

# Non-extensive Boltzmann equation & thermalization

Miklós Horváth (PhD student)  
Supervisor: Tamás Biró



Zimányi 2012

# Content

---

- Non-extensivity & kinetic equations
- Motivations
- Toy-model (colliding point-particles)
  - Details of kinetics
  - Detailed balance no-go
  - Numerical results about long-time evolution

Goal:

Long-time behaviour of two-component **modified** kinetic equations:



**NESS (Non-equilibrium steady state)**

# Non-Extensivity

- Why can NOT we just sum up micro-/mesoscopic quantities?



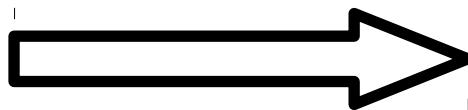
- Pair-interaction in general:

$$\frac{E}{N} = \frac{1}{N} \sum_i E_i + \frac{1}{N} \sum_{ij} V_{ij}$$

- in TDL:

# Non-Extensivity

- Why can NOT we just sum up micro-/mesoscopic quantities?



**INTERACTION**

- Pair-interaction in general:

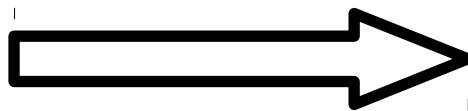
$$\frac{E}{N} = \frac{1}{N} \sum_i E_i + \frac{1}{N} \sum_{ij} V_{ij}$$

- in TDL:

small: weak interaction (few pairs)

# Non-Extensivity

- Why can NOT we just sum up micro-/mesoscopic quantities?



**INTERACTION**

- Pair-interaction in general:

$$\frac{E}{N} = \frac{1}{N} \sum_i E_i + \frac{1}{N} \sum_{ij} V_{ij}$$

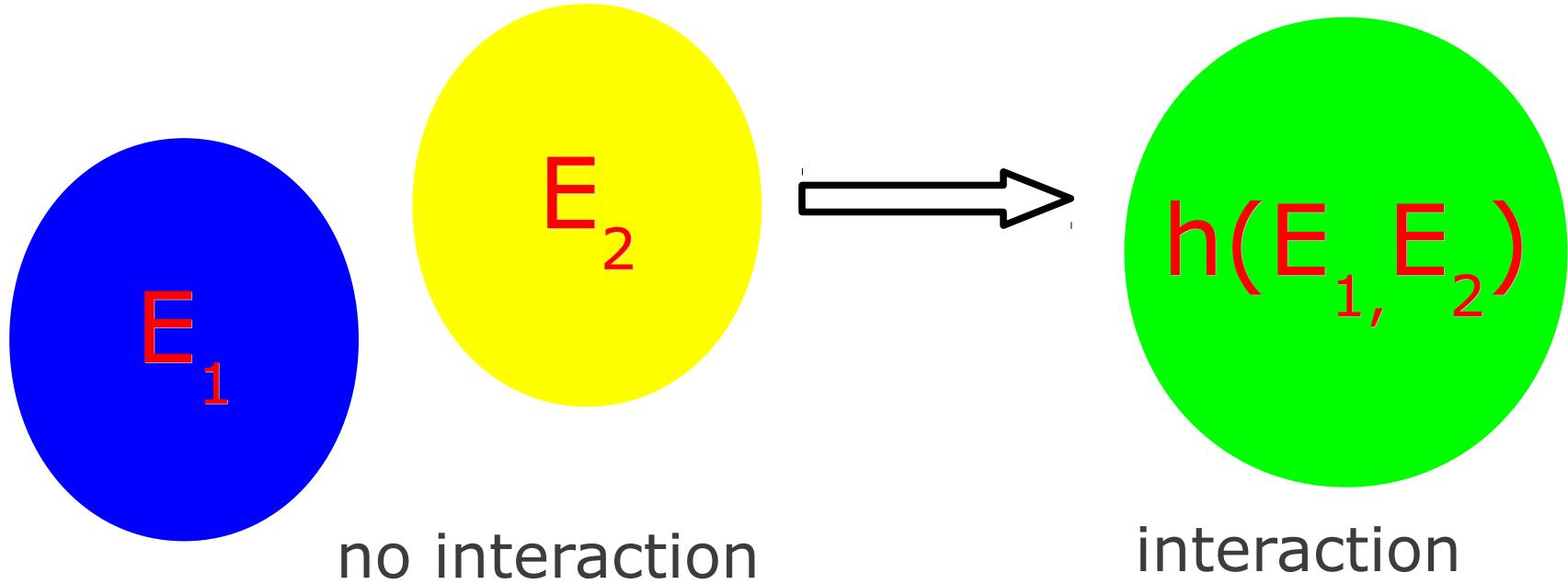
- in TDL:

small: weak interaction (few pairs)

$\sim N^{\delta-1}$

big: strong interaction (a lot of pairs)

# Non-Extensivity



**Trick to use additive quantities**

(if  $h$  is associative):

$$h(E, E') = K \Rightarrow L(E) + L(E') = L(K) = \text{const.}$$

$$L(E) = \int_0^E \frac{d\epsilon}{\partial_2 h(\epsilon, 0)}$$

# Kinetic equations

---

$$\dot{f}_1^\alpha = \sum_\beta \iiint_{234} \mathcal{W}_{1234}^{\alpha\beta} (f_3^\alpha f_4^\beta - f_1^\alpha f_2^\beta)$$

- symmetries:  $1234 \leftrightarrow 3412, 1234 \leftrightarrow 2134$   
 $\alpha\beta \leftrightarrow \beta\alpha$
- constraints: 2-p. (quasi-)energy, 2-p. total momentum
- conserving quantities:  
total momentum, number of particles

$$\mathcal{W}_{1234}^{\alpha\beta} = \omega_{1234}^{\alpha\beta} \delta(h(E_1, E_2) - h(E_3, E_4)) \delta^{(3)}(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{p}_3 - \mathbf{p}_4)$$

# Kinetic equations

$$\dot{f}_1^\alpha = \sum_\beta \iiint_{234} \mathcal{W}_{1234}^{\alpha\beta} (f_3^\alpha f_4^\beta - f_1^\alpha f_2^\beta)$$

- symmetries:  $1234 \leftrightarrow 3412, 1234 \leftrightarrow 2134$   
 $\alpha\beta \leftrightarrow \beta\alpha$
- constraints: 2-p. (quasi-)energy, 2-p. total momentum  
$$\mathcal{W}_{1234}^{\alpha\beta} = \omega_{1234}^{\alpha\beta} \delta(h(E_1, E_2) - h(E_3, E_4)) \delta^{(3)}(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{p}_3 - \mathbf{p}_4)$$
- conserving quantities:  
total momentum, number of particles

# Kinetic equations

$$\dot{f}_1^\alpha = \sum_\beta \iiint_{234} \mathcal{W}_{1234}^{\alpha\beta} (f_3^\alpha f_4^\beta - f_1^\alpha f_2^\beta)$$

- **symmetries:**  $1234 \leftrightarrow 3412, 1234 \leftrightarrow 2134$   
 $\alpha\beta \leftrightarrow \beta\alpha$
- **constraints:** 2-p. (quasi-)energy, 2-p. total momentum  
$$\mathcal{W}_{1234}^{\alpha\beta} = \omega_{1234}^{\alpha\beta} \delta(h(E_1, E_2) - h(E_3, E_4)) \delta^{(3)}(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{p}_3 - \mathbf{p}_4)$$
- conserving quantities:
  - total momentum, number of particles

