

BGV detector design studies

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Number of reconstructed tracks and stat. precision

- The requirement to have at least X tracks per reconstructed vertex has relation both to:
 - The statistical precision of the measured σ_{beam} / the gas pressure
 - The primary vertex (PV) resolution (and the related systematic error)

Rate of inelastic beam-gas interactions per bunch:

$$R_{\text{inel}} = \int_{z=z_1}^{z=z_2} \rho(z) dz \cdot \sigma_{\text{pA}}(E) \cdot N \cdot f_{\text{rev}}$$

- Not all inelastic events will be useful
- The vertex resolution will be sufficiently good only for events with at least N_{Tr} reconstructed tracks. The fraction of these events, F_{good} , depends on:
 - the geometrical distributions (η) of the beam-gas interaction products
 - the detector geometry
- $N_{\text{good}} = R_{\text{inel}} \Delta t F_{\text{good}}$ determines the statistical precision of the measured σ_{beam} :

$$\frac{\delta\sigma_{\text{beam}}}{\sigma_{\text{beam}}} = \frac{1}{\sqrt{2 N_{\text{good}}}}$$
- The knowledge of F_{good} is essential in the **detector** + **gas-target** design

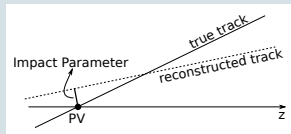
Number of reconstructed tracks and PV resolution

- As discussed previously, the primary vertex (PV) resolution in a given event depends on the number of measured tracks
- Assuming that all tracks in the event have the same impact parameter (IP) resolution σ_{IP} , the PV resolution scales approximately as $\sigma_{IP}/\sqrt{N_{Tr}}$
 - N_{Tr} : number of tracks making up the vertex
 - Many effects: different track p_T , different z_{vtx} , ...

Reminder of the IP definition and formulas:

The impact parameter resolution is determined by:

- σ_{MS} – IP induced by multiple scattering (MS)
- σ_{extrap} – IP induced by detector hit resolution



$$\sigma_{IP}^2 = \sigma_{MS}^2 + \sigma_{extrap}^2$$

In each transverse coordinate:

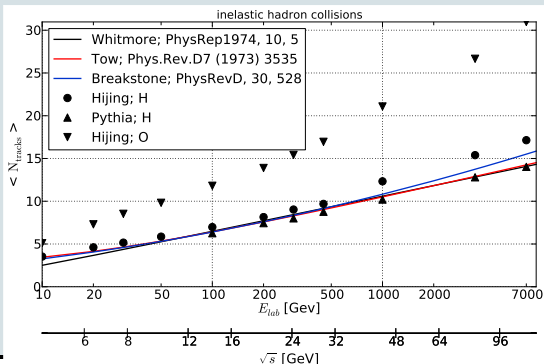
$$\sigma_{MS} = r_1 \frac{13.6 \text{ MeV}}{p_T} \sqrt{\frac{x}{X_0}} \left(1 + 0.038 \log \frac{x}{X_0} \right) \approx r_1 \frac{13.6 \text{ MeV}}{p_T} \sqrt{\frac{x}{X_0}}$$

$$\sigma_{extrap} = \sqrt{\frac{z_2^2 \sigma_1^2 + z_1^2 \sigma_2^2}{(z_2 - z_1)^2}} \cos^2 \theta$$

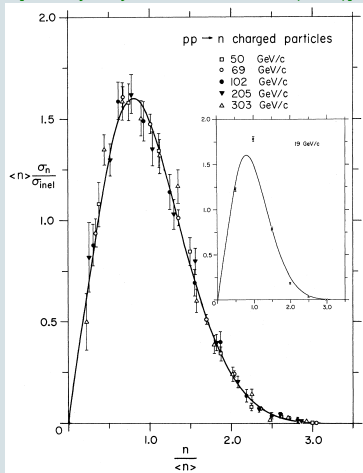
σ_1 and σ_2 are the detector hit resolutions, which in principle can depend on the radial position

