

Study of jet energy loss and the LHC ALICE data

Péter Lévai

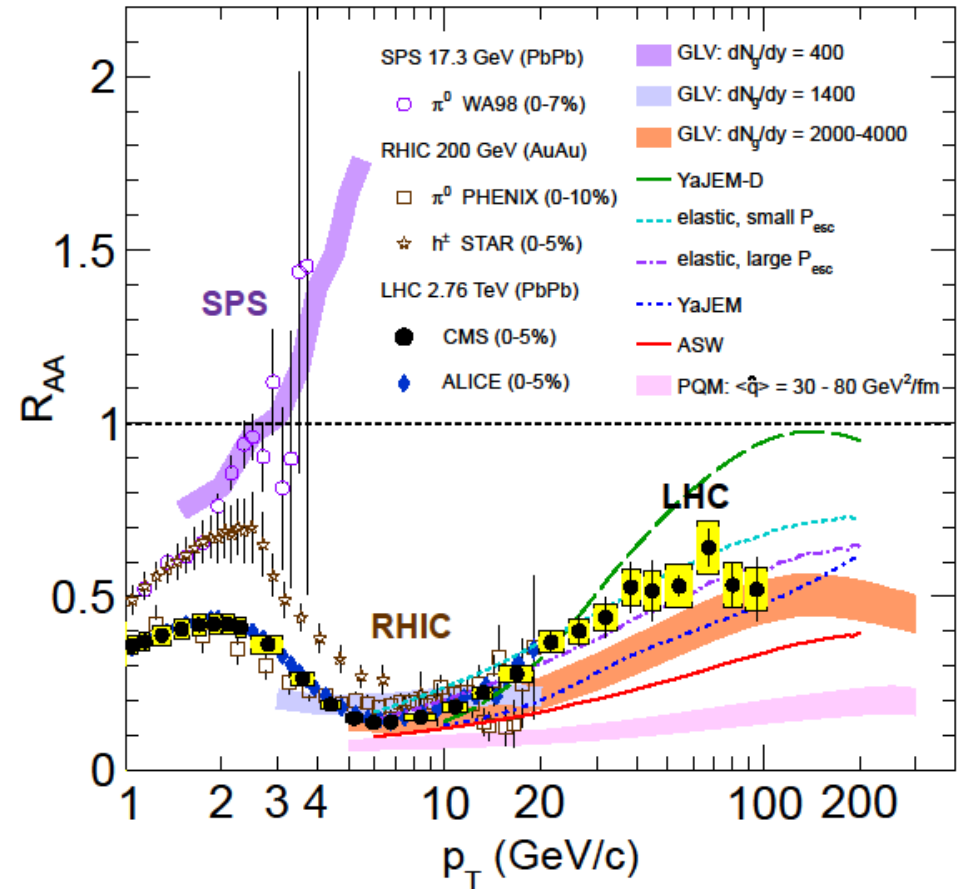
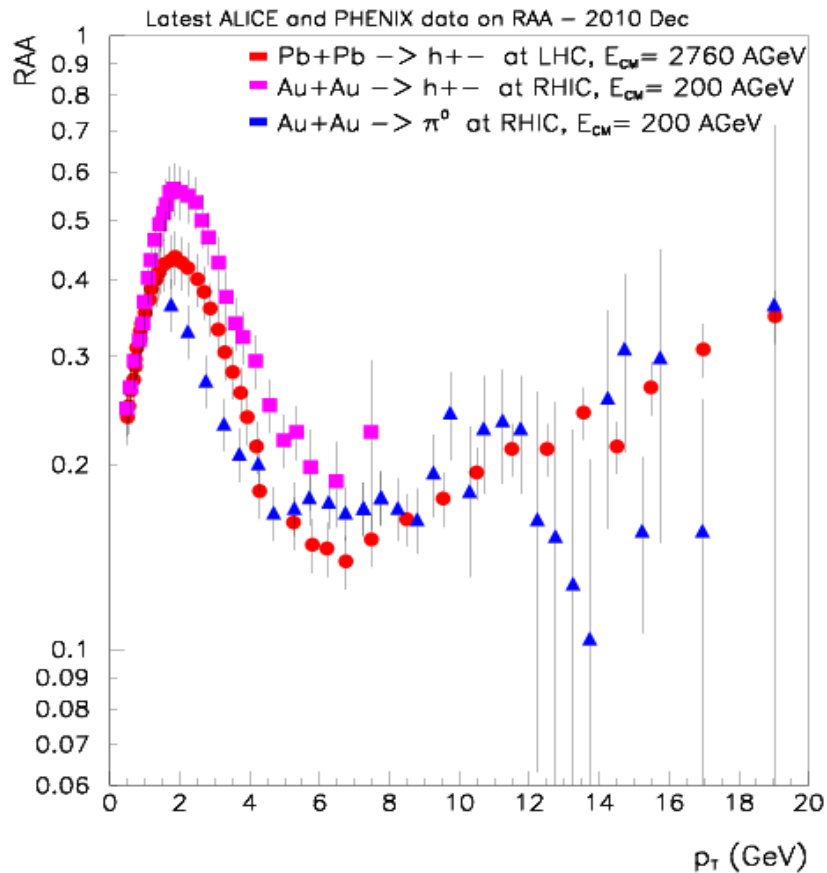
WIGNER RCP, Budapest, Hungary

12th Zimanyi Winter School

2 December 2012, Budapest

ICPAQGP'2010: New result on RAA

in Pb+Pb at 2760 AGeV ALICE at LHC --- J. Schukraft



First glance:

RAA overlaps at RHIC and LHC !!!

Second glance in 2012: NOT!

(we have good microscope)

Contents

1. Introduction

--- particle production, pQCD, jets, pp, pA

2. Jet energy loss

--- mechanism, description

3. Data vs. Theory at RHIC and LHC energies

--- answers and new questions

4. A new channel for extra proton production

--- coherent gluon field and diquark coalescence

5. Conclusion

1. Introduction

--- particle production, pQCD, jets, pp, pA

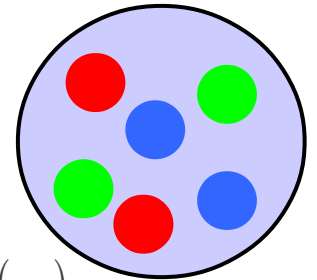
Particle production mechanisms in high energy HI collisions:

I. Dilute parton gas limit as initial condition + parton cascade:

PDF(p,n) + pQCD + Glauber + [Shad; Multisc; Quench; Fluct; ...]

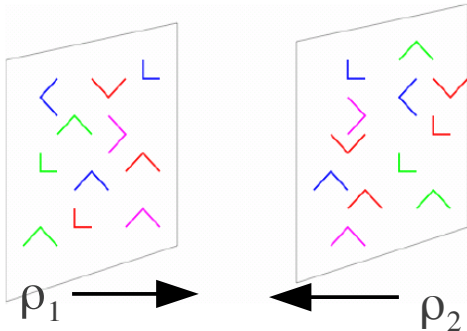
$$E_\pi \frac{d\sigma^{pp}}{d^3 p_\pi} = \int dx_1 \int dx_2 \int dz_c f_{a/p}(x_a, Q^2) f_{b/p}(x_b, Q^2) \frac{d\sigma}{d\hat{t}} \frac{D_c^\pi(z_c)}{\pi z_c^2}$$

$$E_\pi \frac{d\sigma^{AB}}{d^3 p_\pi} = \int d^2 b d^2 r t_A(\vec{r}) t_B(|\vec{b} - \vec{r}|) E_\pi \frac{d\sigma^{pp}}{d^3 p_\pi} \otimes S(\dots) \otimes M(\dots) \otimes Q(\dots) \otimes F(\dots)$$



Dilute gas

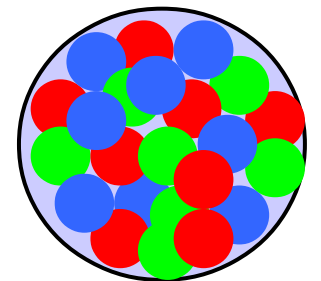
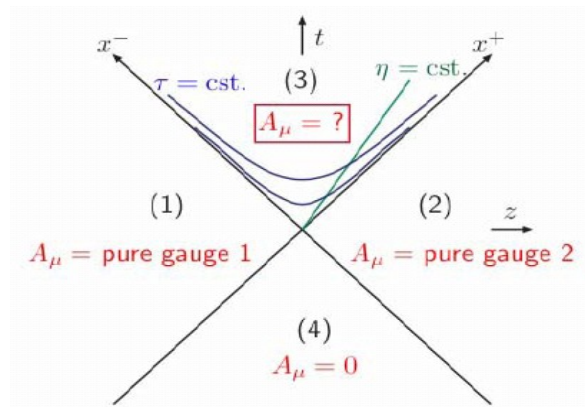
II. Dense gluon matter limit as initial condition + hydro:



CGC initial condition:

$$J^\mu = \delta^{\mu+} \delta(x^-) \rho_1(\mathbf{x}_T) + \delta^{\mu-} \delta(x^+) \rho_2(\mathbf{x}_T)$$

where $-D_i \alpha_{(m)}^i = \rho_{(m)}(\mathbf{x}_\perp)$. and α_1, α_2 gluon fields of nuclei

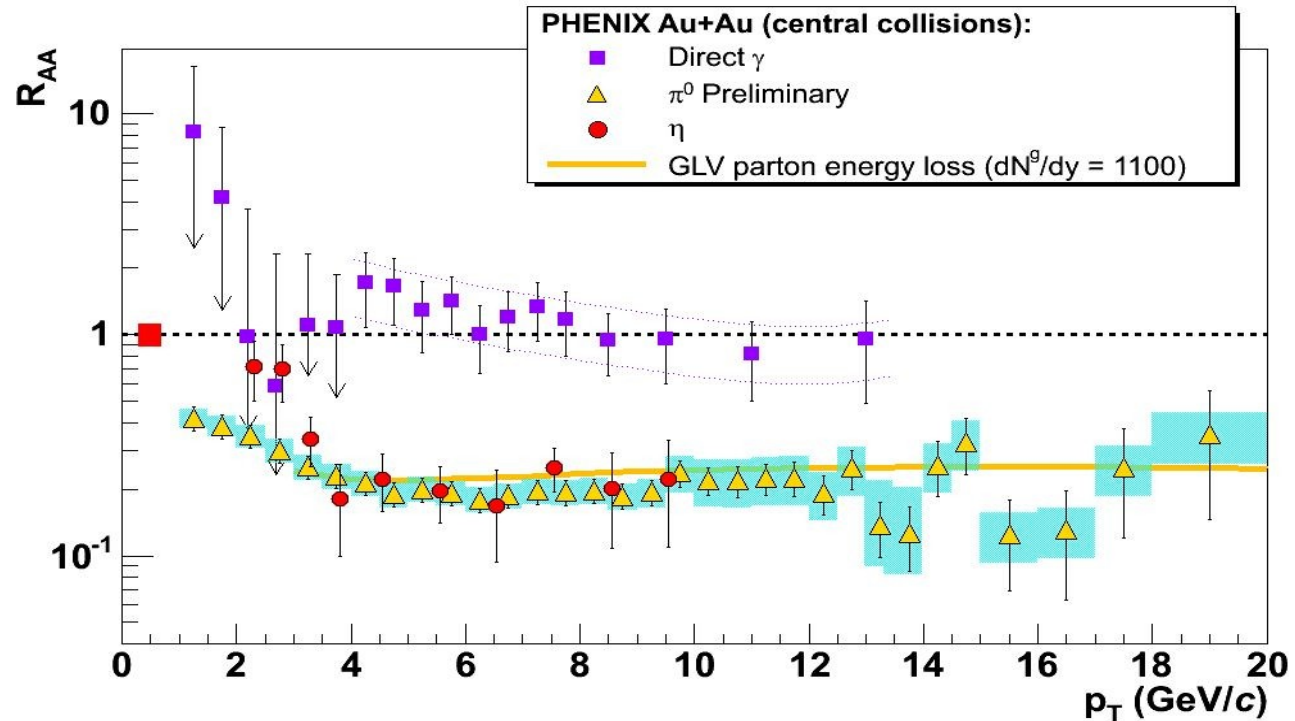


CGC: high density gluons

Successful applications of I and II:

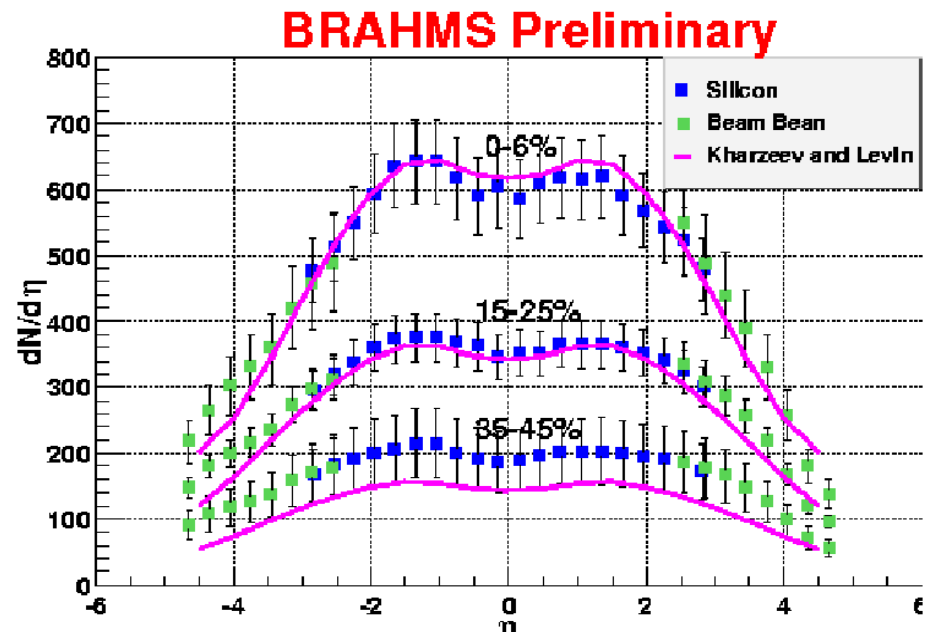
I. pQCD model:

- hard probes
- high- p_T physics
- jets
- h-h correlations
- ...

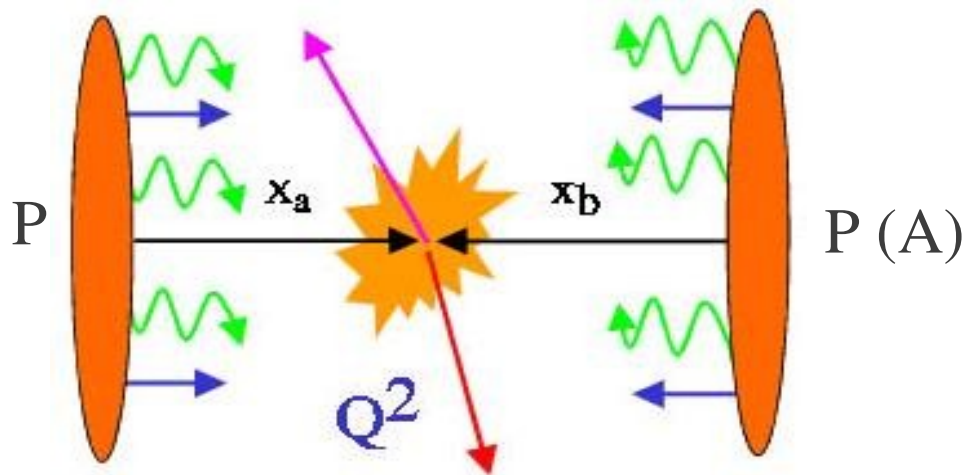


II. CGC model:

- soft physics
- multiplicities
- centrality dependence
- E_T production
- rapidity distributions
- ...



DILUTE PARTON GAS LIMIT: Jets in pp and in pA collisions:



**Jet production in pp collision
("in vacuum"):
➔ pQCD description**

Jet production in pA collision ("in cold dilute matter")

➔ **modified PQCD description:**

--- **SHADOWING** inside A

--- **MULTISCATTERING/BROADENING** penetrating A and the "cold matter"

--- **ENERGY LOSS** penetrating the "cold dilute matter"

Can we separate these mechanisms?

Can we determine them separately during theoretical data analysis ?

We could learn a lot from high precision RHIC and LHC data on dAu & pPb!

Hard physics: pion production in pp collision at high- p_T

Perturbative QCD calculations in LO/NLO for $p+p \rightarrow \pi + X$ process with finite - k_T
NLO : M. Aversa et al. NPB327,105; P. Chiappetta et al. NPB412,3; P. Aurenche et al. NPB399,34; ...)
+ intrinsic kT: G. Papp, P. Levai, G.G. Barnaföldi, G. Fai, hep-ph/0212249, EPJC33(2004)609

$$E_\pi \frac{d\sigma^{pp}}{d^3 p_\pi} = \frac{1}{S} \sum_{abc} \int_{vW/z_c}^{1-(1-v)/z_c} \frac{dv}{v(1-v)} \int_{vW/vz_c}^1 \frac{dw}{w} \int^1 dz_c$$

$$\int d^2 k_{Ta} \int d^2 k_{Tb} f_{a/p}(x_a, k_{Ta}, Q^2) f_{b/p}(x_b, k_{Tb}, Q^2)$$

$$\left[\frac{d\sigma^{BORN}}{dv} \delta(1-w) + \frac{\alpha_s(Q_R)}{\pi} K_{ab,c}(s, v, w, Q, Q_R, Q_F) \right] \frac{D_c^\pi(z_c)}{\pi z_c^2}$$

An approximation for the unintegrated parton distribution functions (PDFs) :

$$f_{a/p}(x_a, k_{Ta}, Q^2) = f_{a/p}(x_a, Q^2) g(k_{Ta})$$

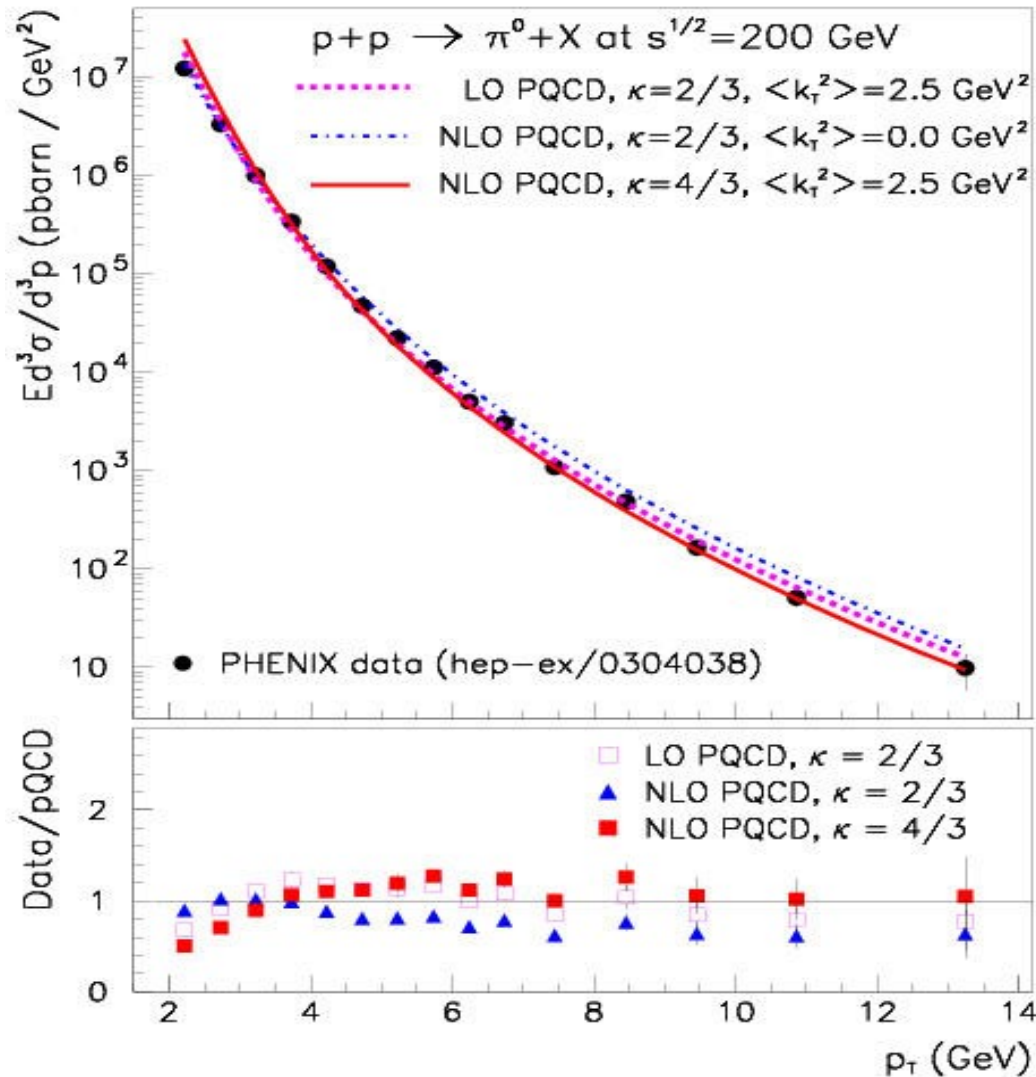
Where we use gaussian

$$g(k_{Ta}) = \frac{1}{\pi \langle k_T^2 \rangle} e^{-k_T^2 / \langle k_T^2 \rangle}$$