# Kinetic equations & entropy growth

$$S = - \sum_{\alpha} \int_1 f_1^{\alpha} \ln f_1^{\alpha}$$

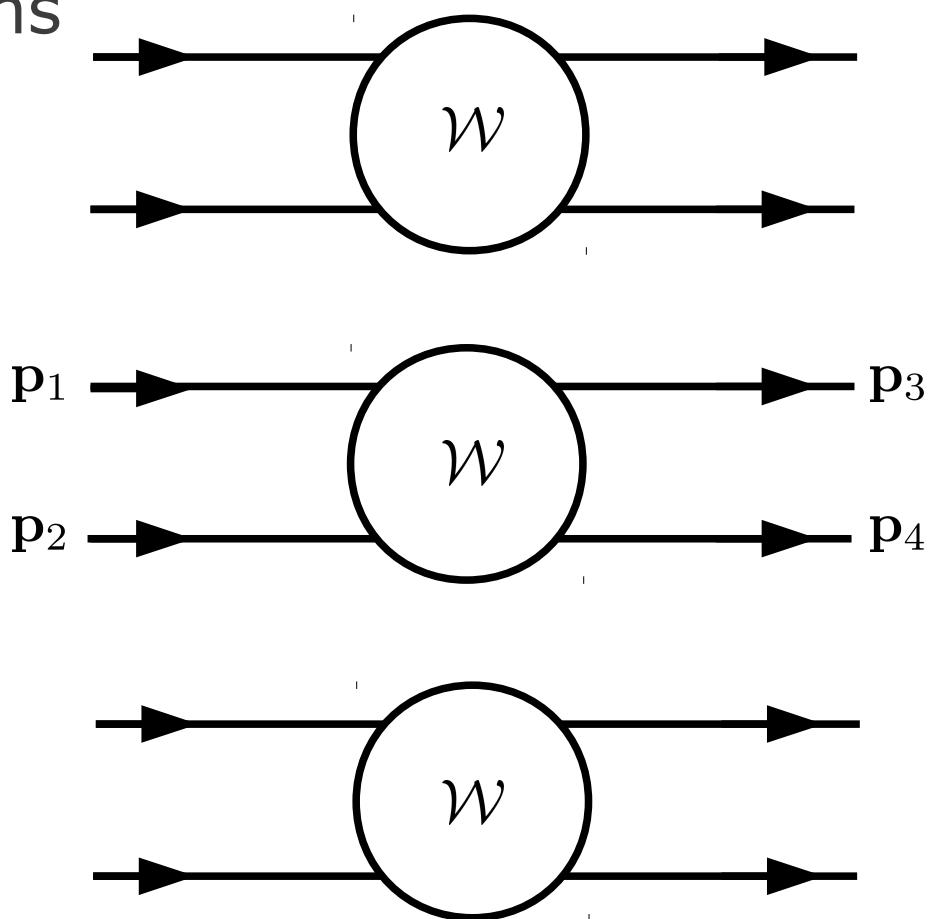
- H-theorem:  $S$  is increasing monotonically
- approaches its maximum in detailed balance (if exists):

$$f^{\alpha}(E_1)f^{\beta}(E_2) = f^{\alpha}(E_3)f^{\beta}(E_4)$$

$$L^{\alpha\beta}(E_1) + L^{\alpha\beta}(E_2) = L^{\alpha\beta}(E_3) + L^{\alpha\beta}(E_4)$$

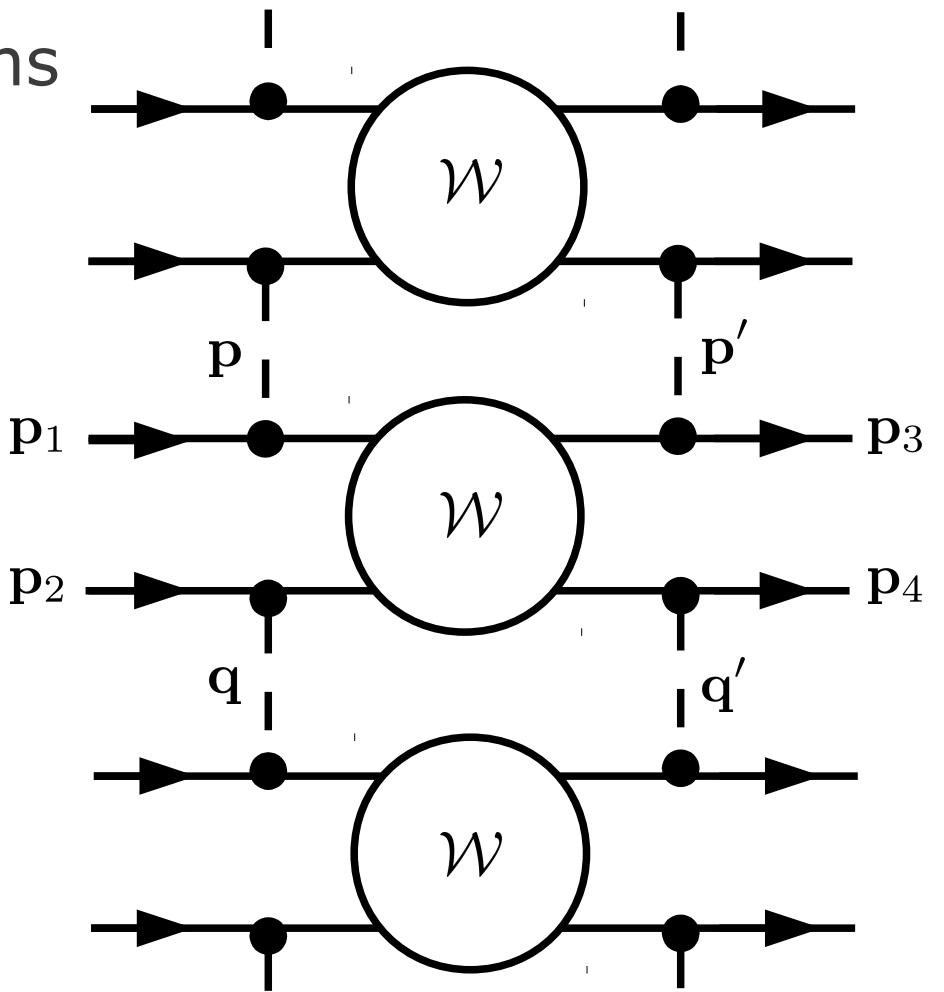
# Motivations

- power law tails in spectra
- not independent collisions in dense plasma
- modified dispersion relation



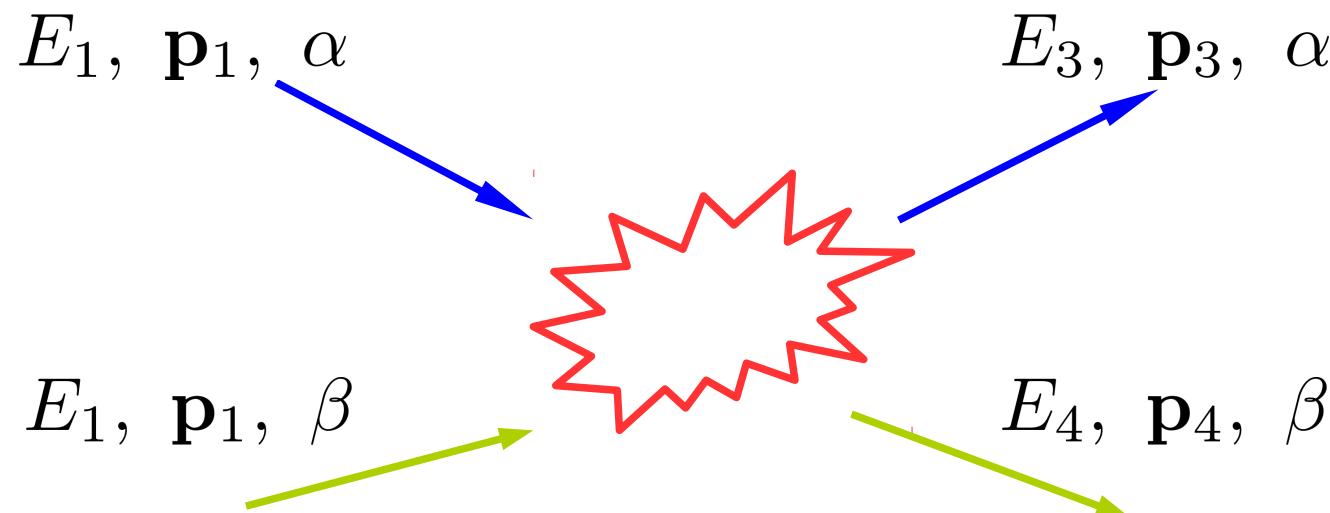
# Motivations

- power law tails in spectra
- not independent collisions in dense plasma
- modified dispersion relation