Average charged particle multiplicities

- The average number of charged particles produced in pp collisions have been measured by different experiments at different center-of-mass energy \sqrt{s}
 - Later I refer to this number as $\langle N_{\text{Tr}} \rangle$
 - Different parametrizations exist as $f(s)$. E.g.:
 - $\langle N_{\text{Tr}} \rangle (s) = A + B/\sqrt{s} + C \ln s$ [Tow; Phys.Rev.D7 (1973) 3535]
- To-do: make comparisons with experimental data of pA collisions
- Generated beam-gas interactions with PYTHIA (target = H) and HIJING (target = H, O or Xe), using the LHCb computing framework
 - Compare $\langle N_{\text{Tr}} \rangle$ in the simulated samples with the parametrizations based on experimental data



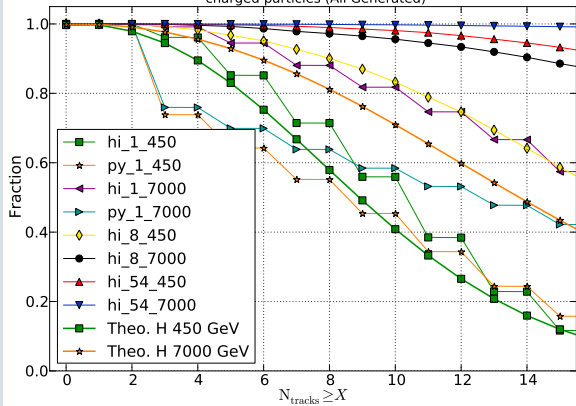
[Slattery; Phys.Rev.Lett.29, 1624 (1972)]



- For different \sqrt{s} the distributions of N_{Tr} are different (larger energy \Rightarrow larger $\langle N_{Tr} \rangle$)
- However, it was found that the distribution of $\langle N_{Tr} \rangle \frac{\sigma_n}{\sigma_{inel}}$ vs $\frac{N_{Tr}}{\langle N_{Tr} \rangle}$ does not depend on \sqrt{s} , at least in the \sqrt{s} range we are interested in (29 – 114 GeV).
 - σ_n is the so-called “topological” cross-section: cross-section for an interaction to produce exactly n charged particles.
- For a fixed \sqrt{s} we can use the parametrization of this distribution to determine the fraction of events producing at least n tracks
 - This is our expectation, which we compare with the results from the MC generated samples

Fraction of events with at least X tracks

Fraction of events with at least X charged particles (All Generated)

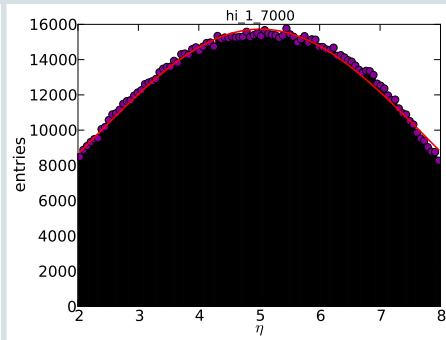
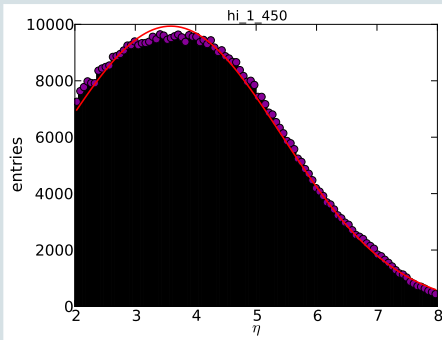


- ▶ hi/py : Hijing or Pythia MC generators
- ▶ 1/8/54: H or O or Xe target
- ▶ 450/7000: beam energy in GeV

- The step-like behavior in the pH samples is attributed to el. charge conservation
- The agreement between PYTHIA and HIJING is not so good (elastic interactions!)
- The agreement between the expectation and the MC samples is not so bad
- Overall, acceptable agreement between MC simulation and parametrization based on real data

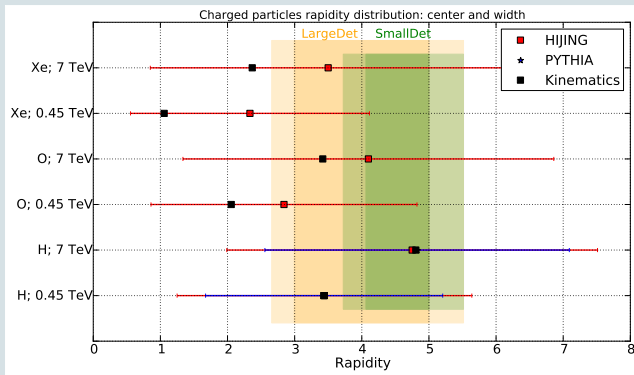
η distributions & detector acceptance (1)

