

Cosmic Strings post Planck

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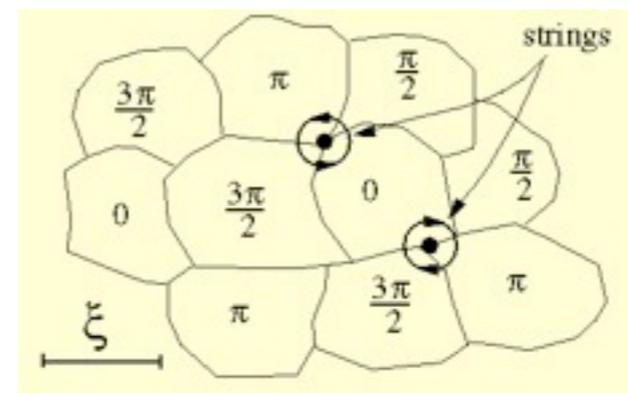
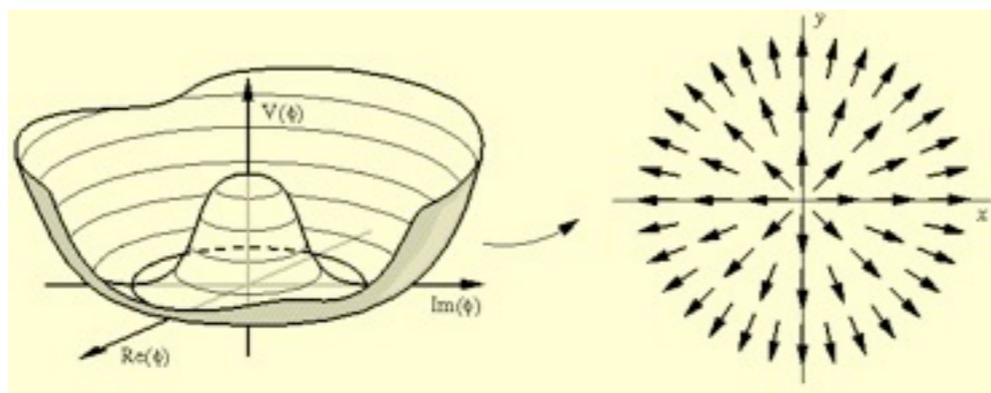
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Cosmic Strings



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- ▶ Strings can form during phase transitions in early Universe

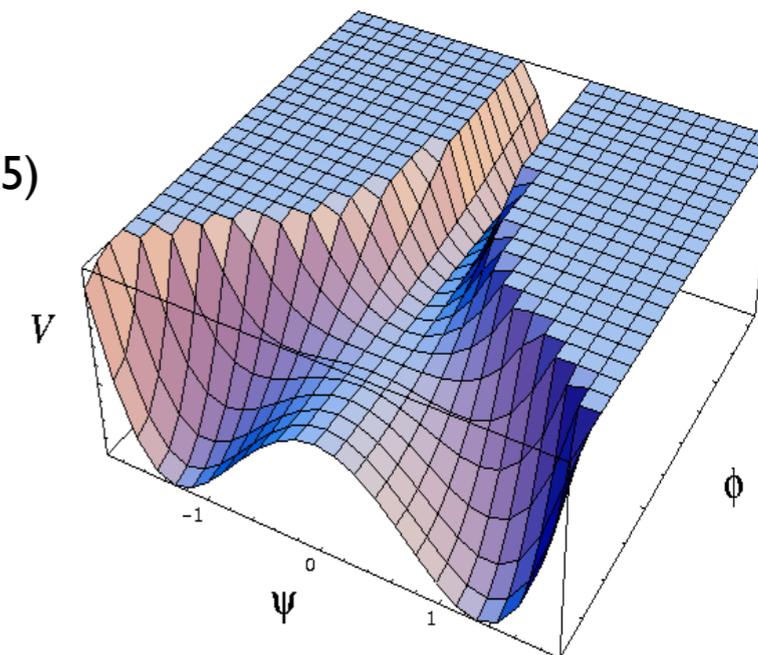


(Kibble 1976)

- ▶ Ruled out as primary mechanism to produce anisotropies

(E.g. Polchinski 2005)

- ▶ Formed in hybrid, brane inflation etc models
- ▶ GUT scale phase transition gives $\Delta T/T \approx 10^{-5}$
- ▶ Potential window to high energy physics



Observables



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- ▶ CMB power spectrum. Includes recombination and post-recombination physics. Strongest limit on allowed fraction of strings, improved by Planck
- ▶ Non Gaussianity in CMB maps. Search for signatures of post-recombination Doppler shift induced by moving strings (Kaiser-Stebbins effect). Strongest limit also from Planck
- ▶ CMB B modes. Defects produce comparable scalar, vector and tensor fluctuations.
- ▶ Gravitational lensing. Search for microlensing events (not discussed here)
- ▶ Pulsar timing/gravitational waves. Stochastic background from loops, waves from cusps (not discussed here)

String Evolution



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- ▶ Perturbations from strings are **active**, so their evolution is key to understanding CMB anisotropies
- ▶ Main idea is strings evolve toward self-similar **scaling** regime
- ▶ Average properties of network are (nearly) constant with time
- ▶ Dynamics can be studied using numerical simulations
- ▶ Two approaches - **Nambu** and **Abelian-Higgs** models
- ▶ Both have advantages and disadvantages
- ▶ Main issue is **dynamical range** - assumptions have to be made in either case
- ▶ Will present Planck constraints for each case

Nambu Model



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- ▶ Thin string approximation
- ▶ Ignore radiation back reaction
- ▶ Impose reconnection by hand
- ▶ Network characterised by **correlation length** L
- ▶ Energy density is $\rho = \frac{\mu}{L^2}$,
- ▶ Observationally, string tension μ is main quantity of interest
- ▶ Find scaling solution $L \sim t$
- ▶ Make measurements of correlation length, velocity, small scale structure (wiggleness)

(Martin and Shellard)

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(Martin and Shellard)

VOS Model



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- ▶ VOS = Velocity One-Scale model (Kibble 1985, Martin and Shellard 1996/2000)
- ▶ Expansion stretches strings, long strings reconnect and loops decay
- ▶ From Nambu action

Loop production, calibrated from simulations Curvature term, proportional to RMS velocity v

$$\frac{1}{L} \frac{dL}{dt} = (1 + v^2)H + \frac{\tilde{c}v}{2L}, \quad \frac{dv}{dt} = (1 - v^2) \left(\frac{\tilde{k}}{L} - 2Hv \right),$$

- ▶ For power law expansion $a(t) \propto t^\beta$, find attractor solution with scaling

$$L = \epsilon t, \quad \epsilon = \sqrt{\frac{\tilde{k}(\tilde{k} + \tilde{c})}{4\beta(1 - \beta)}}, \quad v = \sqrt{\frac{\tilde{k}(1 - \beta)}{\beta(\tilde{k} + \tilde{c})}}$$

- ▶ More complicated VOS models can be constructed for superstrings, networks with junctions etc.

- ▶ Will use comoving correlation length $l = \frac{L}{a} = \xi \tau,$

VOS Parameters



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- ▶ Simulations find network parameters vary between matter and radiation eras
- ▶ Density in radiation era is greater than in matter era
- ▶ Simulations of Martins and Shellard find

$$\xi_{\text{rad}} = 0.13 \rightarrow \xi_{\text{mat}} = 0.21 \quad v_{\text{rad}} = 0.65 \rightarrow v_{\text{mat}} = 0.60$$

- ▶ Ringeval et al find

$$\xi_{\text{rad}} = 0.16 \rightarrow \xi_{\text{mat}} = 0.19$$

- ▶ Strings also have small scale structure, or 'wiggleness' (Carter 2000)
- ▶ Effective coarse grained energy momentum tensor gives rescaled mass per unit length $U = \alpha \mu$,
- ▶ Estimated to be

$$\alpha_{\text{rad}} = 1.5 \rightarrow \alpha_{\text{mat}} = 1.9$$

Abelian Higgs



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- ▶ Lagrangian $\mathcal{L} = (D_\mu \Phi)^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - V(\Phi)$
- ▶ Solve equations of motion on 3D lattice (Bevis et al)
- ▶ Extract unequal time correlator (UETC) of energy momentum tensor
$$\langle \Theta(k, \tau_1) \Theta(k, \tau_2) \rangle$$
- ▶ Radiation into propagating modes included
- ▶ However, simulation box sizes limited to $\sim 300x$ string core width
- ▶ Requires slowing down rate of growth of core
- ▶ Found only very small loop production - most of energy loss to radiation
- ▶ See no observed small scale structure
- ▶ Similar network parameters in matter and radiation eras
$$\xi = 0.3 \quad v = 0.5 \quad (\text{Hindmarsh et al 2009})$$
- ▶ More work required to resolve differences with Nambu model