# A toy-model – construction

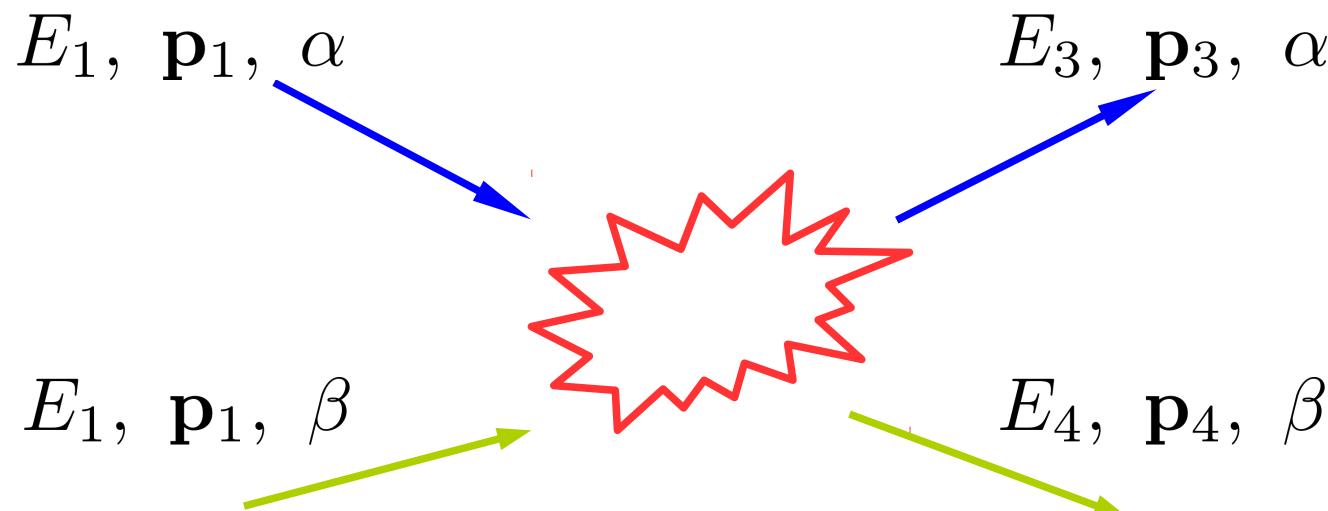
- particles  $\{E_i, \mathbf{p}_i, \alpha_i\}$
- interaction via binary collisions
- dispersion relation:  $E = c|\mathbf{p}|^n$
- collisional invariants:  $\mathbf{p}_1 + \mathbf{p}_2 = \mathbf{p}_3 + \mathbf{p}_4$   
 $h(E_1, E_2) = h(E_3, E_4)$



# A toy-model – construction

- particles  $\{E_i, \mathbf{p}_i, \alpha_i\}$
- interaction via binary collisions
- dispersion relation:  $E = c|\mathbf{p}|^n$

- collisional invariants:  
$$\begin{aligned} \mathbf{p}_1 + \mathbf{p}_2 &= \mathbf{p}_3 + \mathbf{p}_4 \\ E_1 + E_2 + a^{\alpha\beta} E_1 E_2 &= \\ &= E_3 + E_4 + a^{\alpha\beta} E_3 E_4 \end{aligned}$$



# A toy-model – kinetic equation

---

continuum limit:  $(N \rightarrow \infty, \Delta(E, E', \epsilon) \rightarrow 0)$

$$\dot{f}_1^\alpha = \sum_\beta \iiint_{234} \mathcal{W}_{1234}^{\alpha\beta} (f_3^\alpha f_4^\beta - f_1^\alpha f_2^\beta)$$

(Boltzmann-type) kinetic equation

# A toy-model – kinetic equation

continuum limit:  $(N \rightarrow \infty, \Delta(E, E', \epsilon) \rightarrow 0)$

$$\dot{f}_1^\alpha = \sum_\beta \iiint_{234} \mathcal{W}_{1234}^{\alpha\beta} (f_3^\alpha f_4^\beta - f_1^\alpha f_2^\beta)$$

(Boltzmann-type) kinetic equation

$$\mathcal{W}_{1234}^{\alpha\beta} = \rho^{\alpha\beta}(E_3, E_4, K) \underbrace{\delta(h^{\alpha\beta}(E_1, E_2) - h^{\alpha\beta}(E_3, E_4)) \delta^{(3)}(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{p}_3 - \mathbf{p}_4)}_K$$



$$\dot{f}^\alpha(E) = \sum_\beta \int_0^\Lambda dE' \int_0^K d\epsilon \left( f^\alpha(\epsilon) f^\beta(K \ominus \epsilon) - f^\alpha(E) f^\beta(E') \right) \rho^{\alpha\beta}$$

# A toy-model – kinetic equation

continuum limit:  $(N \rightarrow \infty, \Delta(E, E', \epsilon) \rightarrow 0)$

$$\dot{f}_1^\alpha = \sum_\beta \iiint_{234} \mathcal{W}_{1234}^{\alpha\beta} (f_3^\alpha f_4^\beta - f_1^\alpha f_2^\beta)$$

(Boltzmann-type) kinetic equation

$$\mathcal{W}_{1234}^{\alpha\beta} = \rho^{\alpha\beta}(E_3, E_4, K) \delta(h^{\alpha\beta}(E_1, E_2) - h^{\alpha\beta}(E_3, E_4)) \delta^{(3)}(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{p}_3 - \mathbf{p}_4)$$

**modified**

$K$



$$\dot{f}^\alpha(E) = \sum_\beta \int_0^\Lambda dE' \int_0^K d\epsilon \left( f^\alpha(\epsilon) f^\beta(K \ominus \epsilon) - f^\alpha(E) f^\beta(E') \right) \rho^{\alpha\beta}$$

