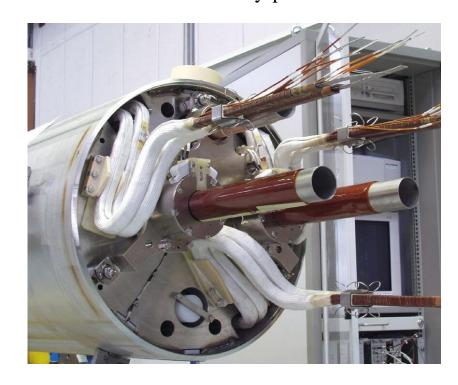
#### It's difficult! - the main challenges are:

#### 1) Series connection of many magnets

- In each octant, 154 dipoles are connected in series. If one magnet quenches, the combined energy of the others will be dumped in that magnet ⇒ vaporization!
- Solution 1: cold diodes across the terminals of each magnet. Diodes normally block  $\Rightarrow$  magnets track accurately. If a magnet quenches, it's diodes conduct  $\Rightarrow$  octant current by-passes.
- Solution 2: open a circuit breaker onto a resistor (several tonnes) so that octant energy is dumped in ~ 100 secs.

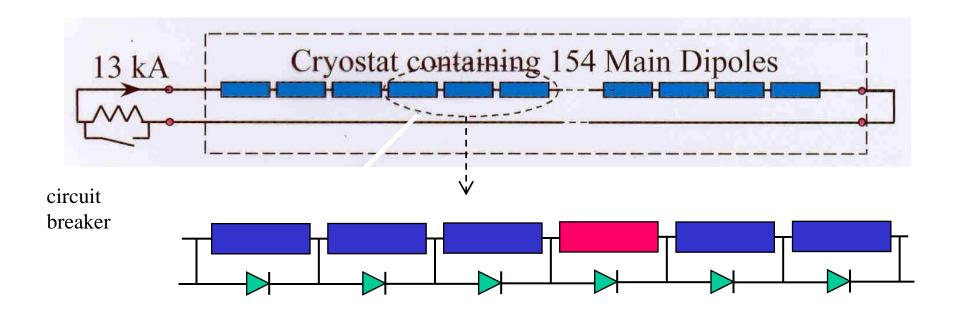
## 2) High current density, high stored energy and long length

- Individual magnets may burn out even when quenching alone.
- Solution 3: Quench heaters on top and bottom halves of every magnet.



Martin Wilson Lecture 5 slide2 JUAS February 2013

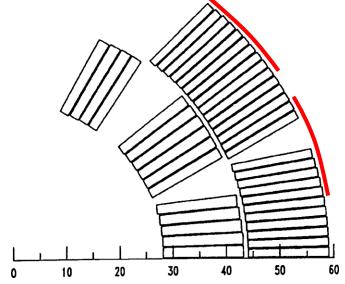
## LHC power supply circuit for one octant

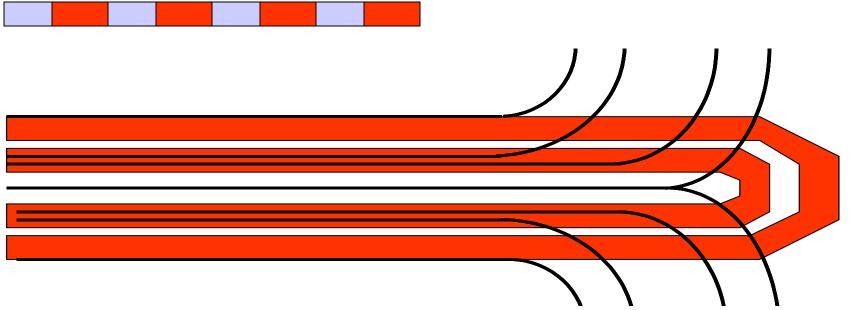


- in normal operation, diodes block ⇒ magnets track accurately
- if a magnet quenches, diodes allow the octant current to by-pass
- circuit breaker reduces to octant current to zero with a time constant of 100 sec
- initial voltage across breaker = 2000V
- stored energy of the octant = 1.33GJ

## LHC quench-back heaters

- stainless steel foil 15mm x 25  $\mu$ m glued to outer surface of winding
- insulated by Kapton
- pulsed by capacitor  $2 \times 3.3 \text{ mF}$  at 400 V = 500 J
- quench delay at rated current = 30msec
  - at 60% of rated current = 50msec
- copper plated 'stripes' to reduce resistance





Martin Wilson Lecture 5 slide4 JUAS February 2013

# Diodes to by-pass the main ring current

Installing the cold diode package on the end of an LHC dipole





Martin Wilson Lecture 5 slide5 JUAS February 2013

#### Current Leads

#### **Optimization**

• want to have low heat inleak, ie low ohmic heating *and* low heat conduction from room temperature. This requires low  $\rho$  and k - but Wiedemann Franz says

$$k(\theta)\rho(\theta) = L_o\theta$$

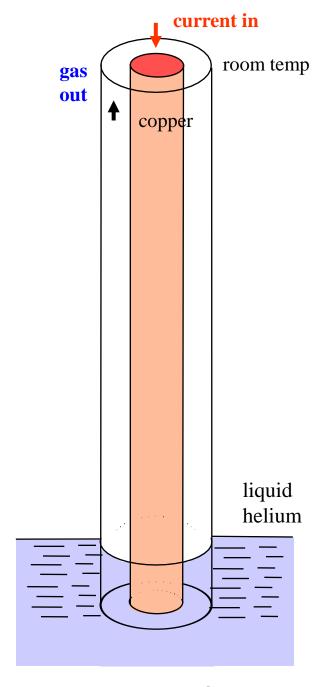
so all metals are the same and the only variable we can optimize is the *shape*