CMB Anisotropies



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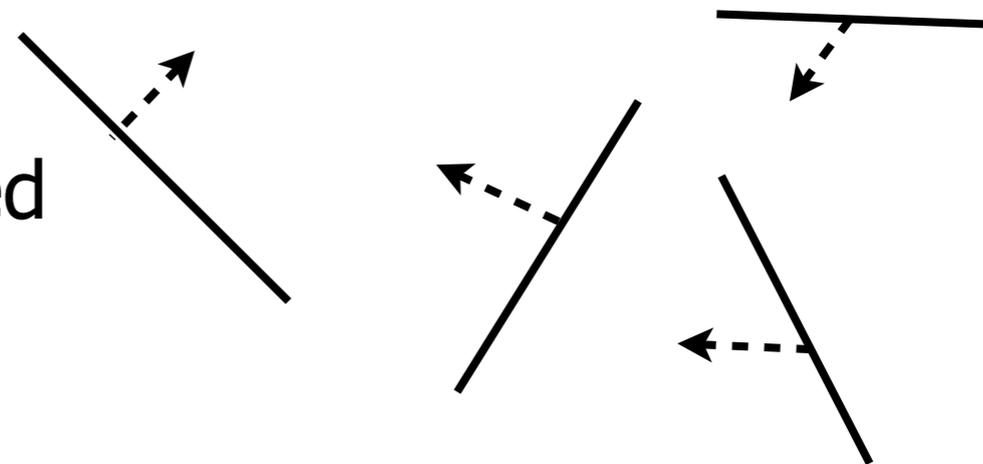
- ▶ **Key** ingredient is the UETC
- ▶ Use stiff source approximation

String energy momentum tensor



$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \quad \delta G_{\mu\nu} = 8\pi G (\delta T_{\mu\nu} + \theta_{\mu\nu})$$

- ▶ Can estimate UETC directly from simulations, and use as sources in CMB codes
- ▶ For thin strings there is a useful intermediate framework called Unconnected Segment Model (USM) (Vincent et al, Albrecht et al, Pogosian and Vachaspati)
- ▶ Model strings are ensemble of uncorrelated straight segments, each moving with random velocity
- ▶ Inputs are correlation length, velocity and wiggleness, as measured from simulations

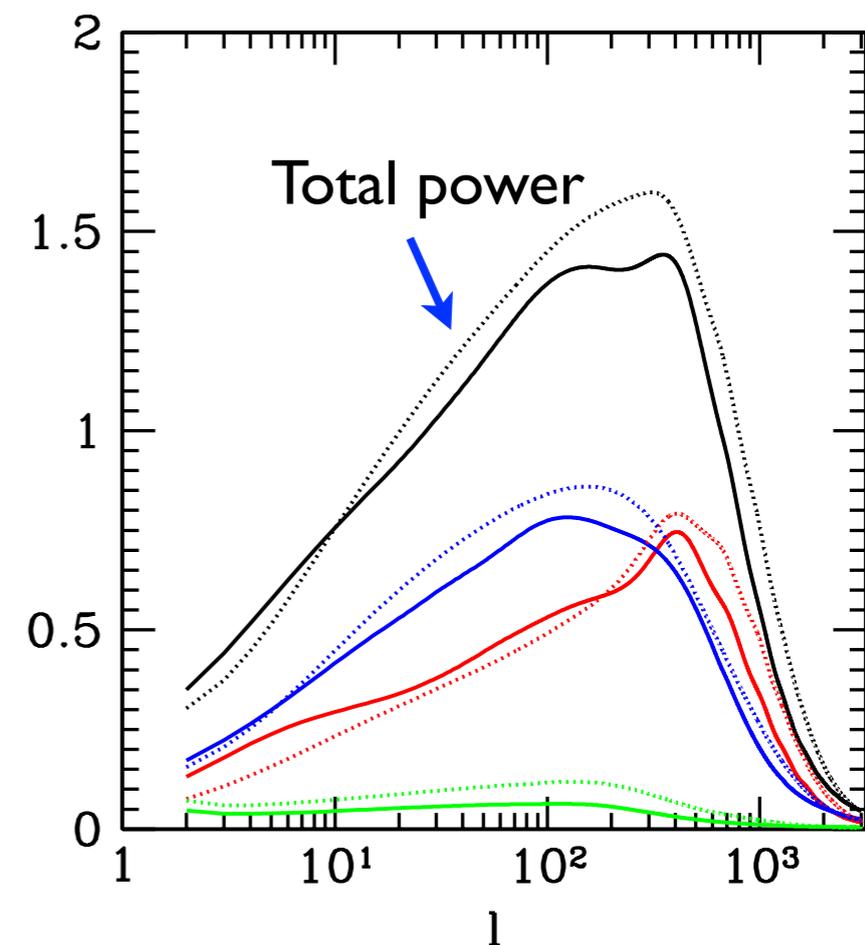


Mimic Model



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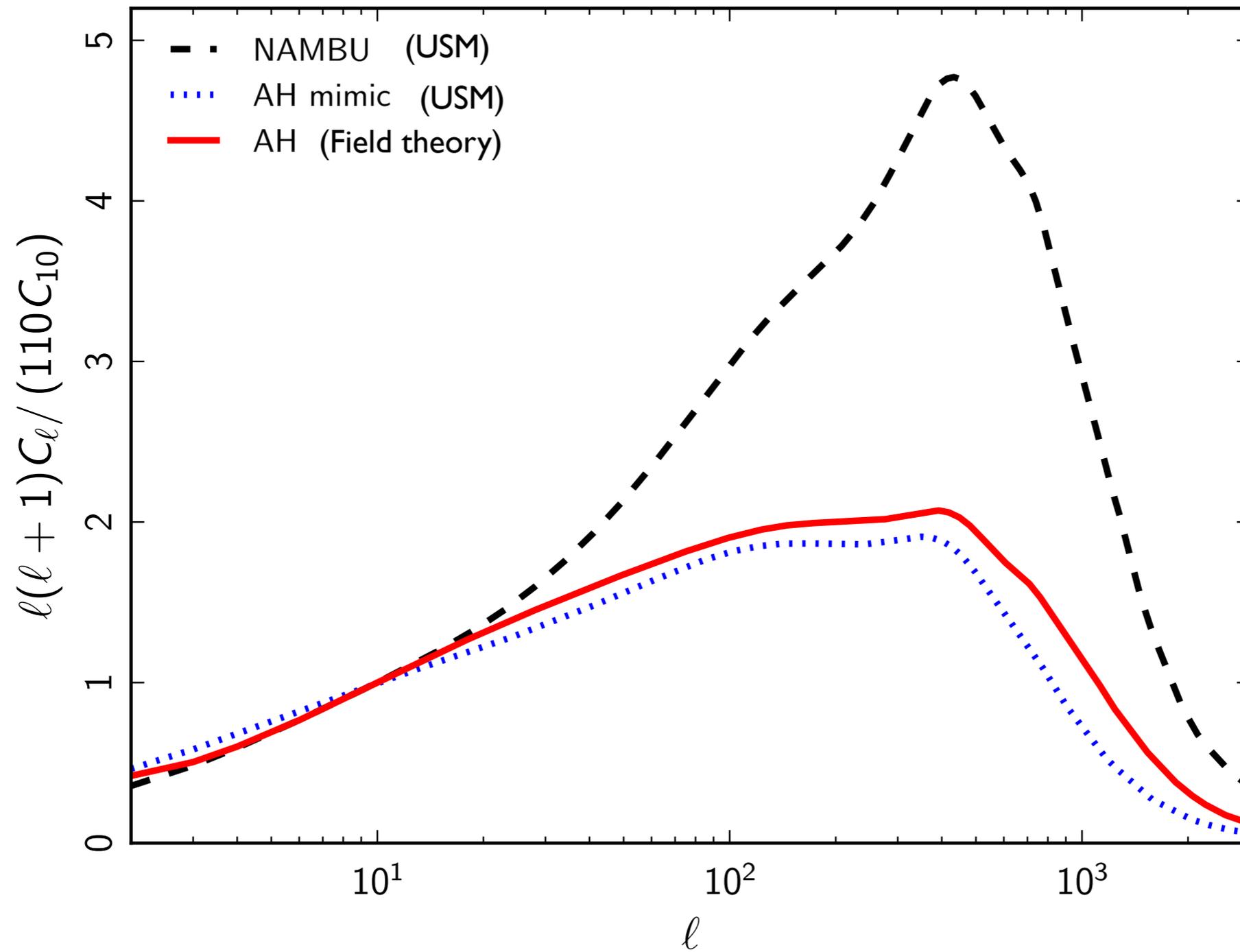
- ▶ Mimic model computed with USM, choosing parameters to more closely resemble Abelian-Higgs
- ▶ Turn off evolution of network parameters between matter and radiation eras
- ▶ Choose $\xi = 0.35$ $v = 0.4$ $\alpha = 1.05$
- ▶ Encouragingly find closer agreement between spectra
- ▶ Remember Nambu suffers less of an issue with dynamical range
- ▶ Abelian-Higgs has advantage of including radiation
- ▶ Even with simple physics USM does a decent job of fitting both!