The width of the gaussian distribution for intrinsic-kT

Hard physics: pion production in pp collision at high- p_T

Perturbative QCD calculations in LO and NLO for pp --- including intrinsic- k_T



LO:

$$Q = \kappa p_T / z_c, \quad Q_F = \kappa p_T$$

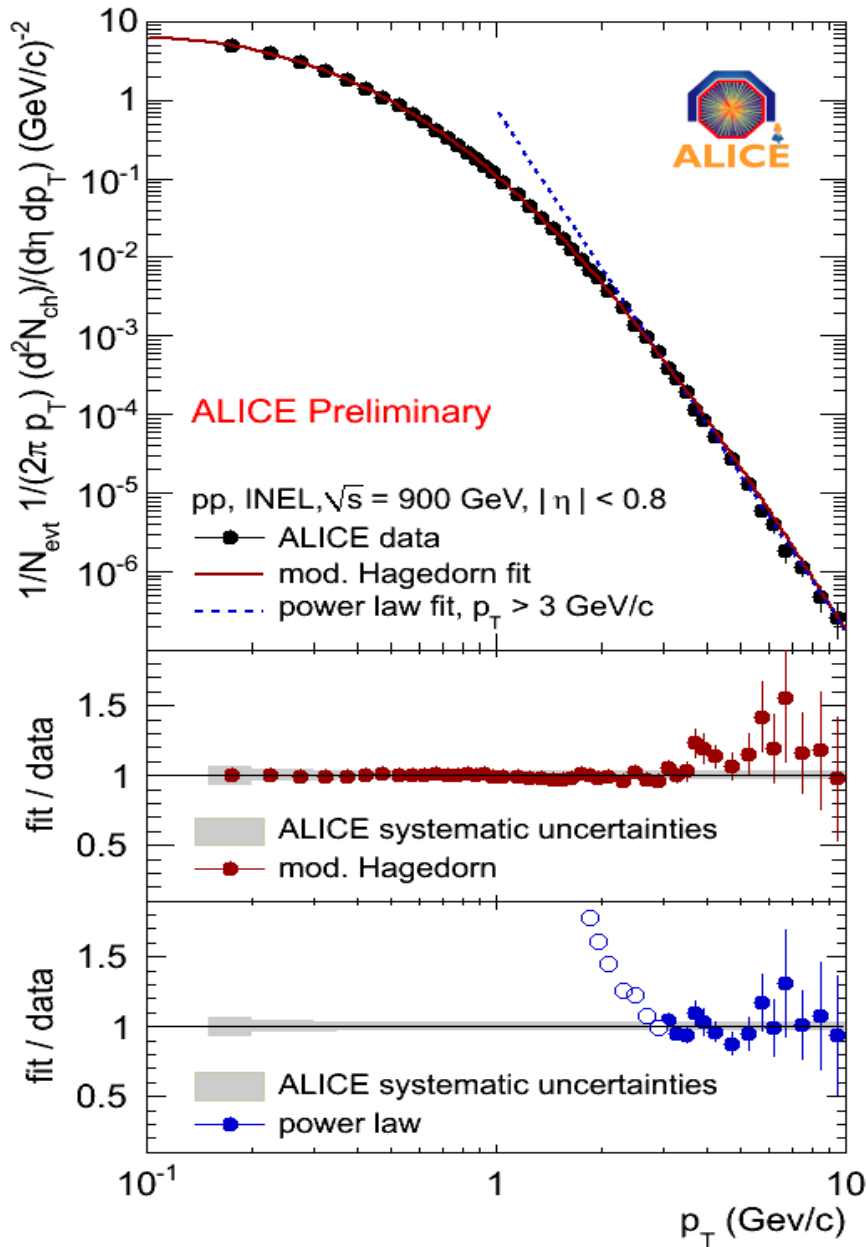
NLO:

$$Q = Q_R = \kappa p_T / z_c, \quad Q_F = \kappa p_T$$

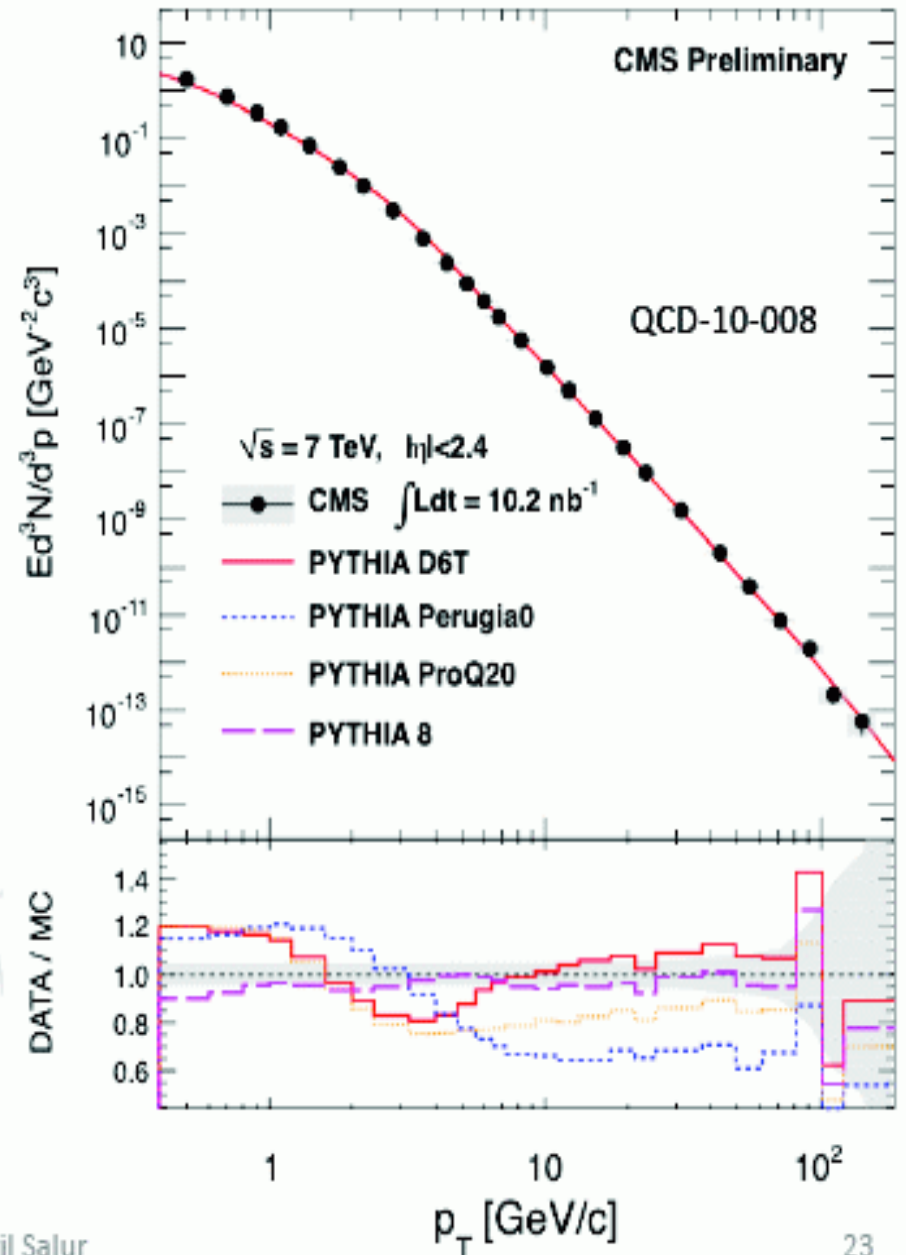
All descriptions with k_T are approx.
good enough for $3 \text{ GeV} < p_T < 15 \text{ GeV}$.

Y. Zhang, G. Fai, G. Papp,
G.G. Barnaföldi, P.L.:
PRC 65 (2002) 034903.
G.G. Barnaföldi et al.
EPJ C33 (2004) 609.

Charged hadron production in pp collisions in the high- p_T region :



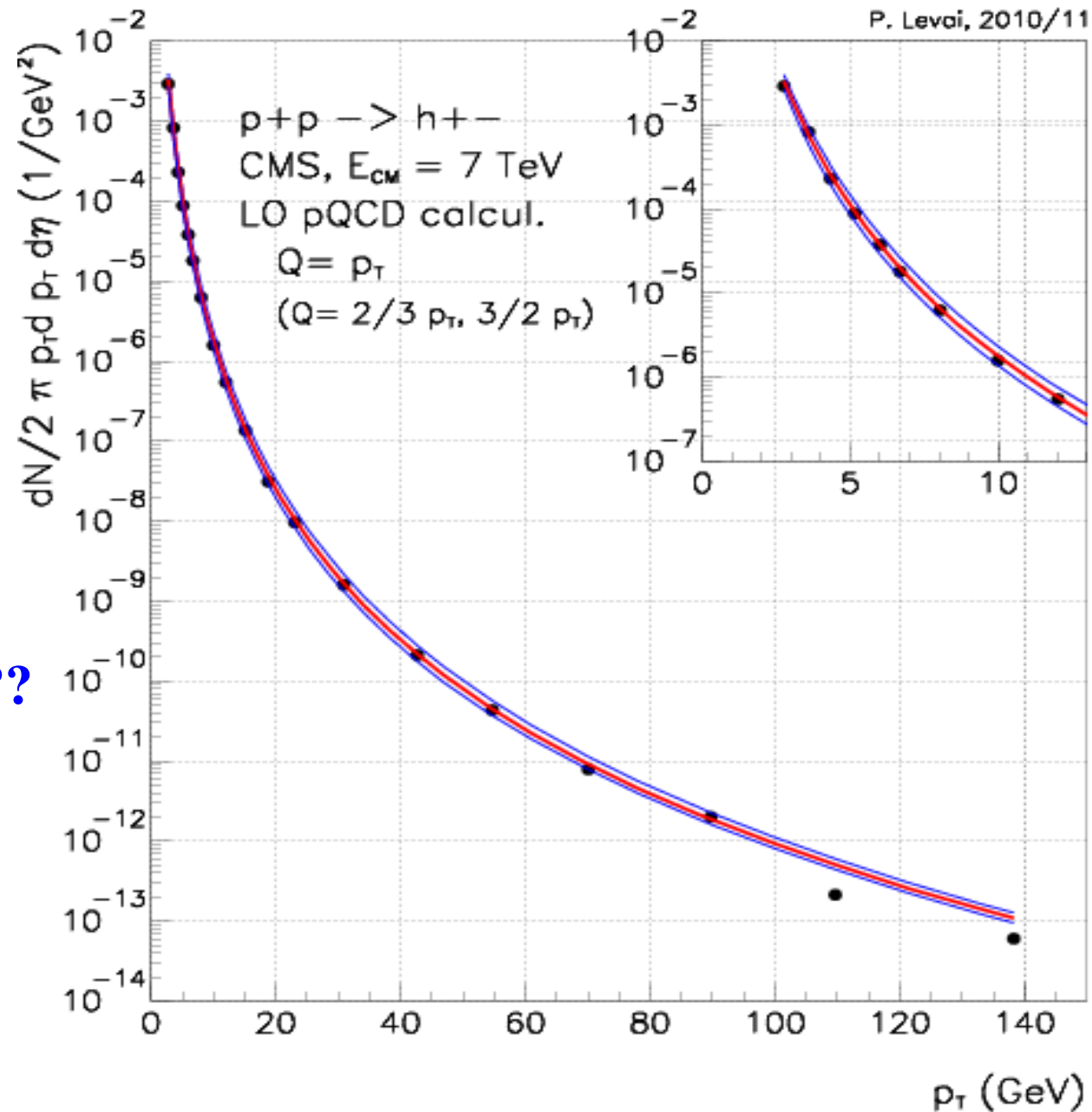
LHC ALICE (Prag'10)



LHC CMS (Prag'10)

Charged hadron production in pp collisions at 7 TeV – CMS data

LO pQCD



No need for
intrinsic-kT !!!?

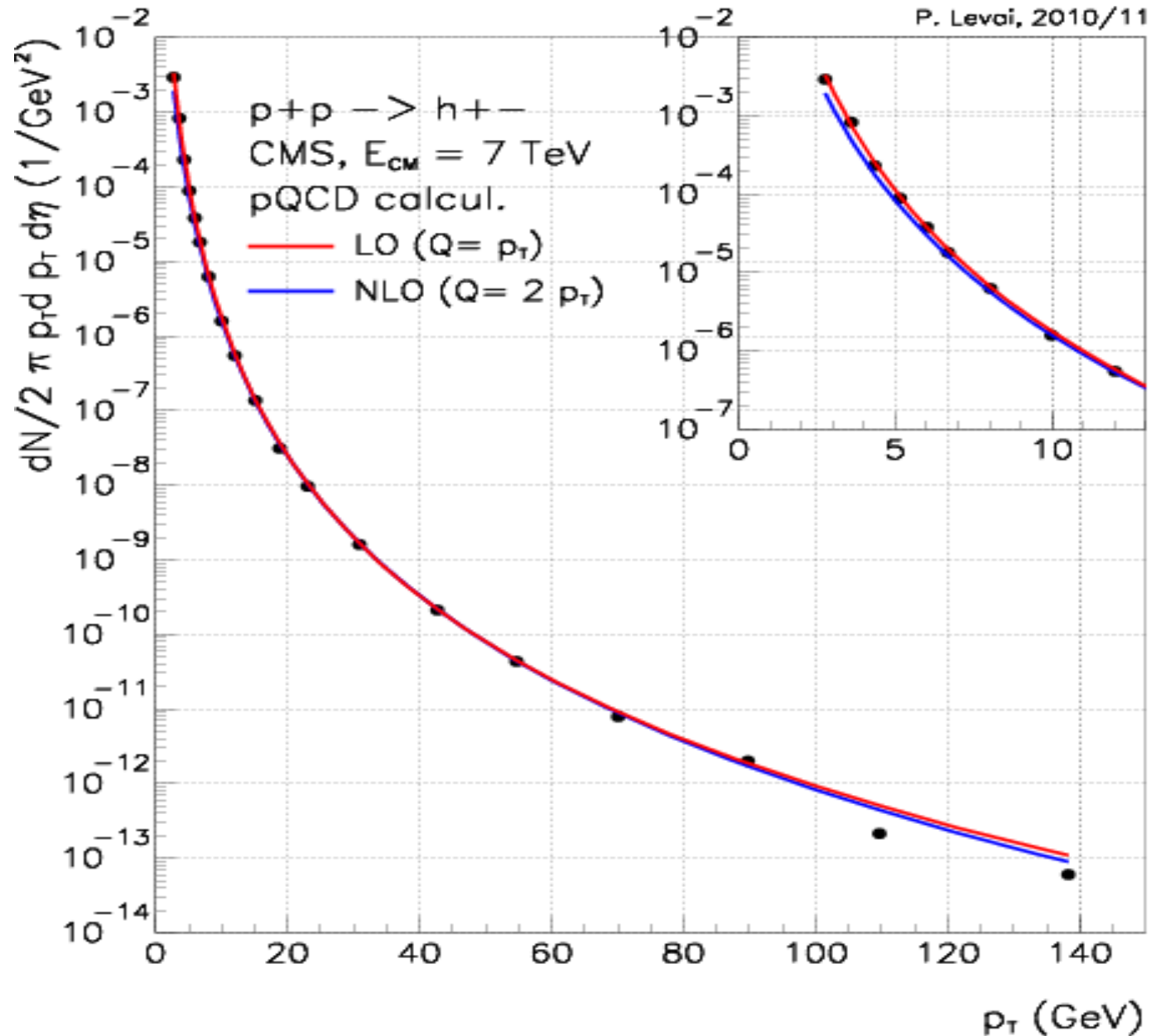
(Neither at
TEVATRON !)

Charged hadron production in pp collisions at 7 TeV – CMS data

NLO pQCD

LO and NLO descriptions are equally good !
(Scale difference)

We will use our pQCD frame at 2.36 ATeV.



Hard physics: pion production in dAu collision at high- p_T

Perturbative QCD calculations in LO and NLO for pp + CRONIN + SHADOWING:

SHADOWING: “New-Hijing” parametrization, **Li & Wang, PLB527 (2002) 85.**

$$f_{a/A}(x_a, Q^2) = A S_a^A(x_a, Q^2) f_{a/N}(x_a, Q^2) \quad 88$$

S.-Y. Li, X.-N. Wang / Physics L

Shadowing function for quarks:

$$S_q^A = 1.0 + 1.19 \log^{1/6} A (x^3 - 1.12 x^2 + 0.21 x) - s_q (A^{1/3} - 1)^{0.6} (1 - 3.5 \sqrt{x}) \exp(-x^2/0.01)$$

Shadowing function for gluons:

$$S_g^A = 1.0 + 1.19 \log^{1/6} A (x^3 - 1.2 x^2 + 0.21 x) - s_g (A^{1/3} - 1)^{0.6} (1 - 1.5 x^{0.35}) \exp(-x^2/0.004)$$

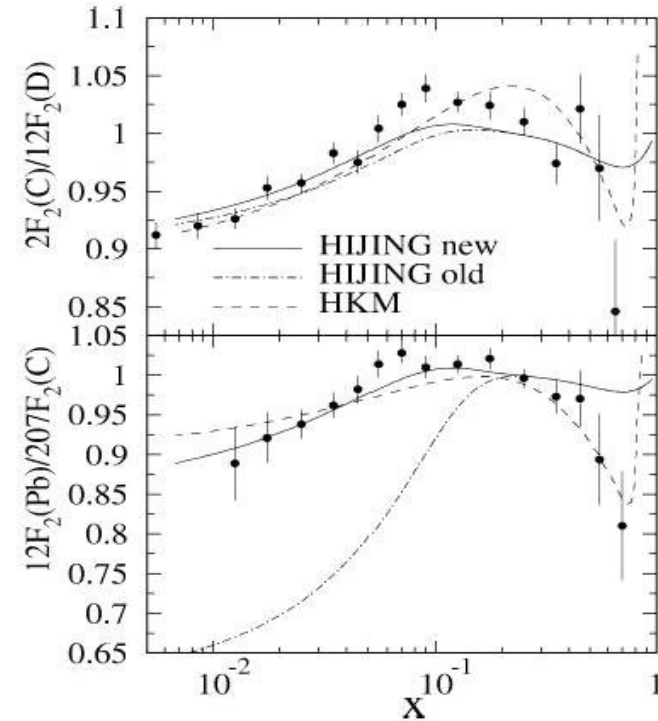


Fig. 2. Ratio of nuclear structure functions as measured in DIS. Solid lines are the new HIJING parameterization (Eq. (8)), dashed lines are the HKM parameterization [32] and dot-dashed lines are the old HIJING parameterization [16]. The data are from Ref. [30].