- These were results for **all** charged particles (indep. of detector acceptance)
 - In practice we care about the number of charged particles **in detector acceptance**
- The pseudorapidity distributions are close to Gaussian; Examples with Hijing, H, 450/7000 GeV:



- Next, compare the fitted Gaussians mean and width for different targets and energy

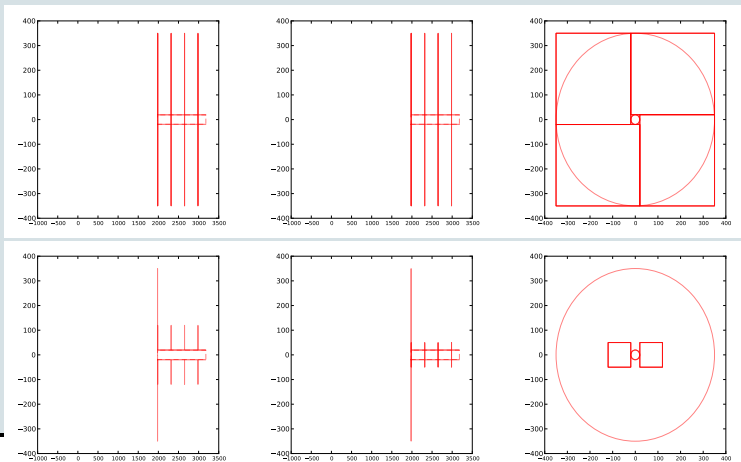
η distributions & detector acceptance (2)



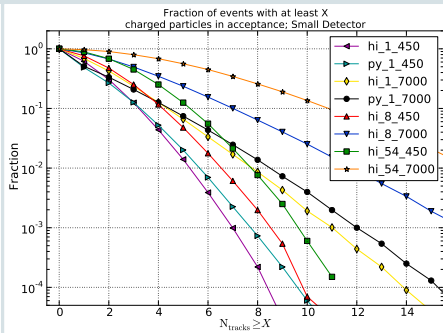
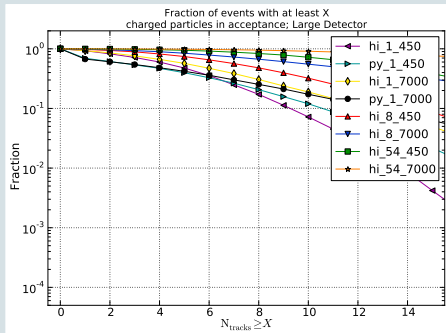
- The rectangles represent the η coverage of the *Large* and *Small* detectors.
 - For each detector type, two rectangles show the coverage for z_{min} (-0.5 m) and z_{max} (0.5 m)
- “Kinematics”: Calculate boost assuming a **rigid** target. The velocity of the center-of-mass in the lab frame is: $\beta_{cm} = p_1^{lab} / (E_1^{lab} + m_2)$. Under a boost in the z -direction to a frame with velocity β : $y \rightarrow y - \tanh^{-1} \beta$
 - Good agreement for Hydrogen; For heavier targets, the MC simulation results are somewhere between a rigid target and a single-nucleon target

Detector geometries used in toy MC simulations

- Both detectors:
 - 4 $x + y$ measuring layers, located at $z = 1984.0, 2317.3, 2650.7,$ and 2984.0 mm ($z = 0$ is the center of the gas target)
 - 1 mm thick exit window
 - beam pipe with inner/outer radius of 19/20 mm
- *Large detector*: covers $20 \leq r \leq 350$ mm
- *Small detector*: made of 100×100 mm sensors on the left/right of the beam pipe