Power Spectrum



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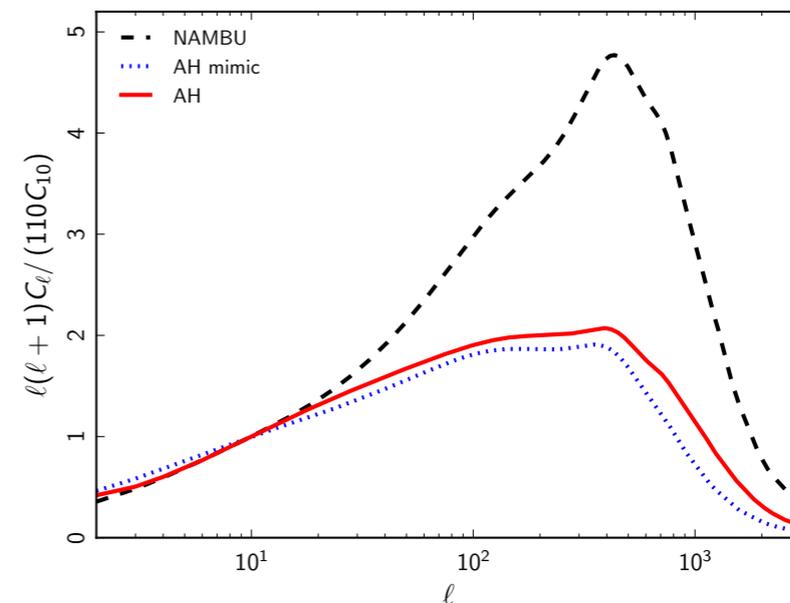


Power Spectrum



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- ▶ Note: Amplitude set by mass per unit length $C_\ell^{\text{string}} \propto (G\mu)^2$
- ▶ Two primary contributions to spectrum - density perturbations at last scattering, Kaiser-Stebbins effect along line of sight
- ▶ Peak position set by correlation length at last scattering
- ▶ Why the difference between Nambu and Abelian-Higgs?
- ▶ Normalisation determined by correlation length
- ▶ Nambu model has more small scale power - smaller string correlation length in radiation era, by factor of ~ 2
- ▶ Different split into scalar, vector and tensors - small scale structure of strings





- ▶ Let's now revisit the USM. In existing codes (CMBACT) the UETC is **not** computed
- ▶ Rather an ensemble of source histories are created, then averaged to find power spectra
- ▶ The EM tensor of a straight string segment is

$$\Theta_{00} = \frac{\mu\alpha}{\sqrt{1-v^2}} \frac{\sin(k\hat{X}_3\xi\tau/2)}{k\hat{X}_3/2} \cos(\chi + k\hat{X}_3v\tau), \quad \Theta_{ij} = \left[v^2\hat{X}_i\hat{X}_j - \frac{1-v^2}{\alpha}\hat{X}_i\hat{X}_j \right] \Theta_{00},$$

- ▶ \hat{X} and $\dot{\hat{X}}$ are randomly orientated unit vectors satisfying $\hat{X} \cdot \dot{\hat{X}} = 0$
- ▶ Phase χ set by location of string
- ▶ Orientated wave vector as $\mathbf{k} = k\hat{k}_3$, perform scalar, vector, tensor split

$$\Theta^S = (2\Theta_{33} - \Theta_{11} - \Theta_{22})/2, \quad \Theta^V = \Theta_{13}, \quad \Theta^T = \Theta_{12}.$$



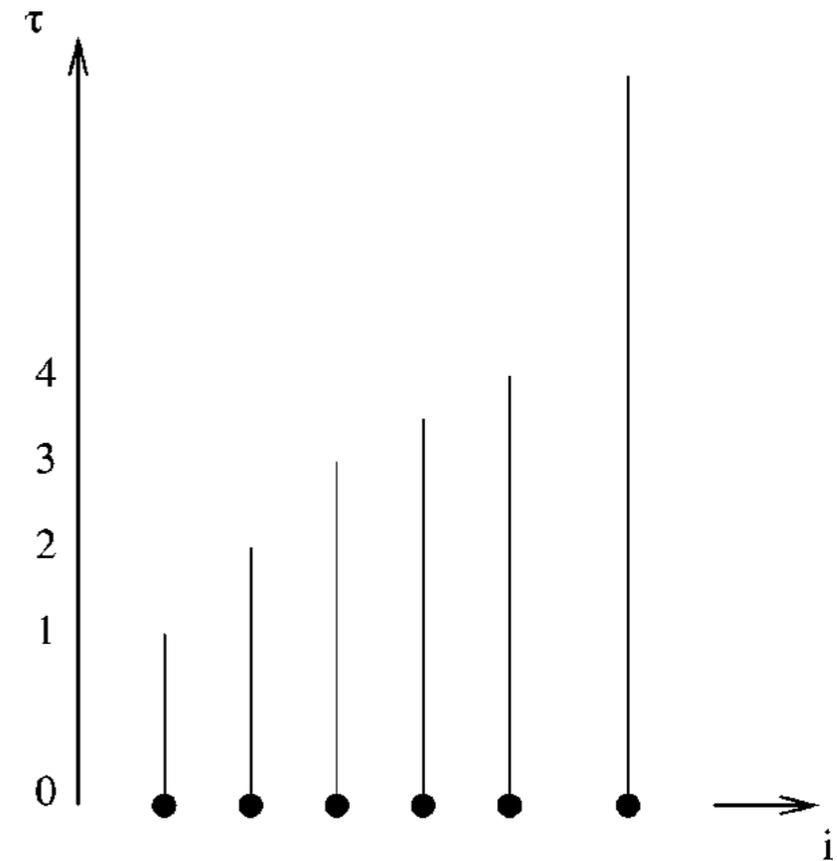
- ▶ During scaling number density of strings scales as $n(\tau) \propto \tau^{-3}$, requires tracking very large number of segments
- ▶ USM consolidates all segments that decay at some time into a **single** segment
- ▶ Number of segments that decay between τ_i and τ_{i-1} is

$$N_d(\tau_i) = V[n(\tau_{i-1}) - n(\tau_i)],$$

- ▶ EM tensor of network (K consolidated segments) is

$$\Theta_{\mu\nu} = \sum_{i=1}^K [N_d(\tau_i)]^{1/2} \Theta_{\mu\nu}^i T^{\text{off}}(\tau, \tau_i, L_f),$$

- ▶ Note consolidated segment has weight $\sqrt{N_d}$
- ▶ T^{off} is segment decay function



Analytic UETC



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- ▶ From these it is possible to work out the USM UETC **analytically**

(Avogoustidis et al 2012)

$$\langle \Theta(k, \tau_1) \Theta(k, \tau_2) \rangle = \frac{2f(\tau_1, \tau_2, \xi, L_f)}{16\pi^3} \int_0^{2\pi} d\phi \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\psi \int_0^{2\pi} d\chi \Theta(k, \tau_1) \Theta(k, \tau_2),$$

- ▶ The scaling function can be computed analytically from the segment decay

$$f(\tau_1, \tau_2, \xi, L_f \rightarrow 1) = \frac{1}{[\xi \text{Max}(\tau_1, \tau_2)]^3}.$$

- ▶ UETC involves several integrals which are doable, e.g.

$$\frac{1}{2} \int_0^\pi d\theta \sin^3\theta \cos(x \cos\theta) J_0(\rho \sin\theta) = \left[1 + \frac{\partial^2}{\partial x^2} \right] \left(\frac{\sin\sqrt{\rho^2 + x^2}}{\sqrt{\rho^2 + x^2}} \right),$$

- ▶ And only **two** which weren't 😞 , e.g.