Hard physics: pion production in dAu collision at high- p_T

Perturbative QCD calculations in LO and NLO for pp + CRONIN + SHADOWING:

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Shadowing function for gluons:

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+ b impact parameter dependence:

$$s_i(\mathbf{b}) = s_i \frac{5}{3} \left(1 - b^2/R_A^2\right) \quad \leftarrow$$

Re-weighting

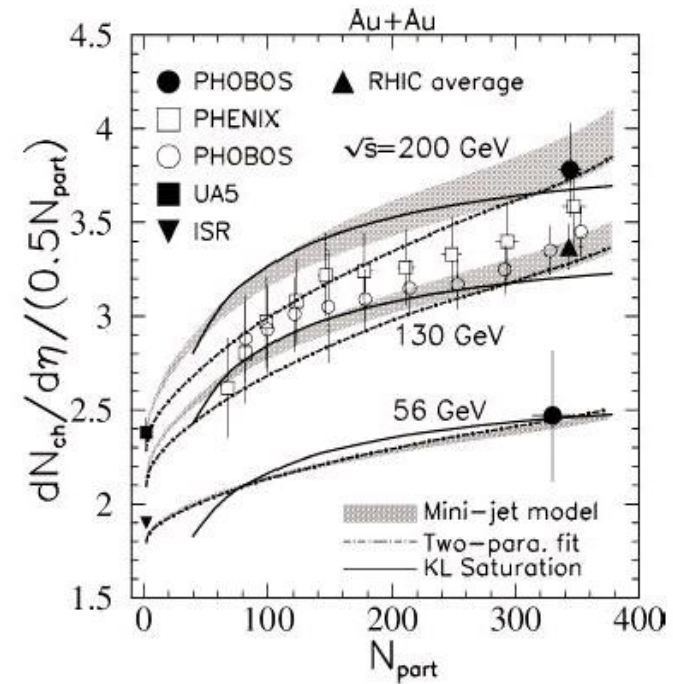
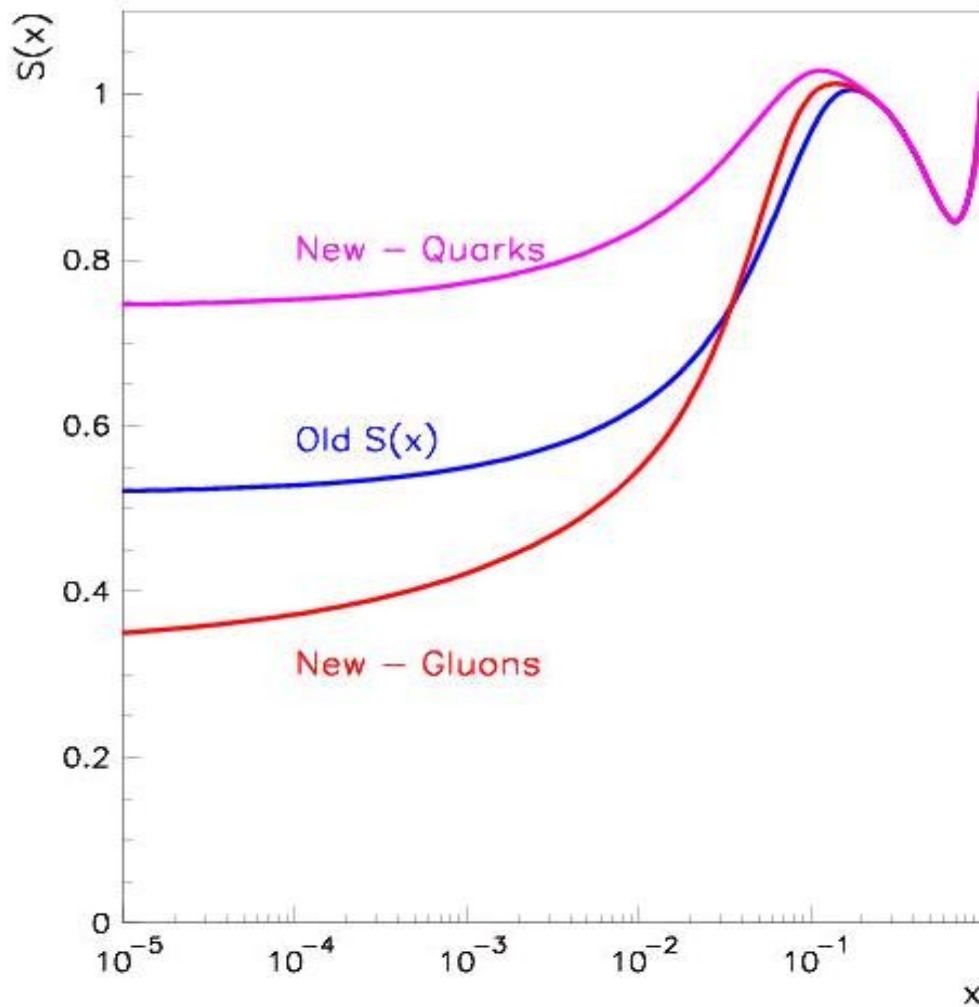


Fig. 4. The charged hadron central rapidity density per participant nucleon pair as a function of the averaged number of participants from the two-component model (shaded lines), two-parameter fit (Eq. (11)) (dot-dashed lines) and parton saturation model [9] as compared to experimental data [3,5,20,21].

Shadowing effect in dAu collisions:

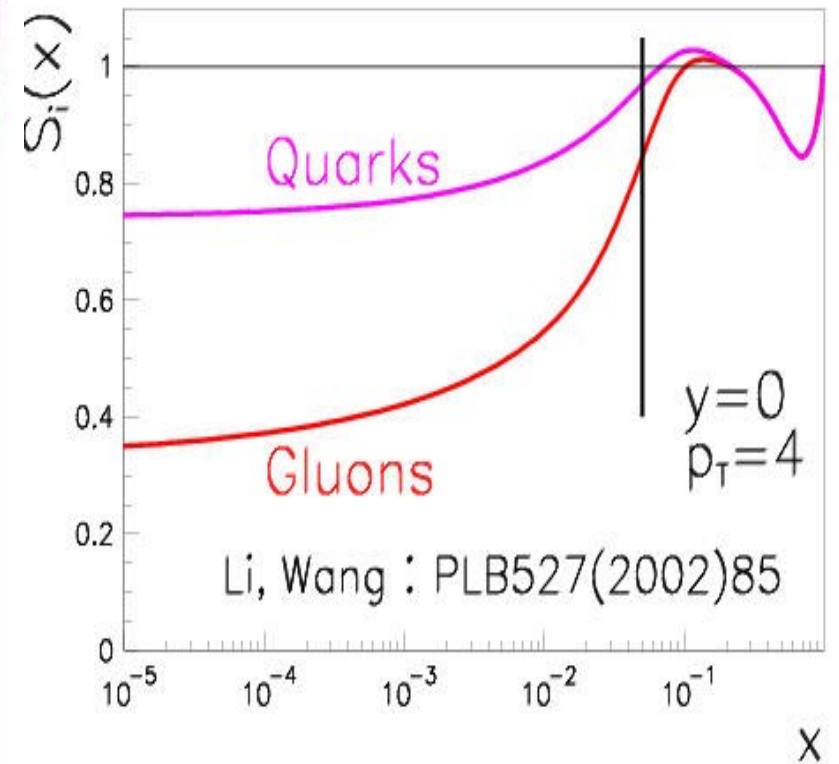
$S(x)$: shadowing function

Shadowing functions – HIJING old and new



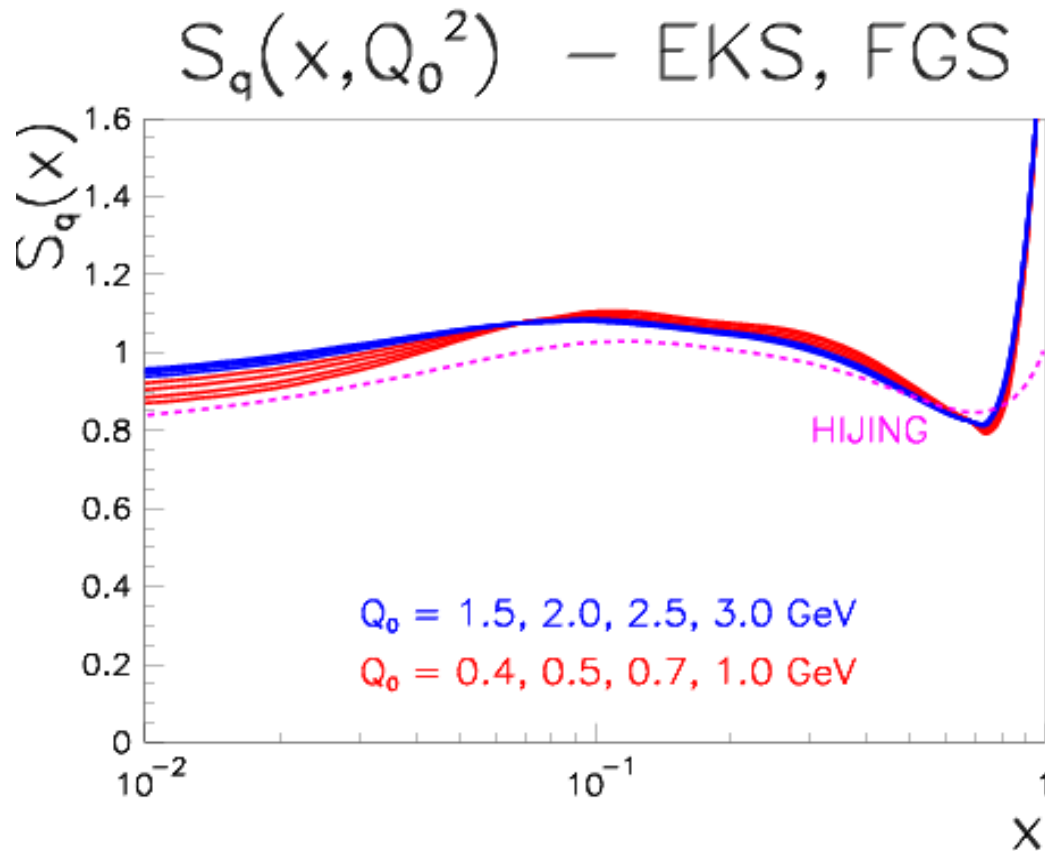
$S(x)$ at $y=0$ & $p_T = 4$ GeV

$S_i(x)$ shadowing functions



$X \rightarrow 0$: $S(x)$ is small

EKS99 shadowing function with enhanced antishadowing



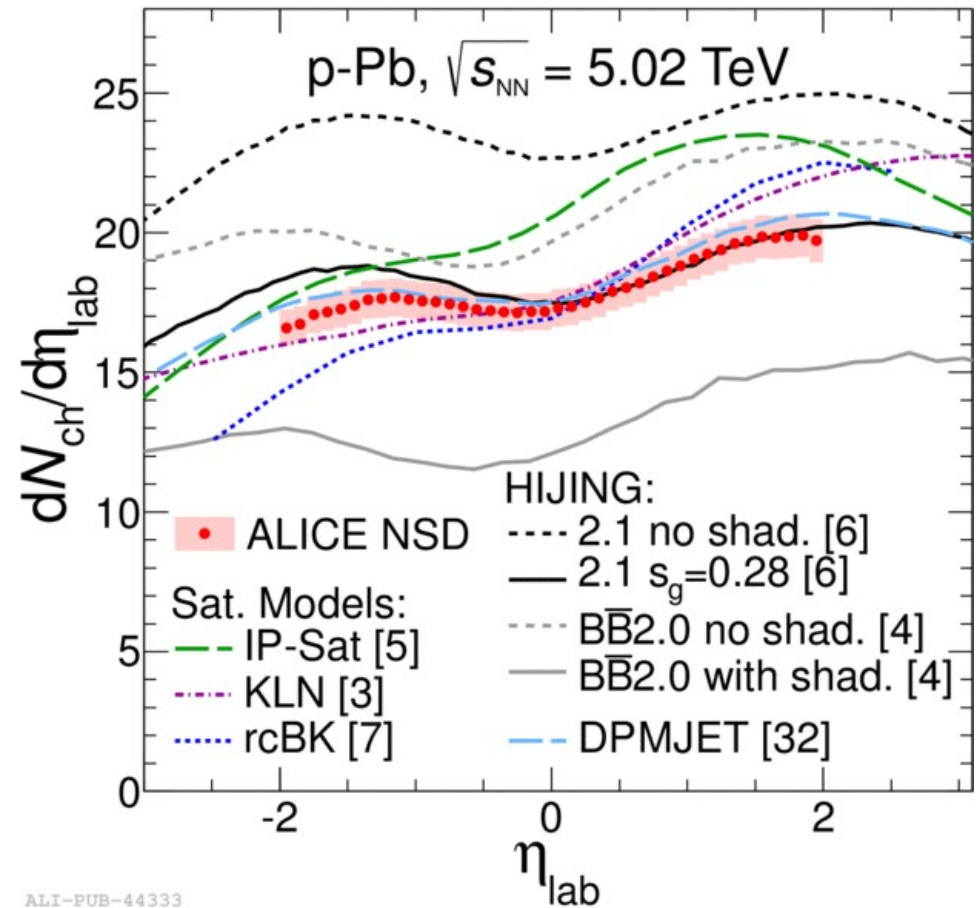
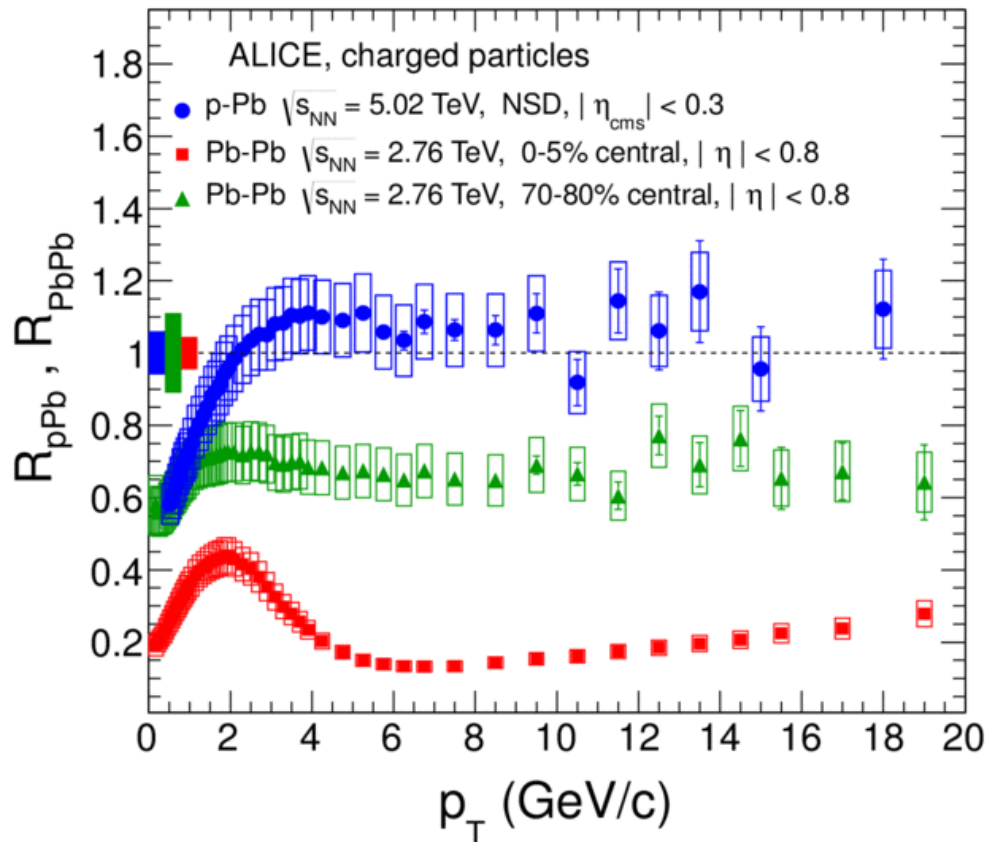
**K.J. Eskola,
V.J. Kolhinen,
C.A. Salgado**

EPJ C9, 61 (1999)

**EKS: antishadowing effect
for valence quarks**

**stronger than HIJING
at large-x
weaker than HIJING
at small-x**

Results at LHC in pPb collisions at $(s_{NN})^{1/2} = 5.02$ TeV:



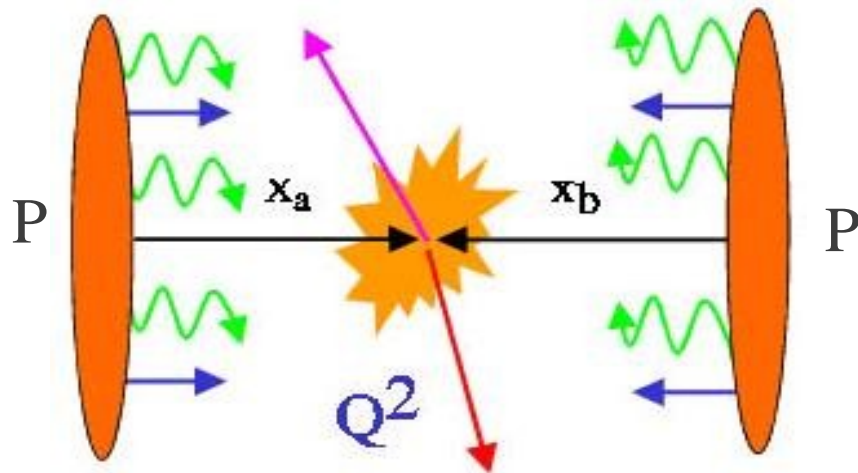
Shadowing effects are there, but they can be neglected in R_{pA}
 → in R_{AA} ?? My answer: YES, neglect it.

ALICE: arXiv: 1210.4520 and arXiv: 1210.3615

2. Jet energy loss

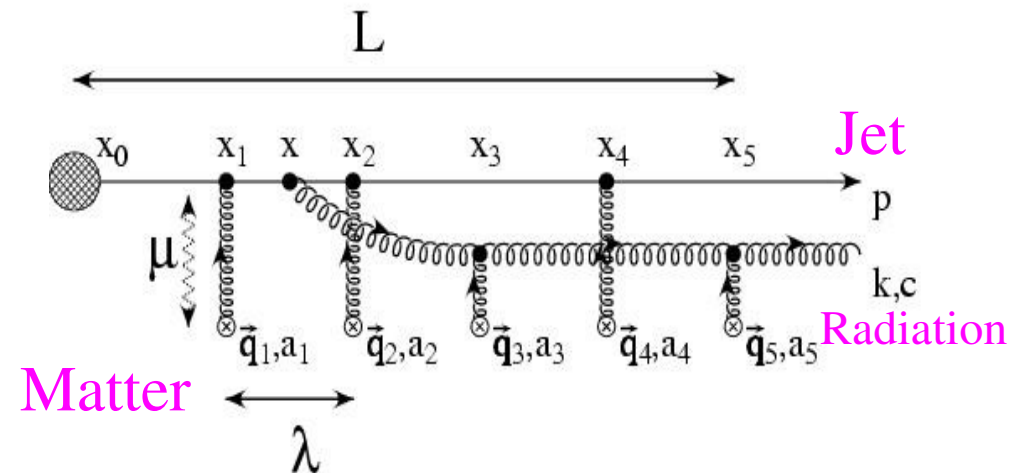
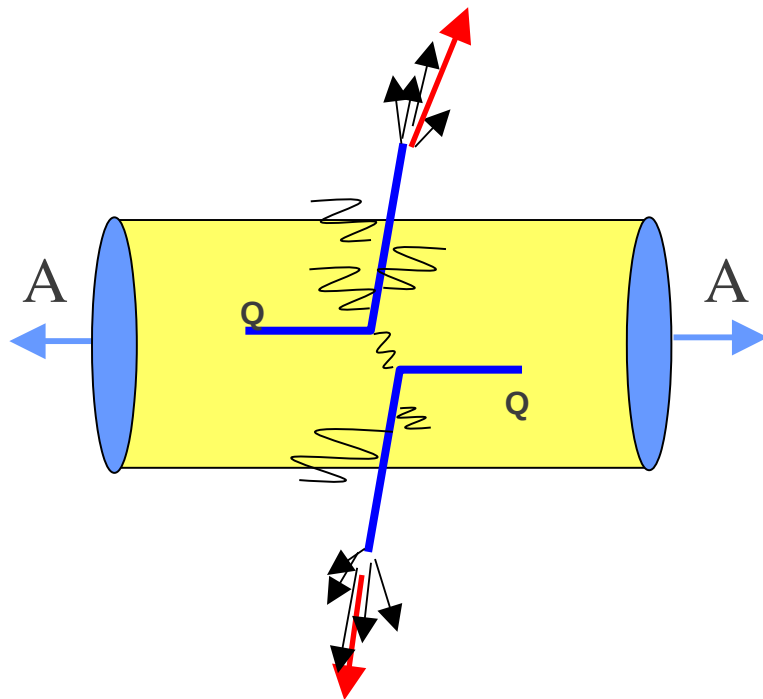
--- mechanism, description

Jets in pp and in AA collisions:

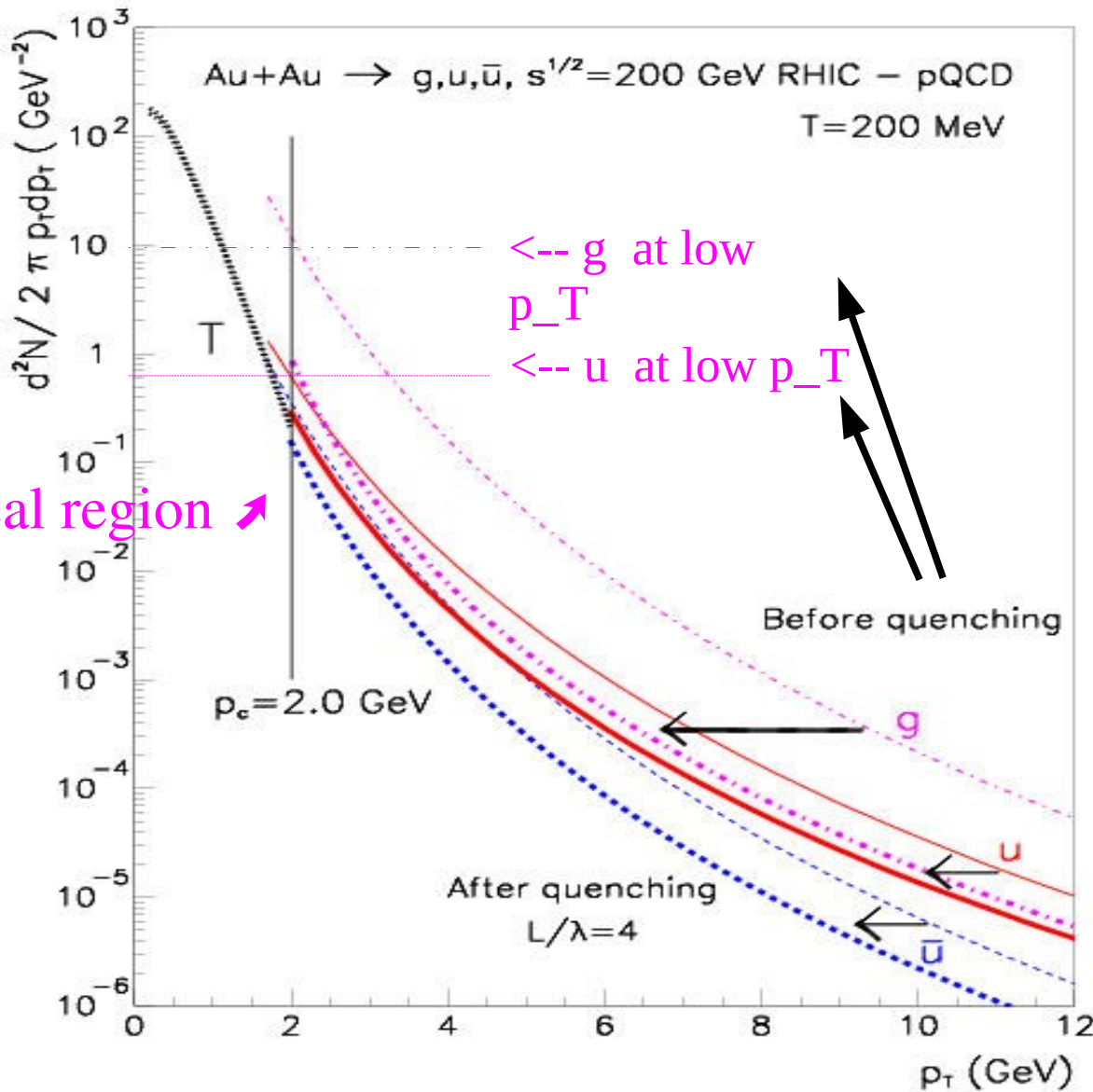


**Jet production in pp collision
(in vacuum):
→ pQCD description**

**Jet production and propagation
in AA collision (inside hot dense matter)
→ induced gluon radiation in a
modified pQCD description
JET-TOMOGRAPHY**