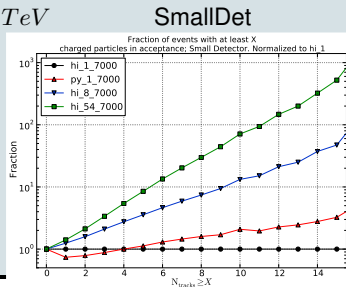
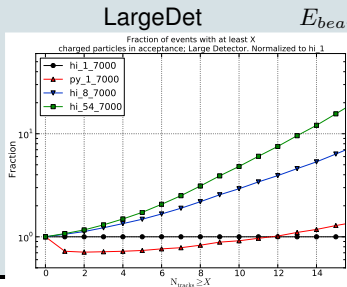
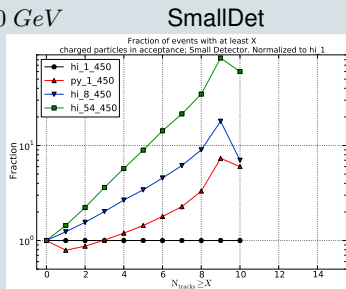
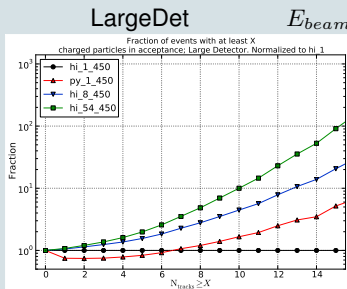
Fraction of events with at least X tracks in Acceptance



- Reminder: F_{good} and the gas pressure determine the time needed to achieve certain statistical precision on the measured σ_{beam}
- F_{good} itself depends on several parameters, including detector geometry and cut on N_{Tr}

Fraction of events with at least X tracks in Acceptance

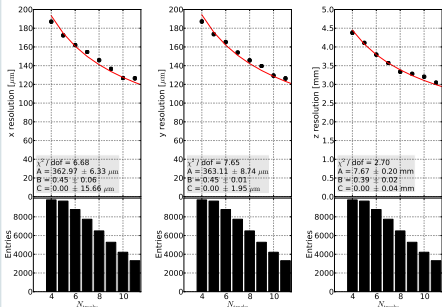
- Check the effect on F_{good} from using heavier targets



- Determine PV resolution as function of the cut on N_{Tr} , using the toy MC detector simulation tool

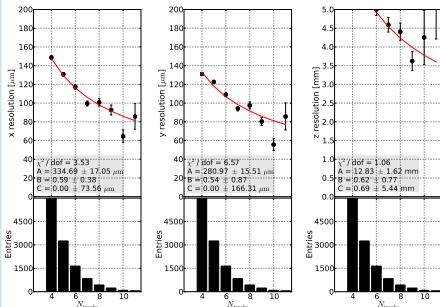
LargeDet

PV resolutions fitted to $A / N_{tracks}^B + C$



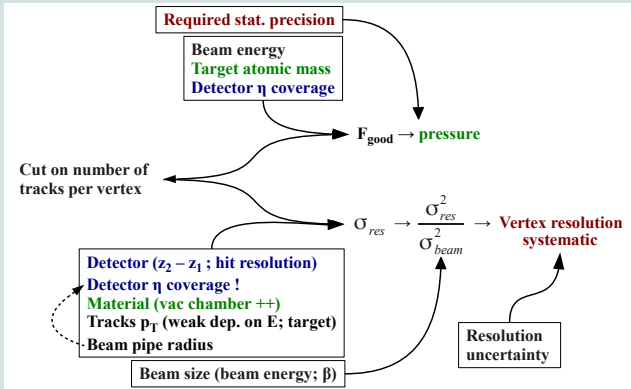
SmallDet

PV resolutions fitted to $A / N_{tracks}^B + C$



- These are 7 TeV H collisions. The results at 450 GeV are very similar
- However, we see a significant difference with the two detectors: the small one is more precise in x and y , and less precise in z
 - We attribute this difference to the different average η of the measured tracks: lower η tracks provide better constraint on the transverse coordinates. A calculation is needed for confirmation