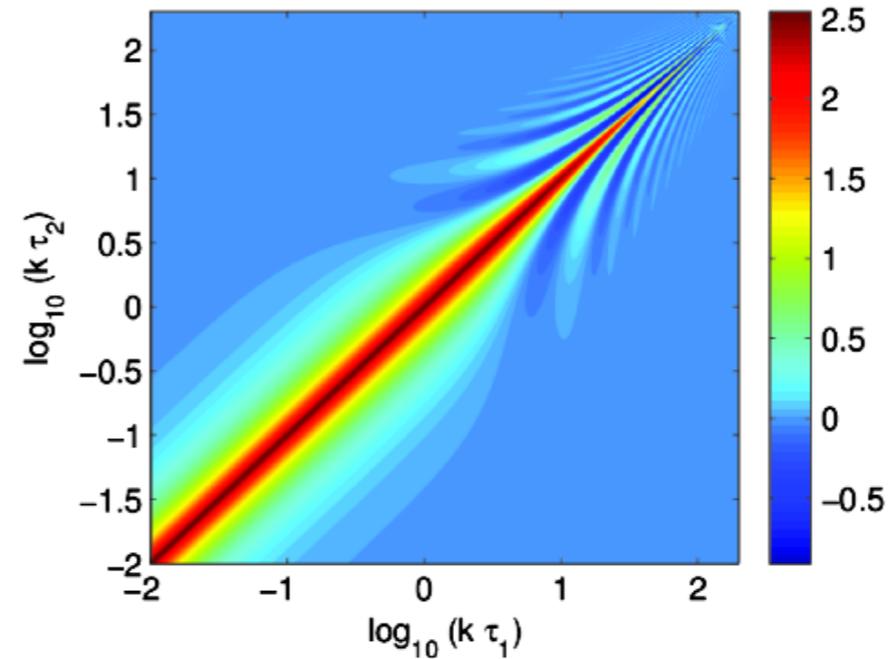
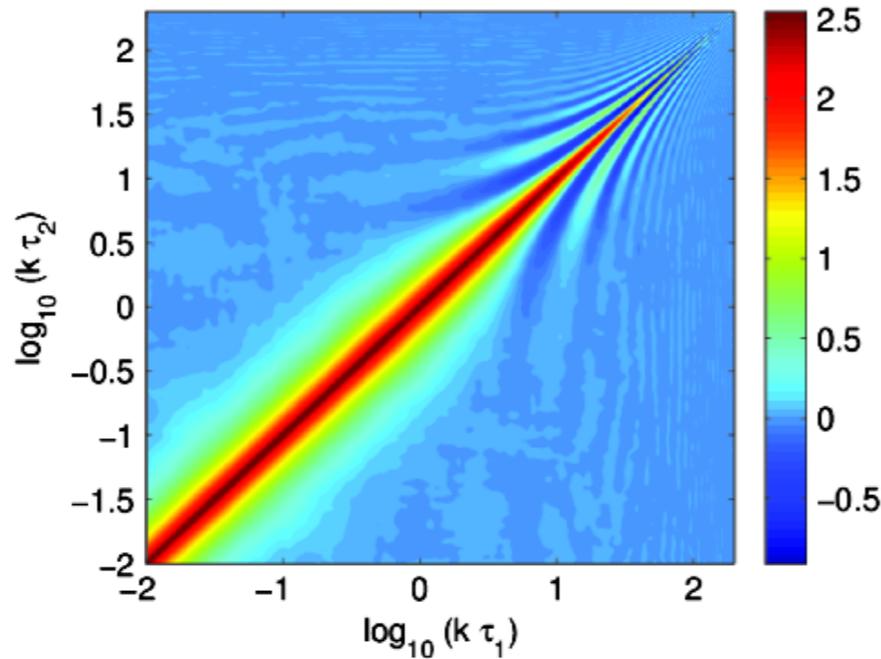
$$\frac{1}{2} \int_0^\pi d\theta \sin\theta \cos(x \cos\theta) J_0(\rho \sin\theta) \sec^2\theta = \sum_{c=0}^{\infty} \frac{1}{c!} \frac{\rho}{(2c-1)} \left(-\frac{x^2}{2\rho} \right)^c j_{c-1}(\rho),$$

Comparison



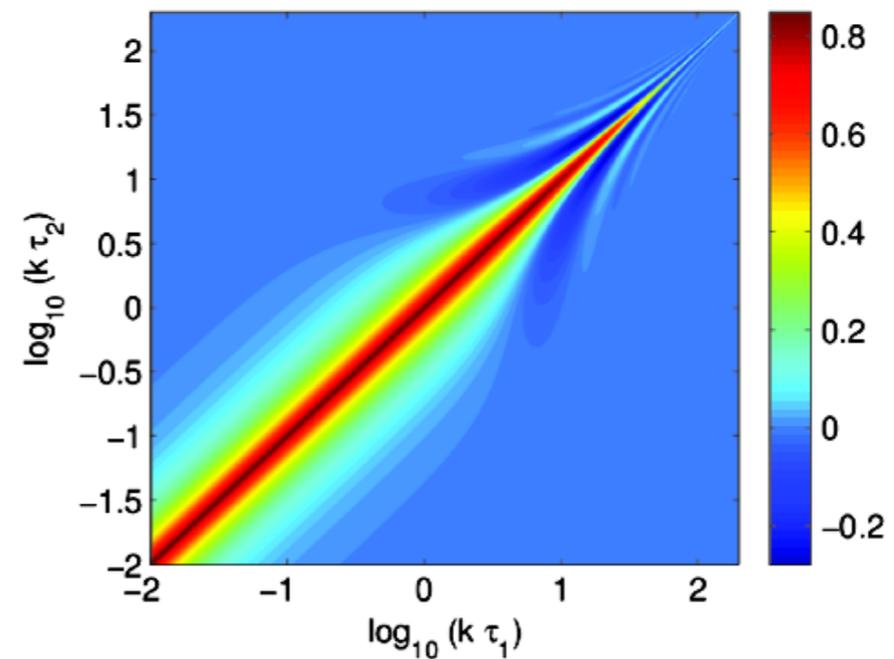
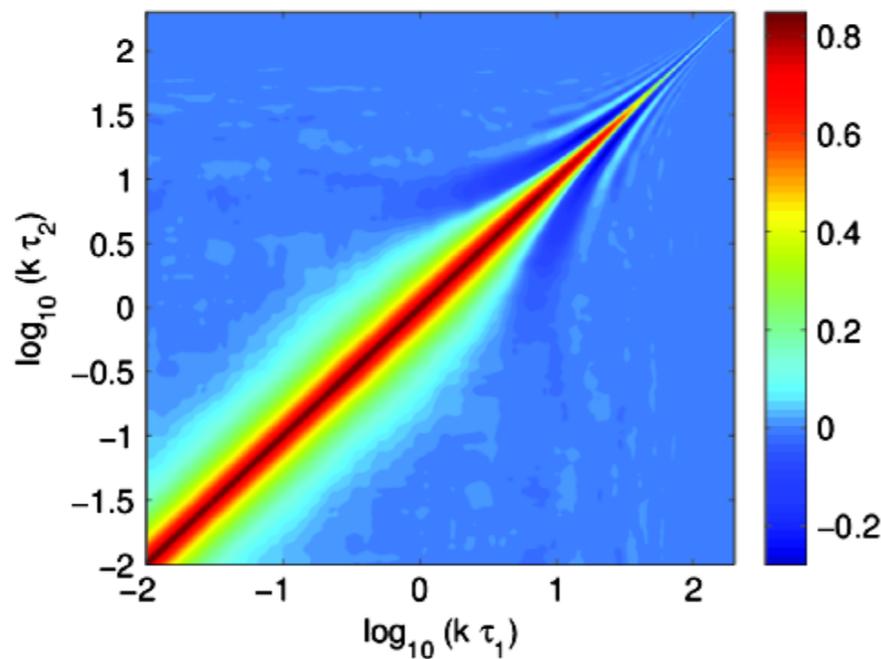
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Scalar
anisotropic
stress,
numerical
simulations



Scalar
anisotropic
stress,
analytic

Vector
anisotropic
stress,
numerical
simulations



Vector
anisotropic
stress,
analytic

(Hours)

(Seconds)

Eigenmodes



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- ▶ UETC can be decomposed into eigenmodes - each of these are **coherent**, can be used as source functions in CMB code

(Pen et al 1997)

$$(k^2 \tau_1 \tau_2)^\gamma (\tau_1 \tau_2)^{1/2} \langle \Theta(k, \tau_1) \Theta(k, \tau_2) \rangle = \sum_{i=1}^n \lambda_i u_i(k\tau_1) \otimes u_i(k\tau_2),$$

- ▶ Diagonalization introduces a change of basis, but since modes are orthogonal and perturbations are linear, CMB sources are

$$\Theta(k\tau) \rightarrow \frac{u(k\tau)}{(k\tau)^\gamma \tau^{1/2}}.$$

- ▶ Power spectra found by summing over eigenmodes (ordered from highest to lowest) and truncated at some number