# A toy-model – simulation

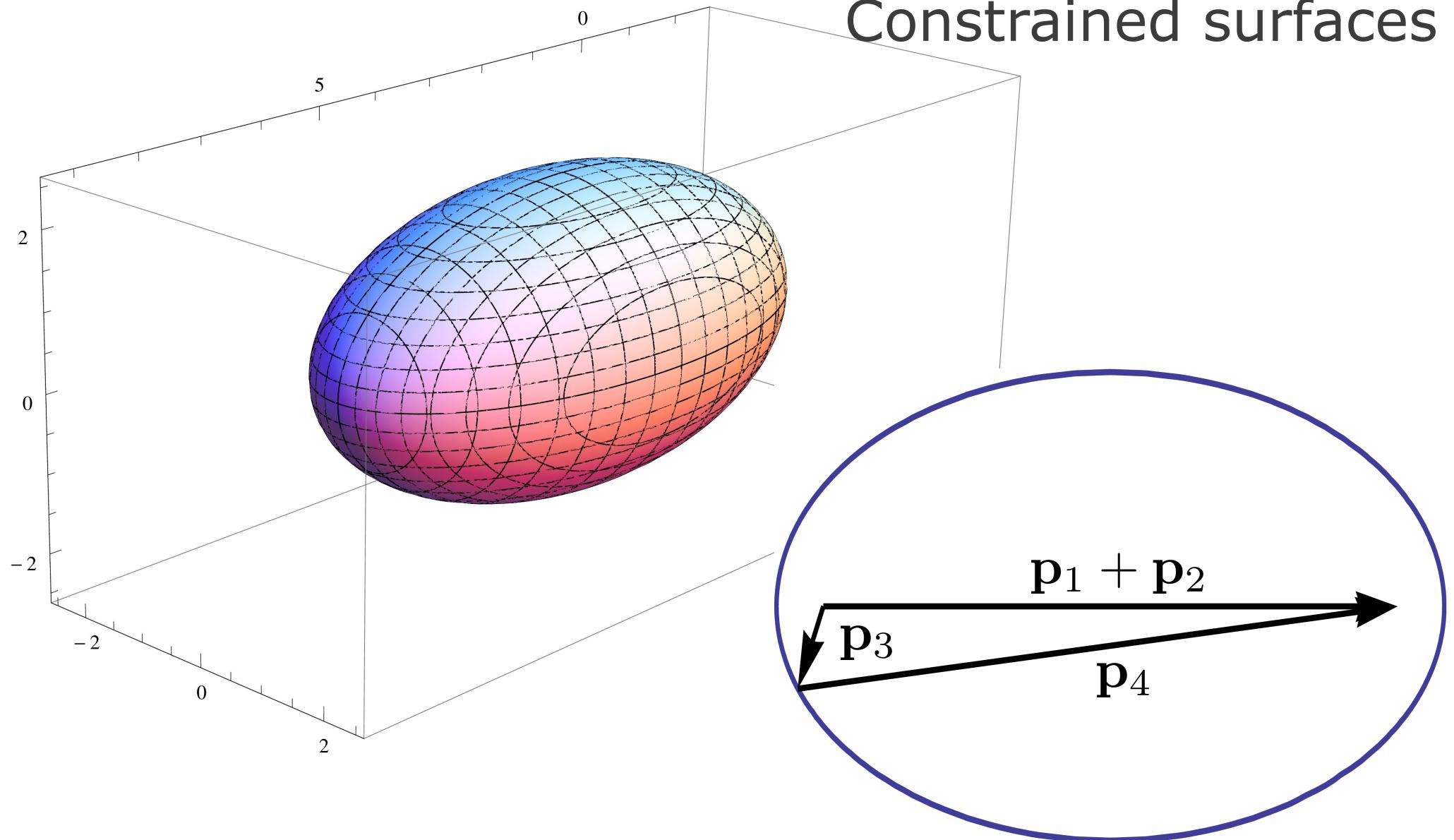
---

## Numerical implementation:

- ensemble described by  $\{(\mathbf{p}_i, E_i = c|\mathbf{p}|^n), i = 1 \dots N\}$
- collision-by-collision evolution, NOT in real time
- random sampling on the constrained phase space
- random variables:
  - 1 outgoing energy
  - angle between incoming & outgoing planes

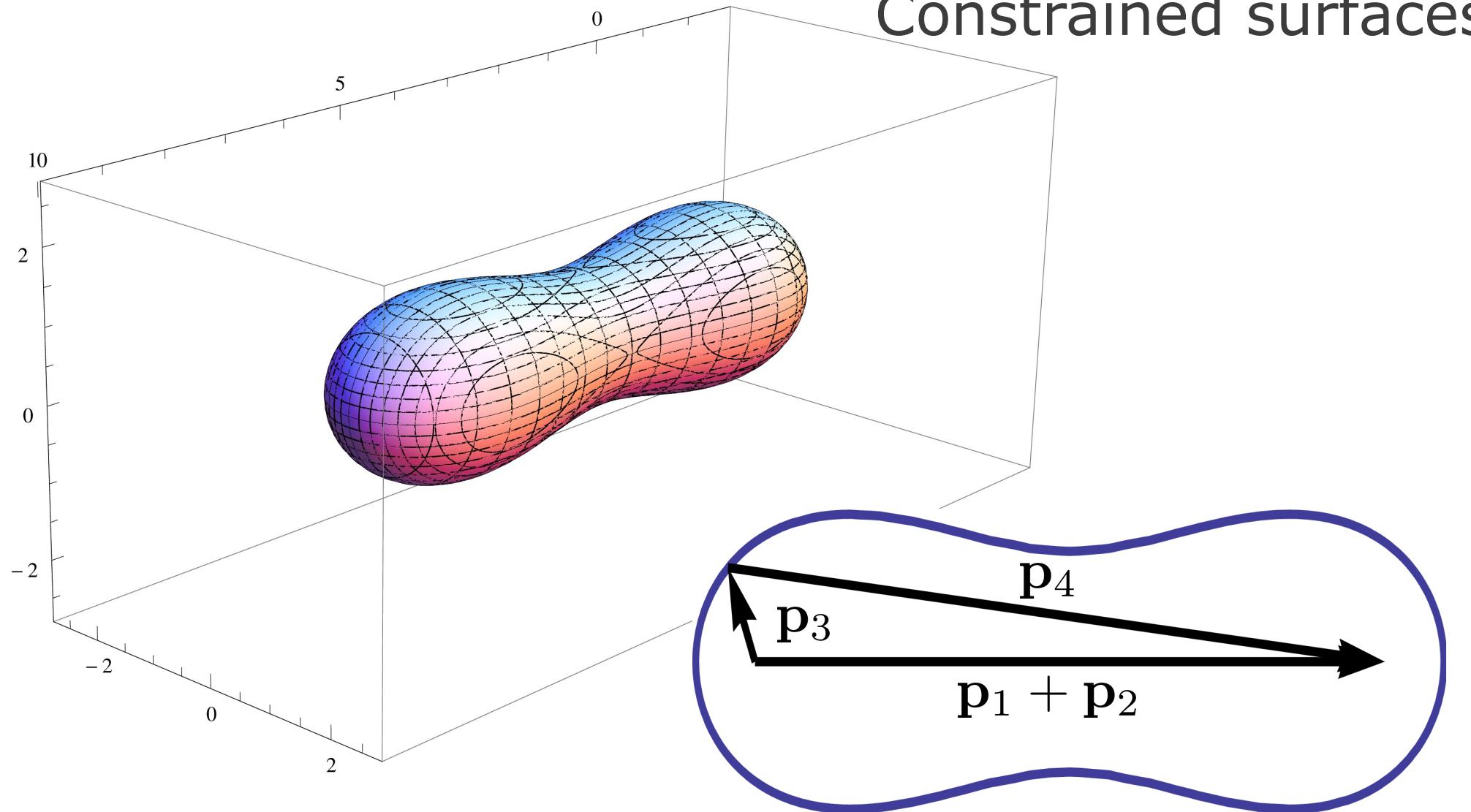
# A toy-model – simulation

Constrained surfaces



# A toy-model – simulation

Constrained surfaces



# A toy-model – detailed balance



$$\dot{f}_1^\alpha = \sum_\beta \iiint_{234} \underbrace{\mathcal{W}_{1234}^{\alpha\beta} (f_3^\alpha f_4^\beta - f_1^\alpha f_2^\beta)}_{=0}$$

$$f^\alpha(E_1) f^\beta(E_2) = f^\alpha(E_3) f^\beta(E_4)$$

$$h^{\alpha\beta}(E_1, E_2) = h^{\alpha\beta}(E_3, E_4)$$

# A toy-model – detailed balance



$$\dot{f}_1^\alpha = \sum_\beta \iiint_{234} \underbrace{\mathcal{W}_{1234}^{\alpha\beta} (f_3^\alpha f_4^\beta - f_1^\alpha f_2^\beta)}_{=0}$$

$$f^\alpha(E_1)f^\beta(E_2) = f^\alpha(E_3)f^\beta(E_4)$$

$$L^{\alpha\beta}(E_1) + L^{\alpha\beta}(E_2) = L^{\alpha\beta}(E_3) + L^{\alpha\beta}(E_4)$$

$n+n(n-1)/2$  independent equations

# A toy-model – detailed balance

for  $\alpha = \beta$

$$f^\alpha(E_1)f^\alpha(E_2) = f^\alpha(E_3)f^\alpha(E_4)$$

$$L^{\alpha\alpha}(E_1) + L^{\alpha\alpha}(E_2) = L^{\alpha\alpha}(E_3) + L^{\alpha\alpha}(E_4)$$



$$f^\alpha \sim e^{-\gamma_{\alpha\alpha} L^{\alpha\alpha}(E)}$$

# A toy-model – detailed balance

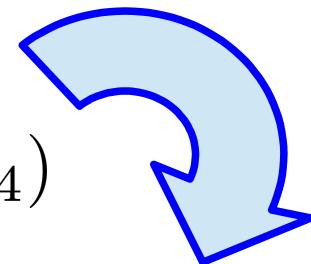
for  $\alpha = \beta$

$$f^\alpha(E_1)f^\alpha(E_2) = f^\alpha(E_3)f^\alpha(E_4)$$

$$L^{\alpha\alpha}(E_1) + L^{\alpha\alpha}(E_2) = L^{\alpha\alpha}(E_3) + L^{\alpha\alpha}(E_4)$$



$$f^\alpha \sim e^{-\gamma_{\alpha\alpha} L^{\alpha\alpha}(E)}$$



$$f^\alpha(E_1)f^\beta(E_2) = f^\alpha(E_3)f^\beta(E_4)$$

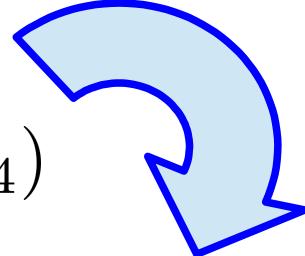
$$\gamma_{\alpha\alpha} L^{\alpha\alpha}(E_1) + \gamma_{\beta\beta} L^{\beta\beta}(E_2) = \gamma_{\alpha\alpha} L^{\alpha\alpha}(E_3) + \gamma_{\beta\beta} L^{\beta\beta}(E_4)$$