#### Gas cooling helps (recap helium properties Lecture 4)

•  $\Delta$ enthalpy gas / latent heat of boiling = 73.4 - lots more cold in the boil off gas

$$\Delta H = \int_{4.7}^{293} C(\theta) d\theta$$

- so use the enthalpy of the cold gas which is boiled off to cool the lead
- we make the lead as a heat exchanger



Martin Wilson Lecture 5 slide6 JUAS February 2013

# Current lead theory

equation of heat conduction

$$\frac{d}{dx}\left(k(\theta)A\frac{d\theta}{dx}\right) - f\dot{m}C_{p}\frac{d\theta}{dx} + \frac{I^{2}\rho(\theta)}{A} = 0$$

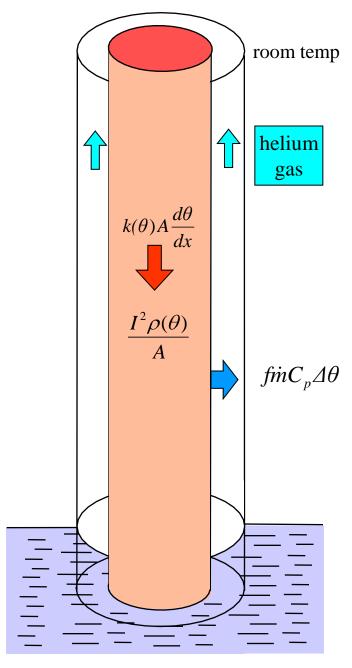
#### where:

f = efficiency of heat transfer to helium gas

 $\dot{m}$  = helium mass flow

 $C_p$  = specific heat of gas

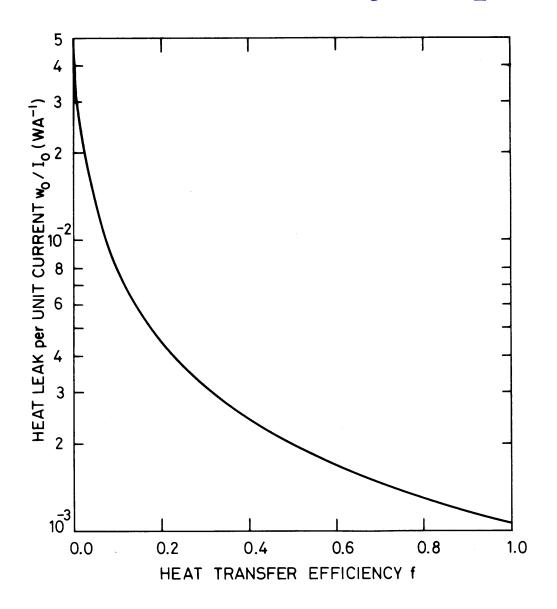
- solution to this equation in 'Superconducting Magnets p 257.
- there is an optimum shape (length/area) which gives the minimum heat leak
  - 'Watts per Amp per lead'
- heat leak is a strong function of the efficiency of heat transfer f to the cold gas



Martin Wilson Lecture 5 slide7

JUAS February 2013

## Heat leak of an optimised lead



 with optimum shape and 100% efficient heat transfer the heat leak is

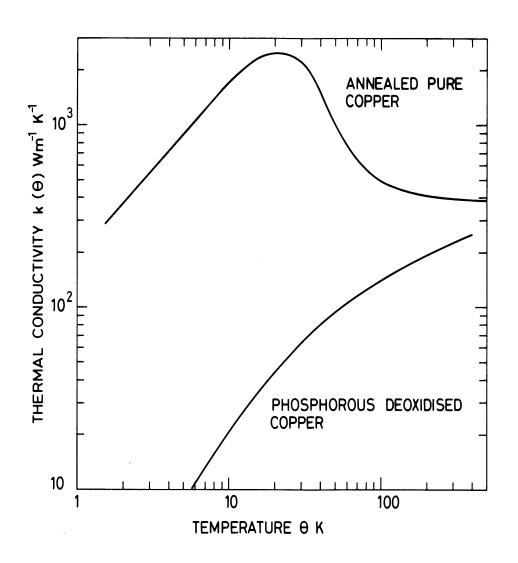
#### 1.04 mW/Amp

per lead

- with optimum shape and no heat transfer the heat leak is
   47 mW/Amp
- Note the optimum shape varies with the heat transfer efficiency

Martin Wilson Lecture 5 slide8 JUAS February 2013

# Optimum shape of lead



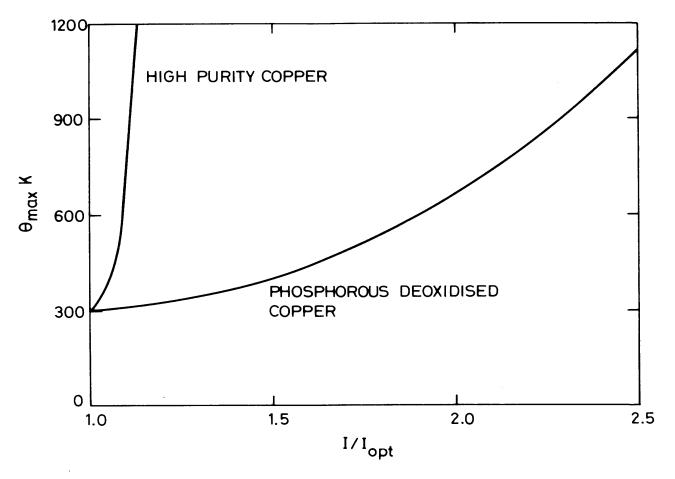
- the optimum shape depends on temperature and material properties, particularly thermal conductivity.
- for a lead between 300K and 4.2K the optimum shape is
- for a lead of annealed high purity copper