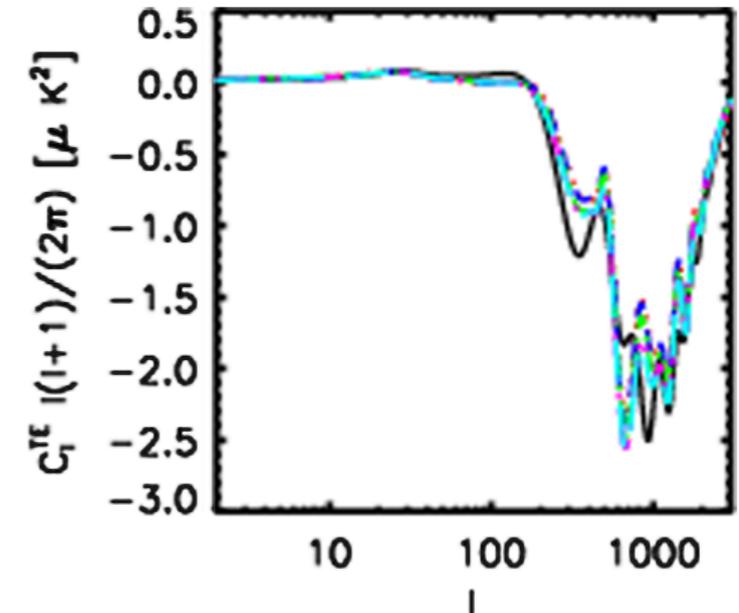
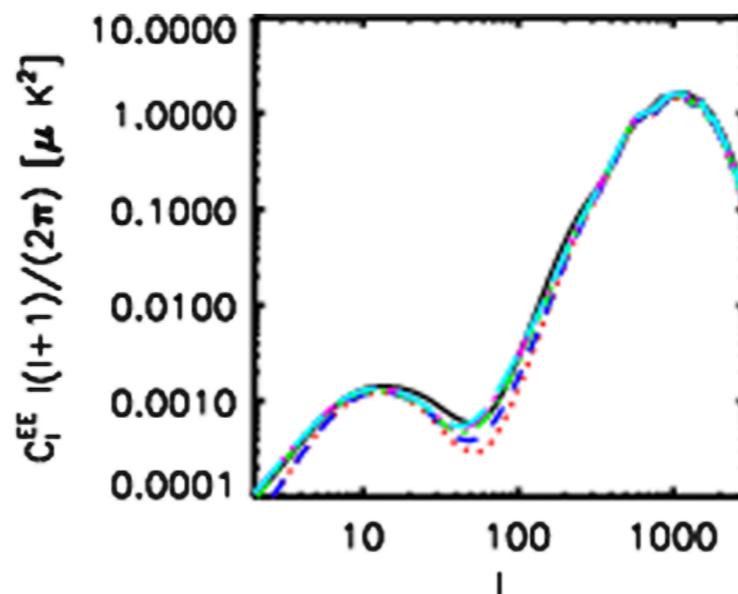
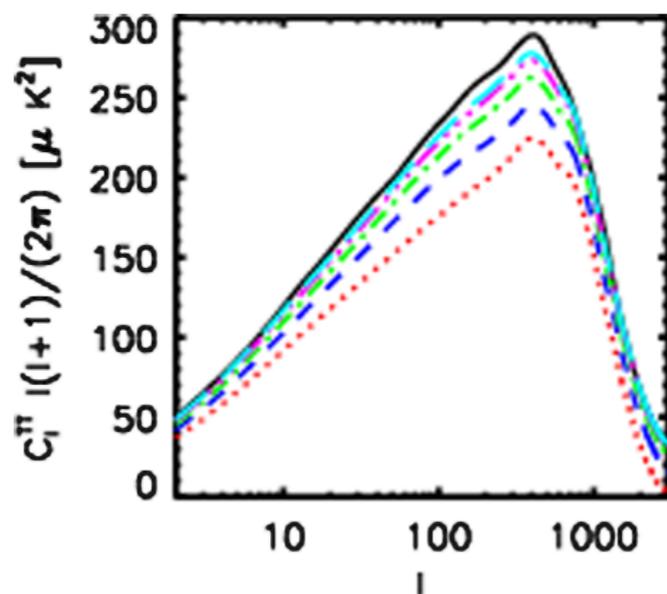
$$C_\ell = \sum_{i=1}^n \lambda_i C_\ell^i,$$

CMB Spectra



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- ▶ CMB spectra are $C_{\ell}^{i(I)} = \frac{2}{\pi} \int k^2 dk \Delta_{\ell}^{i(I)}(k, \tau_0) \Delta_{\ell}^{i(I)}(k, \tau_0)$,
- ▶ Incorporated into CAMB - CAMBACT (Avogoustidis et al 2012)
- ▶ E.g. scalar spectra



$$C_{\ell}^{\text{string}} \propto (G\mu)^2$$

Black - CMBACT, 2000 realizations
 Red - CMBACT, 16 eigenmodes
 Green - CMBACT, 64 eigenmodes
 Magenta - CMBACT, 128 eigenmodes

Some Numbers



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- ▶ Running CMBACT for 2000 source realisations takes ~ 20 hours
- ▶ Obviously this makes it impossible to run any sort of MCMC with Planck (other than fitting for the overall amplitude of a given spectrum)
- ▶ With CAMBACT we need somewhere between 50-100 eigenmodes for reasonable accuracy
- ▶ Each mode requires running CAMB for scalars, vectors and tensors (i.e. 150-300 CAMB evaluations)
- ▶ CAMB takes ~ 1 second to run, hence total CAMBACT computational time is several minutes
- ▶ A big improvement, but currently still rather prohibitive for Planck MCMC

Abelian Higgs



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- ▶ So far focused on USM case. Idea is to use simulations (both Nambu and Abelian-Higgs) to inform network parameters
- ▶ State of the art Abelian-Higgs simulations have also been used for the Planck analysis (Bevis et al)
- ▶ Fields evolved on 1024^3 grid, starting from random initial conditions designed to mimic a phase transition
- ▶ Brief diffusive period ensures system rapidly reaches scaling
- ▶ String cores are partially fattened to enlarge dynamical range
- ▶ Various runs performed to check results insensitive to string fattening parameter
- ▶ UETC's calculated at regular intervals and used in CMBEASY code

Planck Constraints



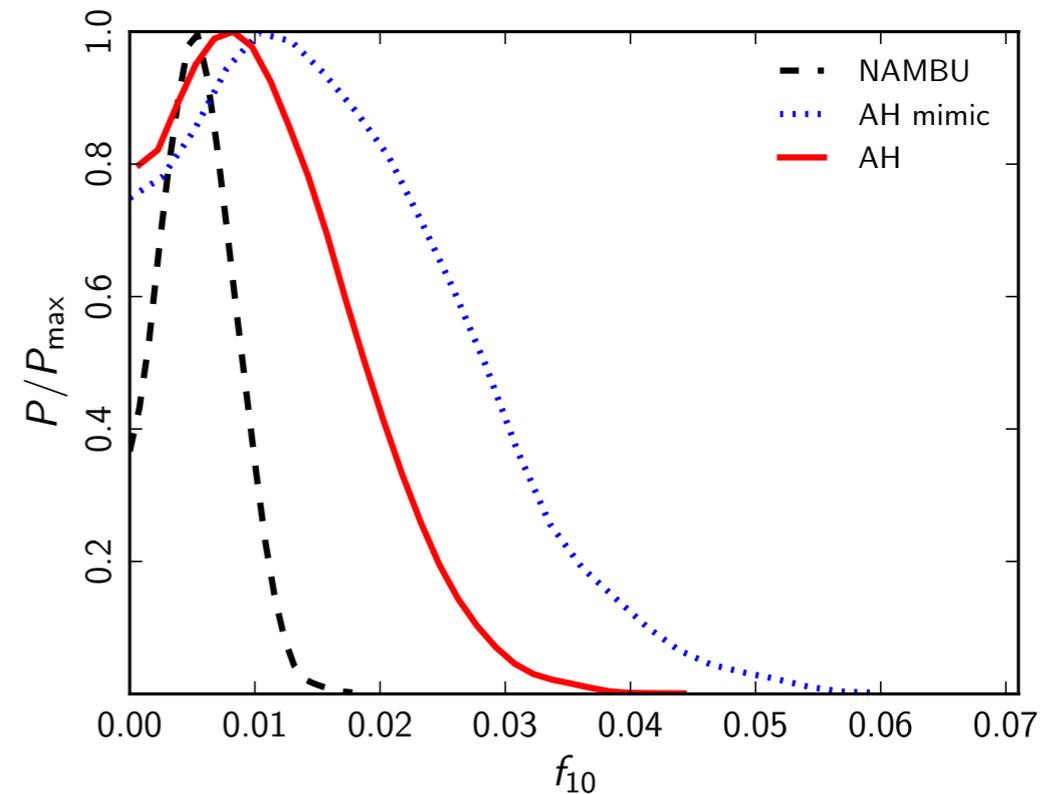
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- ▶ Additional parameter added to LCDM, fraction of strings at $\ell = 10$

$$f_{10} = \frac{C_l^{\text{string}}}{C_l^{\text{string}} + C_l^{\text{inflation}}}$$