# A toy-model – detailed balance

for  $\alpha = \beta$

$$f^\alpha(E_1)f^\alpha(E_2) = f^\alpha(E_3)f^\alpha(E_4)$$

$$L^{\alpha\alpha}(E_1) + L^{\alpha\alpha}(E_2) = L^{\alpha\alpha}(E_3) + L^{\alpha\alpha}(E_4)$$

$$f^\alpha \sim e^{-\gamma_{\alpha\alpha} L^{\alpha\alpha}(E)}$$


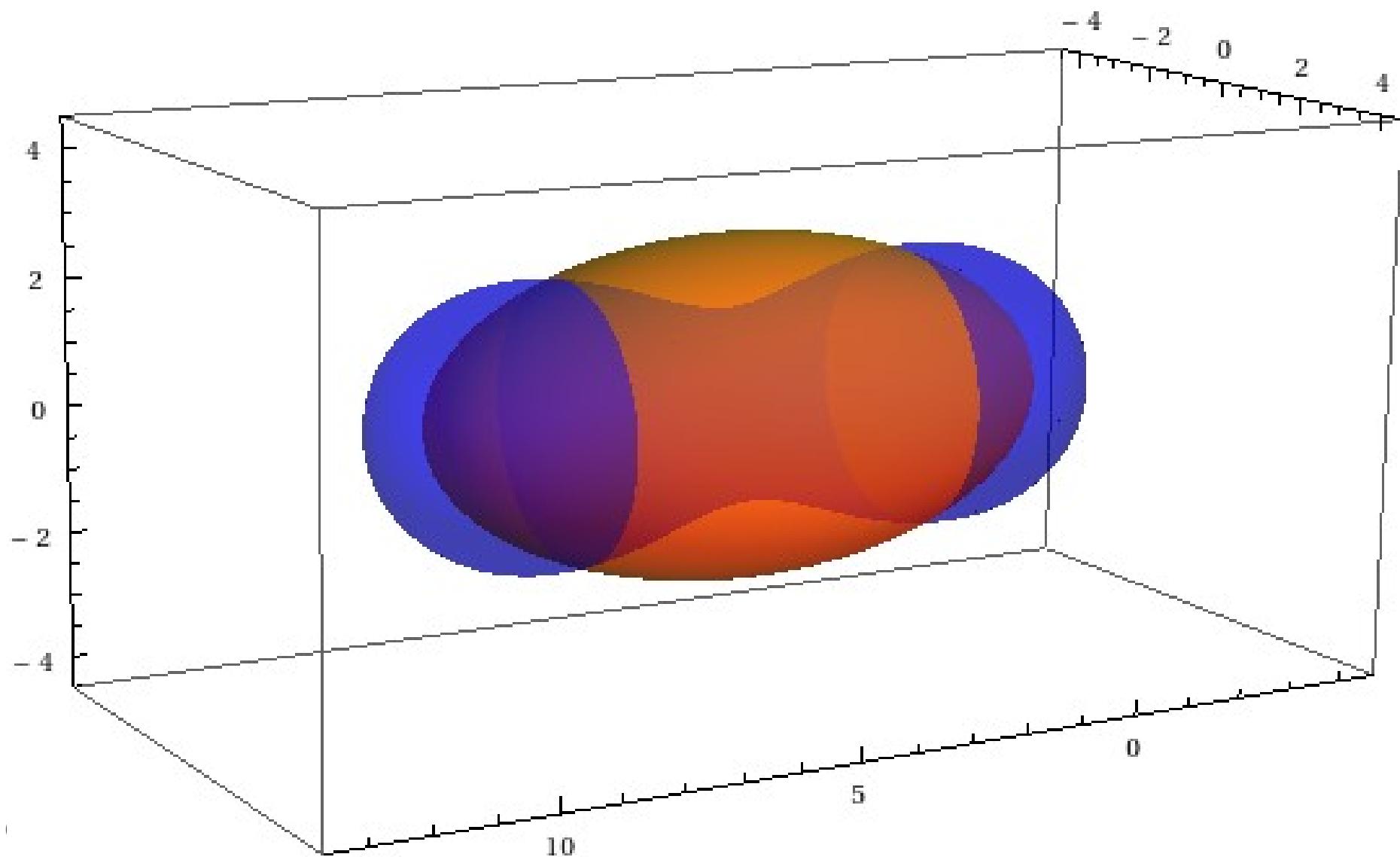
$$\gamma_{\alpha\alpha} L^{\alpha\alpha}(E_1) + \gamma_{\beta\beta} L^{\beta\beta}(E_2) = \gamma_{\alpha\alpha} L^{\alpha\alpha}(E_3) + \gamma_{\beta\beta} L^{\beta\beta}(E_4)$$

AND for  $\alpha \neq \beta$

$$L^{\alpha\beta}(E_1) + L^{\alpha\beta}(E_2) = L^{\alpha\beta}(E_3) + L^{\alpha\beta}(E_4)$$

# A toy-model

$d=3, n=2$ : closed curves



# A toy-model

---

Lot of constraints...

$$L^{\alpha\beta}(E_1) + L^{\alpha\beta}(E_2) = L^{\alpha\beta}(E_3) + L^{\alpha\beta}(E_4) \quad \alpha, \beta = 1 \dots n$$

cannot be fulfilled for every pairs



detailed balance does not exist



no global conservations

**Long time behaviour?**

# A toy-model – numerical results

---

- collisional invariants:

$$E_1 + E_2 + \frac{a^{\alpha\beta}}{\langle E \rangle^n} E_1 E_2 = E_3 + E_4 + \frac{a^{\alpha\beta}}{\langle E \rangle^n} E_3 E_4$$

- scaling distributions (NESS):  $f^\alpha(E) = \frac{1}{\langle E \rangle} \phi(E/\langle E \rangle)$

$$\phi(x) \sim (1 + Ax)^{-B}$$

# A toy-model – numerical results

- collisional invariants:

$$E_1 + E_2 + \frac{a^{\alpha\beta}}{\langle E \rangle^n} E_1 E_2 = E_3 + E_4 + \frac{a^{\alpha\beta}}{\langle E \rangle^n} E_3 E_4$$