$$\left\{\frac{L}{A}\right\}_{optimum} = \frac{2.6x10^7}{I}$$

 for a lead of impure phosphorous deoxised copper (preferred)

$$\left\{\frac{L}{A}\right\}_{optimum} = \frac{3.5x10^6}{I}$$

## Impure materials make more stable leads

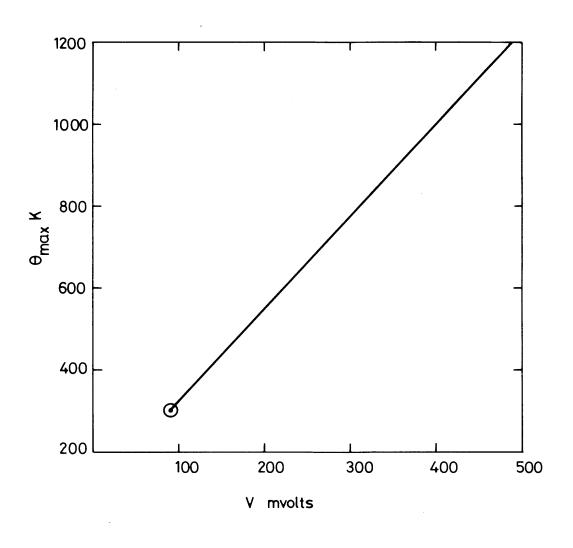


if current lead burns out ⇒ magnet open circuit ⇒ large voltages ⇒ disaster

- for an optimized lead, the maximum temperature is room temperature (at the top of the lead)
- when the lead is not optimized, the temperature of an intermediate region rises above room temperature
- the optimum for pure metals is more sensitive than for impure metals

Martin Wilson Lecture 5 slide10 JUAS February 2013

## Health monitoring



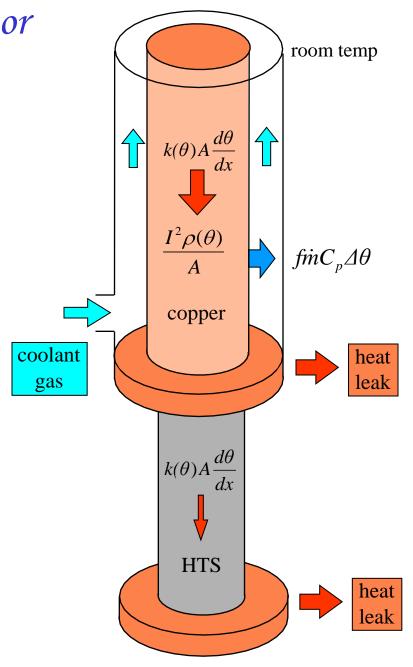
- all leads between the same temperatures and with the same cooling efficiency drop the same voltage at optimum
- for a lead between 300K and 4.2K with with 100% cooling efficiency, the voltage drop at optimum is 75mV
- measure the volts across your lead to see if it is optimised
- if a lead burns out, the resulting high voltage and arcing (magnet inductance) can be disastrous
- monitor your lead and trip the power supply if it goes too high

# High temperature superconductor HTS Current leads

- at temperatures below 50 -70K can use HTS
- material has very low thermal conductivity
- no Ohmic heat generation
- but from room temperature to 50 70 K must have copper leads
- the 50 70 K junction must be cooled or its temperature will drift up and quench the HTS

#### For the HTS section beware of

- overheating if quenches
- fringe field from magnet



Martin Wilson Lecture 5 slide12 JUAS February 2013

## HTS (high temperature superconductor) current leads



- HTS materials have a low thermal conductivity
- make the section of lead below ~ 70K from HTS material
- heat leak down the upper lead is similar, but it is taken at a higher temperature
  - $\Rightarrow$  less refrigeration power
- LHC uses HTS leads for all main ring magnets
- savings on capital cost of the refrigerator > cost of the leads
- reduced running cost is a continuing benefit

*⇐13kA lead for LHC* 

600A lead for LHC  $\Rightarrow$ 



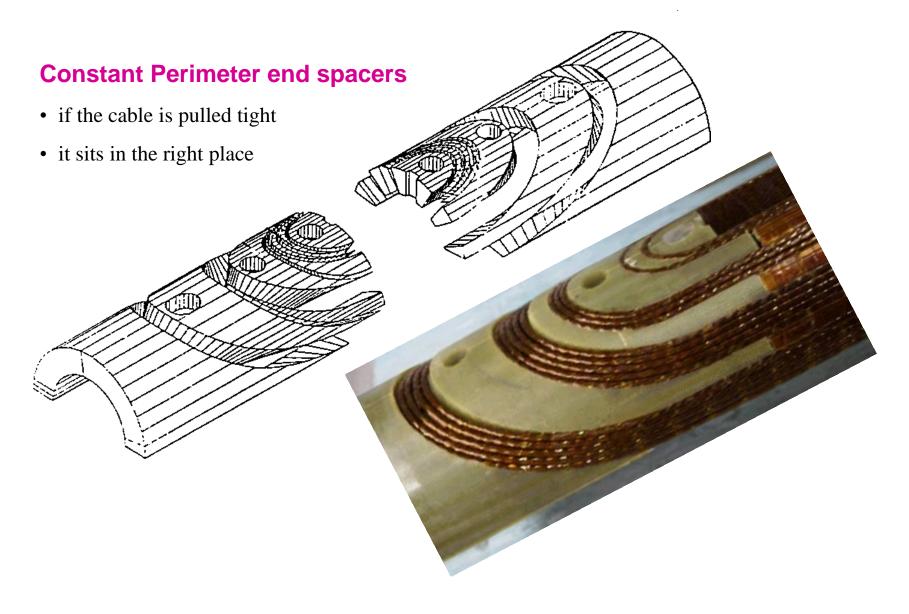
Martin Wilson Lecture 5 slide13 JUAS February 2013