Initial parton distributions before and after quenching



Before quenching:

AuAu spectra

\equiv binary scaled pp spectra

primary partons \rightarrow

$p_T > p_c$: polinomial

$p_T < p_c$: suppressed or saturated

$p_T > p_c$:

80 % of dE_T/dy ($y=0$)

Initial parton distributions before and after quenching

After quenching:

radiated gluons →
 thermal bath ($T=200$ MeV)
 ≡ **hot dense matter formation**

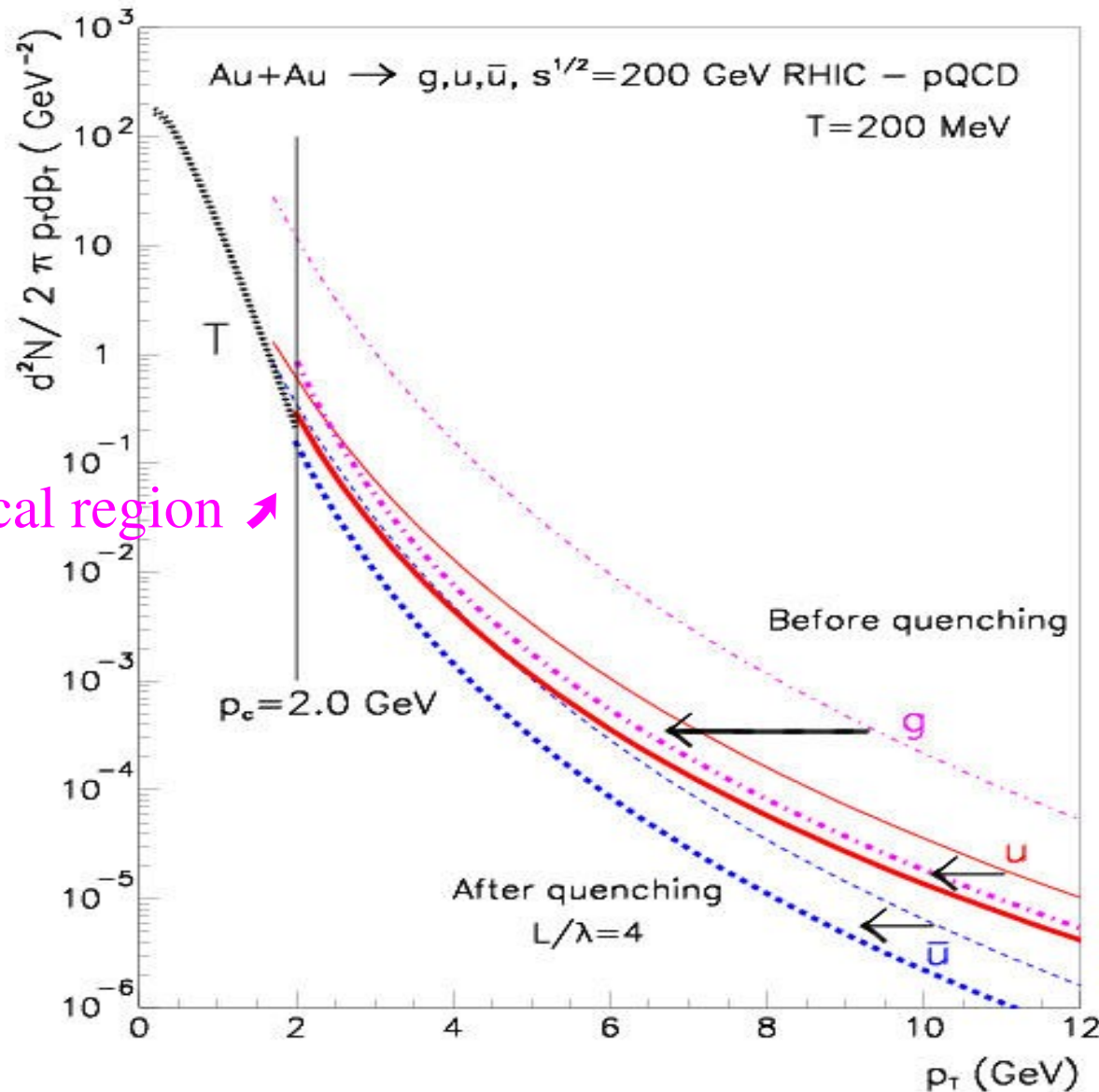
$dE_T/dy (y=0) = 570$ GeV
 80 % is soft ($p_T < p_c$)
 20 % is hard ($p_T > p_c$)

Modification in parton spectra
is shifted to smaller hadronic
 p_T in case of independent
jet fragmentation.

JET-QUENCHING



SUPPRESSION in
HADRON SPECTRA



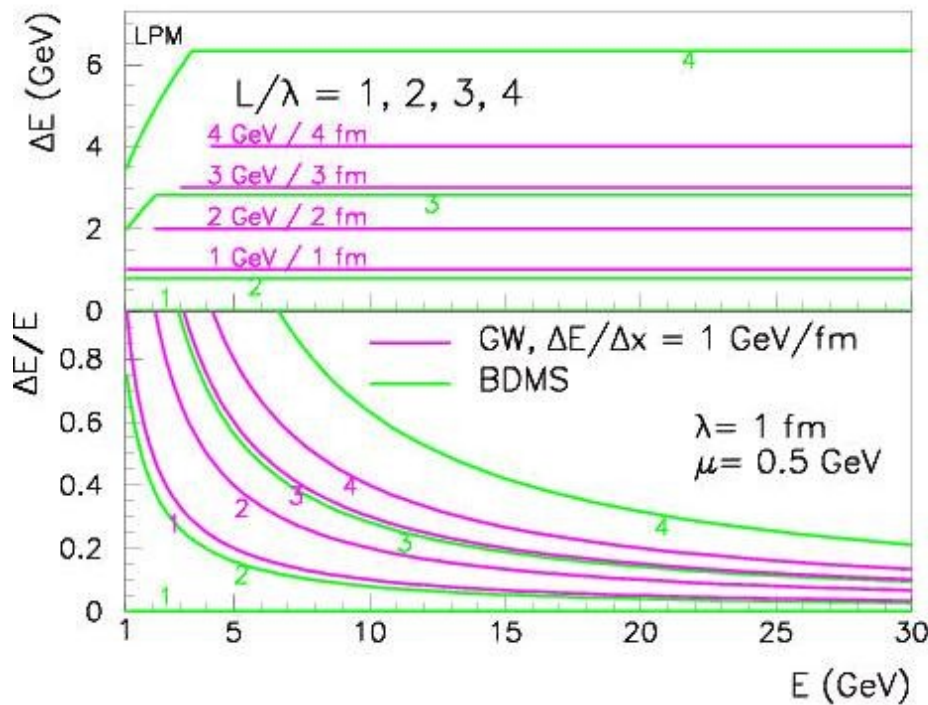
Critical region ↗

A+A collision \rightarrow induced jet energy loss : thick plasma $L \gg \lambda$

QED analogy (Gyulassy-Wang): 1-2 GeV/fm B. Betz: $dE/dx = c E^0$

Time ordered pQCD diagrams in infinite matter limit
 large number of very soft interaction ($L \gg \lambda_g$)
 (energy loss \rightarrow diffusion equation)

BDMS 97-98, Zakharov 96-98



$$\Delta E_{BDMS} \approx \frac{C_R \alpha_s}{4} \frac{L^2 \mu^2}{\lambda_g} \log \frac{L}{\lambda_g}$$

E-independent ΔE energy loss

**Strong E-dependence of $\Delta E/E$
 in the window
 $3 < E < 10$ GeV**

'Jet-quenching' : induced jet energy loss --- in thin colored matter

M. Gyulassy, P. Levai, I. Vitev,

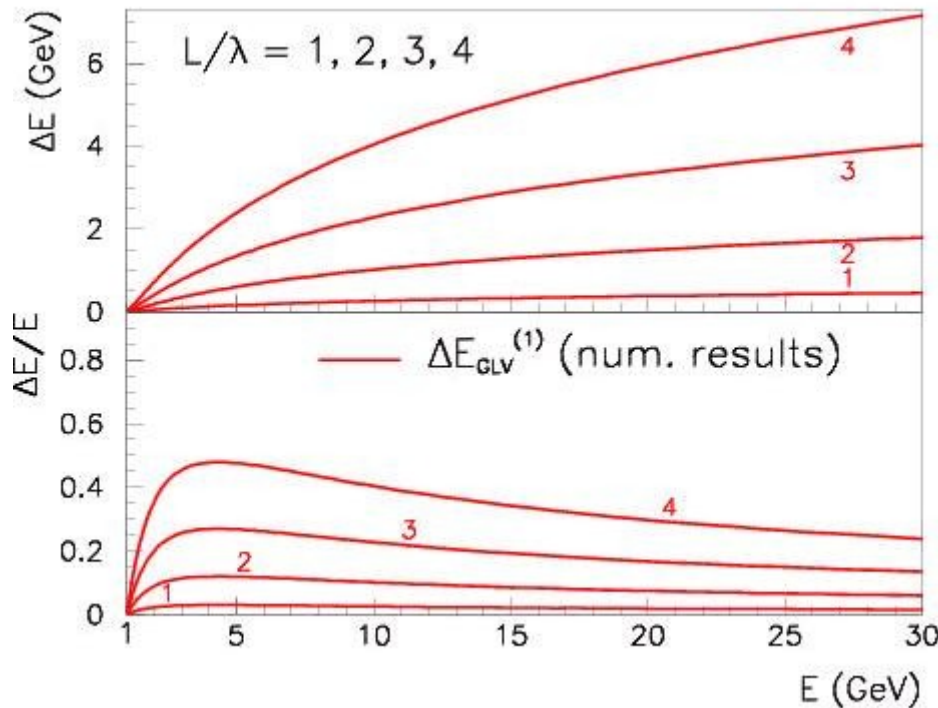
PRL85,5535(2000), NPB594,371(2001)

GLV: time-ordered pQCD (Feynman diagrams)

+ OPACITY expansion ($N = 1, 2, 3, \dots$)

+ kinematical cuts

$$\Delta E_{GLV} \approx \frac{C_R \alpha_s}{N(E)} \frac{L^2 \mu^2}{\lambda_g} \log \frac{E}{\mu}$$



E-dependent ΔE energy loss

**E-independent $\Delta E/E$
in the window**

$3 < E < 10-15$ GeV

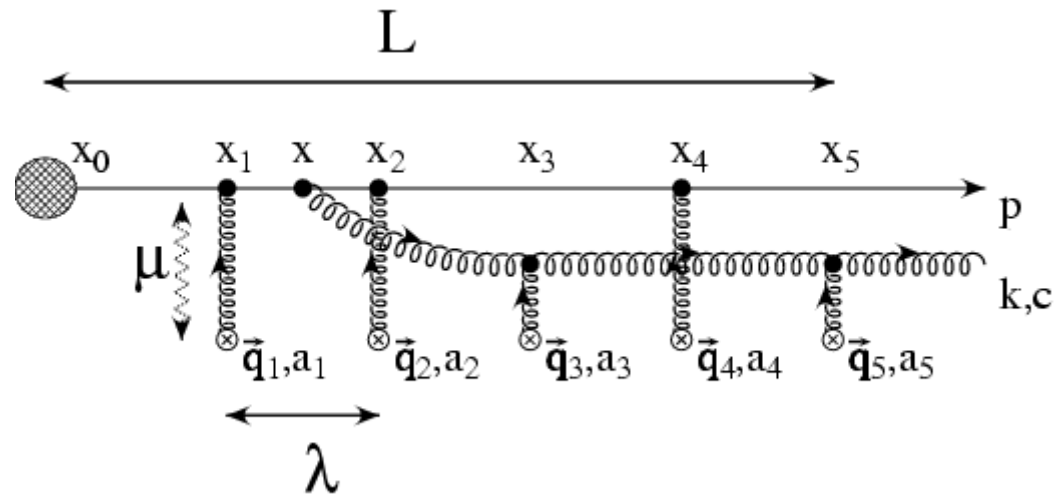
$$L/\lambda \rightarrow \int \tau \rho(\tau) d\tau$$

Opacity \rightarrow Density

Induced jet energy loss --- agreements and disagreements:

BDMS, GW, GLV, Zakharov, Wiedemann, Salgado, ...

1. $\Delta E_{loss} \sim L^2$ **non-abelian nature**
2. $\Delta E_{loss} \sim \hat{q}$ **transport coefficient: $\hat{q} \approx \mu^2/\lambda \approx \int d^2 q_T q_T^2 d\sigma/dq_T^2$**
3. $\Delta E_{loss} = C_R \alpha_S \hat{q} L^2 F[...]$ **where $F[...]$ depends on theories**



Coherence & Interference

Induced jet energy loss in expanding matter: D. Molnar (Wuhan)

1. Averaged opacity \rightarrow time dependent color density:

$$1/\lambda_{col} = \sigma_{el} \rho_{col} \rightarrow \frac{9/2 \pi \alpha_s^2}{\mu^2} \frac{2}{L^2} \int_0^L \tau \rho_{col}(\tau) d\tau$$

2. 1-DIM Bjorken expansion:

$$\rightarrow \frac{9/2 \pi \alpha_s^2}{\mu^2} \frac{2}{L^2} \frac{1}{A_T} \frac{dN^{col}}{dy} L$$

3. Energy loss with rapidity density:

$$\Delta E_{GLV}^{1DIM} \approx \frac{9 C_R \pi \alpha_s^3}{4} \frac{1}{A_T} \frac{dN^{col}}{dy} L \log \frac{2 E}{\mu^2 L}$$

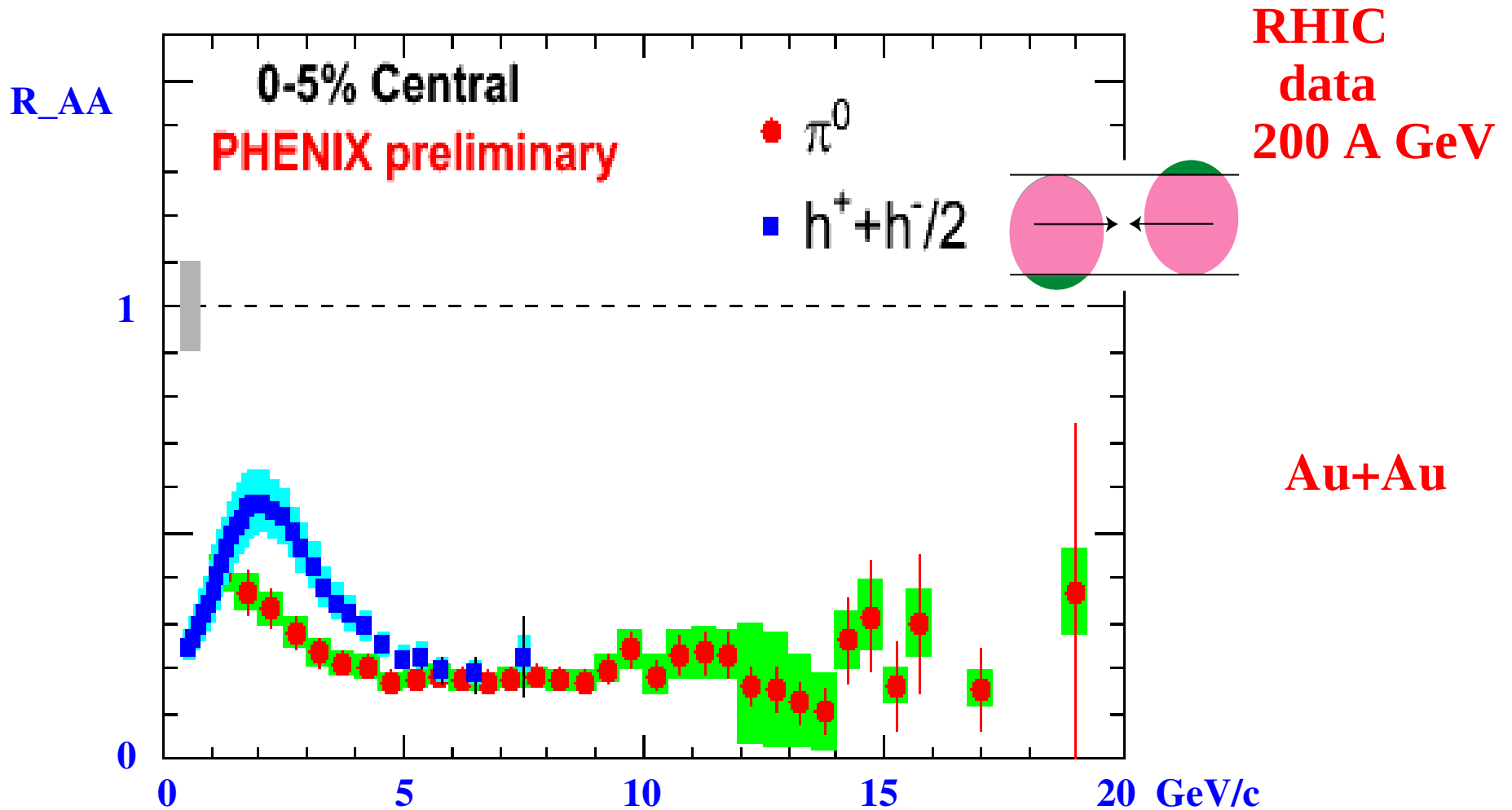
3. Data vs. Theory at RHIC and LHC energies
--- answers and new questions

Hard physics: pion production in AA collision at high- p_T

Perturbative QCD calculations in NLO for heavy ion collisions:

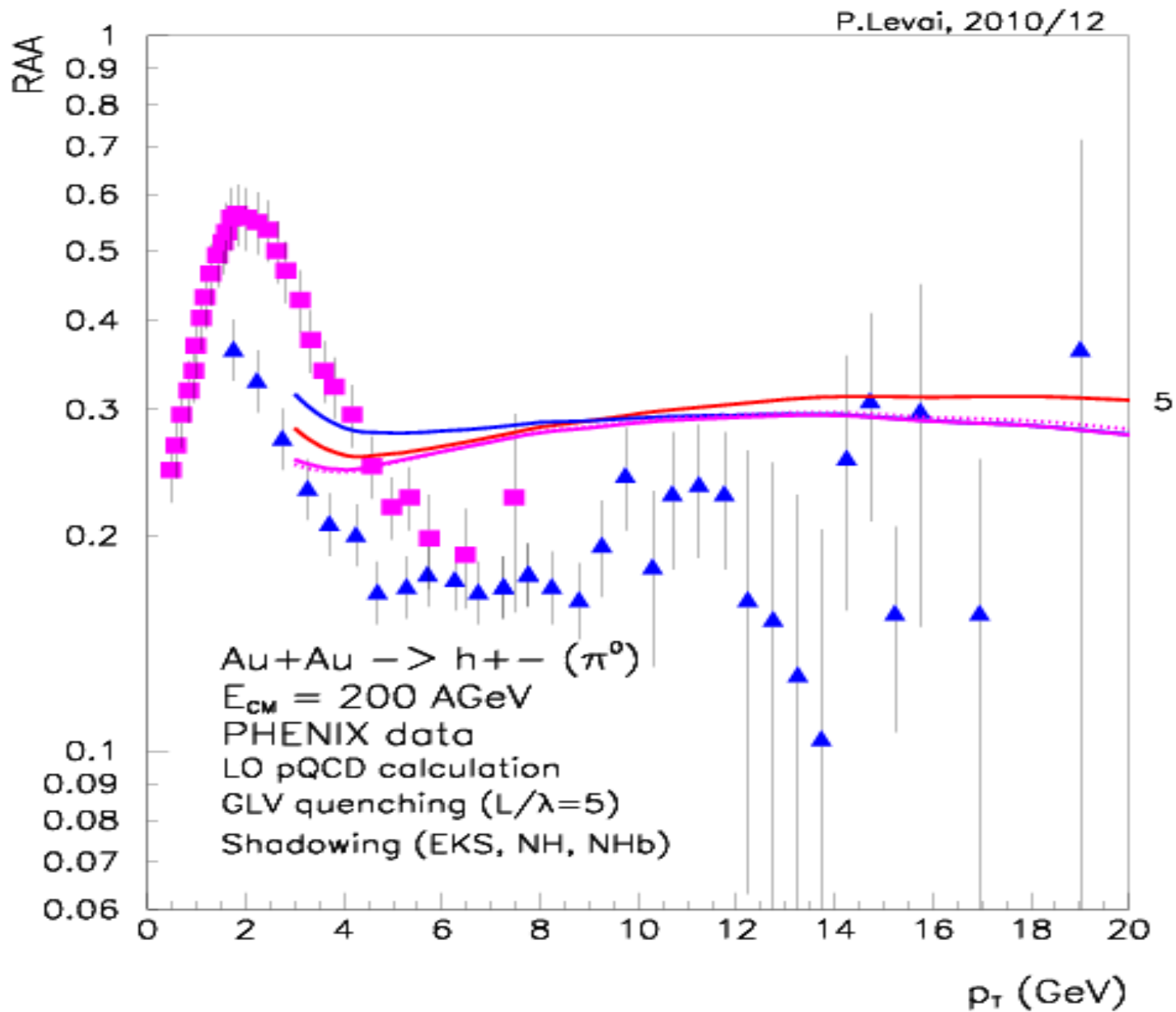
geometrical overlap + shadowing, multiscattering, jet-quenching, ...