- ▶ Use same dataset/priors as in main cosmology paper

Defect type	Planck+WP		Planck+WP+highL	
	f_{10}	$G\mu/c^2$	f_{10}	$G\mu/c^2$
NAMBU	0.015	1.5×10^{-7}	0.010	1.3×10^{-7}
AH-mimic	0.033	3.6×10^{-7}	0.034	3.7×10^{-7}
AH	0.028	3.2×10^{-7}	0.024	3.0×10^{-7}
SL	0.043	11.0×10^{-7}	0.041	10.7×10^{-7}
TX	0.055	10.6×10^{-7}	0.054	10.5×10^{-7}



Planck + WP + highL

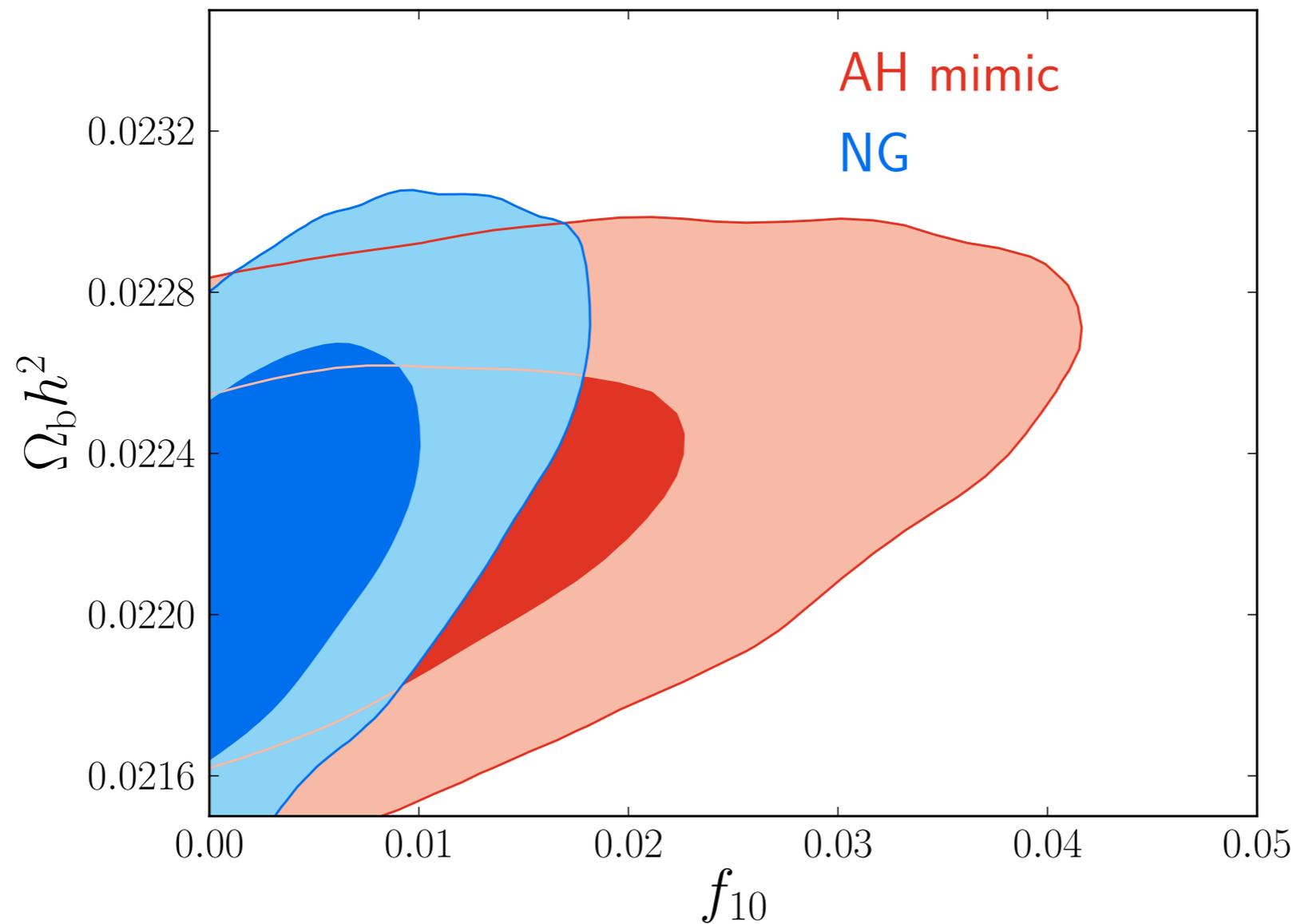
- ▶ Limits on string tension significantly improved, e.g. in Nambu model string fraction $< 1\%$

Planck Constraints



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- ▶ Strings exhibit no significance correlations with any other cosmological parameter (as for WMAP)

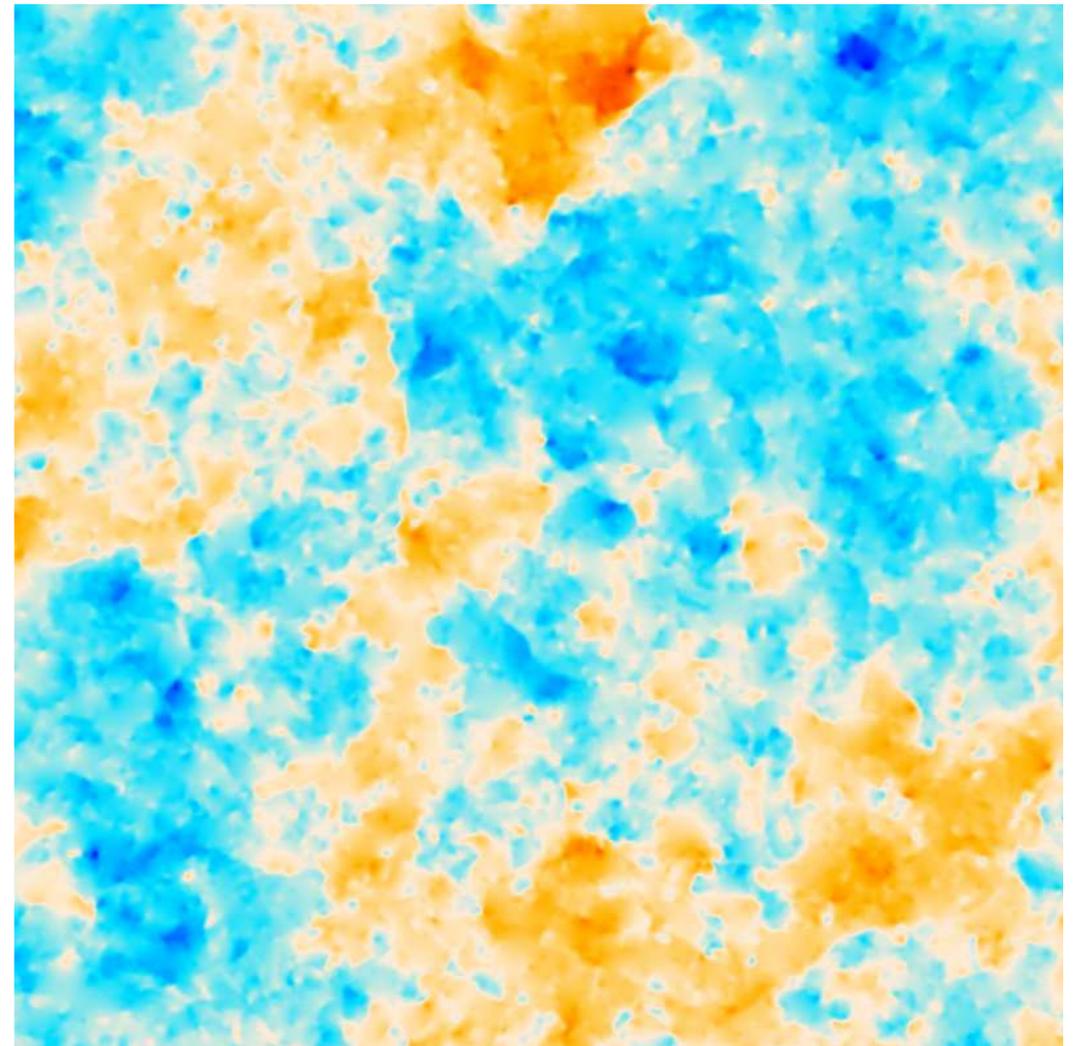


Non Gaussianity



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- ▶ To go beyond 2-point function Planck has used Nambu simulations to compute post-recombination string maps (Ringeval and Bouchet 2012)
- ▶ Maps expected to be accurate and large and small scales, will underestimate power at intermediate scales
- ▶ Maps produced at 0.85' and 1.7' resolution
- ▶ For lower resolution maps aliasing effects carefully taken into account
- ▶ Map has Gaussian component, but non Gaussian step like features clearly apparent