# Winding the LHC dipoles



Martin Wilson Lecture 5 slide14 JUAS February 2013

# End turns

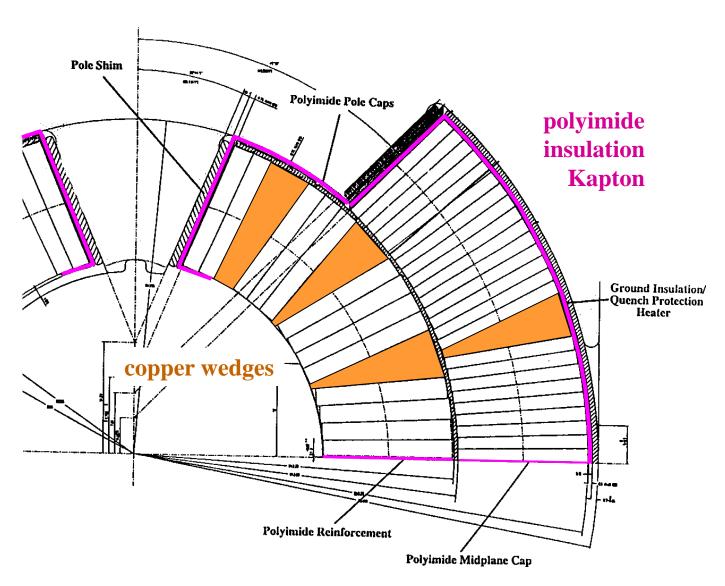


Martin Wilson Lecture 5 slide15

JUAS February 2013

# Spacers and insulation

- copper wedges between blocks of winding
- beware of voltages at quench
- care needed with insulation, between turns and ground plane
- example: FAIR dipole quench voltage = 340V over 148 turns



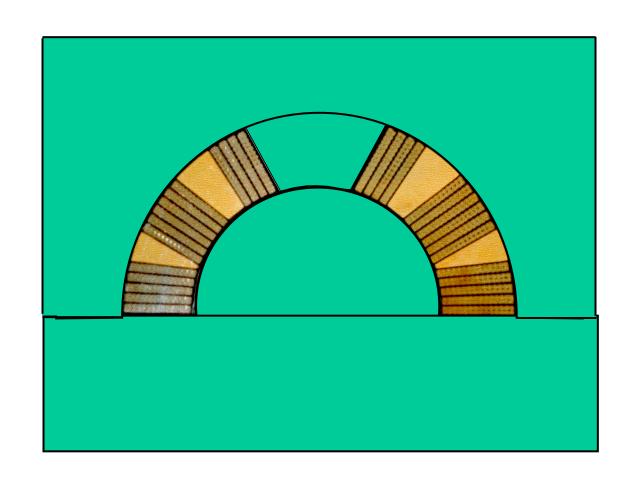
Martin Wilson Lecture 5 slide16

JUAS February 2013

# Compacting and curing

 After winding, the half coil, (still very 'floppy') is placed in an accurately machined tool

- Tool put into a curing press, compacted to the exact dimensions and heated to 'cure' the polyimide adhesive on the Kapton insulation.
- After curing, the half coil is quite rigid and easy to handle

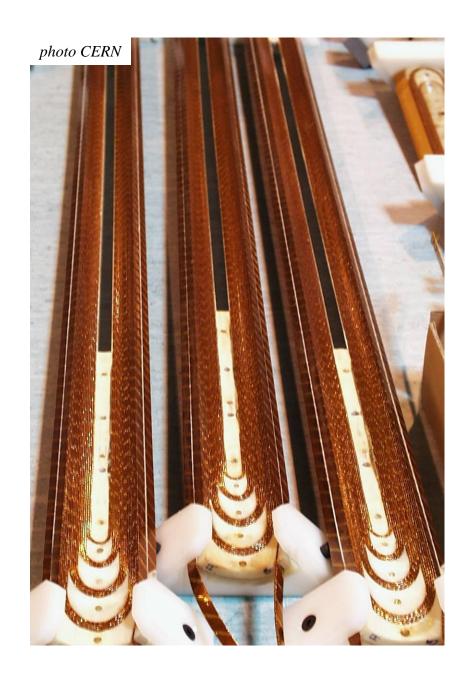


Martin Wilson Lecture 5 slide17 JUAS February 2013

# Curing press



Martin Wilson Lecture 5 slide18 JUAS February 2013



### Finished coils

after curing, the coil package is rigid and relatively easy to handle



Martin Wilson Lecture 5 slide19

JUAS February 2013

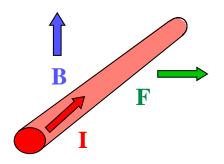
# Coils for correction magnets



On a smaller scale, but in great number and variety, many different types of superconducting correction coils are needed at a large accelerator