$$E_\pi \frac{d\sigma^{AB}}{d^3p_\pi} = \int d^2b d^2r t_A(\vec{r}) t_B(|\vec{b}-\vec{r}|) E_\pi \frac{d\sigma^{pp}}{d^3p_\pi} \otimes S(\dots) \otimes M(\dots) \otimes Q(\dots)$$

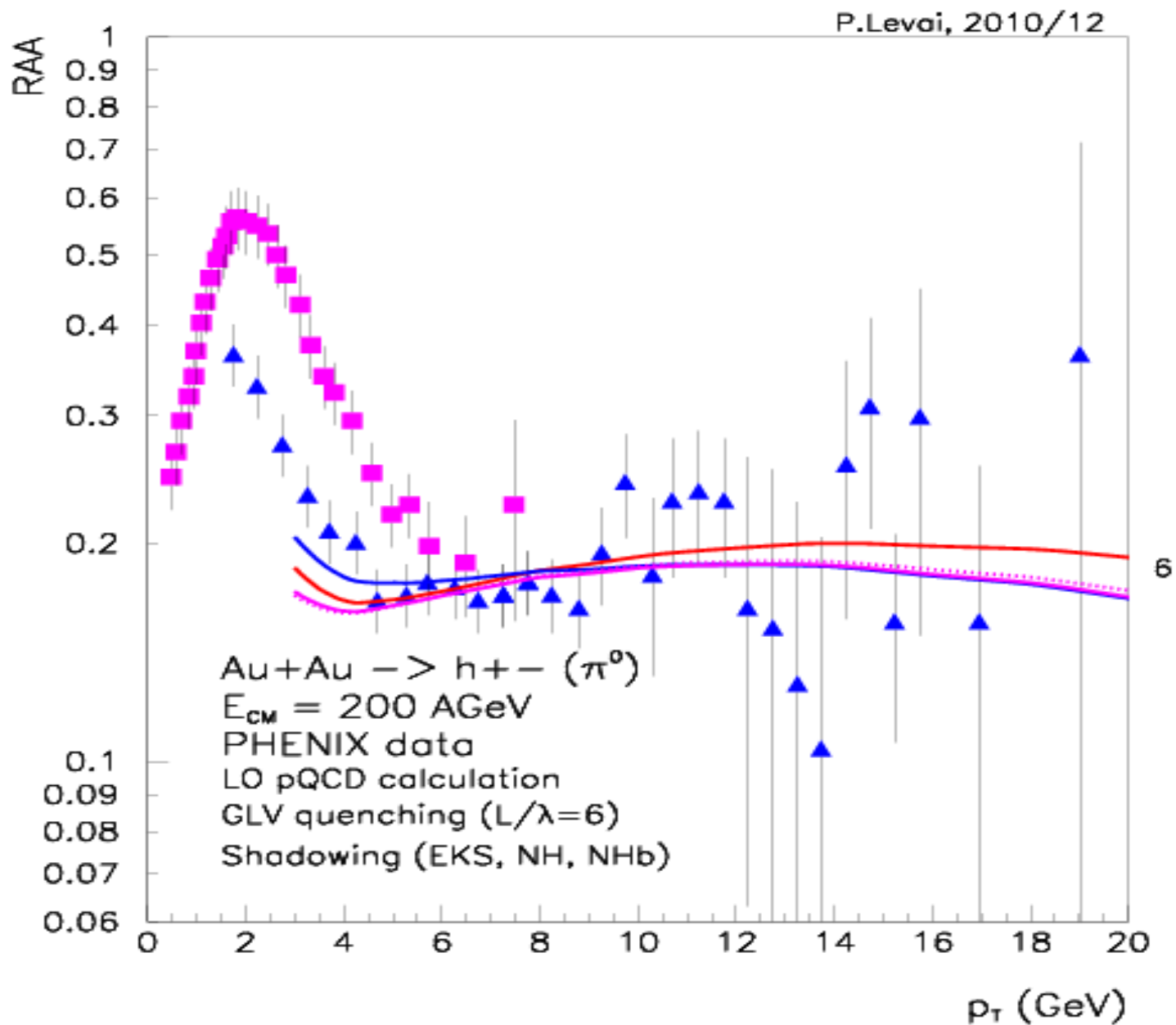


Most central Au+Au collisions (5%) at RHIC 200 AGeV

“Quenching at $L/\lambda=5$ + Shadowing”



Most central Au+Au collisions (5%) at RHIC 200 AGeV “Quenching at $L/\lambda=6$ + Shadowing”



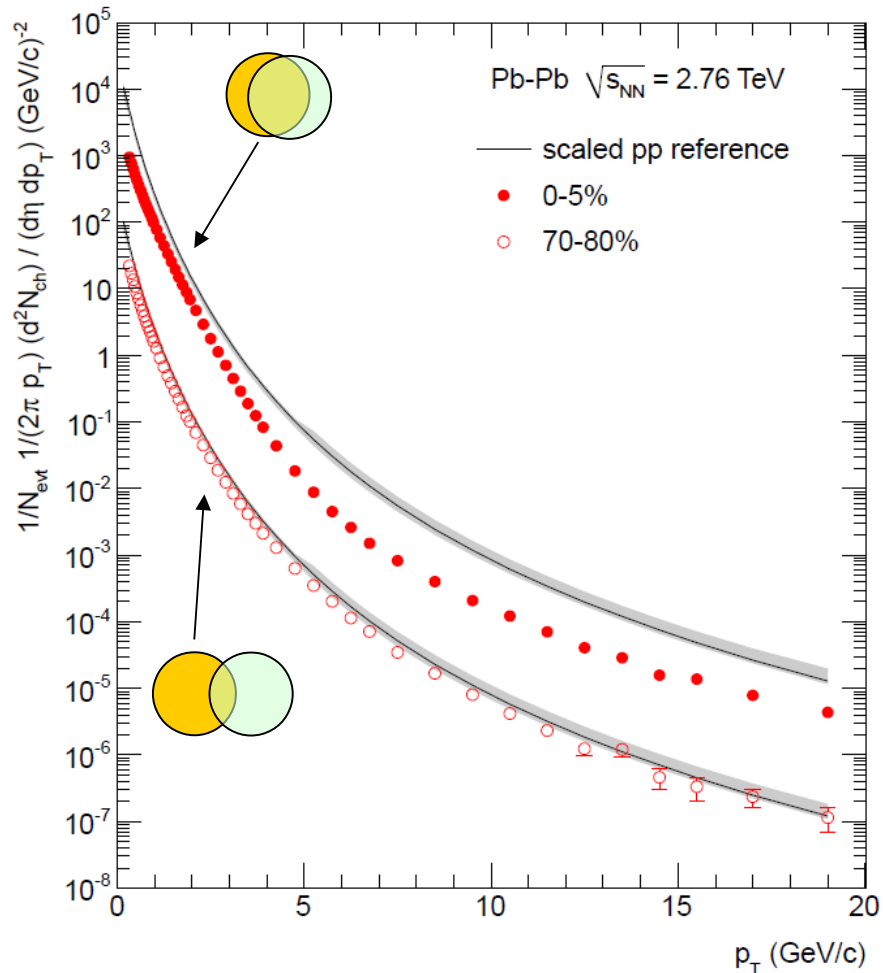
**Conclusion from light quark quenching
at RHIC energy:**

**$L/\lambda = 5.5$ will describe RHIC data well
Shadowing has small influence**

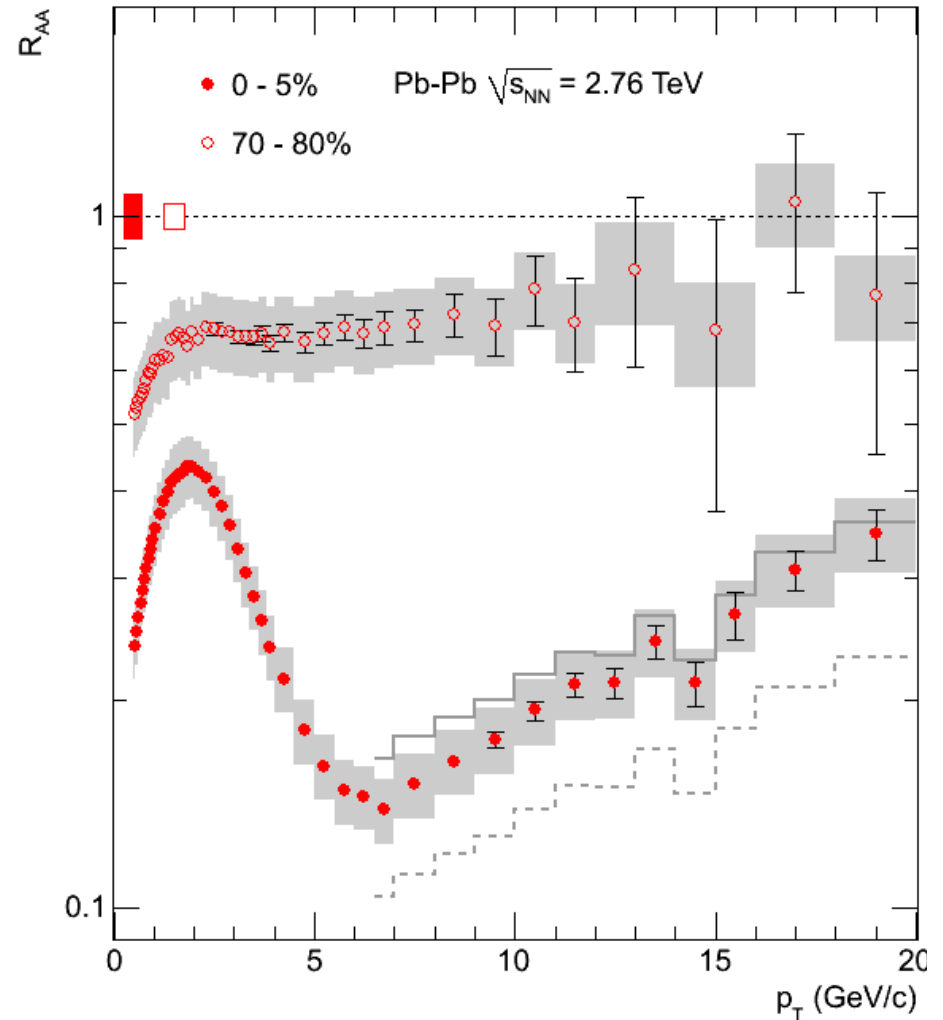
Results at LHC energies:

Pb+Pb at 2.76 ATeV

Charged hadron production
in PbPb collision at 2.76 TeV
--- first ALICE data ---
(PLB696,30,2011)

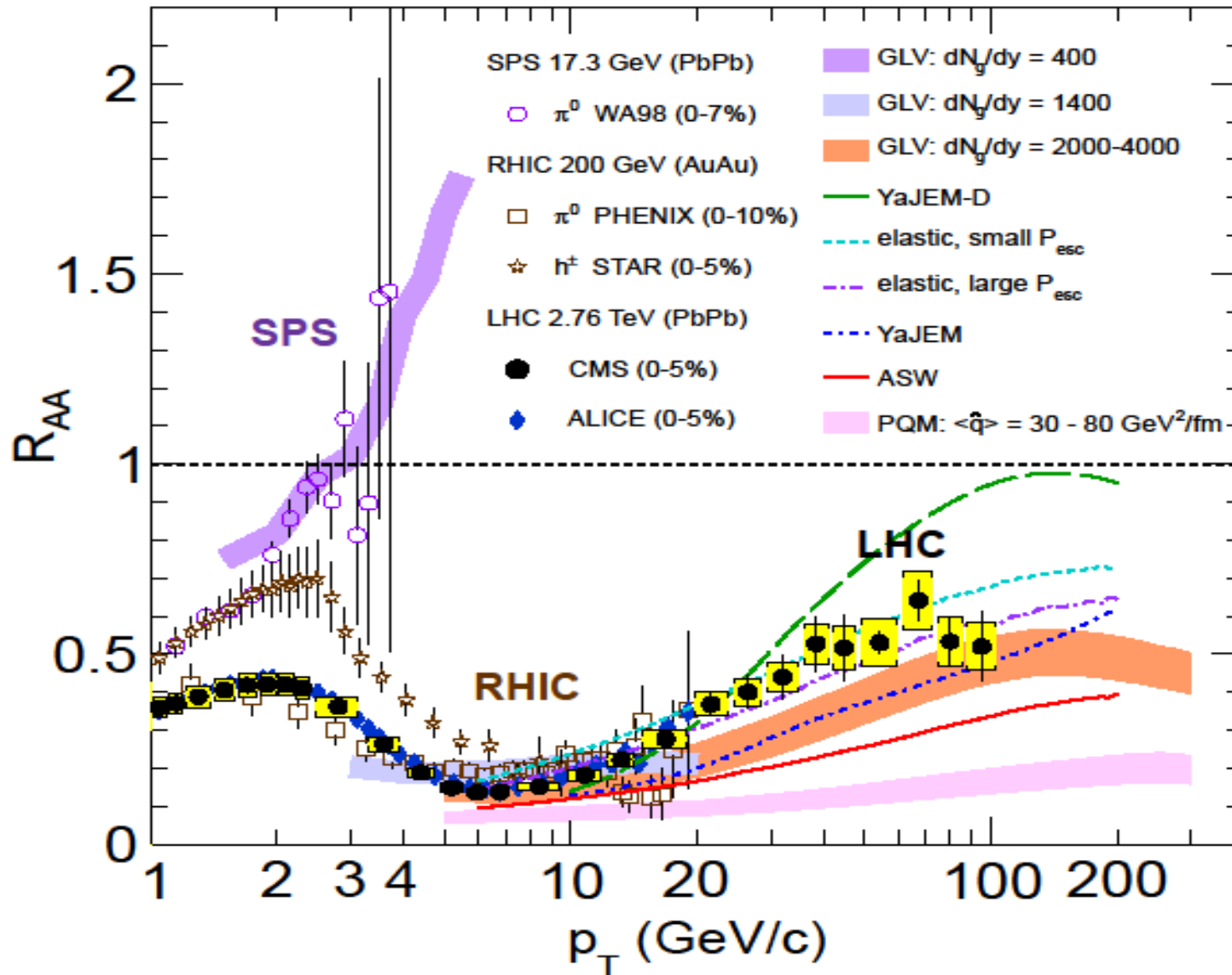


$$R_{AA}(p_T) = \frac{(1/N_{evt}^{AA}) d^2 N_{ch}^{AA} / d\eta dp_T}{\langle N_{coll} \rangle (1/N_{evt}^{pp}) d^2 N_{ch}^{pp} / d\eta dp_T}$$



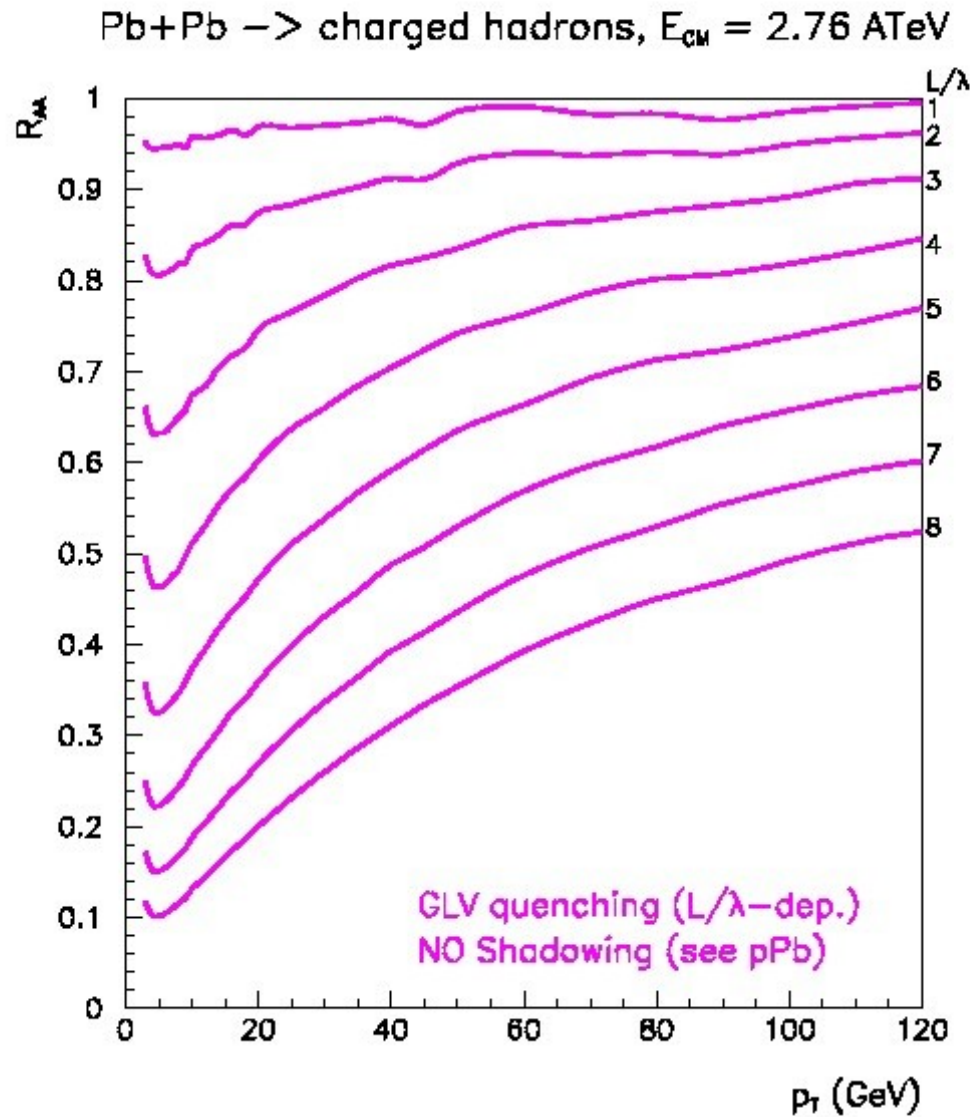
Most central Pb+Pb collisions (5%) at LHC 2.76 ATeV

Data from ALICE and CMS (EPJ C72, 1945,2012)



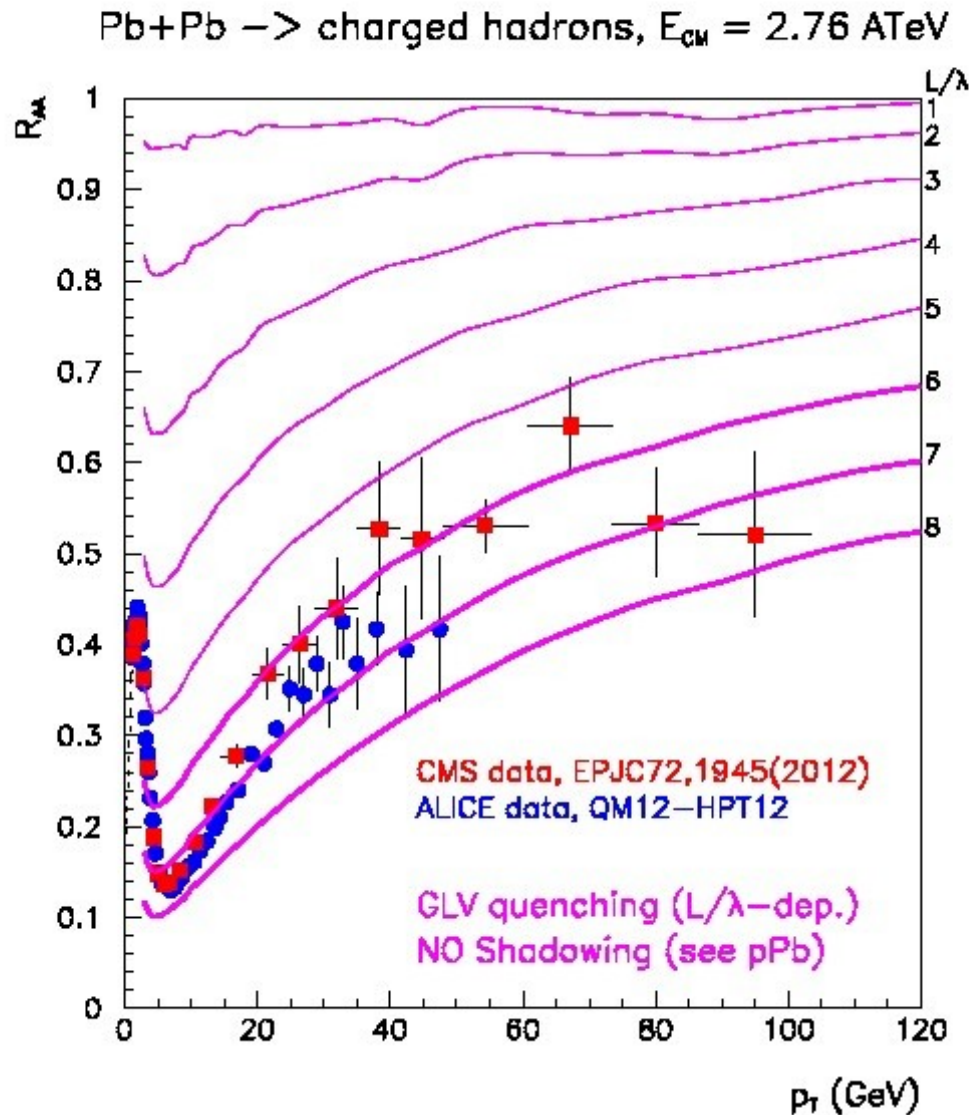
Most central Pb+Pb collisions (5%) at LHC 2.76 ATeV