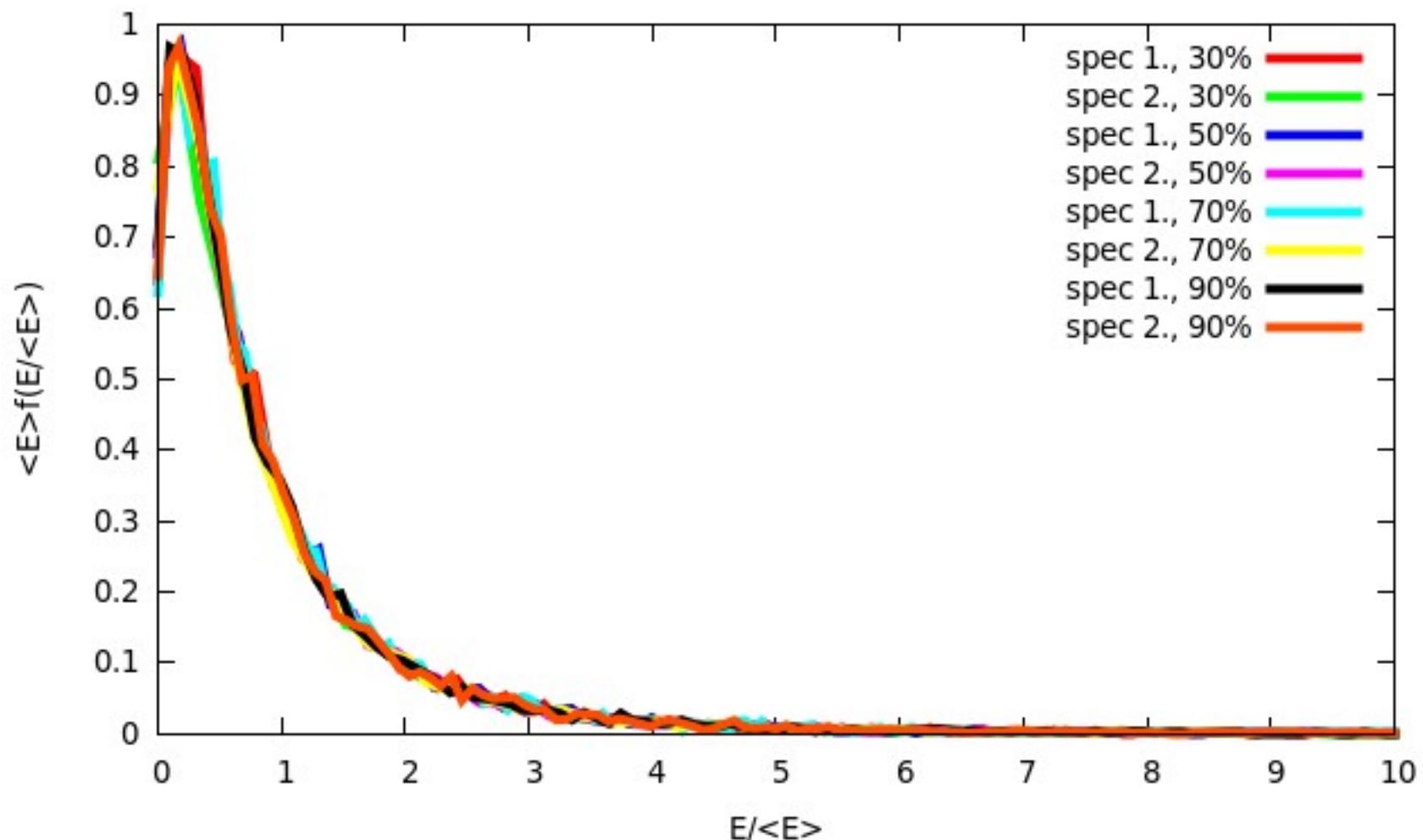
- scaling distributions (NESS):  $f^\alpha(E) = \frac{1}{\langle E \rangle} \phi(E/\langle E \rangle)$

$$\phi(x) \sim (1 + Ax)^{-B}$$

determined by initial values

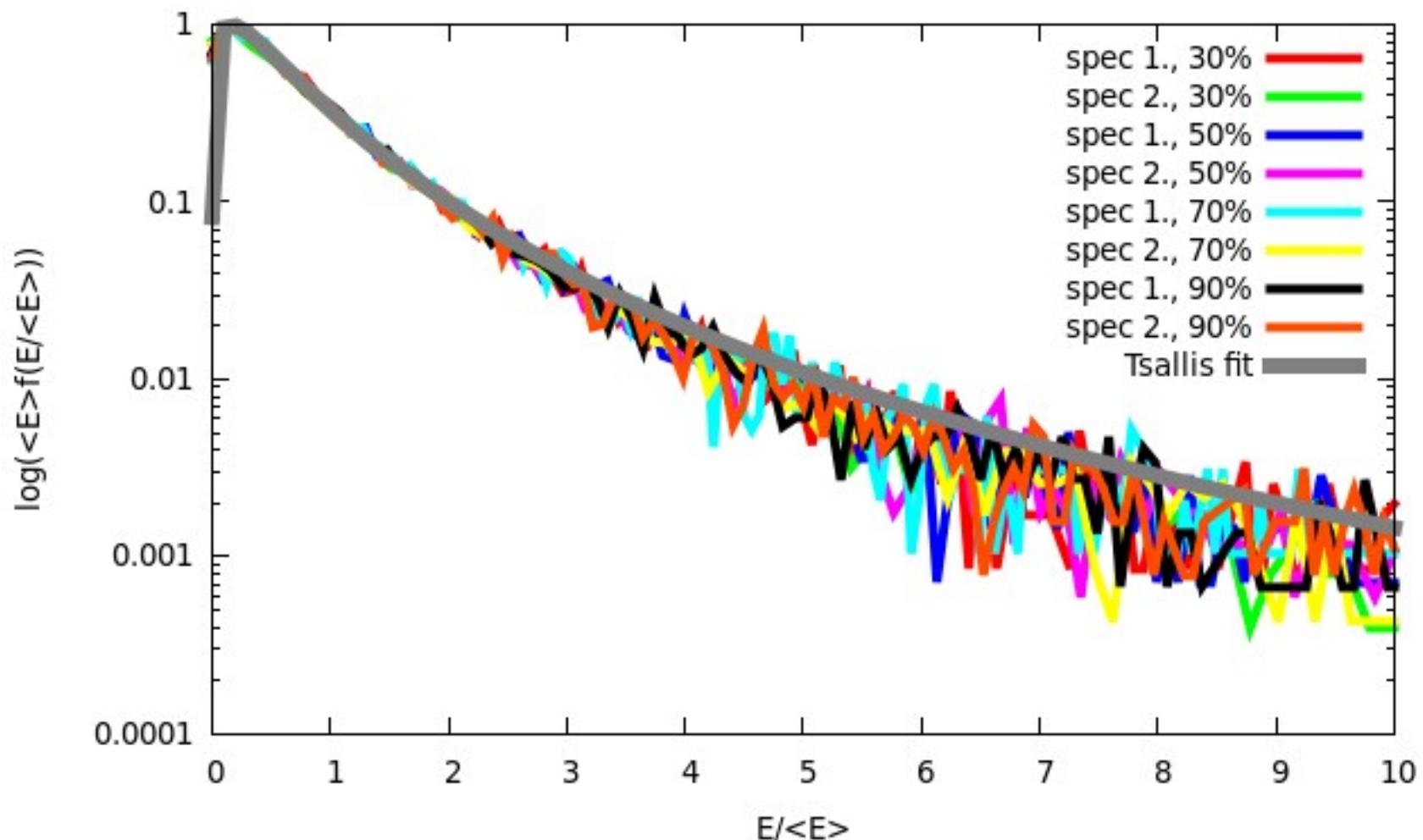
# A toy-model – numerical results

- scaling distributions:  $\langle E \rangle f^\alpha(E) = \phi(E/\langle E \rangle)$