Martin Wilson Lecture 5 slide20 JUAS February 2013

# Electromagnetic forces in dipoles

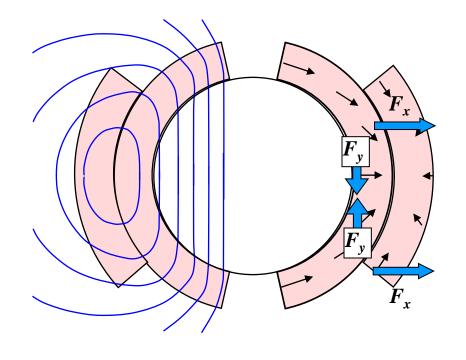


$$\underline{F} = \underline{B} \wedge \underline{I}$$

- forces in a dipole are horizontally outwards and vertically towards the median plane
- recap lecture 2 slide 12, for a *thin* winding

total outward force *per quadrant* 

$$F_x = \frac{B_i^2}{2\mu_o} \frac{4a}{3}$$



LHC dipole  $F_x \sim 1.6 \times 10^6 \,\text{N/m} = 160 \,\text{tonne/m}$ 

total vertical force *per quadrant* 

$$F_y = -\frac{B_i^2}{2\mu_o} \frac{4a}{3}$$

- the outward force must be supported by an external structure
- $F_x$  and  $F_y$  cause compressive stress in the conductor and insulation
- apart from the ends, there is no tension in the conductor

for thick winding take ~ mean radius or use MNW eq 4.24

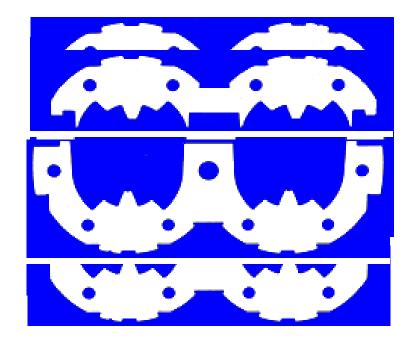
### Collars

**Question:** how to make a force support structure that

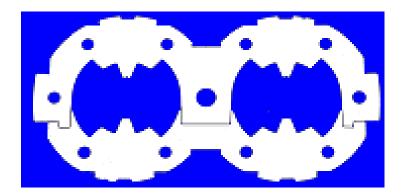
- fits tightly round the coil
- presses it into an accurate shape
- has low ac losses laminated
- can be mass produced cheaply

**Answer:** make collars by precision stamping of stainless steel or aluminium alloy plate a few mm thick

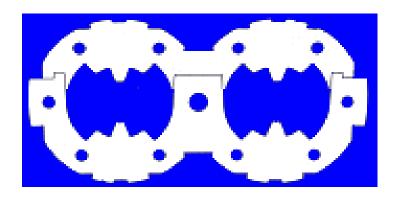
- inherited from conventional magnet laminations



press collars over coil from above and below



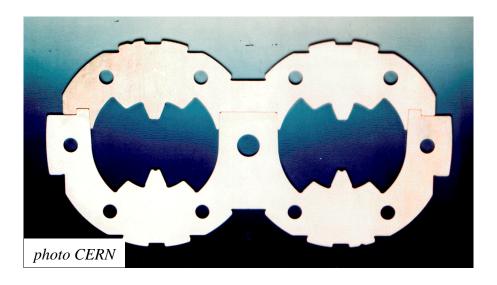
invert alternate pairs so that they interlock



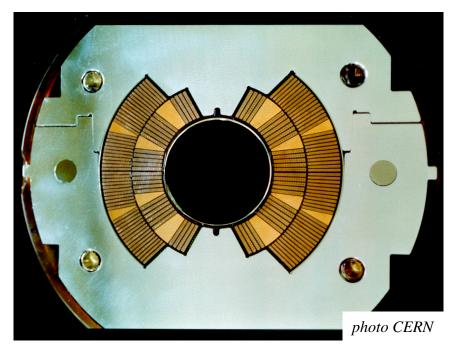
push steel rods through holes to lock in position