Design scheme



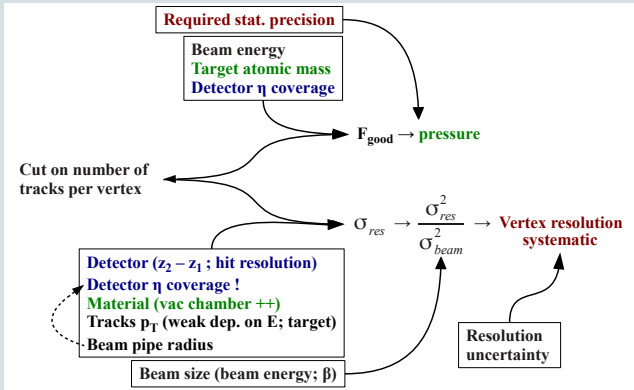
Fix stat. precision (5 % per bunch) and time (3 min) \rightarrow need ~ 1 Hz per bunch. Determine needed pressure, given certain F_{good} \rightarrow see Massi's talk

Vertex resolution systematic:

$$\frac{\delta\sigma_{beam}}{\sigma_{beam}} = \frac{\sigma_{res}^2}{\sigma_{beam}^2} \cdot \frac{\delta\sigma_{res}}{\sigma_{res}}$$

- Note that certain properties (e.g. detector η coverage) have implications both on the stat. and syst. precisions
- It is a complex inter-connected system: fixing parameters like r_{pipe} and beam size (optics β) will facilitate greatly the identification of an optimal design

Design scheme



Fix stat. precision (5 % per bunch) and time (3 min) → need ~ 1 Hz per bunch. Determine needed pressure, given certain F_{good} → see Massi's talk

Vertex resolution systematic:

$$\frac{\delta\sigma_{beam}}{\sigma_{beam}} = \frac{\sigma_{res}^2}{\sigma_{beam}^2} \cdot \frac{\delta\sigma_{res}}{\sigma_{res}}$$

- Why a measurement at 7 TeV is more challenging than at 0.45 TeV:
- When we go from 7 TeV (where low σ_{res} is the challenging requirement) to 450 GeV:
 - F_{good} decreases significantly, but if we choose lower N_{Tr} cut (3/4/5 ?) then:
 - F_{good} gets larger by about a factor of 10 (2.5) for small (large) detector, in comparison to 7 TeV → can operate at lower pressure
 - σ_{res} increases by about a factor of 1.5, but the ratio $(\sigma_{res}/\sigma_{beam})^2$ gets 7 times smaller!

Additional Slides

Default definition

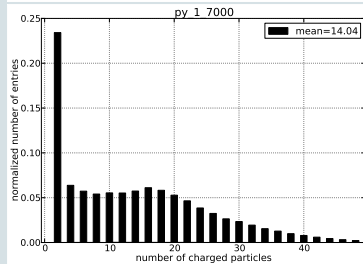
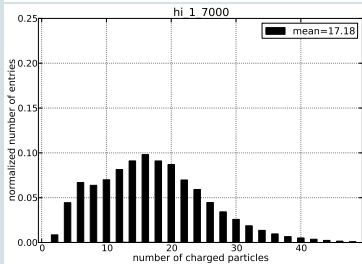
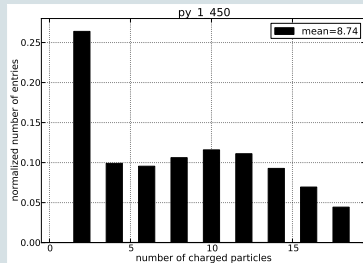
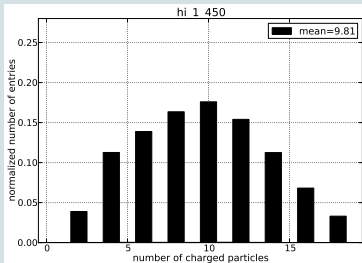
- Count as charged particles all generated particles that:
 - 1. Have electric charge = ± 1
 - 2. Are produced at most 5 mm away from the PV (distance in z)
 - 3. Have no decay vertices or have exactly 1 decay vertex, which is located at least 10 mm away from the PV (distance in z)

Alternative definition

- Count as charged particles all generated particles that:
 - 1. Same as above
 - 2. Same as above
 - 3. Have no decay vertices or have exactly 1 decay vertex, and the distance between the *end* and *origin* vertex is at least 2 m (distance in z)
- Both definitions give the same results within 1 %. The reason is that there are very few particles that have decay length between 10 mm and 2 m ($K_S, \Lambda, ?$)