NG Searches



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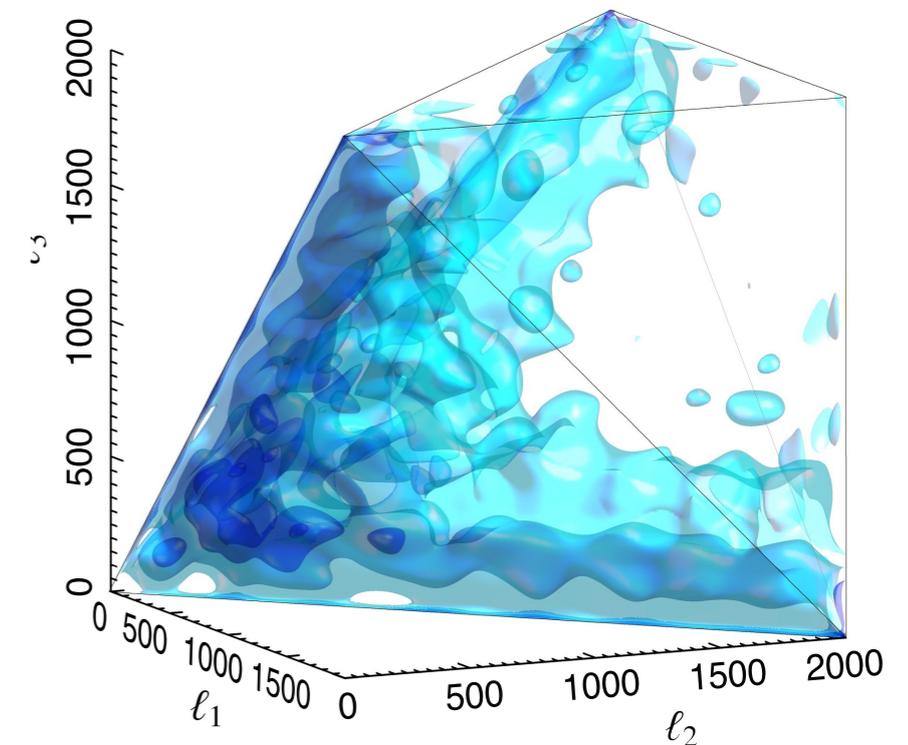
- ▶ Use 4 different Planck foreground cleaned maps
- ▶ Check results are consistent with each map, insensitive to choice of mask

(Fergusson and Shellard)

- ▶ Bispectrum computed using **modal** method. Expands reduced bispectrum $b_{\ell_1 \ell_2 \ell_3}^{\text{string}}$ in separable modes

$$\frac{b_{\ell_1 \ell_2 \ell_3}^{\text{string}}}{\sqrt{C_{\ell_1} C_{\ell_2} C_{\ell_3}}} = \sum_n \beta_n^Q Q_n(\ell_1, \ell_2, \ell_3),$$

- ▶ Express bispectrum amplitude estimate in terms of modes $Q_n(\ell_1, \ell_2, \ell_3)$
- ▶ Modal reconstruction performed on post-recombination string maps
- ▶ Correlation with primordial shapes is since it does not contain an oscillatory feature



NG Result



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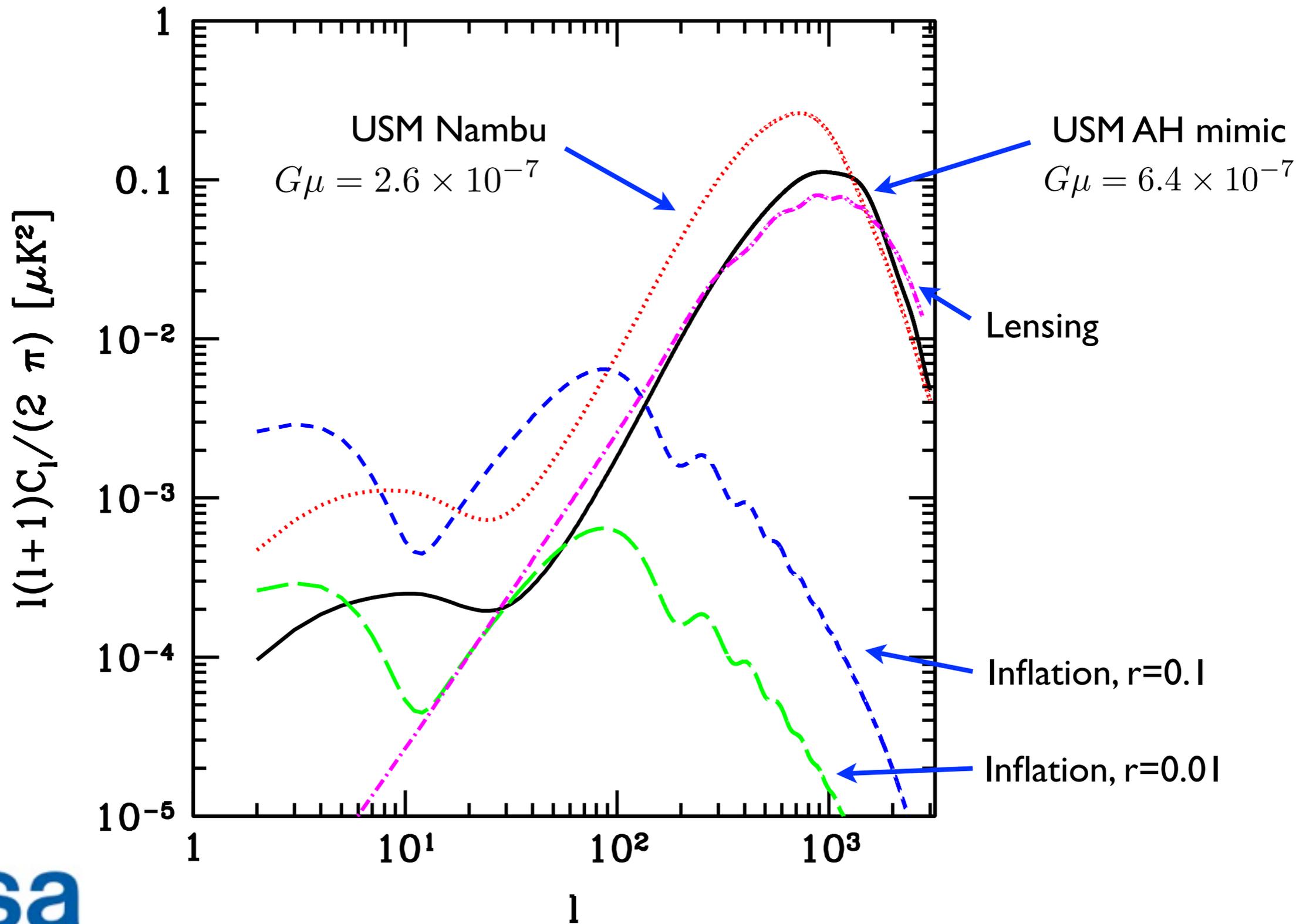
Method	Bispectrum signal type	Independent analysis f_{NL}	ISW subtract Joint f_{NL}	ISW/FG res. Joint f_{NL}
SMICA	Lensing ISW	0.75 ± 0.37	–	–
	Diff. PS $\times 10^{28}$	1.05 ± 0.32	1.35 ± 0.34	1.40 ± 0.34
	Cosmic strings	0.19 ± 0.20	0.50 ± 0.21	0.37 ± 0.21
	$G\mu/c^2$ (95%)	8.4×10^{-7}	9.7×10^{-7}	9.3×10^{-7}
NILC	Lensing ISW	0.91 ± 0.36	–	–
	Diff. PS $\times 10^{28}$	1.16 ± 0.32	1.44 ± 0.34	1.44 ± 0.34
	Cosmic strings	0.13 ± 0.20	0.46 ± 0.21	0.23 ± 0.21
	$G\mu/c^2$ (95%)	8.1×10^{-7}	9.6×10^{-7}	8.7×10^{-7}
SEVEM	Lensing ISW	0.6 ± 0.36	–	–
	Diff. PS $\times 10^{28}$	1.07 ± 0.35	1.33 ± 0.38	–
	Cosmic strings	0.10 ± 0.20	0.38 ± 0.21	–
	$G\mu/c^2$ (95%)	7.9×10^{-7}	9.3×10^{-7}	–