### Collars

LHC dipole collars support the twin aperture coils in a single unit



12 million produced for LHC

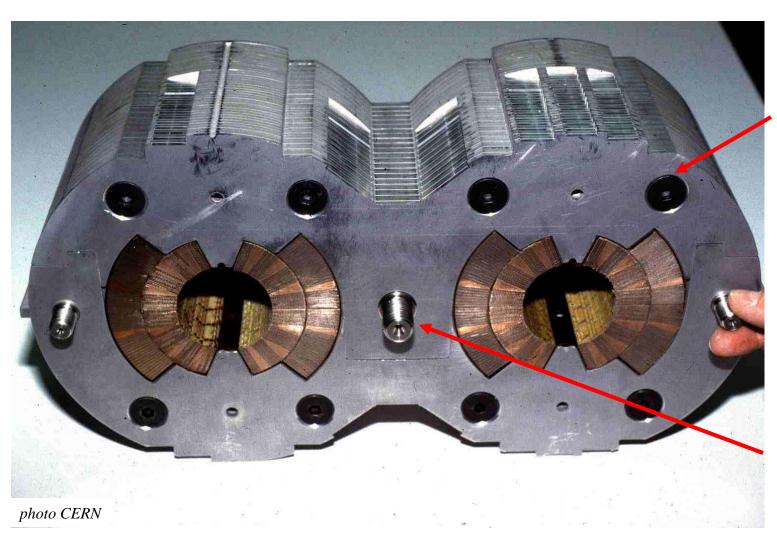




Martin Wilson Lecture 5 slide23

JUAS February 2013

# LHC dipole collars

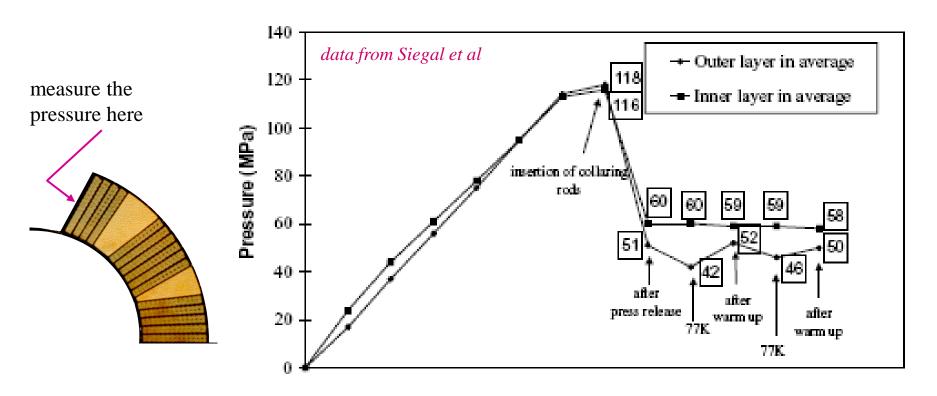


sub-units of several alternating pairs are riveted together

stainless rods lock the subunits together

Martin Wilson Lecture 5 slide24 JUAS February 2013

# Pre-loading the coil



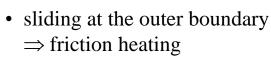
#### CERN data during manufacture and operation

#### data from Modena et al

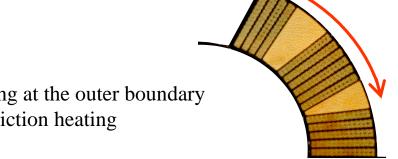
	after collaring at 293K		after yoking at 293K		at 1.9K		at 1.9K and 8.3T	
	inner	outer	inner	outer	inner	outer	inner	outer
MBP2N2	62Mpa	77Mpa	72Mpa	85Mpa	26MPa	32MPa	2MPa	8Mpa
MBP2O1	51MPa	55MPa	62MPa	62MPa	24MPa	22MPa	0MPa	2MPa

Collars and end plate (LHC dipole)



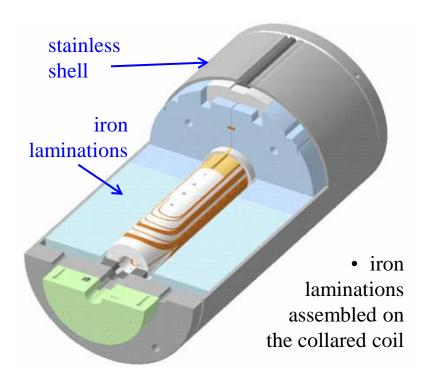


use kapton layers



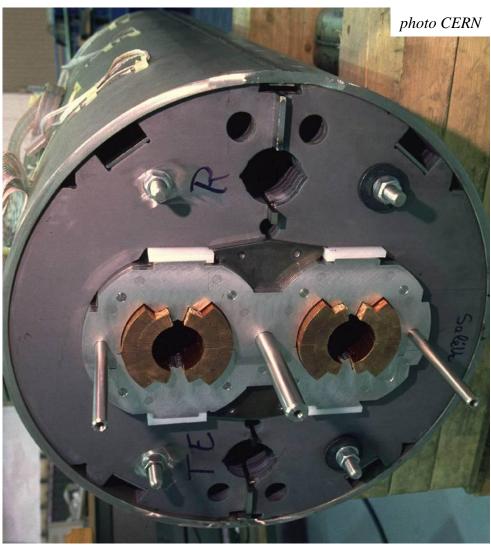
JUAS February 2013 Martin Wilson Lecture 5 slide26

photo CERN



- pushed into place using the collaring press
- **BUT** pure iron becomes brittle at low temperature
- tensile forces are therefore taken by a stainless steel shell which is welded around the iron, while still in the press
- stainless shell also serves as the helium vessel