Theory from GLV – energy loss only



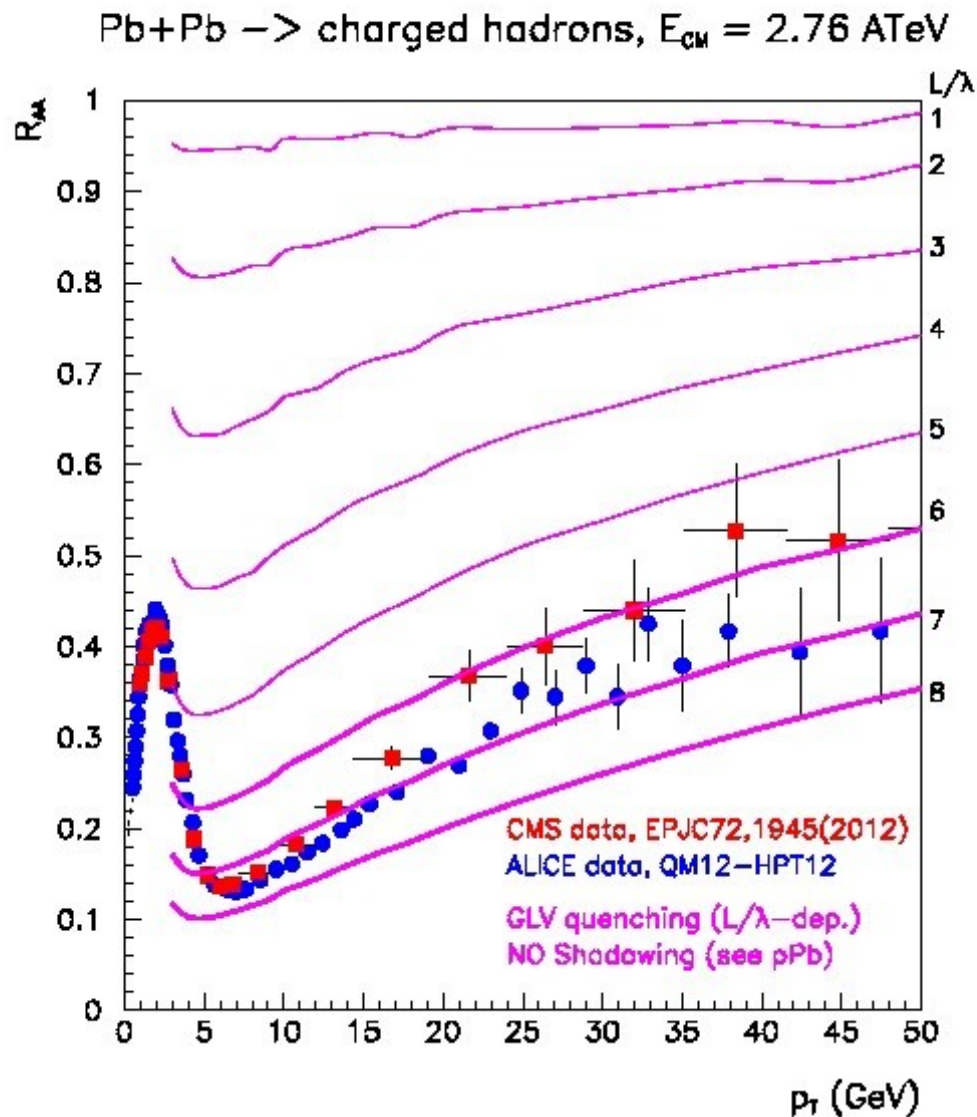
Most central Pb+Pb collisions (5%) at LHC 2.76 ATeV

ALICE and CMS Data vs. Theory from GLV – energy loss only



Most central Pb+Pb collisions (5%) at LHC 2.76 ATeV

ALICE and CMS Data vs. Theory from GLV – energy loss only



$L/\lambda = 7$ is needed

**Conclusion from light quark quenching
in PbPb collisions at 2760 AGeV LHC energy:**

**$L/\lambda = 7$ is needed without shadowing
(or with EKS shadowing)**

**This means 2x larger color density
at LHC energy w.r.t RHIC energy**

What about the q and g details?

Jet energy loss for quark and gluon jets

ALICE and CMS Data vs. Theory from GLV – energy loss only

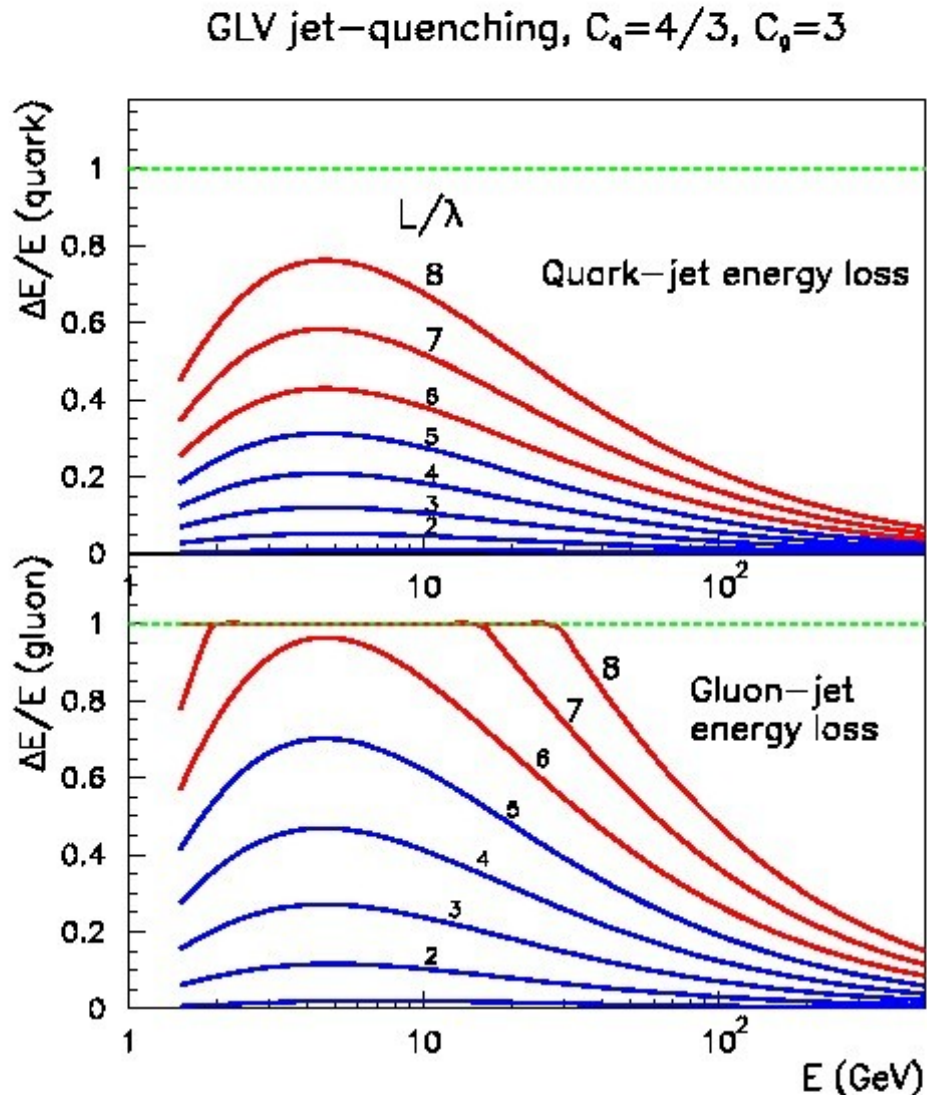
All gluon jet suffer quenching in the windows:

$$L/\lambda = 7$$

$$2 < p_T < 12 \text{ GeV}/c$$

$$L/\lambda = 8$$

$$1 < p_T < 20 \text{ GeV}/c$$



What will happen with them?

- 1, Thermalization, QGP**
- 2, Coherent gluon field & quark dominance**

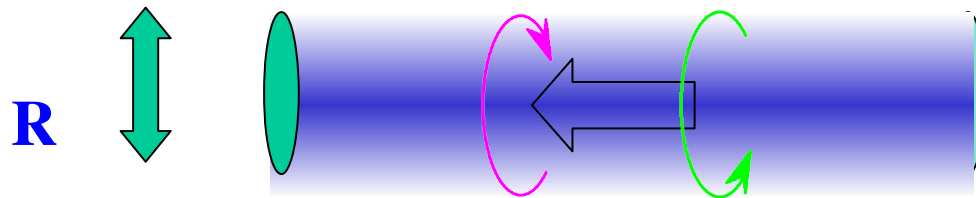
- 4. A new channel for extra proton production**
 - coherent gluon field and diquark coalescence**

A further (good) model for particle production:

III. Non-perturbative, non-asymptotic color transport:

“confined flux tube formation and breaking”

- phenomenological approximations are known (string, rope)
- phenomenology is applied successfully in string-based codes
- FRITIOF, PYTHIA, HIJING 1.0/2.0 are using strings
- URQMD, HIJING-BB is using ropes (melted strings)
- good agreement with data at different energies



- formal QCD-based equations are known (Heinz, Mrowczynski)
- YM-field evolution in 3+1 dim, collision (Poschl, Müller)
- lattice-QCD calculations have been started (Krasnitz, Lappi)
- time-dependent strong fields, color kinetic eq. (Skokov, PL)
- ...

A further model for particle production:

III. Non-perturbative, non-asymptotic color transport: “pair-creation in strong fields”

--- strong (Abelian) **static** E field: Schwinger mechanism
probability of pair-creation:

$$P(p_T) d^2 p_T = -\frac{eE}{4\pi^3} \ln\left(1 - \exp\left[-\pi \frac{m^2 + p_T^2}{eE}\right]\right) d^2 p_T$$

integrated probability at mass m:

$$P_m = \frac{(eE)^2}{4\pi^3} \sum_{n=1}^{\infty} \frac{1}{n^2} \exp\left[-\pi \frac{nm^2}{eE}\right]$$

ratio of production rates (e.g. strange to light)

$$\gamma_s = \frac{P(s\bar{s})}{P(q\bar{q})} = \exp\left[-\pi \frac{m_s^2 - m_q^2}{eE}\right] \quad eE = 0.9 \text{ GeV} / \text{fm}$$

--- strong time dependent SU(N) color fields:

Kinetic Equation for the color Wigner function

A.V. Prozokevich, S.A. Smolyansky, S.V. Ilyin, hep-ph/0301169.

How could we decide between the two scenarios:

**1, Gluon energy loss → thermal bath of QGP →
hadronization with quark coalescence**

**2, Gluon energy loss → coherent gluon field →
quark pair production from strong field →
quark-antiquark dominated plasma (QAP) →
hadronization with quark coalescence**

Answer:

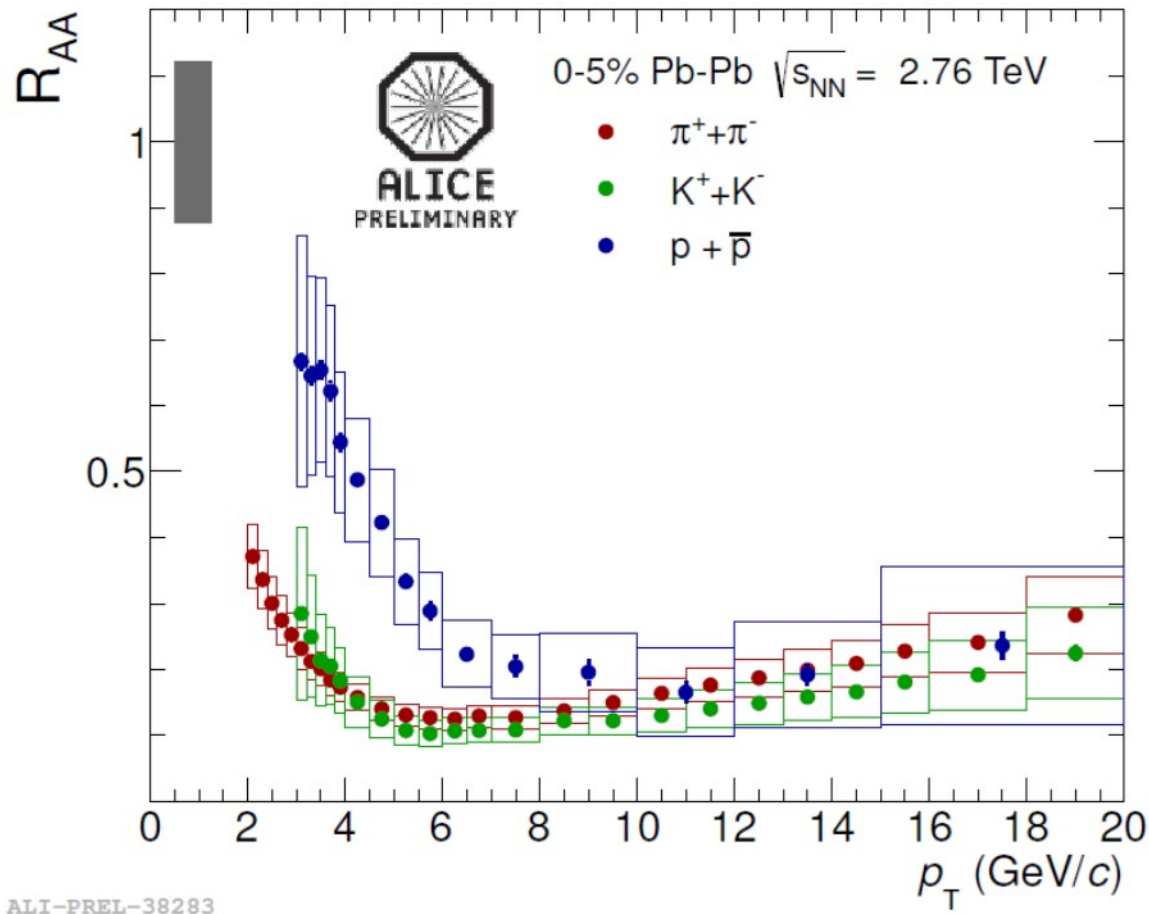
In 1: Thermal coalescence of quarks and antiquarks

**In 2: Schwinger production of quark-antiquark pairs
and diquark-antidiquark pairs**

**Coalescence of (Schwinger) quark, antiquark
and diquark, anti-diquark**

→ proton/pion ratio (R_{AA}) can be different !!!

New data from ALICE: R_{AA} of identified charged hadrons in PbPb



Theoretical investigations:

D. Berényi, A. Pasztor, V.V. Skokov, P.L.

J. Phys. G36 (2009) 064068; G38 (2011) 124155.

arXiv: 1208.0448

Kinetic equation for fermion pair production:

Wigner function: $W(k_1, k_2, k_3)$

Color decomposition: $W = W^s + W^a t^a$, where $a = 1, 2, \dots, N_c^2 - 1$

Spinor decomposition: $W^{s;a} = a^{s;a} + b_\mu^{s;a} \gamma^\mu + c_{\mu\nu}^{s;a} \sigma^{\mu\nu} + d_\mu^{s;a} \gamma^\mu \gamma^5 + i e^{s;a} \gamma^5$

Color vector field (longit.): $A_\mu^a = (0, -\vec{A}) = (0, 0, 0, A_3^a)$

Kinetic equation for Wigner function:

$$\begin{aligned} \partial_t W + \frac{g}{8} \frac{\partial}{\partial k_i} \left(4 \{ W, F_{0,i} \} + 2 \{ F_{i\nu}, [W, \gamma^0 \gamma^\nu] \} - [F_{i\nu}, \{ W, \gamma^0 \gamma^\nu \}] \right) = \\ = i k_i \{ \gamma^0 \gamma^i, W \} - i m [\gamma^0, W] + i g [A_i, [\gamma^0 \gamma^i, W]]. \end{aligned}$$

for details see V.V. Skokov, PL: PRD71 (2005) 094010 for U(1)
PRD78 (2008) 054004 for SU(2)
in preparation for SU(3)

Distribution function for fermions with mass m:

$$f_f(\vec{k}, t) = \frac{m a^s(\vec{k}, t) + \vec{k} \vec{b}^s(\vec{k}, t)}{\omega(\vec{k})} + \frac{1}{2}$$

Time dependent external field, $E(t)$ and neglected mass, $m=0$:

A, Pulse field (dotted):

$$E_{pulse}(t) = E_0 [1 - \tanh^2(t/\delta)]$$

B, Constant field (dashed):

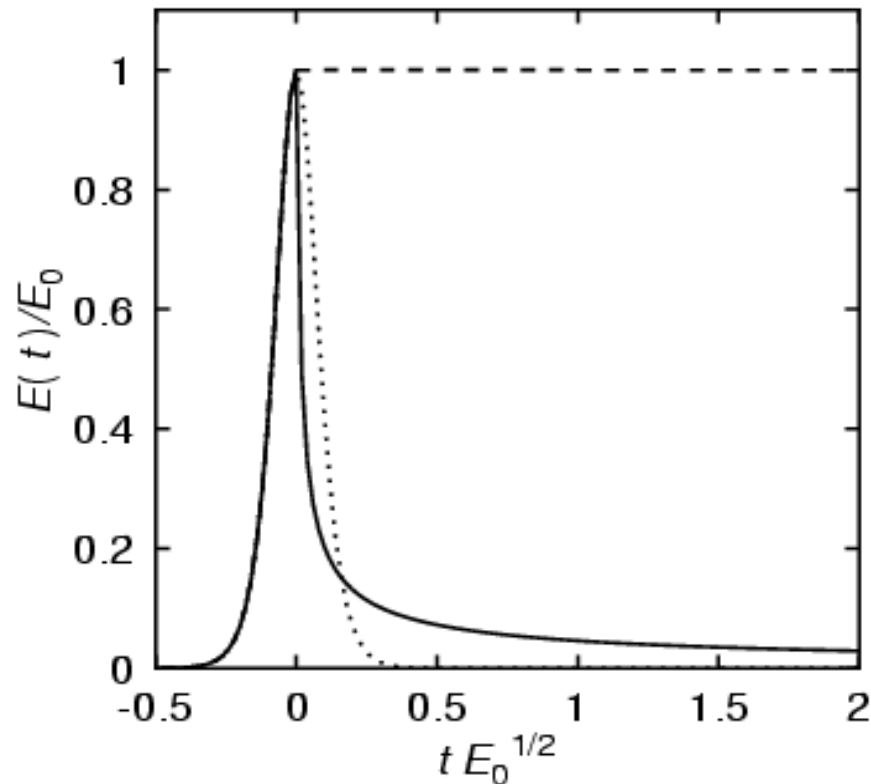
$$E_{const}(t) = E_{pulse}(t) \quad \text{at } t < 0$$

$$E_{const}(t) = E_0 \quad \text{at } t > 0$$

C, Scaled field (solid):

$$E_{scaled}(t) = E_{pulse}(t) \quad \text{at } t < 0$$

$$E_{scaled}(t) = \frac{E_0}{(1+t/t_0)^\kappa} \quad \text{at } t > 0$$



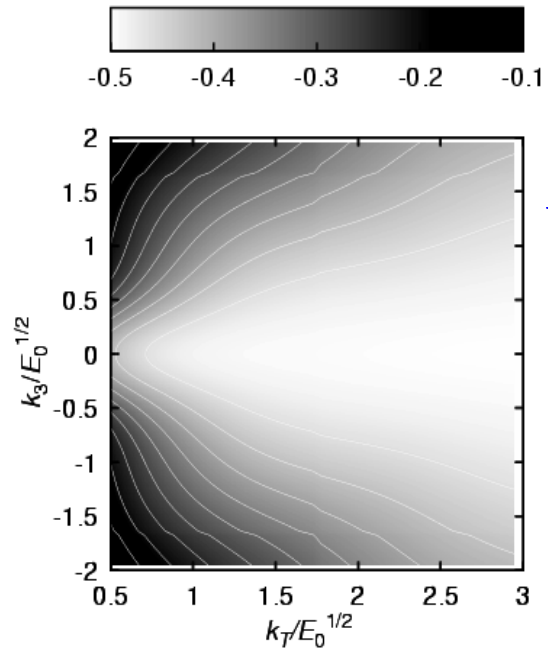
$$\delta = 0.1 / E_0^{1/2} \quad \text{at RHIC energy}$$

$$\kappa = 2/3 \quad \text{for scaled Bjorken expansion}$$

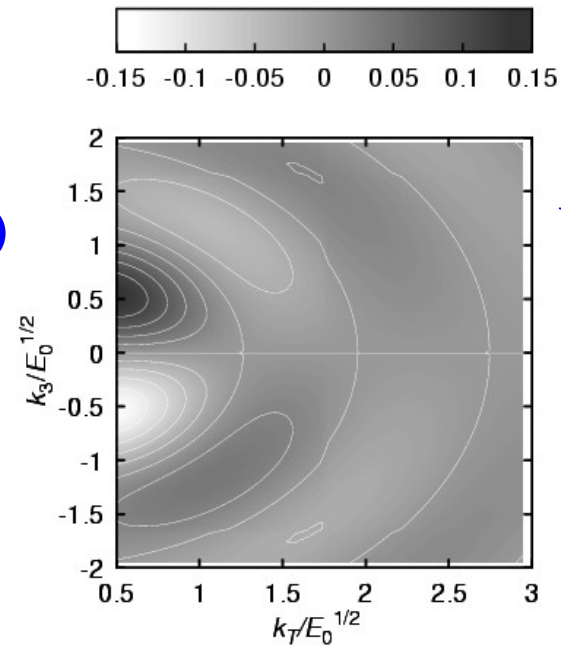
$$\text{with } t_0 = 0.01 / E_0^{1/2}$$

Numerical results (b^i) for the Bjorken expansion at $t=2/\sqrt{E_0}$ in $SU(2)$:

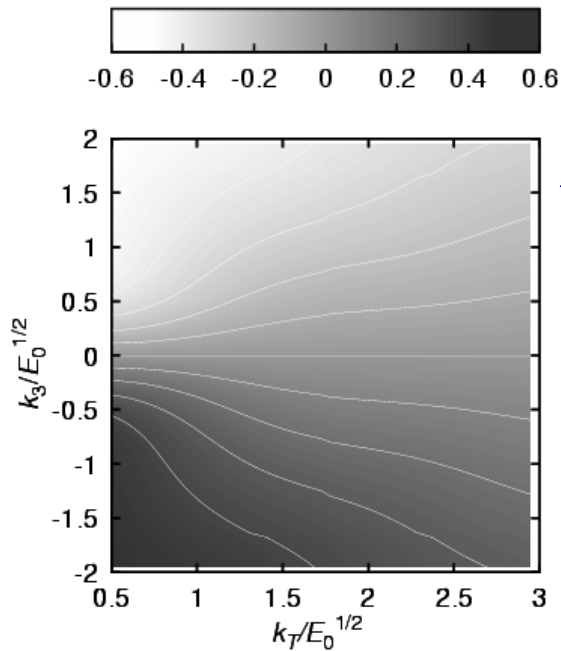
$m = 0$



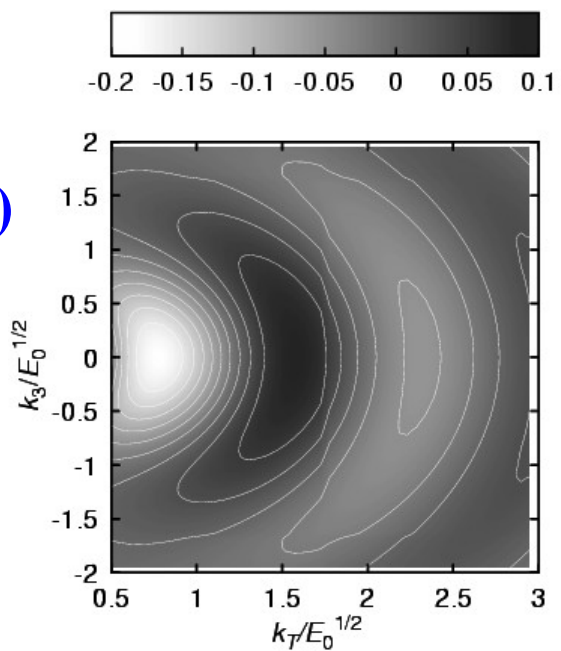
$b^s_T(k_T, k_3)$



$b^a_T(k_T, k_3)$

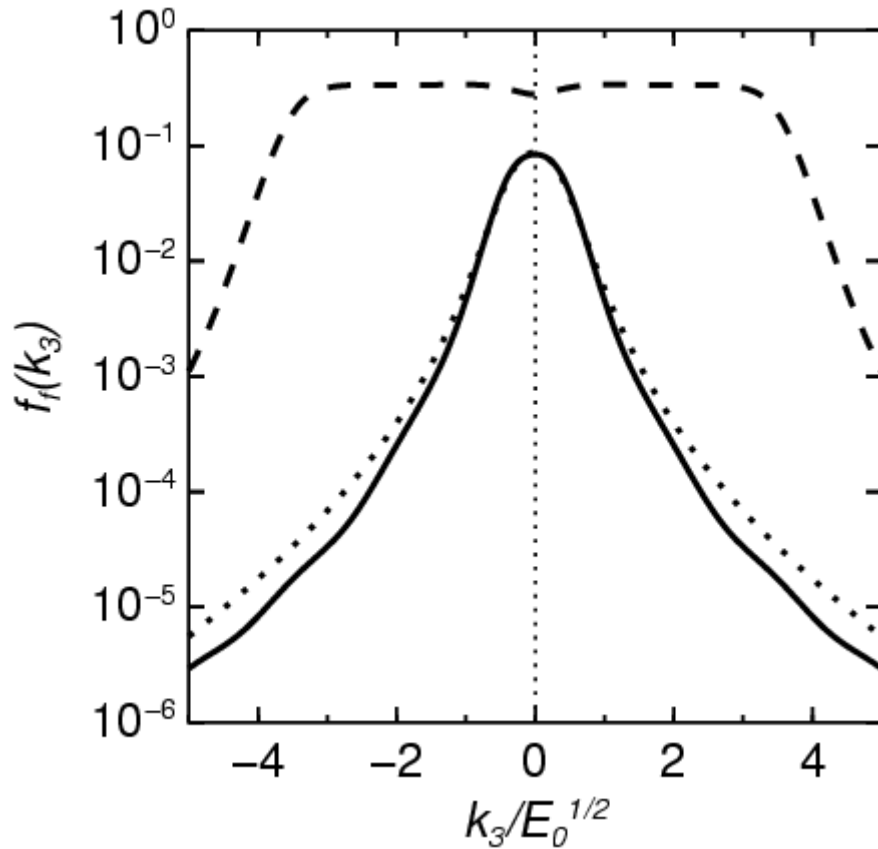


$b^s_3(k_T, k_3)$



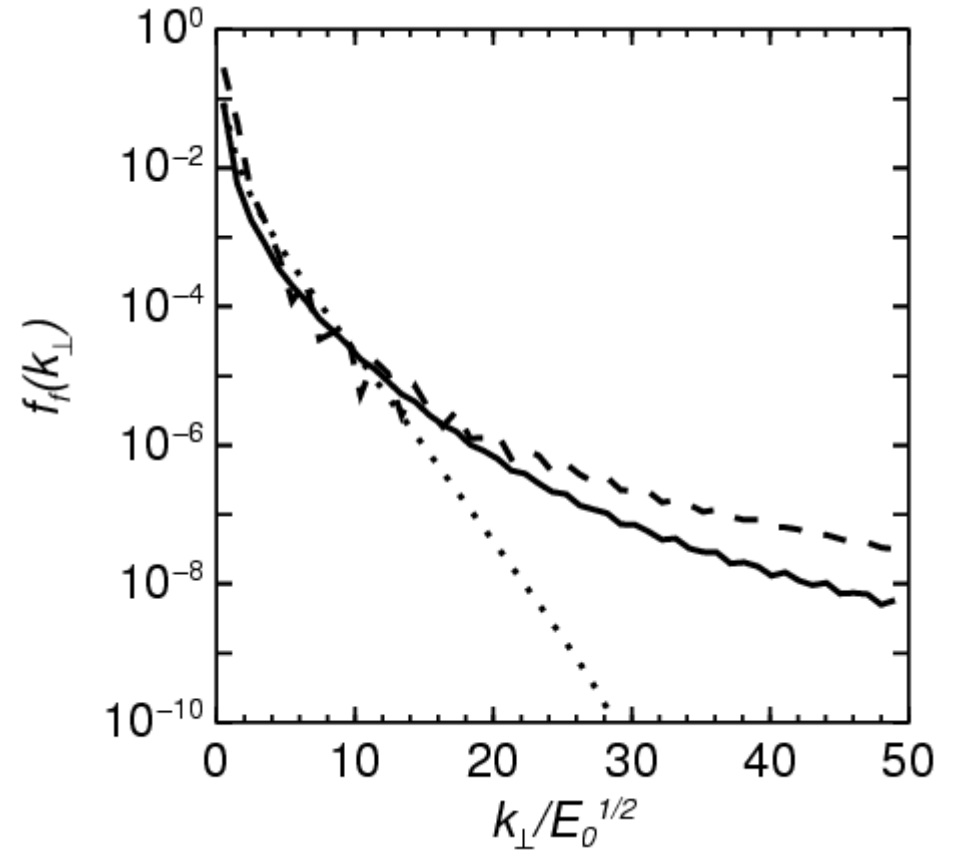
$b^a_3(k_T, k_3)$

Numerical results for fermion distributions at $t=2/\sqrt{E_0}$ in $SU(2)$:



$f_f(k_3)$: longitudinal mom. distr.

$k_T/\sqrt{E_0} = 0.5$



$f_f(k_T)$: transv. mom. distr.

$k_3 = 0$

\Rightarrow exponential (pulse)

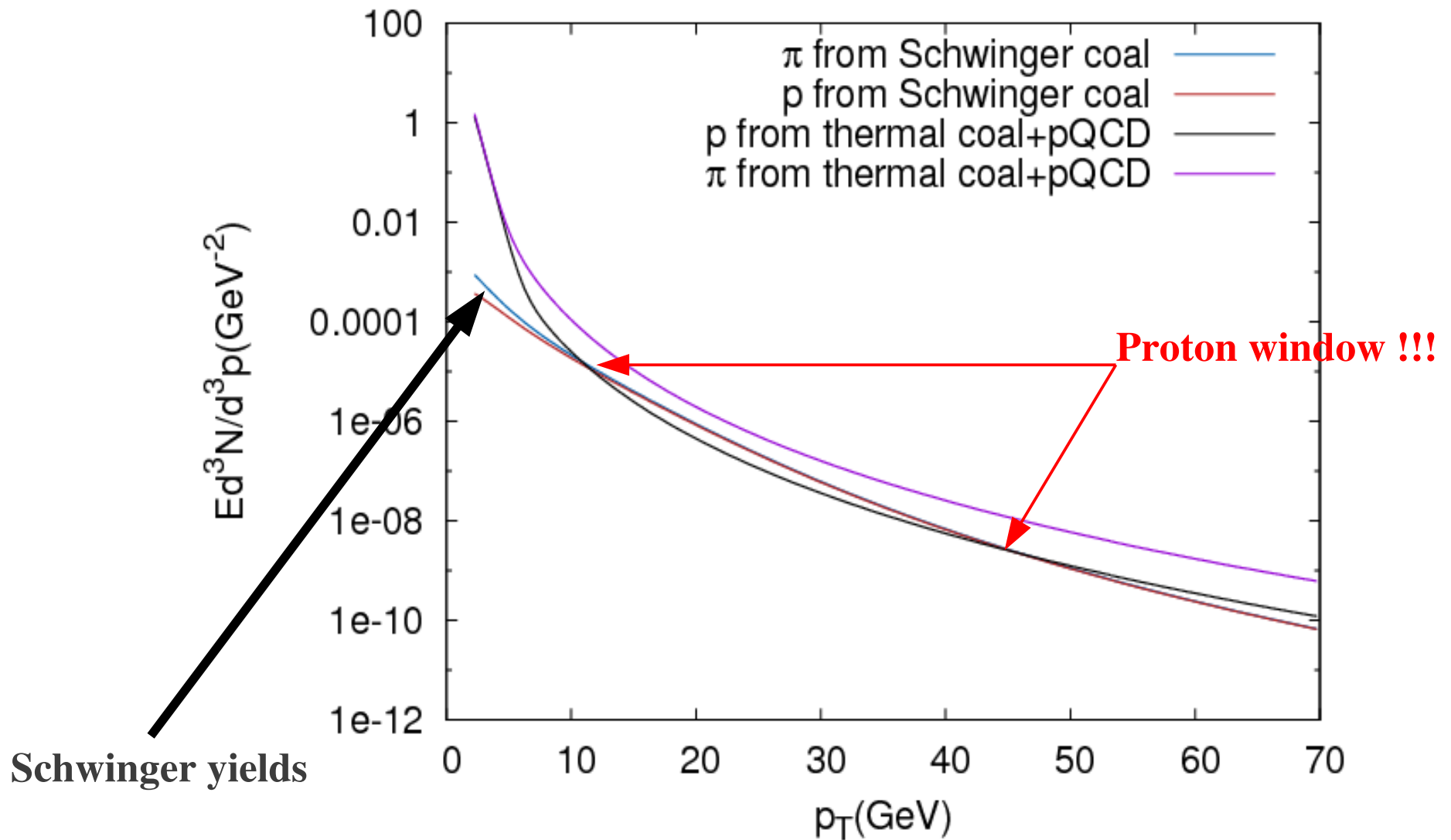
\Rightarrow polinomial (scaled)

Typical result: Proton and pion production from

jet fragmentation (FF)

coalescence of thermal quarks

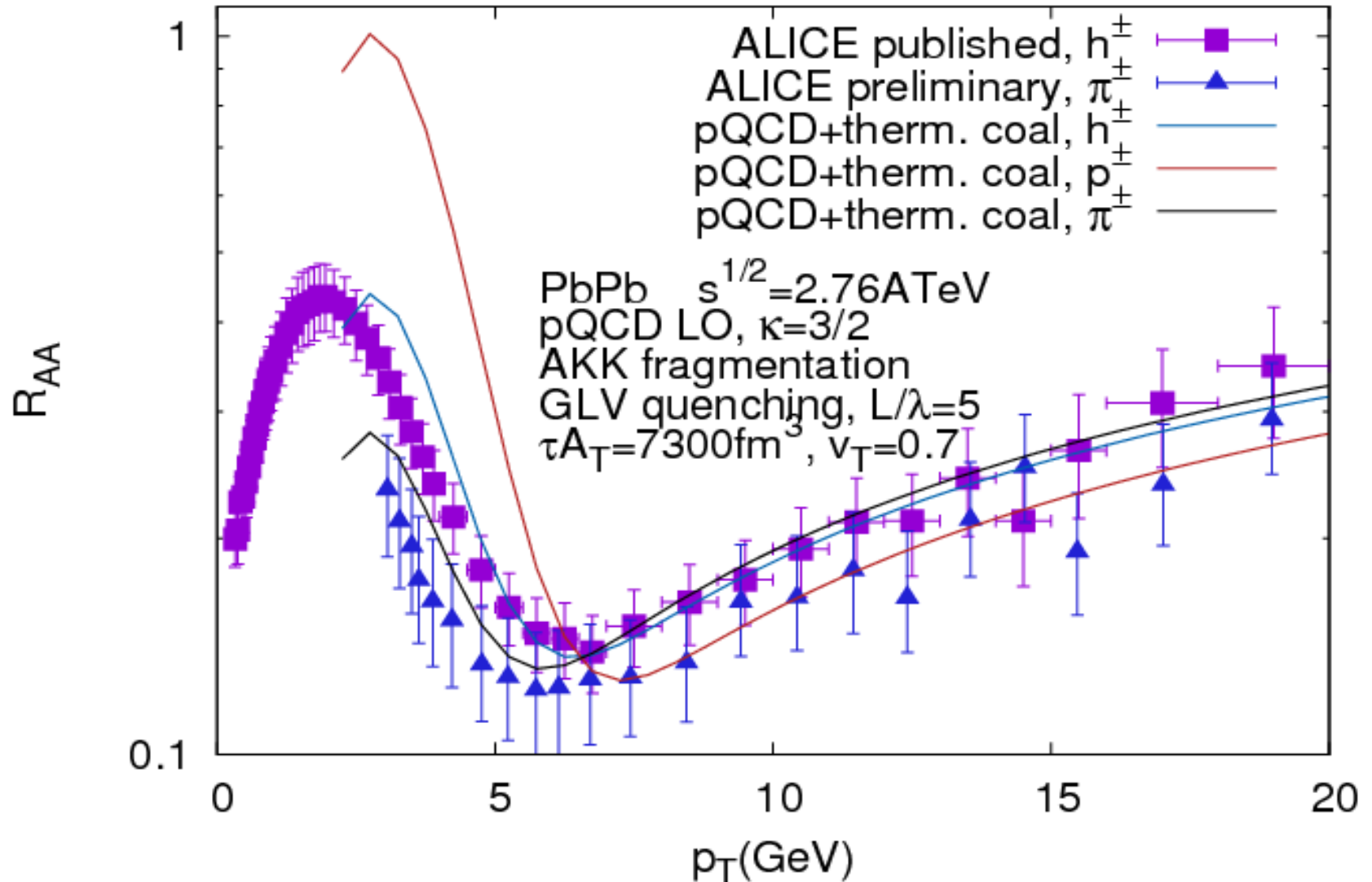
coalescence of Schwinger quarks/diquarks



Latest calculations on charged hadron R_{AA} in Pb+Pb at 2.76 ATeV

pQCD + Quenching + Thermal coalesc.

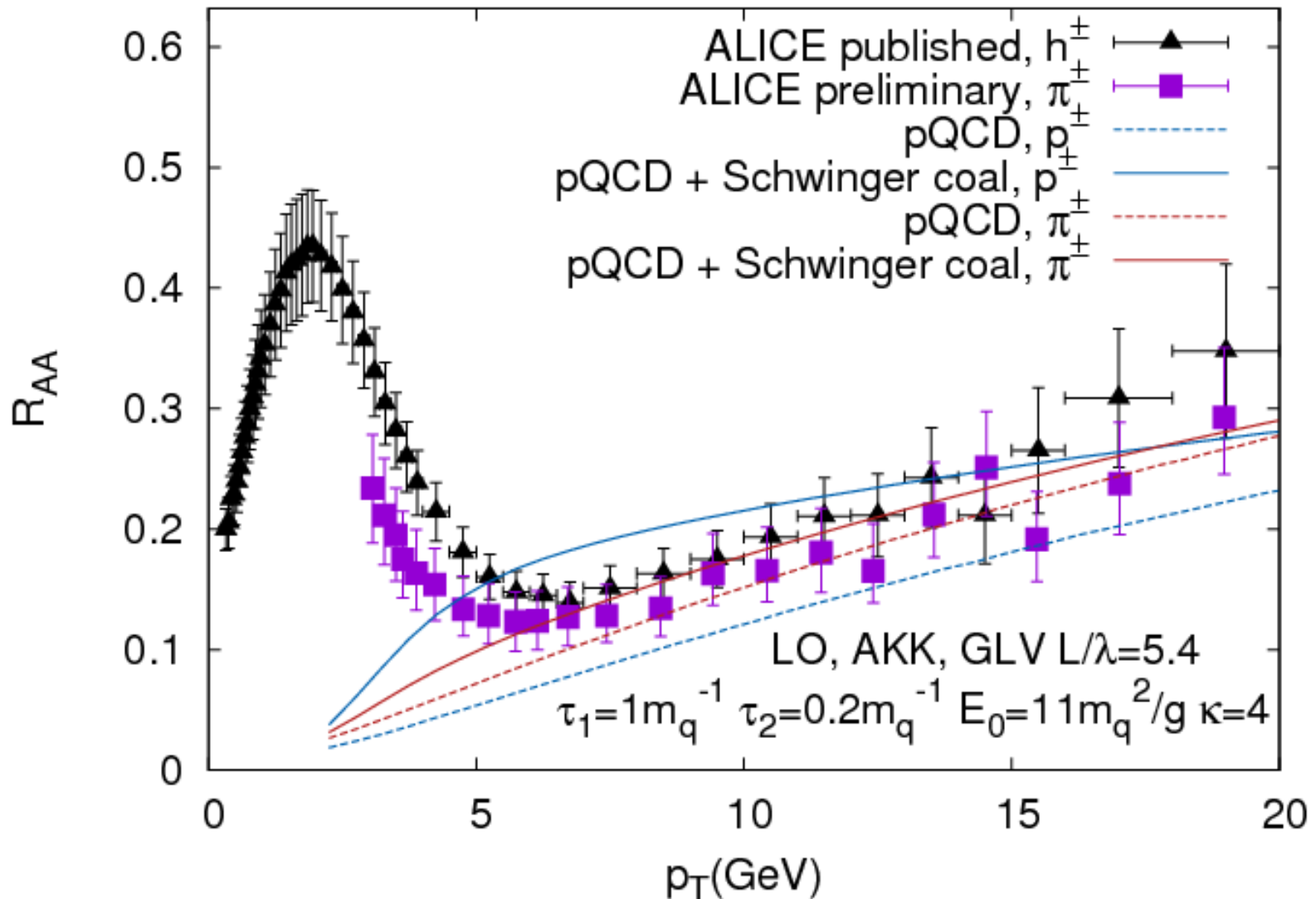
P. Lévai et al. 2011



Latest calculations on charged hadron R_{AA} in Pb+Pb at 2.76 ATeV

pQCD + Quenching + Schwinger diquark coalesc.