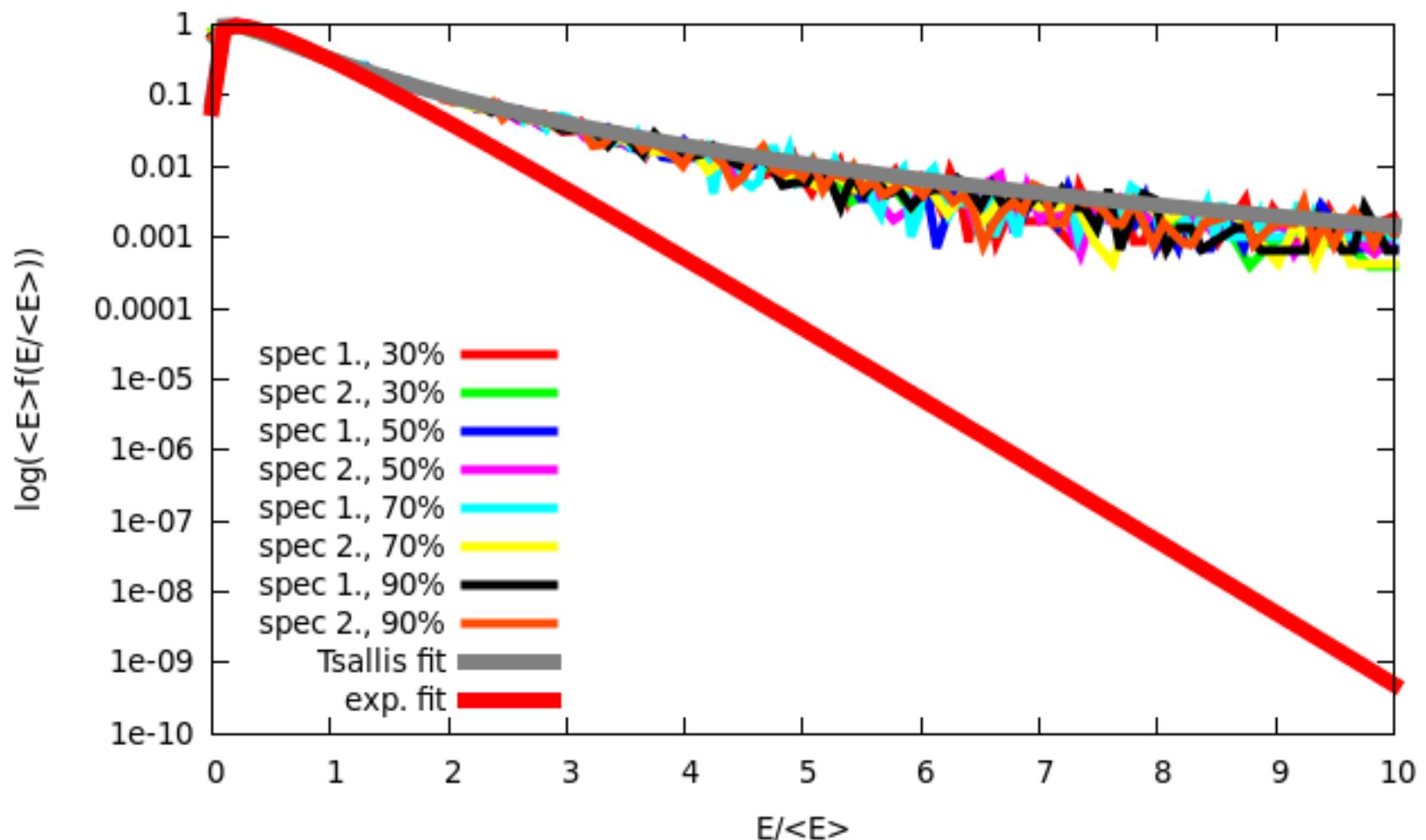
# A toy-model – numerical results

- scaling distributions:  $\langle E \rangle f^\alpha(E) = \phi(E/\langle E \rangle)$



# A toy-model – numerical results

- scaling distributions:  $\langle E \rangle f^\alpha(E) = \phi(E/\langle E \rangle)$



## A toy-model – $a \sim 1/\langle E \rangle$

$$\frac{d}{d\tau} \langle E \rangle = \sum_{\alpha} \int_0^{\Lambda} dE \dot{f}^{\alpha}(E) E = \langle E \rangle \sum_{\alpha\beta} I^{\alpha\beta}$$

# A toy-model – $a \sim 1/\langle E \rangle$

$$\frac{d}{d\tau} \langle E \rangle = \sum_{\alpha} \int_0^{\Lambda} dE \dot{f}^{\alpha}(E) E = \langle E \rangle \sum_{\alpha\beta} I^{\alpha\beta}$$

$\langle E \rangle$  scaled out: no time dependence



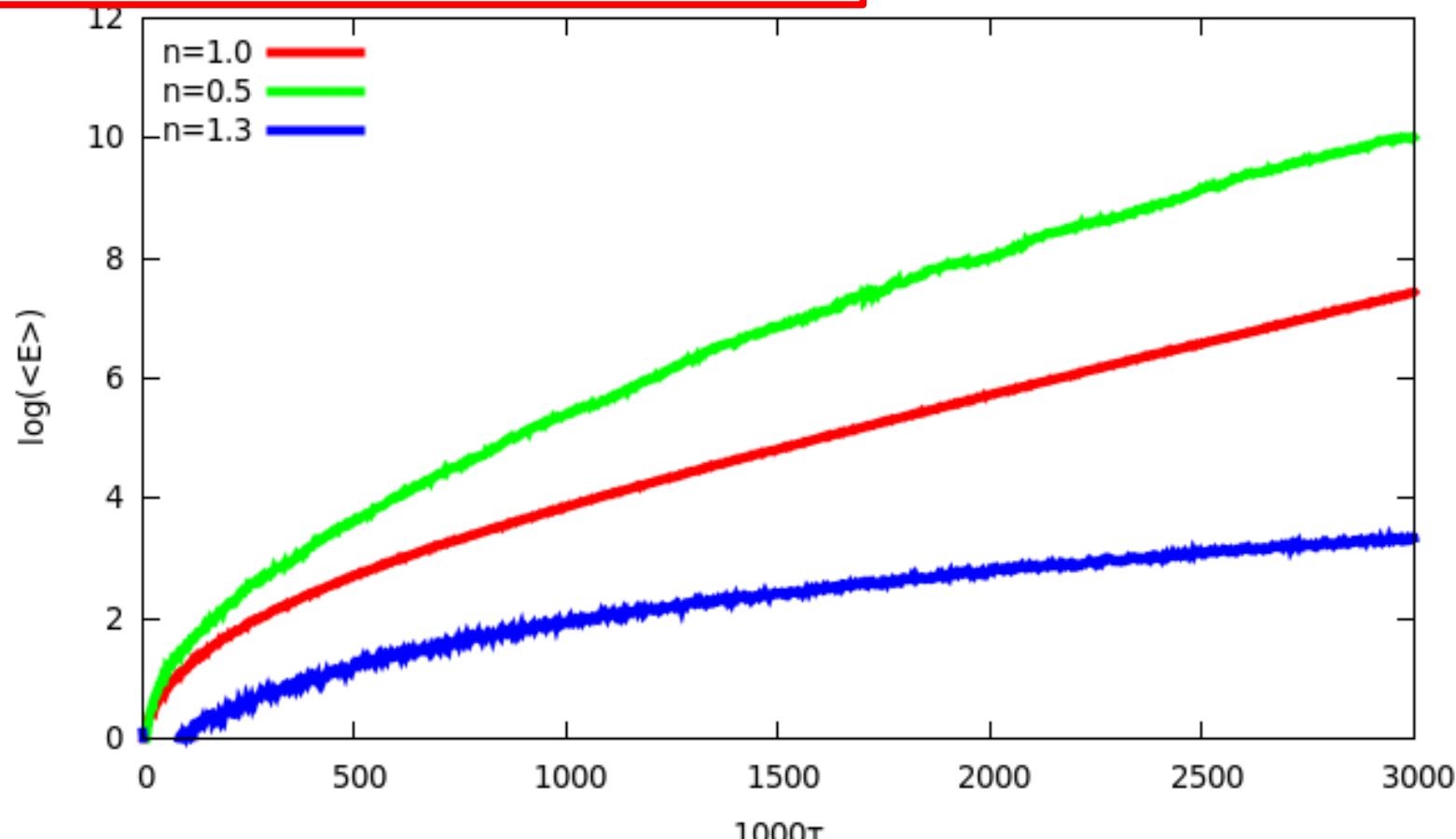
$$\langle E \rangle \sim e^{\gamma\tau}$$

# A toy-model – $a \sim 1/\langle E \rangle$

$$\frac{d}{d\tau} \langle E \rangle = \sum_{\alpha} \int_0^{\Lambda} dE \dot{f}^{\alpha}(E) E = \langle E \rangle \sum_{\alpha\beta} I^{\alpha\beta}$$