- ▶ No evidence (currently) for NG string signal assuming post-recombination string maps

B Modes



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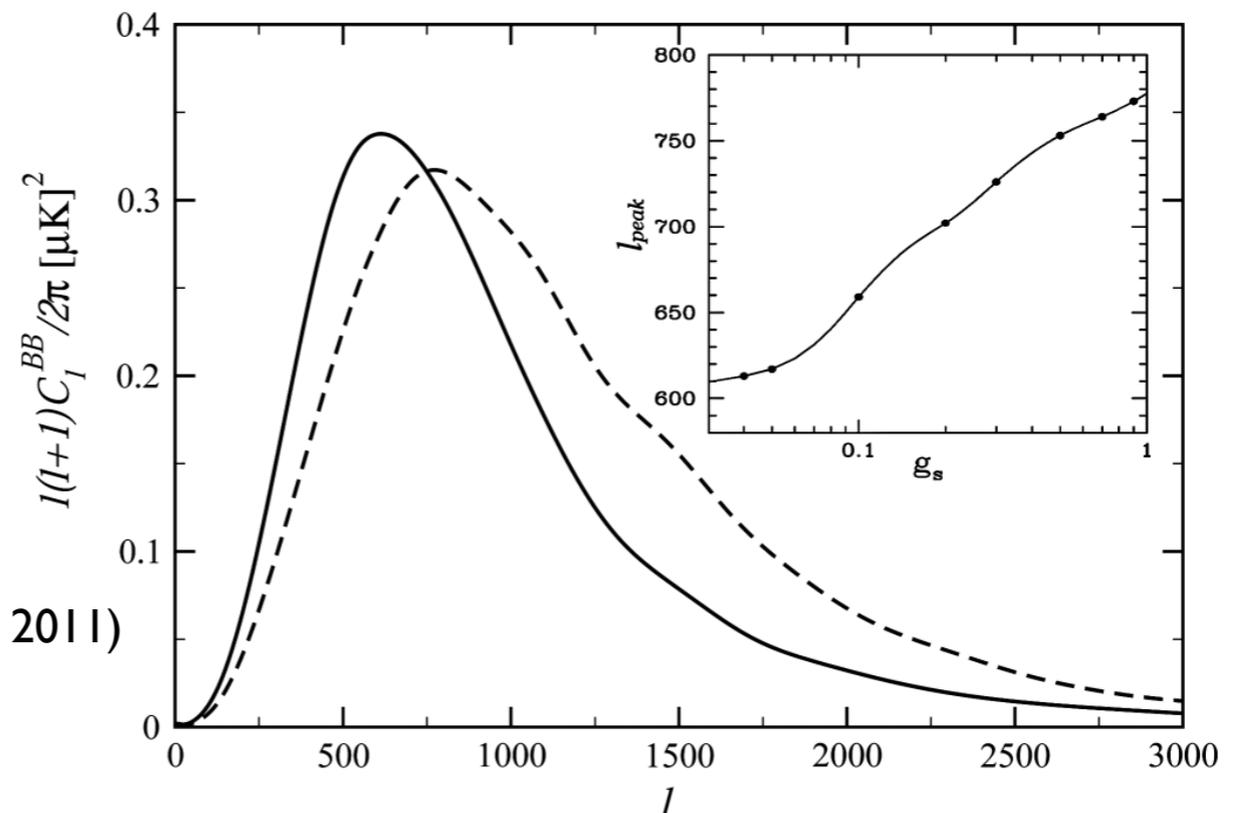


Superstrings



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- ▶ Strings have a fascinating connection with high energy physics
- ▶ It is possible more complicated networks may arise from superstrings
- ▶ Observational properties depend on fundamental string coupling g_s
- ▶ Generally now a spectrum of string tensions determined by g_s and charges which they carry
- ▶ Intercommutation properties modified
- ▶ Junctions can form between strings of different charge
- ▶ Modified VOS equations have been developed
- ▶ Find potentially observable shift in B mode spectrum



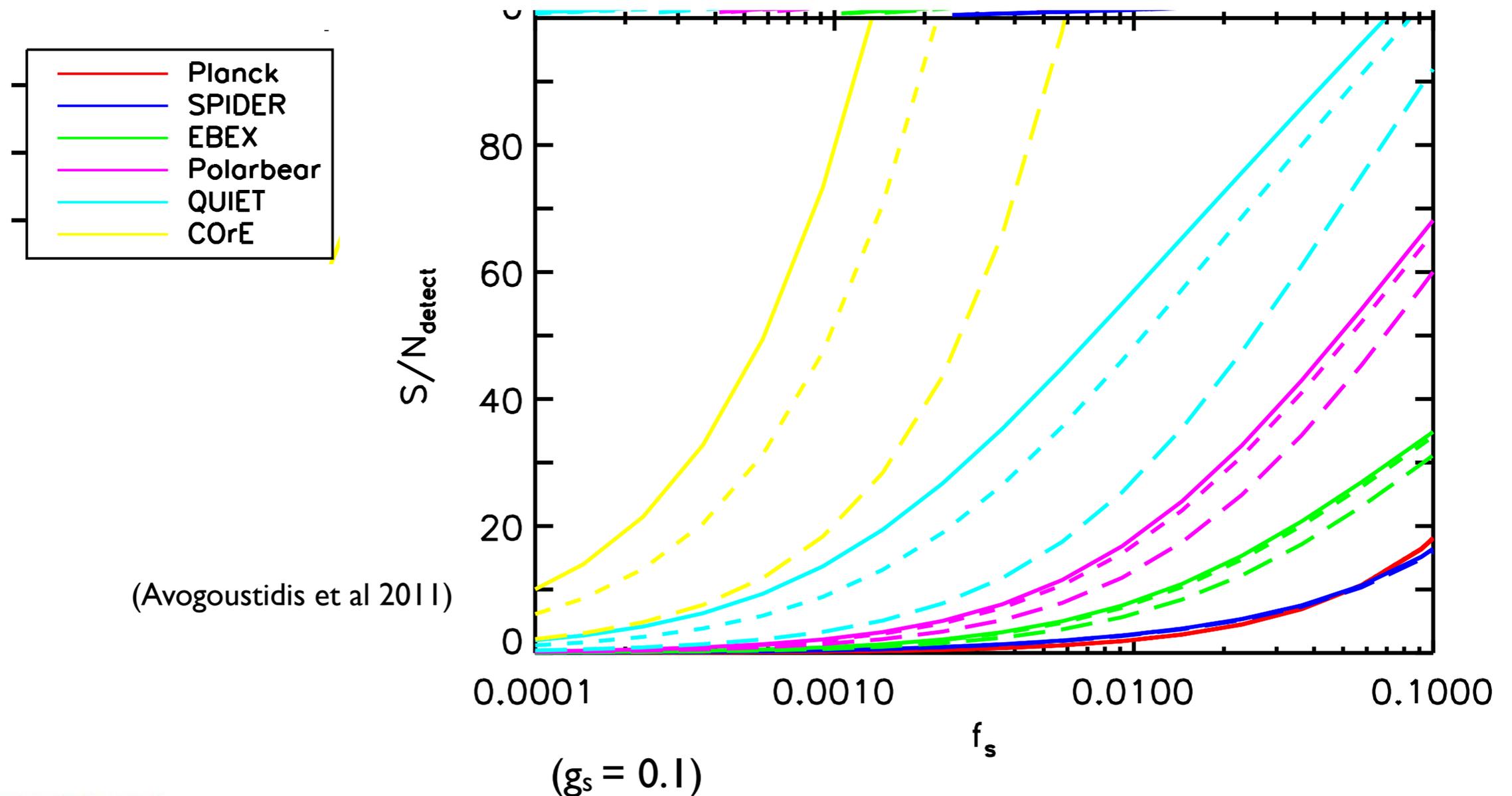
(Avogoustidis et al 2011)

Superstrings



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- ▶ B mode experiments have the potential to measure a fundamental property of string theory



Conclusions



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- ▶ Planck has provided stringent limits on strings
- ▶ Power spectrum gives

$$G\mu/c^2 < 1.5 \times 10^{-7}, \quad f_{10} < 0.015, \quad (\text{Nambu})$$

$$G\mu_{\text{AH}}/c^2 < 3.2 \times 10^{-7}, \quad f_{10} < 0.028. \quad (\text{Abelian-Higgs})$$

- ▶ Non Gaussian constraints give $G\mu/c^2 < 7.8 \times 10^{-7}$ (95 % CL).
- ▶ String evolution is crucial for understanding CMB anisotropies
- ▶ Still some uncertainty in modelling approaches - main problem is dynamical range
- ▶ Possible that improved USM model could enable MCMC search of string parameter space
- ▶ (Optimistic) hope is strings could link to high energy theory