P. Lévai et al. 2012

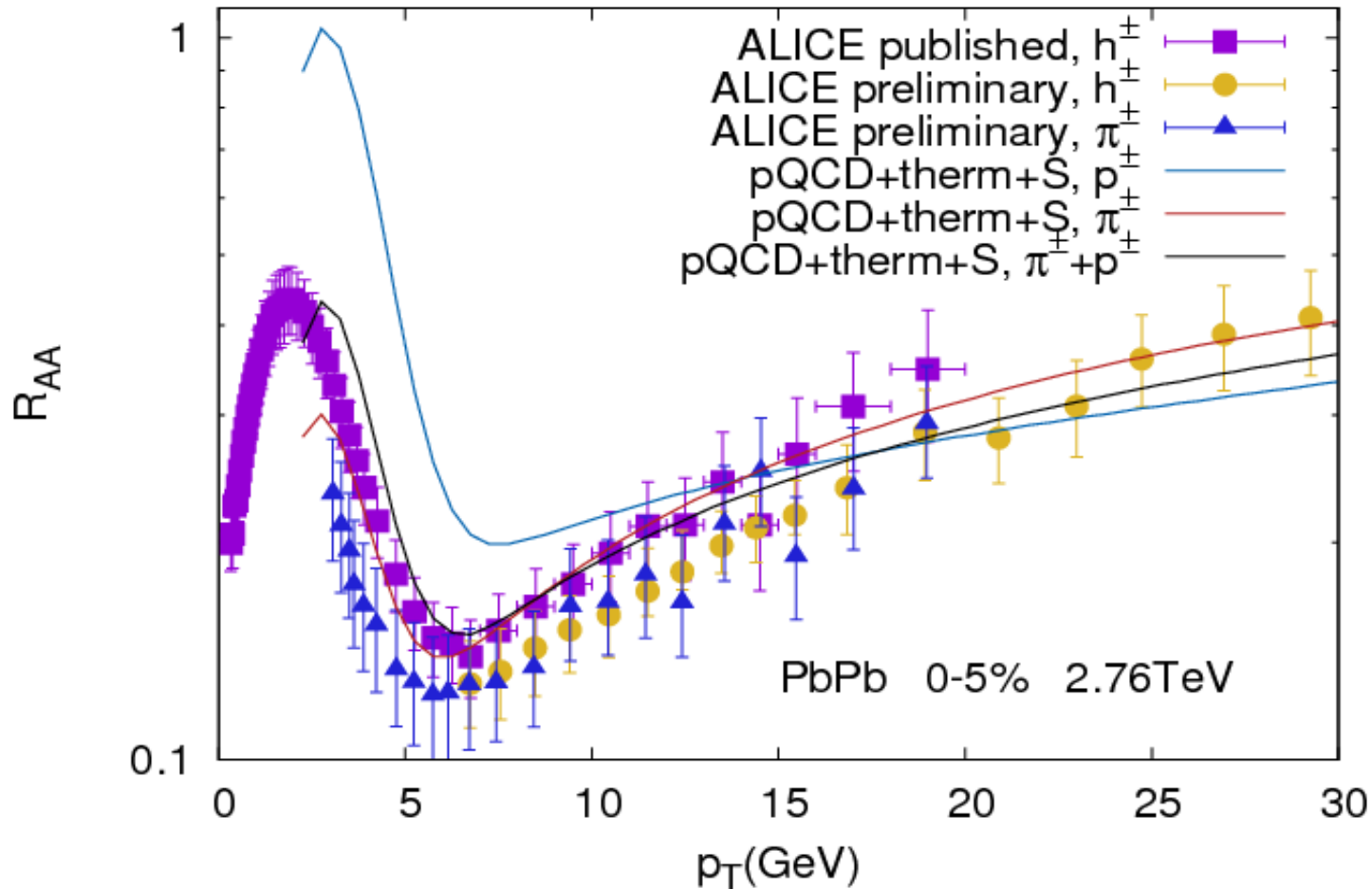


Latest calculations on charged hadron R_{AA} in Pb+Pb at 2.76 ATeV

pQCD + Quenching + (Thermal + Schwinger diquark) coalesc.

P. Lévai et al. 2012

ArXiv: 1208.0448



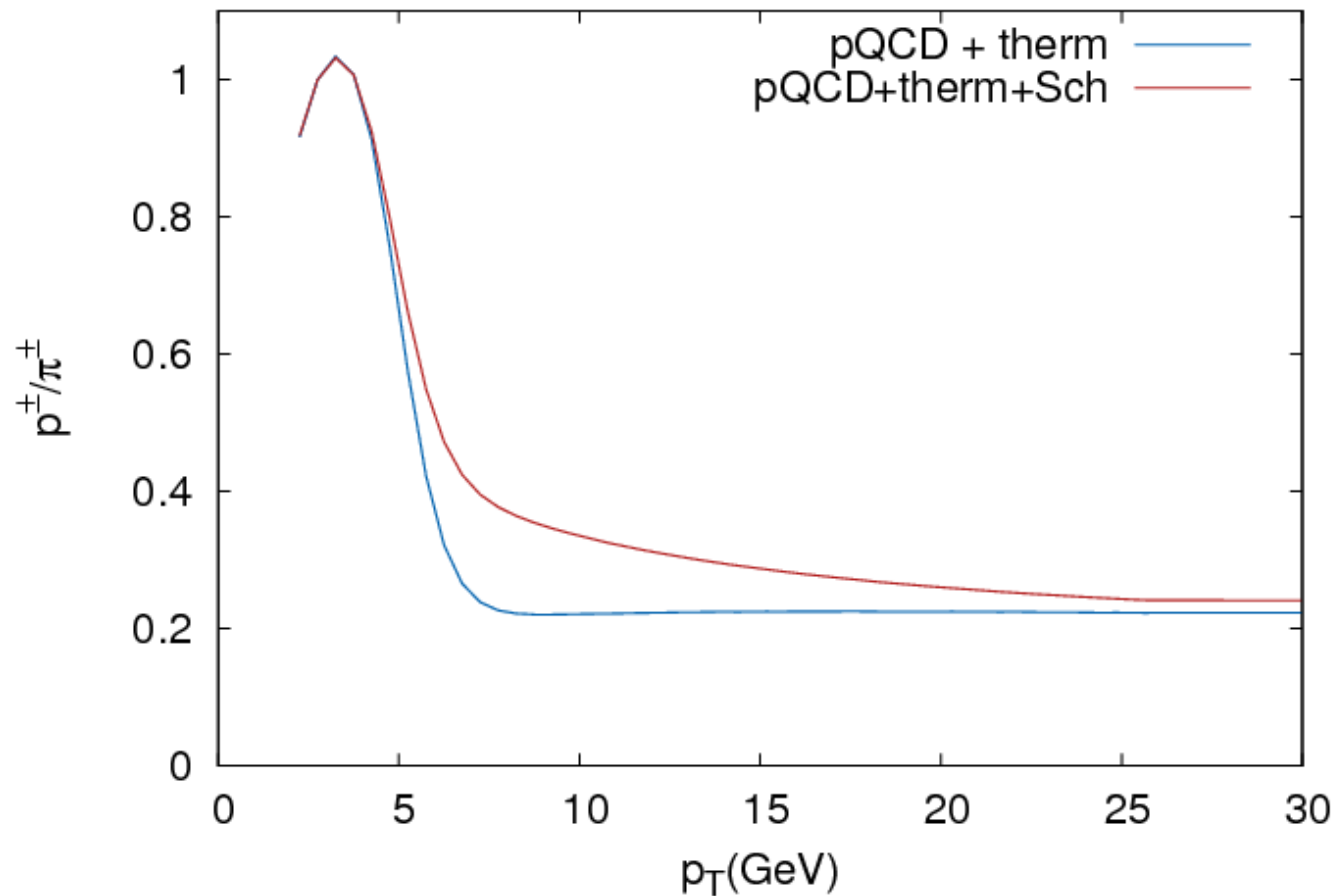
$R_{AA}(p) > R_{AA}(\text{pion})$ for $2 < p_T < 15 \text{ GeV}/c$

Latest calculations on proton/pion ratio in Pb+Pb at 2.76 ATeV

pQCD + Quenching + (Thermal + Schwinger diquark) coalesc.

P. Lévai et al. 2012

ArXiv: 1208.0448



**Extra proton yield at intermediate- and high- p_T
In the window: $5 < p_T < 20$ GeV/c**

Conclusions:

**1. Jet energy loss is large at RHIC energies
and even larger at LHC energies.**

Since R_{pPb} approx 1, then jet energy loss is needed.

In the window $5 < pT < 100$ GeV seems to be all right.

**2. Gluon energy loss is much larger, gluons will be quenched
at LHC energies in the window $5 < pT < 20$ GeV.**

**3. Showers can create a thermal bath (QGP)
or melting into a coherent gluon field**

Which may appear in the early phase of the AA collision

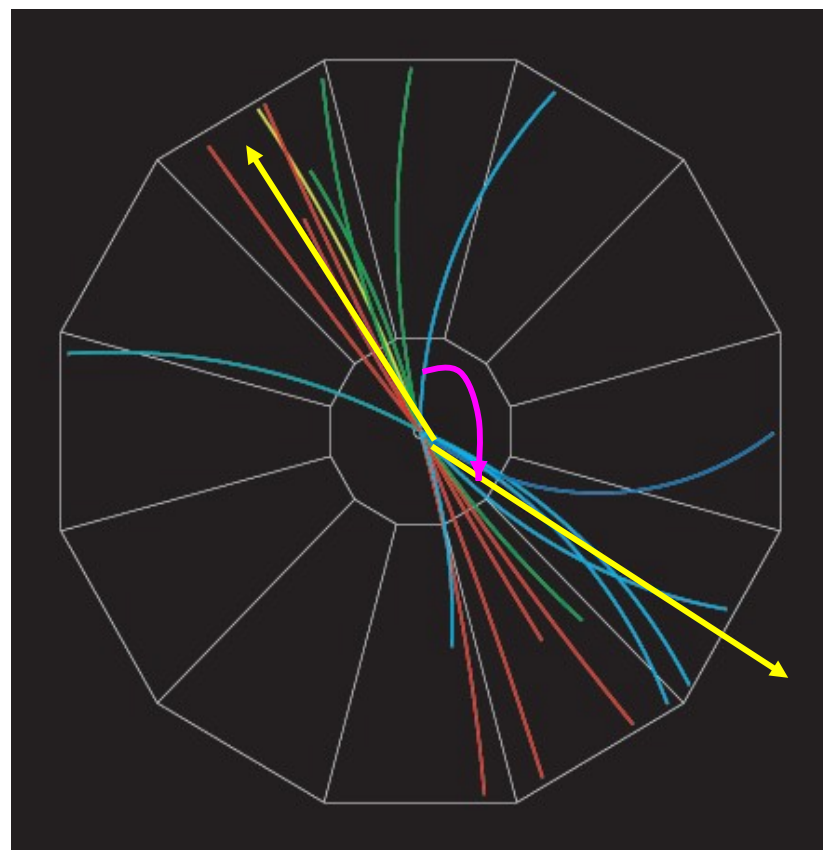
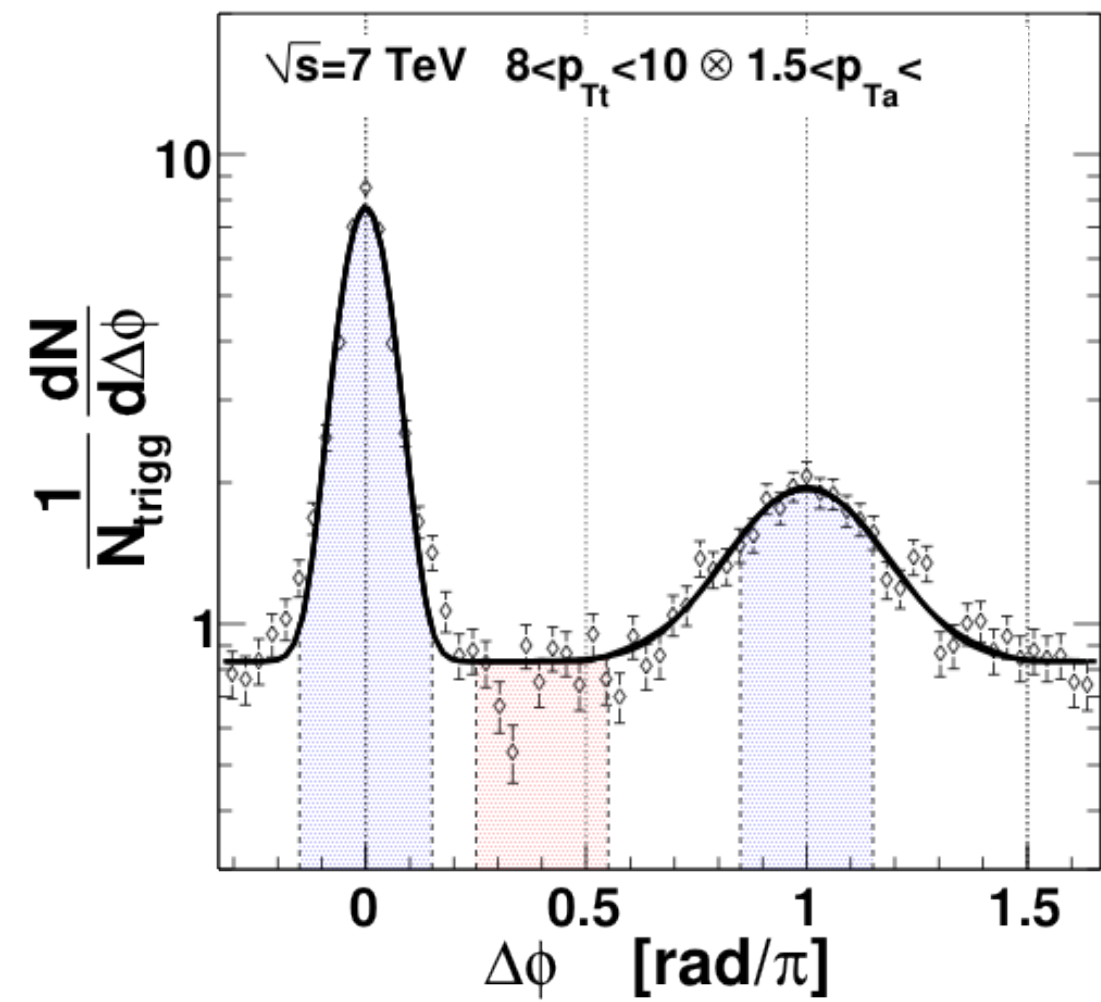
**4. Baryon yields and baryon/meson ratios can decide
if coherent field appears or not.**

(Extra diquark-antidiquark production and coalescence.)

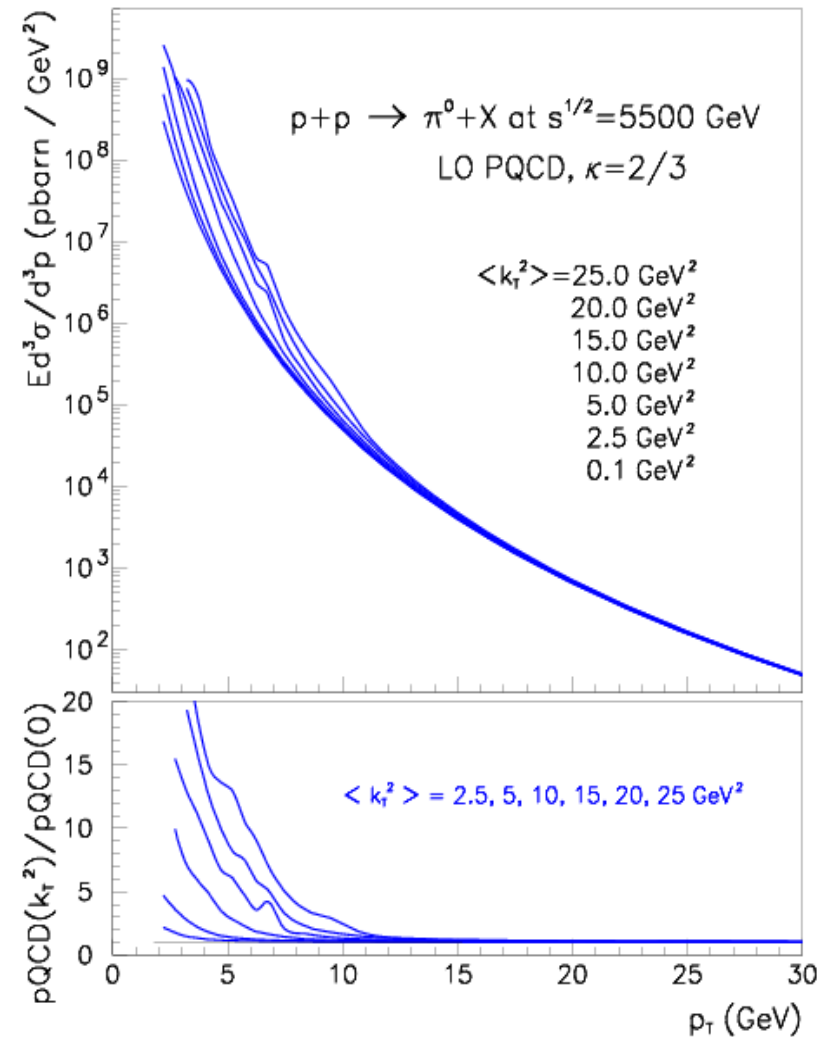
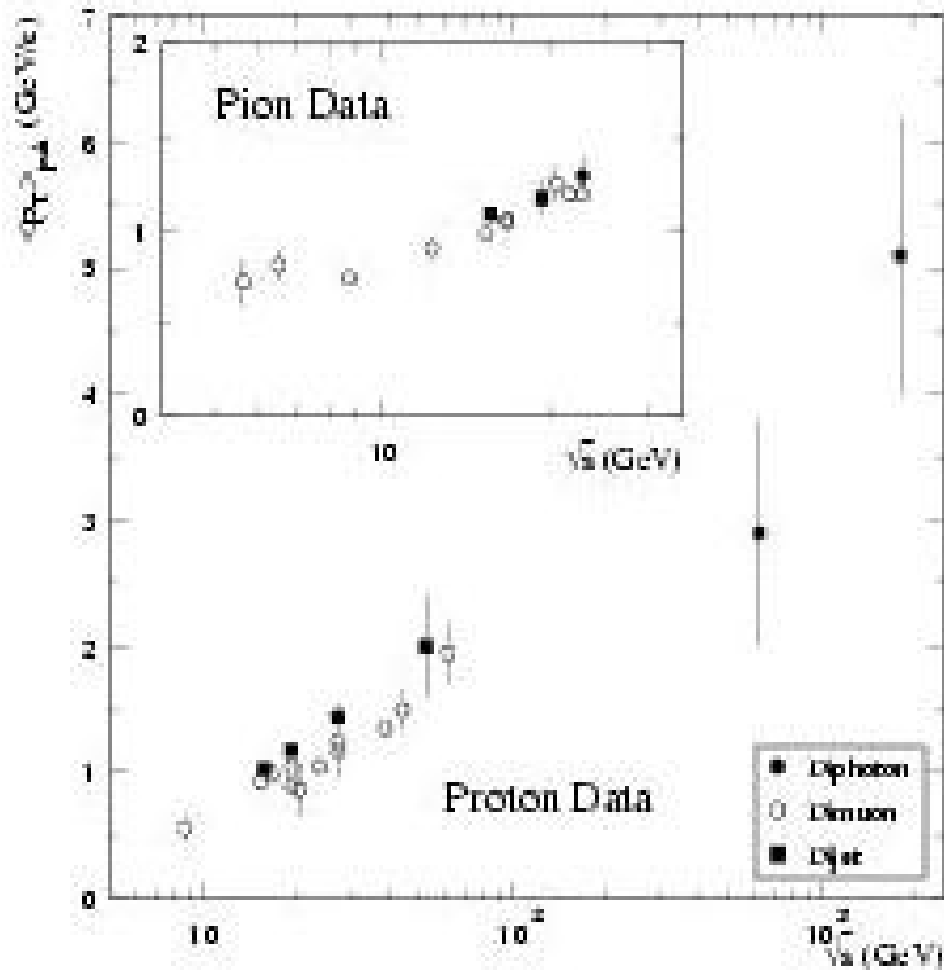
Back-up slides on 2-particle correlations:

where is the kT -imbalance in the 1-particle spectra?????

Why LO pQCD is working so nicely without intrinsic- kT ???



k_T -imbalance parameter --- extracted from 2-hadron correlation and applied in 1-particle distribution



But no room for intrinsic- k_T in CDF and CMS data !! ???