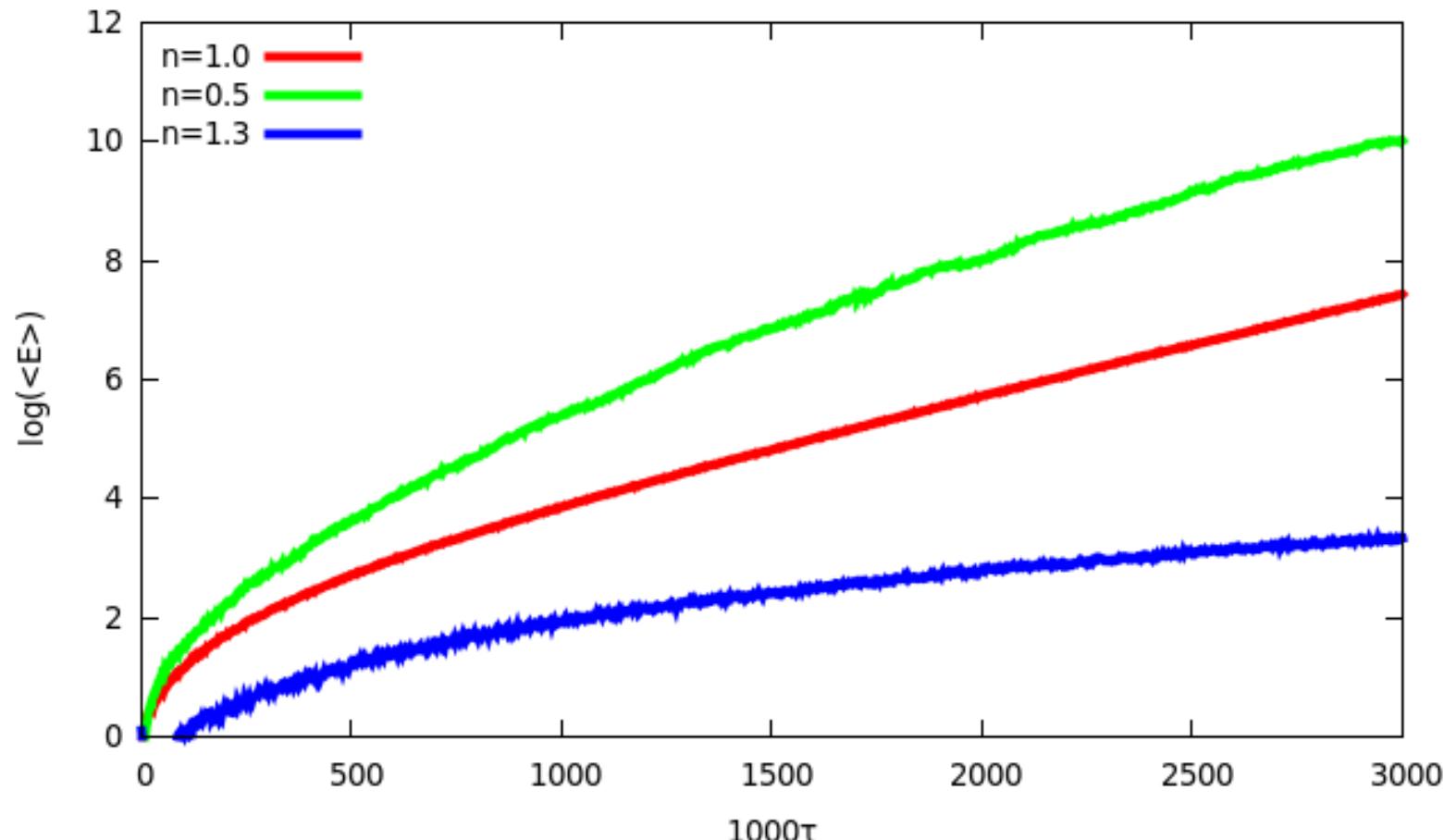
$\langle E \rangle$  scaled out: no time dependence

$$\langle E \rangle \sim e^{\gamma\tau}$$

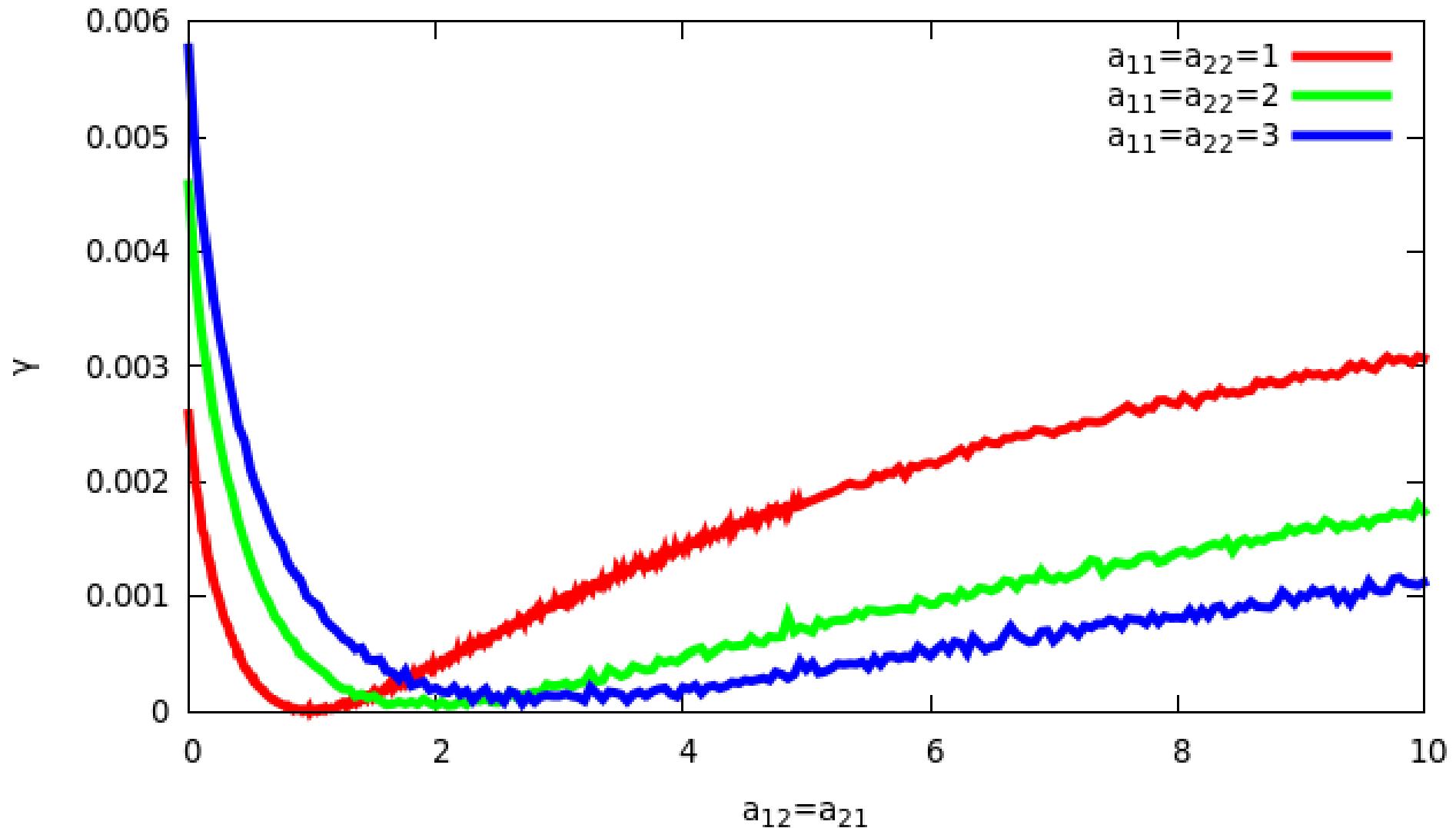


# A toy-model – $a \sim 1/\langle E \rangle$

$$S_{en} = - \sum_{\alpha} \int_0^{\Lambda} dE f^{\alpha}(E) \ln f^{\alpha}(E) = \text{const.} + \ln(\langle E \rangle)$$



# A toy-model – $a \sim 1/\langle E \rangle$



# Conclusions & summary

---

- Simple kinetic toy-model with non-additive coll. invariants:
  - two (or more) component: frustration (disappears continuously)  
(balance states will be different for each component)
  - no detailed balance, but NESS in long time
  - thermal description: heating system (driving)
  - stable equation of states
- Motivation for energetically non-extensive models:
  - dense plasmas,
  - dealing with >2-part. collisions in a mean-field way

# Acknowledgement

---

- Antal JAKOVÁC (Eötvös University)
- Péter VÁN (Wigner RCP)
- Károly ÜRMÖSSY (Wigner RCP)

Thank you for your attention!

---

Questions?