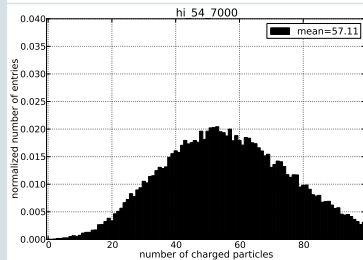
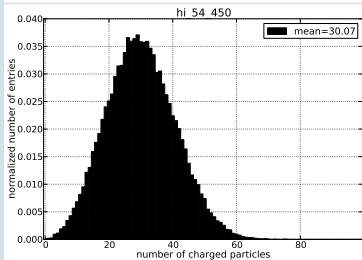
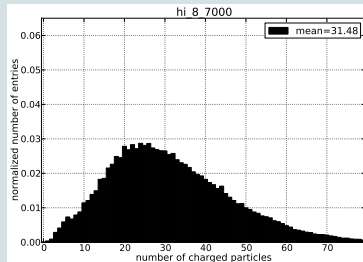
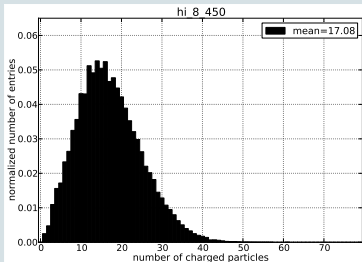
Distribution of charged particles multiplicity (1)

- 4 generated samples: Hijing/Pythia; Hydrogen; 450/7000 GeV
 - no selection cuts



Distribution of charged particles multiplicity (2)

- 4 generated samples: Hijing; Oxygen/Xenon; 450/7000 GeV
 - no selection cuts



Geometry definition: *Large detector*

- x and y hit resolution = $58 \mu\text{m}$

```

# typ ODW   IDH     L       posx   posy   posz   X0
# exit window (very short large diameter tube)
  1 700.0   40.0       1.0     0.0    0.0  1980.0  89.0
# beam pipe (long small diameter tube)
  1 40.0    38.0   1200.0    0.0    0.0  2580.0  89.0
# first plane of XY sensors (4 rectangles) covering 20 <= r <= 350 mm
  2 370.0   330.0       1.0  165.0  185.0  1984.0  93.0
  2 330.0   370.0       1.0  185.0 -165.0  1984.0  93.0
  2 370.0   330.0  1.0 -165.0 -185.0  1984.0  93.0
  2 330.0   370.0  1.0 -185.0  165.0  1984.0  93.0
# second plane of XY sensors (4 rectangles) covering 20 <= r <= 350 mm
  2 370.0   330.0       1.0  165.0  185.0  2317.3  93.0
  2 330.0   370.0       1.0  185.0 -165.0  2317.3  93.0
  2 370.0   330.0  1.0 -165.0 -185.0  2317.3  93.0
  2 330.0   370.0  1.0 -185.0  165.0  2317.3  93.0
# third plane of XY sensors (4 rectangles) covering 20 <= r <= 350 mm
  2 370.0   330.0       1.0  165.0  185.0  2650.7  93.0
  2 330.0   370.0       1.0  185.0 -165.0  2650.7  93.0
  2 370.0   330.0  1.0 -165.0 -185.0  2650.7  93.0
  2 330.0   370.0  1.0 -185.0  165.0  2650.7  93.0
# fourth plane of XY sensors (4 rectangles) covering 20 <= r <= 350 mm
  2 370.0   330.0       1.0  165.0  185.0  2984.0  93.0
  2 330.0   370.0       1.0  185.0 -165.0  2984.0  93.0
  2 370.0   330.0  1.0 -165.0 -185.0  2984.0  93.0
  2 330.0   370.0  1.0 -185.0  165.0  2984.0  93.0