# Adding the iron

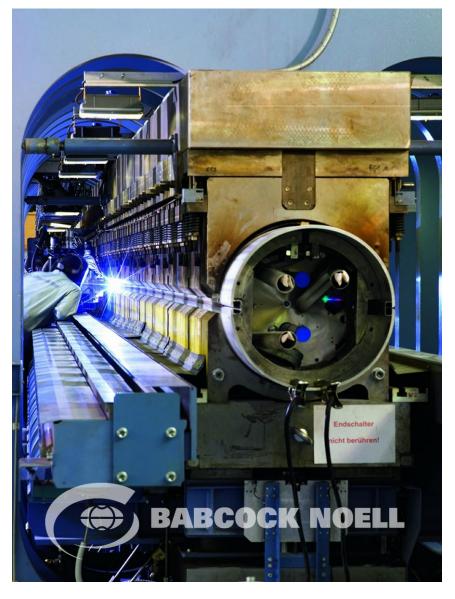


Martin Wilson Lecture 5 slide27

JUAS February 2013

# Compressing and welding the outer shell

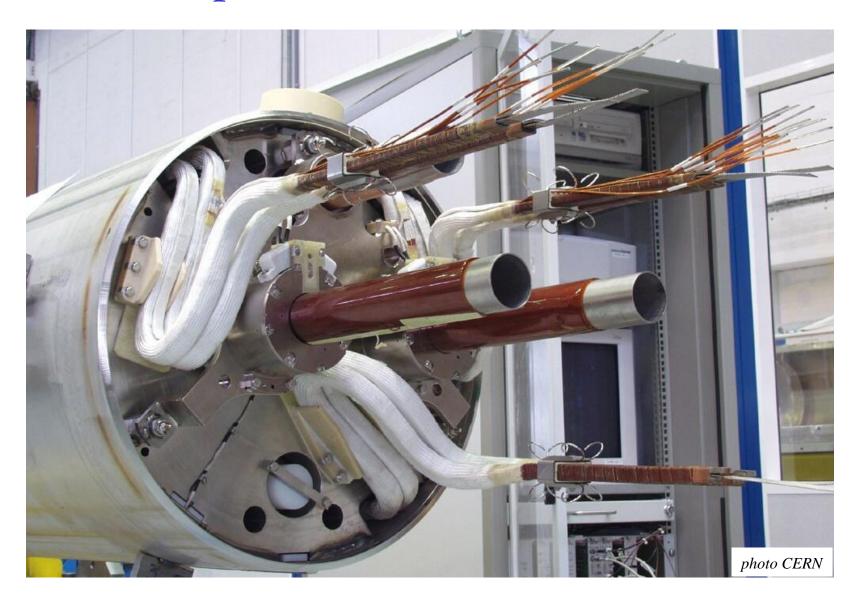




Martin Wilson Lecture 5 slide28

JUAS February 2013

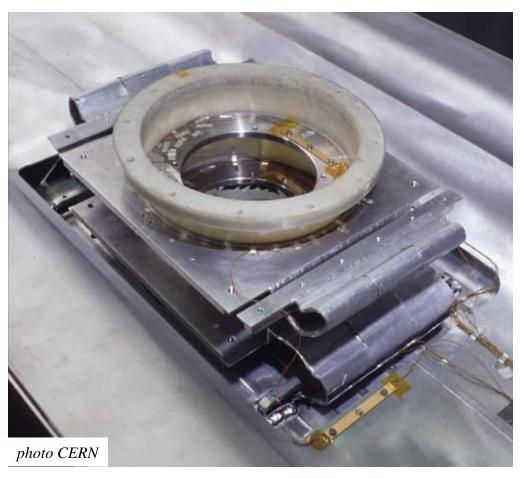
# Dipole inside its stainless shell



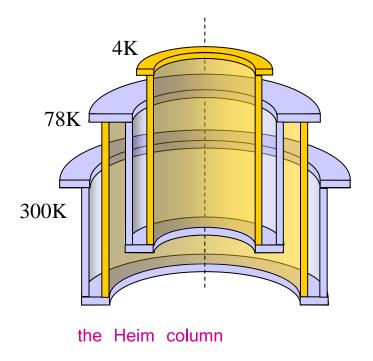
Martin Wilson Lecture 5 slide29

JUAS February 2013

## Cryogenic supports



'feet' used to support cold mass inside cryostat (LHC dipole)

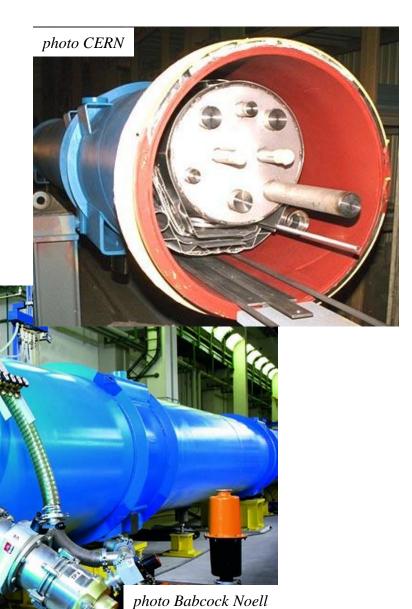


- long path length in short distance
- mechanical stiffness of tubes
- by choosing different material contractions can achieve zero thermal movement

Martin Wilson Lecture 5 slide30 JUAS February 2013



## Complete magnet in cryostat

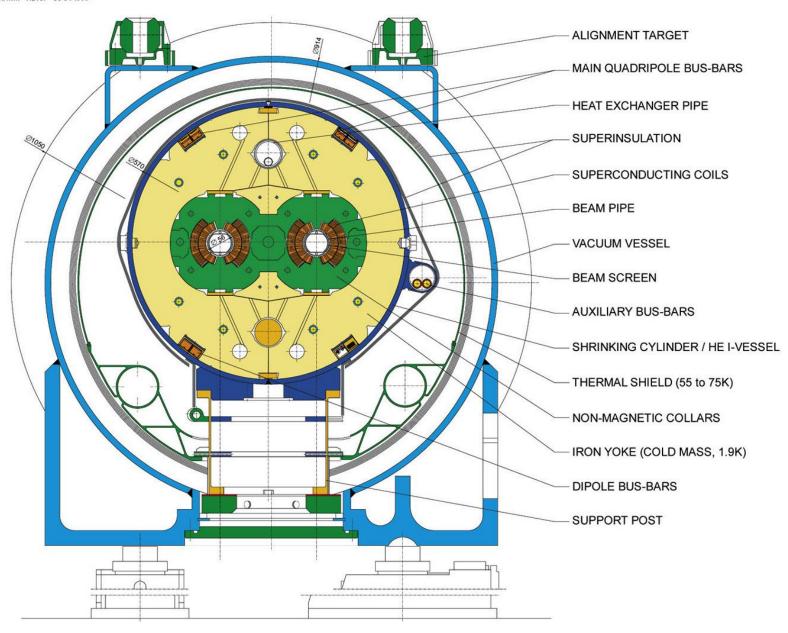


Martin Wilson Lecture 5 slide31

JUAS February 2013

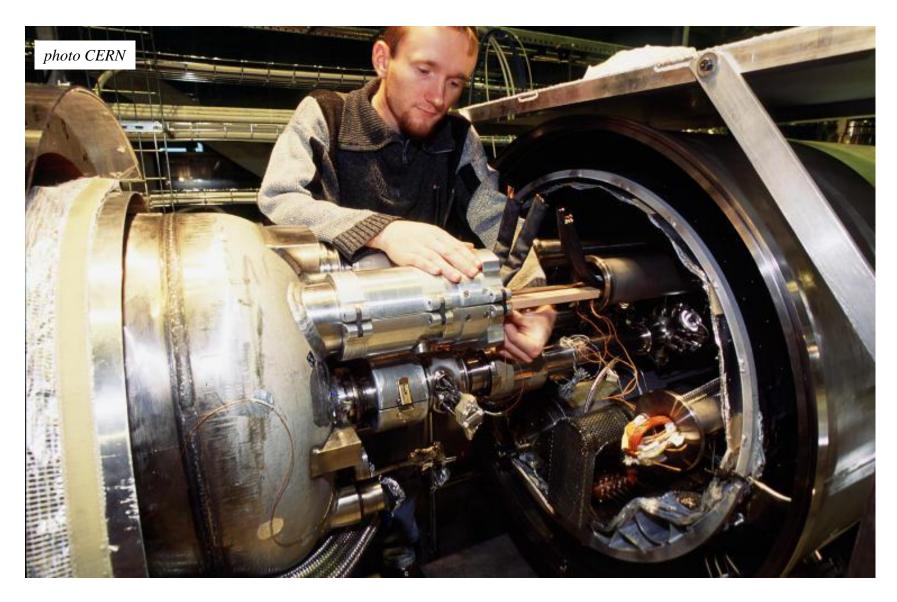
#### LHC DIPOLE: STANDARD CROSS-SECTION

CERN AC/DI/MM - HE107 - 30 04 1999



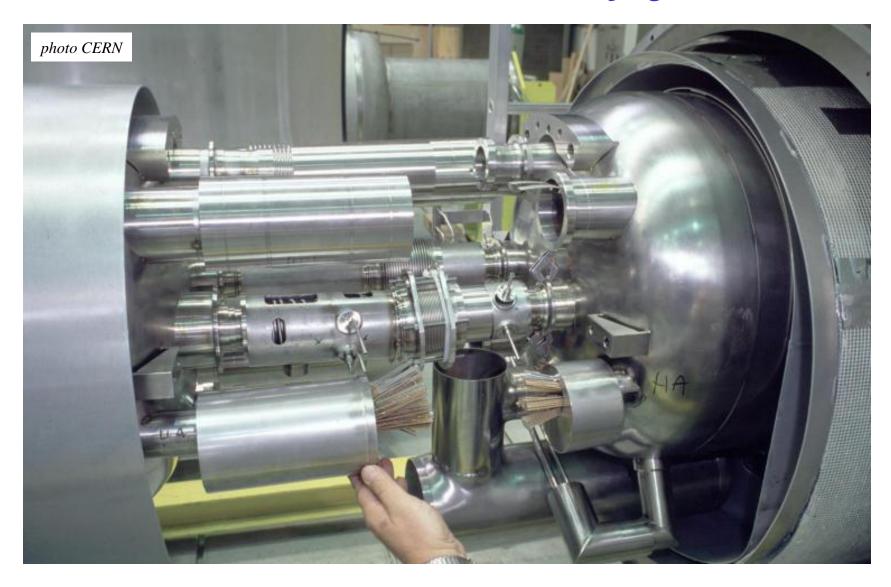
Martin Wilson Lecture 5 slide32 JUAS February 2013

# Make the interconnections - electrical



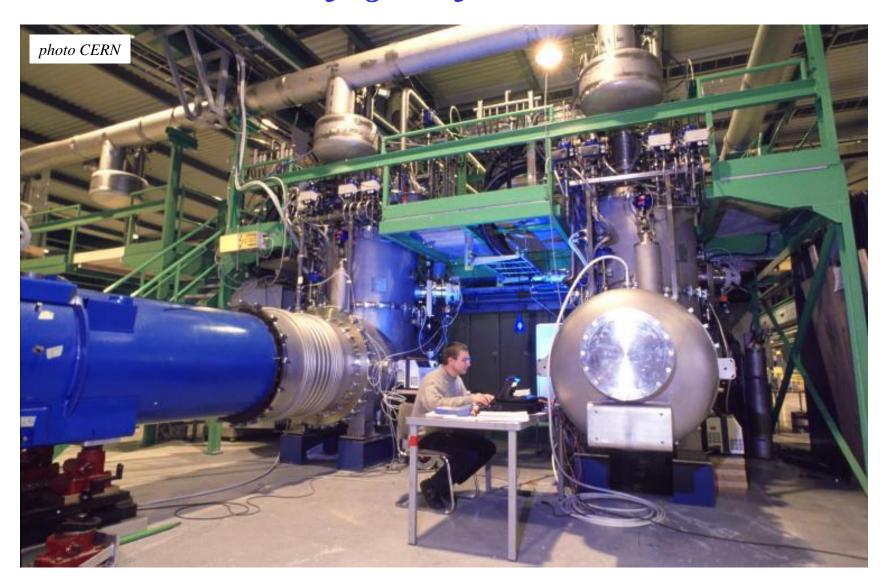
Martin Wilson Lecture 5 slide33 JUAS February 2013

# Make interconnections - cryogenic



Martin Wilson Lecture 5 slide34 JUAS February 2013

# Connect to the cryogenic feed and current leads



Martin Wilson Lecture 5 slide35 JUAS February 2013