```

Geometry definition: *Small detector*

- x and y hit resolution = $58 \mu\text{m}$

```

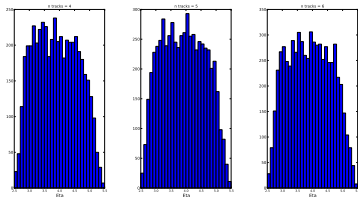
# typ ODW      IDH      L      posx   posy   posz      X0
# exit window (very short large diameter tube)
  1 700.0   40.0      1.0    0.0    0.0  1980.0  89.0
# beam pipe (long small diameter tube)
  1 40.0    38.0  1200.0    0.0    0.0  2580.0  89.0
# first plane of XY sensors: one rectangle (10x10 cm) on the left/right of the beam
  2 100.0   100.0      1.0   70.0    0.0  1984.0  93.0
  2 100.0   100.0      1.0  -70.0    0.0  1984.0  93.0
# second plane of XY sensors: one rectangle (10x10 cm) on the left/right of the beam
  2 100.0   100.0      1.0   70.0    0.0  2317.3  93.0
  2 100.0   100.0      1.0  -70.0    0.0  2317.3  93.0
# third plane of XY sensors: one rectangle (10x10 cm) on the left/right of the beam
  2 100.0   100.0      1.0   70.0    0.0  2650.7  93.0
  2 100.0   100.0      1.0  -70.0    0.0  2650.7  93.0
# fourth plane of XY sensors: one rectangle (10x10 cm) on the left/right of the beam
  2 100.0   100.0      1.0   70.0    0.0  2984.0  93.0
  2 100.0   100.0      1.0  -70.0    0.0  2984.0  93.0

```

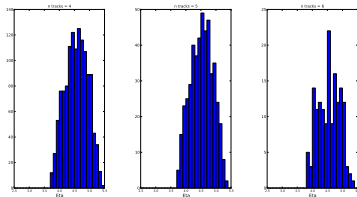
PV reconstruction: η distributions

- 0.45 TeV Hydrogen collision

LargeDet



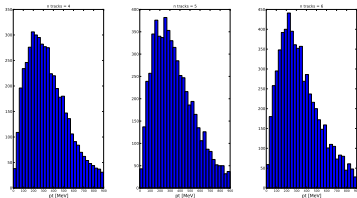
SmallDet



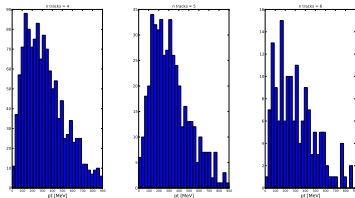
Detector Type	$N_{Tr} > X$	$\langle \eta \rangle$
<i>Large</i>	4	3.90
	5	3.92
	6	3.91
<i>Small</i>	4	4.57
	5	4.53
	6	4.56

- 0.45 TeV Hydrogen collision

LargeDet



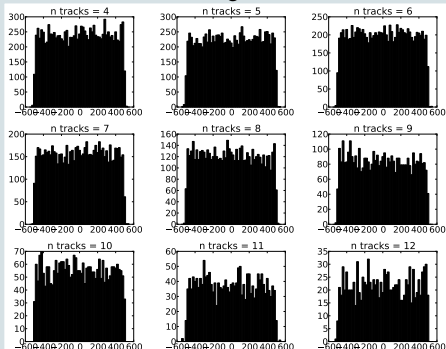
SmallDet



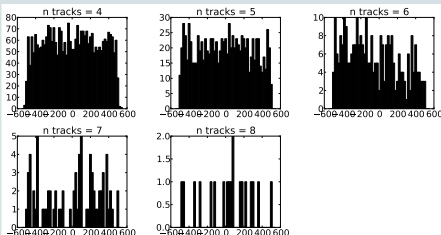
Detector Type	$N_{Tr} > X$	$\langle p_T \rangle$ [MeV]
<i>Large</i>	4	366.2
	5	365.4
	6	358.2
<i>Small</i>	4	345.6
	5	350.8
	6	324.8

- 450 GeV Hydrogen collision

LargeDet

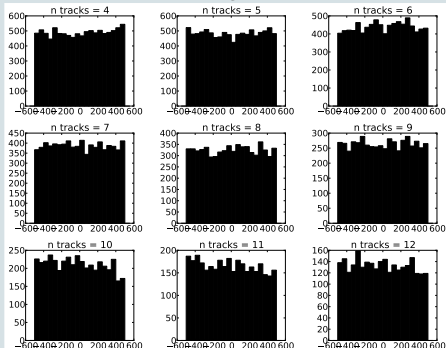


SmallDet



- 7 TeV Hydrogen collision

LargeDet



SmallDet

